

Theoretical predictions for $b \rightarrow s\mu^+\mu^-$

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Based on
arXiv: 2011.09813, 2206.03797, 2305.06301, 2312.14146
in collaboration with
T. Feldmann, D. van Dyk, J. Virto, and M. Reboud

58th Rencontres de Moriond
Electroweak Interactions & Unified Theories
La Thuile, 27-March-2024



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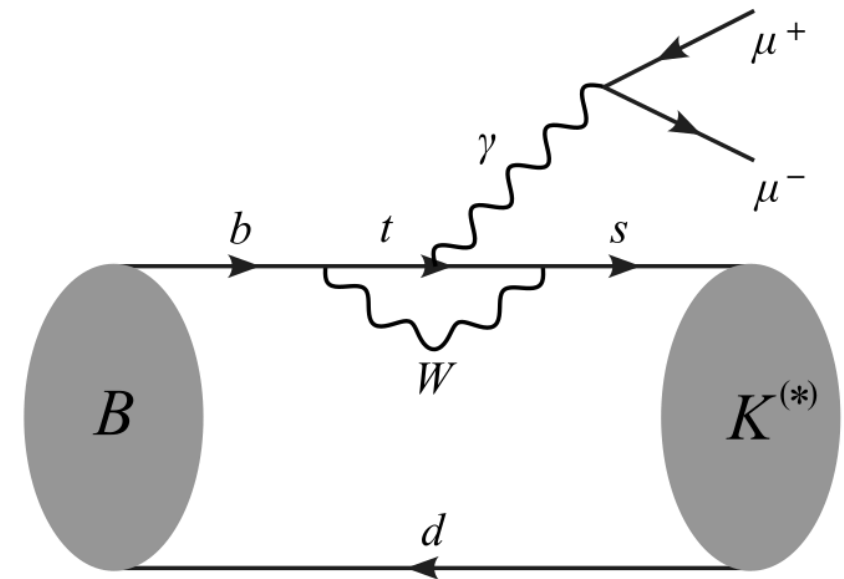
Introduction

$B \rightarrow K^{(*)} \mu^+ \mu^-$ decays

flavour physics \Rightarrow probe the SM through **indirect searches**

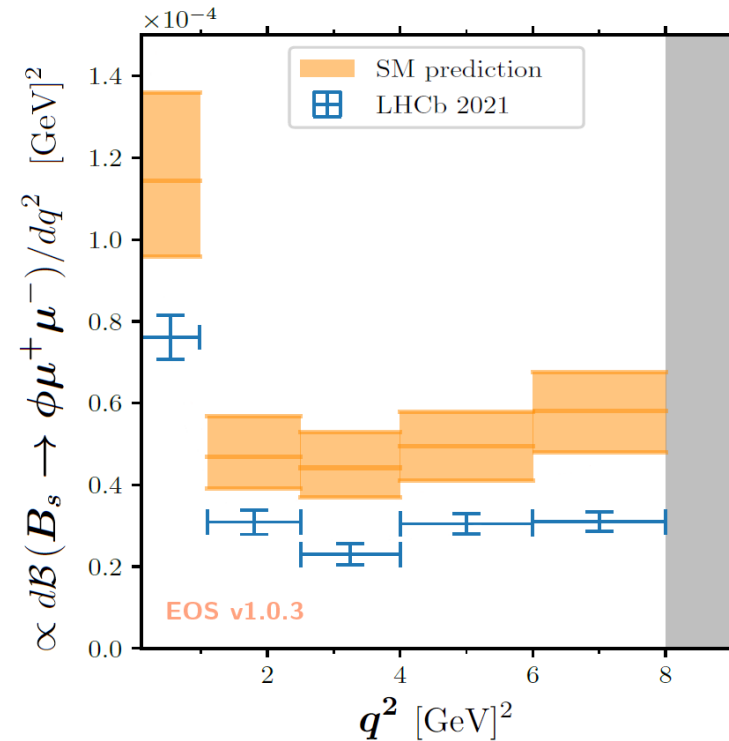
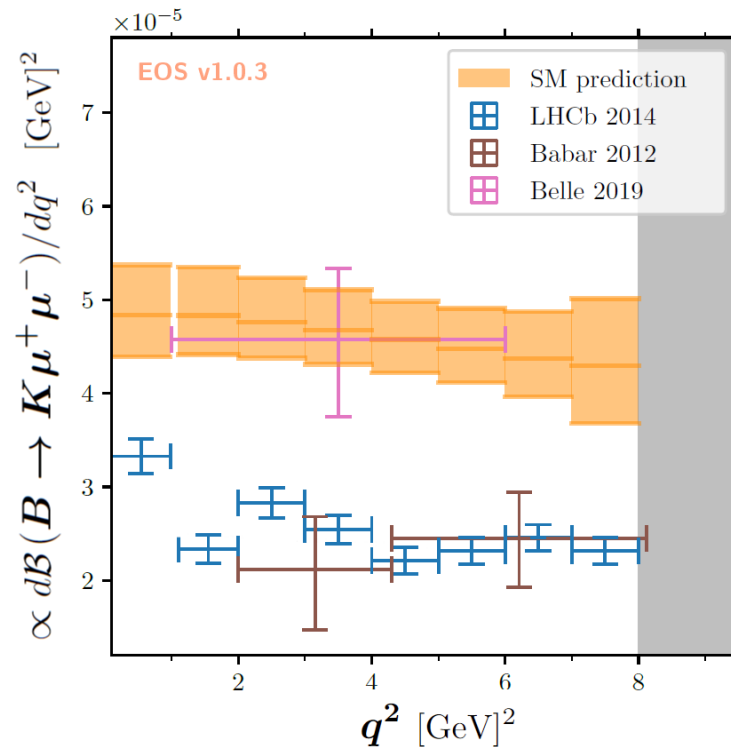
1. **measure** physical observables
2. **calculate** the observables in the SM
3. **compare** measurements and calculations
 \Rightarrow obtain constraints on NP (or new discovery?)

