



Exploring t-channel models for Dark Matter

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T-channel Portal

We have provided a complete matching of the t-channel portal, with a complete analysis focused on Direct Detection and Relic Density

From High energy

$$\mathcal{L}_{scalar} = \Gamma_L^{f_i} \bar{f}_i P_R \Psi_{f_i} \Phi_{DM} + \Gamma_R^{f_i} \bar{f}_i P_L \Psi_{f_i} \Phi_{DM} + \text{h.c.} \\ + \lambda_{1H\Phi} (\Phi_{DM}^\dagger \Phi_{DM}) (H^\dagger H) + \lambda_{2H\Phi} (\Phi_{DM}^\dagger T_\Phi^a \Phi_{DM}) (H^\dagger \frac{\sigma^a}{2} H)$$

$$\mathcal{L}_{fermion} = \Gamma_L^{f_i} \bar{f}_i P_R \Phi_{f_i} \Psi_{DM} + \Gamma_R^{f_i} \bar{f}_i P_L \Phi_{f_i} \Psi_{DM} + \text{h.c.} \\ + \lambda_{1H\Phi} (\Phi_{f_i}^\dagger \Phi_{f_i}) (H^\dagger H) + \lambda_{2H\Phi} (\Phi_{f_i}^\dagger T_\Phi^a \Phi_{f_i}) (H^\dagger \frac{\sigma^a}{2} H)$$

where:

$$\mathcal{O}_{\mu\nu}^q \equiv \bar{q} i \left(\frac{D_\mu \gamma_\nu + D_\nu \gamma_\mu}{2} - \frac{1}{4} g_{\mu\nu} \not{D} \right) q,$$

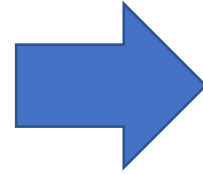
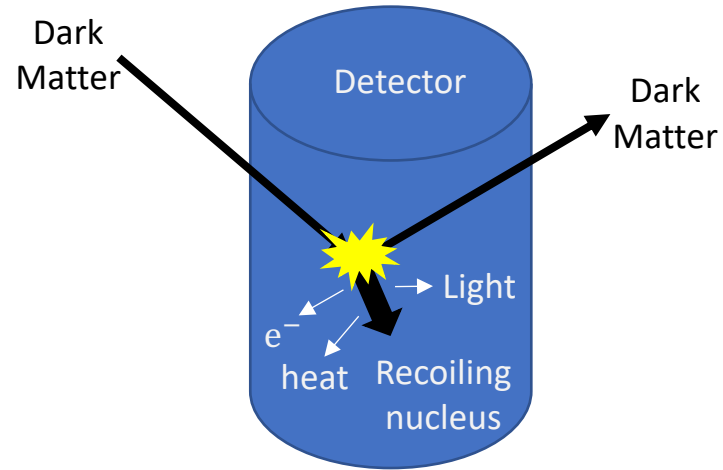
$$\mathcal{O}_{\mu\nu}^g \equiv G_\mu^{a\rho} G_{\nu\rho}^a - \frac{1}{4} g_{\mu\nu} G_{\rho\sigma}^a G^{a\rho\sigma},$$

To Low energy

$$\mathcal{L}_{eff}^{Scalar,q} = \sum_{q=u,d} c^q \left(\Phi_{DM}^\dagger i \overleftrightarrow{\partial}_\mu \Phi_{DM} \right) \bar{q} \gamma^\mu q + \sum_{q=u,d,s} d^q M_q \Phi_{DM}^\dagger \Phi_{DM} \bar{q} q + d^g \frac{\alpha_s}{\pi} \Phi_{DM}^\dagger \Phi_{DM} G^{a\mu\nu} G_{\mu\nu}^a \\ + \sum_{q=u,d,s} g_1^q \frac{\Phi_{DM} (i\partial^\mu) (i\partial^\nu) \Phi_{DM} \mathcal{O}_{\mu\nu}^q}{M_{\Phi_{DM}}^2} + g_1^g \frac{\Phi_{DM} (i\partial^\mu) (i\partial^\nu) \Phi_{DM} \mathcal{O}_{\mu\nu}^g}{M_{\Phi_{DM}}^2},$$

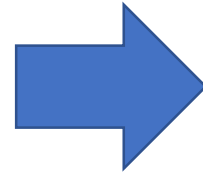
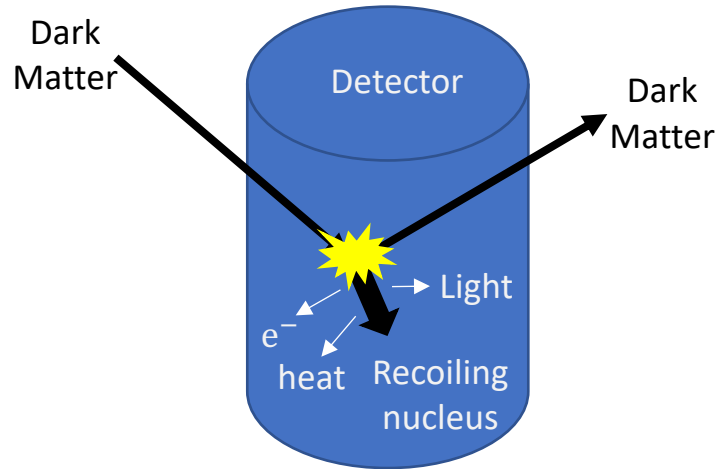
$$\mathcal{L}_{eff}^{Dirac,q} = \sum_{q=u,d} c^q \bar{\Psi}_{DM} \gamma_\mu \Psi_{DM} \bar{q} \gamma^\mu q + \sum_{q=u,d,s} \tilde{c}^q \bar{\Psi}_{DM} \gamma_\mu \gamma_5 \Psi_{DM} \bar{q} \gamma^\mu \gamma_5 q \\ + \sum_{q=u,d,s} d^q m_q \bar{\Psi}_{DM} \Psi_{DM} \bar{q} q + \sum_{q=c,b,t} d^g \bar{\Psi}_{DM} \Psi_{DM} G^{a\mu\nu} G_{\mu\nu}^a \\ + \sum_{q=u,d,s} \left(g_1^q \frac{\bar{\Psi}_{DM} i \partial^\mu \gamma^\nu \Psi_{DM} \mathcal{O}_{\mu\nu}^q}{M_{\Psi_{DM}}} + g_2^q \frac{\bar{\Psi}_{DM} (i\partial^\mu) (i\partial^\nu) \Psi_{DM} \mathcal{O}_{\mu\nu}^q}{M_{\Psi_{DM}}^2} \right) \\ + \sum_{q=c,b,t} \left(g_1^{g,q} \frac{\bar{\Psi}_{DM} i \partial^\mu \gamma^\nu \Psi_{DM} \mathcal{O}_{\mu\nu}^g}{M_{\Psi_{DM}}} + g_2^{g,q} \frac{\bar{\Psi}_{DM} (i\partial^\mu) (i\partial^\nu) \Psi_{DM} \mathcal{O}_{\mu\nu}^g}{M_{\Psi_{DM}}^2} \right).$$

Experimental bounds

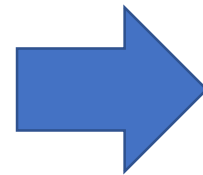
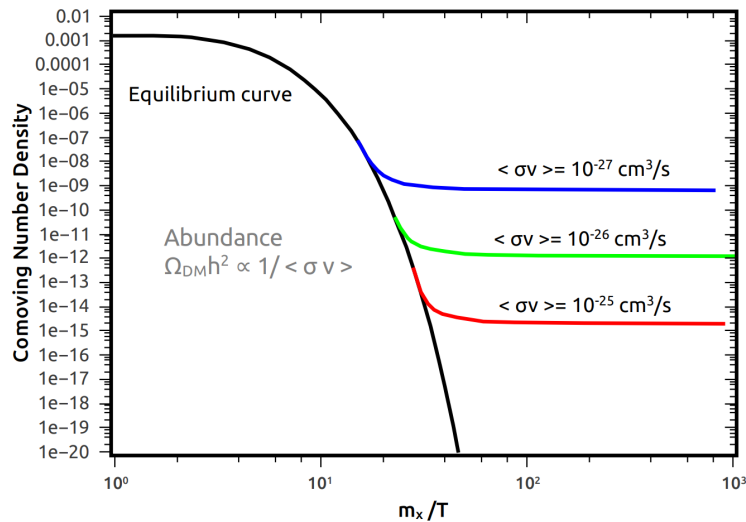


$$\sigma_{\Phi_{\text{DM}}}^{\text{SI},p} = \frac{\mu_{\Phi_{\text{DM}}p}^2}{\pi} \frac{[Zf_p + (A - Z)f_n]^2}{A^2}$$

