



Disentangling LIV from intrinsic lag? Sami Caroff (LAPP) 16 July 2024

Based on Bolmont et al 2022 ApJ 930 75



















Intrinsic time delays ?



Intrinsic delay

 $\Delta t_{total}(z) = \Delta t_{LIV}(z) + \Delta t_{int}(1+z) \qquad \Delta \tau_{total}(z) = \Delta \tau_{LIV} \kappa(z) + \Delta \tau_{int}(1+z)^2$





-

LIV
$$\frac{dP}{dtdE}(t + \lambda_n \kappa_n(z)E^n, E) = \phi_E(E)\phi_t(t + \lambda_n \kappa_n(z)E^n)$$

1.0

0.8

0.6

0.4

0.2 ·

0.0

-40

-20

Flux

$$\frac{dP}{dtdE}(t+\lambda(1+z)^2E,E) = \phi_E(E)\phi_t(t+\lambda(1+z)^2E)$$







100

z = 0.1

z = 0.5





- We can use the two dimension of the time lag to distinguish it with intrinsic
 - Redshift
 - Energy
- What is the problem ?
 - We have poor understanding of how both time lag vary versus redshift and energy for both intrinsic and LIV !
 - For Quantum Gravity induced lag, we simply don't know !
 - Better understanding for intrinsic time lag, but probably source dependent and even worst flare dependent
- Will try to show you under some approximations and simple scenarios, we can try to tackle this challenge
- But this possible approximations need to be corroborated with models, and even better with data



LIV redshift dependance (reminder)

$$E^{2} = \frac{p^{2}c^{2}}{a^{2}(t)} \left(1 + \frac{h_{QG}(E, E_{QG})}{h_{QG}(E, E_{QG})}\right)$$
Cosmological scale factor
$$E^{2} = \frac{p^{2}c^{2}}{a^{2}(t)} \left(1 + \sum_{1}^{n} \left(\frac{E}{E_{QG,n}}\right)^{n}\right)$$

$$v(E,t) = \frac{\partial E}{\partial p} \simeq \frac{c}{a(t)} \left(1 + \frac{n+1}{2} \left(\frac{E(t)}{E_{QG,n}} \right)^n \right) \qquad \left(\frac{E(t)}{E_{QG,n}} \right)^n \simeq \left(\frac{pc}{a(t)E_{QG,n}} \right)^n$$
$$\Delta t(E_h, E_l, z) = \frac{n+1}{2} \frac{E_h^n - E_l^n}{E_{QG}^n} \int_0^z \frac{(1+z')^n}{H(z')} dz'$$



QG perturbation term

$$E^{2} = \frac{p^{2}c^{2}}{a^{2}(t)} \left(1 + h_{QG}(p, E_{QG})\right)$$
$$E^{2} = \frac{p^{2}c^{2}}{a^{2}(t)} \left(1 + \sum_{1}^{n} \left(\frac{pc}{E_{QG,n}}\right)^{n}\right)$$

$$\Delta t(E_h, E_l, z) = \frac{n+1}{2} \frac{E_h^n - E_l^n}{E_{QG}^n} \int_0^z \frac{1}{H(z')} dz'$$







 $\Delta t($

LIV redshift dependence

$$E^{2} = \frac{p^{2}c^{2}}{a^{2}(t)} (1 + h_{QG}(E, p, E_{QG}))$$

$$E^{2} = \frac{p^{2}c^{2}}{a^{2}(t)} \left(1 + \alpha \sum_{1}^{n} \left(\frac{E}{E_{QG,n}}\right)^{n} + \beta \sum_{1}^{n} \left(\frac{pc}{E_{QG,n}}\right)^{n}\right)$$

$$E_{h}, E_{l}, z) = \frac{n+1}{2} \frac{E_{h}^{n} - E_{l}^{n}}{E_{QG}^{n}} \int_{0}^{z} \frac{\alpha + \beta(1+z)^{n}}{H(z')} dz'$$







- In simple scenarios (pure energy or momentum), redshift dependence is growing monotonously
- Not true anymore if you mix energy and momentum in dispersion
- But two important (and usefull) features appears whatever scenario we are using :
 - This is a propagation effect, lag is 0 at z=0
 - Whatever the model, the function is continuous







$$E^{2} = \frac{p^{2}c^{2}}{a^{2}(t)} \left(1 + h_{QG}(E, p, E_{QG})\right)$$

$$E^{2} = \frac{p^{2}c^{2}}{a^{2}(t)} \left(1 + \alpha \sum_{1}^{n} \left(\frac{E}{E_{QG,n}}\right)^{n} + \beta \sum_{1}^{n} \left(\frac{pc}{E_{QG,n}}\right)^{n}\right)$$

1

- Choosing an order n is only a simplification, the true shape is unknown
- Growing (in absolute) with Energy if we suppose on term n dominates
- If not, it can be whatever... but still for E < E_qg we expect nothing



ASTROVIBE - Sami Caroff - 16/07/2024



Simplification ?

