

Talk at CEA September 2024  
Journées Itzykson.

1) What do we mean by probabilistic construction of CFT in 2d?

Correlations functions

$(\alpha, \gamma) = (\alpha_1, \dots, \alpha_m, \gamma_1, \dots, \gamma_n)$

The diagram shows a Riemann surface  $\Sigma$  with three points  $x_1, x_2, x_3$  marked. Above each point is a weight  $\alpha_i$ . The points are labeled  $(x_1, \alpha_1)$ ,  $(x_2, \alpha_2)$ , and  $(x_3, \alpha_3)$ .

$x_i \in \Sigma$  Riemann surface  
 $\alpha_i$  weight (say  $\in \mathbb{C}$ ).

$$\langle V_{\alpha_1}(x_1) \dots V_{\alpha_n}(x_n) \rangle_{\Sigma} = \mathbb{E} \left[ F_{(\alpha, \gamma)}(x) \right]$$

↑ primary field  
Vertex operator
↑ random variable

which satisfy the CFT axioms

in particular if  $g$  metric compatible with complex structure

$(g = e^{\sigma} |dz|^2 \text{ locally})$

$$\langle \prod_i V_{\alpha_i}(x_i) \rangle_{\Sigma, g e^{\sigma}} = \langle \prod_i V_{\alpha_i}(x_i) \rangle_{\Sigma, g} e^{-\frac{c}{24\pi} \int_{\Sigma} |dz|^2 g + 2K g \omega}$$

← central charge  
 Scalg.  
 Weyl anomaly.

\* physicist have formulas for certain models

Goal : ① link proba construction to explicit formulas

② Construct probabilistically Hilbert space  $\mathcal{H}$   
and a rep. of Virasoro  $(L_n)_n$  as operators  
in  $\mathcal{H}$  + decompose  $\mathcal{H}$  as sums of rep. of  
 $(L_n)$  (or Kac-Moody).

③ Prove bootstrap formulas

$$\langle \prod V_{\alpha_i}(n_i) \rangle = \int_{(\text{Spec})^N} \prod C_j(\alpha, p) \left( \sum_{\alpha, p} |n_i|^2 \right) dp$$

$$N = 3g_{\Sigma} - 3 + n$$

moduli space  
of complex st  
with marked pts.



prove convergence of conf blocks

Conformal blocks

Advantage :

- Crossing symmetry is for free
- Conv of blocks & of formulas in bootstrap should be consequence.
- Concrete objects, can be possibly related to scaling limits
- Application for new formulas (X. Sun & collab)

2) Examples that we know how to construct.

A] Liouville  $c_L > 25$

$$\langle \prod_{i=1}^n V_{\alpha_i}(n_i) \rangle = \int \prod e^{\alpha_i \phi(n_i)} e^{-S_{\Sigma}^L(\phi)} \mathcal{D}\phi$$

path integral

with  $S_{\Sigma}^L(\phi) = \frac{1}{4\pi} \int_{\Sigma} (d\phi)_g^2 + Q K_g \phi + p e^{\gamma \phi} dV_g$

where  $Q = \frac{c_L - 25}{2}$

$\gamma \in (0, 2)$ ,  $p > 0$

$$(DKRV, GRV, DRV)$$

$$\Phi = c + X_g \quad \begin{array}{l} \leftarrow \text{massless Gaussian} \\ \text{Free Field} \\ \leftarrow 0\text{-mode} \end{array}$$

$$\langle \prod_{i=1}^n V_{\alpha_i}(x_i) \rangle_{\Sigma, g} = \int \prod_{i=1}^n e^{\alpha_i(c + X_g(x_i))} e^{-\frac{1}{4\pi} \int_{\Sigma} \kappa_g(c + X_g) + p e^{\frac{c}{2}} e^{\frac{X_g}{2}} dV} \quad \text{GMC.}$$

