V_{ud} : Recent theoretical progress and experimental opportunities

Leendert Hayen V_{ud} workshop GANIL 5 November 2024





Introduction

Theory progress in the last 5 years

Experimental opportunities

Summary & Outlook

Three out of four fundamental forces (no gravity):

Standard Model



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18 free parameters



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Great (annoyingly so), consistent with constraints at $\sim 10^{0-2}~\text{TeV}$



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Open questions: dark matter, gravity, neutrino masses, ...



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SM tests @ low energy: sensitive to off-shell exotic physics (footprints rather than actual beast)

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Besides precision QED $(a_{e,\mu}, r_p, \ldots)$, weak interactions probe

- (C)P violation
- CKM unitarity
- Lorentz structure

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Today: CKM unitarity

Introduction: Weak interaction & CKM matrix

Cabibbo-Kobayashi-Maskawa matrix relates weak and mass eigenstates

$$\left(\begin{array}{c} d\\s\\b\end{array}\right)_{w} = \left(\begin{array}{ccc} V_{ud} & V_{us} & V_{ub}\\V_{cd} & V_{cs} & V_{cb}\\V_{td} & V_{ts} & V_{tb}\end{array}\right) \left(\begin{array}{c} d\\s\\b\end{array}\right)_{m}$$

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(nuclear) eta decay, meson decay (π , K), $|V_{ub}|^2 \sim 10^{-5}$

Violations are sensitive to TeV scale new physics!

CKM unitarity: Current status

Signs of non-unitarity at few σ level...

Disagreement between K/2 and K/3 $|V_{us}|$ 'Cabibbo angle anomaly'



SM has V-A structure, but more generally

$$\begin{split} \mathcal{L}_{\text{eff}} &= -\frac{G_{\text{F}} \,\tilde{V}_{ud}}{\sqrt{2}} \bigg\{ \bar{e} \gamma_{\mu} \nu_{L} \cdot \bar{u} \gamma^{\mu} [c_{V} - (c_{A} - 2\epsilon_{R}) \gamma^{5}] d + \epsilon_{\text{S}} \, \bar{e} \nu_{L} \cdot \bar{u} d \\ &- \epsilon_{P} \, \bar{e} \nu_{L} \cdot \bar{u} \gamma^{5} d + \epsilon_{\text{T}} \, \bar{e} \sigma_{\mu\nu} \nu_{L} \cdot \bar{u} \sigma^{\mu\nu} (1 - \gamma^{5}) d \bigg\} + \text{h.c.}, \end{split}$$

at the quark level

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at the quark level

All ϵ_i are proportional to $(M_W/\Lambda_{BSM})^2$, change kinematics $\epsilon_i \lesssim 10^{-4} \rightarrow \Lambda_{BSM} \gtrsim 15$ TeV assuming natural couplings

CKM unitarity: V_{ud}

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 $\Gamma \propto {\it G}_{\it F}^2 |V_{\it ud}|^2 (1+{\it RC})|\langle {\it O}_{\sf hadr}
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Things you need to know

- G_F (μ lifetime)
- Radiative corrections
- Hadronic theory
- For each β transition: $t_{1/2}, Q_{\beta}, BR, (GT/F \text{ mixing})$

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Master formula

$$ft(1+\delta_R')(1+\Delta_R^V)(1+\delta_{NS}-\delta_C) = \frac{K}{G_F^2 V_{ud}^2 M_{\text{tree}}^2}$$

CKM unitarity: V_{ud} precision

Nuclear sandbox \rightarrow make hadronic theory easy

- Pion
- Neutron

- $\bullet~$ Superallowed $0^+ \rightarrow 0^+$
- T = 1/2 mirrors

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 $\pi^+
ightarrow \pi^0 e^+
u_e$ very hard (BR $\sim 10^{-8}$), SA new nuclear corrections!

