



Precision measurements: paths to New Physics at low-energies and high-intensities

Ana M. Teixeira

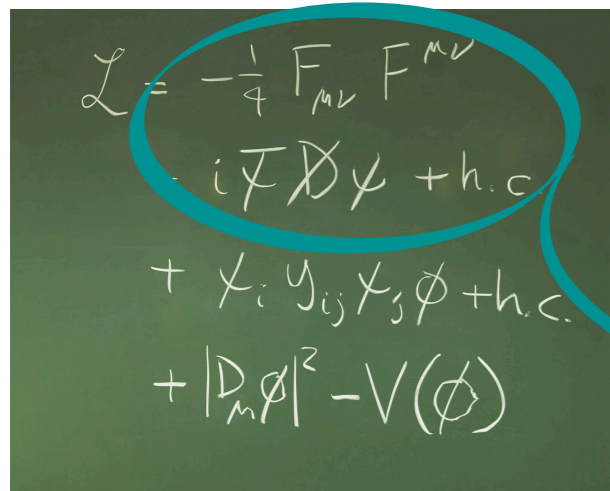
Laboratoire de Physique de Clermont - LPC



CS IN2P3 - IPHC, 24 June 2024

"Tests de précision à basse énergie des interactions
fondamentales"




$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - i\bar{\psi}\not{D}\psi + \text{h.c.} + \chi_i Y_{ij} \chi_j \phi + \text{h.c.} + |D_\mu \phi|^2 - V(\phi)$$

A history of success!

Minimal formulation allowing to understand (and predict) most of the phenomena in particle physics

Electroweak & strong interactions in the **Standard Model**: tested with **impressive precision** since first dedicated searches

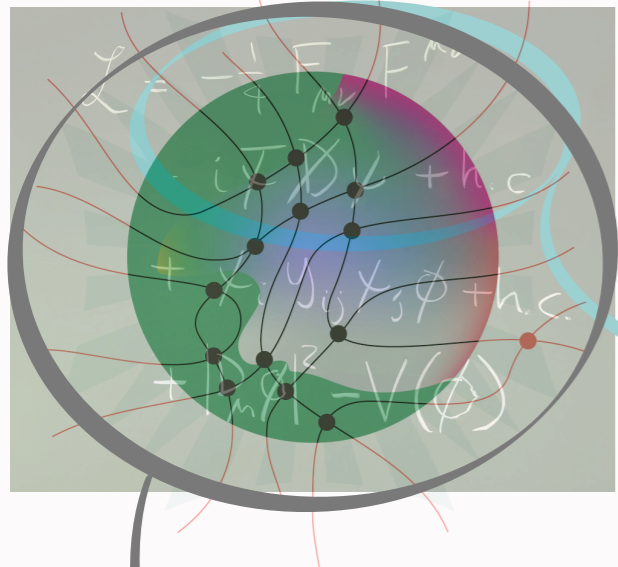
Theoretically, not "the perfect" picture: unification !?

$SU(3)_c \otimes SU(2)_L \otimes U(1) \leftrightarrow$ Unified interactions?

At which scale ? gravitation @ Λ_{Planck} ... **Desert of scales...**

Strong interactions - confinement of quarks in matter (nucleons, nuclei, atoms...)
non-perturbative effects very hard to deal with

The SM: a history of success ... and 3+n caveats



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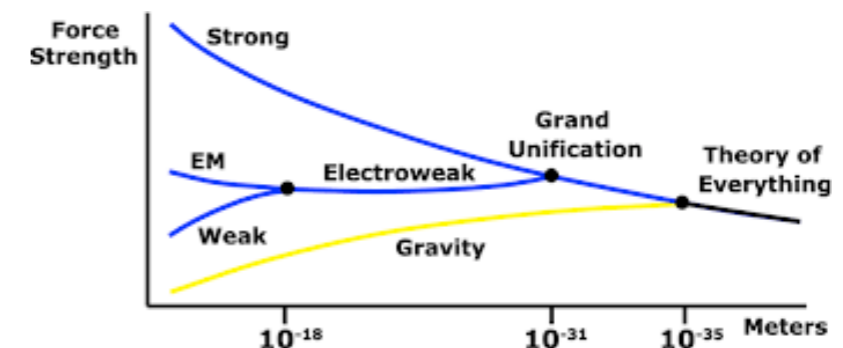
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(4th) missing interaction: **Gravity!** - not included in the SM formulation
non-negligible effects above the Planck scale



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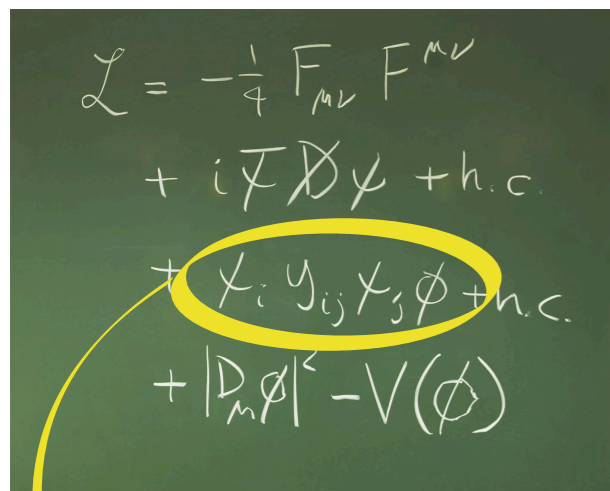
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(nucleons, nuclei, atoms...)

non-perturbative effects very hard to deal with

Higgs: CP even scalar doublet, in good agreement with expectation
Verification of the mechanism of **electroweak symmetry breaking**

Further *theory woes*: **fine-tuning** issues (hierarchy problem),
implications for **vacuum stability**, ...



$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}\not{D}\psi + h.c. + y_{ij} \psi_i \psi_j \phi + h.c. + |D_\mu \phi|^2 - V(\phi)$$

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Flavour in the Standard Model:

interactions between *fermion* families (and the *Higgs*)

Flavour has paved the way to the SM!

From prediction of *charm* quark ... to the existence of 3 *families*!

Y_{ij}^u, Y_{ij}^d and $Y_{ij}^\ell \rightsquigarrow$ encode flavour dynamics
(masses, mixings & CP violation)

⇒ **Flavour-universal** gauge interactions

⇒ **Lepton sector:** 3 massive ℓ^\pm ; massless ν ; no leptonic mixing...

⇒ **Quark sector:** 6 massive states, $V_{CKM}^{ij} W^\pm \bar{q}_i q_j'$

SM flavour & CP: *accidental symmetries* (lepton & baryon number conservation, conservation of lepton flavours, lepton flavour universality of gauge interactions) and the “**CKM paradigm**”

A number of **theoretical** caveats... and observations unaccounted for in the SM:
baryon asymmetry of the Universe, viable dark matter candidate, ν oscillations

- **Matter dominated Universe:** explaining **the baryon asymmetry of the Universe (BAU)**
- (i) **initial asymmetric composition** ✗ (incompatible with inflation)
 - (ii) **statistical fluctuations during evolution** ✗ (negligible effects)
 - (iii) **large scale spatial separation** ✗ (incompatible with evolution of primordial Universe)
- ... **Dynamical generation! "Baryon-genesis"** ✓

Sakharov's conditions for a (successful) BAU

a priori, all are present in the SM! (**electroweak baryogenesis**)

- If originally symmetric Universe, **baryon number violation**

Sphaleron production \Rightarrow ***B & L number violation***

- Differentiate matter from antimatter, **CP violation**

CPV from CKM mechanism ***highly suppressed...***

- Suppress inverse processes, **out of (thermal) equilibrium**

Strong 1st order EW phase transition ? soft crossover for a "heavy Higgs" (125 GeV)

Explain the BAU \Rightarrow **BSM physics** is also required!

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► **Matter dominated Universe: Explaining the baryon asymmetry of the Universe (BAU)**

Explain the **BAU** \Rightarrow **BSM physics** is also required!

► **Dark Matter in the Universe**

PLANCK, WMAP, ... & Galactic dynamics \Rightarrow most **matter is "dark"** $\Omega_{\text{CDM}} = 0.259 \pm 0.006$

"ordinary (SM) matter" - a tiny fraction of mass-energy density $\Omega_{\text{b}} = 0.049 \pm 0.001$

Dark matter candidate: massive, non-luminous, no strong interactions...
(at best) weakly interacting, stable!

No such candidate in the Standard Model!

The need for New Physics



IN2P3
Les deux infinis

A number of **theoretical caveats...** and observations unaccounted for in the **SM**:
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▶ **Matter dominated Universe: Explaining the baryon asymmetry of the Universe (BAU)**

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▶ **Dark Matter in the Universe**

Dark matter candidate: necessarily from **New Physics!**

▶ **Neutrino oscillations** \Rightarrow **massive neutrinos** and **non-trivial leptonic mixing!**

1st "laboratory" discovery of *physics beyond the SM (BSM)*

New (**Majorana**) fields? New sources of **CP violation?**

(LNV & CPV \rightarrow crucial for BAU!)

The need for New Physics

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New Physics is indeed needed - but which new physics model?

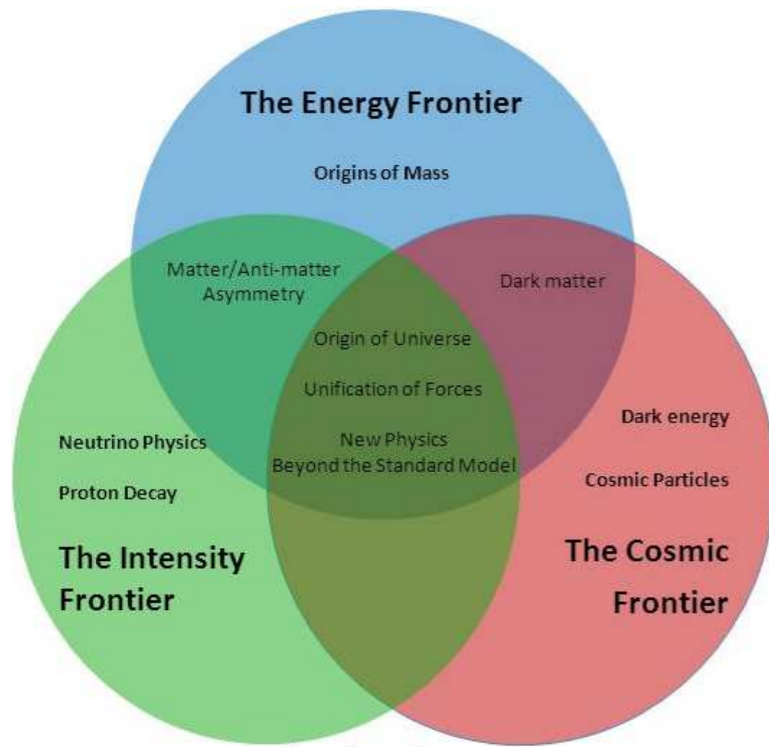
Models of New Physics generically introduce (at high scales):

- (i) new sources of **CP** and **flavour violation**
- (ii) new **Lorentz structure** (beyond V-A)
- (iii) new (heavy) **propagators**

How do we search for these new ingredients?