$B \rightarrow K^{(*)} \mu^+ \mu^-$ decays excellent to perform indirect searches since they are suppressed in the SM \Rightarrow **sensitive to NP** (loop, GIM and CKM suppressed)



Tensions in $b \rightarrow s\mu^+\mu^-$ decays

compare exp. measurements and theory predictions for branching ratios and angular observables $B \rightarrow K^{(*)}\mu^+\mu^-$ and $B_s \rightarrow \phi\mu^+\mu^-$



tension (or anomalies) in $b \rightarrow s\mu^+\mu^-$ observables \Rightarrow need to understand this tension (independent from the LFU ratios R_K and R_{K^*})

Theoretical framework

Weak effective theory for FCNC

flavour changing neutral currents (FCNC)
are absent at tree level in the SM

integrate out DOF heavier than the b

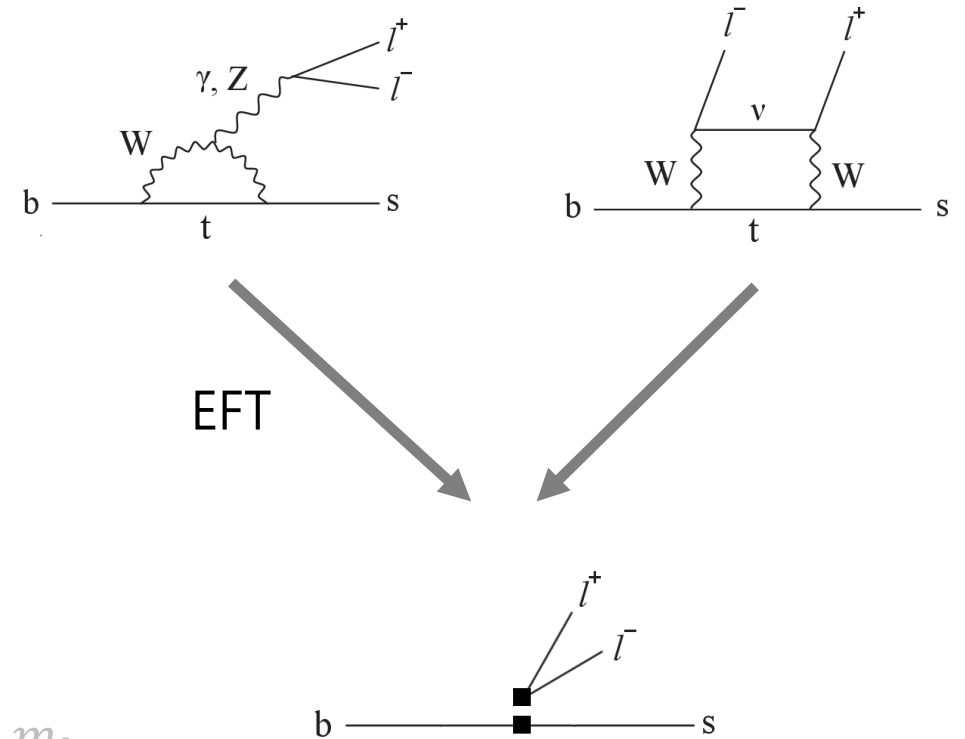


weak effective field theory

transitions described by the **effective Hamiltonian**

$$\mathcal{H}(b \rightarrow s \ell^+ \ell^-) = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_{i=1}^{10} C_i(\mu) O_i(\mu) \quad \mu = m_b$$

C_i Wilson coefficients, O_i effective operators



Decay amplitude for $B \rightarrow K^{(*)} \ell^+ \ell^-$

calculate decay amplitudes precisely to probe the SM

$b \rightarrow s \mu^+ \mu^-$ anomalies: NP or underestimated QCD uncertainties?

$$\mathcal{A}(B \rightarrow K^{(*)} \ell^+ \ell^-) = \mathcal{N} \left[\underbrace{(C_9 L_V^\mu + C_{10} L_A^\mu)}_{\text{Wilson coefficients, leptonic matrix elements}} \mathcal{F}_\mu - \frac{L_V^\mu}{q^2} \underbrace{(C_7 \mathcal{F}_{T,\mu} + \mathcal{H}_\mu)}_{\text{constants } \alpha, V_{CKM} \dots} \right]$$