• Let's explore various scenarios (those scenarios would need to be more physically motivated than what I am doing here) :

Same for all sources

$$\Delta t_{int}(E,z) = (1+z)(\Delta t_{mean} + \Delta t_{stoch})$$

Stochastic for each source (each flare?)

3 scenarios :

$$\Delta t_{mean}(E) >> \Delta t_{stoch}(E)$$
$$\Delta t_{mean}(E) << \Delta t_{stoch}(E)$$
$$\Delta t_{mean}(E) \sim \Delta t_{stoch}(E)$$



$$\Delta t_{mean}(E) >> \Delta t_{stoch}(E)$$

• Most simple scenario (but likely wrong right ?)

$$\begin{aligned} \text{LIV} \quad & \frac{dP}{dtdE}(t + \lambda_n \kappa_n(z)E^n, E) = \phi_E(E)\phi_t(t + \lambda_n \kappa_n(z)E^n) \\ \text{Intrinsic} \quad & \frac{dP}{dtdE}(t + \lambda(1+z)^2E, E) = \phi_E(E)\phi_t(t + \lambda(1+z)^2E) \end{aligned}$$

- Low redshift are particularly usefull for distinction between intrinsic and LIV
- Absence of delay \rightarrow we can already derive a limit on Lambda_s
- Detection of delays for source population :
 - LIV versus intrinsic can be tested and compared (Likelihood ratio)
 - In those two models, Lambda can be computed
 - The most constraining sources will probably be low redshift



- In this scenario, we expect lag to be fully stochastic (between sources, and even more between flares)
- Two major way to distinguish this scenario from LIV :
 - Continuity : delay at redshift z is fully uncorrelated of delay at z+dz, we simply expect a random distribution
 - Still possible to have very similar shape for complex LIV scenario (with momentum + energy mixed in complicated way), but simple simple scenario would be excluded
 - Significant variability between flares or events at a similar redshift are a smocking gun (between scenario 1 and 2, and with LIV as well)
- Can be constrained per source and combined as a mean stochasticity (or limits on mean stochasticity)
- Can as well derive standard deviation of this stochasticity (or limits on it)

 $(t + \lambda_i (1+z)^2 E, E) = \phi_E(E)\phi_t(t + \lambda_i (1+z)^2 E)$

Per source Lambda_i



- General trend + fluctuation around this trend
- Same than before but mixed
- We should still be able to distinguish between the two type of delay (particularly if stochastic lag is strong)
- It would probably need an iterative process to avoid degeneracy (first only mean lag, than adding lag per source for outliers only)

$$\frac{dP}{dtdE}(t + (\lambda_i + \lambda_{mean})(1+z)^2 E, E) = \phi_E(E)\phi_t(t + (\lambda_i + \lambda_{mean})(1+z)^2 E)$$





- This model can be implemented in LIVelihood without lot of efforts
- Already many data to constrain the different terms
- Model is probably degenerated but not completely
 - LIV effect for close sources is 0 (kappa(0) = 0 while (1+z=0)^2 = 1)
 - Stochastic delay very degenerated with mean delay :
 - Can be handled with an iterative procedure...
 - ... or by input from theory (energy dependence different ? Expected value?)
- Some inputs from theory would be needed
 - Does this model make sense ?
 - Can we simplify it (scenario 1 or 2 ?)
 - Can we have an estimation of lambda_i and lambda_mean?
 - Variability versus energy of lambda_i and lambda_mean ?
- But it can as well feed the theory :
 - We can already put a limit on lambda_mean based on absence of time dependent delay
 - Can it be used to constrain the models ? (Pulsar, AGN, GRBs ?)



- I believe that a simple model, handling multiple sources can be implemented and tested on data (currently available and future data)
- This model can be more robust by having theoretical input...
- ... and can give back to theory constraints
- Current status is that no lag is observed at TeV



- Redshift variability of the limit is ruled by three processes :
 - **Distance** (reduce events by D_L^2)
 - **EBL absorption** (high energy events absorbed)
 - Delay increase with redshift



Gamma 2022 - Sami Caroff - LIV consortium - 07/07/2022