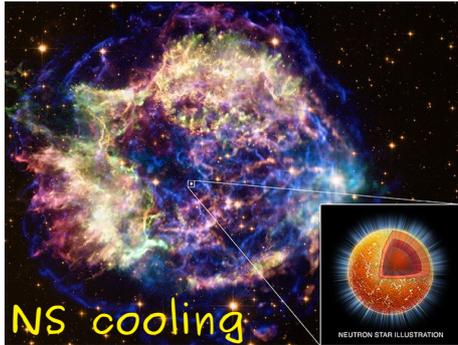
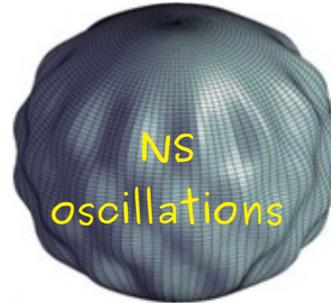


# Macroscopic superfluidity in neutron stars

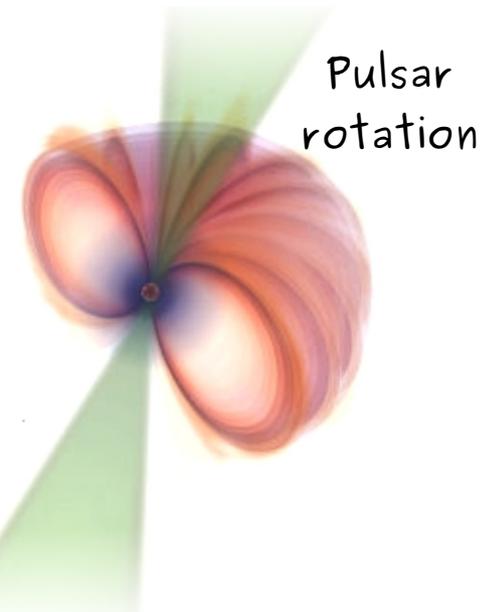
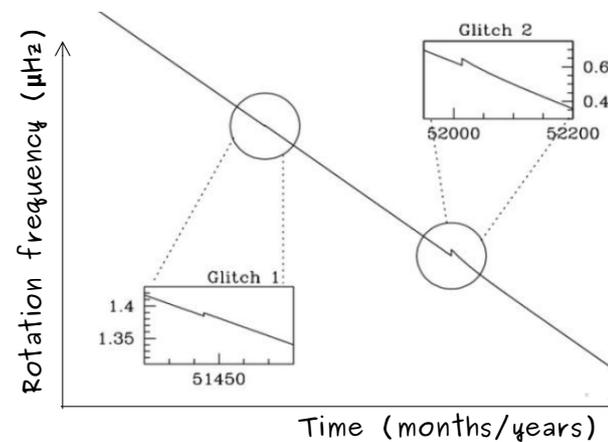


Superfluidity affects neutrino emission and heat capacity

Pulsar glitches (and maybe timing noise) are related to the presence of superfluid neutrons  
Review: [arXiv:2301.12769](https://arxiv.org/abs/2301.12769) (2022)



Superfluidity allows for more and different oscillation modes (like second sound in Helium-II)



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# Pairing channels

Total angular momentum operator:  $\mathbf{J} = \mathbf{L} + \mathbf{S}$

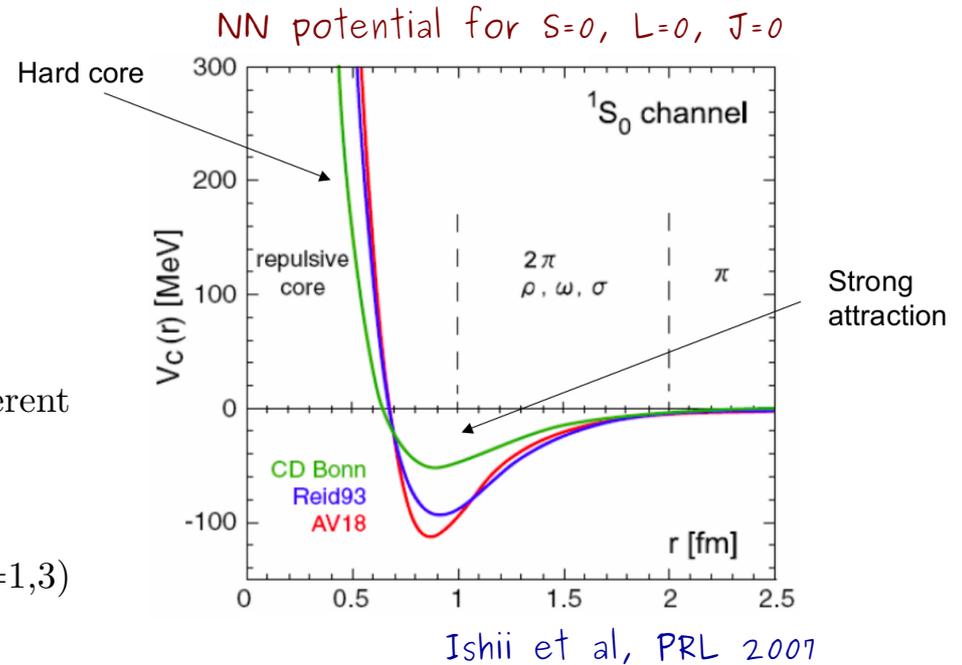
Notation:  $^{2S+1}L_J$  ( $L=0,1,2,3\dots \rightarrow S,P,D,F\dots$ )

$^1S_0$  isotropic pairing:  $\Delta =$  “energy gap”  $\sim 0.57 T_c$

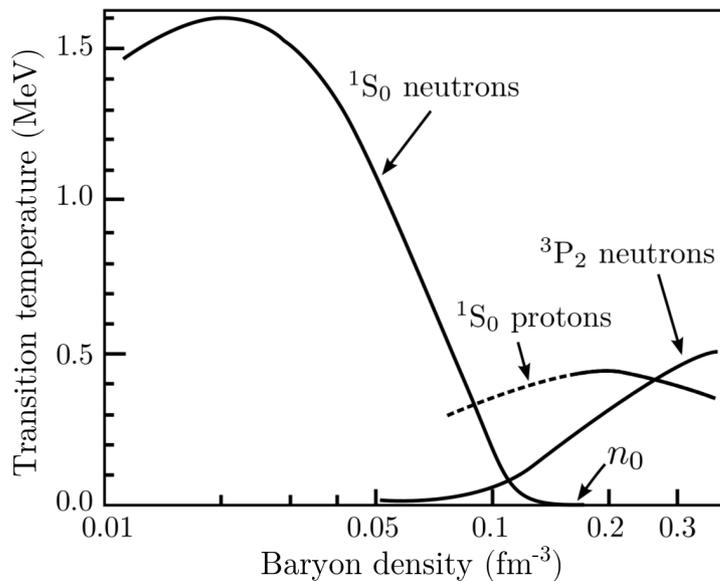
$^3S_1$ – $^3D_1$  binds the deuteron: but in NS  $\mathbf{n}$  and  $\mathbf{p}$  have very different Fermi surfaces  $\rightarrow$  **no n-p pairing**

$^3PF_2$  partial-wave channel ( $\Delta$  has contributions from both  $L=1,3$ ) preferred at larger Fermi momenta, where  $^1S_0$  is repulsive.

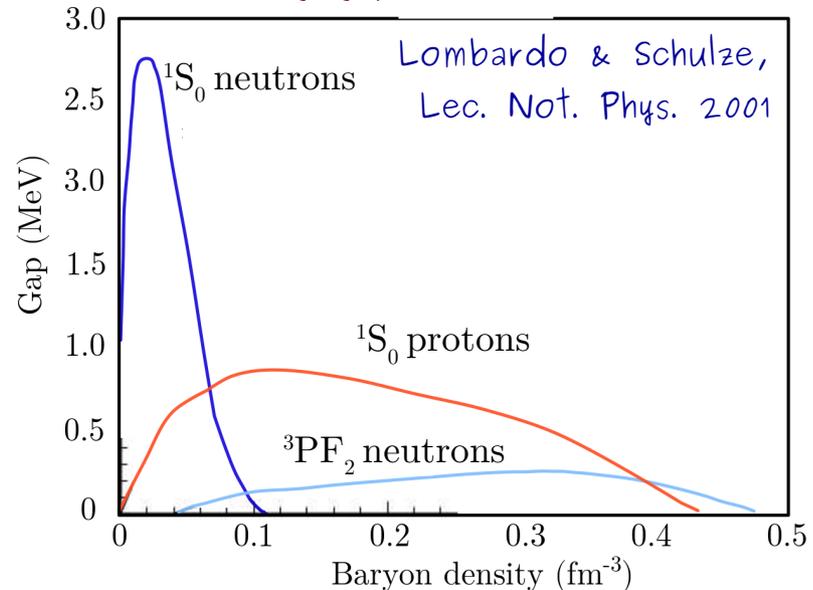
Uncertain gap, usually treated as free **parameter** in cooling simulations.