Wilson coefficients, leptonic matrix elements (and constants $\alpha, V_{CKM} \dots$)

perturbative objects, **small uncertainties**

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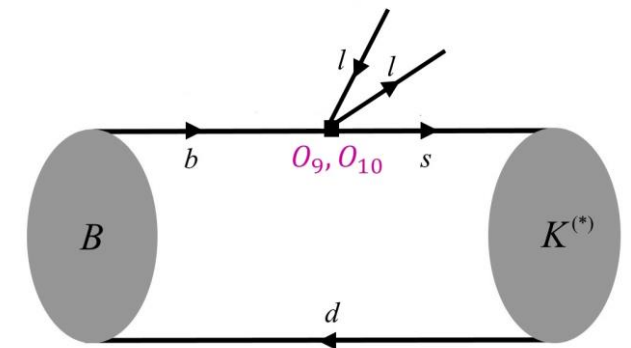
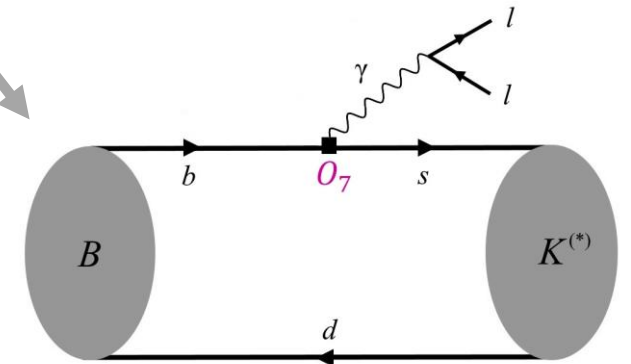
local hadronic matrix elements (MEs)

$$\mathcal{F}_\mu = \langle K^{(*)} | O_{7,9,10}^{\text{had}} | B \rangle \quad O_{7,9,10}^{\text{had}} = (\bar{s} \Gamma b)$$

leading hadronic contributions

non-perturbative QCD objects

moderate uncertainties (3% – 15%)



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non-local hadronic MEs

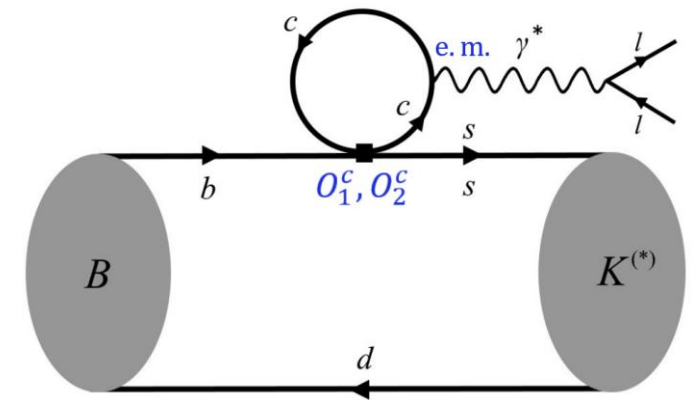
$$\mathcal{H}_\mu = i \int d^4x e^{iq \cdot x} \langle K^{(*)} | T \{ j_\mu^{\text{em}}(x), O_{1,2}^c(0) \} | B \rangle$$

$$O_{1,2}^c = (\bar{s} \Gamma b)(\bar{c} \Gamma c)$$

subleading (?) hadronic contributions

non-perturbative QCD objects

large uncertainties



Theoretical calculations

Local matrix elements calculations

MEs are functions of the momentum transfer squared q^2
non-perturbative techniques are needed to compute MEs \mathcal{F} :

1. lattice QCD (LQCD)

$B \rightarrow K^{(*)}$ and $B_s \rightarrow \phi$ at high q^2

[HPQCD 2013/2023] [FNAL/MILC 2015]

[Horgan et al. 2015] [HPQCD 2023]

small and reducible uncertainties

2. light-cone sum rules (LCSRs)

$B \rightarrow K^{(*)}$ and $B_s \rightarrow \phi$ at low q^2

[Bharucha et al. 2015] [Khodjamirian/Rusov 2017]

[NG/Kokulu/van Dyk 2018]

moderate uncertainties

$B \rightarrow K$ MEs excellent status (need independent calculation at low q^2)

more LQCD results needed for vector states (for high precision K^* width cannot be neglected)

how to combine different calculations and obtain result **whole** semileptonic region?

Local MEs parametrizations

we propose a new parametrization

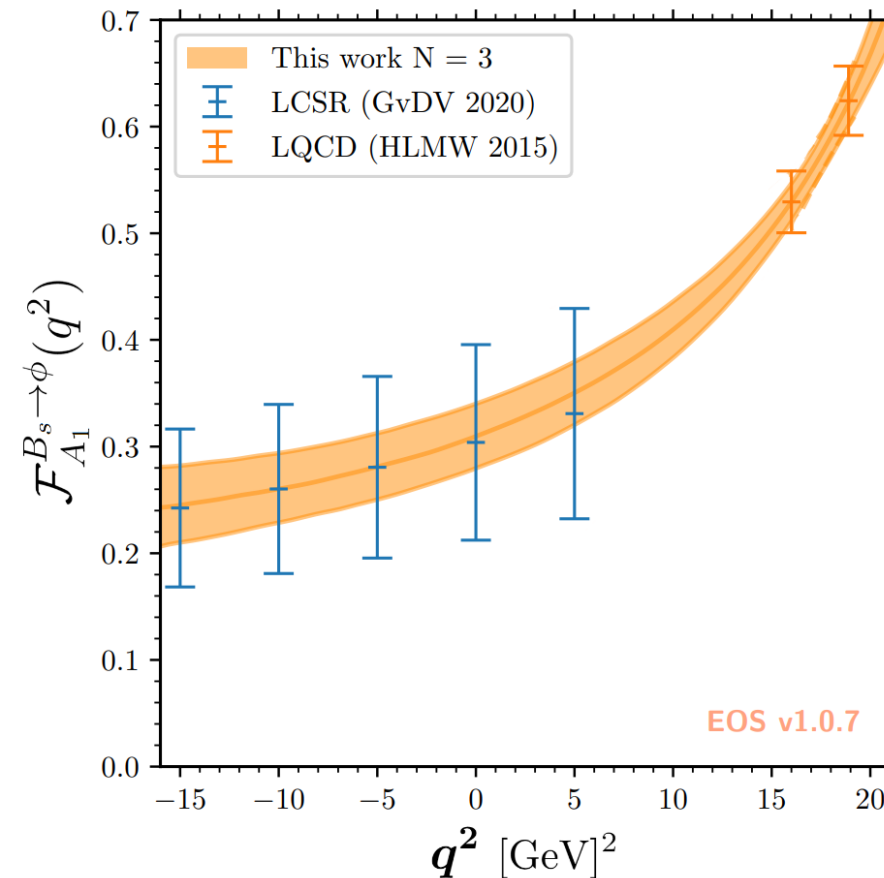
$$\mathcal{F}(q^2) \propto \sum_{k=0}^{\infty} \alpha_k p_k(q^2) \quad \sum_{k=0}^{\infty} |\alpha_k|^2 < 1$$