Experimental bounds



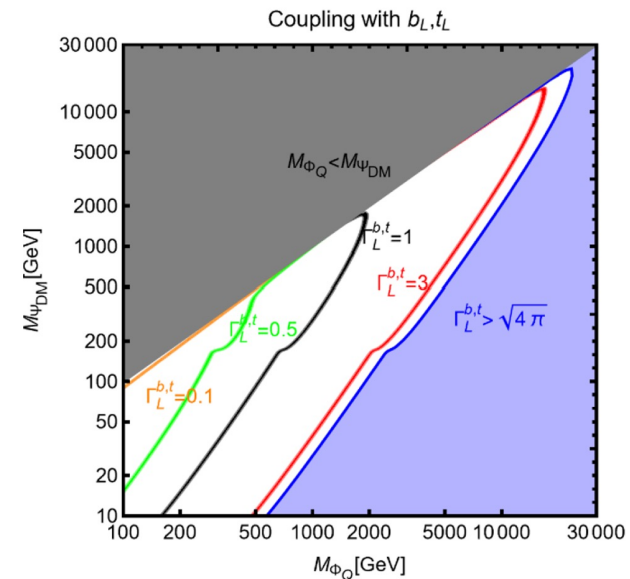
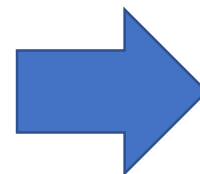
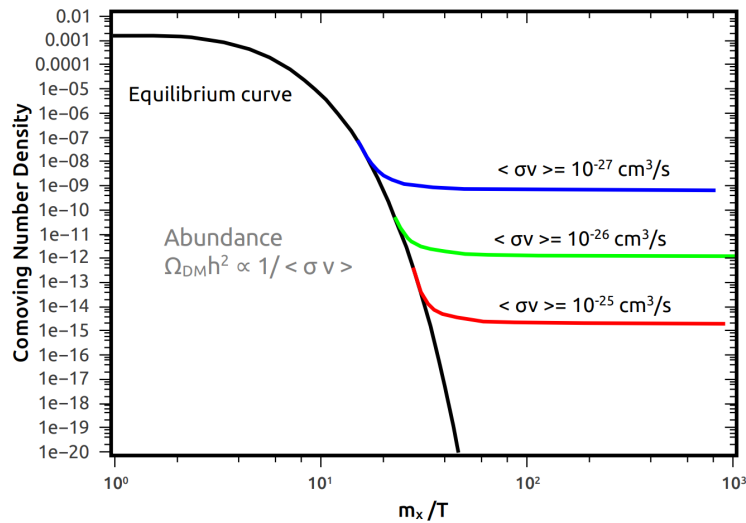
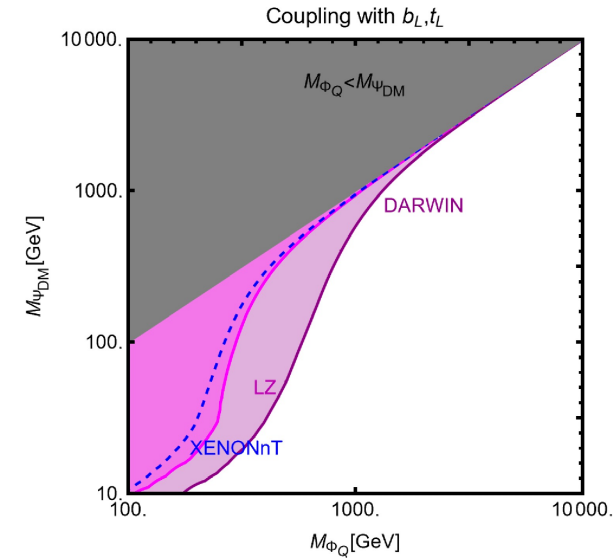
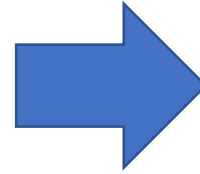
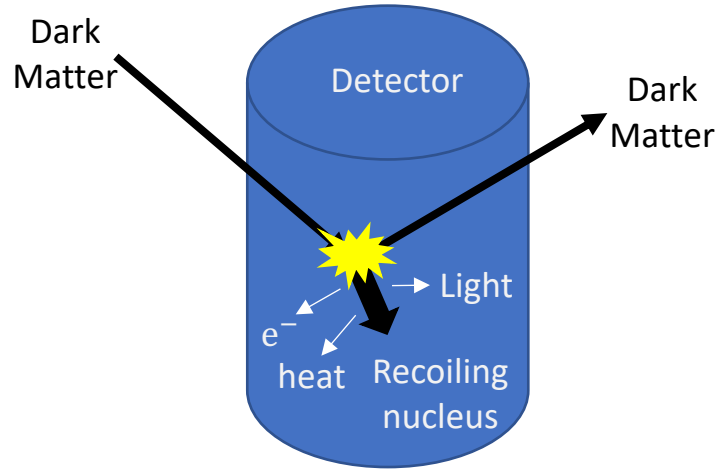
$$\sigma_{\Phi_{\text{DM}}}^{\text{SI},p} = \frac{\mu_{\Phi_{\text{DM}}p}^2}{\pi} \frac{[Zf_p + (A-Z)f_n]^2}{A^2}$$



$$\Omega_{\text{DM}} h^2 \approx 8.76 \times 10^{-11} \text{ GeV}^{-2} \left[\int_{T_{\text{f.o.}}}^{T_0} g_*^{1/2} \langle \sigma v \rangle_{\text{eff}} \frac{dT}{M_{\text{DM}}} \right]^{-1}$$

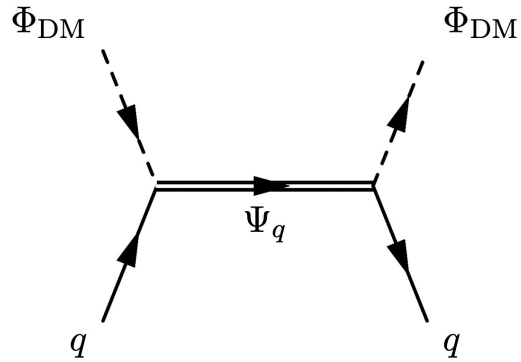
arXiv:hep-ph/9704361

Experimental bounds

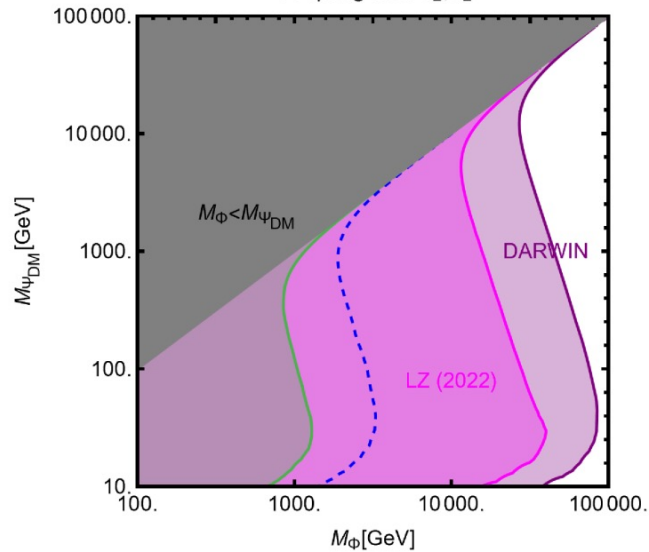


Why do we need radiative corrections?