New Physics searches at three "experimental" frontiers:



- ▶ **Cosmic frontier:** cosmological impact, evolution of the Universe (observation)
- ▶ **High-energy frontier:** new heavy states produced if sufficiently large collision energy (lepton or hadron beams)
- ▶ **High-intensity frontier:** indirect "virtual" effects of NP states (feebly coupled or very rare processes) for high "luminosities"

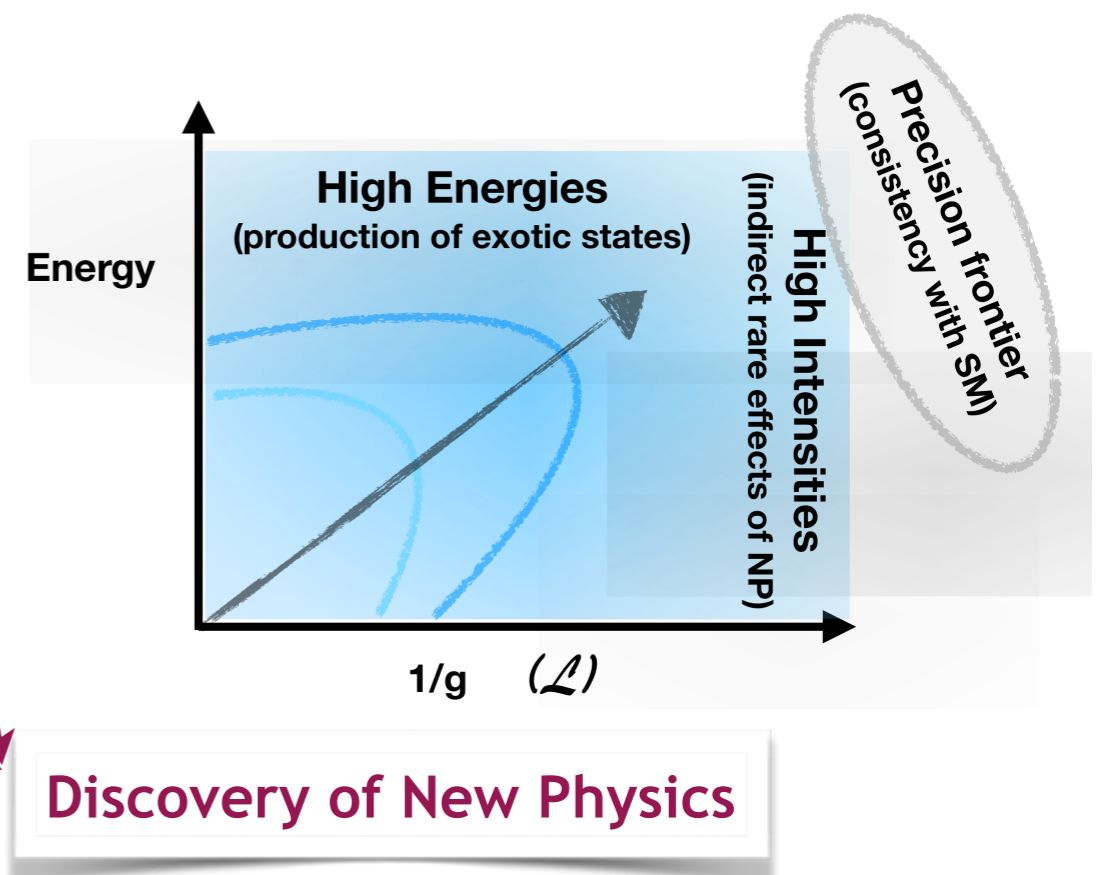
Precision tests of fundamental laws (SM, ...)

at high- and low-energies (& "tabletop energies")

⇒ test predictions to unprecedented precision

⇒ reveal tensions and inconsistencies

⇒ challenge "null" expectations
(conservation, forbidden processes, ...)



Heavy flavours

(K, D and B meson oscillations and decays, ...)

Nucleons, nuclei and atoms

(EDM, weak decays, neutral currents, strong interaction tests, ...)

Charged leptons

(cLFV processes, $(g - 2)_\ell$, EDMs, ...)

Precision tests of the SM

Proton decay

Light weakly coupled particles

(axions and ALPs, dark γ , ...)

Neutrinos

(oscillations, nature, mass interaction with matter, ...)

Weak equivalence principle

Nucleons, nuclei and atoms
(EDM, weak decays, neutral currents,
strong interaction tests, ...)

**Precision tests
of the SM**

Charged leptons
(cLFV processes, $(g - 2)_\ell$,
EDMs, ...)

*Weak equivalence
principle*

- The need for New Physics - searches at three frontiers
- Theory approaches: **effective** approach (model-independent)
- **cLFV rare decays**: overview & impact for NP searches of cLFV in the **muon** sector
- CP violation in the hadron sector: **neutron EDM**
- **EW precision tests in nuclear β decays**: superallowed transitions & **CKM** anomaly searches for **non (V-A) interactions**
CP violation in beta decays
- Testing the weak equivalence principle: **gravity effects on antimatter**

Extensive range of topics! Covered by a particle physics phenomenologist!

Highlight most relevant aspects - unorthodox approach

A tiny subset... subject to time constraints and personal "bias" 😊

Complementary to subsequent presentations

Precision tests of the SM: constraining New Physics



EFT approach to New Physics

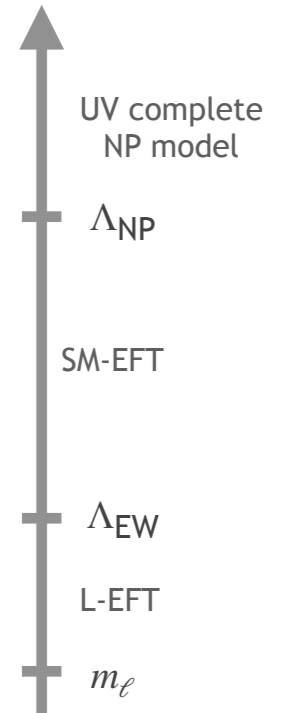
SM interpreted as a **low-energy limit** of a (complete, yet unknown) NP model
 \Rightarrow **Model-independent, effective field theory approach (EFT)**

$$\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \sum_{n \geq 5} \frac{1}{\Lambda^{n-4}} \mathcal{C}^n(g, Y, \dots) \mathcal{O}^n(\ell, q, H, \gamma, \dots)$$

effective operators

(unknown) NP scale effective coefficients

$\mathcal{O}^5 \rightsquigarrow$ Weinberg operator (m_ν)
 $\mathcal{O}^6 \rightsquigarrow$ flavoured contributions
 (among many others!)



Derive the new "effective" interactions (vertices, ...), and compute **contributions to observables**
 Agnostic approach, allowing to generically parametrise NP effects
 on observables **forbidden** in SM and/or observables **suggesting deviations** from SM

$$\mathcal{A} \sim \mathcal{A} \left(\frac{\mathcal{C}^6}{\Lambda_{\text{NP}}^2} \right) + \dots$$

$$\mathcal{A} \sim \mathcal{A}^{\text{SM}} + \mathcal{A} \left(\frac{\mathcal{C}^6}{\Lambda_{\text{NP}}^2} \right) + \dots$$

\Rightarrow master SM prediction!

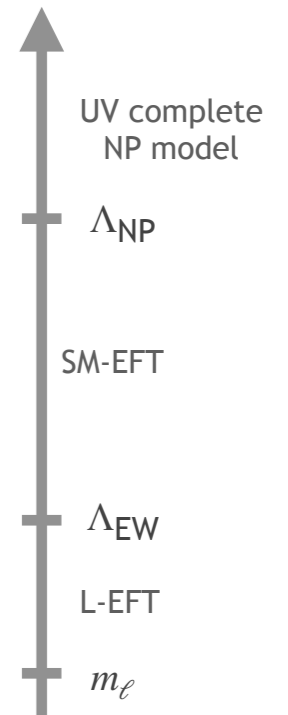
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(unknown) NP scale
effective coefficients

effective operators



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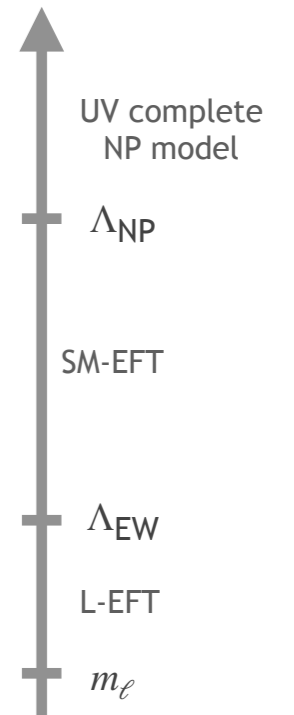
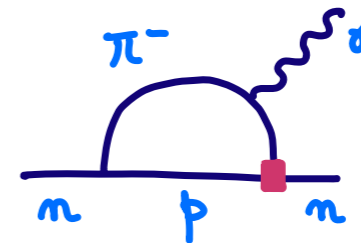
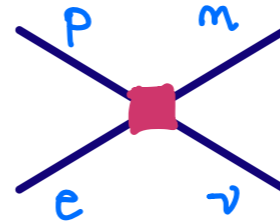
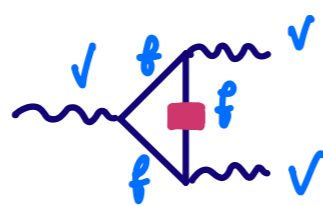
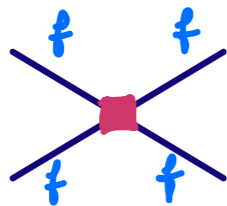
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Λ^{n-4} \rightarrow (unknown) NP scale
 \mathcal{C}^n \rightarrow effective coefficients
 \mathcal{O}^n \rightarrow effective operators



Cast **current data** (limits, ...) in terms of \mathcal{C}_{ij}^n and Λ_{NP}
 and attempt at **inferring info** on the **dominating operator**, and **scale of NP**
 \Rightarrow Beyond $(V - A)$ structure? New vector/axial, (pseudo)scalar or tensor currents?
 Flavour violation beyond SM flavour paradigm?
 \Rightarrow But **many unknowns**: minimal assumptions must be made, e.g.