Transition temperature

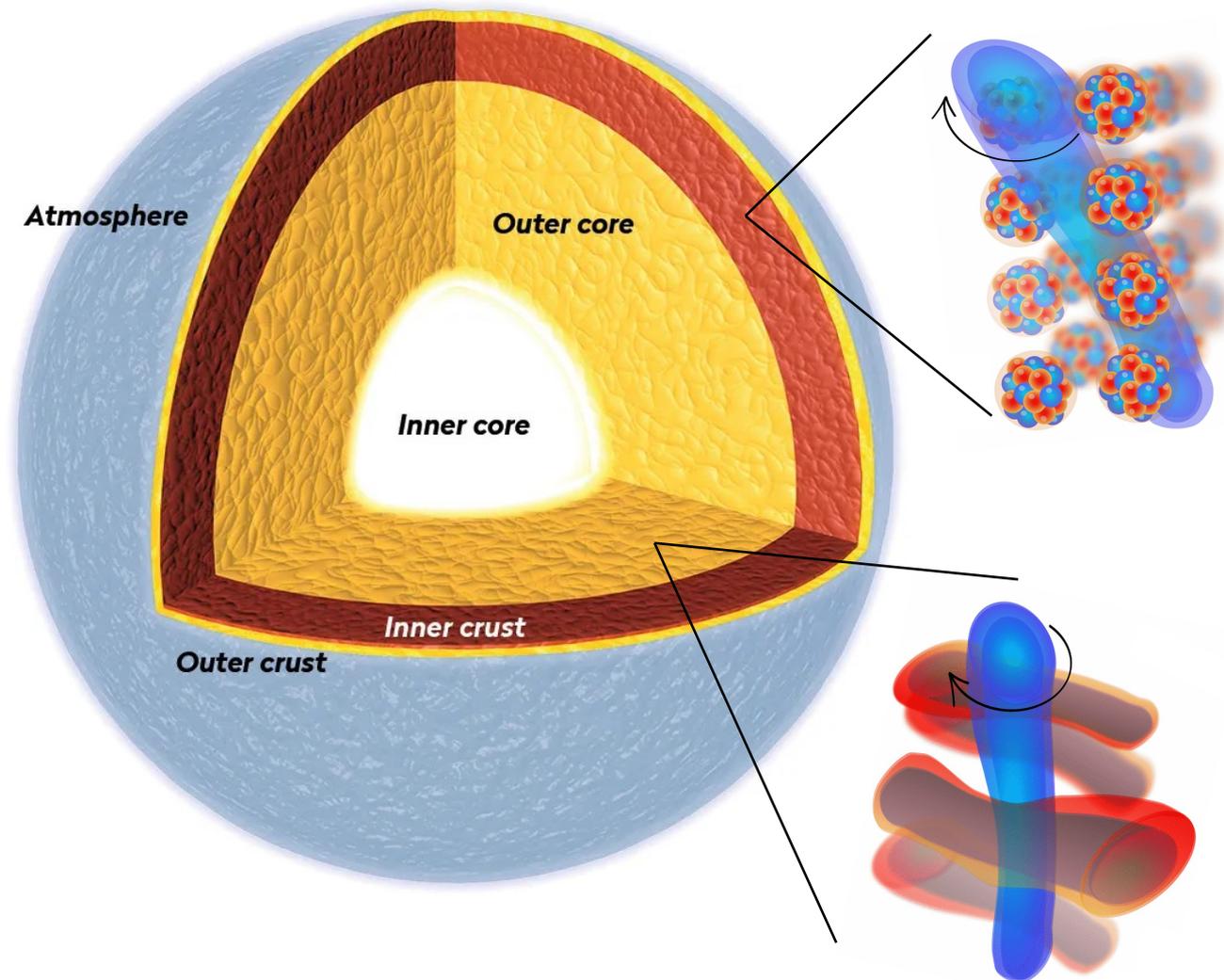


Pairing gap  $\Delta \sim 1 \text{ MeV} \sim 10^{10} \text{ K}$



# Continuum description?

Not a simple fluid but many layers with different inhomogeneities, defects, currents



**Inner crust**  
(Visco-)elastic lattice  
Neutron “scalar” superfluid  
Excitations (entropy/heat)  
Ideal gas of electrons  
B field  
Frame of the lattice: 3 currents

**Outer core**  
Neutron superfluid  
Proton superconductor  
Excitations (entropy/heat)  
Ideal gas of electrons  
B field in fluxtubes (type II?)

Different phenomena may be well described with a smaller number of “fluids”

e.g. cooling, glitches, oscillation modes, “mountains”, B evolution...  
...it is however true that the “magneto-thermo-rotational” evolution is coupled

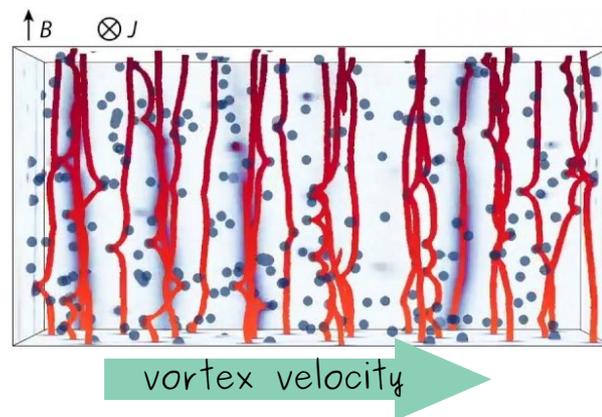
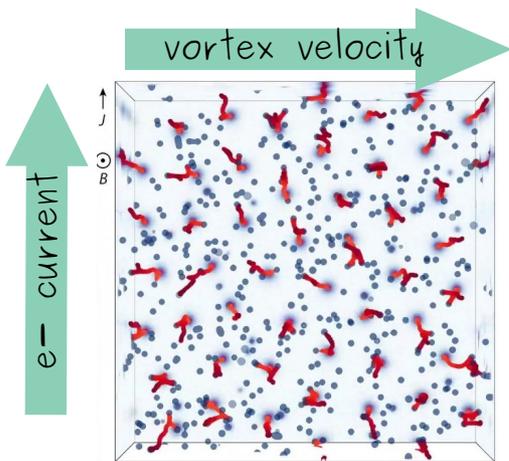
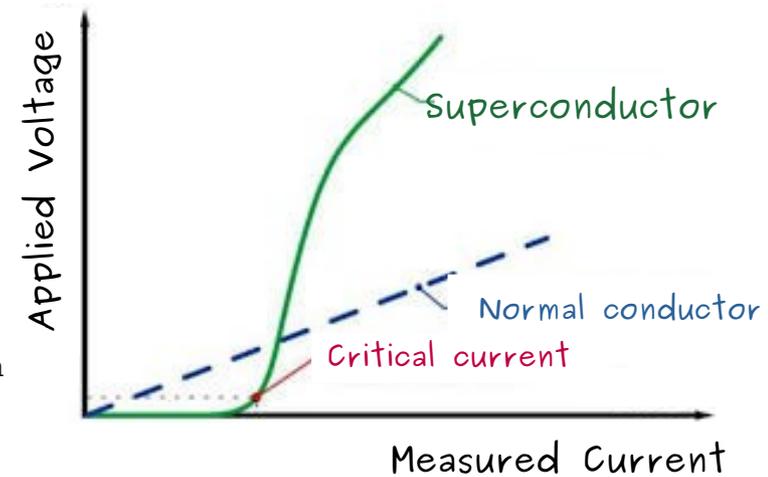
# Role of vortices in the continuum description?

Many coexisting species and independent currents, but **what is the role of the topological defects?**

Hint: consider laboratory type-II superconductors

Vortex-defect interaction associated with a **critical current** in the V-I relation

The slope is the resistance:  
Flat region  $\rightarrow$  no dissipation



Willa et al. (2018)

## Nutron Stars – Lab Superconductor

Crust's frame of reference – Frame of the lattice

Delocalized neutrons – Electron sea

Vorticity – magnetic field

Vortices – Flux-tubes

Neutron current – electric current

Magnus force – Lorentz force

*Conservative regime:*

Pinned vortices – pinned flux-tubes

*Dissipative regime:*

Moving vortices – Moving flux-tubes

# Superfluid hydrodynamics

Relativistic formulation: [Gavassino, MA](#)

[ArXiv: 2012.10288, 2001.08951](#)

Review: [arXiv:2301.12769 \(2022\)](#)



**Vortex core scale:**  
 $\sim 10$  fm in a NS

**Microscopic** models needed:

**Helium:** stochastic GPE, mean field...

**NS:** phenomenological Ginzburg-Landau+GPE, TDLDA, HFB...

Conserved baryon charge:

$$\# \text{ Baryons} \sim n_0 (10 \text{ fm})^3 \sim 10^3$$



**Inter-vortex scale:**  
 $\sim 10^{-3}$  cm in a NS

**Mescoscopic** models needed:

**Helium:** Vortex filament model  
(K. Schwarz's 80s papers)

**NS:** the same but needs extension to non-homogeneous environment

Conserved baryon charge:

$$\# \text{ Baryons} \sim n_0 (10^{-3} \text{ cm})^3 \sim 10^{29}$$



**Fluid element:**  
from  $\sim mm$  in a NS

**Macroscopic** hydro:

**Helium:** HVBK hydro  
(Hall & Vinen 1956)

**NS:** extensions/decorations of HVBK hydro (more species, GR)

Conserved baryon charge:

Not so relevant, need to include many vortices in the fluid element

# Superfluid hydrodynamics

Relativistic formulation: [Gavassino, MA](#)

[ArXiv: 2012.10288, 2001.08951](#)

Review: [arXiv:2301.12769 \(2022\)](#)



**Vortex core scale:**

$\sim 10$  fm in a NS



**Inter-vortex scale:**

$\sim 10^{-3}$  cm in a NS



**Fluid element:**

from  $\sim mm$  in a NS

We can not take into account each vortex ( $\sim 10^{16}$  in a pulsar)  $\rightarrow$  HVBK hydrodynamics

2 Euler-like equations + **entrainment** + **mutual friction**

$$\partial_t \rho_x + \nabla_i (\rho_x v_x^i) = 0$$

$$(\partial_t + v_x^j \nabla_j) (v_i^x + \varepsilon_x w_i^{yx}) + \nabla_i (\tilde{\mu}_x + \Phi) + \varepsilon_x w_{yx}^j \nabla_i v_j^x = f_i^x / \rho_x$$

The **dynamics of vortices** in a fluid element gives the form and strength of **mutual friction**

$x, y = n \rightarrow$  superfluid neutrons     $x, y = p \rightarrow$  normal component (electrons, excitations, protons, lattice...)