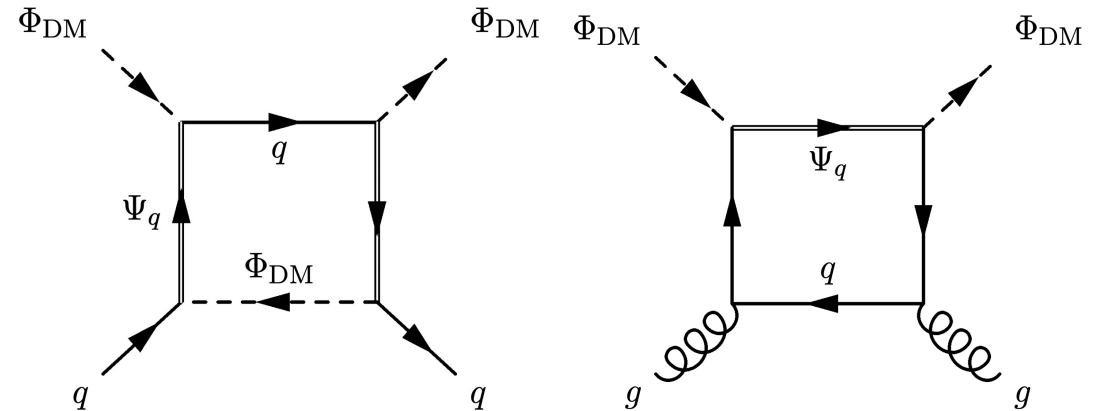
Coupling 1^o generation



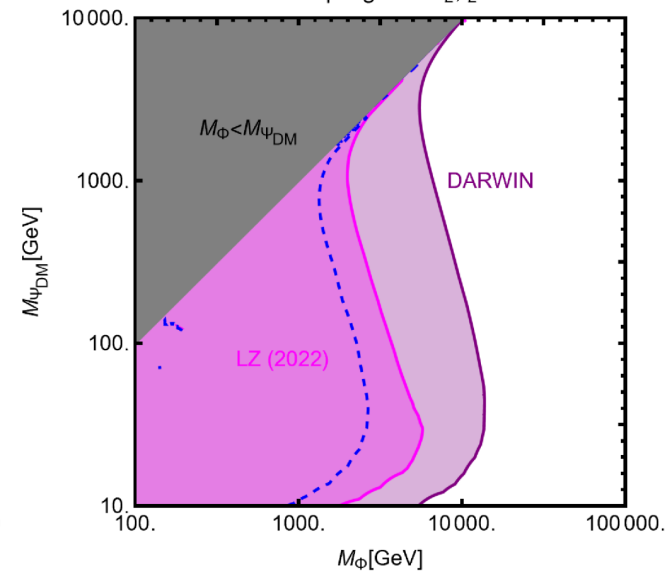
Coupling with u_L, d_L



Coupling 2^o & 3^o generation

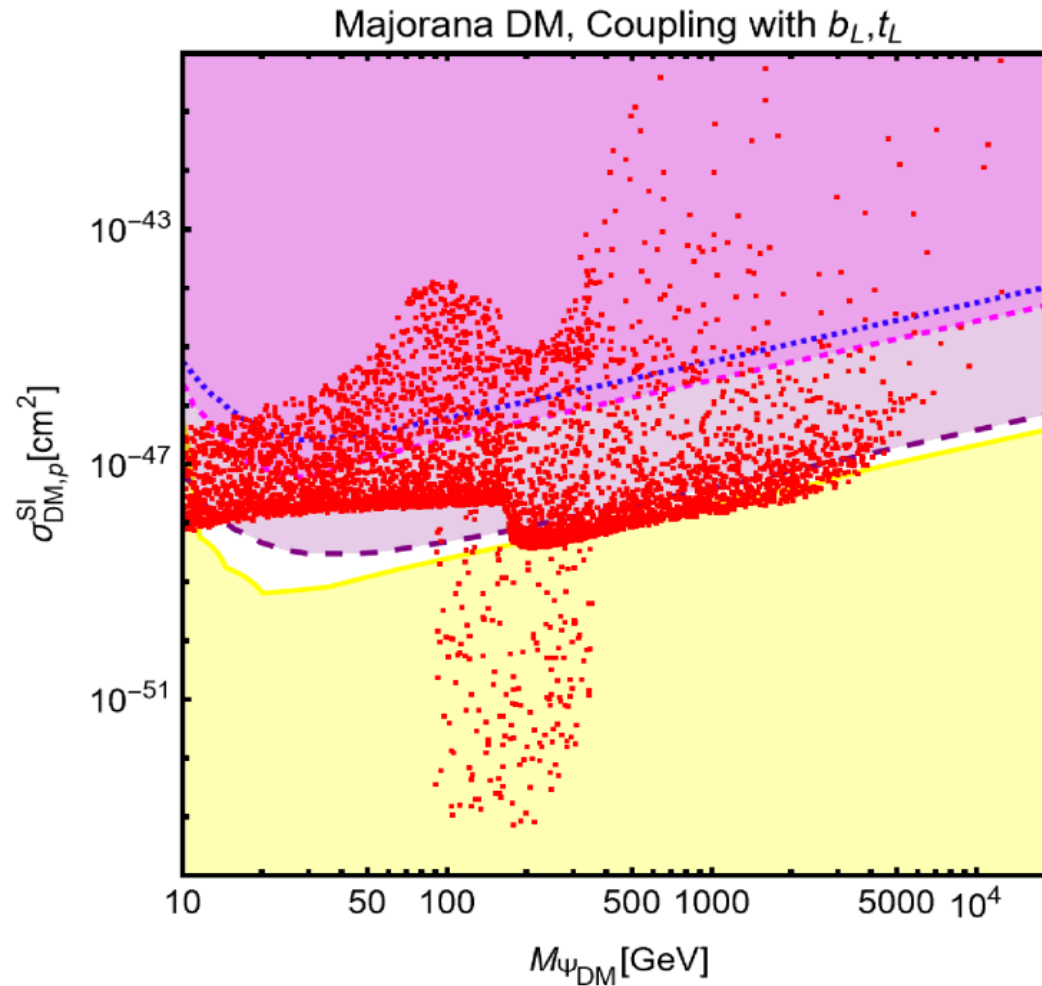


Coupling with b_L, t_L



Scan with RD & DD bounds

$$M_{\Phi_{\text{DM}}, \Psi_{\text{DM}}} \in [10, 10^5] \text{ GeV} \quad M_{\Phi_f, \Psi_f} \in [100, 10^5] \text{ GeV} \quad \Gamma_{L,R}^f \in [10^{-3}, \sqrt{4\pi}]$$



Conclusions

- Complete matching for both scalars and fermionic DM candidates to DD EFT Lagrangian.
- Strong bounds for the Complex case and Dirac case, with the exception of the very fine-tuned coannihilation region.
- Real DM weaker DD constrains but very suppressed annihilation cross-section which also strongly constrain the candidate.
- Majorana DM results the most favored among the ones considered in this work and the only allowing for viable masses of order or below 100 GeV.

Thank You!



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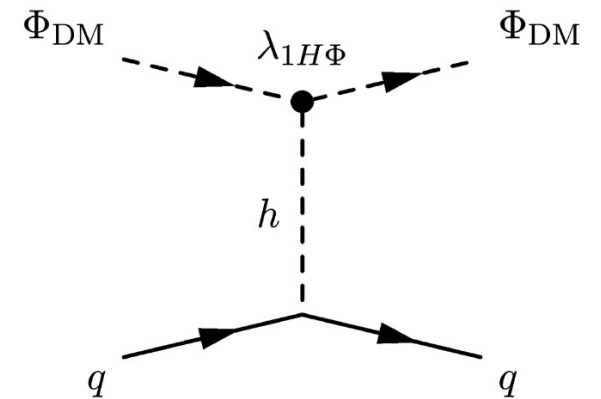
Back up

Hidden Higgs portal

$$\lambda_{1H\Phi}^{(0)} = \lambda_{1H\Phi}(\mu) + \sum_f \frac{g^2 m_f^2 |\Gamma_{L,R}^f|^2}{16\pi^2 m_W^2} \left(\frac{1}{\epsilon} + \log \frac{\mu^2}{m_f^2} \right)$$

Without resummation:

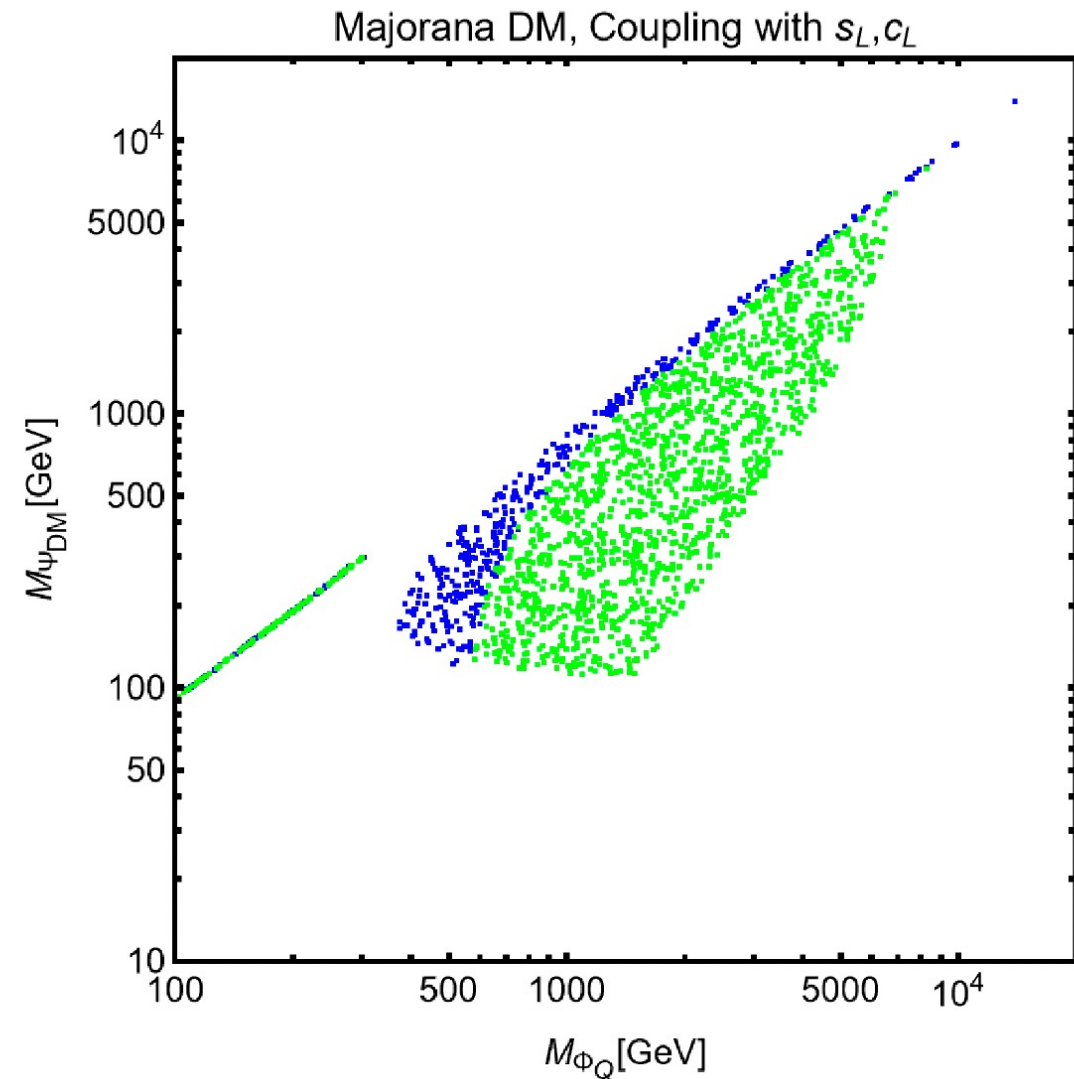
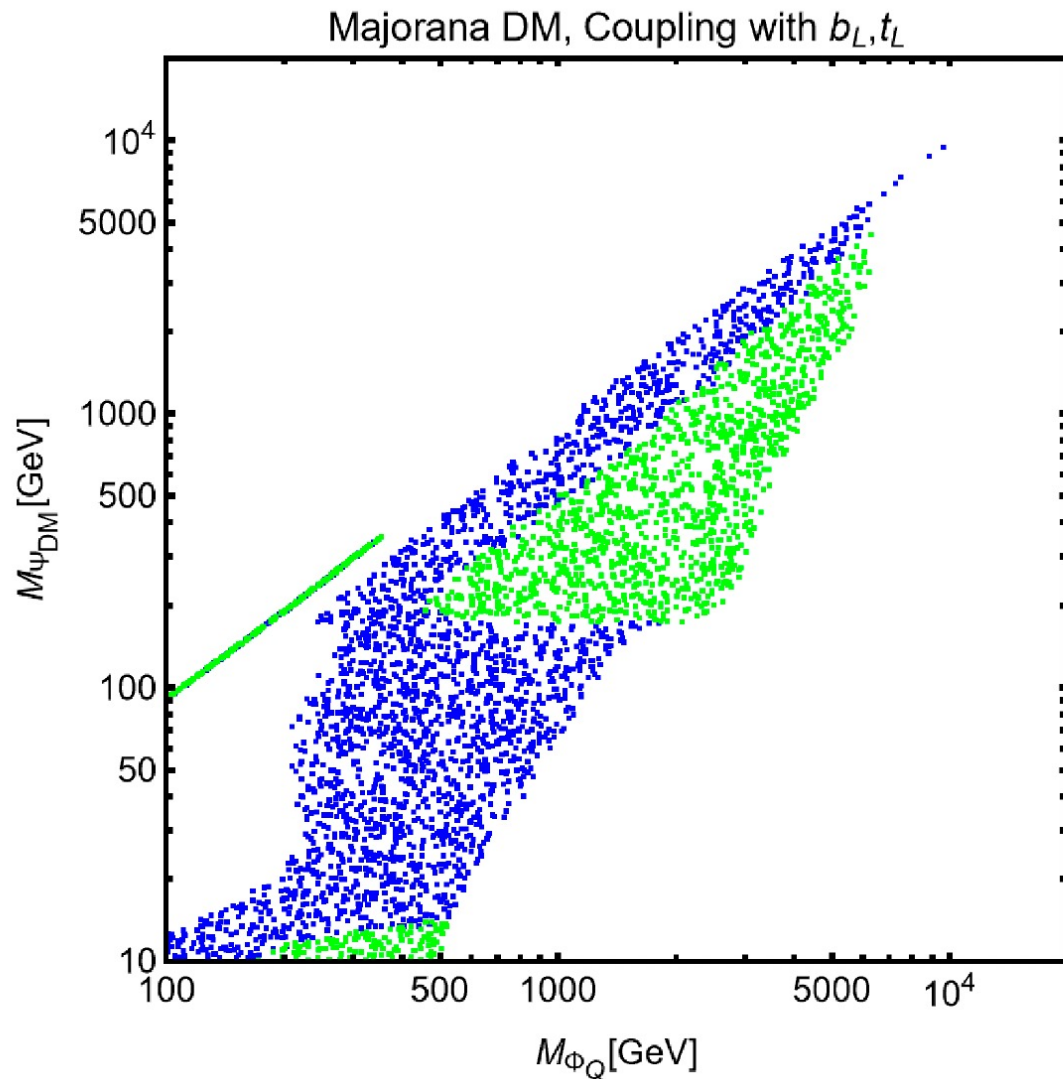
$$\lambda_{1H\Phi}(\mu) = \lambda_{1H\Phi}(M) - \log \frac{\mu^2}{M^2} \sum_f \frac{g^2 m_f^2 |\Gamma_{L,R}^f|^2}{16m_W^2 \pi^2}$$