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Status of $0^+ \rightarrow 0^+$ great nuclear structure triumph

LH, ARNPS 74 (2024) 497

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Recall master equation:

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Separate into tree level & loop level

Summary:

- δ_C : Isospin symmetry breaking of M_F
- f: phase space factor
- δ'_R : 'outer' radiative corrections
- Δ_R^V : single-nucleon 'inner' radiative corrections
- δ_{NS} : Changes in Δ_R^V due to nuclear structure

All except for Δ_R^V are **open questions** to this day!

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Summary of changes:

- δ_C : Ab initio & data-driven methods
- f: weak radii & Fermi function
- δ'_R : RGE methods find differences in $\mathcal{O}(\alpha^2 Z^3)$ and beyond
- Δ_R^V : Dispersion and lattice QCD confirm 'inner' RC change
- δ_{NS} : Focus on coherent quasielastic nuclear response

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$$M_F^2 = (M_F^0)^2 (1 - \delta_C)$$

with $\delta_{C} \sim 0.1 - 1\%$ for **nuclei**.

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with $\delta_{\it C} \sim 0.1-1\%$ for <code>nuclei</code>. Traditional approaches separate into

- δ_{C1} : isospin-mixing meaning $\langle \pi | a_{p,\alpha} | \phi_i \rangle^* \neq \langle \phi_f | a_{n,\alpha}^{\dagger} | \pi \rangle$
- δ_{C2} : radial mismatch, i.e. proton and neutron orbits are not the same

but conceptual issues already noted 15 years ago (Miller & Schwenk)

See talk by N. Smirnova

New proposal to use charge radii & ab initio theory. May write

$$\delta_{\mathcal{C}} \simeq \sum_{T=0,1,2} rac{\langle a; T || V_{\mathrm{ISB}} || g; 1
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over all states a and ground state g, assuming $V_{\rm ISB}$ is isovector.

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Interesting development: In single nucleon, ISB was assumed negligible $(\delta_C \sim (m_u - m_d)^2 / \Lambda_{\rm QCD}^2)$, but recently challenged and can be $\mathcal{O}(10^{-4})!$

PLB 838 (2023) 137654; PLB 846 (2023) 138259

Phase space updates

Integrating over β spectrum in usual expression $f = m_e^{-5} \int_{m_e}^{E_0} dE \ pE(E_0 - E)^2 F(Z, E) C(Z, E) K(Z, E)$

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$$f = m_e^{-5} \int_{m_e}^{E_0} dE \ pE(E_0 - E)^2 F(Z, E) C(Z, E) K(Z, E)$$

but contains subtleties

- Depends on nuclear wave functions in $C(Z, E) \rightarrow$ weak charge density ρ_{cw}
- Special place for Fermi function F(Z, E) → is this a nice QFT object?

First point was long known (C_I in Wilkinson, shell model in H&T), but model-dependent. Data-driven treatment using charge radii

$$\rho_{\rm cw} = \rho_{\rm ch,0} + \frac{Z_{-1}}{2} (\rho_{\rm ch,-1} - \rho_{\rm ch,1})$$

uncertainties $\mathcal{O}(10^{-3-4})$

Seng & Gorchtein, PRC 109 045501 (2024)

Phase space and δ'_R updates

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F(Z, E) probes e^{\pm} density at r = 0, but since solution $\rightarrow \infty$, introduce UV cutoff *avant la lettre*: *R*, the nuclear radius.

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Special place of Fermi function is artefact of traditional segmented calculations (actually long-wavelength photon exchange)

F(Z, E) probes e^{\pm} density at r = 0, but since solution $\rightarrow \infty$, introduce UV cutoff *avant la lettre*: R, the nuclear radius.

Currently called into question: is F(Z, E) a good QFT object? When including higher-order corrections $(F \rightarrow F(1 + \delta_R))$, things become more complicated.

