

# ROXAS: a new spectral code for isolated neutron star gravitational wave signal.

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Work carried in LUTH

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- Core collapse supernovae and binary merger simulations in General Relativity (GR) have gone under tremendous progress during the last  $\sim 25$  years.
- Shibata and Uryu: first 3D binary neutron star (NS) merger simulation [Shibata and Uryū, 2000].
- Valencia formulation [Banyuls et al., 1997; Font, 2008]: conservative formulation of GR-hydro equations.
- Development of high-resolution shock capturing (HSRC) methods.
- Numerical Relativity (NR) simulation of compact binary mergers: waveforms of inspiral + merger + ringdown.
- Next two decades: post-O5 runs of LIGO-Virgo-Kagra + Einstein Telescope & Cosmic Explorer, sensitivity increased in the kHz band.

- Most hydrodynamical codes rely on conserved variables [Banyuls et al., 1997; Font, 2008].
- Different discretisation: finite volume [Cipolletta et al., 2021], Smooth Particle Hydrodynamics [Rosswog and Diener, 2021], spectral methods [Hébert et al., 2018].
- Allows the treatment of shocks.
- Inspiral + merger simulations possible.
- Conserved  $\leftrightarrow$  physical (*primitive*) variables computationally expensive.
- Only  $\sim 500$  merger simulations in two decades: too few for a parametric study.

## Writing a full-GR, full non-linear neutron star evolution code for post-merger phase.

- Final goal: run simulations of a post-merger hypermassive neutron star with 3-parameter equations of state (EoS), perform a parametric study by varying the EoS.
- Test beds in spherical symmetry with barotropic EoS.
- Axisymmetric and non-axisymmetric frequency extraction with barotropic EoS.
- Waveform extraction (gravitational waves signal).

# Framework

## 3+1 formulation of GR

We consider a perfect fluid of baryons carried by a timelike unitary 4-velocity  $u^\nu$ :

$$T^{\mu\nu} = (e + p)u^\nu u^\mu + pg^{\mu\nu}$$

with  $e$  the total internal energy,  $p$  the fluid pressure,  $g^{\mu\nu}$  the metric which is decomposed in the 3+1 formalism:

$$g_{\mu\nu}dx^\mu dx^\nu := -N^2 dt^2 + \gamma_{ij}(dx^i + \beta^i dt)(dx^j + \beta^j dt)$$

with  $N$  the lapse function,  $\beta^i$  the shift vector and  $\gamma_{ij}$  the 3-metric.

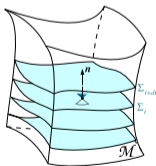


Figure: 3+1 foliation in spacelike hypersurfaces

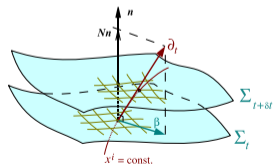


Figure: Illustration of the lapse and the shift

Physical variables are called "primitive" variables:

- $e$  the energy density,  $p$  the pressure
- $n_B$  the baryon number density
- $m_B$  a baryon mass
- $H = \ln\left(\frac{e+p}{m_B n_B}\right)$  the log-enthalpy
- $U_i$  the Eulerian velocity field
- $v_i$  the coordinate velocity field
- $c_s$  the sound speed



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The variables  $D = m_B n_B \Gamma^2$ ,  $S_j = (e + p) \Gamma^2 U_j$  and  $\tau = (e + p) \Gamma^2 - p$ , with  $\Gamma = (1 - U_i U^i)^{-1/2}$  the Lorentz factor, are conserved in the sense that  $\mathbf{u} = (D, S_j, \tau)$  obeys an equation that looks like

$$\partial_t \mathbf{u} + \text{div}(F(\mathbf{u})) = \text{source}$$

but the knowledge of  $e$ ,  $p$ ,  $U_i$  is compulsory to solve Einstein equations to compute the metric. **Recovery procedures are needed to recover the primitive variables in multiple steps of NR codes and account for a significant part of the computation:**

- Solving for the metric.
- Solving for the Riemann problems in each grid cell at each time step.