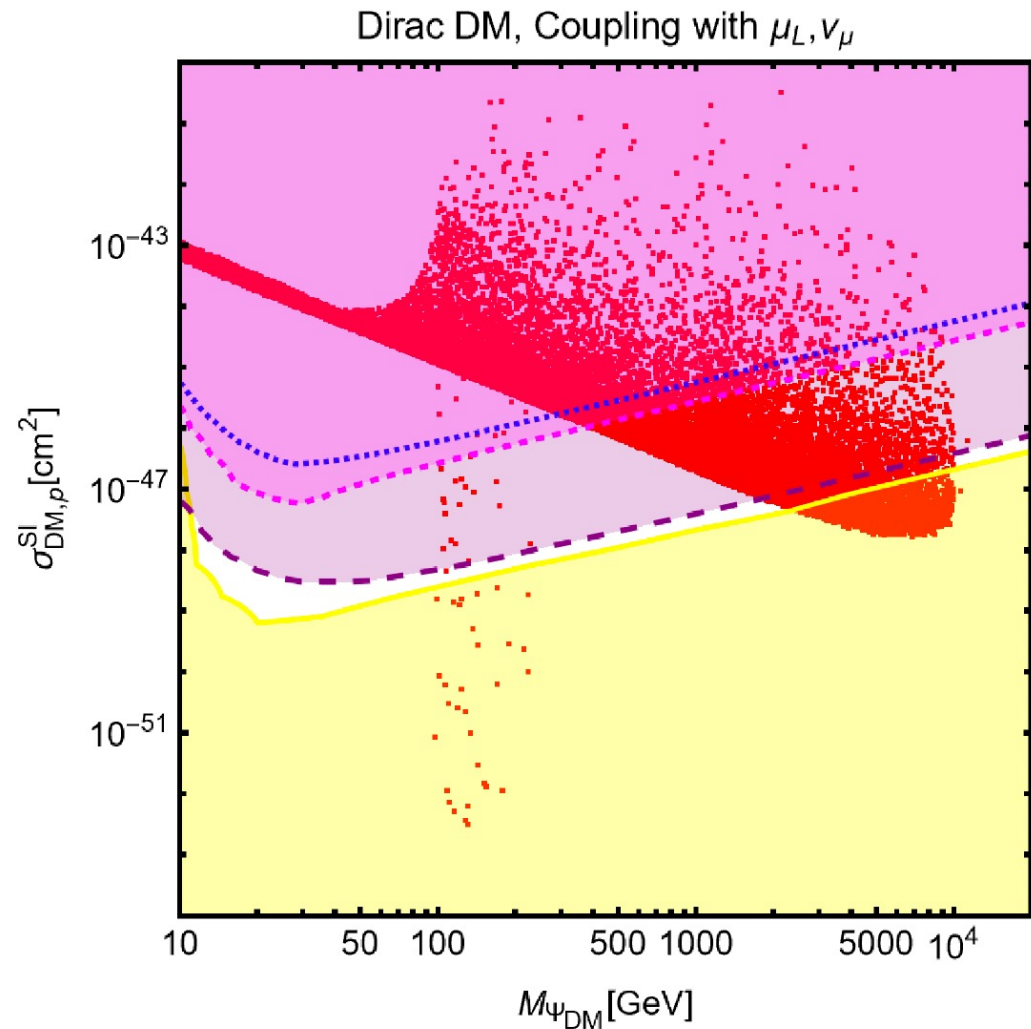
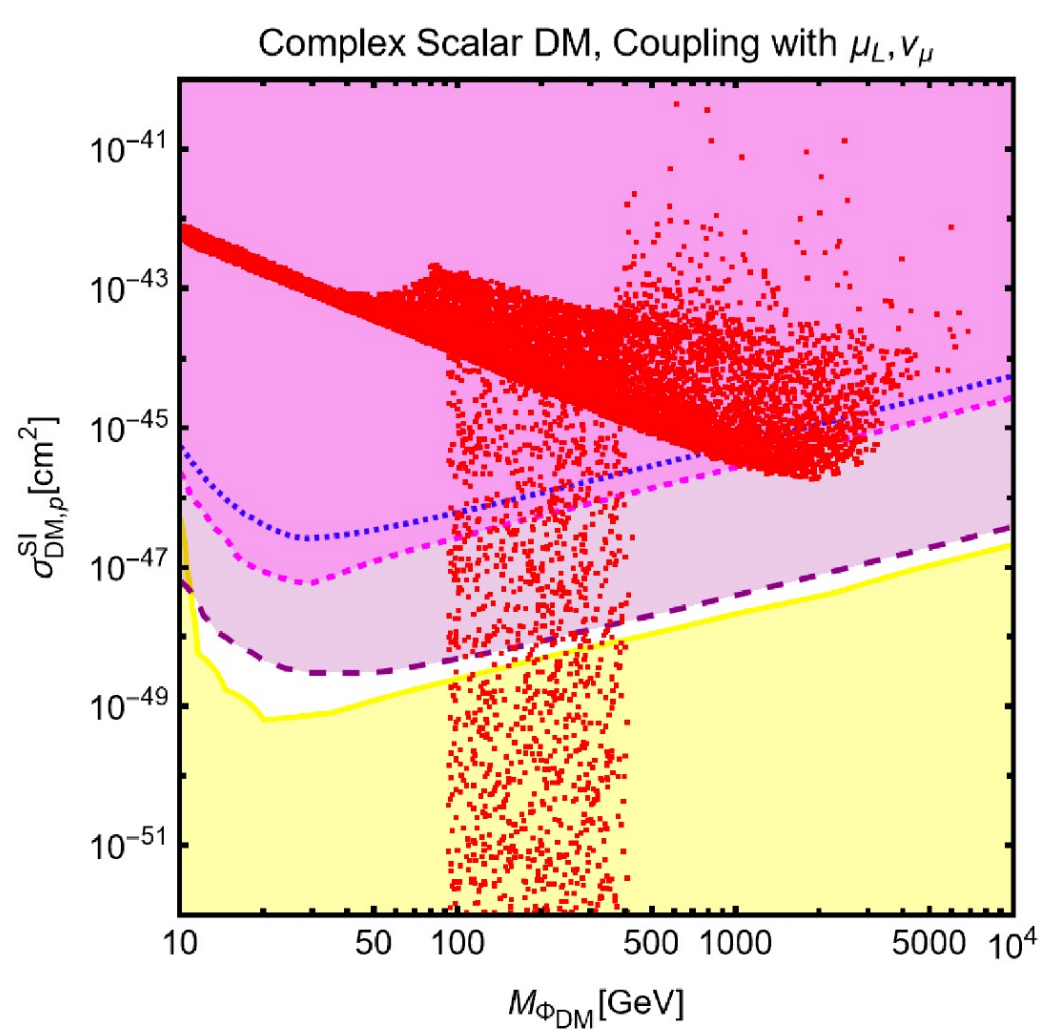
Velocity expansion

$$\begin{aligned}
 \langle \sigma v \rangle_{\text{DMDM}}^{\text{Majorana}} &= \sum_f N_c^f \frac{|\Gamma_{L,R}^f|^4 M_{\Psi_{\text{DM}}}^2 (M_{\Psi_{\text{DM}}}^4 + M_{\Phi_f}^4) v^2}{48\pi (M_{\Psi_{\text{DM}}}^2 + M_{\Phi_f}^2)^4} & \langle \sigma v \rangle_{\text{DMDM}}^{\text{Dirac}} &= \sum_f N_c^f \frac{|\Gamma_{L,R}^f|^4 M_{\Psi_{\text{DM}}}^2}{32\pi (M_{\Psi_{\text{DM}}}^2 + M_{\Phi_f}^2)^2} \\
 \langle \sigma v \rangle_{\text{DMDM}}^{\text{Real}} &= \sum_f N_c^f \frac{|\Gamma_{L,R}^f|^4 M_{\Phi_{\text{DM}}}^6 v^4}{60\pi (M_{\Phi_{\text{DM}}}^2 + M_{\Psi_f}^2)^4} & \langle \sigma v \rangle_{\text{DMDM}}^{\text{Complex}} &= \sum_f N_c^f \frac{|\Gamma_{L,R}^f|^4 M_{\Phi_{\text{DM}}}^2 v^2}{48\pi (M_{\Phi_{\text{DM}}}^2 + M_{\Psi_f}^2)^2},
 \end{aligned}$$