"natural" $\Lambda_{\text{NP}} \rightarrow$ constrain \mathcal{C}_{ij}^n

"natural" $\mathcal{C}_{ij}^n \approx 1 \rightarrow$ hint on Λ_{NP}

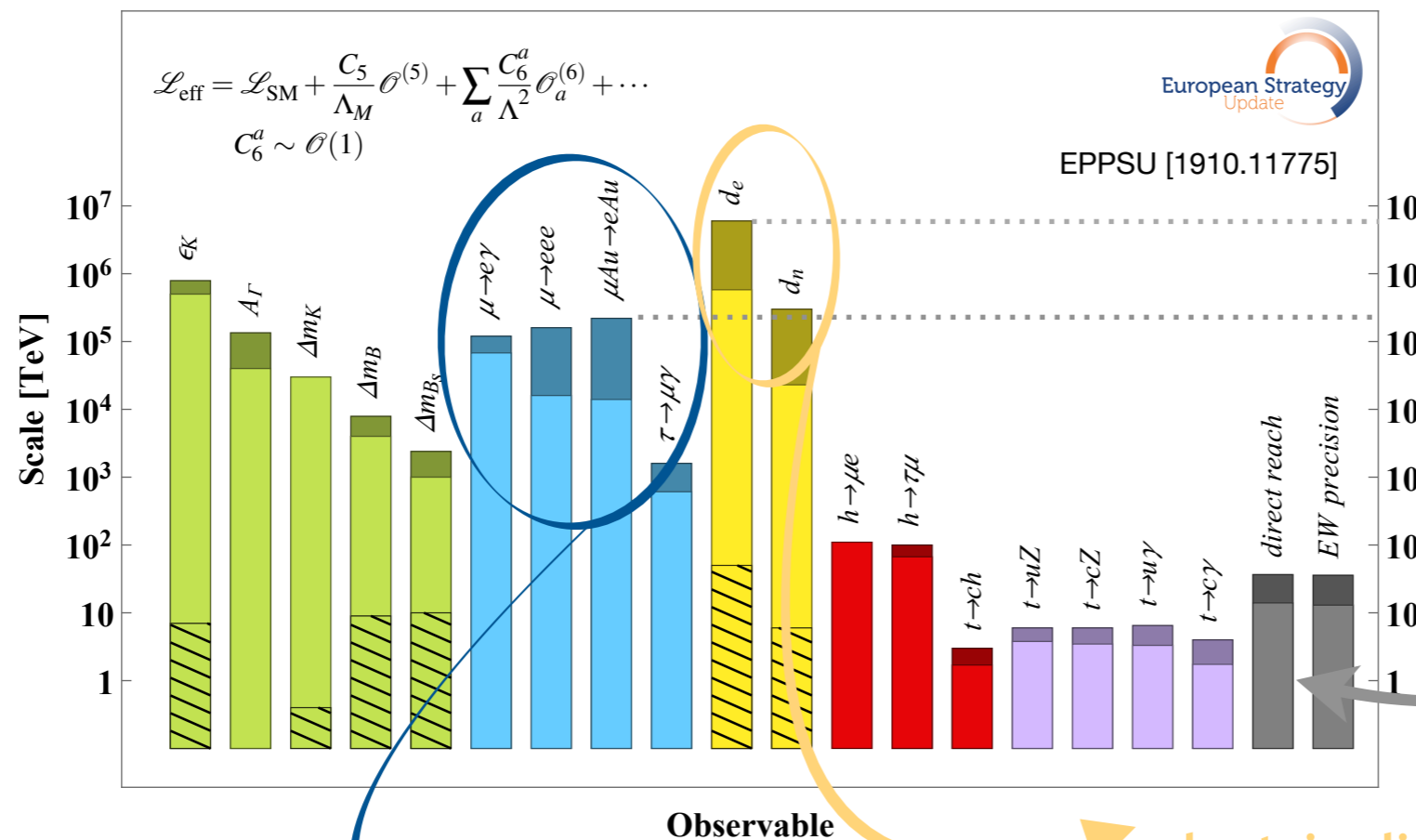
The probing power of flavour & CPV

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Cast **current "flavoured" data** in terms of \mathcal{C}_{ij}^6 and Λ_{NP} : $\mathcal{C}_{ij}^6 \approx 1 \Rightarrow$ bounds on Λ_{NP}



Flavour & CPV observables:
probes sensitive to very high NP scales

$$\Lambda_{\text{NP}} \sim \mathcal{O}(10^5 \text{ TeV})$$

well beyond collider's reach!

electric dipole moments
charged lepton flavour violating observables!

New Physics searches at the high-intensity frontier: (lepton) flavours and CP violation



Charged lepton flavour violation: muon sector opportunities



Lepton flavours: from ν oscillations...

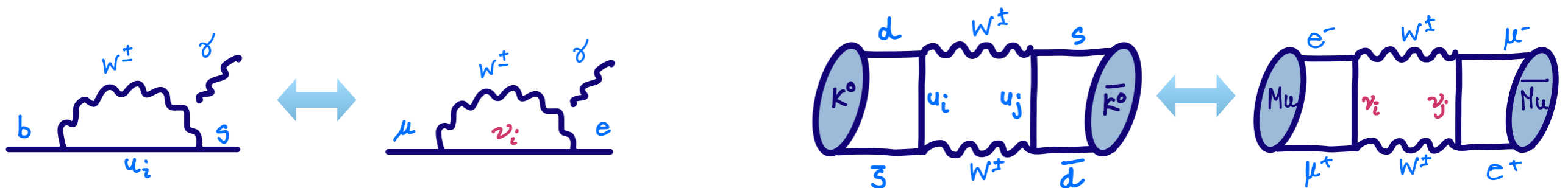
SM hadron sector: plethora of transitions and decays (and CPV!); CKM paradigm
 theoretical predictions increasingly more precise (under control)

SM lepton sector: (strictly) massless neutrinos
 conservation of total lepton number and lepton flavours
 lepton flavour universality preserved (only broken by Yukawas)
 tiny leptonic EDMs (4-loop... $d_e^{\text{CKM}} \leq 10^{-38} e \text{ cm}$)

Neutrino oscillations: SM description insufficient! Added complexity to the flavour problem...

In propagation $\nu_\alpha \rightarrow \nu_\beta \rightarrow \dots \Rightarrow$ oscillations signal the violation of neutral lepton flavours

\Rightarrow Violation of lepton flavour in neutral lepton sector opens a wide door
 to flavour violation in the charged lepton sector



Similar flavour violating transitions in quarks & leptons ?!

Lepton flavours: from ν oscillations...

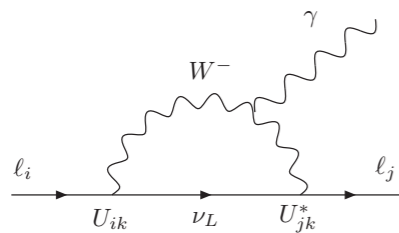
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Neutrino oscillations: SM description insufficient! Added complexity to the flavour problem...
Extend the SM to accommodate $\nu_\alpha \leftrightarrow \nu_\beta$: assume most minimal extension SM_{m_ν}

[SM_{m_ν} = “ad-hoc” m_ν (Dirac), U_{PMNS}]

In SM_{m_ν} : total lepton number still conserved (LNC)



BUT! flavour violation in neutral leptons \Rightarrow charged leptons as well!

cLFV possible... but not observable!! $\text{BR}(\mu \rightarrow e\gamma) \sim 10^{-54}$

lepton EDMs still beyond observation (2-loop contributions from δ_{CP})

cLFV, LNV, lepton EDMs, ...: observation of SM-forbidden leptonic modes

\Rightarrow **Discovery of New Physics!** (possibly before direct signal @ LHC)

Lepton flavours: from ν oscillations...

SM hadron sector: plethora of transitions and decays (and CPVI) • CKM paradigm

SM l

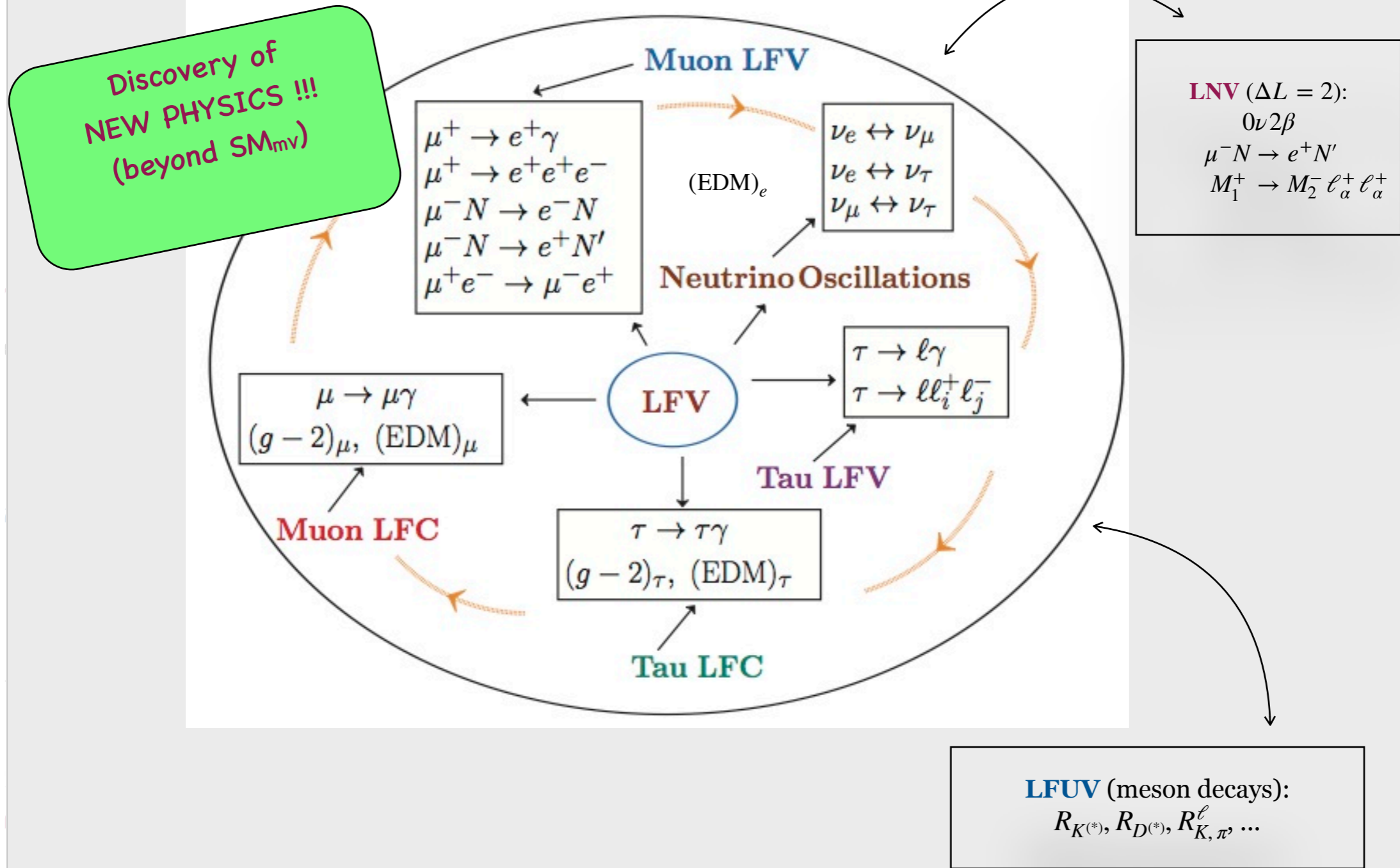
Neutr

E

In SM

l_i

cLF



!!