# New set of equations using primitive variables [Servignat et al., 2023]

From the principles of stress-energy and baryon number conservation:

$$\nabla_{\mu}(n_B u^{\mu}) = 0, \quad \nabla_{\mu} T^{\mu\nu} = 0,$$

the following holds for a barotropic, non-reactive perfect fluid with  $K_{ij}$  the extrinsic curvature tensor:

$$\begin{aligned} \partial_t U_i &= -v^j D_j U_i - D_i N - \frac{N}{\Gamma^2} \left( D_i H - \frac{\Gamma^2(1 - c_s^2)}{\Gamma^2 - c_s^2(\Gamma^2 - 1)} U_i U^j D_j H \right) \\ &\quad + U_j D_i \beta^j + U_i U^j D_j N \\ &\quad + \frac{N c_s^2}{\Gamma^2 - c_s^2(\Gamma^2 - 1)} U_i D_j U^j + \frac{N \Gamma^2 (c_s^2 - 1)}{\Gamma^2 - c_s^2(\Gamma^2 - 1)} U_i U^j U^l K_{jl} \\ \partial_t H &= -v^i D_i H - c_s^2 N \frac{\Gamma^2}{\Gamma^2 - c_s^2(\Gamma^2 - 1)} \left[ U^i U^j K_{ij} - \frac{U^i}{\Gamma^2} D_i H + D_i U^i \right] \end{aligned}$$

This new set of equations is **covariant** within the 3+1 formalism.

$e$ ,  $p$ ,  $c_s$  are recovered in a single call to the EoS (**no iteration**).

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## ROXAS: Relativistic Oscillations of non-axisymmetric neutron star

- Implementation based upon LORENE<sup>a</sup> (spectral methods).
- Two distinct grids (see next slide):
  - 1 Metric grid: spherical domains with nucleus, shells, CED.
  - 2 Hydro grid: one nucleus + shells.
- Hydro grid adapted to the surface of the star: deformed domains.
- Multi-domain matching.
- Filters.
- Well-balanced formulation of the equations.
- EoS: polytrope ( $p = \kappa n_B^\gamma$ )
- Dirac gauge + Maximal slicing.
- Einstein equations : fully constrained formalism (Poisson-like PDEs).
- Conformal Flatness Condition (CFC) = approximation.
- Time evolution with Adams-Bashforth explicit scheme (finite differences).
- The hydro grid external border is comoving with the star's surface.

<sup>a</sup><https://lorene.obspm.fr>

# Deformed domains

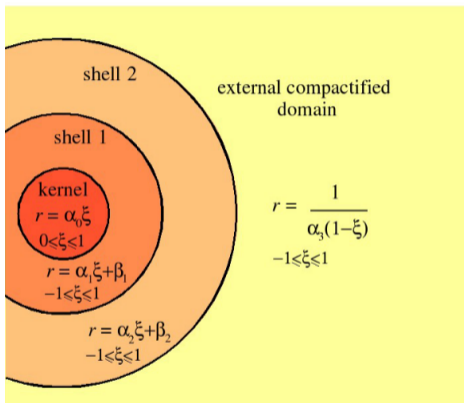


Figure: Metric grid: spherical numerical domains in LORENE.

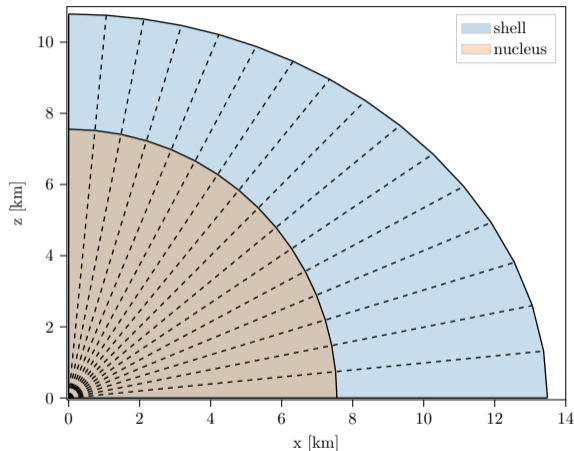


Figure: Hydro grid: deformed external domain (rotating star).

Test beds in spherical symmetry [Servignat et al., 2023]:

- Migration test.
- Collapse to a black hole.
- Spherically symmetric oscillations + frequency extraction: Fourier transform of  $R(t)$  (not shown here).

Non-spherically symmetric simulations:

- BU sequence: set of rigidly rotating polytropic NS + pseudo-polytropic representation of SLy4 EoS [Gulminelli and Raduta, 2015; Servignat et al., 2024].
- GW extraction with quadrupole formula (see next slides).
- Frequency extraction on  $Y_{\ell m}$  decomposition of  $R(t)$ .



