

nEDM limits on ALP couplings to fermions



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arXiv:2403.12133



Axion-like particles (ALPs)

Axions and ALPs are:

- pseudo-Goldstone bosons of some new $U(1)$
- well motivated NP candidates
- targeted by an extensive experimental program

| | Axion | ALPs |
|-------------------|--------------|-------------|
| Strong CP problem | ✓ | ✗ |
| Dark Matter | ✓ | ✓ |
| Cosmic Inflation | ✓ | ✓ |
| Baryogenesis | ✓ | ✓ |

The ALP Effective Field Theory

ALP couplings to up- and down-type quarks:

$$\begin{aligned}\mathcal{L}_a \supset & \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{1}{2} m_a^2 a^2 \\ & + (\bar{u}_L \mathbf{M}_u u_R + \bar{d}_L \mathbf{M}_d d_R + \text{h.c.}) + \theta \frac{\alpha_s}{8\pi} G_{\mu\nu} \tilde{G}^{\mu\nu} \\ & + \frac{\partial_\mu a}{f_a} (\bar{Q}_L \gamma^\mu \mathbf{C}_Q Q_L + \bar{u}_R \gamma^\mu \mathbf{C}_{u_R} u_R + \bar{d}_R \gamma^\mu \mathbf{C}_{d_R} d_R)\end{aligned}$$

[Georgi, Kaplan, Randall, *Phys. Lett. B* 169 (1986) 73-78]

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The ALP Effective Field Theory

ALP couplings to up- and down-type quarks:

$$\mathcal{L}_a \supset (\bar{u}_L \mathbf{M}_u u_R + \bar{d}_L \mathbf{M}_d d_R + \text{h.c.}) + \theta \frac{\alpha_s}{8\pi} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

related by anomalous $U(1)_{\text{Axial}}$

\implies Physical combination is

$$\bar{\theta} = \theta + \text{Arg det}(\mathbf{M}_u \mathbf{M}_d)$$

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ALP couplings to up- and down-type quarks:

$$\mathcal{L}_a \supset \frac{\partial_\mu a}{f_a} \left(\bar{Q}_L \gamma^\mu \underbrace{C_Q Q_L + \bar{u}_R \gamma^\mu C_{u_R} u_R + \bar{d}_R \gamma^\mu C_{d_R} d_R}_{\text{CP-violation in flavor-nondiagonal entries}} \right)$$

CP-violation in flavor-nondiagonal entries

[Georgi, Kaplan, Randall, *Phys. Lett. B* 169 (1986) 73-78]

arXiv:2403.12133

The ALP Effective Field Theory

ALP couplings to up- and down-type quarks:

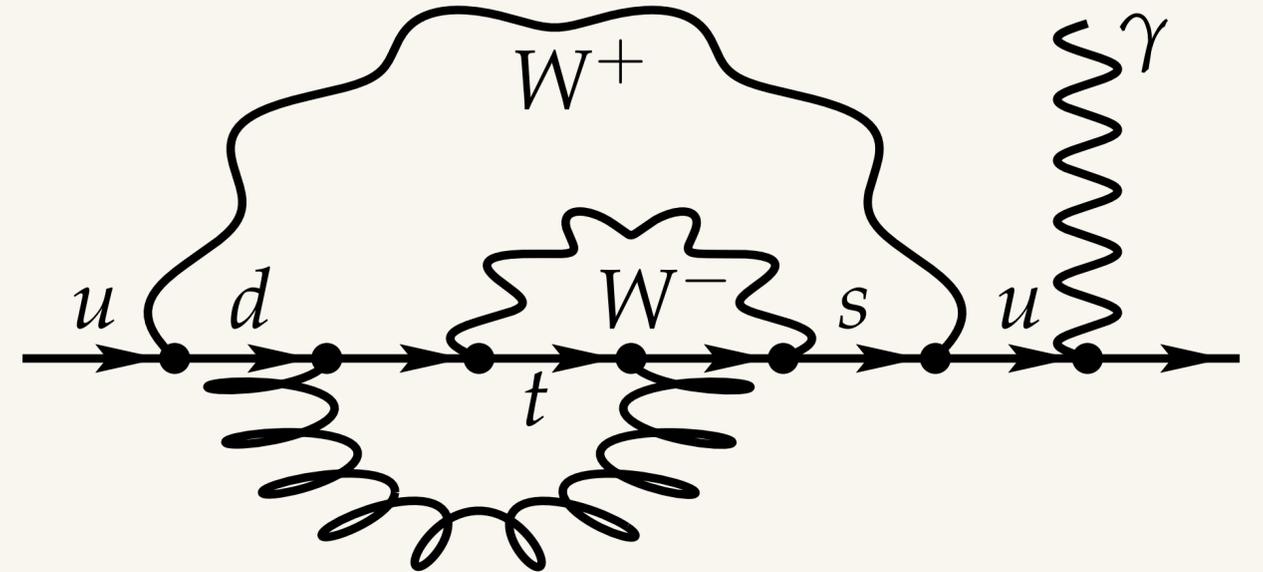
$$\mathcal{L}_a \supset \frac{\partial_\mu a}{f_a} \left(\bar{Q}_L \gamma^\mu \mathbf{C}_Q Q_L + \bar{u}_R \gamma^\mu \mathbf{C}_{u_R} u_R + \bar{d}_R \gamma^\mu \mathbf{C}_{d_R} d_R \right)$$