Still an **open question**, confusion due to disagreements with older calculations (1980's Jaus & Rasche)

PRL 133, 021803 (2024), PRD 109, 056006 (2024), PRD 108 (2023) 053003
Δ_R updates

Loop contribution that is (\sim)solved: Δ_R^V

Single-nucleon RC in β -decay can (\sim) be separated into

- 1. Energy-dependent, QCD-independent part: δ_R
- 2. Energy-independent, QCD-dependent part: Δ_R

Δ_R updates

Loop contribution that is (\sim)solved: Δ_R^V

Single-nucleon RC in eta-decay can (\sim) be separated into

1. Energy-dependent, QCD-independent part: δ_R

2. Energy-independent, QCD-dependent part: Δ_R

 δ_R mainly originates from real photon emission

$$\Delta_R^{V,A}$$
 renormalizes $g_{V,A}$

$$g_i^2
ightarrow g_i^2 (1 + \Delta_R^i)$$





The culprit for Δ_R^V : famous γW box



Specifically, axial-vector contribution \rightarrow symmetries don't save you & QCD at intermediate effects

Recent breakthrough using dispersion relations

2006: Marciano & Sirlin $\Delta_R^V = 0.02361(38)$, but heuristic uncertainty from 'intermediate' energy scale

2018: Seng, Gorchtein, Patel, Ramsey-Musolf $\Delta_R^V = 0.02467(22)$ 4 σ shift

Beginning of our CKM debacle!



Seng, Gorchtein, Ramsey-Musolf PRD 100 (2019) 013001

Δ_R^V updates

Number of different calculations performed, convergence



Small differences remain, neutron experimental uncertainty too large to distinguish LH, ARNPS 74 (2024) 497

δ_{NS} : status 2018

Nuclear medium changes nuclear response, but also spectrum





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Paradigm shift in analysis, two major effects Quasi-elastic contributions Nuclear polarization

$$\delta^{A}_{NS} = \frac{\alpha}{\pi} [-0.47 \pm 0.14]^{\text{QE}}$$
 $\delta^{A}_{NS}(E) \sim (1.6 \pm 1.6) \times 10^{-4} \left(\frac{E}{\text{MeV}}\right)^{-4}$

Estimated using free Fermi gas Current $0^+ \rightarrow 0^+$ bottleneck

Seng et al., PRD 100 013001

Current status on δ_{NS}

More sophisticated picture, first ab initio calculations emerging (See talk by Mehdi Drissi)



Energy-dependent effects might be detectable, **nuclear** shadowing effects largely unknown Seng, Gorchtein ARNPS 74 (2024) 1

Situation is analogues but more complicated than $0^+ \to 0^+.$ Significant questions on:

- How do energy-dependent terms enter for axial transitions?
- What about nuclear shadowing for spin-dependent transitions?

Mirror decays extract $\rho = g_A M_{GT}/g_V M_F$ from angular correlations $(a_{\beta\nu}, A_{\beta})$, but both effects may mean $\rho^{\text{corr}} \neq \rho^{\text{Ft}}$.

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Happened before: double counting was resolved and V_{ud}^{mirror} now agrees with $V_{ud}^{0^+ \rightarrow 0^+}$ LH, PRD 103, 113001; LH, ARNPS 74 (2024) 497



Progress in nuclear ab initio theory

Field is charging full steam ahead on nuclear ab initio



H. Hergert, Frontiers in Physics (2020)

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Now, let's talk experiment

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Several science drivers in parallel

- New experimental techniques to sidestep common systematics
- Spectrum shape measurements for $\delta_{\it NS}$ validation
- Precision measurements on low-mass isotopes for nuclear ab initio ladder benchmark

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Necessary push in neutrons, opportunities in mirrors, validation in superalloweds

Progress in neutron experiments

Situation is complicated in both τ_n and $\lambda = g_A/g_V$ determinations



'Outlier' measurements agree with most precise for CKM unitarity

Status on λ

Current PDG average



aSPECT $(a_{\beta\nu})$ is in tension with other recent measurements (A_{β}) Falkowski et al., JHEP04(2021)126