p_k are known polynomials

fit α_k coefficients to LQCD (and LCSR) results
impose unitarity bounds

first parametrization that consistently implements
analyticity and unitarity bounds
 \Rightarrow control systematic uncertainties

obtain numerical results for $B \rightarrow K^{(*)}$ and $B_s \rightarrow \phi$
local MEs in the **whole semileptonic region**

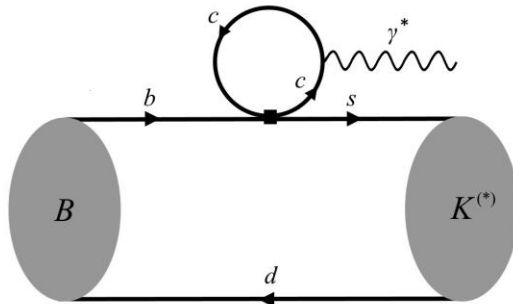


Non-local MEs calculations

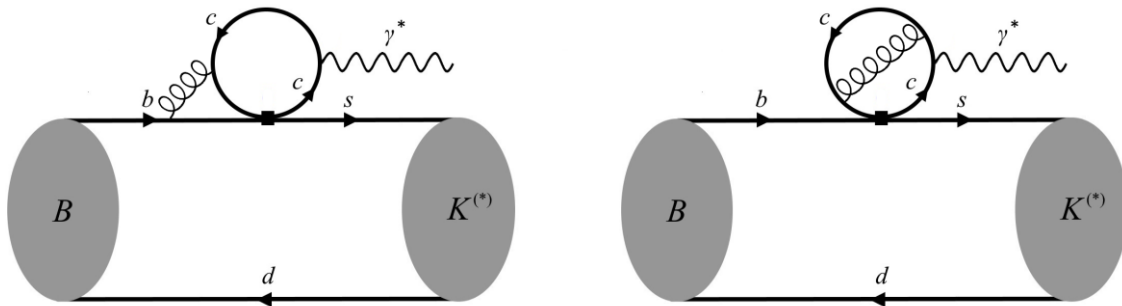
- compute the \mathcal{H} using a light-cone OPE at **low q^2**

$$\mathcal{H}(q^2) = C(q^2)\mathcal{F}(q^2) + \tilde{C}(q^2)\mathcal{V}(q^2) + \dots$$

leading power (LO in α_s)

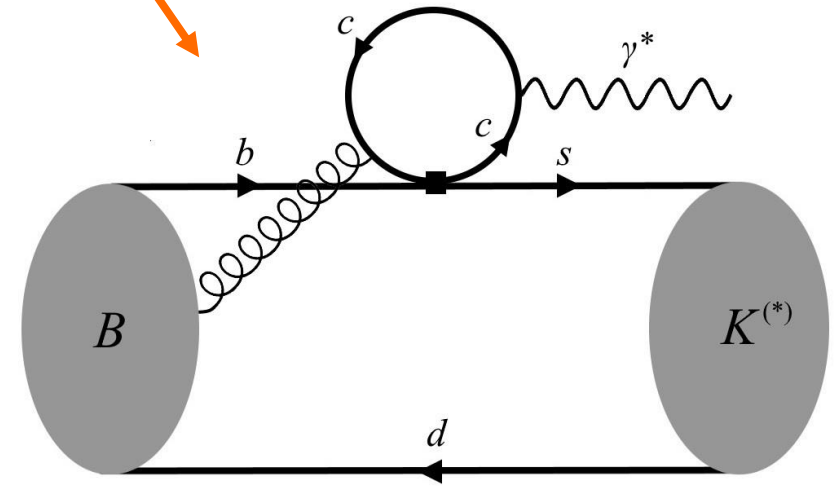


+ hard gluons (α_s) corrections



[Bell/Huber 2014] [Asatrian/Greub/Virto 2019]

soft gluon correction
non-perturbative
 \Rightarrow not α_s suppressed



[Khodjamirian et al. 2010]
[NG/van Dyk/Virto 2020]