CP-violation in flavor-nondiagonal entries

Will source CP-violating observables e.g. EDMs

The neutron EDM

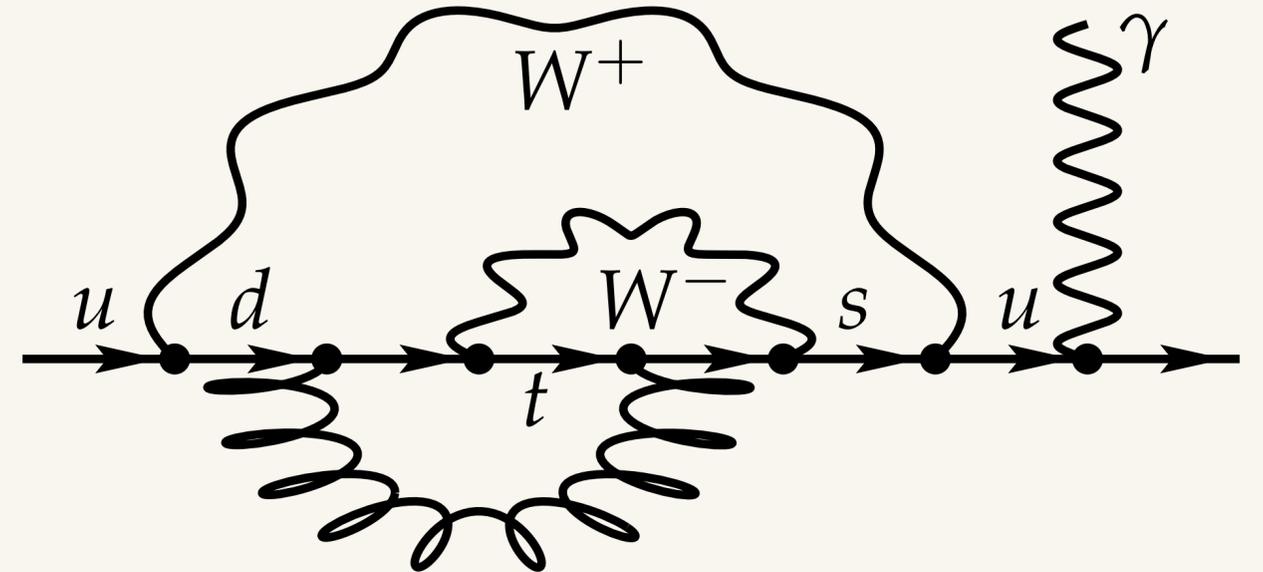
SM prediction 3-loop suppressed!!



nEDM extremely well constrained: $d_n^{\text{exp}} \lesssim 2.6 \times 10^{-26} \text{ [e}\cdot\text{cm]}$
[Abel et al., 2001.11966] [Pendlebury *et al.*, 2001.11966]

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How can C_Q , C_{u_R} , C_{d_R} contribute?