Progress in neutron experiments

Several campaigns worldwide (see talk by Bastian Märkisch)



Nab is only $a_{\beta\nu}$ experiment aiming at 0.1%, crucial input

Nab - overview

Measurement of β - ν angular correlation $d\Gamma \propto d\Gamma_0 \left[1 + a_{\beta\nu}\beta \hat{p}_e \cdot \hat{p}_\nu\right]$ in **neutron** β **decay** @ SNS (ORNL)





Nab - overview

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To leading order in SM

$$a_{eta
u}=rac{1-g_A^2}{1+3g_A^2}$$

and

$$\frac{\delta a}{a} \approx 5 \frac{\delta g_A}{g_A}$$

meaning factor 5 sensitivity enhancement!



Nab progress

Measure p⁺ instead of ν , $\vec{p}_p = -(\vec{p}_e + \vec{p}_\nu)$



Nab progress

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First physics runs at ORNL are promising!



Sensitivity enhancement in mirror transitions

Several mirror isotopes have well-known ft values to rival SA



Strong enhancements depending on the system

LH, ARNPS 74 (2024) 497

Community is investi(gati)ng in different ideas



with new spectroscopy techniques & traps

with additional great progress in the A = 8 system

New technology: CRES

Cyclotron Radiation Emission Spectroscopy

$$f = \frac{|q|}{2\pi} \frac{B}{m_e + E_{kin}}$$

 $^{6}\mathrm{He}$ and $^{19}\mathrm{Ne}$





New technology: CRES

Use ratio method: ⁶He and ¹⁹Ne have opposite b_F sign



Ratio means many systematic effects cancel to first order

Physical Review Letters 131 (2023), 082502

Direct recoil spectroscopy

Richness in pure recoil spectra, but experimentally very difficult!



Enabled by novel superconducting tunnel junctions

SALER@FRIB: First STJ online measurements



- Acceptance testing complete
- Commissioning started
- Will continue through 2024





Open STJs up to all ISOL beams, precision spectroscopy



Reduce systematic effects by 2-3 orders of magnitude!

(See talk by Mohamad Kanafani)

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New spectroscopy techniques incoming, recoil spectroscopy with quantum sensors is highly promising!

Thank you

Thank you!



 β decay symmetries according to Stable Diffusion

Aside: recent progress on Δ_R^A

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$$\Delta_R^A-\Delta_R^V=0.13(13) imes 10^{-3}$$

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$$\Delta_R^A - \Delta_R^V = 0.13(13) \times 10^{-3}$$

but only first half of the story... also here large ISB effects



First time: $\delta_{\text{RC}}^{(\lambda)} \in \{1.4, 2.6\} \cdot 10^{-2}$ LH, PRD 103 113001; Seng, Particles 2021, 397; Gorchtein & Seng, JHEP 10 53; PRL 129 121801

Superconducting tunnel junctions (Slide by Kyle Leach)



- Pulsed 355 nm (3.49965(15) eV) laser at 5 kHz fed through optical fiber to 0.1 K stage
- Illumination of STJ provides a comb of peaks at integer multiples of 3.5 eV
- Intrinsic resolution of our Ta-based devices is between ~1.5 and ~2.5 eV FWHM at ~10 – 200 eV
- Stable response and small quadratic nonlinearity (10⁻⁴ per eV)



The BeEST experiment (Slide by Kyle Leach)

∂TRIUMF

Rare-isotope implantation at TRIUMF-ISAC





A. Samanta et al., Phys. Rev. Mat. (in press) (2022) S. Friedrich et al., J. Low Temp. Phys. (in press) (2022) C. Bray et al., J. Now Temp. Phys. (in press) (2022) K.G. Leach and S. Friedrich, J. Low Temp. Phys. (in press) (2022) S. Friedrich et al., Phys. Rev. Lett. **126**, 021803 (2021) S. Friedrich et al., Phys. Rev. Lett. **127**, 032701 (2020) S. Friedrich et al., J. Low Temp. Phys. **200**, 200 (2021)