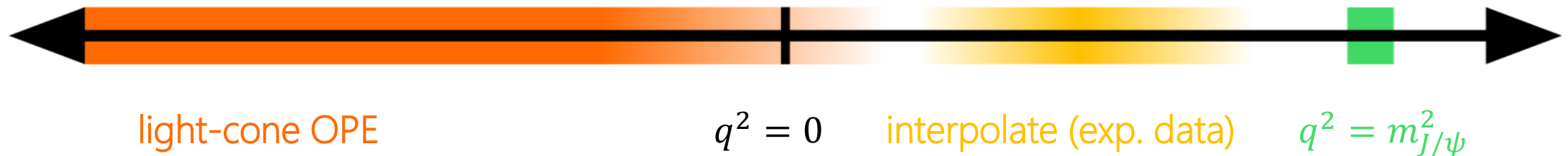
Non-local MEs calculations

1. compute the \mathcal{H} using a light-cone OPE at **low q^2**

$$\mathcal{H}(q^2) = C(q^2)\mathcal{F}(q^2) + \tilde{C}(q^2)\mathcal{V}(q^2) + \dots$$

2. extract \mathcal{H}_λ at $q^2 = m_{J/\psi}^2$ from $B \rightarrow K^{(*)}J/\psi$ and $B_s \rightarrow \phi J/\psi$ measurements (decay amplitudes independent of the local MEs)

3. **new approach: interpolate** these two results to obtain theoretical predictions in the **low q^2 ($0 < q^2 < 8 \text{ GeV}^2$)** region \Rightarrow compare with experimental data



Non-local MEs predictions

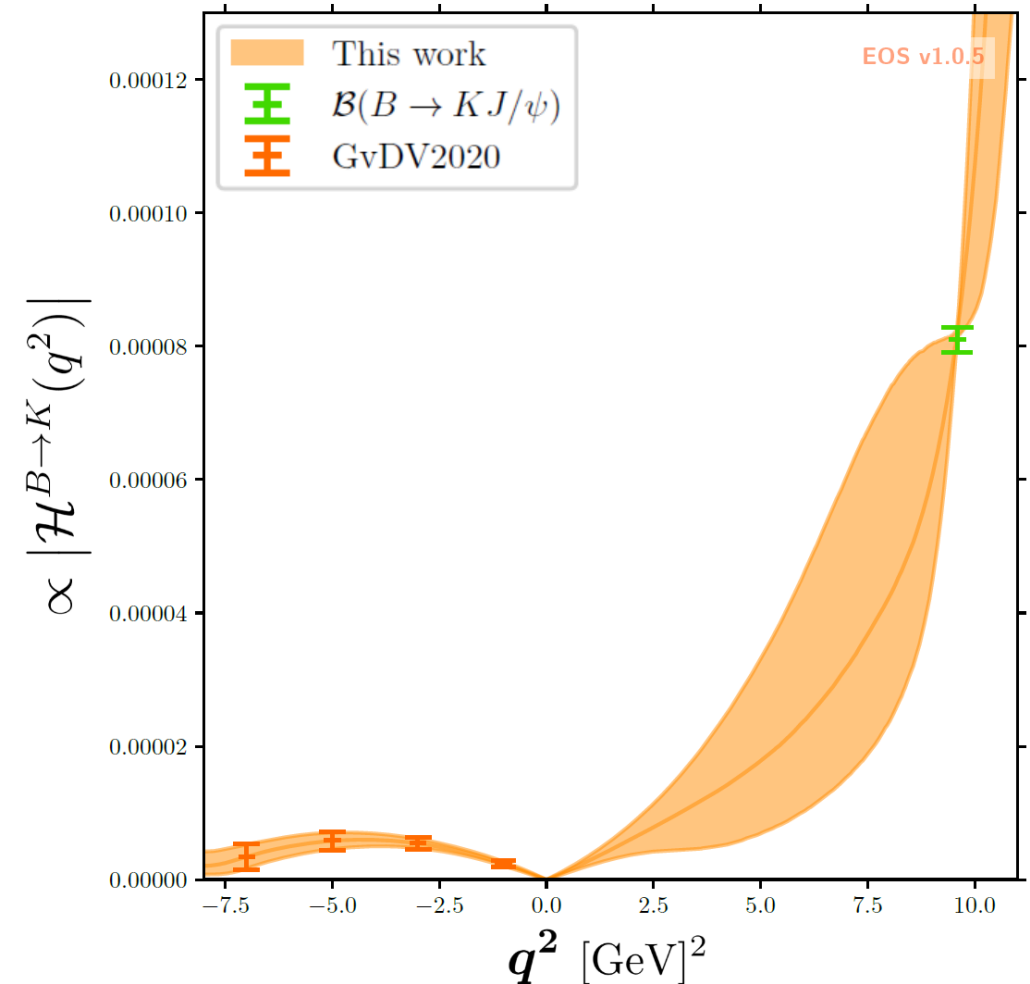
similar approach to local MEs \mathcal{F}_λ

$$\mathcal{H}(q^2) \propto \sum_{n=0}^{\infty} \beta_n p_n(q^2) \quad \sum_{k=0}^{\infty} |\beta_k|^2 < 1$$

fit β_k coefficients to OPE and $B \rightarrow K^{(*)}J/\psi$ data
impose unitarity bounds

first unitarity bounds for non-local MEs \mathcal{H}
 \Rightarrow control systematic uncertainties

obtain numerical results for $B \rightarrow K^{(*)}$ and $B_s \rightarrow \phi$
non-local MEs **below 8 GeV²**