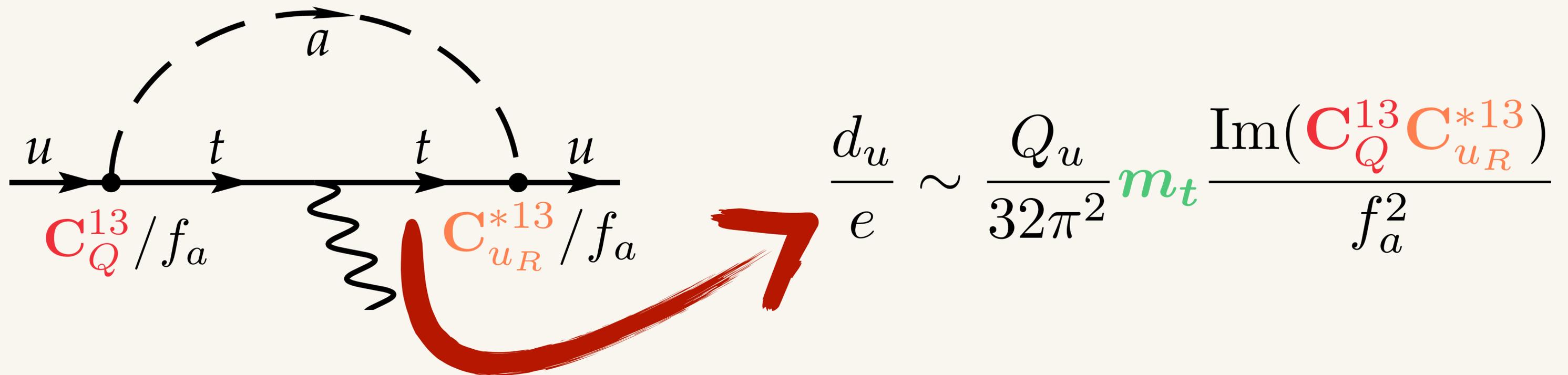
nEDM sourced by $\left\{ \begin{array}{l} \bullet \text{ Quark EDMs and CEDMs} \\ \bullet \text{ The } \bar{\theta} \text{ parameter} \end{array} \right.$

[Baluni, *Phys. Rev. D* 19, 2227]

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ALP contributions appear at 1-loop

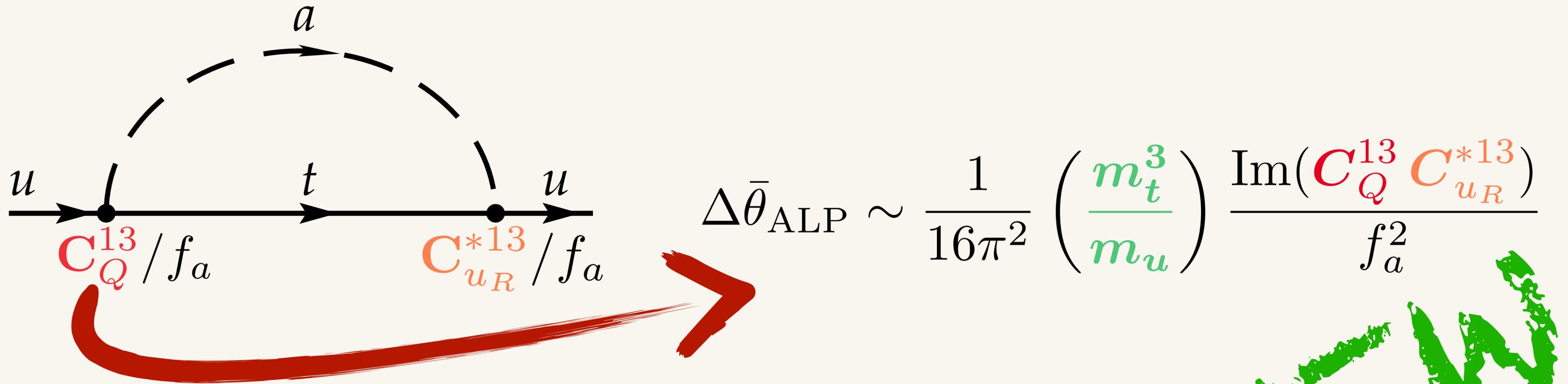
ALP contributions to the nEDM



Corrections to the quark EDMs and CEDMs

[Di Luzio et al., 2010.13760]

