Heavy lons: Experiments

Francesco Prino – Sezione di Torino

EUROPEAN PHYSICAL

HEP2025

MARSEILLE



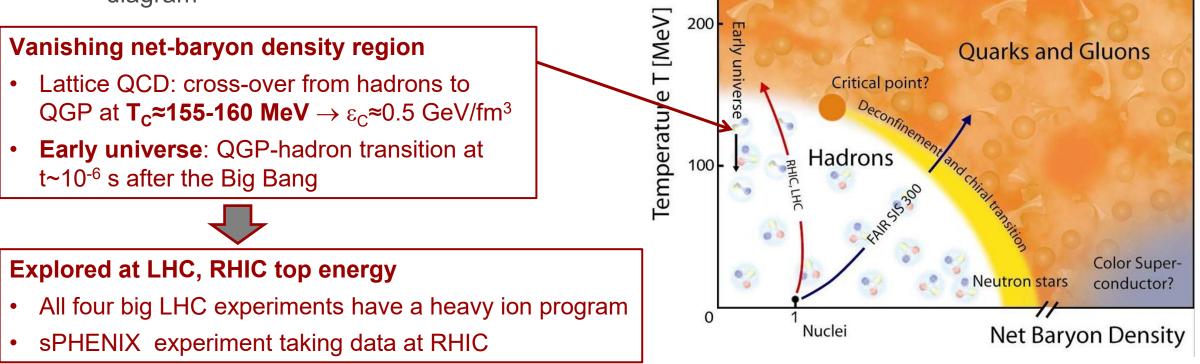
Marseille, July 11th 2025



Marseille, July 11th 2025

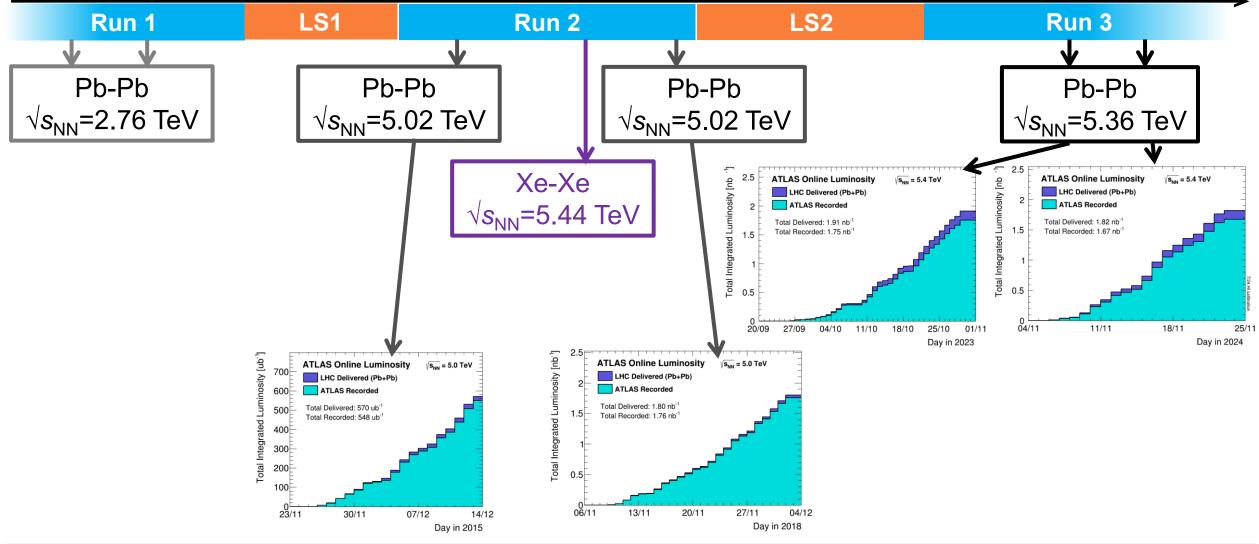
Heavy-ion collisions and QCD

- Goal: study the properties of strongly-interacting matter at extreme conditions of temperature and energy density
 - Explore the rich phase diagram of QCD matter
 - Transition to a state where quarks and gluons are deconfined (Quark Gluon Plasma, QGP) and chiral symmetry is restored
 - Heavy-ion collisions at different energies allow us to explore different regions of QCD phase diagram



Heavy-ion collisions at the LHC

2010 2011 2012 2013 2014 2015 2016 2017 2018 2019 2020 2021 2022 2023 2024 2025



Colliding systems

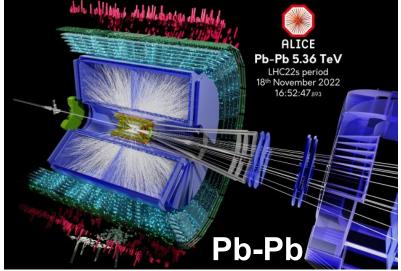
• Heavy-ion (A-A) collisions:

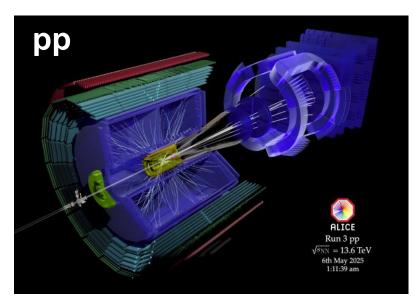
Produce extremely hot and dense QCD matter
 Access the regime of QCD phase transition and study QGP properties

• And reference systems:

➡Proton-nucleus (p-A) collisions

- ✓ Initially aimed at studying "cold nuclear matter" effects and disentangle them from QGP effects
- ➡Proton-proton (pp) collisions
 - ✓ Initially aimed at providing the necessary reference for A-A collisions





Colliding systems

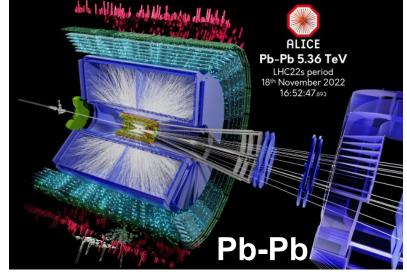
• Heavy-ion (A-A) collisions:

Produce extremely hot and dense QCD matter
 Access the regime of QCD phase transition and study QGP properties

• And reference systems:

➡Proton-nucleus (p-A) collisions

- ✓ Initially aimed at studying "cold nuclear matter" effects and disentangle them from QGP effects
- ➡Proton-proton (pp) collisions
 - ✓ Initially aimed at providing the necessary reference for A-A collisions

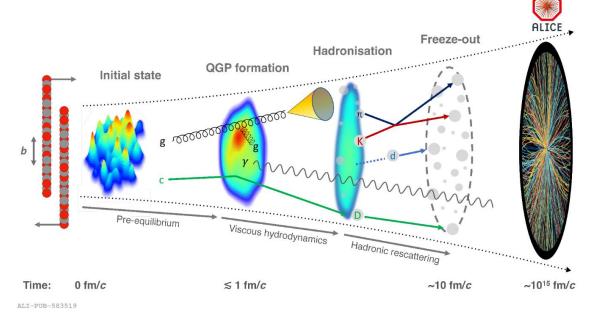


• Not just reference systems:

➡ Discovery at the LHC: high-multiplicity pp and p-A collisions show "features" resembling those observed in A-A collisions and understood as due to the collective expansion of the QGP

Paradigm shift: chasing the small-size limit of the QGP

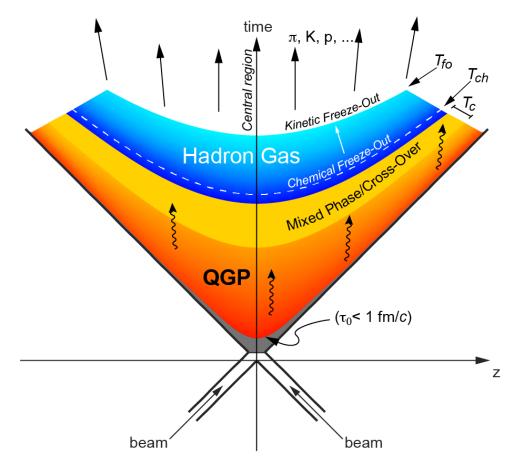
Space-time evolution of A-A collisions



- Heavy-ion collisions lead to large energy deposition in the collision region
- Formation of a "strongly coupled" QGP
 - ⇒ Expands (and cools down) under strong pressure gradients according to viscous hydrodynamics
- Transition to hadrons when the temperature drops below the (pseudo)critical value
- Hadron gas phase after QGP hadronization until decoupling / freeze-out
- System lifetime ~ a few fm/*c*

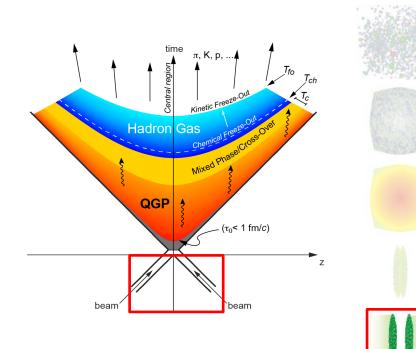
 \Rightarrow Long with respect to elementary collisions in high-energy particle physics \rightarrow space-time picture

Get insight into QGP properties



- How to investigate the QGP which stays "beyond the horizon" of the freeze-out?
- Three classes of observables (multimessengers from the QGP):
 - Soft particles: Hadrons with small momentum, decoupling from the equilibrated system at the kinetic freeze-out
 - Electromagnetic probes: Photons and dileptons produced in the hot phase and not influenced by later stages
 - ➡ Hard probes: Energetic and/or massive particles / jets produced in hard-scattering processes in the early stages of collision and traversing the QGP

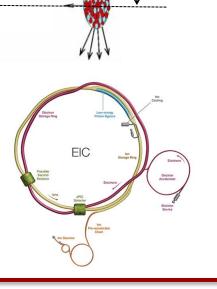
Before diving in the QGP: initial state



• What is the structure of the colliding objects?

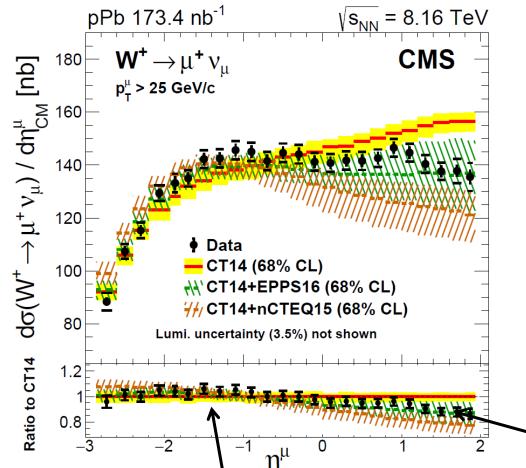
Spatial distribution of incoming nucleons inside nuclei
 Momentum distribution of partons inside nucleons
 Modification of the PDFs in bound nucleons (nPDF)
 Gluon saturation at small Bjorken-x

- Insight into initial state via:
 - ⇒p-A collisions
 - ➡ Ultra-peripheral A-A collisions
 - ✓ The impact parameter of the two colliding nuclei is greater than the sum of the two nuclear radii
 - The strong electromagnetic fields surrounding the nuclei give rise to γA interactions
 - Future measurements in e-A interactions at the electron-ion collider



b > R1+R2

Initial state: W bosons in p-Pb collisions



Negative rapidity (Pb-going):

ElectroWeak bosons

rightarrow Produced on short timescales + lifetime ~ 0.1 fm/c

➡ Decay into energetic leptons

 Do not interact strongly with other particles produced in the collision

⇒ Encode information about the early stages

 Measured cross sections in p-Pb collisions described by predictions using nPDFs better than with those using the free proton PDFs

Results incorporated into global fits to extract the parton densities in heavy nuclei

CMS, Phys.Rept. 1115 (2025) 219

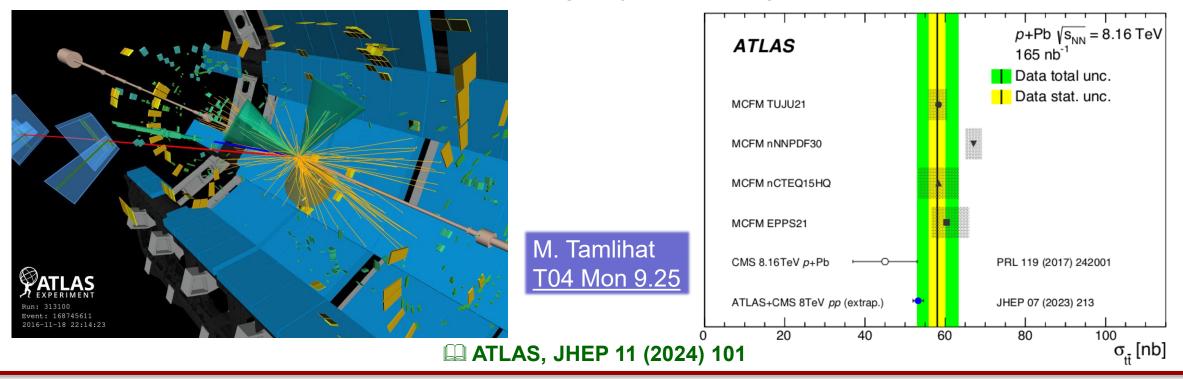
Positive rapidity (p-going):

Higher-*x* partons in Pb nucleus \rightarrow antishadowing $\|$ Low-*x* partons in Pb nucleus \rightarrow shadowing

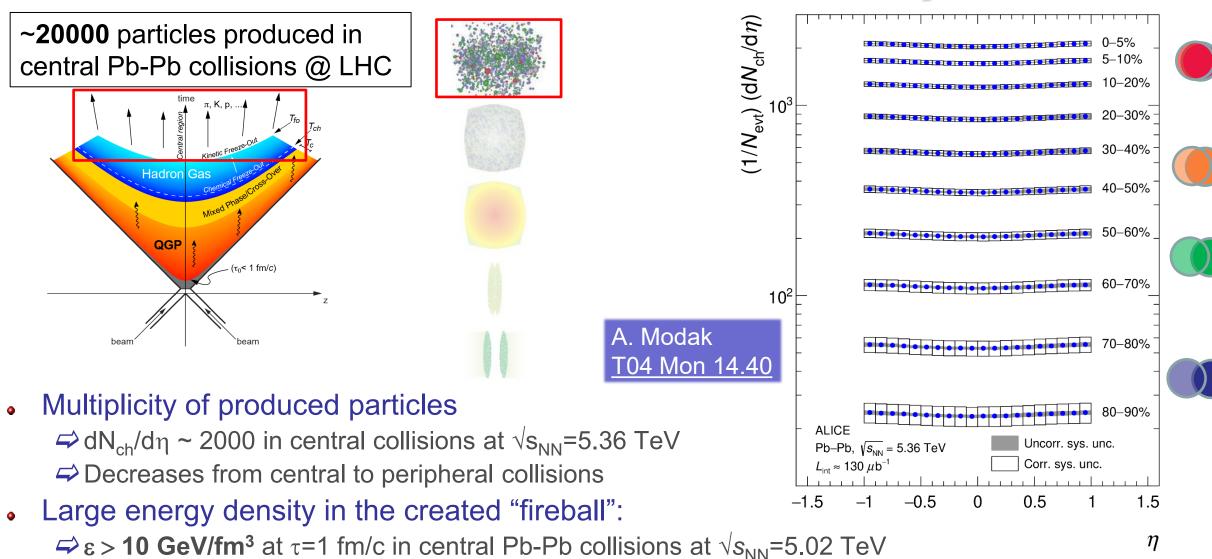
Initial state: tt in p-Pb collisions

- Top quark production measured from W+b decay \rightarrow sensitive probe of gluon nPDFs \Rightarrow Cover antishadowing and EMC regions at large Bjorken-*x* (3×10⁻³ < *x* < 0.5)
 - ✓ Complementary to EW bosons
- Measured cross sections described by NNLO calculations with different nPDF sets

 ⇒ Largest discrepancy for nNNPDF30, which does not include Run 2 LHC data in the tuning
- Data can further constrain nPDFs in the high Bjorken-*x* region