# Validation of the relativistic code: migration test [Servignat et al., 2023]

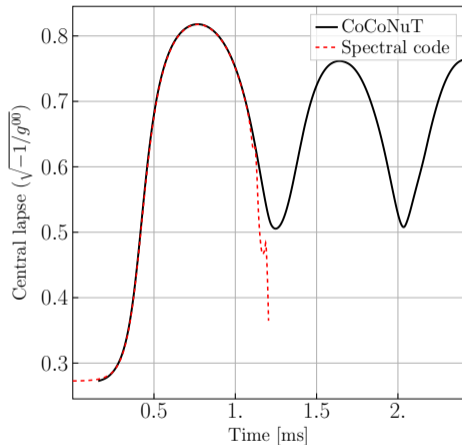


Figure: Central lapse.

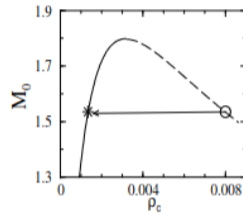
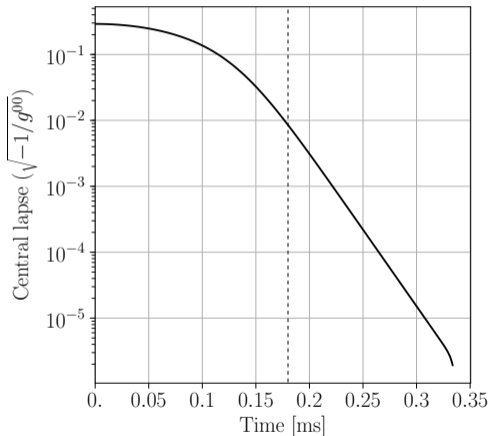


Figure: Migration test principle.

- Shock formation during the oscillations
- Gibbs phenomenon: the spectral code terminates.
- CoCoNuT: finite volume, high-resolution shock capturing techniques.
- 10s (spectral code) vs 6min (CoCoNuT).



→ tests the strong gravity regime.

- Central lapse down to a few  $10^{-6}$ .
- Freezing of the hydro quantities (maximal slicing).
- Apparent Horizon appearing around 0.18ms (dashed vertical line).

- Quadrupole formula:

$$h_{ij}^{TT}(\mathbf{x}, t) = \frac{2\mathcal{G}}{c^4 r} P_{ij}^{kl}(\mathbf{n}) \ddot{Q}_{kl} \left( t - \frac{r}{c} \right)$$

- Stress formula:

$$\ddot{I}_{ij} = \frac{\mathcal{G}}{c^4} \int_{\mathbb{R}^3} \rho(\mathbf{x}, t) (2v_i v_j - x_i \partial_j \Phi - x_j \partial_i \Phi) d^3x$$

- Plus/cross polarisation decomposition:

$$h_{ij}^{TT}(\mathbf{x}, t) = \frac{1}{R} (A_+ \hat{\mathbf{e}}_+ + A_\times \hat{\mathbf{e}}_\times)$$

where

$$\hat{\mathbf{e}}_+ = \hat{\mathbf{e}}_\theta \otimes \hat{\mathbf{e}}_\theta - \hat{\mathbf{e}}_\varphi \otimes \hat{\mathbf{e}}_\varphi,$$

$$\hat{\mathbf{e}}_\times = \hat{\mathbf{e}}_\theta \otimes \hat{\mathbf{e}}_\varphi + \hat{\mathbf{e}}_\varphi \otimes \hat{\mathbf{e}}_\theta.$$

- Expression of  $A_+$  and  $A_\times$ :

$$A_+^e = \ddot{I}_{zz} - \ddot{I}_{yy}, \quad A_\times^e = -2\ddot{I}_{yz}.$$

# GW extraction: stress formula, improvements

- Weight method [Dimmelmeier et al., 2002].  $\Phi_1$  pure Newtonian gravitational potential:

$$\Delta\Phi_1 = 4\pi\mathcal{G}\rho.$$

Relativistic gravitational potential in the Newtonian limit (exact up to first PN order):

$$\Phi_2 = \frac{1}{2}(1 - \Psi^4).$$

Weighted potential:

$$\Phi_w = \frac{\Phi_1 + a\Phi_2}{1 + a}, \quad a = \frac{1}{2}.$$

- Relativistic density

$$\rho^* = \Psi^6\Gamma\rho$$

- Replacing Newtonian velocity field  $v_i$  with relativistic Eulerian velocity field  $U_i$ .
- Final formula:

$$A_+^e = \frac{2\mathcal{G}}{c^4} \int_{\mathbb{R}^3} \rho^*(\mathbf{x}, t) (U_z^2 - U_y^2 + y\partial_y\Phi_w - z\partial_z\Phi_w) d^3x,$$

# Numerical oscillations

BU4 model: rigidly rotating polytrope,  $f_{\text{rot}} = 673$  Hz.