# Superfluid hydrodynamics

Relativistic formulation: [Gavassino, MA](#)

[ArXiv: 2012.10288, 2001.08951](#)

Review: [arXiv:2301.12769 \(2022\)](#)



**Vortex core scale:**  
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$$p_x^i = m_x [(1 - \epsilon_x)v_x^i + \epsilon_x v_y^i]$$

Hydrodynamic canonical momenta

Uniform matter, external force  
applied to a single species  
(e.g. the neutrons)

$$\mathbf{F}_n^{ext} = n_n \dot{\mathbf{p}}_n = \rho_n \dot{\mathbf{v}}_n$$

No entrainment

$$\mathbf{F}_n^{ext} = n_n \dot{\mathbf{p}}_n = \rho_{nn}^* \dot{\mathbf{v}}_n + \rho_{pn}^* \dot{\mathbf{v}}_p$$

With entrainment

Entrainment is a non-dissipative coupling (but you can not see this at this level, entropy is needed)

# Superfluid hydrodynamics

Relativistic formulation: [Gavassino, MA](#)

[ArXiv: 2012.10288, 2001.08951](#)

Review: [arXiv:2301.12769 \(2022\)](#)



**Vortex core scale:**  
~ 10 fm in a NS



**Inter-vortex scale:**  
~  $10^{-3}$  cm in a NS



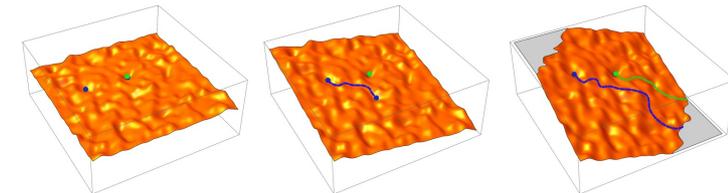
**Fluid element:**  
from ~mm in a NS

We can not take into account each vortex ( $\sim 10^{16}$  in a pulsar)  $\rightarrow$  HVBK-like hydrodynamics

2 Euler-like equations + **entrainment** + **mutual friction**

$$\begin{aligned} \rho_n D_t \mathbf{v}_n + \dots &= \mathbf{F}_{MF} \\ \rho_p D_t \mathbf{v}_p + \dots &= -\mathbf{F}_{MF} \end{aligned}$$

The **dynamics of vortices** in a fluid element gives the form and strength of the macroscopic “**mutual friction**”



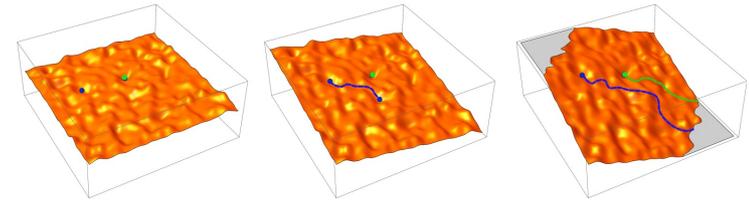
$x, y = n \rightarrow$  superfluid neutrons     $x, y = p \rightarrow$  normal component (electrons, excitations, protons, lattice...)

# Vortex motion $\rightarrow$ Mutual friction (with pinning)

Antonelli & Haskell, arXiv:2007.11720 (2020)

It's a "kinetic approach" but with point vortices instead of particles

- Fix a background "lag" (background current of superfluid neutrons)
- Assign random position of a vortex in the pinning landscape and solve the trajectory



disordered pinning force field

$$\hat{\mathbf{k}} \times (\dot{\mathbf{x}}(t) - \mathbf{v}_{np}) - \mathcal{R} \dot{\mathbf{x}}(t) + \mathbf{f} = 0$$

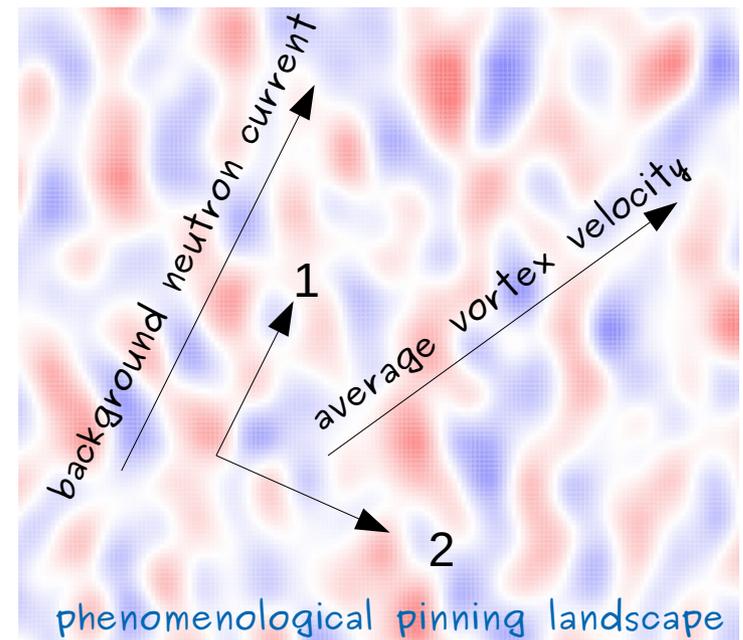
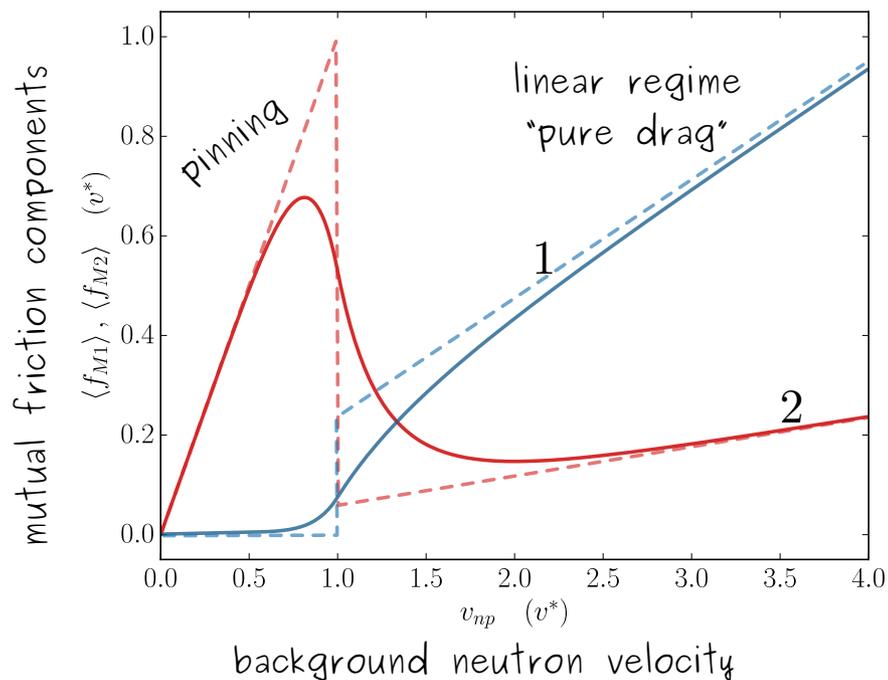
- Repeat many times and find the average vortex velocity for the given "lag"

average vortex velocity

- The mutual friction is given by

$$\mathbf{F}_n = -\kappa n_v \hat{\mathbf{k}} \times (\langle \dot{\mathbf{x}} \rangle - \mathbf{v}_{np})$$

assigned "lag"



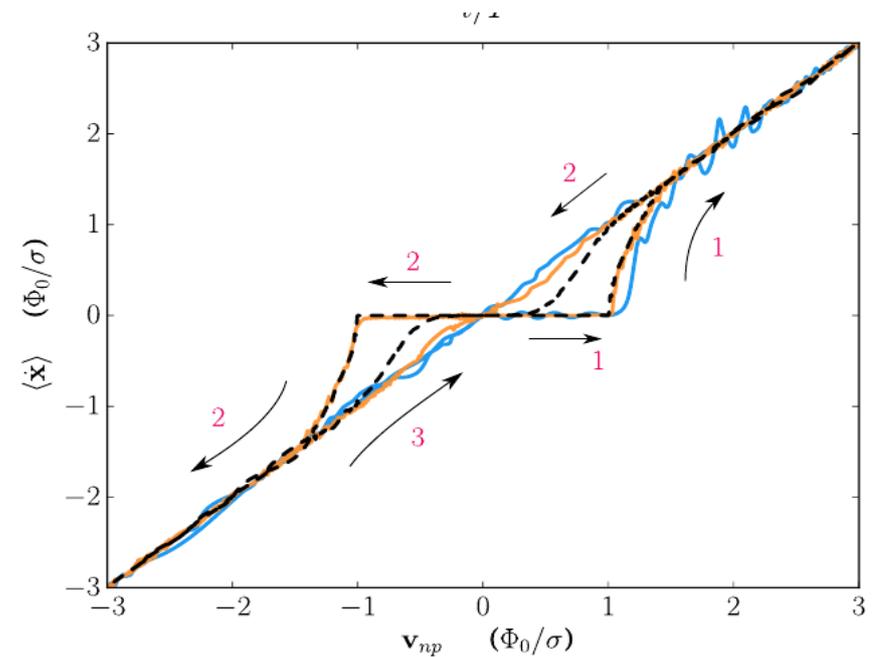
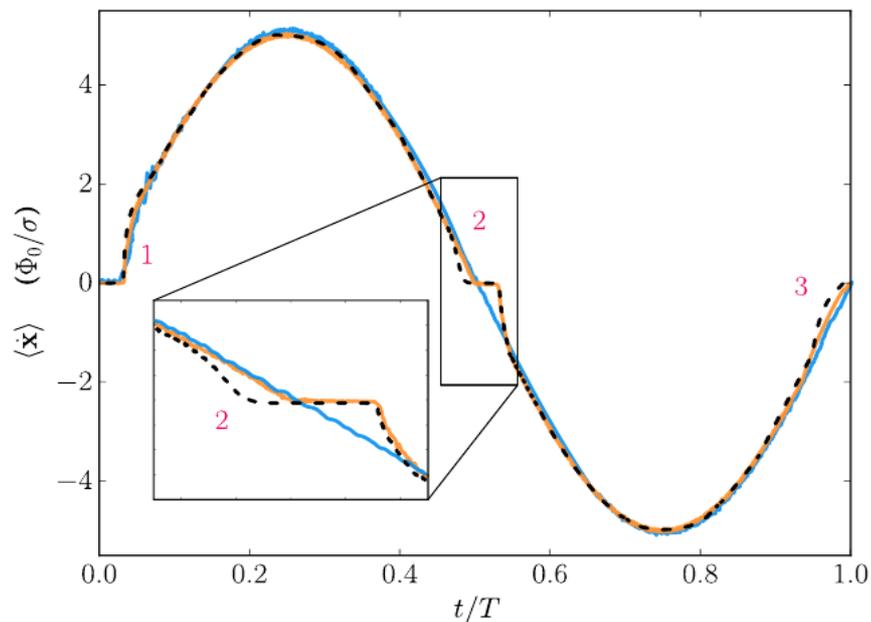
# Beyond hydrodynamics: hysteresis

**Rate-dependent hysteresis:** lag between an input and an output that disappears if the input is varied more slowly. If the input is reduced to zero, the output continues to respond for a finite time.