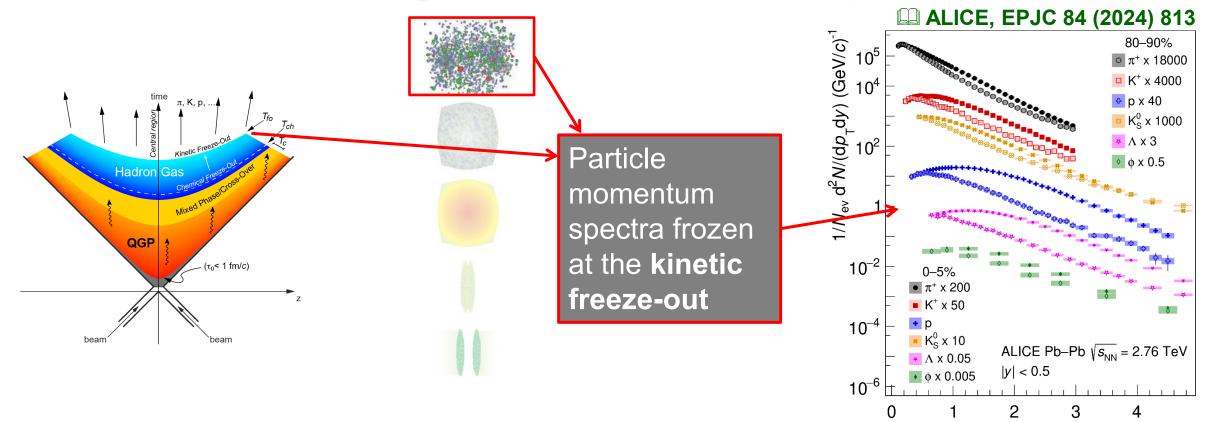
Final state: the "bulk" of soft particles



12

ALICE, arXiv:2504.02505

The flowing "bulk" of soft particles



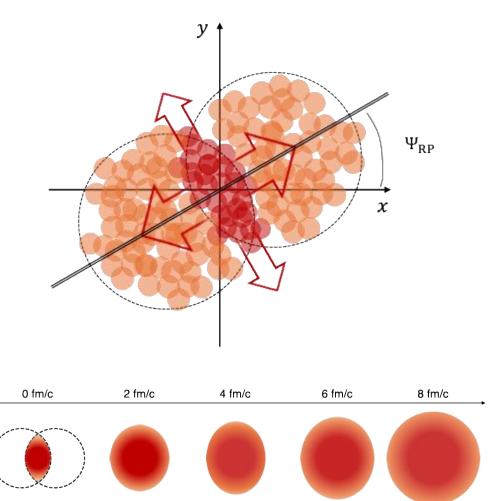
- Hardening of spectral shapes with increasing centrality and particle mass

 *p*_T
 i ∈ Effect of random thermal motion + collective expansion with common radial flow velocity
- Described by hydrodynamic expansion of the medium

 ⇒ Radial flow velocity β_T~0.5-0.6c at freeze-out temperature T_{fo}~100 MeV

*p*_{_} (GeV/*c*)

Anisotropic transverse flow



- Initial geometrical anisotropy in non-central collisions
 - Impact parameter selects preferred direction in the transverse plane
- Hydrodynamical evolution converts the initial geometrical anisotropy into an observable final-state particle momentum anisotropy

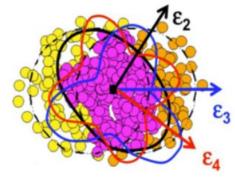
 \Rightarrow Characterized by anisotropic flow coefficients v_n :

$$\frac{dN}{d\varphi} = \frac{N_0}{2\pi} \left\{ 1 + 2\sum_{n=1}^{\infty} v_n(p_T) \cos[n(\varphi - \Psi_{RP})] \right\}$$

Fourier coefficient
$$v_2$$
 = elliptic flow

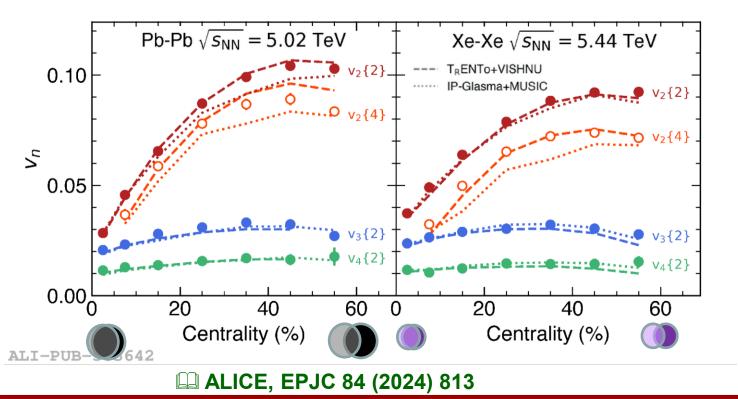
- \Rightarrow Dominant effect from initial geometry (v_2 > other harmonics) in non-central collisions
- Sensitive to the equation of state of the system in the early stages of the medium

Anisotropic flow: higher order harmonics



• Fluctuations and lumpiness of initial geometry (participant nucleons) and initial energy density profile give rise to non-zero values of higher harmonics (v₃, v₄, ...), if they are not damped

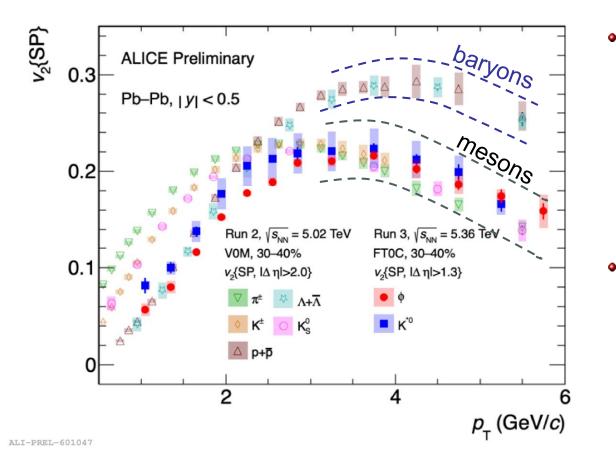
Connection between initial-state fluctuations and final-state particle v_n sensitive to "slow modes" in the hydrodynamics, which depend on the transport coefficients of the QGP (like shear and bulk viscosity)



- QGP produced in the collision is strongly coupled
 - Quarks and gluons form a collective medium that flows as a relativistic fluid with exceptionally low viscosity/entropy ratio

Open question: how does a strongly coupled liquid emerge from QCD which is asymptotically free at short distances?

Elliptic flow of light-flavour hadrons



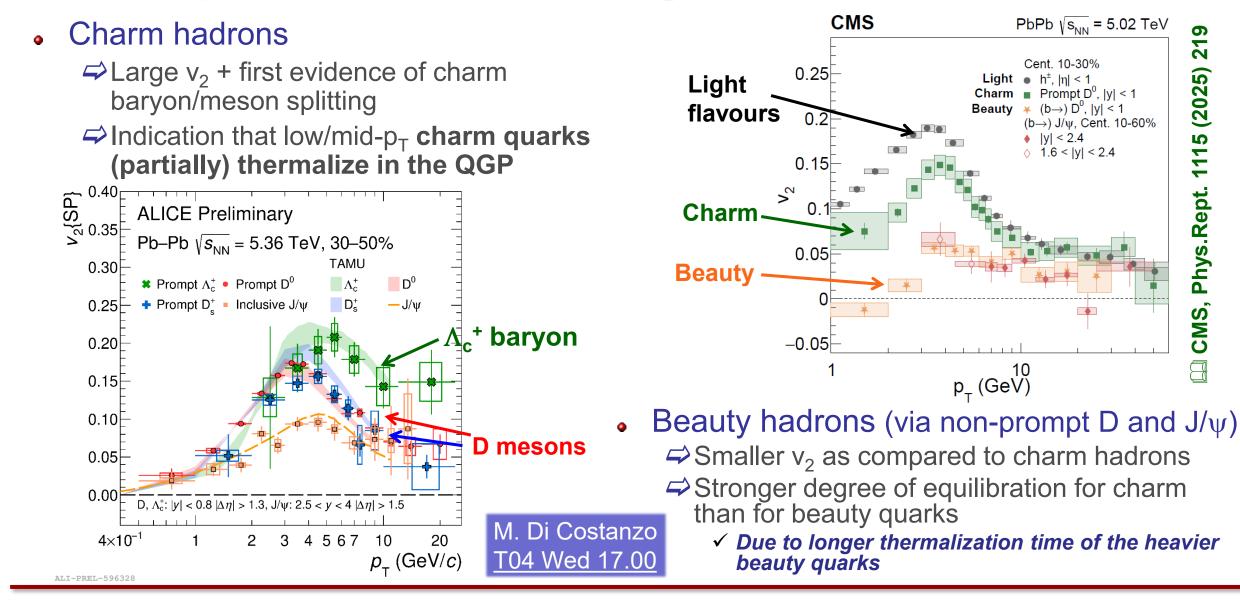
Mass ordering of v₂ at low p_T

- Originating from interplay between radial flow and anisotropic expansion of the fireball
- Mass ordering develops not only during the partonic evolution of the medium but also in the late hadronic re-scattering phase
- → Described by hydrodynamic models

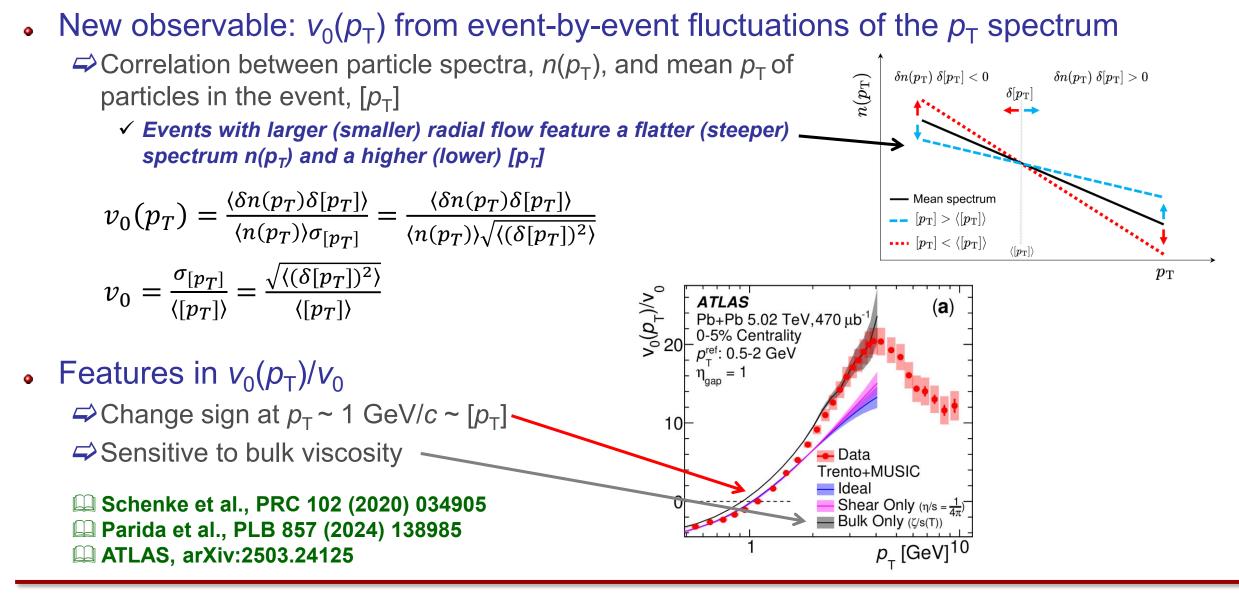
Baryon/meson grouping at intermediate p_T

- Anisotropic flow is driven by quark content rather than by mass in this momentum range
- Pattern expected in the case of hadron formation via quark recombination at the QGP hadronization
- Evidence of partonic collectivity

Elliptic flow of heavy-flavour hadrons



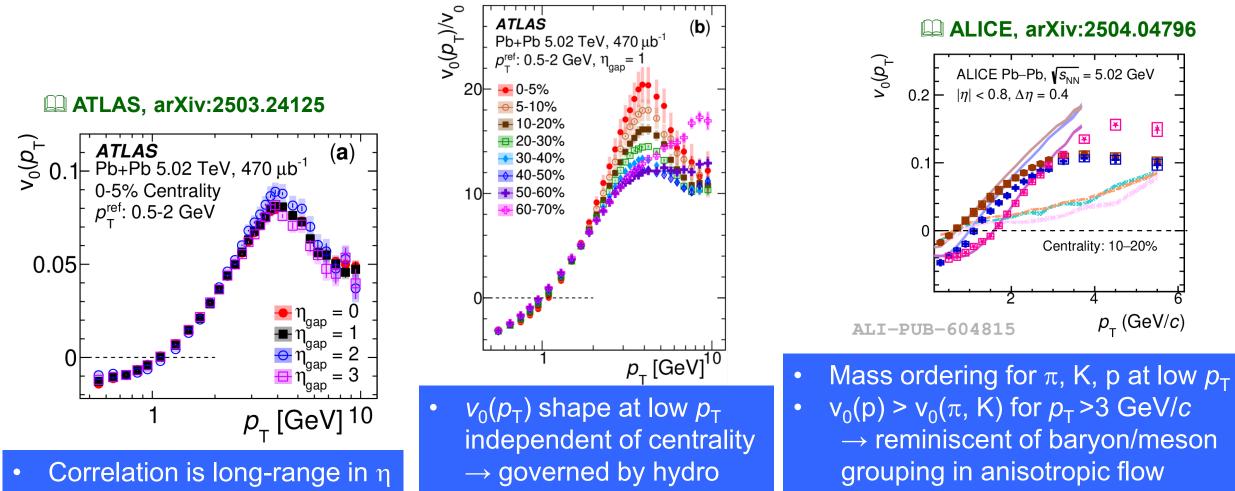
Back to radial flow



Back to radial flow

• New observable: $v_0(p_T)$ from event-by-event fluctuations of the p_T spectrum

→ Radial flow shows collective features

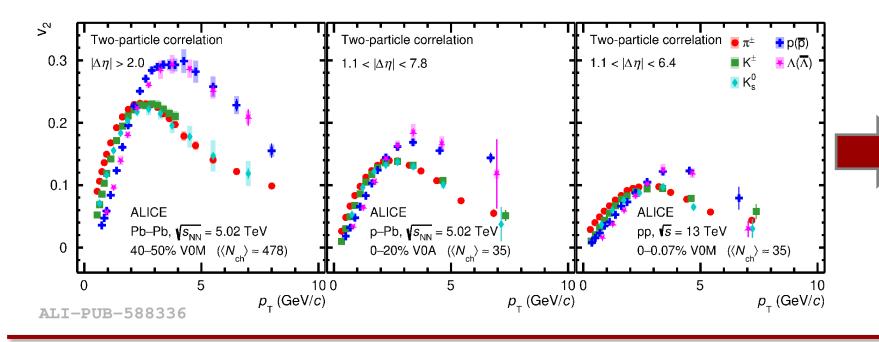


Collectivity in small collision systems

- Long-range angular correlations observed in high-multiplicity pp and p-Pb collisions
 - Similar features as those observed in Pb-Pb and interpreted as due to collective flow
 - \Rightarrow Clear mass ordering of the v₂ coefficients at low p_T
 - ✓ Significant evidence of radial flow in small collision systems

⇒ Distinctive grouping of anisotropic flow for baryons and mesons at intermediate p_T

✓ Anisotropic flow development at the parton level

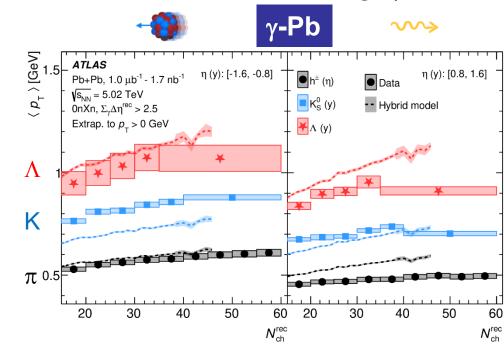


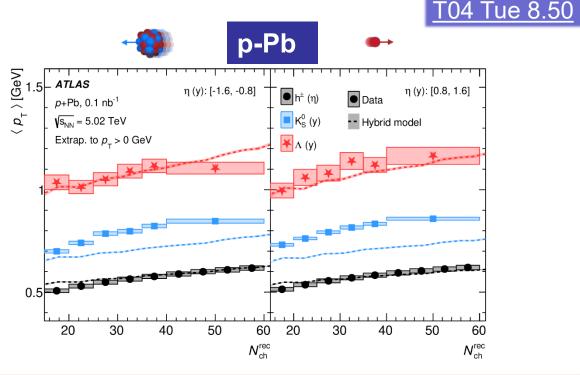
The system created in high-multiplicity pp and p-Pb collisions includes a stage with collectively flowing partons