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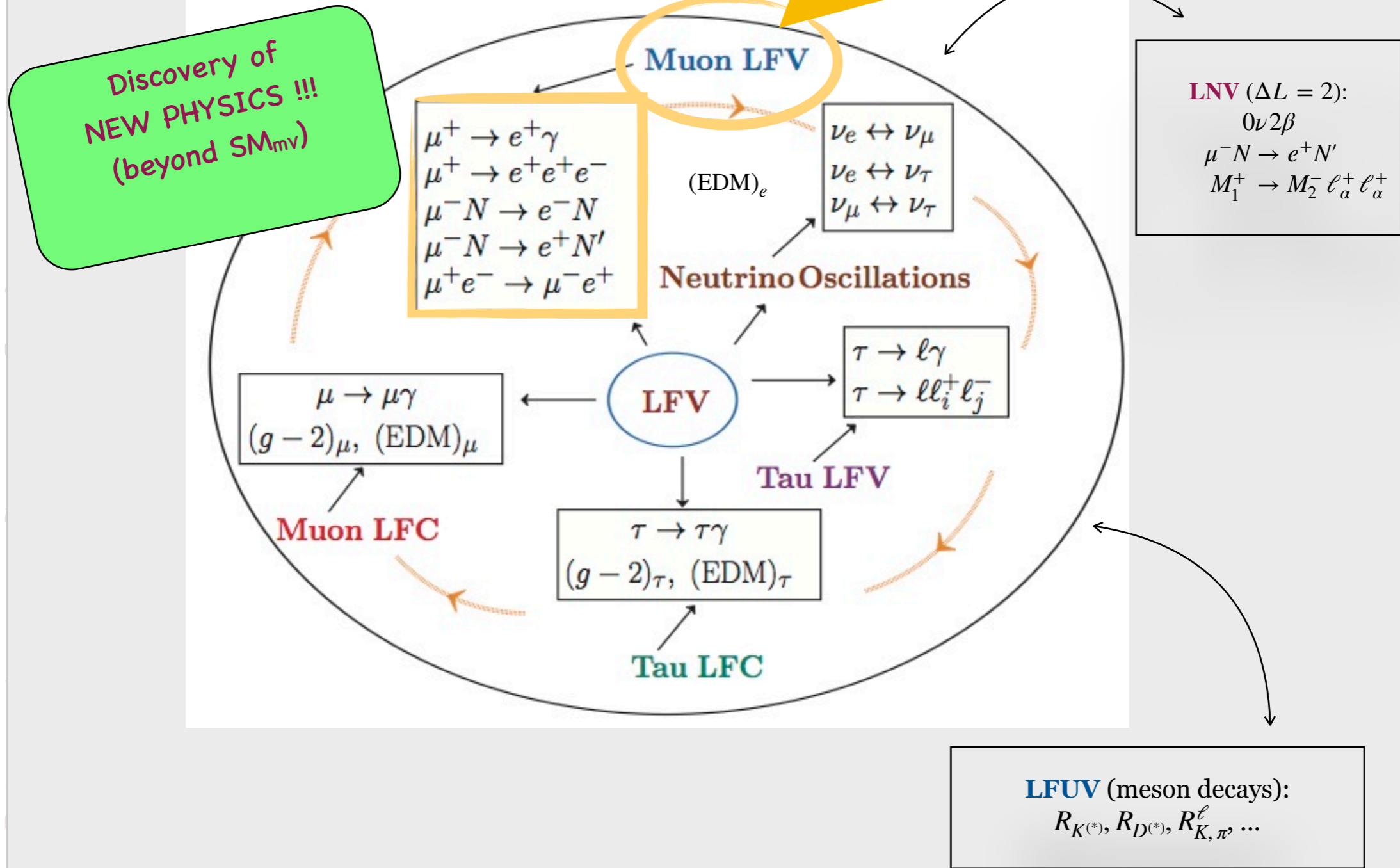
Neutr

E

In SM

l_i

cLF



!!

Muons - ideal **probe for NP**: from lepton flavour universality tests,
to anomalous magnetic moments, ... to **cLFV!**

Muon cLFV - extensive opportunities, numerous observables, relying on **very intense beams**

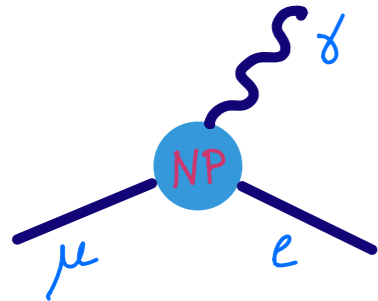
► **Leptonic decays**: radiative $\mu \rightarrow e\gamma$ and three-body $\mu \rightarrow 3e$
muonic atoms $\mu^-(A, Z) \rightarrow e^-(A, Z)$ & LNV $\mu^-(A, Z) \rightarrow e^+(A, Z - 2)^*$
nuclear assisted Coulomb decays $\mu^-e^- \rightarrow e^-e^-$
Muonium oscillations $Mu(\mu^+e^-) - \bar{Mu}(\mu^-e^+)$ and decays $Mu(\mu^+e^-) \rightarrow e^+e^-$
Light "invisible" searches (e.g. $\mu \rightarrow e\phi, \dots$)

► And further! **Semi-leptonic decays**: $M \rightarrow (M')\mu\ell$
And at **colliders**: $Z \rightarrow \mu\tau, H \rightarrow \mu\tau$ (e.g. FCC-ee, CEPC, ...);
high p_T dilepton tails in $pp \rightarrow \mu\ell \dots$
Numerous channels at a **future muon collider!**



Muons: *lightest "unstabiles"* - clean objects, ideal & versatile probes for new physics searches
At the centre of a world-wide comprehensive programme - **experiments and theory**

cLFV muon channels: radiative decays



► cLFV decay: $\mu^+ \rightarrow e^+ \gamma$

► Event signature: $E_e = E_\gamma = m_\mu/2$ (~ 52.8 MeV)

Back-to-back $e^+ - \gamma$ ($\theta \sim 180^\circ$); Time coincidence

► Backgrounds \Rightarrow prompt physics & accidental

Prompt: radiative μ decays ($\mu \rightarrow e \bar{\nu}_e \nu_\mu \gamma$, very low E_ν)

$[\propto R_\mu]$

Accidental: coincidence of γ with positron from Michel decays $\mu \rightarrow e \bar{\nu}_e \nu_\mu$:

photon from $\mu \rightarrow e \bar{\nu}_e \nu_\mu \gamma$; γ from in-flight e^+e^- annihilation

$[\propto R_\mu^2]$

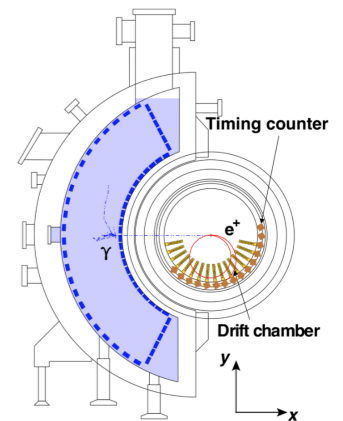
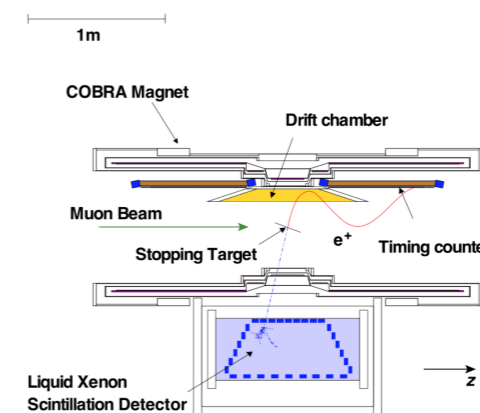
► Experimental status:

First searches (!) in 1940's

[MEG Coll., 1605.05081]

Advent of intense muon beams in 2000's MEG @ PSI

$BR(\mu^+ \rightarrow e^+ \gamma) \leq 4.2 \times 10^{-13}$ (90% CL)



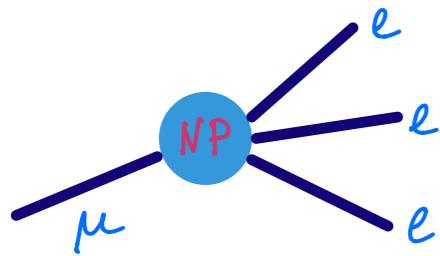
► Future prospects:

[MEG II Coll., 2201.008200]

MEG II (@ PSI): $BR(\mu^+ \rightarrow e^+ \gamma) \leq 6 \times 10^{-14}$

very hard to go beyond 10^{-15} without conceptually different approach

cLFV muon channels: 3-body decays



► cLFV decay: $\mu^+ \rightarrow e^+ e^- e^+$

► Event signature: $\Sigma E_e = m_\mu; \Sigma \vec{P}_e = \vec{0}$

common vertex; Time coincidence

► Backgrounds \Rightarrow physics & accidental

Physics: multi-body μ decays ($\mu \rightarrow e \bar{\nu}_e \nu_\mu e^+ e^-$, very low E_ν)

Accidental: Bhabha scattering of Michel e^+ from $\mu \rightarrow e \bar{\nu}_e \nu_\mu$ decays with atomic $e^+ e^-$

Michel positrons with $e^+ e^-$ from γ conversion

► Experimental status:

SINDRUM @ PSI

[SINDRUM Coll., '88]

$$BR(\mu^+ \rightarrow e^+ e^- e^+) \leq 1.0 \times 10^{-12} \quad (90\% \text{ CL})$$

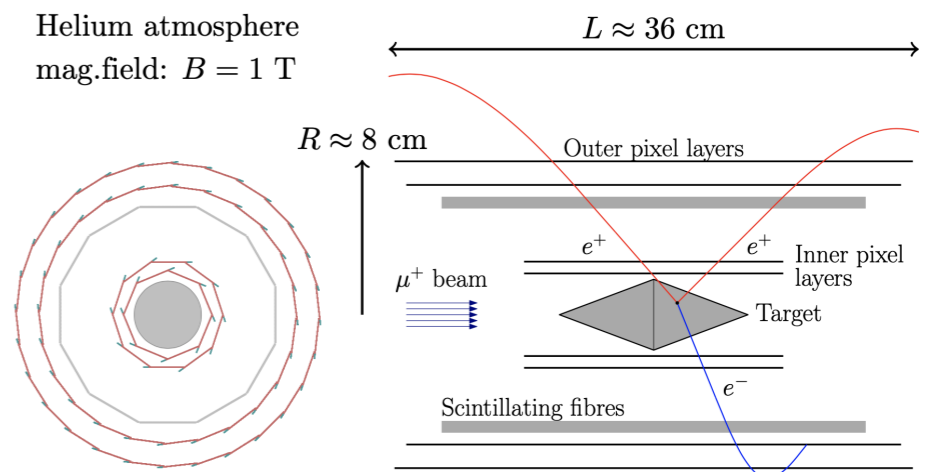
► Future prospects:

[Mu3e Coll., 2009.11690]

Mu3e (@ PSI): expected sensitivity $\mathcal{O}(10^{-15})$ for Phase I

with HIMB, $\mathcal{O}(10^{-16})$ for Phase II

[Aiba et al, 2111.05788]



cLFV in muonic atoms: $\mu - e$ conversion

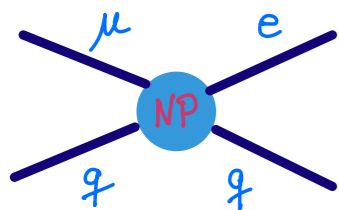
► **Muonic atoms:** 1s bound state formed when μ^- stopped in target

SM allowed processes: decay in orbit (DIO) $\mu^- \rightarrow e^- \nu_\mu \bar{\nu}_e$

nuclear capture $\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1)$

► In the presence of New Physics - **cLFV neutrinoless $\mu^- - e^-$ conversion**

$$\mu^- + (A, Z) \rightarrow e^- + (A, Z)$$



► **Event signature:** single mono-energetic electron

$$E_{\mu e} = m_\mu - E_B(A, Z) - E_R(A, Z)$$

For Aluminium, Lead, Titanium $\sim E_{\mu e} \approx \mathcal{O}(100 \text{ MeV})$

Which target?*** For coherent conversion, maximal rates for $30 \leq Z \leq 60$

► **Backgrounds** \Rightarrow **Only physics!** μ decay in orbit, beam purity, cosmic rays, ...