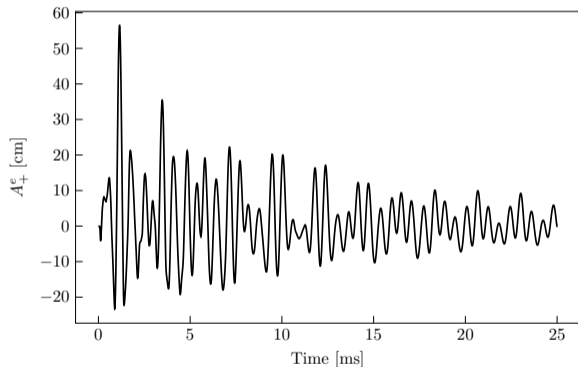


Figure: Waveform (BU4 model).

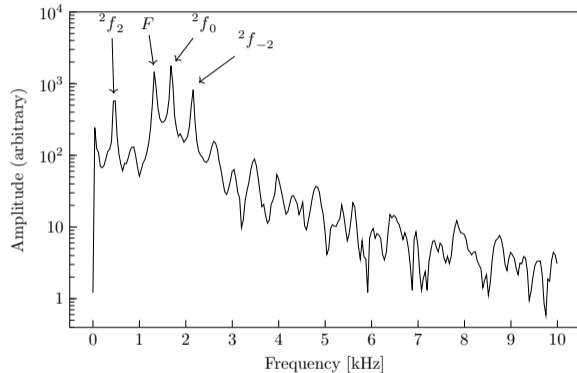


Figure: Spectrum (BU4 model).

# Axisymmetric modes [Gaertig and Kokkotas, 2008; Dimmelmeier et al., 2006]

Model	$f_{\text{rot}}$ [Hz]	$M [M_{\odot}]$	$r_e$ [km]	${}^2f_0$ [kHz]		Relative difference (%)	
				COWLING	CFC	COWLING	CFC
BU0	0	1.400	11.99	1.881	1.596	0.48	0.63
BU1	347	1.431	12.30	1.885	1.609	0.27	0.12
BU2	487	1.465	12.64	1.889	1.642	0.90	0.43
BU3	590	1.502	13.03	1.883	1.680	0.63	0.65
BU4	673	1.542	13.49	1.877	1.715	0.64	0.99
BU5	740	1.585	14.03	1.837	1.723	0.38	0.52
BU6	793	1.627	14.70	1.762	1.744	1.82	0.86
BU7	831	1.665	15.55	1.681	1.751	1.31	1.77
SLy4	191	1.361	9.34	2.400	1.964	×	×

# Non-axisymmetric modes [Krüger et al., 2010; Krüger and Kokkotas, 2020]

Model	$f_{\text{rot}}$ [Hz]	$M [M_{\odot}]$	$r_e$ [km]	${}^2f_{-2}$ [kHz]		Relative difference (%)	
				COWLING	CFC	COWLING	CFC/GR
BU0	0	1.400	11.99	1.881	1.565	0.00	0.83
BU1	347	1.431	12.30	2.259	1.953	0.09	0.56
BU2	487	1.465	12.64	2.364	2.052	0.09	0.10
BU3	590	1.502	13.03	2.443	2.123	0.61	0.24
BU4	673	1.542	13.49	2.473	2.153	0.57	0.09
BU5	740	1.585	14.03	2.480	2.158	0.44	1.44
BU6	793	1.627	14.70	2.477	2.143	0.40	1.82
BU7	831	1.665	15.55	2.396	2.105	1.88	3.94
SLy4	191	1.361	9.34	2.640	2.185	×	3.7