ALICE, arXiv:2411.09323

Collectivity in smaller collision systems

- Photo-nuclear events show similar features as p-Pb collisions for observables sensitive to anisotropic and radial flow (+ baryon-to-meson ratios, strangeness enhancement)
 - \Rightarrow E.g.: distinct mass ordering in $\langle p_{\rm T} \rangle$ + larger $\langle p_{\rm T} \rangle$ in Pb-going direction in γ -Pb events
 - Hybrid model (assuming QGP formation and hydrodynamic expansion) captures qualitatively some features of the data
 - Supports the vector meson dominance picture: γ fluctuates into a vector meson \rightarrow a subset of events are, e.g., ρ -A collisions





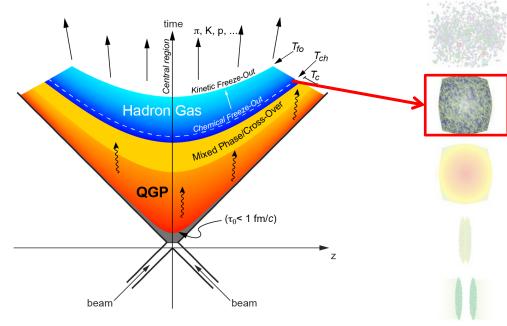
R. Schotter

S. Ragoni

T05 Wed 9.50

"Chemical" composition

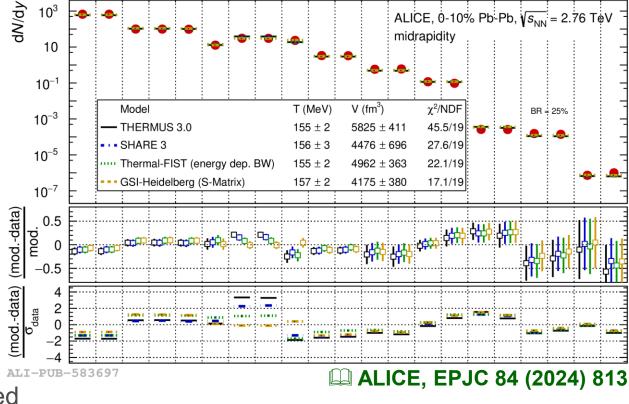
10³



At the chemical freeze-out

 \Rightarrow Inelastic collisions cease

- Abundances of different hadron species fixed
- Hadron yields (dominated by low- p_T particles) described by statistical/thermal models •
 - \Rightarrow Hadron abundances follow equilibrium populations \rightarrow thermal origin of particle production
 - A Hadron abundances depend on hadron masses (and spins), chemical potentials, and temperature
 - \Rightarrow Estimate temperature at the chemical freeze-out: T_{chem} =155 MeV

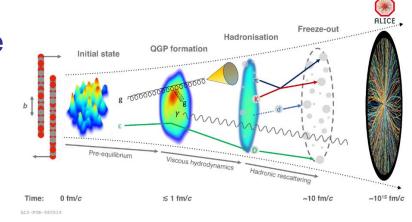


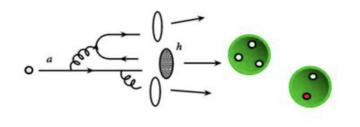
 \overline{d} ³He ³He ³ \overline{He} ³_AH ³ \overline{H} ⁴He ⁴He

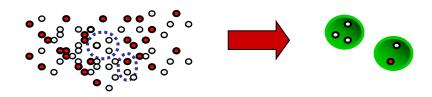
ALICE, 0-10% Pb-Pb, √s_{NN} = 2.76 TeV

Hadronization of the QGP

- Hadronization of the QGP at the (pseudo)critical temperature
 - Transition from a deconfined medium composed of quarks, antiquarks and gluons to color-neutral hadronic matter
 - Hadronic phase space populated following "maximal entropy"
- How does the hadronization occur at microscopic level?
- Non-perturbative process, requires phenomenological models
 - Parton fragmentation ("vacuum-like" dynamics)
 - ✓ Creation of $q\bar{q}$ and/or diquark pairs in the hadronization process
 - One parton fragments into many hadrons, each of them taking a fraction z of the parton momentum
 - ✓ Modelled via fragmentation functions or string fragmentation (PYTHIA)
 - - ✓ Phase space at the QGP hadronization filled with (thermalized) partons
 - \checkmark No need to create $q\overline{q}$ pairs via splitting / string breaking
 - ✓ Partons that are <u>"close" to each other in phase space</u> (position and momentum) can recombine into hadrons



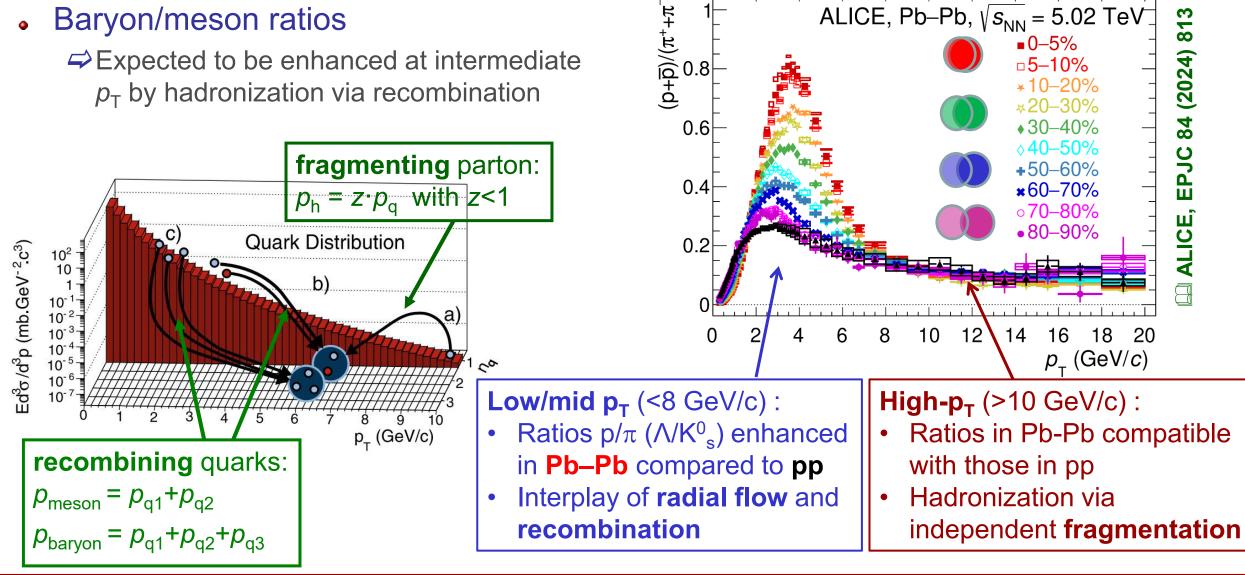




Hadronization of the QGP

Baryon/meson ratios •

 \Rightarrow Expected to be enhanced at intermediate $p_{\rm T}$ by hadronization via recombination



3

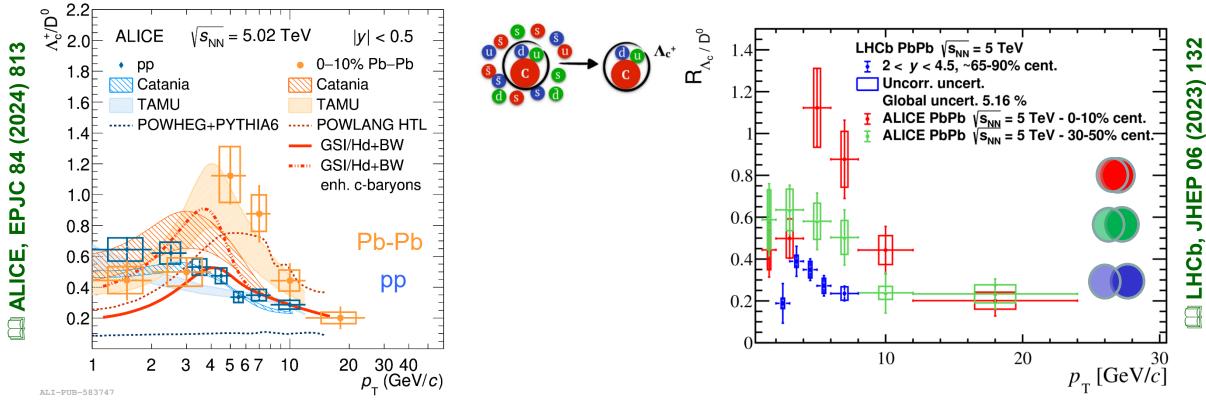
<u></u>

ALICE, Pb–Pb, $\sqrt{s_{NN}} = 5.02 \text{ TeV}^-$

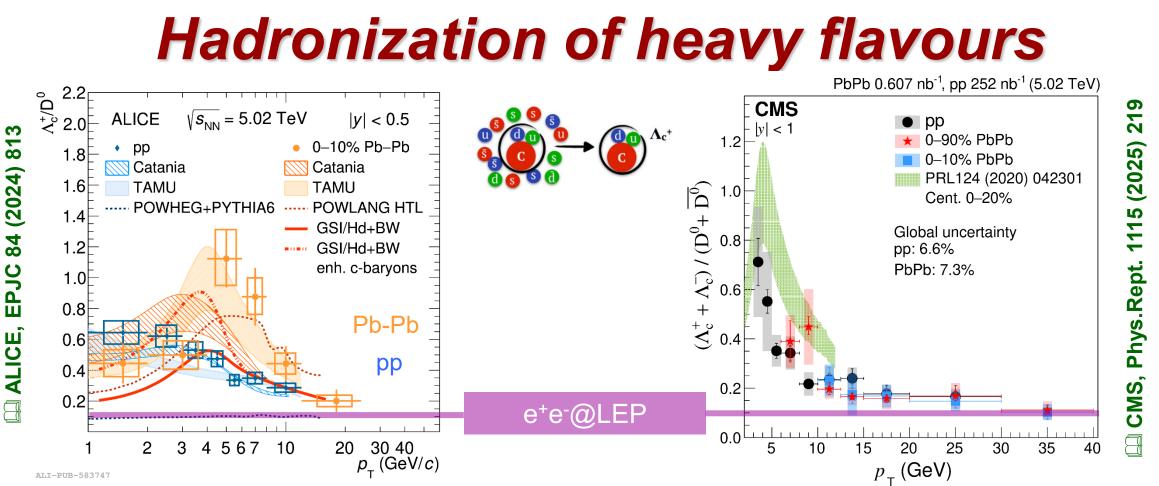
■0–5%

□5–10% ×10-20%

Hadronization of heavy flavours



- Λ_c/D^0 ratios at intermediate p_T (<10 GeV/c) enhanced in Pb-Pb compared to pp
 - ⇒ Enhancement increases from peripheral to central collisions
 - Described by models including hadronization via recombination with light quarks from the bulk



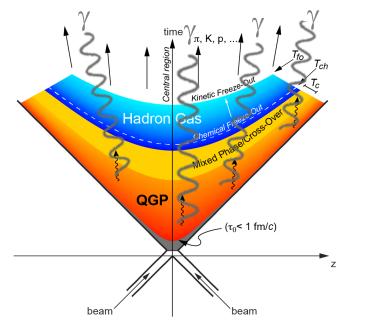
- Λ_c/D^0 ratio in pp collisions at the LHC substantially higher than in e⁺e⁻ collisions
 - Indicates a breakdown of the universality of charm quark fragmentation functions
 - ⇒ Captured in theoretical models with different approaches:
 - ✓ Color reconnections beyond-leading-color in PYTHIA8
 - ✓ Small QGP droplet + hadronization via recombination / statistical hadronization 105 Tue 18.10

J. Aichelin

T05 Wed 8.50

M. Karwoska

Electromagnetic probes

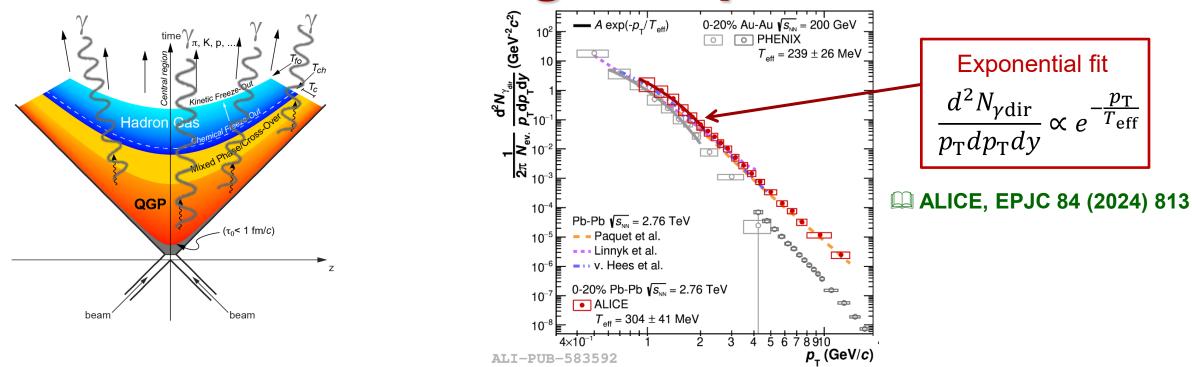


- Electromagnetic radiation emitted throughout the system evolution by a variety of sources in the form of:
 - ⇒ Real "direct" photons
 - ✓ Direct photons = do not originate from parton fragmentation nor hadronic decays
 - \Rightarrow Virtual photons, measurable via their internal conversion into e⁺e⁻ or $\mu^+\mu^-$ pairs



- ✓ Dilepton measurements at the LHC not yet sensitive to possible thermal signals 🥌
- Photons and dileptons escape the fireball unaffected, delivering information on the QGP conditions

Electromagnetic probes

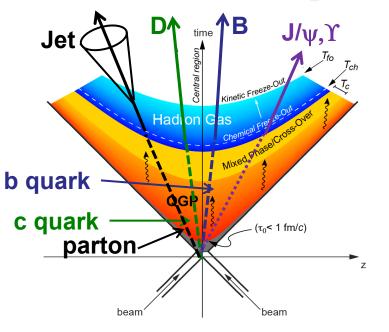


• Estimate an effective temperature of the fireball $T_{\rm eff}$ via exponential fit to direct photon $p_{\rm T}$ spectrum at low $p_{\rm T}$

 \Rightarrow Low- p_T region (<3 GeV/c) dominated by **thermal photons** emitted by the QGP during its evolution $\Rightarrow T_{eff} \sim 304$ MeV in central collisions at the LHC \rightarrow larger than the (pseudo)critical temperature

- Relating $T_{\rm eff}$ to the initial temperature of the fireball is challenging
 - Requires models that incorporate the QGP evolution and account for the radial expansion of the medium, which causes a blue-shift of the emitted photons

Hard probes of the QGP medium

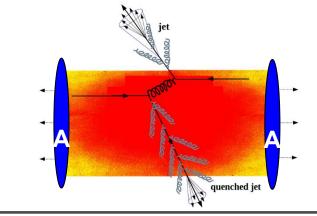


- Produced at the very early stages of the collision in partonic scattering processes with large momentum transfer
 Produced out-of-equilibrium
- Traverse the hot and dense medium interacting with its constituents

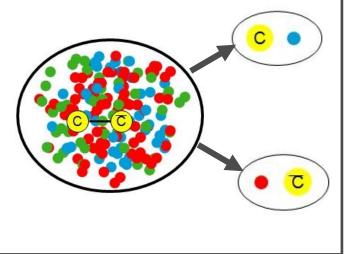
⇒ Unique probes of the properties of the QGP

Jet quenching:

the hard-scattered parton loses energy while crossing the QGP + modification of parton shower and jet properties



Quarkonium dissociation in the QGP



Nuclear modification factor

• Hard processes in nuclear collisions

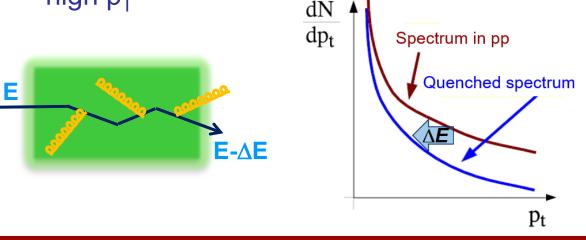
 \rightleftharpoons Production scales with the number of nucleon-nucleon collisions, $N_{\rm coll}$