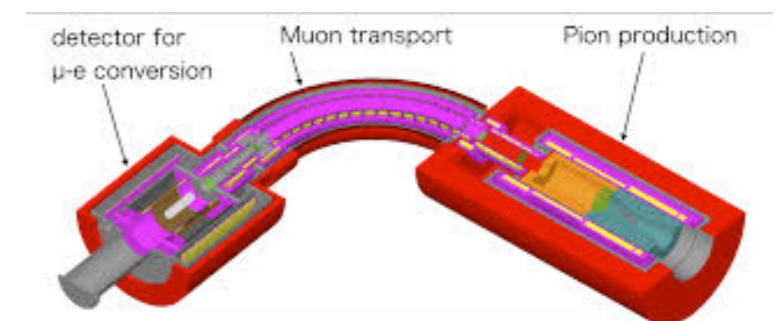
► **Experimental status:** [SINDRUM II Coll., '06]

SINDRUM @ PSI: $CR(\mu^- - e^-, \text{Au}) \leq 7.1 \times 10^{-13}$ (90% CL)

► **Future prospects:**

Mu2e (@ FNAL) - $\mathcal{O}(10^{-17})$, [Bartoszek et al, 1501.05241]

[Abramishvili et al, '20] **COMET (@ JPARC)** - $\mathcal{O}(10^{-15} - 10^{-17})$, ...



COMET Phase-I Layout

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► **Muonic atoms:** 1s bound state formed when μ^- stopped in target

SM allowed processes: decay in orbit (DIO) $\mu^- \rightarrow e^- \nu_\mu \bar{\nu}_e$

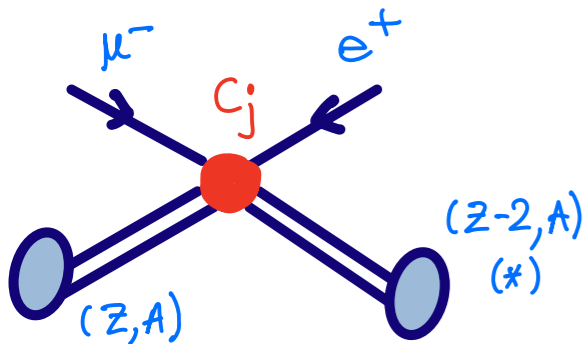
nuclear capture $\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1)$

► In the presence of New Physics - **cLFV & LNV** ($\Delta L = 2$) **neutrinoless $\mu^- - e^+$ conversion**

$$\mu^- + (A, Z) \rightarrow e^+ + (A, Z - 2)^*$$

$\mu^- - e^-$ conversion: coherent process, single nucleon, nuclear ground states

$\mu^- - e^+$ conversion: 2 nucleons ($\Delta Q = 2$), possibly excited final state



► **Event signature:** single positron - but complex energy spectrum

$$E_{\mu e}^{N^*} = m_\mu - E_B(A, Z) - E_R(A, Z) - \Delta_{Z-2}^{(*)}$$

For Aluminium (giant dipole resonance) $\sim E_{\mu-e^+}^{\text{Al, GDR}} \approx \mathcal{O}(83.9 \text{ MeV})$

► **Experimental status:**

Collaboration	year	Process	Bound
PSI/SINDRUM	1998	$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}^*$	3.6×10^{-11}
PSI/SINDRUM	1998	$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}$	1.7×10^{-12}

► **Future prospects:**

Best sensitivity for Ca, S and Ti targets (possibly $\sim \mathcal{O}(\text{few} \times 10^{-15})$); Al@Mu2e?

cLFV in muonic atoms: $\mu - e$ conversion

► **Muonic atoms:** 1s bound state formed when μ^- stopped in target

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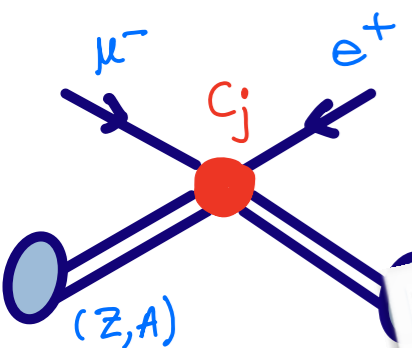
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► **Event signature**

A unique connection between LNV (in association with Majorana nature and possibly, neutrino mass generation) and cLFV

$$R^{(1, Z)} - \Delta_{Z-2}^{(*)}$$

giant dipole resonance) $\sim E_{\mu^-e^+}^{Al, GDR} \approx \mathcal{O}(83.9 \text{ MeV})$

► **Experimental status:**

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► **Future prospects:**

Best sensitivity for Ca, S and Ti targets (possibly $\sim \mathcal{O}(\text{few} \times 10^{-15})$); Al@Mu2e?

► Muonium: $\mu^+ e^-$

Hydrogen-like Coulomb bound state, free of hadronic interactions!
Powerful laboratory for EW tests and cLFV

► In the presence of New Physics - Muonium oscillations and Muonium decays

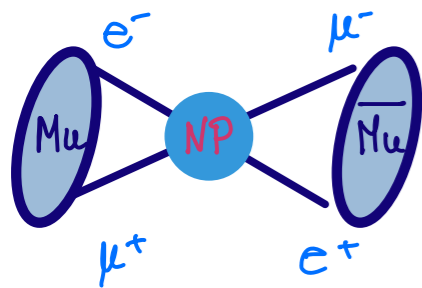
► Mu- $\bar{\text{Mu}}$ oscillation

Spontaneous conversion $\mu^+ e^- \leftrightarrow \mu^- e^+$

Reflects a double (individual) lepton number violation $|\Delta L_e| = |\Delta L_\mu| = 2$

Rate (typically) suppressed by external magnetic fields

Detection: reconstruct Michel electron from μ^- decays and shell positron



Experimental status: MACS - $P(\text{Mu} - \bar{\text{Mu}}) < 8.3 \times 10^{-11}$ [Willmann et al, 1999]

Future prospects: MACE, AMF (@FNAL)

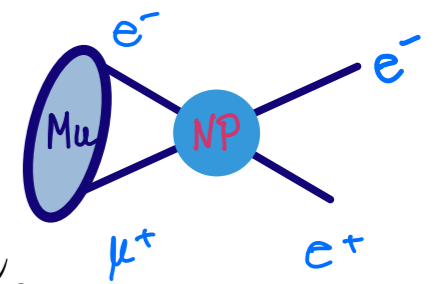
[Bai et al, 2203.11406]

► Mu decays

$$\mu^+ e^- \rightarrow e^+ e^-$$

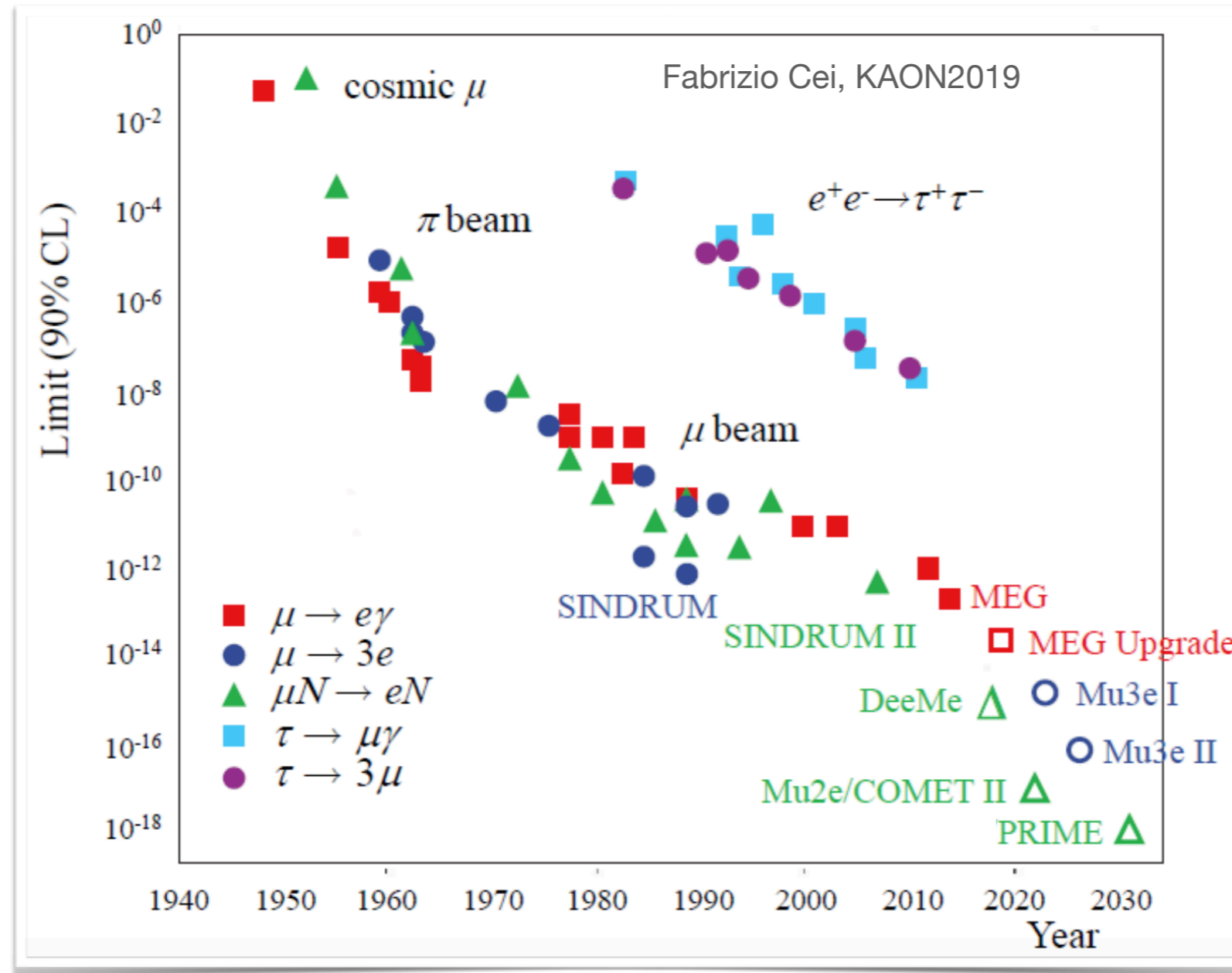
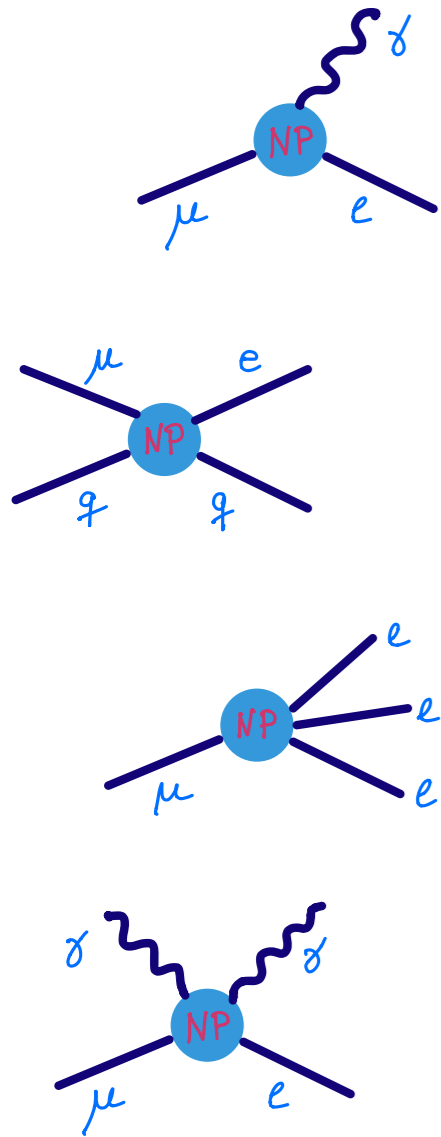
Clear signal compared to SM-allowed muonium decay, $\text{Mu} \rightarrow e^+ e^- \bar{\nu}_\mu \nu_e$