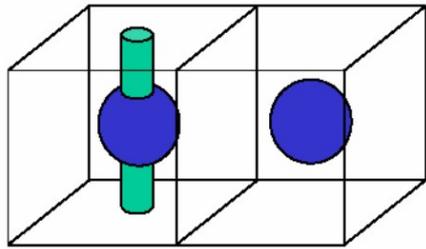
Instantaneous drop to null lag  $\rightarrow$  vortex velocity drops to zero immediately (if NO pinning forces)

Instantaneous drop to null lag  $\rightarrow$  vortex velocity relaxes to zero (with pinning forces)

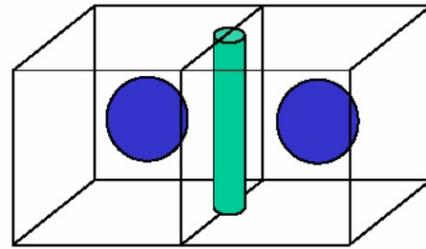
**Rate-independent hysteresis**  $\rightarrow$  vortex-vortex interactions



# Pinning energy (vortex – single nucleus)



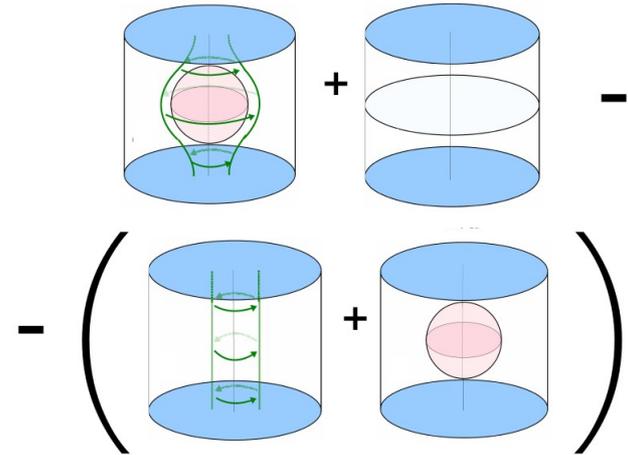
Nuclear pinning



Interstitial pinning

Donati & Pizzochero, Phys Lett B, 640 (2006)

**Semiclassical approach:** static LDA calculation (local Fermi momentum is a function of the neutron number density)



Klausner et al, arXiv:2303.18151 (2023)  
Hartree-Fock-Bogolyubov

Energy contributions to pinning:

- negative condensation energy of the order of  $\Delta^2 / E_F$
- kinetic energy of the irrotational vortex-induced flow
- Fermi energy  $E_F$  of neutrons
- nuclear cluster energy (Woods-Saxon potential)

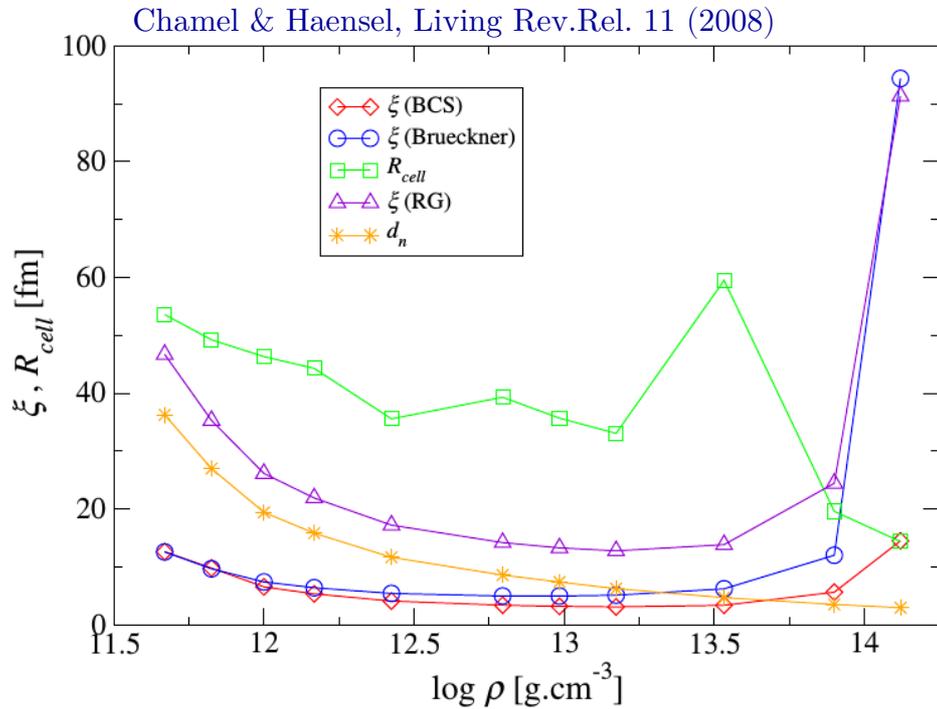
**Uncertain pairing gap  $\Delta$ :** modifies the strength and location of the pinning energies (in-medium effects!)

Maximum pinning energies  $< 3.5$  MeV

Significant pinning occurs only in a restricted range:  $0.07 n_0 < n_B < 0.2 n_0$

Improvements: TD-LDA, Wlazłowski+ (2016), Hartree-Fock-Bogolyubov Klausner+ (2023)

# Pinning forces (inner crust)



Coherence length  $\xi \sim$  vortex core radius

**Strong pinning** when  $\xi <$  lattice spacing

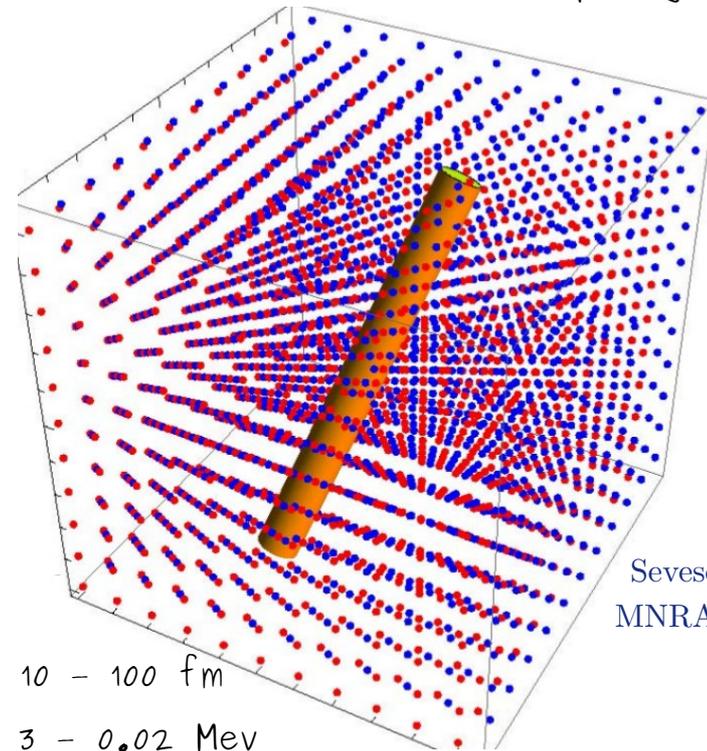
Pinning to **single defects** VS “**collective pinning**”:

Rigid (straight) vortices are “less pinned”

Coherence length  $\xi$  estimates: Mendell, ApJ 38 1991

Lattice spacing: 50-10 fm

Lattice spacing: 50-10 fm



Seveso et al  
MNRAS 2016

$$\xi = 10 - 100 \text{ fm}$$

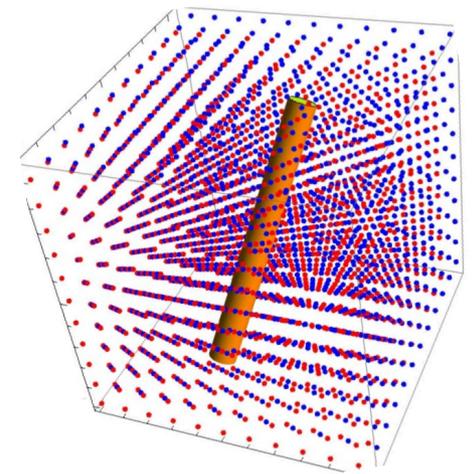
$$E_{\text{pin}} = 3 - 0.02 \text{ Mev}$$

**Inner crust:**

**Problem:** how to calculate the “vortex-lattice” interaction from the “vortex-nucleus” interaction ?

IDEA: consider a segment of vortex line (the length  $L$  is given by the tension) and average over translations and rotations of the total pinning force divided by  $L$

# Stationary test (glitch amplitude)



We can **test** theoretical single-vortex pinning forces with glitches of large amplitude:

- Choose the EOS
- Choose the pinning forces (function of the baryon density)
- Solve the hydrostatic equilibrium in slow rotation (“Hartle”) and calculate:

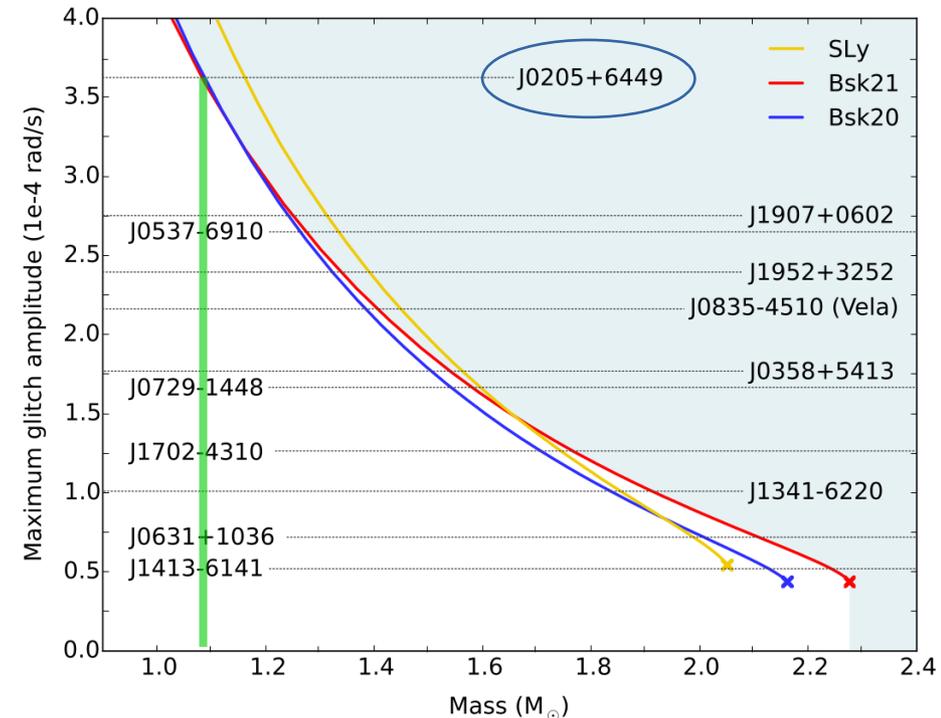
$$\Delta\Omega_{\text{abs}} = \frac{\pi^2}{I\kappa} \int_0^{R_d} dr r^3 e^{\Lambda(r)} \frac{\mathcal{E}(r) + P(r)}{m_n n_B(r) c^2} f_P(r)$$