Nuclear modification factor

 $R_{\rm AA}(p_{\rm T}) = \frac{1}{\langle N_{\rm coll} \rangle} \frac{dN_{\rm AA}/dp_{\rm T}}{dN_{\rm pp}/dp_{\rm T}} \sim \frac{QCD \text{ medium}}{QCD \text{ vacuum}}$

 \Rightarrow If **no nuclear effects** are present: $R_{AA} = 1$

• Interactions in the QGP cause **energy loss** leading to suppression of yield (R_{AA} <1) at high p_T



Nuclear modification factor

• Hard processes in nuclear collisions

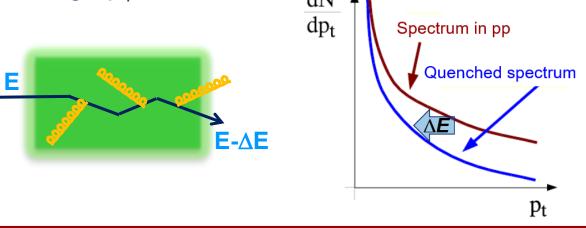
→ Production scales with the number of nucleonnucleon collisions, N_{coll}

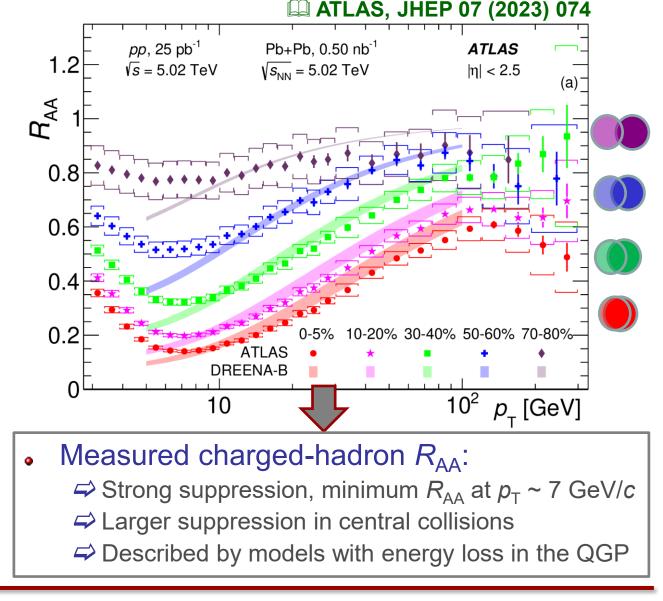
Nuclear modification factor

 $R_{\rm AA}(p_{\rm T}) = \frac{1}{\langle N_{\rm coll} \rangle} \frac{dN_{\rm AA}/dp_{\rm T}}{dN_{\rm pp}/dp_{\rm T}} \sim \frac{QCD \text{ medium}}{QCD \text{ vacuum}}$

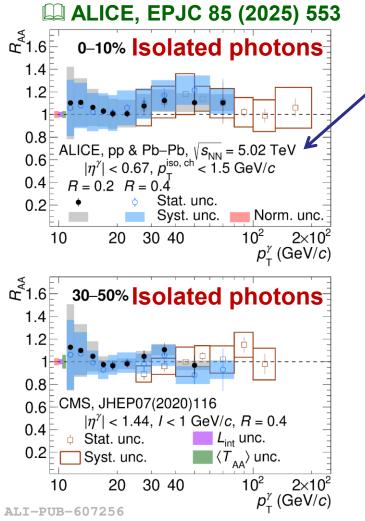
 \Rightarrow If **no nuclear effects** are present: $R_{AA} = 1$

• Interactions in the QGP cause **energy loss** leading to suppression of yield (R_{AA} <1) at high p_T





Control experiments



- Medium-blind probes (γ, Z⁰, W)
 - Production of particles w/o color charge not modified by the QGP
 - → Test binary scaling

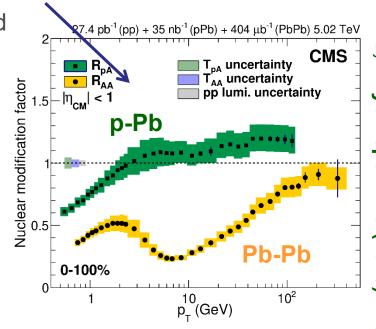
•

0

 \Rightarrow R_{AA} of **isolated (prompt) photons** compatible with 1

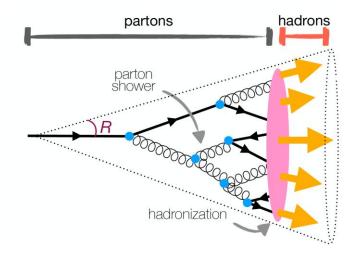
Charged hadrons in p-Pb minimum-bias collisions

- → Large-size QGP phase not expected
- Test cold-nuclear matter effects
- \Rightarrow R_{pPb}~1 for p_T > 3 GeV/c



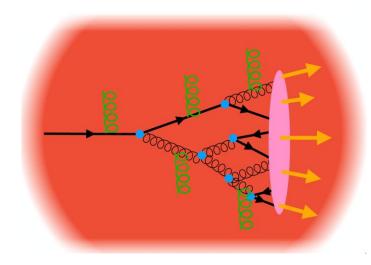
 \rightarrow suppression of high- p_{T} hadrons is due to energy loss of colored partons in the QGP

Probing the QGP with jets



• Jets: "in-vacuum" fragmentation

Collimated sprays of hadrons resulting from fragmentation and subsequent hadronization of energetic partons

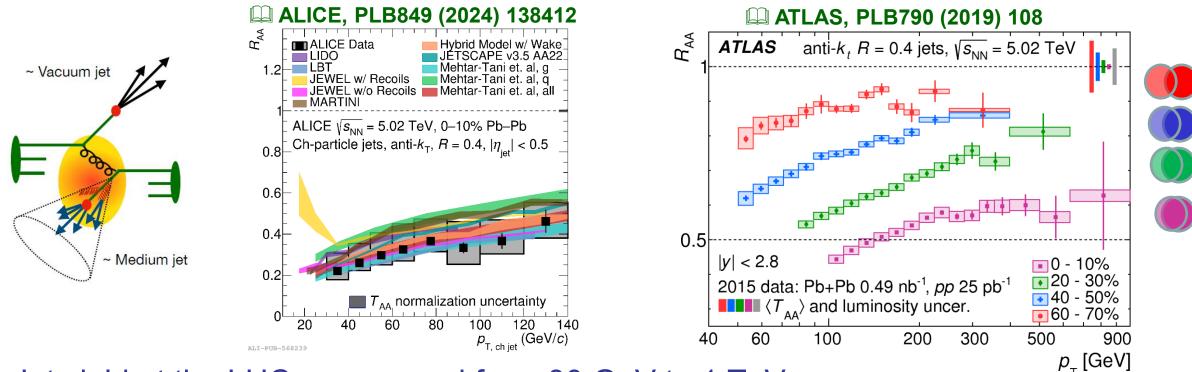


• Jets in A-A collisions

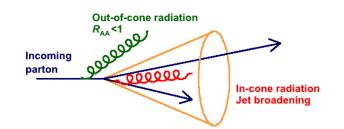
- Parton showers propagate through the QGP: interactions with medium constituents
- \Rightarrow **Jet quenching** $\rightarrow QGP alters the energy and structure of jets passing through it$
- ➡ Medium response → motions induced by the jet energy deposited in the QGP

\rightarrow probe the QGP at short-distance scales





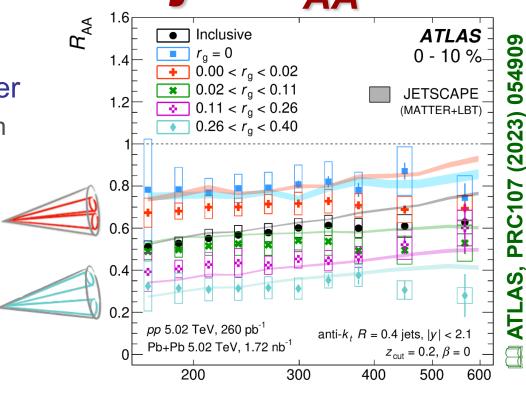
- Jet yield at the LHC suppressed from 30 GeV to 1 TeV
 - Hard-scattered quarks and gluons evolve as parton showers that propagate through the hot and dense medium
 - \Rightarrow Part of energy of parton shower transferred outside the jet cone through interactions with QGP (\rightarrow soft radiation at large angles)
- Suppression decreases from central to peripheral collisions



Substructure dependent jet R_{AA}

- *R*_{AA} vs. angular separation (*r*_g) of the sub-jets defined by the hardest splitting in the parton shower
 ⇒ Hard splitting tagged with soft-drop grooming algorithm
 ⇒ Jets with *r*_g=0 failed the soft-drop requirement and considered as single-prong jets
- Jet quenching depends on structure of parton shower

⇒ Wide jets are more suppressed than narrow jets



 $p_{_{
m T}}^{
m jet}$ [GeV]

• r_{g} -dependent suppression explained in models as arising from coherence effects \rightarrow wider jets are "more resolved" by the medium

 \Rightarrow Large r_g : jet constituents resolved by the medium \rightarrow radiation as multiple color charges \Rightarrow Small r_g : jet constituents unresolved \rightarrow sub-jets radiate coherently as a single-color charge

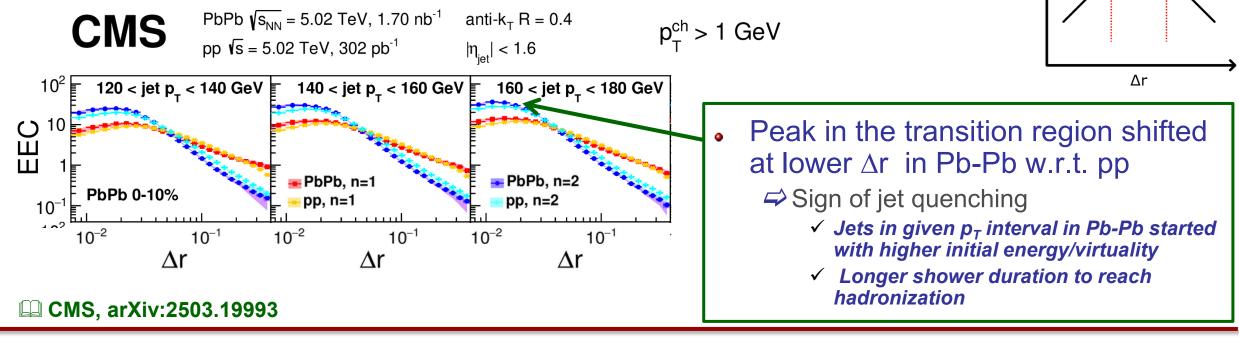
Energy-energy correlators

• EEC = weighted distribution of angular separation Δr of all possible particle pairs within a jet cone

$$\operatorname{EEC}(\Delta r) = \frac{1}{W_{\text{pairs}}} \frac{1}{\delta r} \sum_{\operatorname{jets} \in [p_{\mathrm{T},1}, p_{\mathrm{T},2}]} \sum_{\operatorname{pairs} \in [\Delta r_a, \Delta r_b]} \left(p_{\mathrm{T},i} \, p_{\mathrm{T},j} \right)^n$$

⇒time scales in jet evolution imprinted in different angular scales of the EEC

✓ Early times correspond to large angles; late times to small angles

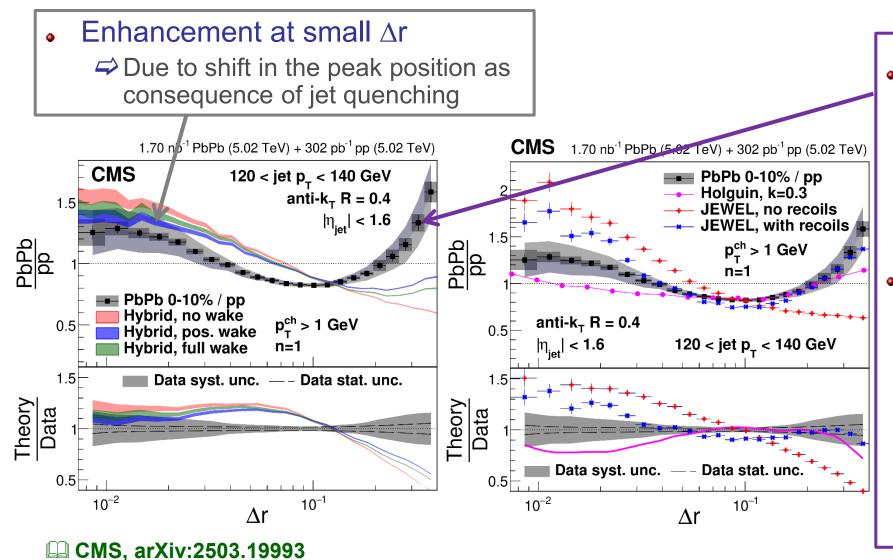


pQCD

npQCD

U U U

EEC ratios: Pb-Pb/pp

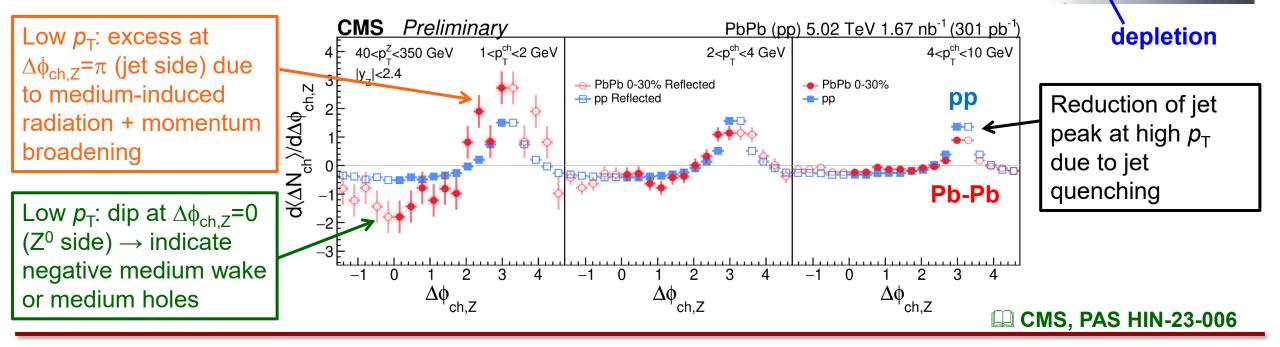


Enhancement at large Δr

- Sensitive to mediuminduced modifications
- Mainly due to soft particles at the periphery of jet cone
- Qualitatively captured by models including
 - Medium response to jets
 ✓ Jet wake / recoils
 - ➡ Color coherence effects
 - Only gluons emitted above a critical angle can independently emit more gluons

Medium response via Z⁰+jet events

- Z⁰ bosons and jets from the same hard scattering
 Z⁰ do not interact strongly in QGP → access to hard-scattered parton momentum
- Explore medium response by measuring yield of Z⁰-tagged charged hadrons vs. azimuthal angle relative to Z⁰ (Δφ_{ch,Z})
 ⇒ Modified in central Pb-Pb as compared to pp collisions
 ⇒ Data better described by models that include medium recoil effects
 ⇒ First evidence of medium response effects caused by a hard probe



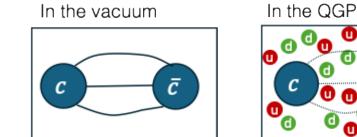
Wake front

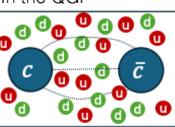
 $\Delta \phi_{ch,Z}$

Quarkonium: in-medium dissociation

- Dissociation of quarkonia in the QGP •
 - \Rightarrow High density of free color charges in the QGP leads to a screening of the QCD force that binds the quarkonium states Quarkonium production suppressed in A-A collisions
 ■
- Sequential pattern expected from the rich spectroscopic • structure of quarkonia
 - Different quarkonium states dissociate in the QGP at different temperatures, depending on their binding energy
 - ✓ The more strongly bound the state is, the hotter must be the medium to dissociate it

Sequential suppression as a QGP thermometer





Matsui, Satz, PLB178 (1986) 416

	J/ψ	Xc	ψ(2S)
Mass (GeV)	3.10	3.53	3.68
E _B (GeV)	0.64	0.20	0.05

	Y(1S)	Y(2S)	Y(3S)	
Mass (GeV)	9.46	10.02	10.36	
E _B (GeV)	1.10	0.54	0.20	

Quarkonium: in-medium dissociation

- Dissociation of quarkonia in the QGP
 - High density of free color charges in the QGP leads to a screening of the QCD force that binds the quarkonium states
 Quarkonium production **suppressed** in A-A collisions
- Sequential pattern expected from the rich spectroscopic structure of quarkonia
 - Different quarkonium states dissociate in the QGP at different temperatures, depending on their binding energy
 - The more strongly bound the state is, the hotter must be the medium to dissociate it