No available bounds, no clear roadmap...



cLFV muon observables: experimental status

Searching for tiny cLFV effects \Rightarrow high-intensity sources for excellent sensitivities



\Rightarrow Need many many (really many!) muons: excellent sensitivity with current sources, **Amazing prospects** with advent of high-intensity beams (PSI, FNAL, JPARC) and beyond?... *Muon facility? Muon collider?*

Lepton flavours: from ν oscillations... EFT!

► Generic New Physics observables in the **lepton sector**:

- Lepton number violation (e.g. neutrino masses, $0\nu 2\beta$ decays, ...)
- Electric and (anomalous) magnetic moments - d_ℓ , $(g - 2)_\ell$
- **charged lepton flavour violation**

Here - a tiny tip of the iceberg!

Back to \mathcal{L}^{eff} : cast observables in terms of \mathcal{C}_{ij} and Λ_{NP}

$$\mathcal{L}^{\text{eff}} = \mathcal{L}^{\text{SM}} + \frac{\mathcal{C}_5 \mathcal{O}^5}{\Lambda_{\text{LNV}}} (m_\nu) + \frac{\mathcal{C}_6 \mathcal{O}^6}{\Lambda_{\text{CLFV}}^2} (\ell_\alpha \leftrightarrow \ell_\beta) + \dots + \frac{\mathcal{C}_7 \mathcal{O}^7}{\Lambda_{\text{LNV}}^3} (0\nu 2\beta) + \dots$$

Majorana ν masses

Kinetic corrections, ...
EW precision, top physics, ...
Electric dipole & anomalous magnetic moments, ...
cLFV (dipole, 3 body, matter assisted, ...)

Lepton number violation, cLFV & **LNV**, ...

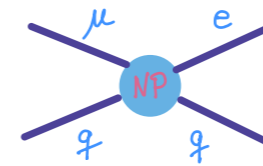
Deceptively simple task... different new physics scales, numerous operators!

For **cLFV**, technically very involved, even if no "SM background"...

Muon cLFV: EFT approach to New Physics

► QED & QCD & NP effective Lagrangian, many involved operators!

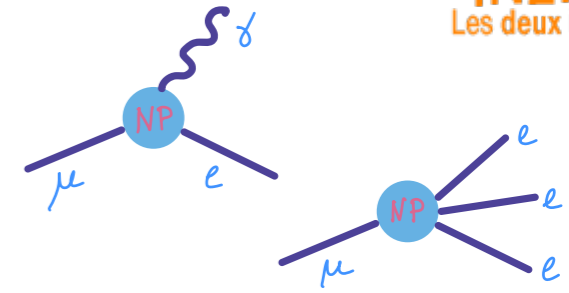
$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{QED}} + \mathcal{L}_{\text{QCD}} + \frac{1}{\Lambda^2} \left\{ C_L^D O_L^D + \sum_{f=q,\ell} (C_{ff}^{V LL} O_{ff}^{V LL} + C_{ff}^{V LR} O_{ff}^{V LR} + C_{ff}^{S LL} O_{ff}^{S LL}) + \sum_{h=q,\tau} (C_{hh}^{T LL} O_{hh}^{T LL} + C_{hh}^{S LR} O_{hh}^{S LR}) + C_{gg}^L O_{gg}^L + L \leftrightarrow R \right\} + \text{h.c.},$$



$$O_L^D = e m_\mu (\bar{e} \sigma^{\mu\nu} P_L \mu) F_{\mu\nu},$$

$$O_{ff}^{V LL} = (\bar{e} \gamma^\mu P_L \mu) (\bar{f} \gamma_\mu P_L f),$$

$$O_{ff}^{V LR} = (\bar{e} \gamma^\mu P_L \mu) (\bar{f} \gamma_\mu P_R f),$$



$$O_{ff}^{S LL} = (\bar{e} P_L \mu) (\bar{f} P_L f),$$

$$O_{hh}^{S LR} = (\bar{e} P_L \mu) (\bar{h} P_R h),$$

$$O_{hh}^{T LL} = (\bar{e} \sigma_{\mu\nu} P_L \mu) (\bar{h} \sigma^{\mu\nu} P_L h),$$

$$O_{gg}^L = \alpha_s m_\mu G_F (\bar{e} P_L \mu) G_{\mu\nu}^a G_a^{\mu\nu}$$

... and further "mixing" effects, from RGE running (including loop effects) ...

► Simple examples: at leading order one has

$$\text{BR}(\mu \rightarrow e\gamma) \simeq 384\pi^2 \frac{v^4}{\Lambda^4} \left(|C_{D,L}|^2 + |C_{D,R}|^2 \right)$$

$$\text{BR}(\mu \rightarrow eee) \simeq \frac{v^4}{\Lambda^4} \left[\frac{1}{8} |C_{S,LL}|^2 + 2 |C_{V,RR} + 4eC_{D,L}|^2 + (64 \ln \frac{m_\mu}{m_e} - 136) e |C_{D,L}|^2 + |C_{V,RL} + 4eC_{D,L}|^2 \right] + (L \leftrightarrow R)$$

CR($\mu - e, N$): far more involved (nuclear target effects, spin (in)-dependent contributions, ...)

$$\approx \frac{1}{\Gamma_{\text{cap}}} \frac{m_\mu^5}{\Lambda^4} \left[\left| e C_L^D D_N + 4 \left(G_F m_\mu m_p \tilde{C}_{(p)}^{SL} S_N^{(p)} + \tilde{C}_{(p)}^{VR} V_N^{(p)} + p \rightarrow n \right) \right|^2 + (L \leftrightarrow R) \right]$$

$D_N, S_N^{(p/n)}, V_N^{(p/n)}$: nuclear "overlap integrals" between lepton wave functions and nucleon densities (*target-dependent*)

Results of a recent EFT approach to muon transitions:

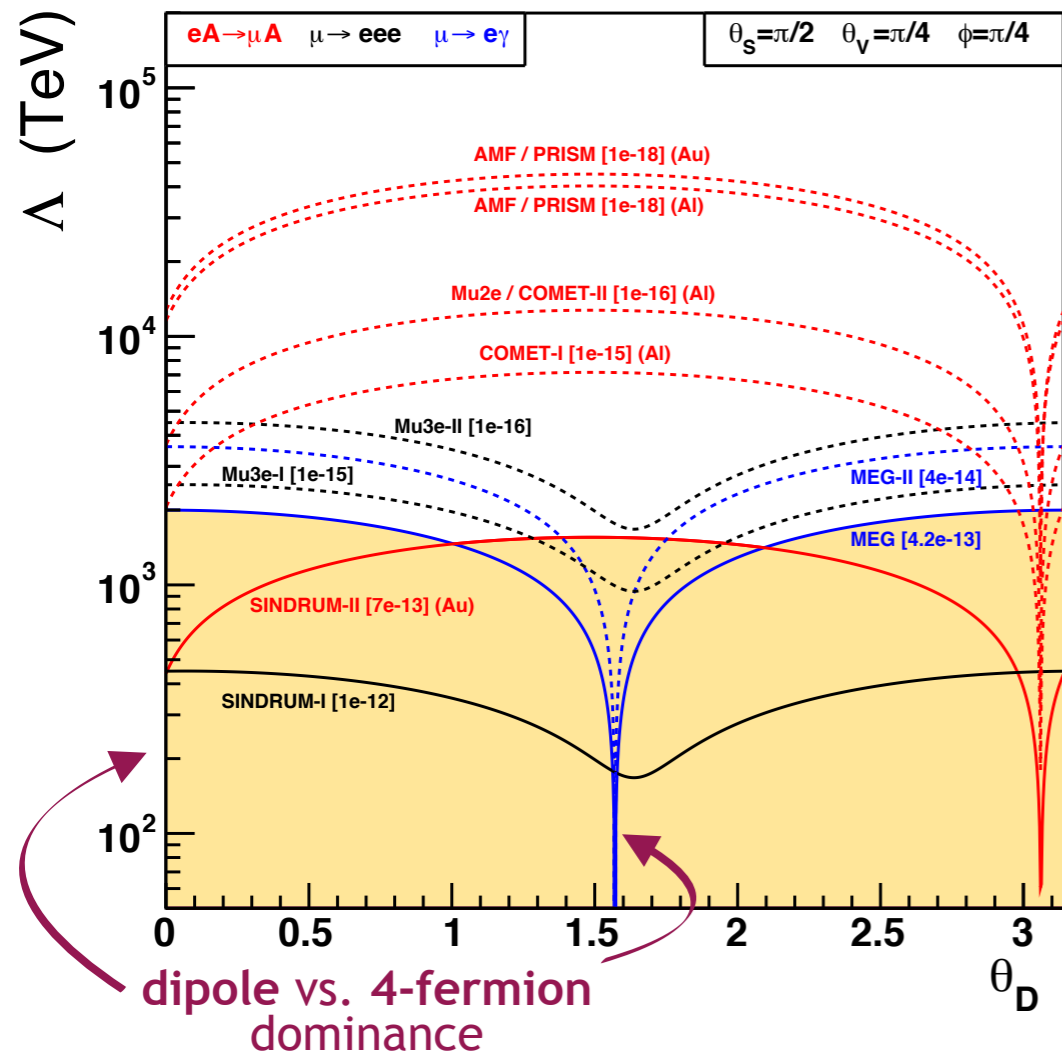
$$\mathcal{L}_{\text{NP, cLFV}}^{\text{eff}} = \frac{1}{\Lambda^2} \left[C_D (\bar{e} \sigma^{\nu\rho} P_R \mu) F_{\nu\rho} + C_S (\bar{e} P_R \mu) (\bar{e} P_R e) + C_{VR} (\bar{e} \gamma^\nu P_L \mu) (\bar{e} \gamma_\nu P_R e) + C_{VL} (\bar{e} \gamma^\nu P_L \mu) (\bar{e} \gamma_\nu P_L e) + C_{\text{N-light}} \mathcal{O}_{\text{N-light}} + C_{\text{N-heavy}\perp} \mathcal{O}_{\text{N-heavy}\perp} \right]$$