Gaussian Multiplicative Chaos

$$e^{\frac{X_g}{2}} dV = \lim_{\varepsilon \rightarrow 0} \varepsilon^{\frac{\alpha^2}{2}} e^{\frac{X_\varepsilon}{2}} dV_g \quad (\text{Kahane, Robert-Vargas, Duplantier-Sheffield})$$

$$X_\varepsilon(x) = \int_{C(x, \varepsilon)} X_g(y) dy.$$

$C(x, \varepsilon) \rightarrow$  geodesic circle of radius  $\varepsilon > 0$

th: Converge if Seiberg bounds  $\sum_i \alpha_i - \chi(\Sigma) Q > 0$   
 $\& \alpha_i < Q \quad \forall i$   
 and satisfy Conformal covariance.  $c_L = 1 + 6Q^2$   
 $\Delta \alpha = \frac{\alpha}{2}(Q - \frac{\alpha}{2})$

## B) Compactified Imaginary Liouville

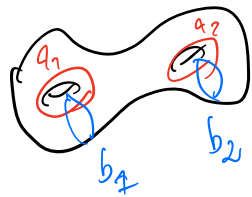
$$(GRV 23) \quad c_{IL} = 1 - 6Q^2$$

$$\langle \prod_j V_{\alpha_j}(x_j) \rangle_{\Sigma, g} = \int_{\text{Maps } \Sigma \rightarrow \mathbb{R}/2\pi\mathbb{Z}} \prod_j e^{i\alpha_j \Phi(x_j)} e^{-S_\Sigma^{IL}(\Phi)} D\Phi$$

$$S_\Sigma^{IL}(\Phi) = \frac{1}{4\pi} \int_{\Sigma} |\partial\Phi|_g^2 - iQ \kappa_g \Phi - p e^{i\beta \Phi} dV$$

$$Q = \beta/2 - \frac{2}{\beta} \quad \text{with } \beta \in (0, \sqrt{2}).$$

$$\text{We set } \Phi^{(n)} = c + X_g^{(n)} + I_{h^{(n)}} \quad \text{with } h \in \mathbb{Z}^{2g_\Sigma}$$



$$\in (0, 2\pi R)$$

GFF

topological term  
multi-valued but

$(a_i, b_i)_{i=1}^{g \in \mathbb{Z}}$  basis of  $H_1(\Sigma, \mathbb{Z}) = \mathbb{Z}^{2g}$

$dI_k =$  harmonic  
1-form on  $\Sigma$

Note:  $e^{i\alpha\phi(x)}$

univalued if

$$\alpha R \in \mathbb{Z}$$

with monodromy  
 $2\pi R k_2, \dots, 2\pi R k_{2g}$   
on  $(a_i, b_i)_i$

Probabilistic definition of correlations:

$$\left\langle \prod_j V_{\alpha_j}(x_j) \right\rangle_{\Sigma, g} =$$

$$\sum_{h \in \mathbb{Z}^{2g \in \mathbb{Z}}} e^{-\frac{\|dI_h\|_2^2}{4\pi}} \int_0^{2\pi R} \left[ \prod_j e^{i\alpha_j(c + x_j + I_h)(x_j)} - \frac{i\alpha}{4\pi} \int_{\Sigma} \gamma(c + x_j + I_h) - p e^{i\beta(c + x_j + I_h)} dv_j \right] dc$$

here  $p \in \mathbb{C}$  cosmological constant

$$e^{i\beta X_g} dv_g = \text{Imaginary GMC}$$

We need  $\beta R \in \mathbb{Z}$

(See: Lacaín - Rhodes - Vargas)

Th: if  $\frac{\beta^2}{4} = \frac{p}{q} \in \mathbb{Q}$ ,  $\alpha_j \in \frac{1}{R}\mathbb{Z}$ ,  $\alpha_j > 0$

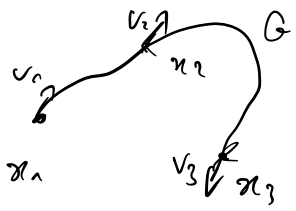
the correlations converge & it defines a CFT with central charge

$$c_{IL} = 1 - 6Q^2, \quad \Delta_\alpha = \frac{\alpha}{2} \left( \frac{\alpha}{2} - Q \right)$$

(Discrete spectrum).

in fact, we can add magnetic charges  $m_j \in \mathbb{Z}$  at  $\alpha_j$

with  $\sum_j m_j = 0$  by setting  $\phi = c + x_j + I_h + J_m$



Graph  $G = \text{cut}$   
for multivalued  
Fct°  $J_m$

$dJ_m = \text{Re}(\mathcal{O}_m)$   
mesomorphic 1-form  
with poles at  $x_j$

Vertex op.  $V_{d_i, m_i} \rightarrow$  Conformal weight  $\frac{\alpha_i}{2} \left( \frac{2i - \alpha}{2} \right) + \frac{m_i^2 R^2}{4}$