Sequential suppression as a QGP thermometer

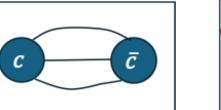
- Several effects lead to a more complex situation:
 - Cold nuclear matter (CNM) effects affecting quarkonia

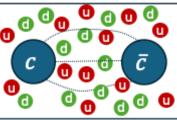
✓ Nuclear PDFs, coherent energy loss in CNM

- Break-up by comoving hadrons in the hadronic phase
- ⇒ **Feed-down** from higher quarkonium states
- Production via recombination of uncorrelated quark pairs originated from different hard scattering processes



In the QGP





📖 Matsui, Satz, PLB178 (1986) 416

	J/ψ	Xc	ψ(2S)
Mass (GeV)	3.10	3.53	3.68
E _B (GeV)	0.64	0.20	0.05

	Y(1S)	Y(2S)	Y(3S)
Mass (GeV)	9.46	10.02	10.36
E _B (GeV)	1.10	0.54	0.20

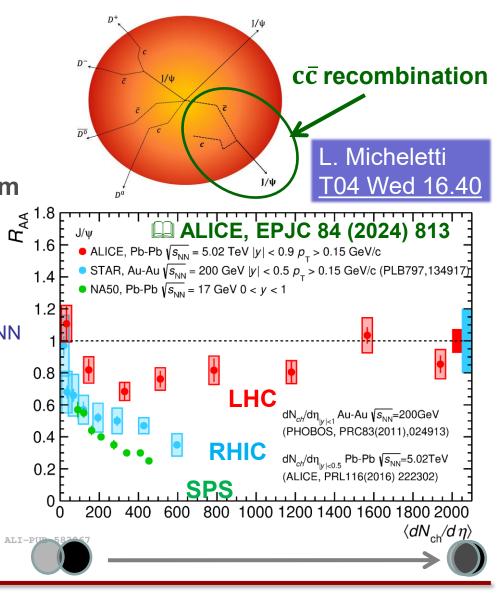
Low-p_T charmonium: recombination

- If QGP is formed, charmonium production can occur also via cc recombination during the QGP expansion or at the phase boundary
 - Charmonia from recombination mainly produced with low momentum
 - Recombination expected to be **negligible for bottomonium** (b quarks are rare)

Braun-Munzinger, Stachel, PLB 490 (2000) 196
 Thews et al., PRC 63 (2001) 054905

- Tested by measuring J/ψ yield in A-A collisions vs. √s_{NN}
 ⇒ Charm quark production cross section increases with √s_{NN}
 ⇒ Probability of cc̄ recombination increases with √s_{NN}
- J/ ψ yield enhanced in A-A collisions at higher $\sqrt{s_{NN}}$ \Rightarrow Strong evidence of J/ ψ formation via $c\bar{c}$ recombination

 \rightarrow implies the presence of a deconfined phase



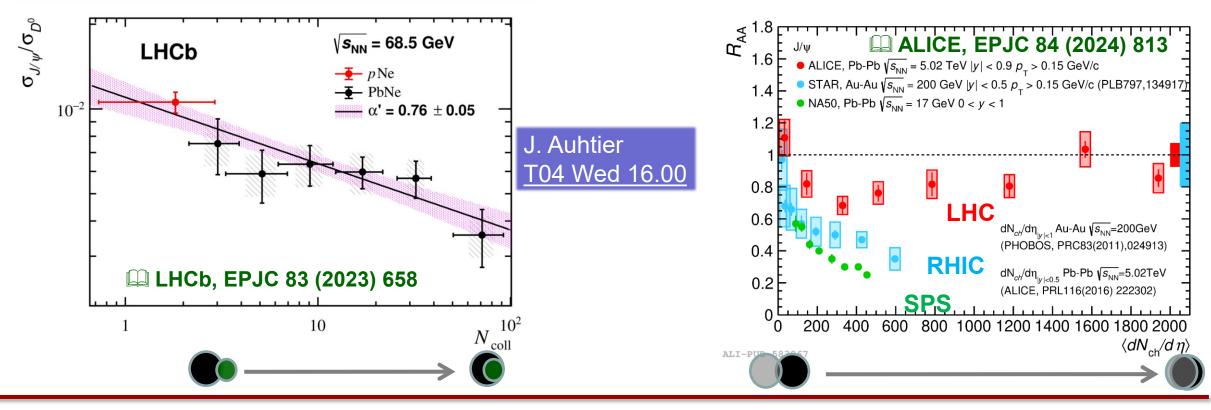
Enriching the J/ψ picture with SMOG@LHCb

• J/ ψ / D⁰ ratio in p-Ne and Pb-Ne using fixed-target configuration (SMOG) of LHCb

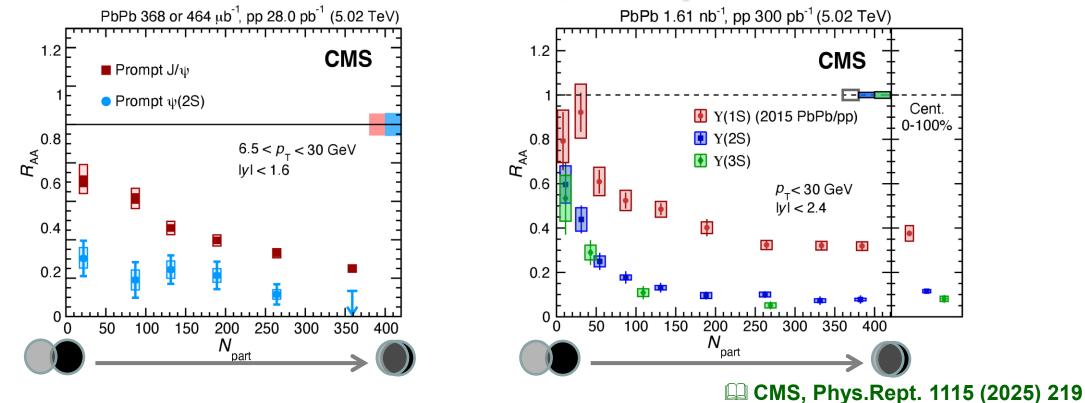
 $\Rightarrow \sqrt{s_{NN}} = 68.5 \text{ GeV}$, in between SPS and RHIC

 \Rightarrow J/ ψ yield suppressed relative to D⁰ \rightarrow J/ ψ affected by additional nuclear effects compared to D⁰

Continuous trend from p-Ne to central Pb-Ne collisions within current uncertainties



Bottomonium and high-p_T charmonium



• High- p_T charmonia: $R_{AA}^{\psi(2S)} < R_{AA}^{J/\psi}$. Bottomonia: $R_{AA}^{\Upsilon(3S)} < R_{AA}^{\Upsilon(2S)} < R_{AA}^{\Upsilon(1S)}$

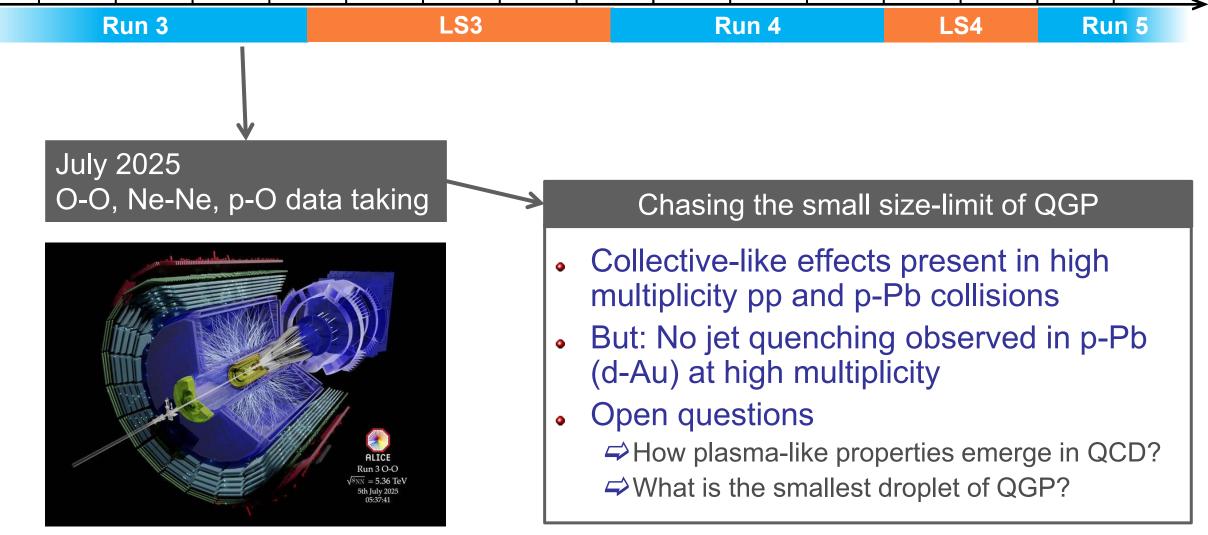
 \Rightarrow The more fragile states are the most suppressed

• Strong indication of sequential suppression effects, ordered by binding energy

For a conclusive interpretation of quarkonia as a QGP thermometer, one needs to account for the feed-down from decays of excited S- and P-wave quarkonium states

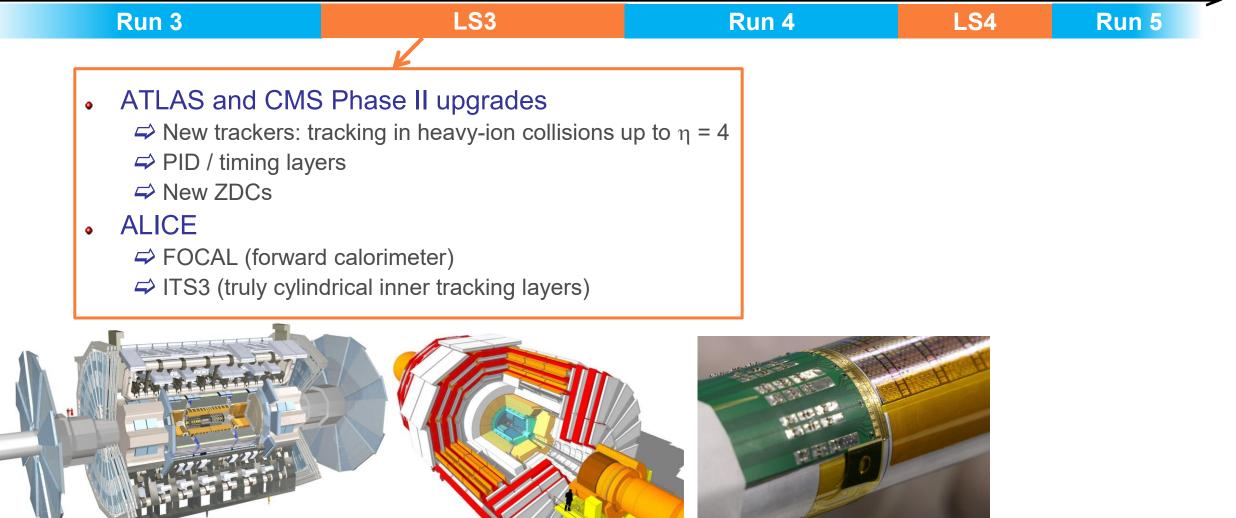
LHC: what's next?

2022 2023 2024 2025 2026 2027 2028 2029 2030 2031 2032 2033 2034 2035 2036 2037



LHC: what's next?

2022 2023 2024 2025 2026 2027 2028 2029 2030 2031 2032 2033 2034 2035 2036 2037

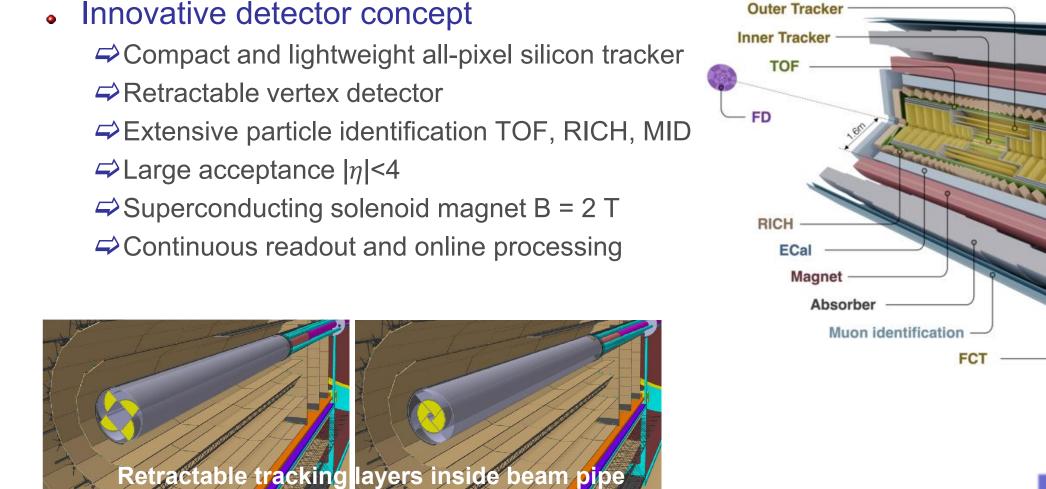


LHC: what's next?

2022 2023 2024 2025 2026 2027 2028 2029 2030 2031 2032 2033 2034 2035 2036 2037

Run 3	LS3	Run 4	LS4	Run 5
		 ALICE 	3	
		🖨 Bra	and new detect	or
		🖨 Exc	cellent pointing	resolution
		🖨 Lar	ge η coverage	,
		 LHCb 	upgrade II	
		🖨 Ful	l centrality rang	ge in Pb-Pb
		Outer Tra Inner Tra TOF FD	acker	ALICE 3
		RICH		

ALICE 3

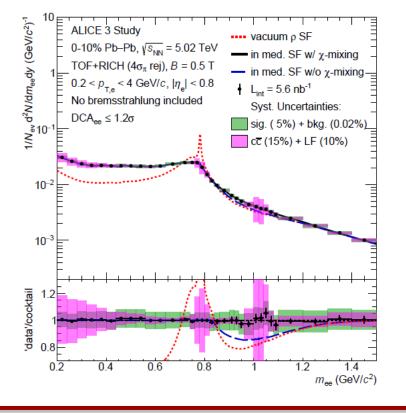


ALICE 3

ALICE 3: new physics opportunities

Low-mass dileptons

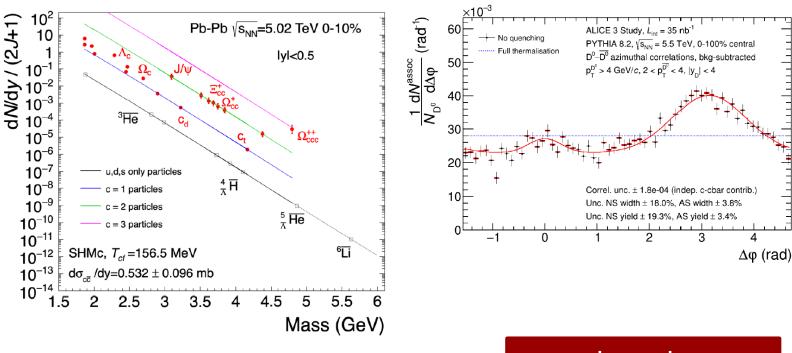
 Access ρ/a₁ mixing to study chiral symmetry restoration
 Sensitive to QGP temperature



- Multi-charmed baryons
 - ➡ Ultimate probe of charm deconfinement and thermalization