Ta, Al, and Nb-based STJ Sensors









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CKM unitarity: Cabibbo Angle Anomaly

Signs of non-unitarity at several σ (Falkowski CKM2021)



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Takeaways assuming Standard Model physics:

- Most precise V_{ud} & V_{us} not consistent with unitarity
- Significant internal inconsistencies within V_{us}
- Taken at face value $\sim 3\sigma$ for new physics

A more modern way of interpreting BSM physics

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Effective field theory: new physics at scale $\Lambda_{BSM} \gg LHC$

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \sum_{i=1} c_i \frac{\mathcal{O}_{4+i}}{\Lambda^i_{BSM}}$$

effective operators O(i). Expansion in parameter $c_i/\Lambda_{BSM}^i \ll 1$

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Phenomenological theories will give different $\{c_i\}$,

but agnostic experimental analysis

SM has V-A structure, but more generally

$$\begin{split} \mathcal{L}_{\text{eff}} &= -\frac{G_{\text{F}}\,\tilde{V}_{ud}}{\sqrt{2}} \bigg\{ \bar{e}\gamma_{\mu}\nu_{L}\cdot\bar{u}\gamma^{\mu}[c_{V}-(c_{A}-2\epsilon_{R})\gamma^{5}]d + \epsilon_{\text{S}}\,\bar{e}\nu_{L}\cdot\bar{u}d \\ &-\epsilon_{P}\,\bar{e}\nu_{L}\cdot\bar{u}\gamma^{5}d + \epsilon_{\text{T}}\,\bar{e}\sigma_{\mu\nu}\nu_{L}\cdot\bar{u}\sigma^{\mu\nu}(1-\gamma^{5})d \bigg\} + \text{h.c.}, \end{split}$$

at the quark level

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at the quark level

All ϵ_i are proportional to $(M_W/\Lambda_{BSM})^2$, change kinematics $\epsilon_i \lesssim 10^{-4} \rightarrow \Lambda_{BSM} \gtrsim 15$ TeV assuming natural couplings

Effective field theory tower Slide by V. Cirigliano



Effective field theory recipe Slide by V. Cirigliano

- In order to build L_{eff}, one needs to specify:
 - * Relevant low-E degrees of freedom: assume SM field content
 - * One Higgs doublet, no light VR and no other light fields
 - * Symmetries: L_{eff} must reflect symmetries of underlying theory
 - ★ Assume underlying theory respects SM gauge group SU(3)_c × SU(2)_W × U(1)_Y
 - ★ But not necessarily SM symmetries that result from keeping only terms of dimension ≤ 4
 - * Power counting in E/A, $v_{EW}/A <<1$ (recall $v_{EW} = G_F^{-1/2}$): organize analysis in terms of operators of increasing dimension (5,6,...)

Recent changes: Δ_R^V

Number of new calculations performed



Now good convergence: uncertainty halved but about 3σ shift

CKM unitarity: V_{ud} precision

Nuclear sandbox \rightarrow make hadronic theory easy

- Pion
- Neutron

- $\bullet~$ Superallowed $0^+ \rightarrow 0^+$
- T = 1/2 mirrors

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Status of $0^+ \rightarrow 0^+$ great nuclear structure triumph

CKM unitarity: V_{ud} precision

Four (\sim)competitive channels of extracting V_{ud}



Status of $0^+ \rightarrow 0^+$ great nuclear structure triumph

2018-2020 reanalysis nuclear structure current bottleneck

Superallowed uncertainties

Experimentally, $T_z = -1$ limited by BR (new ¹⁰C welcome)