- Generous theoretical **upper bound** based only on angular momentum balance
- No superfluid fraction and no entrainment
- No need to solve the internal dynamics
- No dependence on the assumed vortex configuration

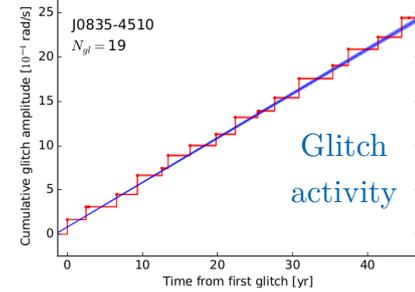
Compare with bounds on the **minimum** mass:

Observed:  $M=1.174 M_{\odot}$  [Martinez+2015 arXiv:1509.08805](#)

CCS simulations:  $M \approx 1.15 M_{\odot}$  [Suwa+2017 arXiv:1808.02328](#)



# Stationary test (glitch activity)

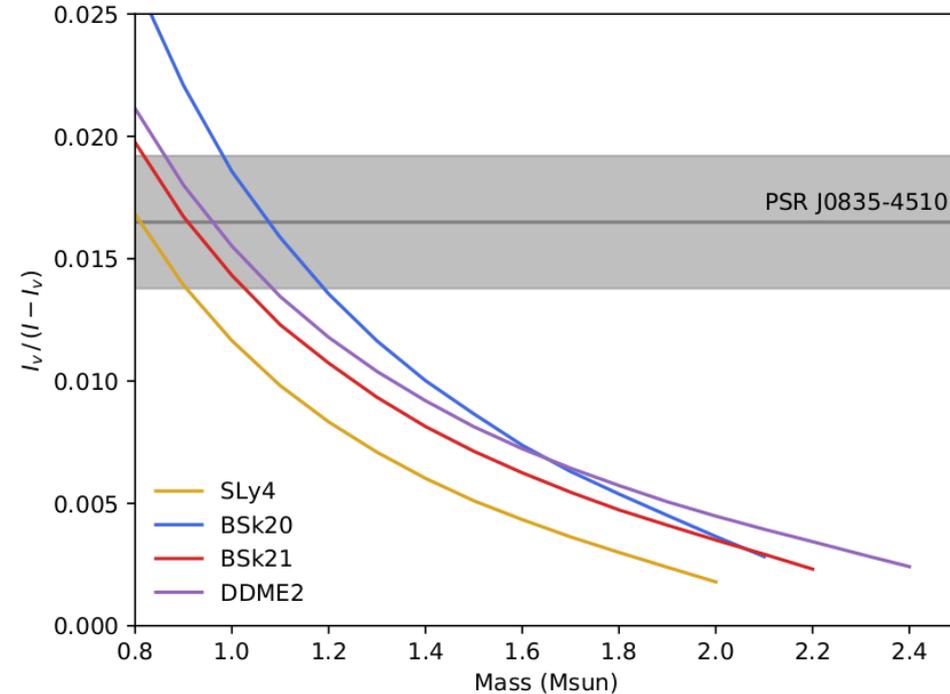


Result: the *superfluid in the crust is not enough* to explain Vela’s activity when strong entrainment is accounted for [Andersson+2012](#) arXiv:1207.0633, [Chamel 2013](#) arXiv:1210.8177

Revised argument: [Montoli+2020](#) arXiv:2012.01539

$$I_v = \frac{8\pi}{3} \int_0^{R_d} dr r^4 \frac{\rho_n(r)}{1 - \epsilon_n(r)} \quad \frac{\mathcal{A}}{|\dot{\Omega}|} < \frac{I_v}{I - I_v}$$

Note: no need to solve the internal dynamics (“stationary”)



“Heteroscedastic” linear regression: uncertainties larger by a factor  $\sim 10$  for the Vela ([Montoli+2020](#) arXiv:2012.01539)

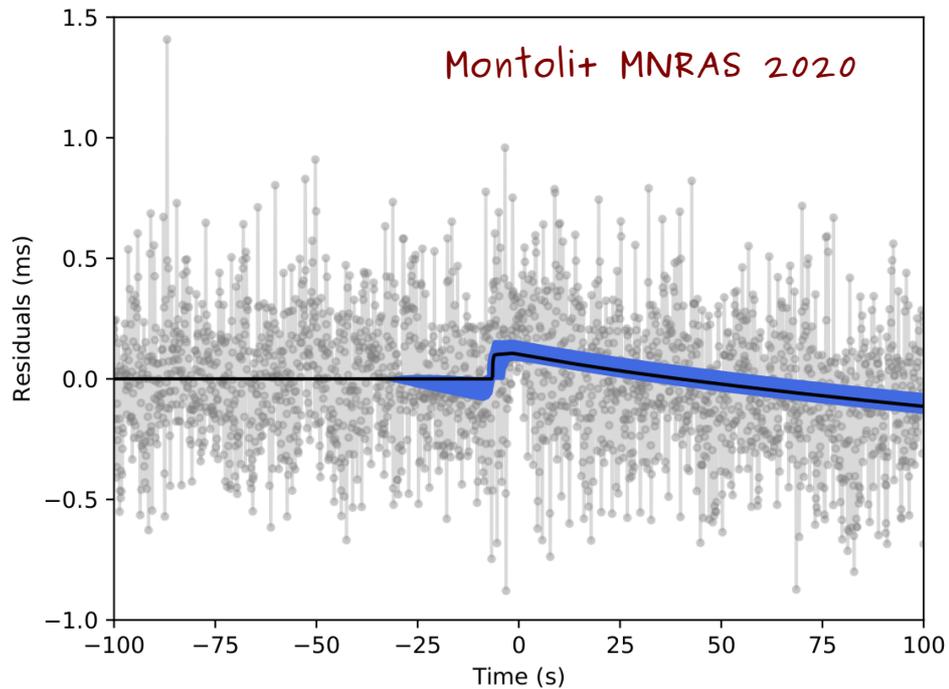
$$\mathcal{A}_a = \frac{\sum_i \Delta\Omega_i}{\sum_i \Delta t_i} = \frac{\sum_i \Delta\Omega_i}{t_{N_{gl}-1} - t_0} \quad \text{Var}(\mathcal{A}_a) = \frac{1}{(N_{gl} - 1)(t_{N_{gl}-1} - t_0)} \sum_i \frac{(\Delta\Omega_i - \mathcal{A}_a \Delta t_i)^2}{\Delta t_i}$$

Unbiased estimators for cumulated data

Compare with bounds on the minimum mass of a NS:

Observed:  $M=1.174 M_\odot$  [Martinez+2015](#) arXiv:1509.08805 CCS simulations:  $M \approx 1.15 M_\odot$  [Suwa+2017](#) arXiv:1808.02328

# Bayesian fit of Vela 2016



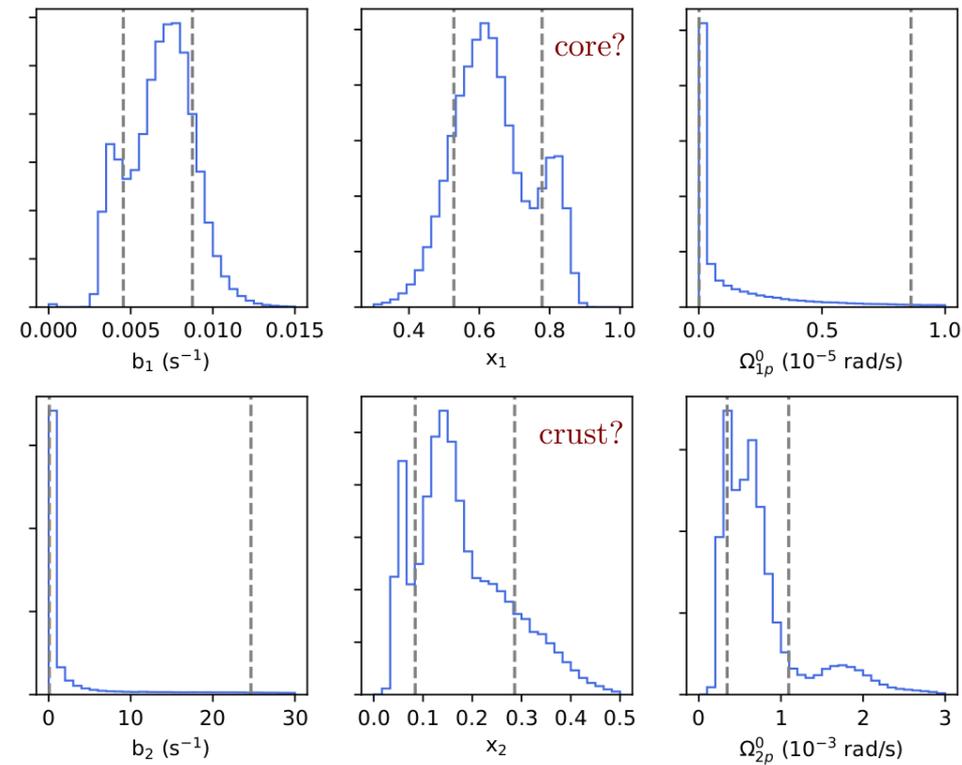
Normal  $x_p \dot{\Omega}_p = -\mathcal{T}_1 - \mathcal{T}_2 - |\dot{\Omega}_\infty|$

Super (core?)  $x_1 \dot{\Omega}_1 = \mathcal{T}_1$

Super (crust?)  $x_2 \dot{\Omega}_2 = \mathcal{T}_2,$

$$\mathcal{T}_i = -x_i b_i (\Omega_i - \Omega_p) \quad i = 1, 2$$

Posteriors for the **physical** parameters



Fit of the TOA residuals of [Palfreyman+2018](#) with a 3-component model:

→ Estimated **moment of inertia fractions** ( $x_2 + x_1 < 1$ ):