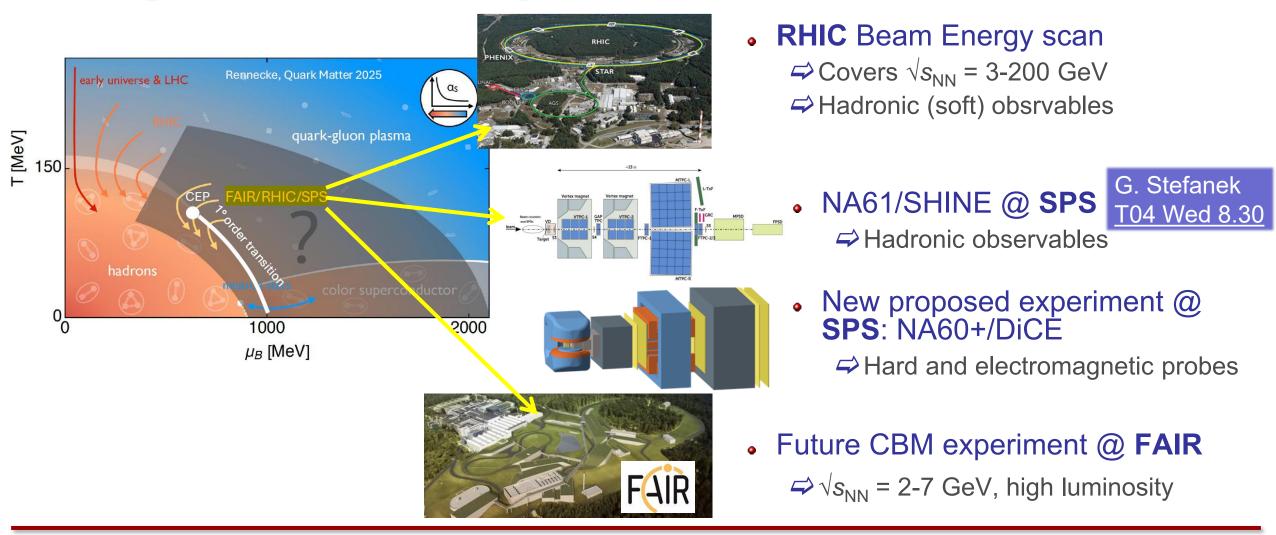
• $D - \overline{D}$ correlations

Measure momentum broadening in QGP



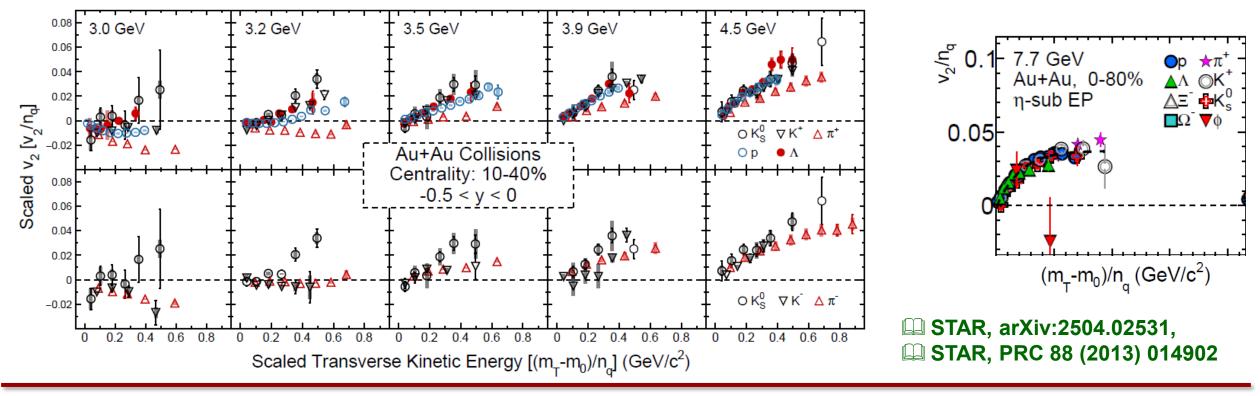
ALICE3 Lol, arXiv:2211.02491

Heavy-ion collisions at low $\sqrt{s_{NN}}$ bridge between early Universe and Neutron Stars



Onset of partonic collectivity at low \sqrt{s}

- Scaling of v_2 with number of constituent quarks (NCQ) expected to reflect the effective degrees of freedom of the medium
- Results from beam energy scan at RHIC:
 - ⇒ NCQ scaling completely broken for $\sqrt{s_{NN}} \le 3.2$ GeV
 - Gradual onset of NCQ scaling at higher energies
 - \Rightarrow Dominance of partonic interactions for $\sqrt{s_{NN}} > 4.5$ GeV \rightarrow emergence of partonic collectivity



Search for the Critical Point

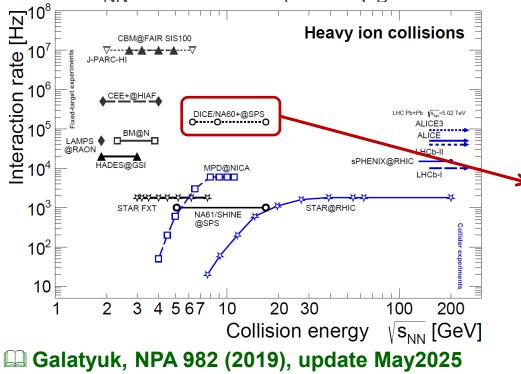
Event-by-event fluctuations of net-proton (net-charge) number N critical point? crossover **QGP** Sensitive to correlation length of the system that diverges at the critical point ∕_{St} \Rightarrow Quantified via cumulants $\kappa_1 = \langle N \rangle$, $\kappa_2 = \langle \delta N^2 \rangle$, $\kappa_3 = \langle \delta N^3 \rangle$, $\kappa_4 = \langle \delta N^4 \rangle - 3\kappa_2^2$, with $\delta N = N - \langle N \rangle$ Hadron gas \Rightarrow Search for **non-monotonic dependence on** $\sqrt{s_{NN}}$ Nuclei The minimum in net-1.5 h⁺-h⁻ No clear critical Data (0-5%) STAR Data (70-80%) 10 ▼ 0-1% Ar+Sc (NA61/SHINE) proton κ_4/κ_2 vs. $\sqrt{s_{NN}}$ Υh point signal in • p+p (NA61/SHINE) seen at RHIC is a net-charge Net-proton characteristic K4/K2 cumulants from feature of the 0.5 NA61/SHINE @ proposed signature SPS of the critical point 0 HRG CE LQCD -2 κ_{2} UrQMD (0-5%) 10 15 200 20 50 100 5 10 The jury is still out √s_{NN} (GeV) Collision Energy √s_{NN} (GeV) **NA61**, arXiv:2503.22484 STAR, arXiv:2504.00817

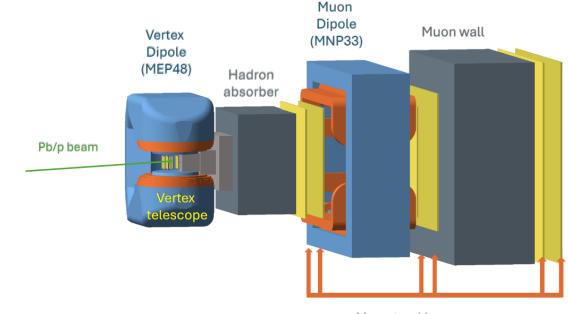
μΒ

Future heavy-ions @ SPS: NA60+/DiCE

New proposed experiment at SPS: NA60+/DiCE

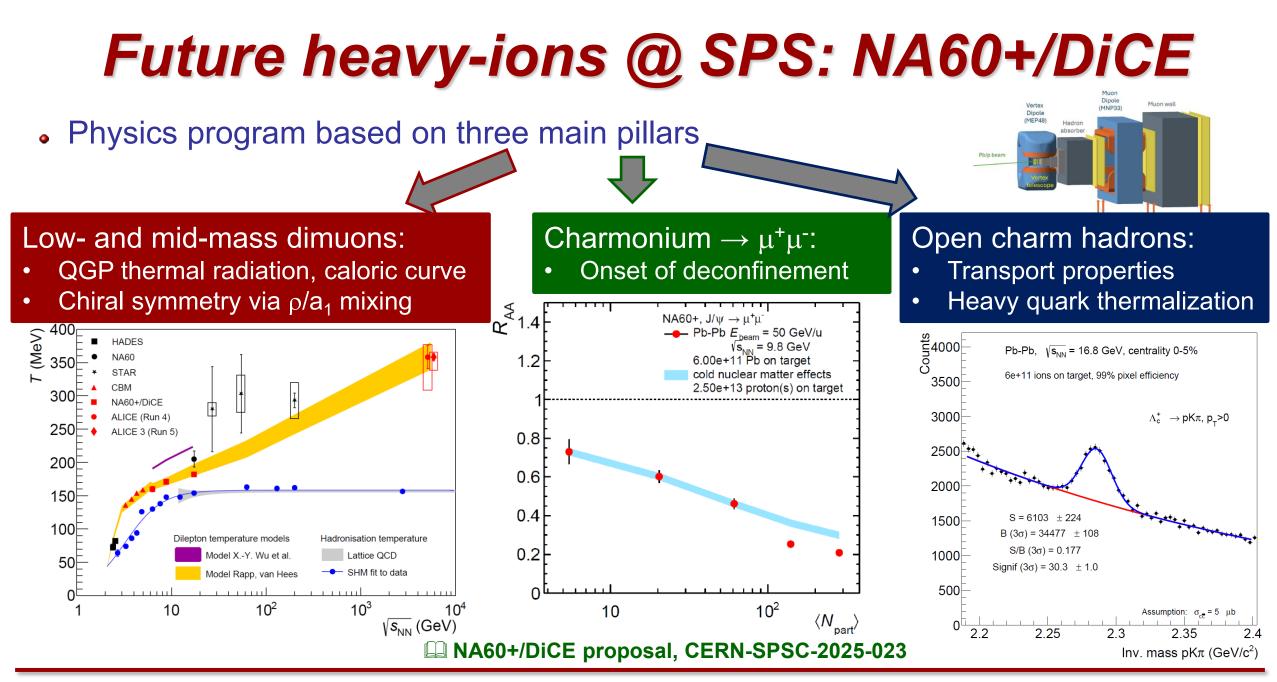
Study hard and electromagnetic probes of the QGP with a beam energy scan in the range √s_{NN} = 6-17 GeV (200 < µ_B < 450 MeV)</p>





Muon tracking

- Study rare probes thanks to large integrated luminosity
- Complementarity with experiments accessing:
 - Different (hadronic) observables in the same energy range (RHIC BES, NA61/SHINE)
 - Similar observables in a lower energy range (CBM at FAIR)



Conclusions and prospects

 Important progresses in the understanding of QCD at extreme conditions of high temperature and energy density

➡ Wealth of results from LHC and RHIC experiments in the last years

 Open questions being actively investigated on the experimental and theoretical sides

⇒ More will come in the next years

Multi-messengers from the QGP will be studied with multiple experiments at different energies in different facilities

 \Rightarrow Rich physics program at the LHC in Run 3, 4 and 5:

- ✓ Major upgrades of ATLAS and CMS in Long Shutdown 3
- ✓ ALICE3 and upgraded LHCb after Long Shutdown 4
- ⇒NA60+/DiCE for a beam energy scan at the **SPS**
- \Rightarrow CBM at **FAIR** to study the very-high- μ_B region
- **EIC** to provide important information about the initial state



Thank you for the attention!

Backup

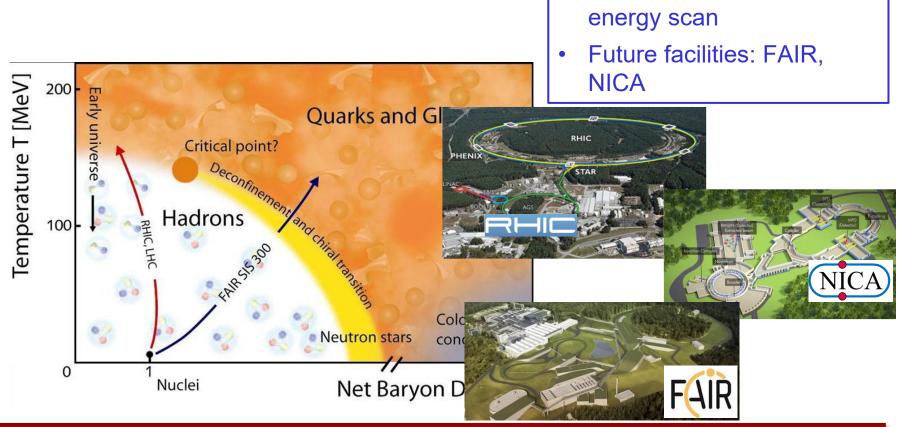
Exploring the QCD phase diagram

• Experimentally, heavy-ion collisions at different energies allow us to explore different regions of the QCD phase diagram

High energy frontier: LHC, RHIC top energy

- All the four main LHC experiments have a heavy ion program
- sPHENIX experiment taking data at RHIC

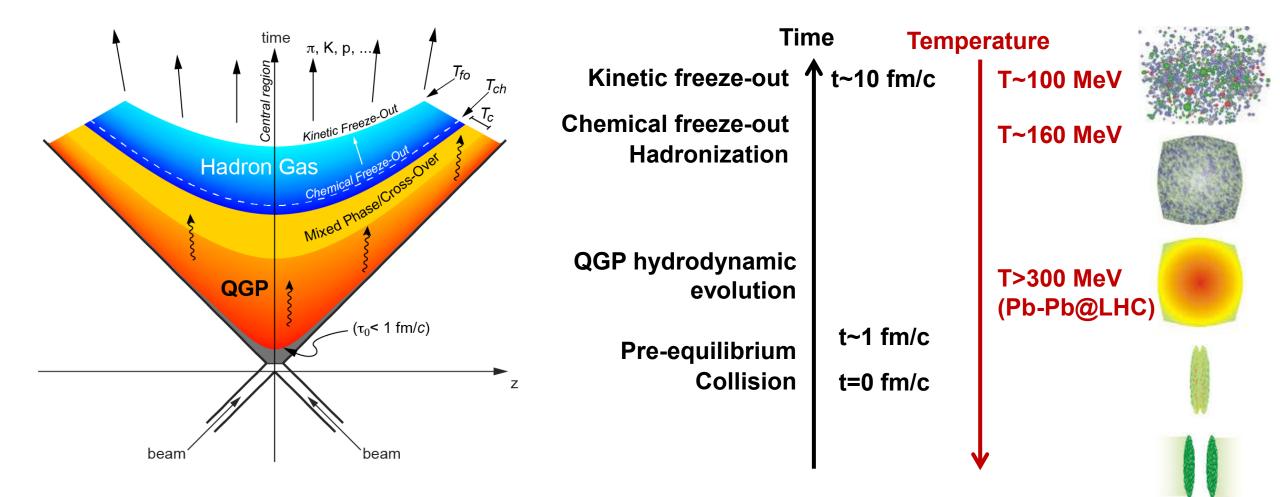




Low energy frontier:

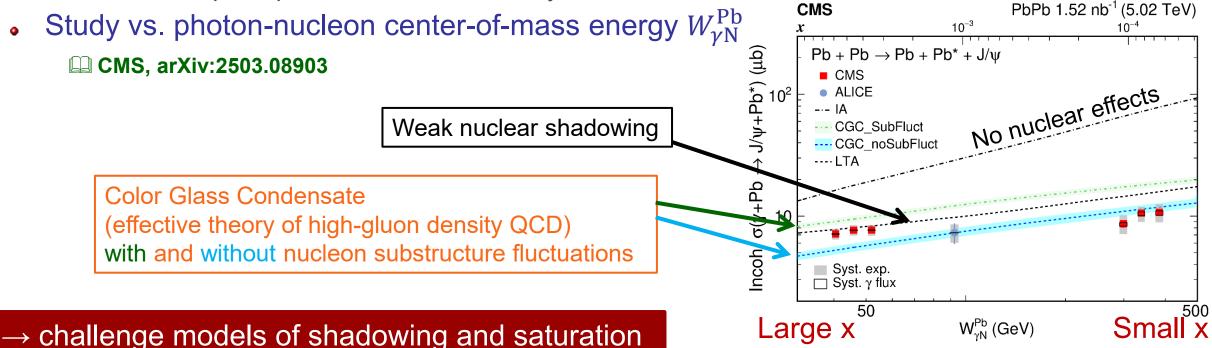
RHIC and SPS beam

Space-time evolution of A-A collisions



UPC: incoherent J/*y* **production**

- Coherent photoproduction: the photon couples to the entire nucleus
- Incoherent photoproduction: the photon couples to individual nucleons
 - Sensitive to localized, fluctuating gluonic hotspots
 - ➡ Often accompanied by nuclear breakup
- Test shadowing / saturation models
 - Nuclear shadowing and gluon saturation can influence coherent and incoherent photoproduction in distinct ways