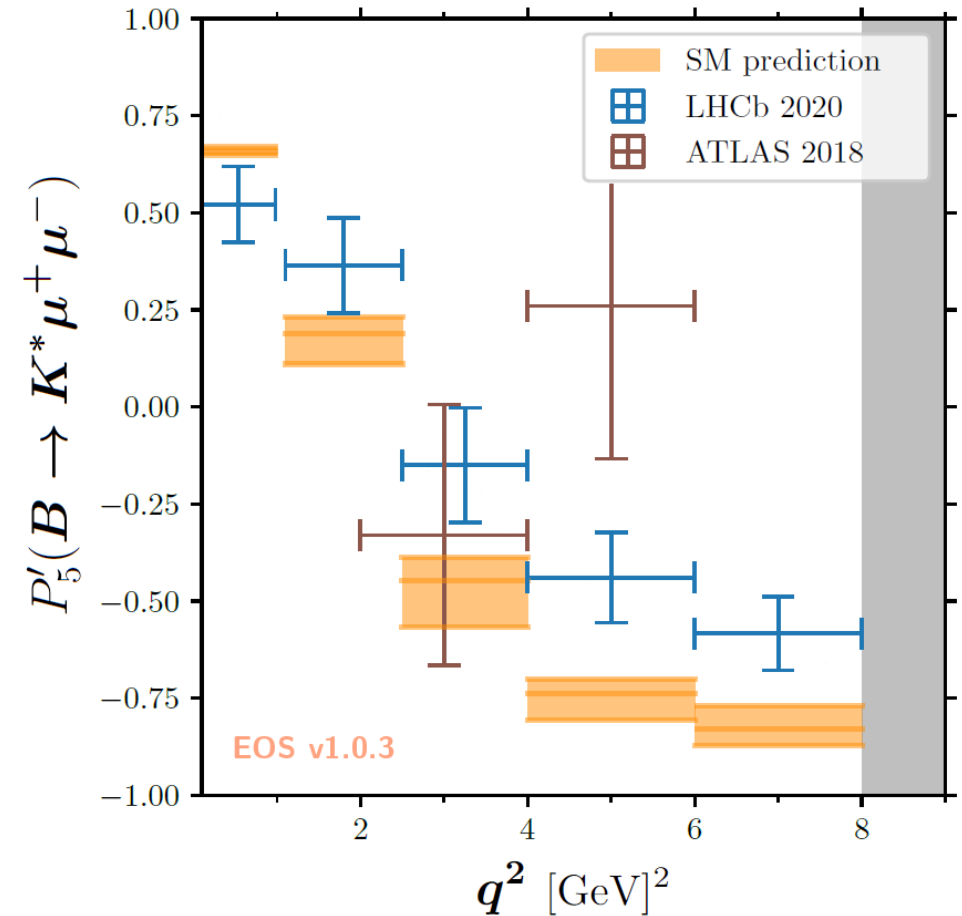
SM predictions and
confrontation with data