# Non-axisymmetric modes [Krüger et al., 2010; Krüger and Kokkotas, 2020]

Model	$f_{\text{rot}}$ [Hz]	$M [M_{\odot}]$	$r_e$ [km]	${}^2f_2$ [kHz]		Relative difference (%)	
				COWLING	CFC	COWLING	CFC/GR
BU0	0	1.400	11.99	1.881	1.565	0.00	0.83
BU1	347	1.431	12.30	1.367	1.080	0.51	0.09
BU2	487	1.465	12.64	1.121	0.840	0.98	0.48
BU3	590	1.502	13.03	0.918	0.640	0.09	0.63
BU4	673	1.542	13.49	0.721	0.451	1.66	1.11
BU5	740	1.585	14.03	0.520	0.280	2.69	4.64
BU6	793	1.627	14.70	0.348	0.130	2.01	12.3
BU7	831	1.665	15.55	0.163	×	11.0	×
SLy4	191	1.361	9.34	2.141	1.675	×	1.2



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# Summary and outlook

- Development of a spectral code to evolve isolated neutron stars.
- New set of equations relying only on primitive variables.
- Non-barotropic, spherically symmetric oscillations studied in [\[Servignat et al., 2024\]](#).
- Extraction of non-axisymmetric mode frequencies with barotropic EoS (polytropes + one realistic EoS).
- GW signals extracted with quadrupole formula (+ improvements).

# Summary and outlook

## Assets:

- Very light code (3D simulations on laptop/office computer).
- Reasonable simulation time (still to be improved).
- No recovery procedure.
- Radius of the star followed by the grid.
- Very good frequency extraction (although CFC).
- Soon to be published/open source.

## Drawbacks:

- In practice: no shock treatment.
- Spectral methods are unforgiving.
- Only oscillations/post-merger phase.

## Outlook:

- Improve interpolation (high priority).
- Getting rid of CFC (low priority).
- Perform longer simulations.
- Implement differentially rotating profiles.

# *Thank you!*



Figure: Roxas, the eponymous cat.

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$$\partial_t \mathbf{u} + \text{div}(F(\mathbf{u})) = \text{source}.$$

This formulation allows to sharply resolve shocks, but the knowledge of  $e$ ,  $p$ ,  $U_i$  is compulsory to solve EINSTEIN equations to compute the metric.

**Recovery procedures are needed to compute the primitive variables and account for a significant part of the computation:**

- Solving for the metric.
- Solving for the Riemann problems in each grid cell at each time step (finite-volume).

# Frequency extraction principle

- Frequency extraction: spectrum computed with FFT algorithm on  $R(t)$  then peak extraction.

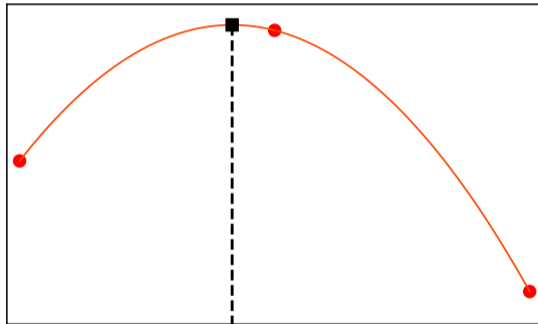


Figure: Frequency extraction principle.

Lapse, shift, and 3-metric entirely determine the geometry of  $\mathcal{M}$ :

$$g_{\alpha\beta} = \begin{pmatrix} g_{00} & g_j \\ g_i & g_{ij} \end{pmatrix} = \begin{pmatrix} -N^2 + \beta_k \beta^k & \beta_j \\ \beta_i & \gamma_{ij} \end{pmatrix}$$

→ The 3+1 formalism allows for a **natural separation of time and space**.

Definition of matter terms:

$$\begin{aligned} E &= T_{\mu\nu} n^\mu n^\nu \text{ the matter energy density,} \\ p_\alpha &= -T_{\mu\nu} n^\mu \gamma_\alpha^\nu \text{ the matter momentum density,} \\ S_{\alpha\beta} &= T_{\mu\nu} \gamma_\alpha^\mu \gamma_\beta^\nu \text{ the matter stress tensor,} \\ \text{and } S_t &= S_{\alpha\beta} \gamma^{\alpha\beta} \text{ and } T = T_{\alpha\beta} g^{\alpha\beta}. \end{aligned}$$

# Conformal decomposition

Let  $\mathbf{f}$  be a flat metric of signature  $(+, +, +)$ .  $(\mathcal{M}, \mathbf{g})$  is asymptotically flat i.e.  $\gamma = \mathbf{f}$  at spatial infinity.