Incoherent

Coherent

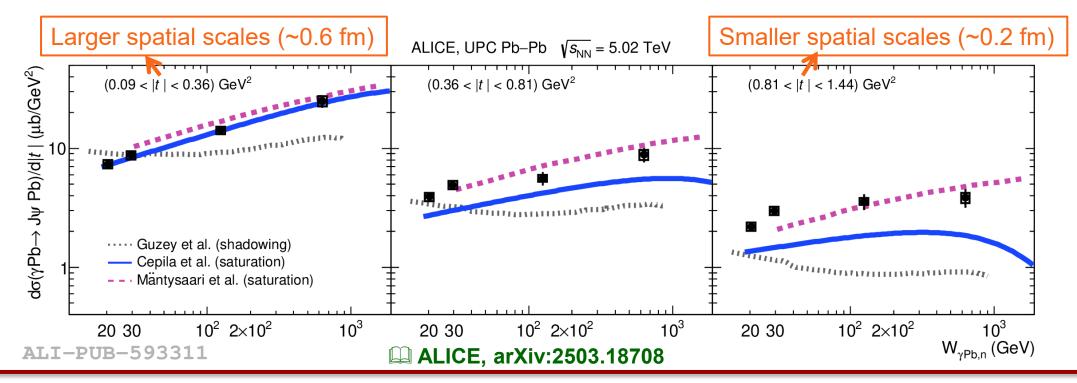
UPC: incoherent J/*y* **production**

• Energy dependence of incoherent J/ ψ photonuclear production in bins of Mandelstam t

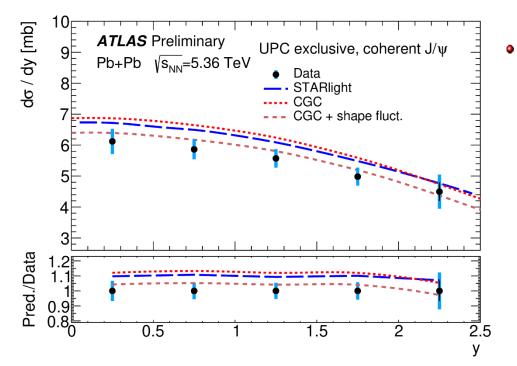
 \Rightarrow From measurements as a function of y and $p_T \rightarrow \text{extract } W_{\gamma Pb}^2 = M_{J/\psi} \sqrt{s_{NN}} e^{\pm y}$ and $|t| = p_T^2$

- Different ranges in Mandelstam-t sensitive to the dynamics of the gluon field at different spatial size scales
- Rate of growth of cross section with increasing energy suppressed at large |t|

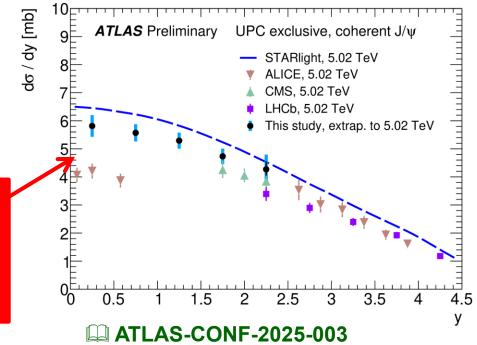
⇒ Data favour saturation-based models over shadowing models



UPC: coherent exclusive J/*\varpsilon* **production**



- Sensitive to nuclear gluon dynamics at low Bjorken-x
 - Test shadowing / saturation models
 - Best description of the data by the color glass condensate (CGC) approach including the effect of nucleon shape fluctuations



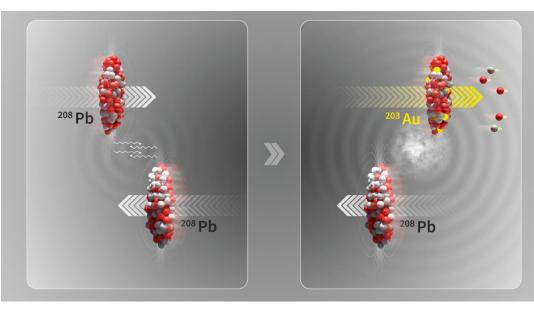
Tension between ALICE and ATLAS results:
 Due to production of additional particle pairs accompanying the J/ψ and violating the exclusivity requirements?

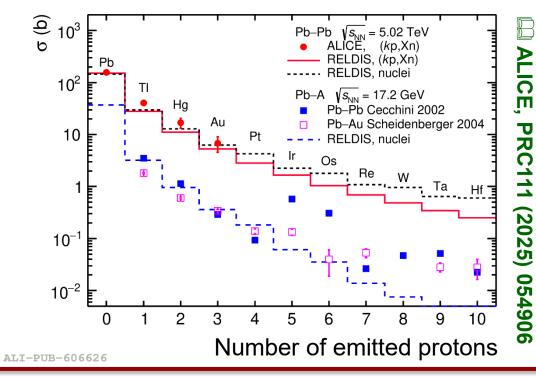
Parenthesis: alchemy at the LHC

- Electromagnetic dissociation processes in ultraperipheral collisions
 - ➡ Dominant de-excitation via neutron emission

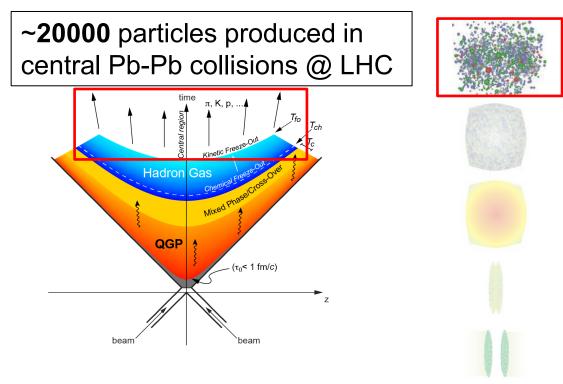
⇒ When high-energy photons are absorbed by the Pb nucleus, also some protons are emitted

- ➡ Detected in Zero Degree Calorimeters
- The emission of 3 protons corresponds to the production of Au nuclei
 ⇒ About 9x10¹⁰ nuclei (~ 2.9 x 10⁻¹¹ g) of Au produced in Run-2 at the LHC
- Lead transmuted into gold by light





Final state: the "bulk" of soft particles



• Multiplicity of produced particles depends on:

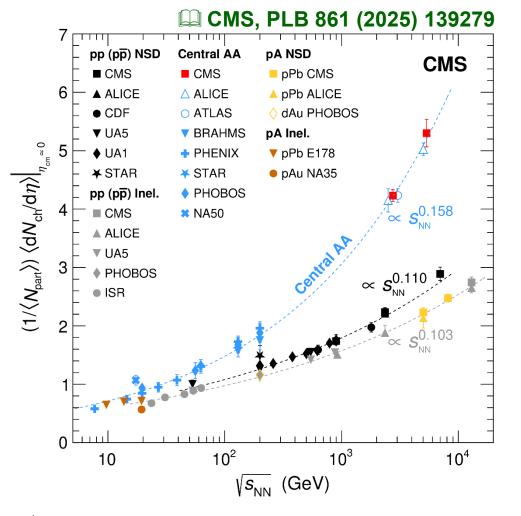
➡ Collision geometry

 \Rightarrow Collision energy: power law dependence on $\sqrt{s_{NN}}$

• Large energy density in the created "fireball":

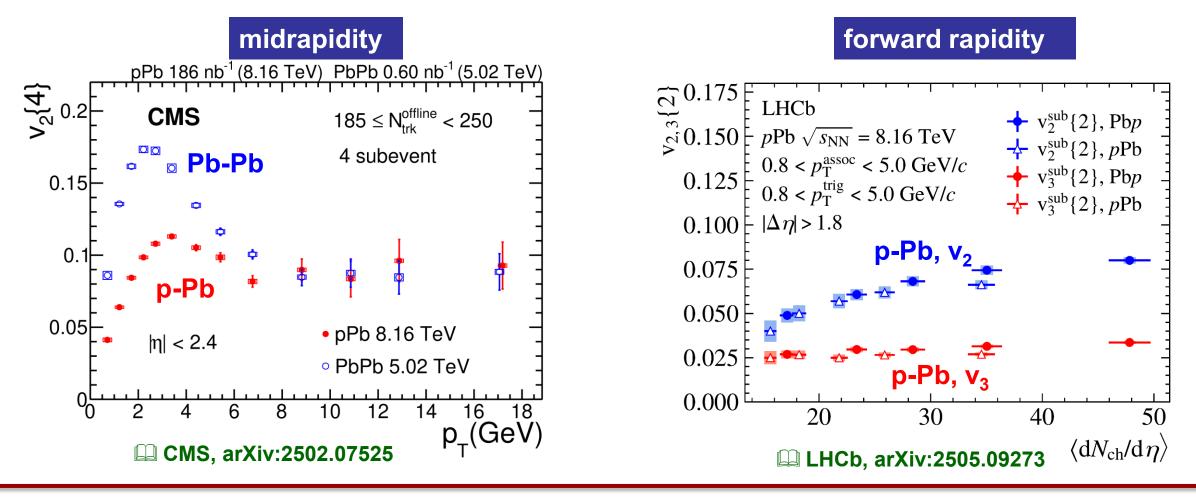
 \Rightarrow ε > 10 GeV/fm³ at τ=1 fm/c in central Pb-Pb collisions at $\sqrt{s_{NN}}$ =5.02 TeV

ALICE, EPJC 84 (2024) 813

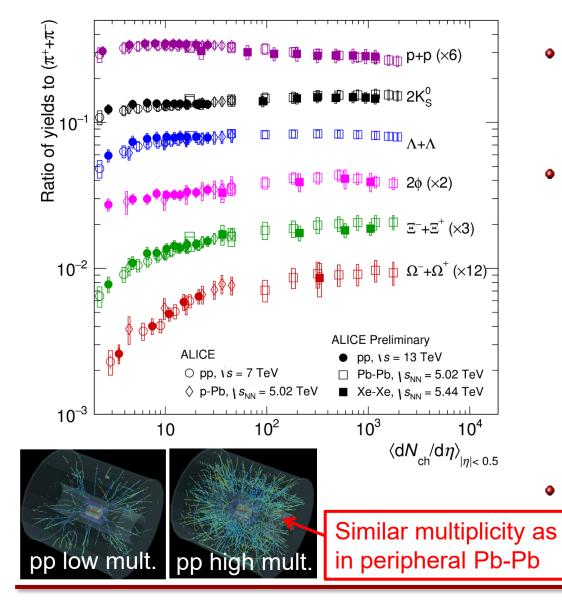


Collectivity in small collision systems

Long-range angular correlations observed in high-multiplicity pp and p-Pb collisions
 Similar features as those observed in Pb-Pb and interpreted as due to collective flow



Strangeness production vs. multiplicity

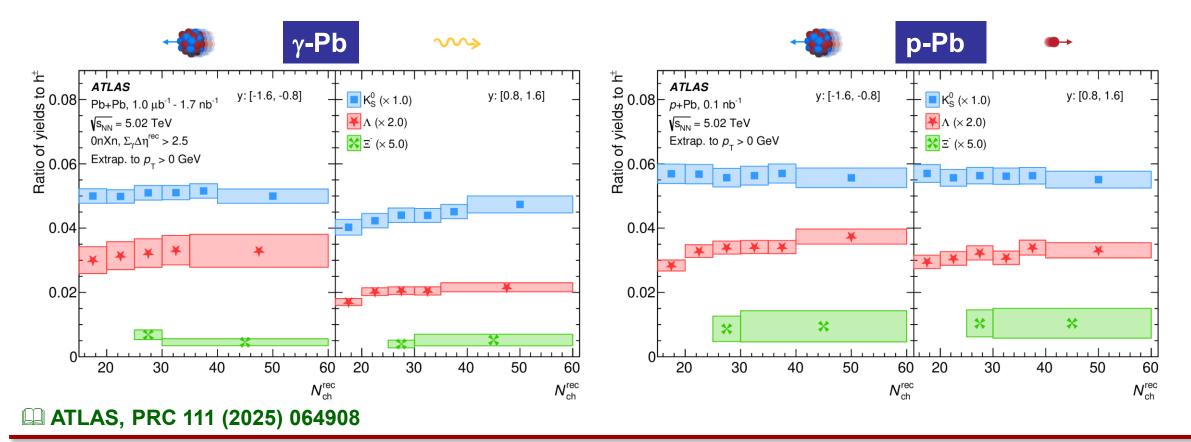


- Smooth evolution of hadrochemistry with multiplicity from small to large collision systems
- Significant enhancement of strange to nonstrange particle production with multiplicity
 - Strangeness production increases with multiplicity in pp and p-Pb collisions
 - ✓ Effect related to strangeness content rather than mass
 - Plateau reached in Pb-Pb collisions at the value expected from statistical hadronisation model with gran-canonical formulation
- Challenge for pp event generators

Strangeness production in UPC

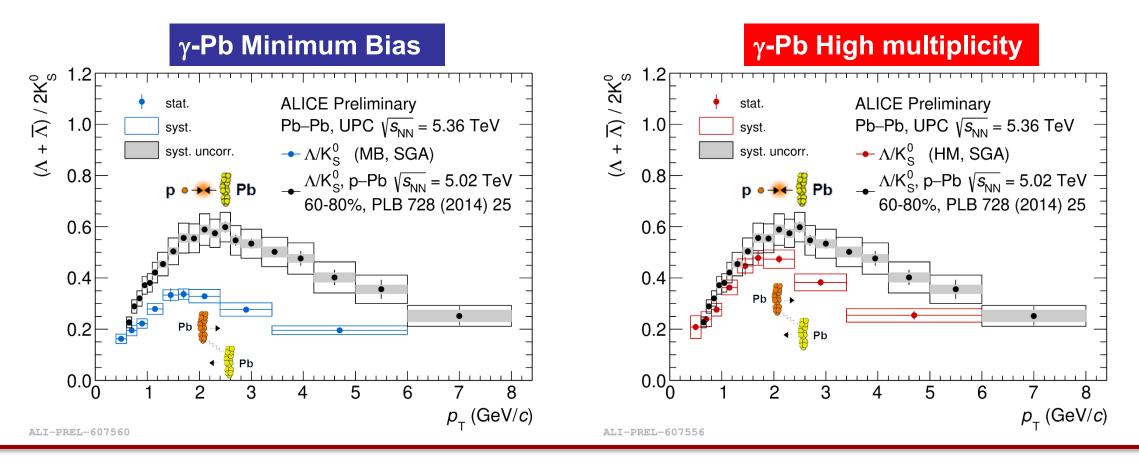
• Yields of hadrons, K_{s}^{0} , Λ and Ξ^{-} in same multiplicity classes in p-Pb and photonuclear events in Pb-Pb collisions (γ -Pb)

 \Rightarrow Similar strange / hadron rations for p-Pb and in the Pb-going side of γ -Pb events \Rightarrow Lower ratios on γ -going side



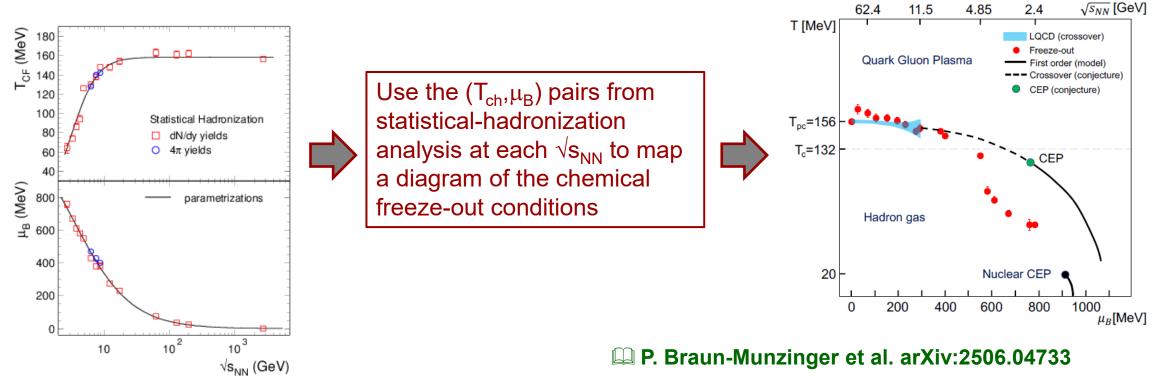
Strangeness production in UPC

- Enhancement of Λ/K_s^0 ratio in photo-nuclear collisions at intermediate p_T reminiscent of that measured in p-Pb collisions
 - In high-multiplicity photo-nuclear events, the Λ/K⁰_s ratio approaches the values measured in low-multiplicity p-Pb