Spin:  $S_{\alpha, m} = \alpha R_m - \alpha R_m.$

Coulomb Gas

Formula:  $(V_{d_1, m_1}(x_1) V_{d_2, m_2}(x_2) V_{d_3, m_3}(x_3))_{\mathcal{G}}$

$$= (2\pi R) \frac{(-1)^l}{l!} \int_{\mathbb{C}^l} \prod_{j=2}^l \frac{d^2 z_j}{\pi} x_j^{\Delta_j - \bar{\Delta}_j} (1-x_j)^{\Delta_2} (1-\bar{x}_j)^{\bar{\Delta}_2} \times \prod_{j < k} \pi |z_j - z_k|^{2\beta} dz_1 \dots dz_l$$

$$l = \frac{2\alpha - \sum \alpha_j}{\beta} \in \mathbb{N}, \quad \Delta_j = \frac{\beta \alpha_j}{2} + \frac{\hbar m_j}{2}, \quad \bar{\Delta}_j = \frac{\beta \alpha_j}{2} - \frac{\hbar m_j}{2}$$

if  $m_j = 0 \Rightarrow = 2\pi R (-1)^l C_{\beta}^{\text{DZZ}}(\alpha_1, \alpha_2, \alpha_3)$

Imaginary DZZ -

In general:

$$= 4\pi^2 R^2 \rho^{\frac{2}{\beta} (2\alpha - \sum \alpha_j)} C_{\beta}^{\text{DZZ}}(\alpha_1 + m_1 R, \alpha_2 + m_2 R, \alpha_3 + m_3 R) \times C_{\beta}^{\text{DZZ}}(\alpha_1 - m_1 R, \alpha_2 - m_2 R, \alpha_3 - m_3 R)$$

### 3) Cutting, gluing, bootstrap.

The probabilistic approach is well designed to prove Segal axioms, where we cut the path integrals on surfaces into pieces (when the classical action and primary fields are local).

- To  $S^2$ , we attach a Hilbert space

$$\mathcal{H} = L^2(H^s(S^2), \mu_0)$$

Sobolev space of order  $s < 0$

$\mu_0$  is the law of the Gaussian free field

on plane  $\hat{\mathbb{C}}$  restricted to  $S^2 = \{z \in \mathbb{C} \mid |z|=1\}$

GFF on circle:

$$\varphi(\theta) = \sum_{n>0} \frac{x_n + iy_n}{2\sqrt{n}} e^{in\theta} + \sum_{n>0} \frac{x_n - iy_n}{2\sqrt{n}} e^{-in\theta} + c$$

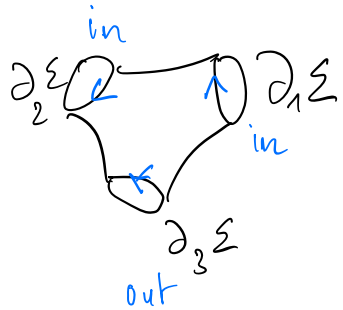
with  $x_n, y_n$  iid Gaussians  $N(0, 1)$

↑  
zero mode

$$\mu_0 = dc \otimes \prod_{n \geq 1} \frac{e^{-\frac{x_n^2 - y_n^2}{2}}}{2\pi} dx_n dy_n$$

- To surfaces with boundary  
( $\Sigma, g$ ) and parametrisations of  $\partial\Sigma$

$$\xi_1: \mathbb{S}^1 \rightarrow \partial_1 \Sigma, \dots, \xi_b: \mathbb{S}^1 \rightarrow \partial_b \Sigma$$



We associate an operator

$$\mathcal{A}_{\Sigma, g, \xi}: \mathcal{H}^{\otimes b_-} \rightarrow \mathcal{H}^{\otimes b_+}$$