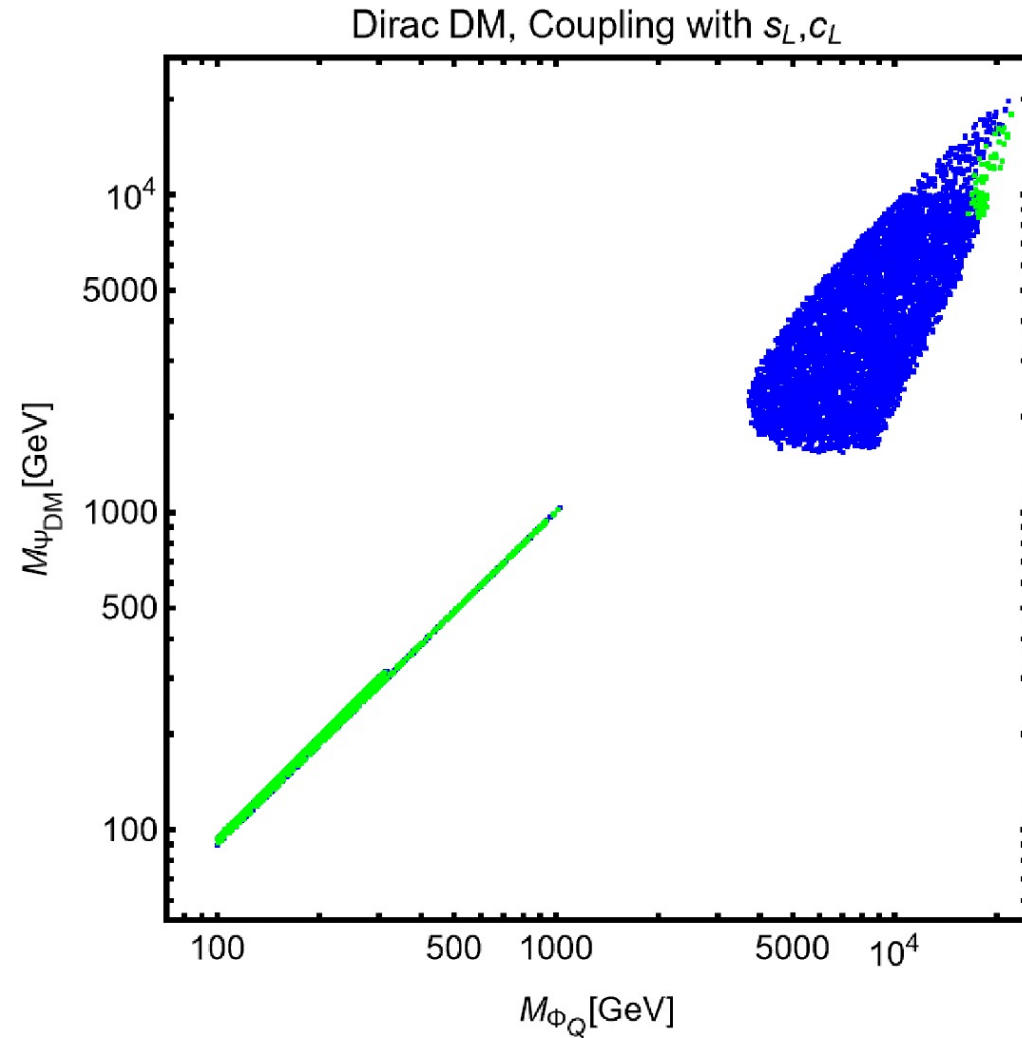
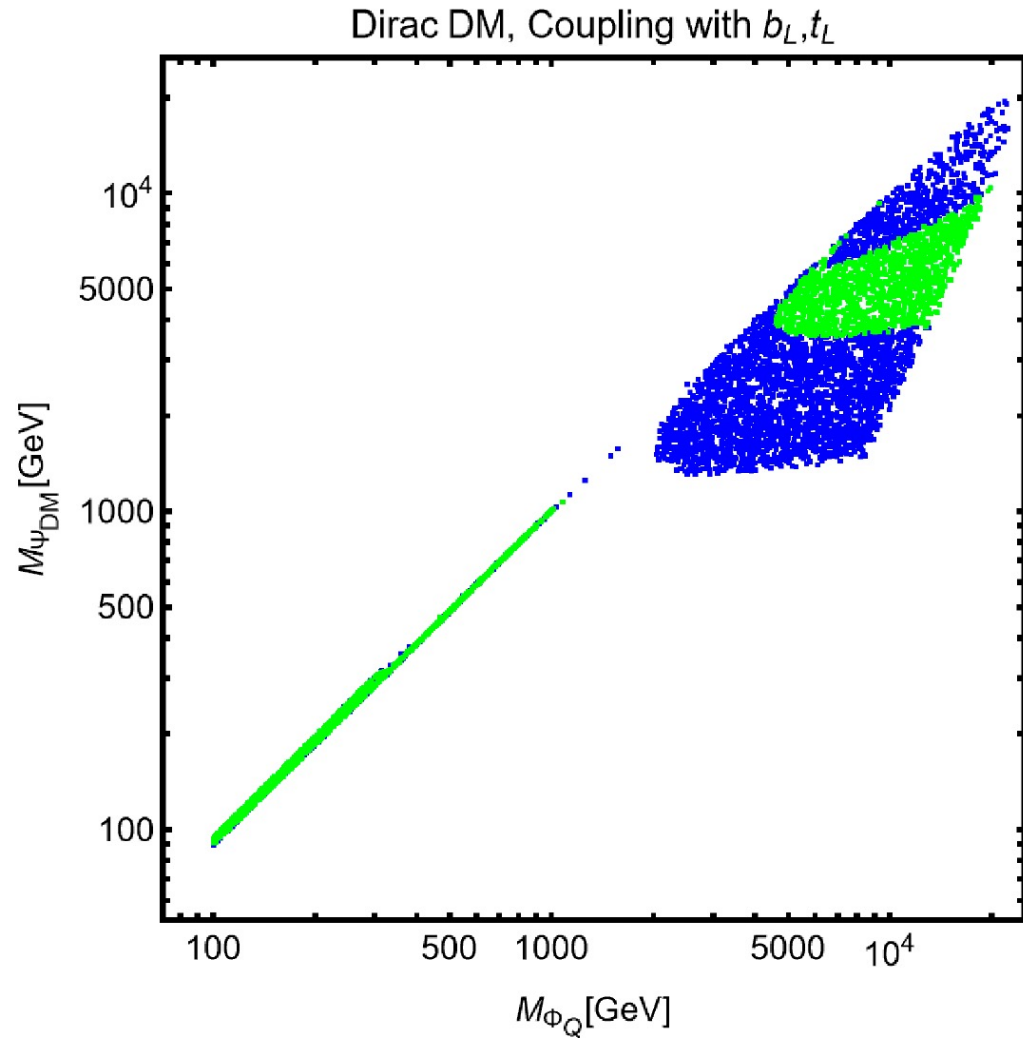
Majorana Bidimensional



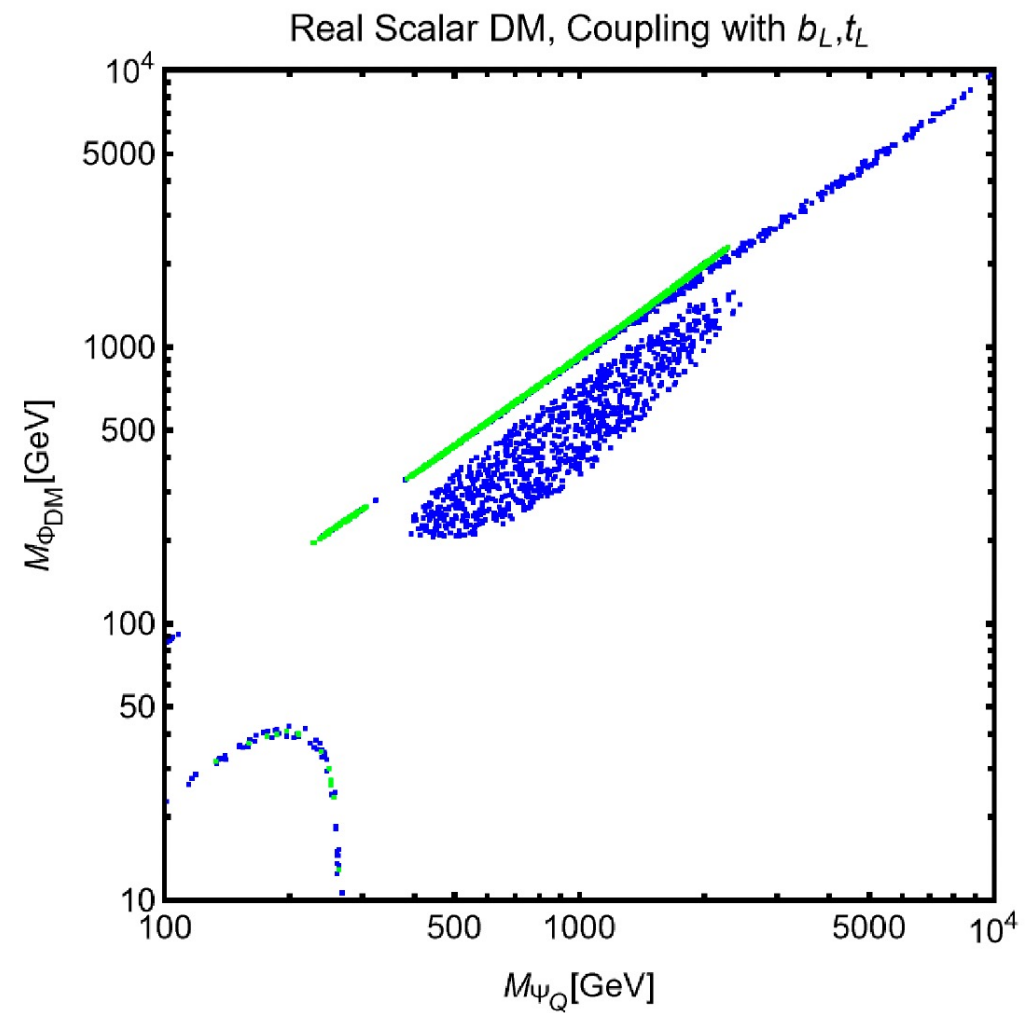
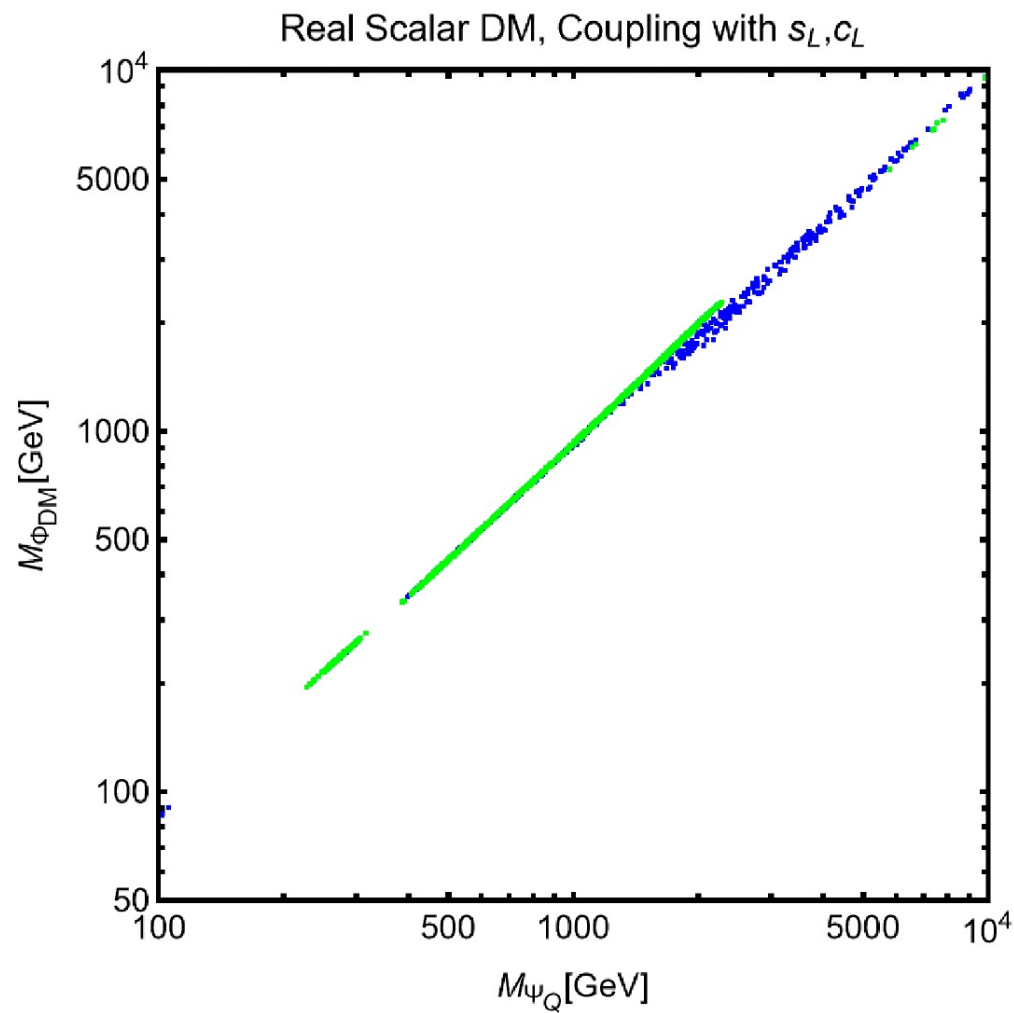
Lepton couplings