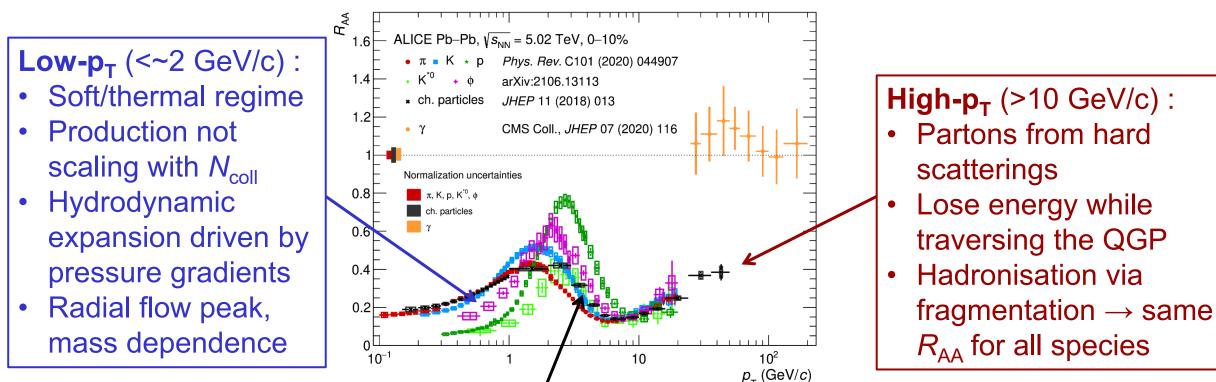


Mapping the phase diagram

- Statistical hadronization analysis of hadron yields at different collision energies
 - ⇒ Chemical freeze-out temperature T_{ch} increases with $\sqrt{s_{NN}}$ at low energies and saturates at T_{ch} ~156 MeV at top SPS energy
- At high $\sqrt{s_{NN}}$: T_{ch} very close to the pseudo-critical temperature from lattice QCD
 - Inelastic interactions between hadrons cease within a narrow temperature interval after QGP hadronization



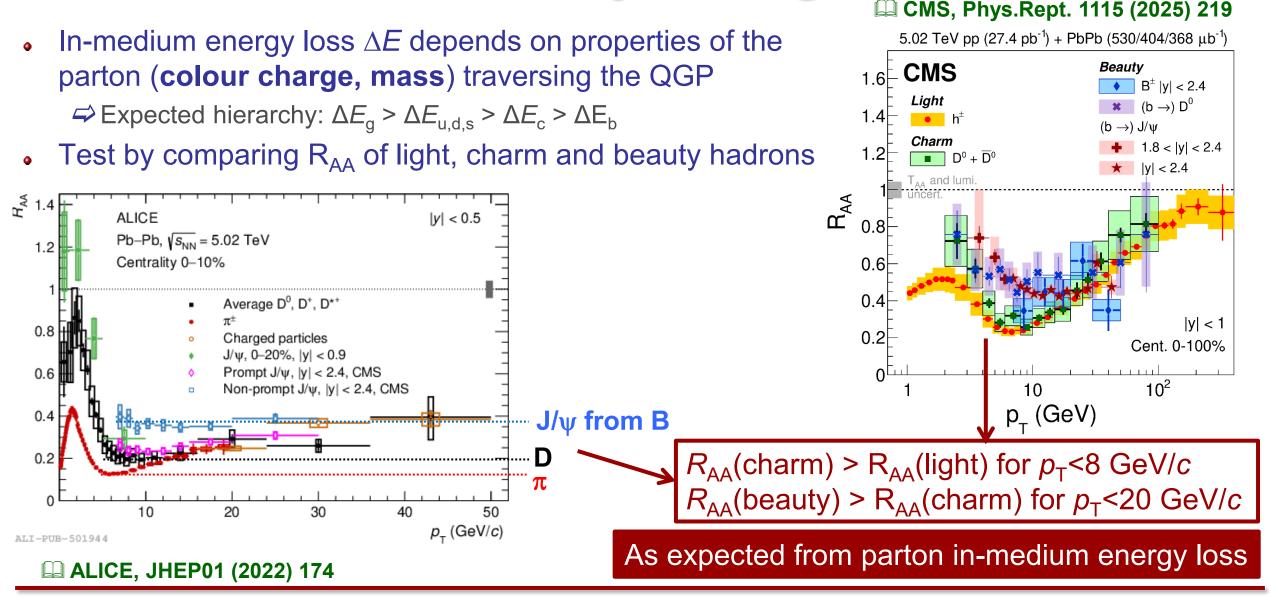
Single particle R_{AA}: family portrait



Intermediate-p_T (ca. $3 < p_T < 8 \text{ GeV/c}$) :

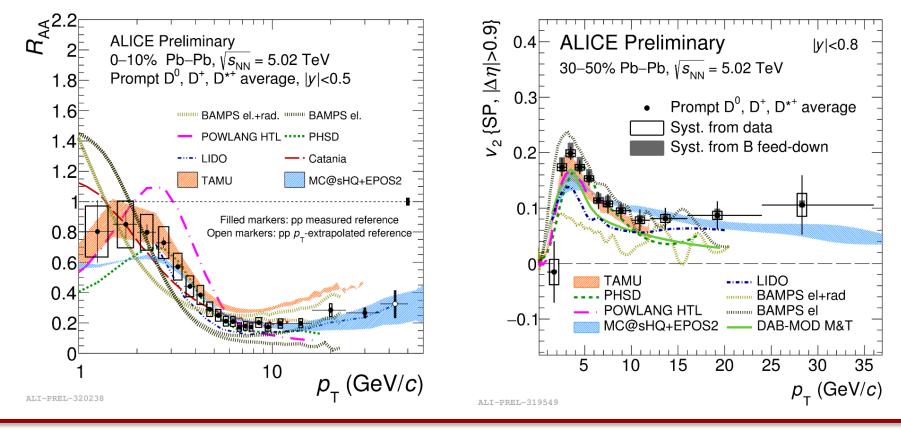
- Kinetic regime (not described by hydro)
- Different R_{AA} for different hadron species
 - Inconsistent with hard partons + energy loss + universal fragmentation
- Features described with in-medium hadronization via quark recombination

Charm vs. beauty vs. light flavours



Charm R_{AA} and v₂ phenomenology

- Simultaneous comparison of R_{AA} and v_2 to models can constrain QGP properties and the description of charm-quark interaction and diffusion in the medium
 - Interplay of CNM effects, collisional and radiative energy loss, hadronisation via coalescence and fragmentation and realistic underlying medium evolution required to describe data



Jet substructure in *γ*-jet events

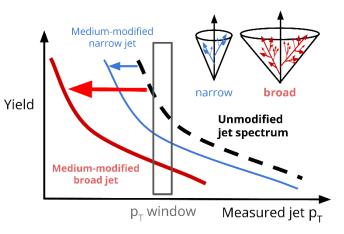
• Jet substructure: narrower jets in Pb-Pb than in pp for inclusive jets with same $p_{\rm T}$

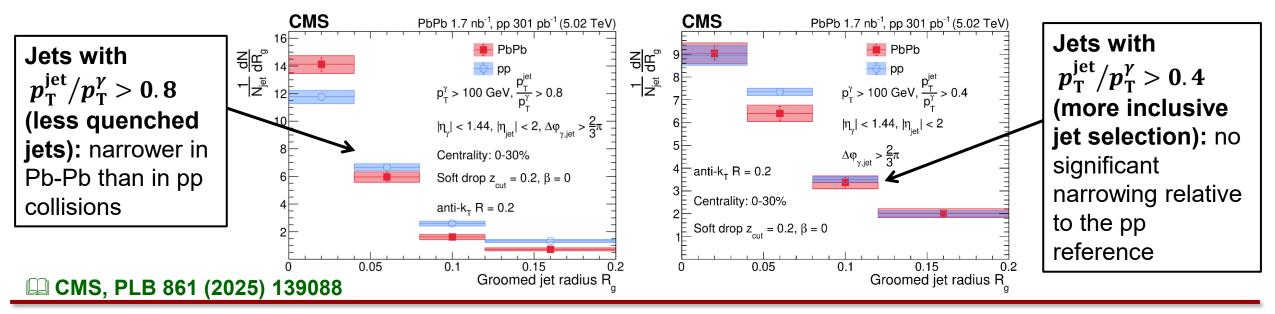
Interpretation in terms of broadening of parton shower require to assess potential selection biases

✓ A given jet p_{τ} interval is preferentially populated by jets that are less quenched

• Measurements of groomed jet radius R_{g} in γ -jet events

 \Rightarrow Photon provides access to p_T of the parton that initiated the recoiling jet \Rightarrow Show that selection biases play an important role





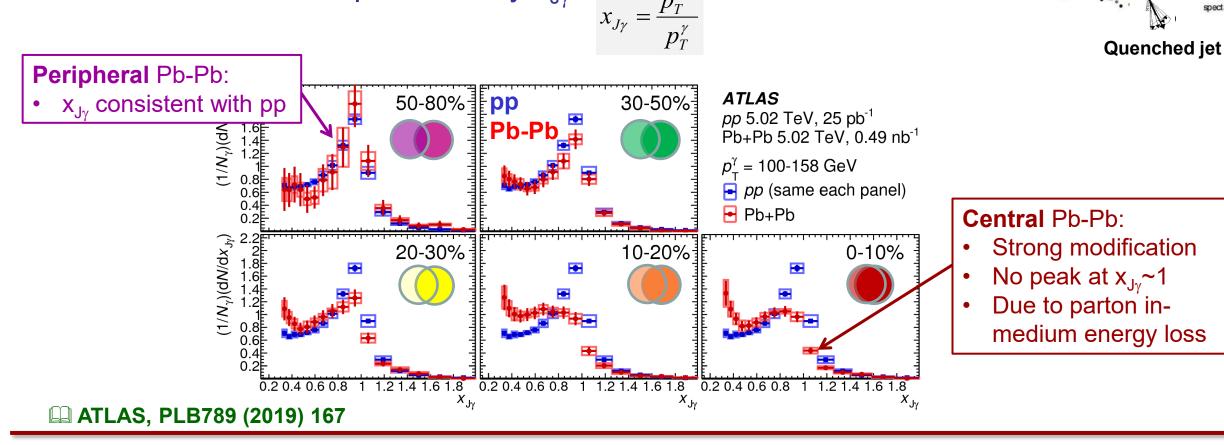
Photon-jet correlations

Photon-jet p_{T} balance: •

➡ Photon does not interact strongly in the QGP

 \Rightarrow Access to p_T of hard-scattered parton before it loses energy in the medium

Momentum balance quantified by $x_{J_{V}}$: 0





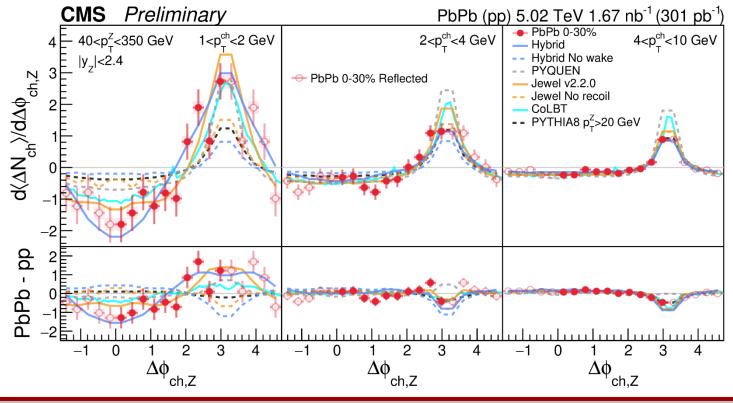
spectator

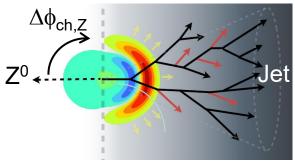
Medium response via Z⁰+jet events

• Explore medium response by measuring yield of Z⁰-tagged charged hadrons vs. azimuthal angle relative to Z⁰ ($\Delta\phi_{ch,Z}$)

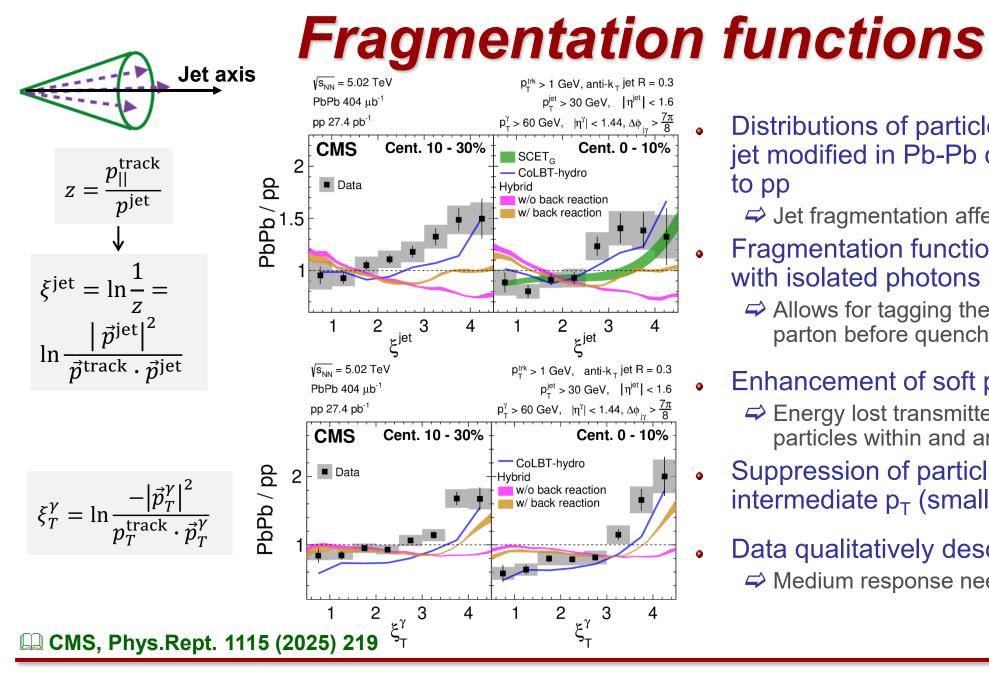
➡ Modified in central Pb-Pb as compared to pp collisions

- Data better described by models that include medium recoil effects
- ⇒ First evidence of medium response effects caused by a hard probe



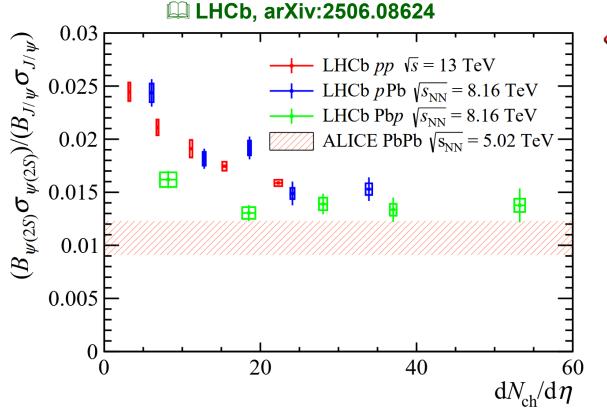


CMS, PAS HIN-23-006



- Distributions of particle momenta within the jet modified in Pb-Pb collisions with respect to pp
 - \Rightarrow Jet fragmentation affected by jet quenching
- Fragmentation functions of jets associated with isolated photons
 - \Rightarrow Allows for tagging the properties of the initial parton before quenching occurs
- Enhancement of soft particles (large ξ)
 - ⇒ Energy lost transmitted to low-momentum particles within and around the jet
- Suppression of particles at high and intermediate p_{τ} (small ξ)
- Data qualitatively described by models Arr Medium response needed in the hybrid model

Charmonia in p-Pb collisions



• Ratio of $\Psi(2S)$ / J/ ψ yield vs. multiplicity

For p-Pb configuration (charmonia measured in the p-going direction)

✓ Decreasing ratio with increasing multiplicity

 Trend compatible with measurements in pp collisions, described by models with dissociation by comoving hadrons

For Pb-p configuration (charmonia measured in the Pb-going direction)

✓ No significant dependence on multiplicity

 ✓ Stronger 𝒱(2S) suppression with flatter behaviour, compatible with ALICE Pb-Pb results

 \rightarrow Additional suppression mechanisms in the Pb-going direction beyond comover effects?