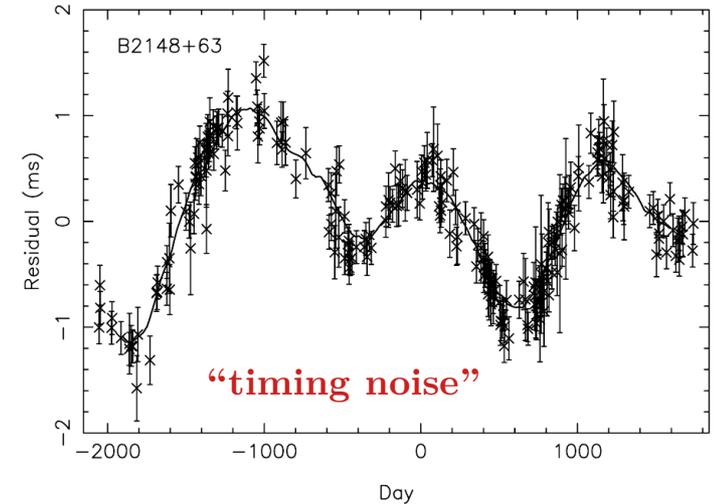
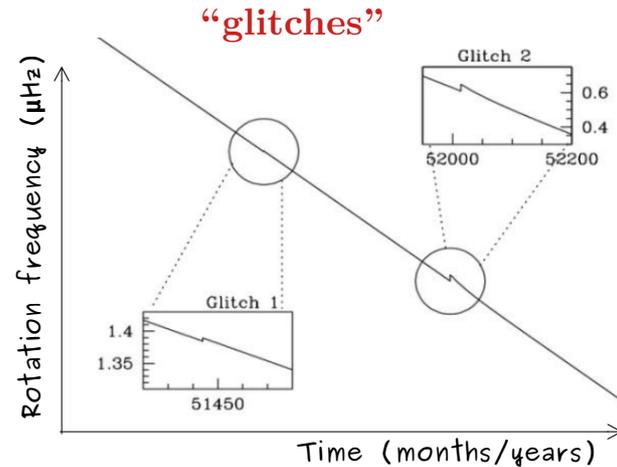
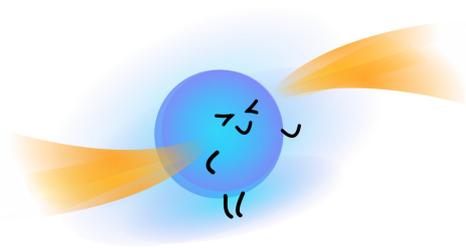
“active”  $x_2 \sim 0.1 - 0.3$ , “passive”  $x_1 \sim 0.5 - 0.7$

→ Confirmed **overshoot** found by [Ashton+2019](#), [Pizzochero+2020](#)

→ Posteriors for **friction parameters** in agreement with the revised estimates of friction by [Graber+2018](#)

→ **First “clue”** that the superfluid in the **outer core** is pinned before the glitch, regardless of entrainment strength in the inner crust

# Pulsar “timing noise”



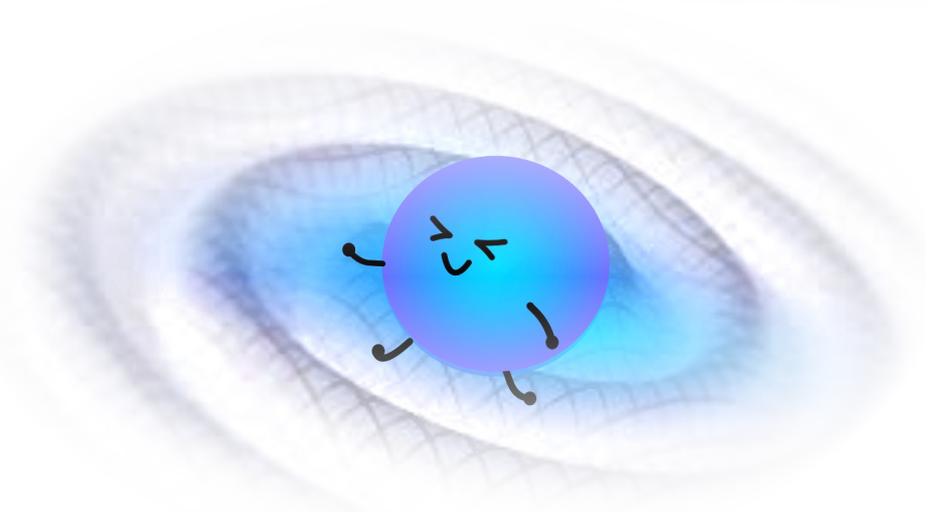
*Antonelli, Basu, Haskell, MNRAS (2023) - arXiv:2206.10416*

→ we import physical ideas from “glitch theory” to model “timing noise”

→ understanding the properties of timing noise is important also for Virgo/Ligo

**NS emitting *continuous* GW is a  
“gravitational pulsar”**

→ need to keep track of pulsar timing (including  
glitches and noise) for “targeted” and “narrow-band”  
searches of continuous GW @Ligo/Virgo



# Stochastic process for timing noise

Non-superfluid component  
(coupled to magnetosphere)

$$x_p \dot{\Omega}_p = -\mathcal{T} - \overset{\text{fluctuating}}{\eta_t^{\mathcal{T}}} + \mathcal{T}_\infty + \overset{\text{fluctuating}}{\eta_t^\infty}$$

Independent standard  
Wiener processes

$$\eta_t^\infty = \sigma_\infty \dot{W}_t^\infty$$

$$\eta_t^{\mathcal{T}} = \sigma_{\mathcal{T}} \dot{W}_t^{\mathcal{T}}$$

Superfluid component  
(not directly observable)

$$x_1 \dot{\Omega}_1 = \mathcal{T} + \overset{\text{fluctuating}}{\eta_t^{\mathcal{T}}} \overset{\text{fluctuating}}{\eta_t^{\mathcal{T}}}$$

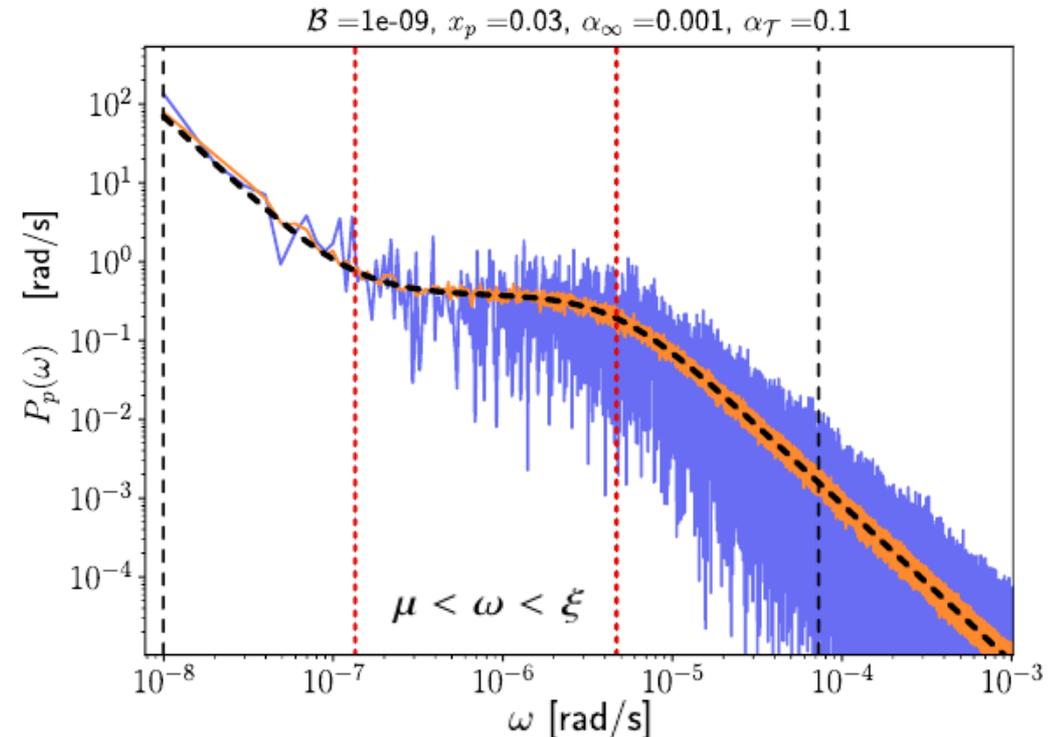
$$\langle \dot{W}_t^i \dot{W}_s^j \rangle = \delta_{ij} \delta(t-s)$$

Only **external** fluctuations: flat region in the PSD shrinks  
→ pure Wiener process + “inertial” corrections

$$P_p(\omega) = \frac{\alpha_\infty^2 \Omega \dot{\Omega}}{x_p^2 \omega^2} \cdot \frac{\omega^2 + x_p^2 \tau^{-2}}{\omega^2 + \tau^{-2}} \approx \frac{\alpha_\infty^2 \Omega \dot{\Omega}}{\omega^2}$$

Only **internal** fluctuations: flat region extends to the origin  
→ typical Lorentzian PSD (Ornstein-Uhlenbeck process)

$$P_p(\omega) = \frac{\alpha_{\mathcal{T}}^2 x_1^2 \dot{\Omega}^2}{2 x_p^2 \mathcal{B} \Omega} \cdot \frac{1}{\omega^2 + \tau^{-2}}$$



# Final considerations

Glitches provide us with some interesting theoretical challenges:

...thank you spinning pulsar!