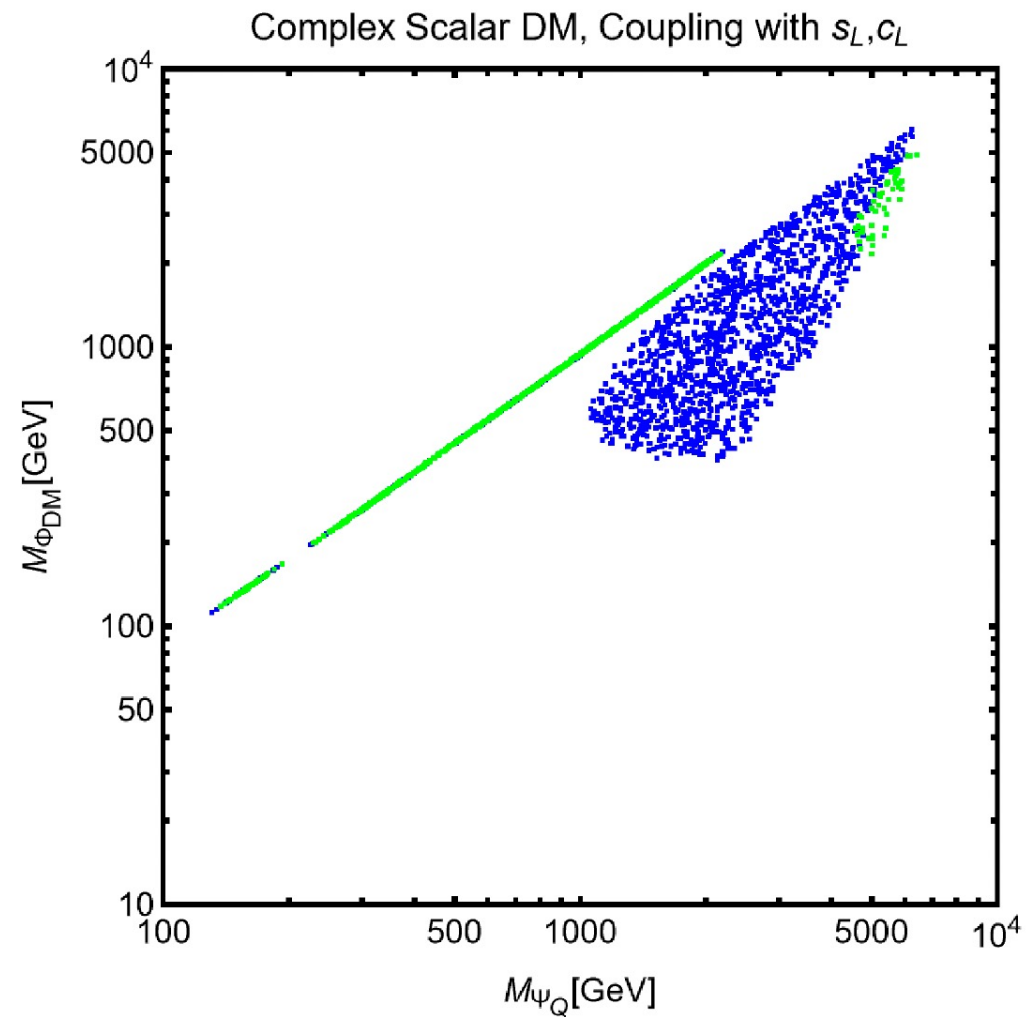
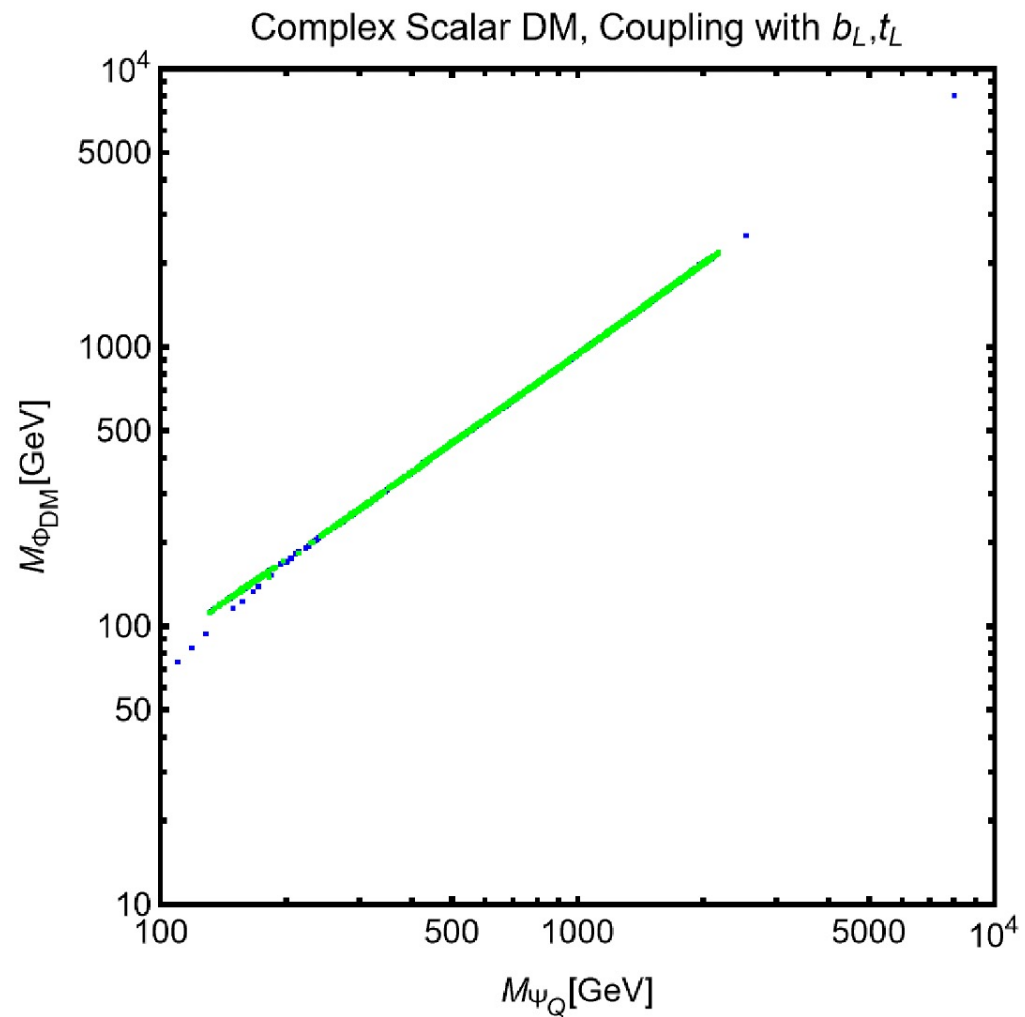
Dirac Fermion