$$\vec{C} = \{C_D, C_S, C_{VR}, C_{VR}, C_{VL}, C_{\text{N-light}}, C_{\text{N-heavy}\perp}\}$$

$$\Rightarrow \text{BR}(\mu \rightarrow e\gamma) \simeq 384\pi^2 \frac{v^4}{\Lambda^4} |\vec{C} \cdot \hat{e}_{DR}|^2 \rightsquigarrow \leq \text{BR}^{\text{exp}}(\text{future})$$

and likewise for other observables

$\text{BR}(\mu \rightarrow 3e)$, $\text{CR}(\mu - e, N)$, Muonium oscillations...



$$2\sqrt{2} C_D \approx \frac{\cos \theta_D}{\Lambda^2}$$

Sensitivity to NP scales (current & future):

MEG ($\mu \rightarrow e\gamma$) $\leftrightarrow \Lambda_{\text{cLFV}} \sim \mathcal{O}(10^3 \text{ TeV})$ [dipole]

SINDRUM II ($\mu - e$, Au) $\leftrightarrow \Lambda_{\text{cLFV}} \sim \mathcal{O}(10^3 \text{ TeV})$ [4f]

Mu2e/COMET II ($\mu - e$, Al) $\leftrightarrow \Lambda_{\text{cLFV}} \lesssim \mathcal{O}(10^4 \text{ TeV})$ [either dipole or 4f]

⇒ **cLFV** data to constrain \mathcal{O}^6 (and infer sensitivity of a process to operator \mathcal{O}^6)

► Fully exploring the potential of atomic (elastic) **muon-electron conversion**, $CR(\mu - e, N)$:

Comparatively more involved theoretical approach!

Explore target-nucleus dependence to distinguish **dominant operator** (hint on NP model!)

[extensive contributions since Kitano et al, 0203110! see Davidson et al, 1810.01884; Heeck et al, 2203.00702, ...]

In the advent of an observation (@ Mu2e, COMET \leadsto using Aluminium targets)

prepare **choice of future targets**

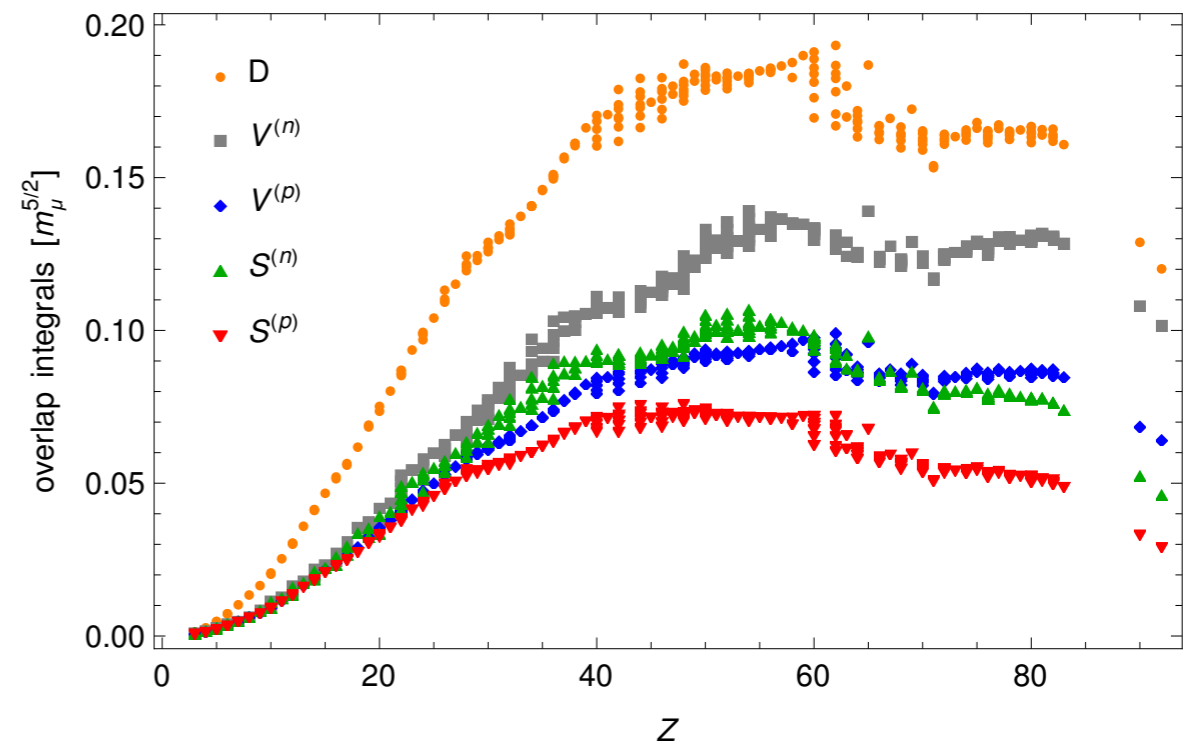
Which offer the **largest complementarity** with respect to Al?

$$BR_{SI}(\mu A \rightarrow eA) = \frac{32G_F^2}{\Gamma_{\text{capture}}} \left[\left| C_{V,R}^{pp} V^{(p)} + C_{S,L}^{pp'} S^{(p)} + C_{V,R}^{nn} V^{(n)} + C_{S,L}^{nn'} S^{(n)} + C_{D,L} \frac{D}{4} \right|^2 + \{L \leftrightarrow R\} \right].$$

Overlap integrals:

more distinguishable at **large Z** !

Better disentangle dominant NP contributions...



[Heeck et al, 2203.00702]

⇒ cLFV data to constrain \mathcal{O}^6 (and infer sensitivity of a process to operator \mathcal{O}^6)

► Fully exploring the potential of atomic (elastic) muon-electron conversion, $CR(\mu - e, N)$:

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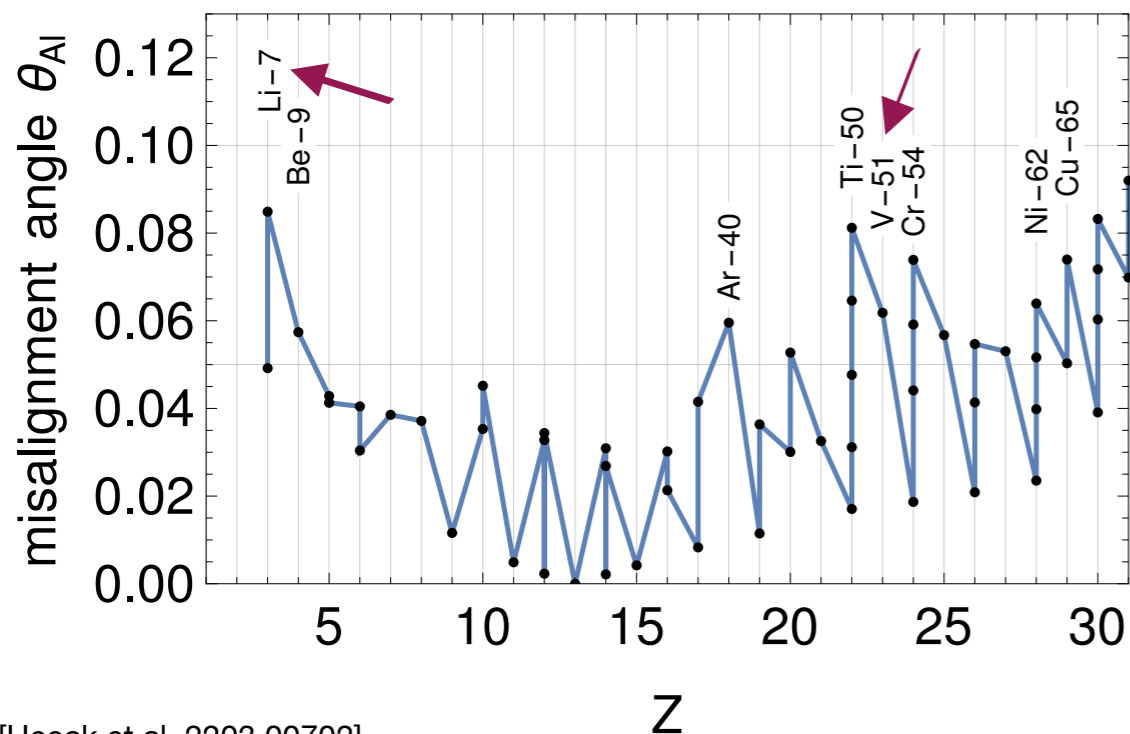
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prepare choice of future targets

Which offer the largest complementarity with respect to Al? θ_{Al}



[Heeck et al, 2203.00702]

- Heavier nuclei (Au, Pb)! ... not feasible... (pulsed beams)
- Among experimental-friendly $Z \leq 25$ targets several (theoretically good) candidates
Li-7, Ti-50, Ti-49, Cr-54, .., V-51