$$\boxed{\Psi = \left(\frac{\gamma}{f}\right)^{1/12}} \quad \gamma = \det(\gamma_{ij}), \quad f = \det(f_{ij})$$

Using

$$\tilde{\gamma}_{ij} = \Psi^{-4} \gamma_{ij}$$

The EINSTEIN equations become (using the CFC:  $\tilde{\gamma} = \mathbf{f}$ ):

$$\begin{aligned} \Delta \Psi &= -2\pi \Psi^5 \left[ E + \frac{1}{16\pi} K_{ij} K^{ij} \right] \\ \Delta(N\Psi) &= 2\pi N\Psi^5 \left[ E + 2S_t + \frac{7}{16\pi} K_{ij} K^{ij} \right] \\ \Delta\beta^i + \frac{1}{3} \bar{D}^i \bar{D}_j \beta^j &= 16\pi N\Psi^4 p^i + 2\Psi^{10} K^{ij} \bar{D}_j \left( \frac{N}{\Psi^6} \right) \end{aligned}$$

# Conformal decomposition

Let  $\mathbf{f}$  be a flat metric of signature  $(+, +, +)$ .  $(\mathcal{M}, \mathbf{g})$  is asymptotically flat i.e.  $\gamma = \mathbf{f}$  at spatial infinity.

$$\boxed{\Psi = \left(\frac{\gamma}{f}\right)^{1/12}} \quad \gamma = \det(\gamma_{ij}), \quad f = \det(f_{ij})$$

Using

$$\tilde{\gamma}_{ij} = \Psi^{-4} \gamma_{ij}$$

The EINSTEIN equations become (using the CFC:  $\tilde{\gamma} = \mathbf{f}$ ):

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Additional conformal decomposition:

$$\hat{A}^{ij} = \Psi^{10} K^{ij}; \quad \hat{A}^{ij} \approx \bar{D}^i X^j + \bar{D}^j X^i - \frac{2}{3} \bar{D}_k X^k f^{ij}.$$

Starred quantities:

$$E^* = \Psi^6 E, \quad S_t^* = \Psi^6 S_t, \quad p_j^* = \Psi^6 p_j.$$

Leads to:

$$\Delta X^i + \frac{1}{3} \bar{D}^i \bar{D}_j X^j = 8\pi f^{ij} p_j^*,$$

$$\Delta \ln \Psi = -2\pi \Psi^{-2} E^* - \Psi^{-8} \frac{f_{il} f_{jm} \hat{A}^{lm} \hat{A}^{ij}}{8} - D_i \ln \Psi D^i \ln \Psi,$$

$$\Delta N = 4\pi N \Psi^{-2} (E^* + S_t^*) + N \Psi^{-8} f_{il} f_{jm} \hat{A}^{lm} \hat{A}^{ij} - 2D_i \ln \Psi D^i N,$$

$$\Delta \beta^i + \frac{1}{3} \bar{D}^i \bar{D}_j \beta^j = \bar{D}_j (2N \Psi^{-6} \hat{A}^{ij}).$$



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**Table:** Comparing the code and the literature.

	$\kappa$	$\gamma$	$H_c [c^2]$	$M [M_\odot]$	$R$ [km]	Fund. [kHz]	1st ov. [kHz]	2nd ov. [kHz]
[?]	100	2	0.2279	1.4	14.15	1.450	3.958	5.935
This work	100	2	0.2279	1.401	14.16	1.442	3.954	5.915
Relative difference						0.6%	0.1%	0.3%
[?]	7.308	5/3	$6.720 \times 10^{-2}$	$\times$	$\times$	0.824	1.94	2.86
This work	7.308	5/3	$6.720 \times 10^{-2}$	0.4866	16.49	0.823	1.95	2.86
Relative difference						0.1%	0.5%	0.0%

**Table:** COWLING approximation (fixing spacetime to initial value).

	$\kappa$	$\gamma$	$H_c [c^2]$	$M [M_\odot]$	$R$ [km]	Fund. [kHz]	1st ov. [kHz]	2nd ov. [kHz]
[?]	100	2	0.2279	1.4	14.15	2.696	4.534	6.346
This work	100	2	0.2279	1.401	14.16	2.685	4.548	6.339
Relative difference						0.4%	0.3%	0.1%