SM predictions vs. data

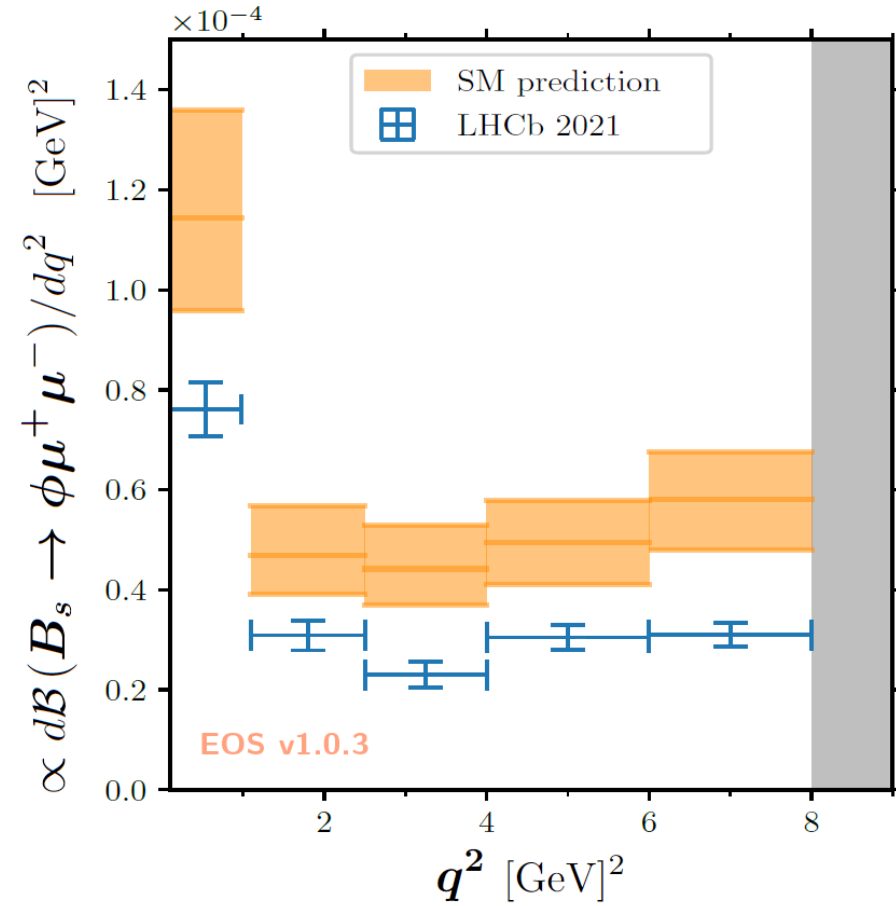
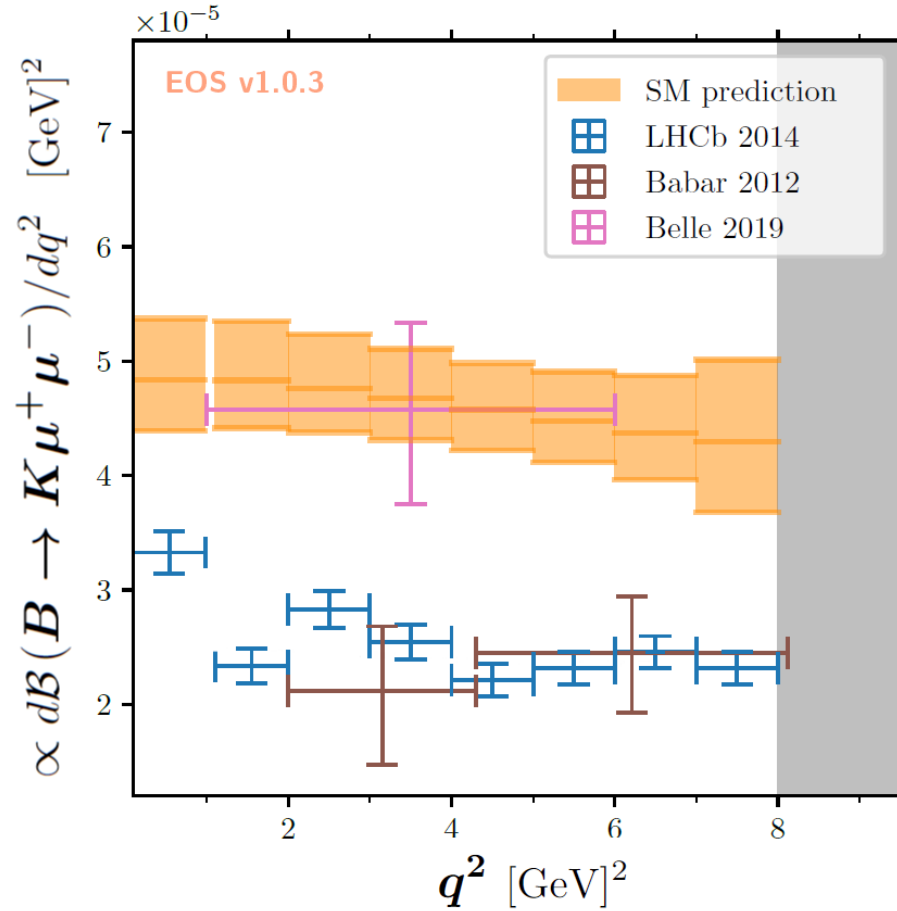
predict observables using our \mathcal{F}_λ and \mathcal{H}_λ results:

BRs and angular observables
for $B \rightarrow K^{(*)} \mu^+ \mu^-$, and $B_s \rightarrow \phi \mu^+ \mu^-$

- theory uncertainties mostly due to \mathcal{F}_λ
- progress in \mathcal{H}_λ calculations urgently needed
- more measurements on the way



SM predictions vs. data



[NG/Reboud/van Dyk/Virto 2022]

coherent tensions between SM predictions and data

WET fits

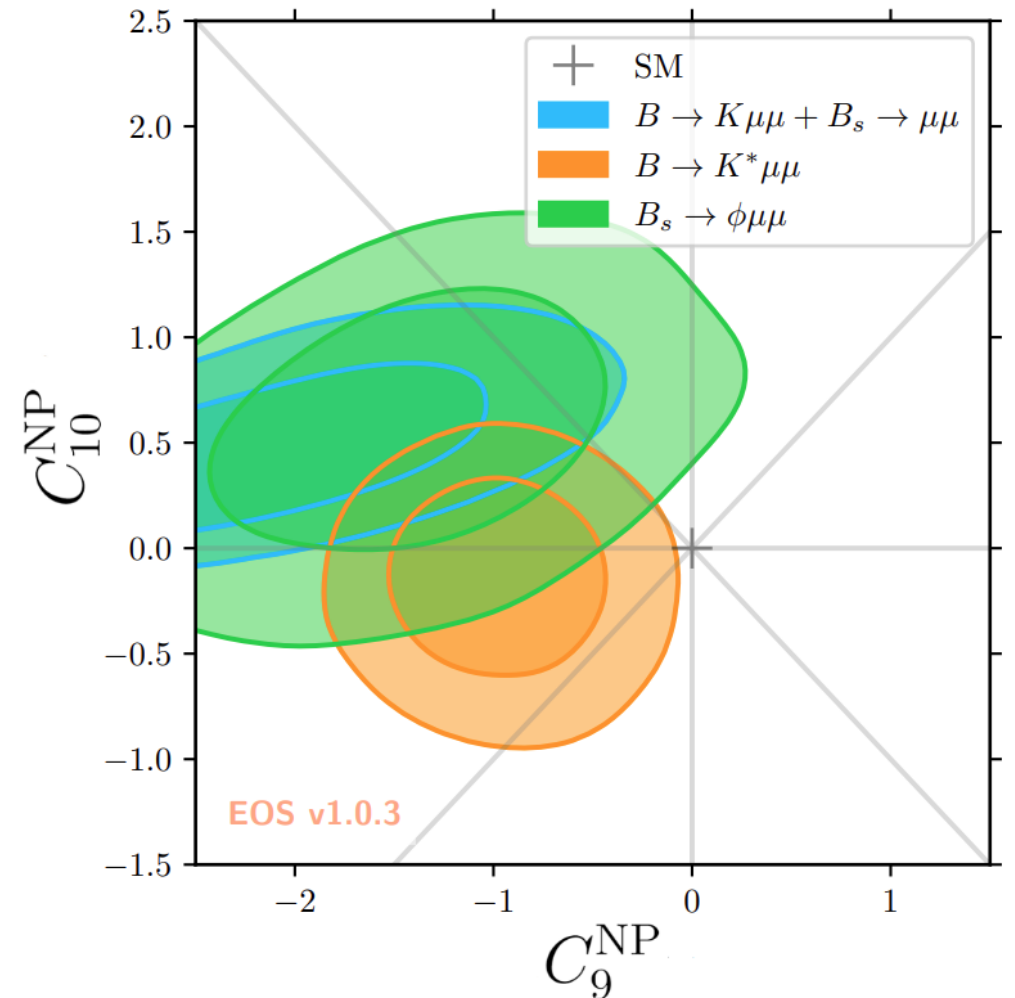
$$C_{9,10} = C_{9,10}^{\text{SM}} + C_{9,10}^{\text{NP}}$$