Moving towards mature ab initio theory evaluation

Talk by Bertram Blank

Hardy & Towner PRC 102 (2020) 045501

Recent changes: δ_{NS}

Nuclear medium changes nuclear response, but also spectrum





Recent changes: δ_{NS}

Nuclear medium changes nuclear response, but also spectrum



Paradigm shift in analysis, two major effects Quasi-elastic contributions Nuclear polarization

$$\delta^{A}_{NS} = \frac{\alpha}{\pi} [-0.47 \pm 0.14]^{\text{QE}}$$
 $\delta^{A}_{NS}(E) \sim (1.6 \pm 1.6) \times 10^{-4} \left(\frac{E}{\text{MeV}}\right)$

Estimated using free Fermi gas Current $0^+ \rightarrow 0^+$ bottleneck

Seng et al., PRD 100 013001

On the radar: δ_C

Proton eq neutron inside nucleus $ightarrow M_F^2 = 2(1-\delta_{\mathcal{C}})$

- 1. Configuration interaction difference initial \leftrightarrow final
- 2. Different radial wave function (Coulomb)

$$\delta_C = \delta_{C1} + \delta_{C2}$$

On the radar: δ_C

Proton \neq neutron inside nucleus $\rightarrow M_F^2 = 2(1 - \delta_C)$

- 1. Configuration interaction difference initial \leftrightarrow final
- 2. Different radial wave function (Coulomb)

$$\delta_{C} = \delta_{C1} + \delta_{C2}$$



Grinyer et al., NIMA 622 (2010) 236

Progress in nuclear ab initio theory



H. Hergert, Frontiers in Physics (2020)

Monte Carlo methods (Slide by Saori Pastore)

Ab initio is providing bottleneck input for spectral measurements



Dominant terms $L_{1^{(0)}}$ and $E_{1^{(0)}}$ have model dependence of ~1% to ~2%

Looking at implementing δ_{NS} for ¹⁰C

Standard Model spectrum for ⁶He



No Core Shell Model (Slide by Michael Gennari)



Going heavier: IM-SRG type methods (Slide by Heiko Hergert)

- IMSRG for closed and open-shell nuclei: IM-HF and IM-PHFB
 - HH, Phys. Scripta, Phys. Scripta 92, 023002 (2017)
 - HH, S. K. Bogner, T. D. Morris, A. Schwenk, and K. Tuskiyama, Phys. Rept. 621, 165 (2016)

• Valence-Space IMSRG (VS-IMSRG)

- S. R. Stroberg, HH, S. K. Bogner, J. D. Holt, Ann. Rev. Nucl. Part. Sci. 69, 165
- In-Medium No Core Shell Model (IM-NCSM)
 - E. Gebrerufael, K. Vobig, HH, R. Roth, PRL 118, 152503
- In-Medium Generator Coordinate Method (IM-GCM)
 - J. M. Yao, J. Engel, L. J. Wang, C. F. Jiao, HH PRC 98, 054311 (2018)
 - J. M. Yao et al., PRL 124, 232501 (2020)

+ Coupled Cluster, \ldots

XYZ

Major advances in last decade, EFT come into its own

Quantifiable theory uncertainties are game-changer for precision FS: paradigm shifts are strong driver of progress in the field

Benefit from 'rigorous' theory overlap at low masses (NCSM, GFMC, QMC)

- $0^+ \rightarrow 0^+$:¹⁰C & ¹⁴O
- Promising isotopes: ${}^{6}\text{He}$, ${}^{11}\text{C}$, ...

to confidently go higher (CC, IM-SRG, IM-GCM, ...)

Path forward for $0^+
ightarrow 0^+ \ V_{ud}$

BeEST implantation



SALER implantation



11 MeV ¹¹C Beam w/ 8µm Al foil

For a given energy, initial beam from <u>ReA</u> can be +/- a few % in spread

- 1% spread gives ~50 nm width in the depth profile
 - Total ¹¹C⁺ to achieve goal: ~10⁷ (< 2 days of beam @ 100 pps)
 - Purity: 1 part in 10⁶

11.1 MeV ¹¹C Beam



10.9 MeV ¹¹C Beam