$b_- = \#$  of incoming parametrized boundary circles

$b_+ = \#$  of outgoing parametrized boundary circles

$\mathcal{A}_{\Sigma, g, \xi}$  is described by some integral kernel on  $(H^{-s}(\Sigma))^b$ ,  $b = b_- + b_+$

$$\mathcal{A}_{\Sigma, g, \xi}(\underbrace{\varphi_+, \varphi_-}_{\varphi \in H^{-s}(\Sigma)^b}) := \left( \det \Delta_{\Sigma, g, D} \right)^{-\frac{1}{2}} \mathcal{A}_{\Sigma, g, \xi}^0(\varphi) \times \mathbb{E}_{\varphi} \left[ e^{-\frac{\alpha}{4\pi} \int_{\Sigma} K_g(X_{g, D} + P\varphi) + \varphi e^{\alpha X_{g, D} + \delta P\varphi} dv} \right]$$

↑  
conditional expectation ( $\varphi$  fixed)

with  $P\varphi =$  harmonic extension of  $\varphi$  in  $\Sigma$

$X_{g, D} =$  Dirichlet GFF.

$$\mathcal{A}_{\Sigma, g, \xi}^0(\varphi) = e^{-\frac{1}{2} \langle (D_{\Sigma} - D)\varphi, \varphi \rangle_{\Sigma}} \quad \text{Free Field amplitude}$$

$$Df(\theta) := \sum_n |n| f_n e^{in\theta} \quad \text{Fourier multiplier}$$

if  $f(\theta) = \sum f_n e^{in\theta}$

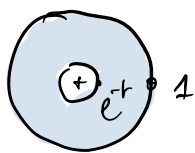
$$D_\Sigma f(\theta) = \partial_n \text{PB}_{1,2\Sigma} \quad \text{Dirichlet-to-Neumann operator}$$

Note:  $D_\Sigma - D$  is smoothing, i.e. it maps  $\mathcal{D}'(\mathbb{S}^1) \rightarrow C^\infty(\mathbb{S}^1)$

Th [GKR12]: Amplitudes are well defined and satisfy Segal gluing/cutting axioms and conformal invariance:

$$\mathcal{A}_{\Sigma, g, \xi}^{\omega}(\varphi) = e^{-\frac{c_L}{24\pi} \int_{\Sigma} (|\nabla\omega|^2 + 2K_g\omega)} \mathcal{A}_{\Sigma, g, \xi}(\varphi)$$

Hamiltonian of Liouville CFT:



$|A_{e^t}$  annulus  $|z| \in [e^{-t}, 1]$

$$g_A = \frac{|dz|^2}{|z|^2}$$

$$\int_1^t (e^{i\theta}) = e^{i\theta}$$

$$\int_{e^{-t}}^1 (e^{i\theta}) = e^{-t+i\theta}$$

$t \rightarrow \mathcal{A}_{|A_{e^t}, g_A, \xi^t} \in \mathcal{L}(\mathcal{H})$

is a contraction semi group.



generator  $\partial_t \mathcal{A}_{A_e^{-1} g_A, \xi^T} \Big|_{t=0} = H$

with  $H = -\frac{1}{2} \partial_c^2 + \frac{Q^2}{2} + \sum_n n (\partial_{x_n}^* \partial_{x_n} + \partial_{y_n}^* \partial_{y_n})$

$+ \rho e^{\delta c} V(\varphi)$

$V(\varphi) = \int_0^{2\pi} e^{\delta \varphi(\theta)} d\theta$

GMC on circle

$\in L^p(H^{-s}(S^1))$

$\forall p < \frac{2}{\delta^2}$

H has continuous spectrum  $[\frac{Q^2}{2}, \infty[$ .

Th [GKR21]:  $\exists$  an analytic family

$\alpha \mapsto \Psi_{\alpha, \nu, \tilde{\nu}}$ ,  $\forall \nu, \tilde{\nu}$  Young diagram.  
 $\mathbb{C}$  of eigenstates of H with eigenvalues

$2\Delta_{\alpha} + |\nu| + |\tilde{\nu}|$  with  $(\Psi_{\alpha+i\rho, \nu, \tilde{\nu}})_{\rho \in \mathbb{R}^+}$