fit the Wilson coefficients C_9^{NP} and C_{10}^{NP}
to the available experimental measurements

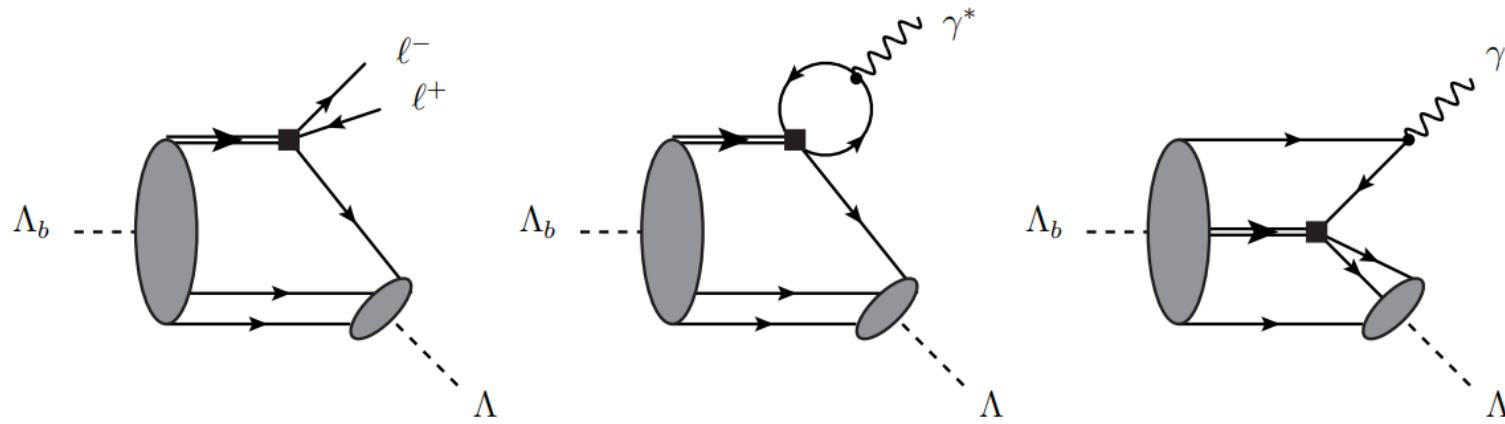
pulls (p value of the SM hypothesis):

- 5.7σ for $B \rightarrow K\mu^+\mu^- + B_s \rightarrow \mu^+\mu^-$
- 2.7σ for $B \rightarrow K^*\mu^+\mu^-$
- 2.6σ for $B_s \rightarrow \phi\mu^+\mu^-$

current predictions for
non-local MEs \mathcal{H}_λ cannot explain this tension



Aside: $\Lambda_b \rightarrow \Lambda \ell^+ \ell^-$ decays



[Feldmann/NG 2023]

if $b \rightarrow s \mu^+ \mu^-$ anomalies are due to New Physics \Rightarrow same shift expected in $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$
but systematic effects are different

already measured by LHCb \Rightarrow new and more precise measurements on the way

progress needed in theory calculations (no estimate of charm-loop beyond naïve factorization)

first calculation of “annihilation” contributions in [Feldmann/NG 2024]

Summary and conclusion

Summary and conclusion

1. improved parametrization for local MEs \mathcal{F}_λ with unitarity bounds

combine LQCD (and LCSR) inputs to get new results for \mathcal{F}_λ in $B \rightarrow K^{(*)}\ell^+\ell^-$ and $B_s \rightarrow \phi\ell^+\ell^-$

2. new theoretical predictions for \mathcal{H}_λ combining our OPE calculation and $B \rightarrow K^{(*)}J/\psi$ data

innovative approach — use unitarity bound to control \mathcal{H}_λ uncertainties

3. new and precise SM predictions for observables in $B \rightarrow K^{(*)}\ell^+\ell^-$ and $B_s \rightarrow \phi\ell^+\ell^-$ decays

coherent deviations between SM and data in $B \rightarrow K^{(*)}\ell^+\ell^-$ and $B_s \rightarrow \phi\ell^+\ell^-$ decays

4. progress on the theory side needed more than ever

Thank you!