- vortex dynamics in non-homogeneous environments
- collective avalanche dynamics
- how to describe pinning at the microscopic scale?
- physics in the core: superfluid-superconductor mixture

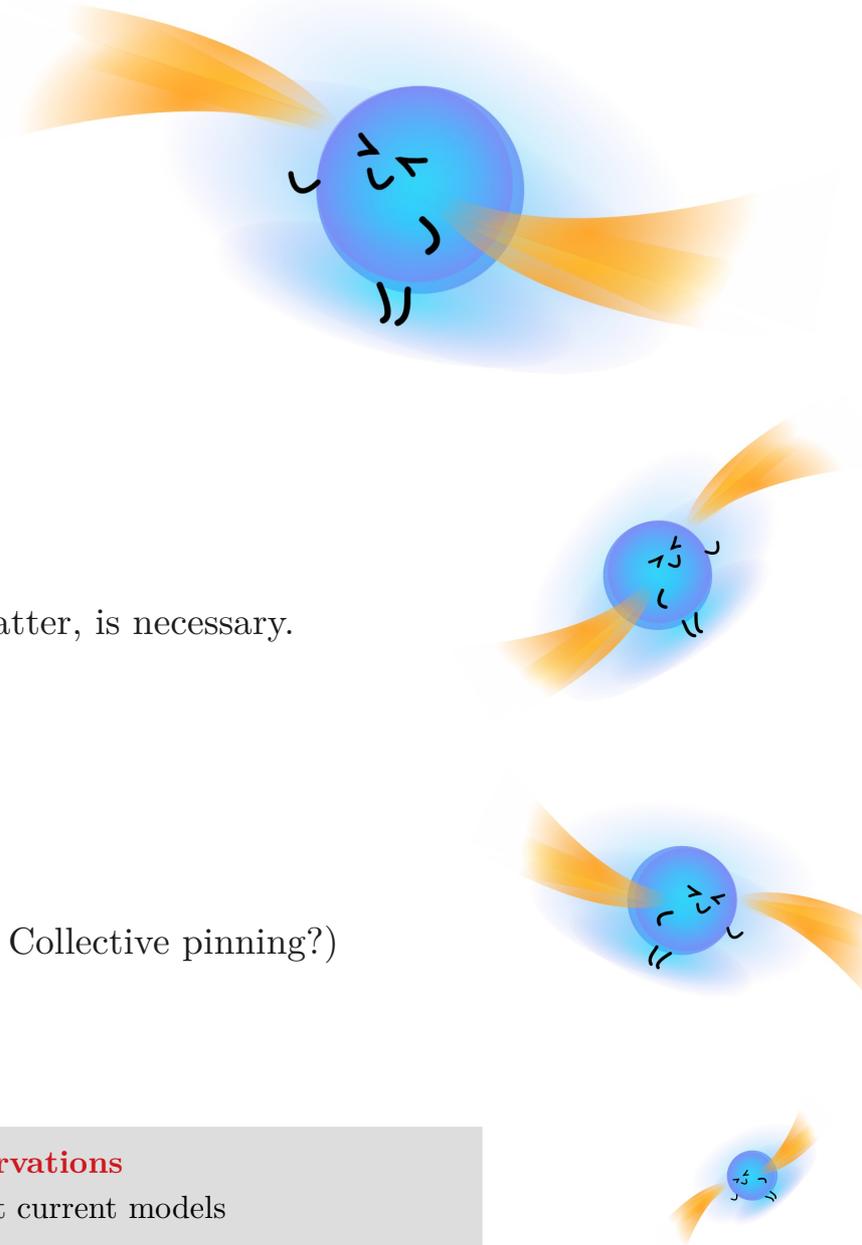
Cross contamination between different fields, especially condensed matter, is necessary.

Some open questions:

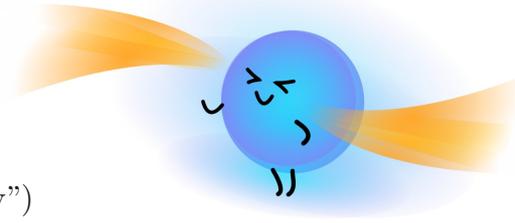
- role of starquakes? (can we really have quakes in a NS?)
- role of entrainment (strong/weak? affected by disorder?)
- better understanding of dissipation at micro/meso scale
- collective aspects of vortex dynamics (Viscoelasticity? Hysteresis? Collective pinning?)

**The most important thing: more and better observations**

Improved timing techniques (and more observation time) → test current models



# (some) References



**Stationary “activity” test** (morally: “effective moment of inertia of pinned region” > “observed activity”)

Link+1999 arXiv:9909146, Andersson+2012 arXiv:1207.0633, Chamel 2013 arXiv:1210.8177, Montoli+2020b arXiv:2012.01539

**Stationary “glitch size” test** (morally: “glitch from maximal critical lag” > “observed glitch amplitude”)

Antonelli+2018 arXiv:1710.05879

**Beyond the stationary tests** (need to solve temporal evolution  $\rightarrow$  model dependent):

Ho+2015 arXiv:1703.00932 (activity+cooling: need to integrate the thermal evolution), Pizzochero+2017 arXiv:1611.10223, Montoli+2020 arXiv:1809.07834 (glitch size+activity: need to integrate the superfluid reservoir evolution)

**General Relativity corrections:** Sourie+2017 arXiv:1607.08213 (2 rigid components), Antonelli+2018 arXiv:1710.05879 (fluid, corrections for the glitch size), Gavassino+2019 arXiv:2001.08951, arXiv:2012.10288 (fluid, corrections for the spin-up timescale)

**Revisiting the starquake paradigm:** Giliberti+2019 arXiv:1902.06345, arXiv:1809.08542 (continuously stratified), Reconret+2021 arXiv:2106.12604 (“starquake is not enough” argument), Bransgrove+2020 arXiv:2001.08658 (quake and null pulses in Vela 2016)

**Glitch overshoot:** virtually any model with more degrees of freedom than Baym+ Nature 224 1969:

Fluid models that are in the same class of Alpar+ ApJ 273 1984: Haskell+2012 arXiv:1107.5295, Antonelli+2017 arXiv:1603.02838, Graber+2018 arXiv:1804.02706 (revised friction parameters)

Models with 3 rigid components: Pizzochero+2020 arXiv:1910.00066 (mathematical condition for the overshoot)

**Vela 2016** (Palfreyman+ Nature 556 2018)

Ashton+2019 arXiv:1907.01124 (agnostic Bayesian fit), Montoli+2020a arXiv:2005.01594 (theory of 3-component model, Bayesian fit of physical parameters), Sourie+2020 arXiv:2001.09668 (pinning in the core), Gügerçinoğlu+2020 arXiv:2003.08724 (fit with the vortex creep model)

**Beyond hydrodynamics:** we are still “stealing” from Anderson&Itoh Nature 256 1975 (short but very dense paper!)

Melatos+2018 arXiv:1809.03064 (correlations in glitches), Khomenko+2018 arXiv:1801.01413, Antonelli+2020 arXiv:2007.11720 (pinning/depinning transition, hysteresis loop), Howitt+2020 arXiv:2008.00365 (point vortex simulations), Haskell+2020 arXiv:2007.02748 (quantum turbulence), Carlin+2021 arXiv:2105.13588

# Entrainment coupling: crust and core

In the inner crust (lattice of ions & S-wave superfluid):

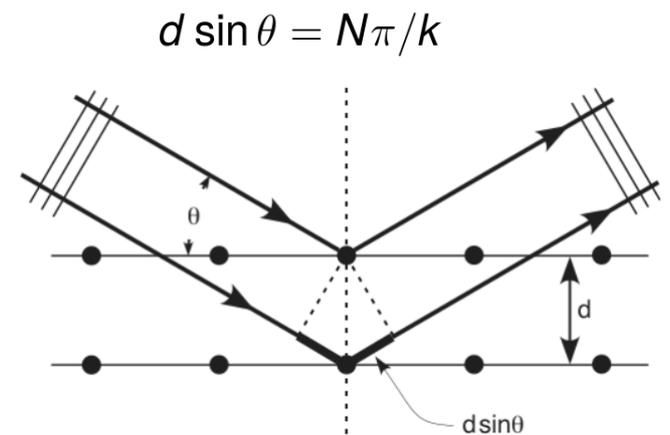
Chamel, PRC 2012

Bragg scattering by crustal lattice entrains the "free" neutrons.

Non-local effect:  $m^* > 1$

→ Consequence: the crustal superfluid is entrained by the normal component: reduced mobility of "free" neutrons is a potential problem for pulsar glitch theory.

Chamel PRL 2013, Montoli, Antonelli et al, Universe 2020



In the core (S-wave superconductor & P-wave superfluid):

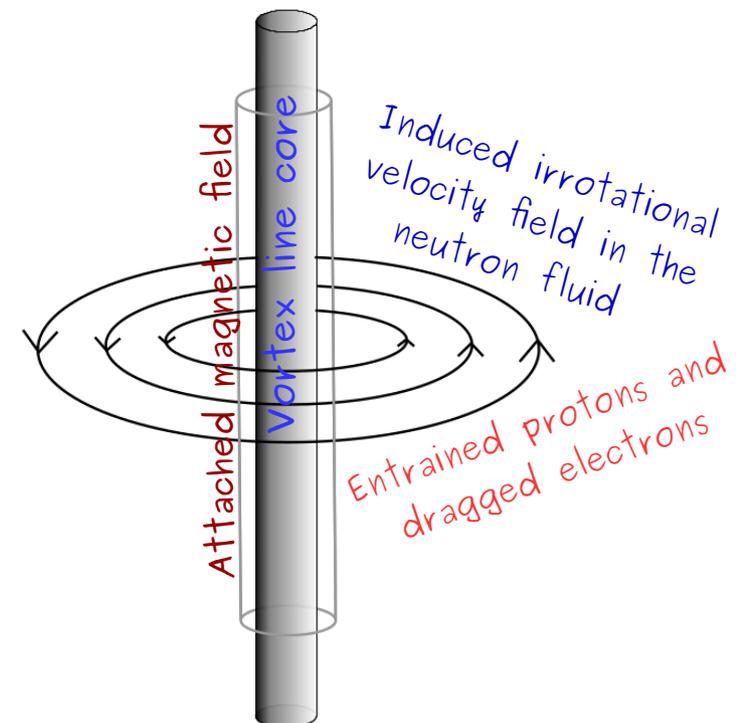
Chamel & Haensel PRC 2006

Entrainment is due to the strong interaction between protons and neutrons. Local effect:  $m^* < 1$

-Consequence #1: Scattering of electrons off vortex cores: the core is coupled to the crust on the timescale of a second

Alpar et al, ApJ 1984

-Consequence #2: Dipole-dipole interaction with flux-tubes (core pinning?)