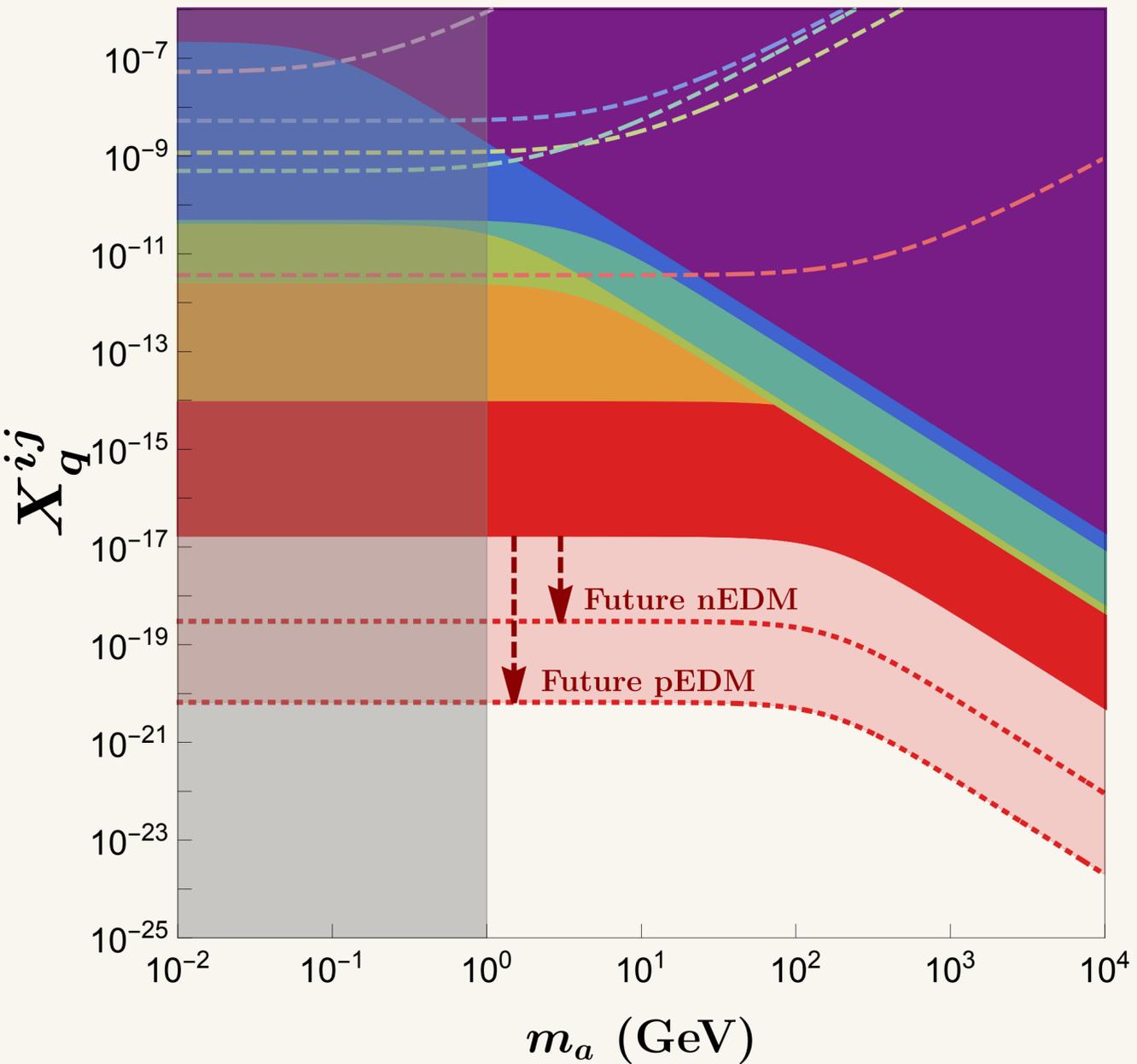
ALP contributions to the nEDM



Corrections to $\bar{\theta} = \theta + \text{Arg det}(M_u M_d)$

NEW

nEDM limits on ALP-fermion couplings



$$X_q^{ij} = \text{Im}(C_L^{ij} C_{qR}^{*ij}) / f_a^2 \text{ (GeV}^{-2}\text{)}$$