$= 2(\alpha - \frac{\alpha}{2})$

being a complete family of eigenfunds in  $\mathcal{H}$

in the sense:  $\forall u, u' \in \mathcal{H}$ :

$$\langle u, u' \rangle_{\mathcal{H}} = \frac{1}{2\pi} \int_0^{\infty} \sum_{\substack{\nu, \tilde{\nu} \\ \nu, \tilde{\nu}'}} \langle u, \Psi_{\alpha+i\rho, \nu, \tilde{\nu}} \rangle_{\mathcal{H}} \langle \Psi_{\alpha+i\rho, \nu, \tilde{\nu}'} | u \rangle F_{\alpha+i\rho}(\nu, \nu') \times F_{\alpha+i\rho}(\tilde{\nu}, \tilde{\nu}') d\rho$$

$F_{\alpha, \nu}$  are Shapovalov coefficients  
 Moreover for  $\alpha < \alpha$ ,  $\mathcal{Y}_\alpha = \mathcal{A}_{\mathbb{D}, |dz|^2, (0, \alpha)}$   
 Disc amplitude

Virasoro representation:

take  $v(z)\partial_z = \sum_{n \in \mathbb{Z}} v_n z^{n+2} \partial_z$  hol. vech field near  $S^2$

s.t its flow  $br(z)$  maps  $S^2$  to  $br(S^2) \subset \mathbb{D}^0$   
 (we call  $v$  Markovian)

then if  $A_{br} = \mathbb{D} \setminus D_{br}$

$S^2$   $D_{br}$  "disc" bounded by  $br(S^2)$ .

Choosing well the metric  $g_r$  on  $A_{br}$   
 we prove with Berezin that

$\mathcal{A}_{A_{br}, g_r, \xi^t}$  is a "projective" semi group and

$$\partial_t|_{t=0} \mathcal{A}_{A_{br}, g_r, \xi^t} = \sum_n v_n L_n + \bar{v}_n \bar{L}_n$$

where  $(L_n), (\bar{L}_n)$  are two commuting

representor of Virasoro with central charge  $1+6Q^2$

and  $\Psi_{Q+p, \nu, \tilde{\nu}} = L_{-\nu(l_1)} \dots L_{-\nu(l_k)} \bar{L}_{-\tilde{\nu}(l_1)} \dots \bar{L}_{-\tilde{\nu}(l_k)} \Psi_{Q+p}$

if  $\nu = (\nu(l_1), \dots, \nu(l_k))$   $\tilde{\nu} = (\tilde{\nu}(l_1), \dots, \tilde{\nu}(l_k))$

Young diagrams

Ward: we can show using this that for a point  $(P, g_P)$

$\langle \mathcal{A}_{P, g_P}, \Psi_{Q+p_1, \nu_1, \tilde{\nu}_1} \otimes \Psi_{Q+p_2, \nu_2, \tilde{\nu}_2} \otimes \Psi_{Q+p_3, \nu_3, \tilde{\nu}_3} \rangle$

$= W(\nu_1, \nu_2, \nu_3, p_1, p_2, p_3) W(\tilde{\nu}_1, \tilde{\nu}_2, \tilde{\nu}_3, l_1, l_2, l_3)$

$\times \langle \mathcal{A}_{P, g_P}, \bigotimes_{j=1}^3 \Psi_{Q+p_j} \rangle$

algebraic coefficients

Structure constant  $\times$  conformal anomaly ( $g_P$  depending)

Th. Bootstrap: [GRV 21+22]

$\alpha = (\alpha_1, \dots, \alpha_m)$

$\langle V_{\alpha_1}(z_1) \dots V_{\alpha_m}(z_m) \rangle_{\Sigma, g}$

product of 3-pt functions on  $S^2$

$= \int_{(\mathbb{R}^+)^{3g-3+m}} \rho(P, \alpha) \left| \mathcal{F}_{\alpha, P}(\Sigma, z) \right|^2 dP$



⇒ Shows convergence of conformal blocks

$$\forall \text{ a.e. } p \in (\mathbb{R}^+)^3 g_{\epsilon^{-3+an}}$$

⇒ Shows crossing symmetry.

Proof relies on Segal cutting/gluing + Spectral resolut<sup>o</sup>  
+ Ward identities.

Remark: ① blocks do depend on cutting curves.

② Changing cutting curve in a fixed homotopy class changes blocks by a cocycle ("Schwarzian cocycle")

③ blocks are global analytic sections on Teichmüller space

This will lead to unitary representation of

Mapping class group in space of conformal blocks

(with Bawez, Kupinen & Rhodes).

also called Modular Functor.