# Pinning – Length scales

**Core** → “Abrikosov lattice” spacing between flux-tubes  $\sim 1000$  fm

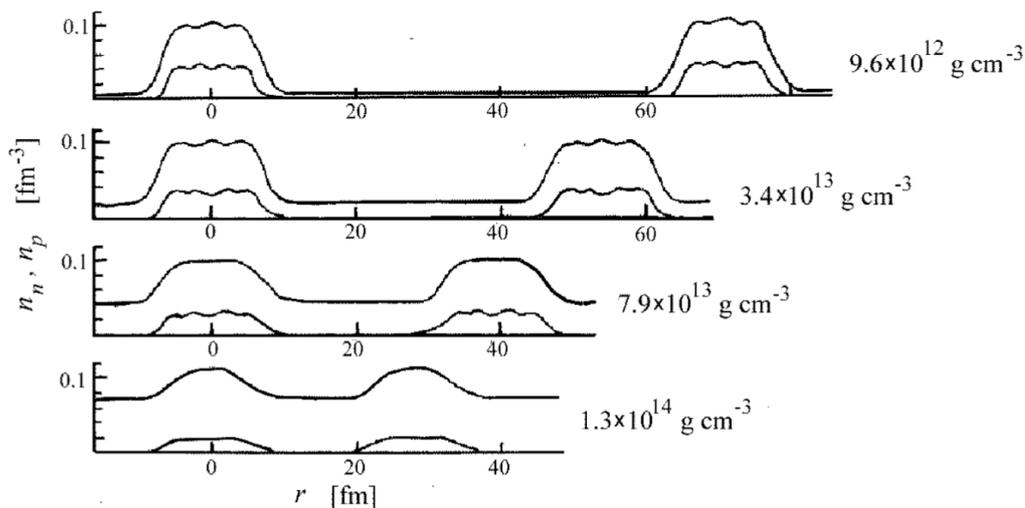
**Crust** → crustal lattice spacing  $\sim 100 - 20$  fm

Vortex-*nucleus* interaction → coherence length  $\sim 10 - 100$  fm

Vortex dynamics and vortex-*lattice* interaction → “mesoscale” (inter-vortex spacing)

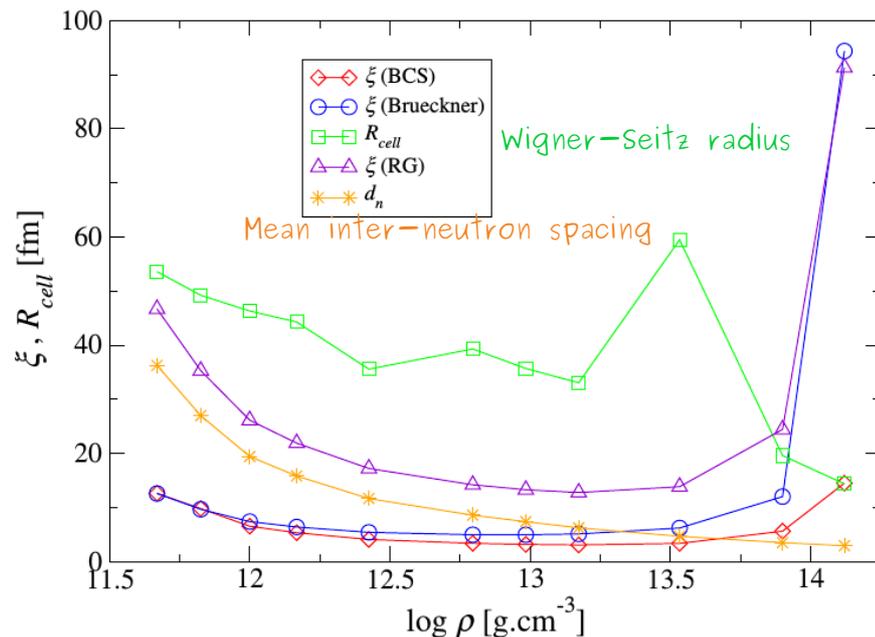
Inter-vortex spacing

$$l_v = \frac{\sqrt{\kappa P}}{2\pi} \approx 7 \times 10^{-3} \sqrt{P} \text{ cm}$$



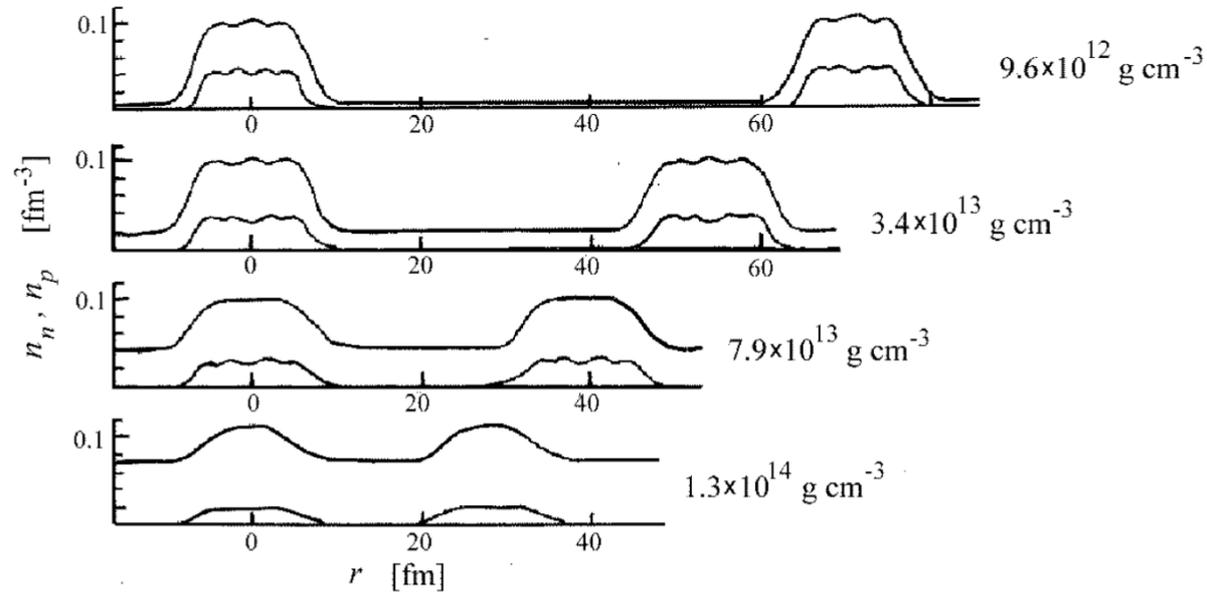
Negele & Vautherin (1973)

Neutron star matter at sub-nuclear densities



Chamel & Haensel, Living Rev. Rel. 11 (2008)

# Inner crust structure



Density profiles of neutron and protons, at several average densities, along a line joining the centers of two adjacent unit cells (HF calculation of the GS in the **inner crust** with effective NN interaction, **no pairing correlations**)

Negele & Vautherin, *Neutron star matter at sub-nuclear densities* (1973)

Include **pairing correlations**: Baldo et al, *The role of superfluidity in the structure of the neutron star inner crust* (2005)

Band theory of solids: Carter et al, *Entrainment Coefficient and Effective Mass for Conduction Neutrons in Neutron Star Crust* (2006)

# Heteroscedastic linear regression

Montoli+2020 arXiv:2012.01539

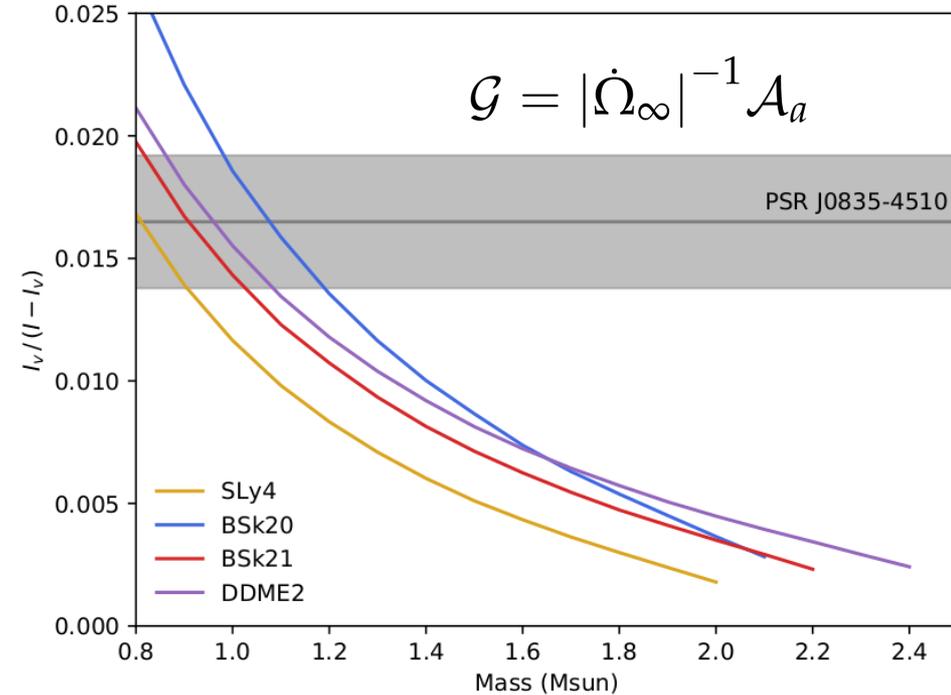
$$\Delta\Omega_i = \mathcal{A}_a \Delta t_i + \varepsilon_i \quad \Omega_i = \mathcal{A}_a t_i + \sum_{j=1}^i \varepsilon_j$$

Usual regression:  
deviations i.i.d.

Cumulated data:  
deviations are  
not i.i.d.

Important: careful inclusion of the “intercept”  
may lower the estimated uncertainty.

Pulsar	$\mathcal{G}_{\text{hom}}(\%)$	$\mathcal{G}_{\text{het}}(\%)$
0534+2200	$0.0079 \pm 0.0007$	$0.008 \pm 0.006$
0537-6910	$0.874 \pm 0.003$	$0.85 \pm 0.15$
0631+1036	$1.77 \pm 0.18$	$2.03 \pm 1.95$
0835-4510	$1.62 \pm 0.02$	$1.6 \pm 0.2$
1341-6220	$1.52 \pm 0.10$	$1.9 \pm 0.6$
1740-3015	$1.22 \pm 0.04$	$1.3 \pm 0.7$



$$\mathcal{A}_a = \frac{\sum_i \Delta\Omega_i}{\sum_i \Delta t_i} = \frac{\sum_i \Delta\Omega_i}{t_{N_{\text{gl}}-1} - t_0}$$

$$\text{Var}(\mathcal{A}_a) = \frac{1}{(N_{\text{gl}} - 1)(t_{N_{\text{gl}}-1} - t_0)} \sum_i \frac{(\Delta\Omega_i - \mathcal{A}_a \Delta t_i)^2}{\Delta t_i}$$

Unbiased  
estimators for  
cumulated data