- X_u^{13}
- X_u^{23}
- X_d^{13}
- X_u^{12}
- X_d^{23}
- X_d^{12}

Dotted lines:

$$\left. \frac{d_n}{e} \right|_{d_q, \tilde{d}_q} \sim \mathcal{O}(1) \times \frac{Q_u}{32\pi^2} m_t \frac{\text{Im}(C_Q^{13} C_{uR}^{*13})}{f_a^2}$$

Solid regions:

$$\left. \frac{d_n}{e} \right|_{\bar{\theta}} \sim \frac{\mathcal{O}(10^{-3})}{16\pi^2} \times \left(\frac{m_t^3}{m_u} \right) \frac{\text{Im}(C_Q^{13} C_{uR}^{*13})}{f_a^2}$$

OLD

NEW

General scalar theory

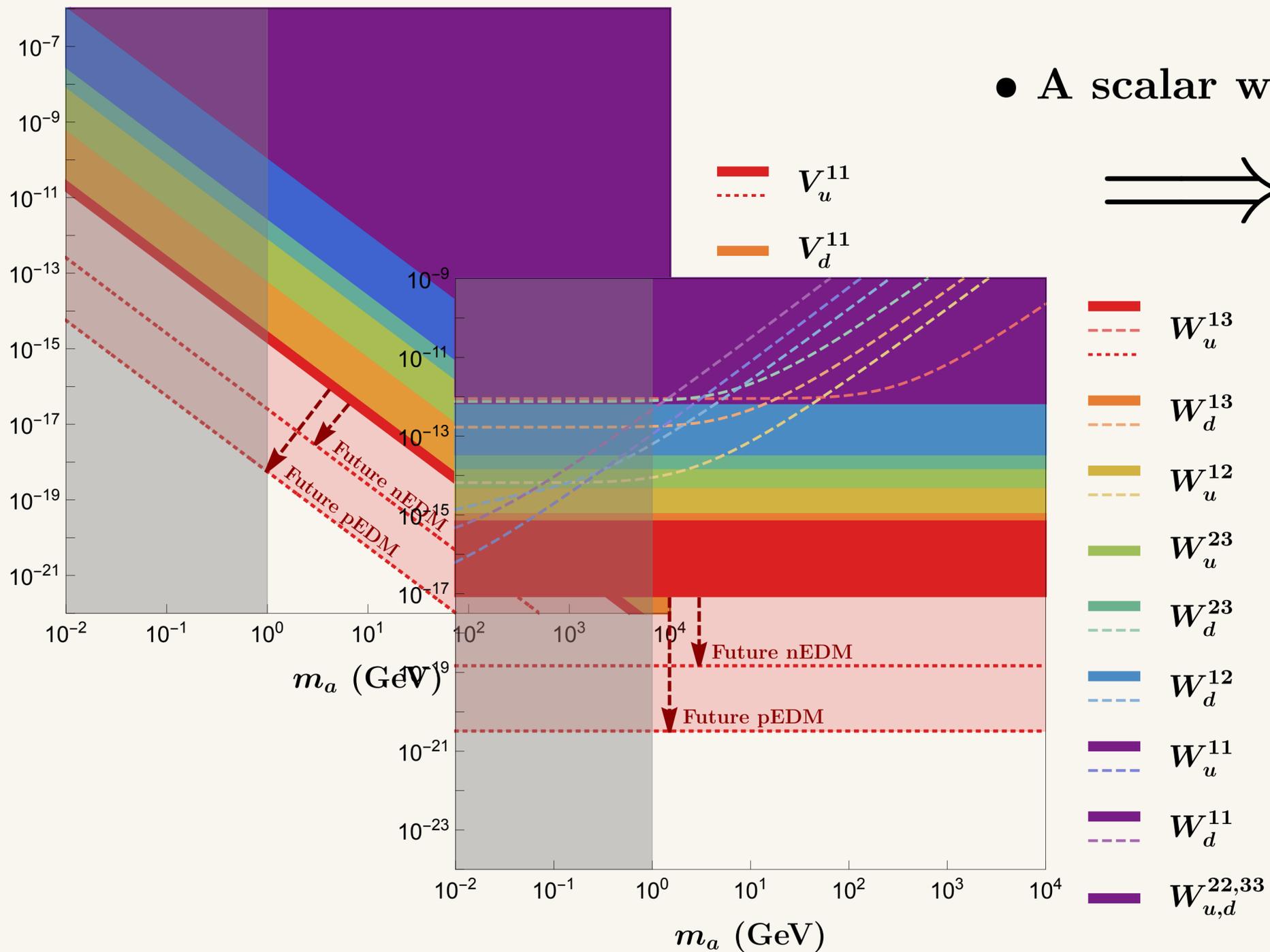
- A scalar which may not be a pseudo-Goldstone