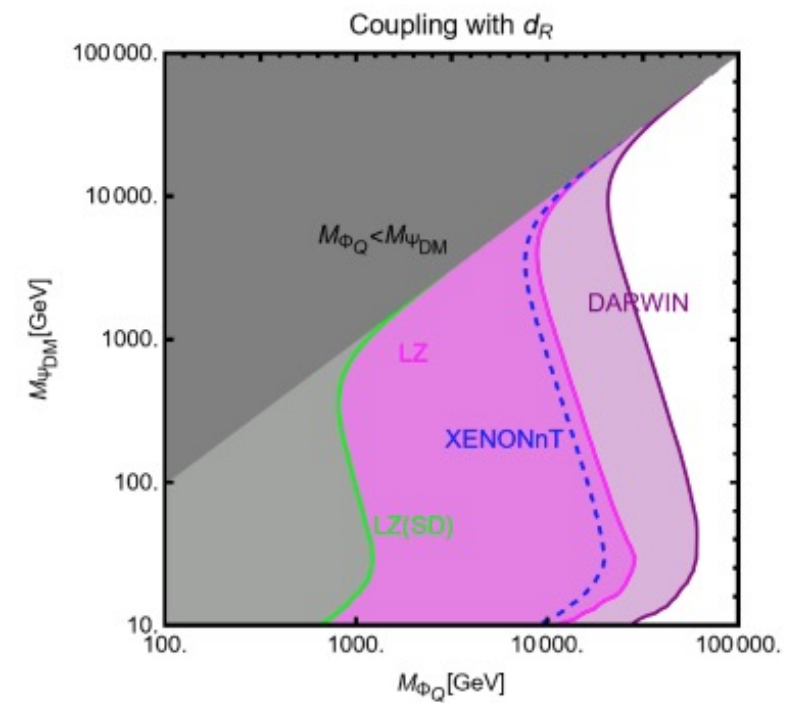
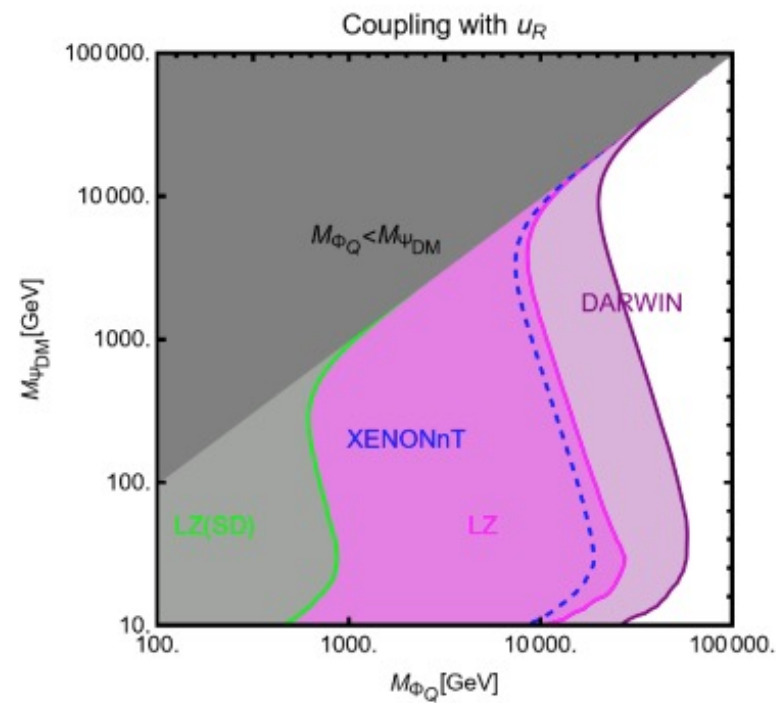
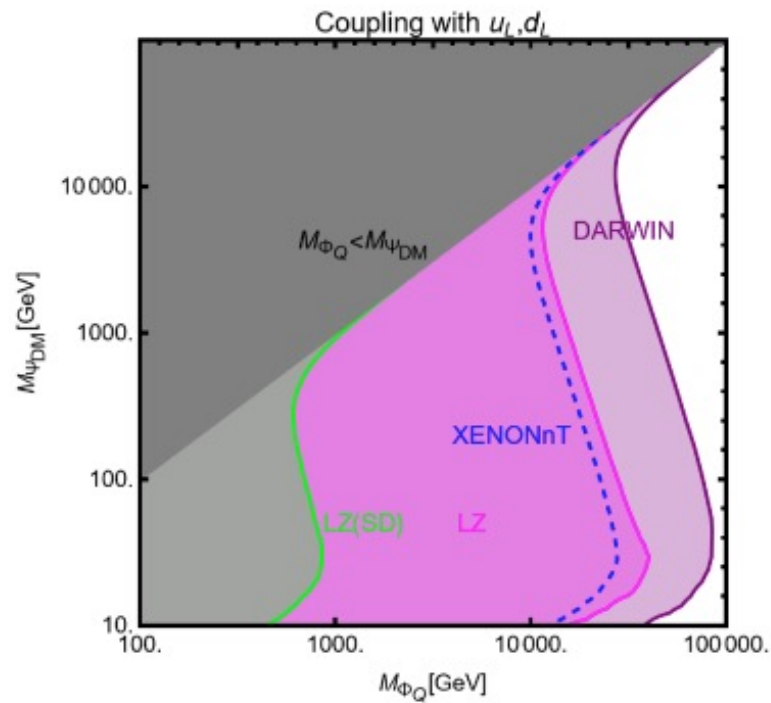
Real Scalar



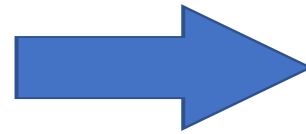
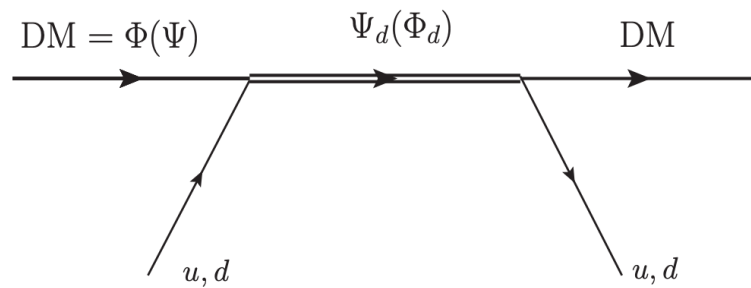
Complex Scalar



Tree-level couplings



Tree level in real DM scenario



$$c^q \left(\Phi_{\text{DM}}^\dagger i \overleftrightarrow{\partial}_\mu \Phi_{\text{DM}} \right) \bar{q} \gamma^\mu q$$

Tree level contribution to the Wilson Coefficient

However,...

This operator doesn't exist for:

- Complex Real
- Majorana Fermion

$$c^q \left(\Phi_{\text{DM}}^\dagger i \overleftrightarrow{\partial}_\mu \Phi_{\text{DM}} \right) \bar{q} \gamma^\mu q$$