More parametric freedom

We also

- Improved existing bounds
- Established new bounds



MORE IN OUR PAPER

Conclusions

- ALP couplings to fermion induce parametrically enhanced corrections to the nEDM at one loop
- We have improved the bounds on CP-odd ALP-fermion couplings by ~ 4 orders of magnitude
- The same kind of improvement applies for a general scalar

Backup

Interactions of the quarks with an ALP:

$$\mathcal{L}_a \supset \frac{\partial_\mu a}{f_a} (\bar{Q}_L \gamma^\mu \mathbf{C}_Q Q_L + \bar{u}_R \gamma^\mu \mathbf{C}_{u_R} u_R + \bar{d}_R \gamma^\mu \mathbf{C}_{d_R} d_R)$$

$$v\mathbf{K}_q \equiv \mathbf{C}_Q \mathbf{M}_q - \mathbf{M}_q \mathbf{C}_{q_R}$$
$$2v\mathbf{F}_q \equiv 2\mathbf{C}_Q \mathbf{M}_q \mathbf{C}_{q_R} - \mathbf{C}_Q^2 \mathbf{M}_q - \mathbf{M}_q \mathbf{C}_{q_R}^2$$

Interactions of the quarks with a general scalar:

$$\mathcal{L} \supset \bar{u}_L v \left[i\mathbf{K}_u \frac{\phi}{\Lambda} + \mathbf{F}_u \frac{\phi^2}{\Lambda^2} \right] u_R + \bar{d}_L v \left[i\frac{\phi}{\Lambda} \mathbf{K}_d + \frac{\phi^2}{\Lambda^2} \mathbf{F}_d \right] d_R + \text{h.c.}$$

Backup

Without a PQ mechanism:

$$\begin{aligned}d_n &= 0.6(3) \times 10^{-16} \bar{\theta} [e \cdot \text{cm}] \\ &- 0.204(11) d_u + 0.784(28) d_d - 0.0028(17) d_s \\ &- 0.32(15) e \tilde{d}_u + 0.32(15) e \tilde{d}_d - 0.014(7) e \tilde{d}_s.\end{aligned}$$

In the presence of a PQ mechanism:

$$\begin{aligned}d_n^{\text{PQ}} &= -0.204(11) d_u + 0.784(28) d_d - 0.0028(17) d_s \\ &- 0.31(15) e \tilde{d}_u + 0.62(31) e \tilde{d}_d\end{aligned}$$

Backup

$$\begin{aligned}\bar{\theta}(\mu_{\text{IR}}) &\simeq \bar{\theta}_0 + \\ &\sum_{u_i=\{u,c,t\}} \frac{m_{u_k} (m_a^2 + \hat{m}_{u_k}^2)}{16\pi^2 f_a^2 m_{u_i}} \text{Im} \left(\mathbf{C}_Q^{ik} \mathbf{C}_{u_R}^{*ik} \right) \log \frac{f_a^2}{\max(m_a^2, m_{u_k}^2)} \\ &+ \sum_{d_i=\{d,s,b\}} \frac{m_{d_k} (m_a^2 + \hat{m}_{d_k}^2)}{16\pi^2 f_a^2 m_{d_i}} \text{Im} \left(\mathbf{C}_Q^{ik} \mathbf{C}_{d_R}^{*ik} \right) \log \frac{f_a^2}{\max(m_a^2, m_{d_k}^2)}\end{aligned}$$