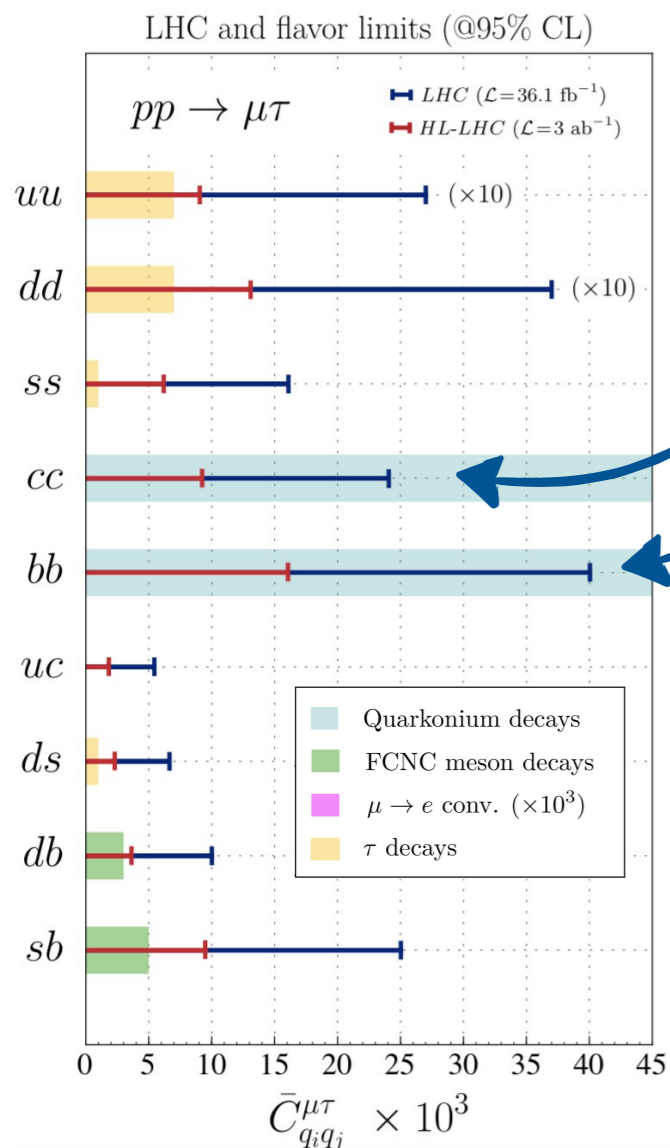
⇒ Li-7 and/or V-51 : preferable "second" targets post $CR(\mu - e, Al)$ observation

$\mu - e$ conversion: "unbeatable" NP probe

▶ Albeit leading to formally different transitions, the same leptonic and semi-leptonic operators can be at the origin of **flavour violating transitions** in very distinct contexts

LHC \rightsquigarrow abundant sources of flavour in pp collisions

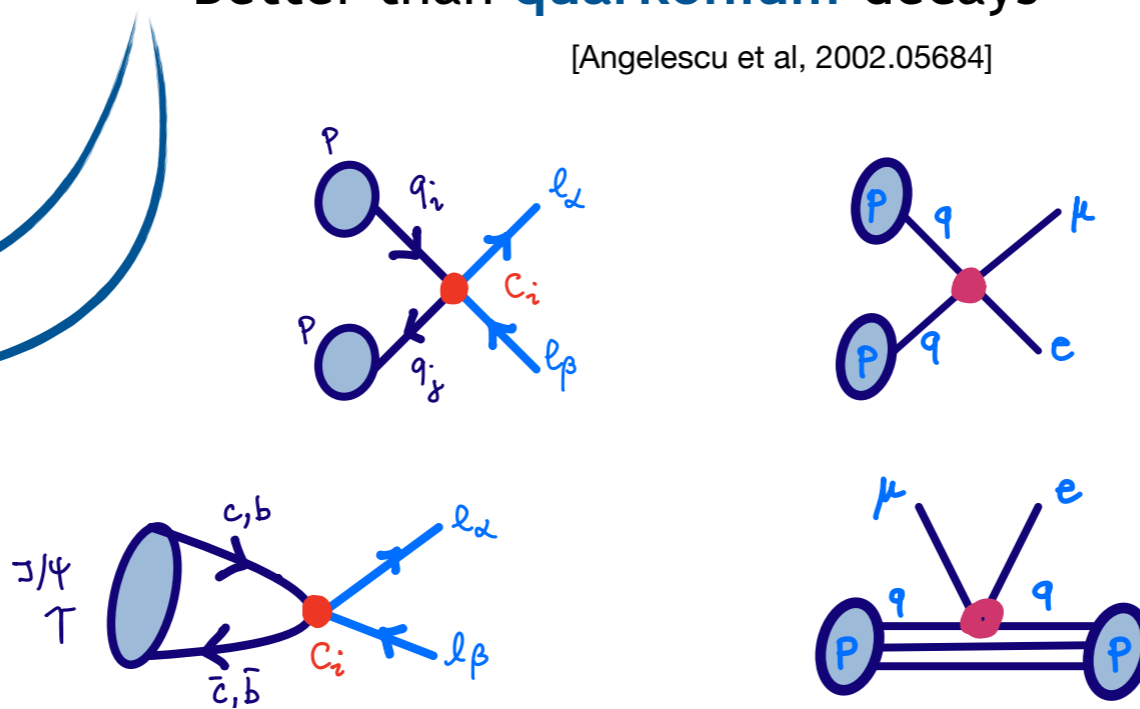
Drell-Yan $q_i \bar{q}_j \rightarrow \ell_\alpha \ell_\beta (\ell_\alpha \nu_\beta)$ probe similar operators, but at **high p_T**



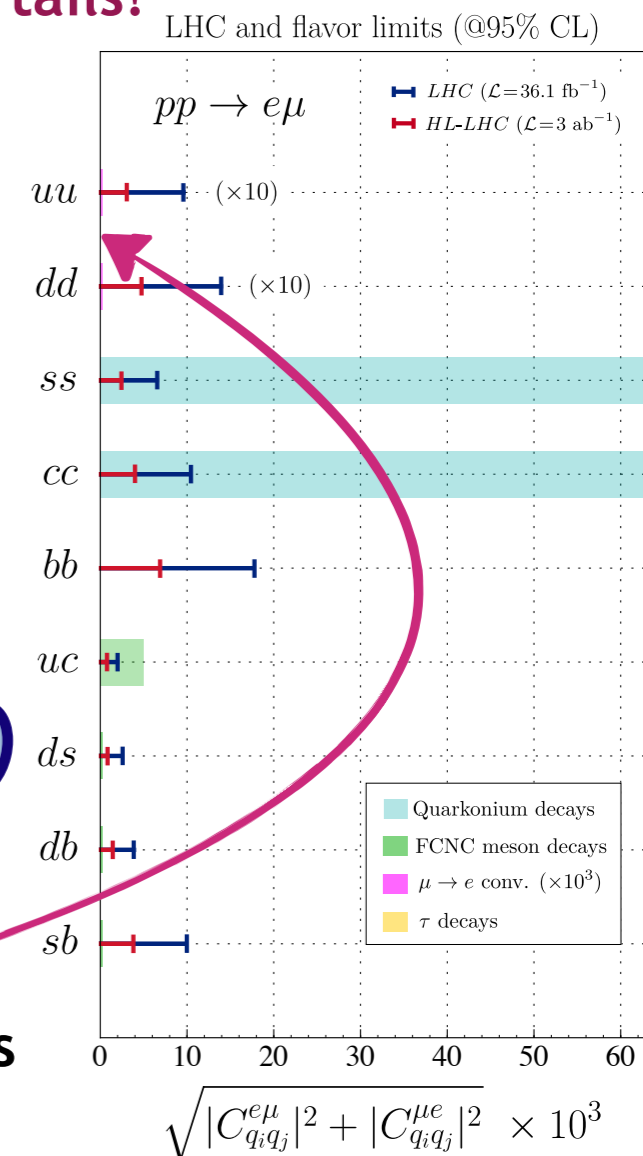
\Rightarrow **LHC limits very competitive for $\mu\tau$ tails!**

Better than **quarkonium-decays**

[Angelescu et al, 2002.05684]



\Rightarrow **Impossible for μe tails to out-perform cLFV searches ($\mu - e, N$) conversion bounds**



[Constraints from quarkonia decays, see also Calibbi et al, 2207.10913]

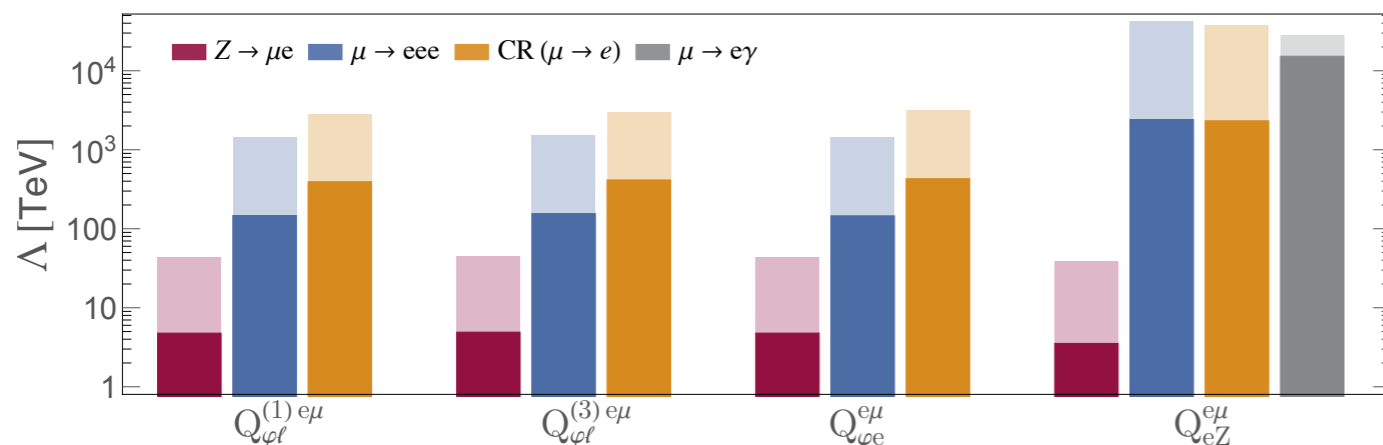
$\mu - e$ conversion: "unbeatable" NP probe

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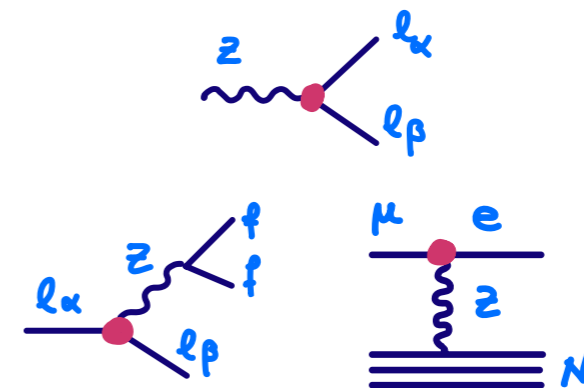
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TeraZ factory (FCC-ee, CEPC) \rightsquigarrow **EW precision & flavour violation**



[Calibbi et al, 2107.10273]



TeraZ factory \rightsquigarrow **cLFV Z decays**

For $Z \rightarrow \mu e$ much better sensitivity of **dedicated (low-energy) cLFV searches**

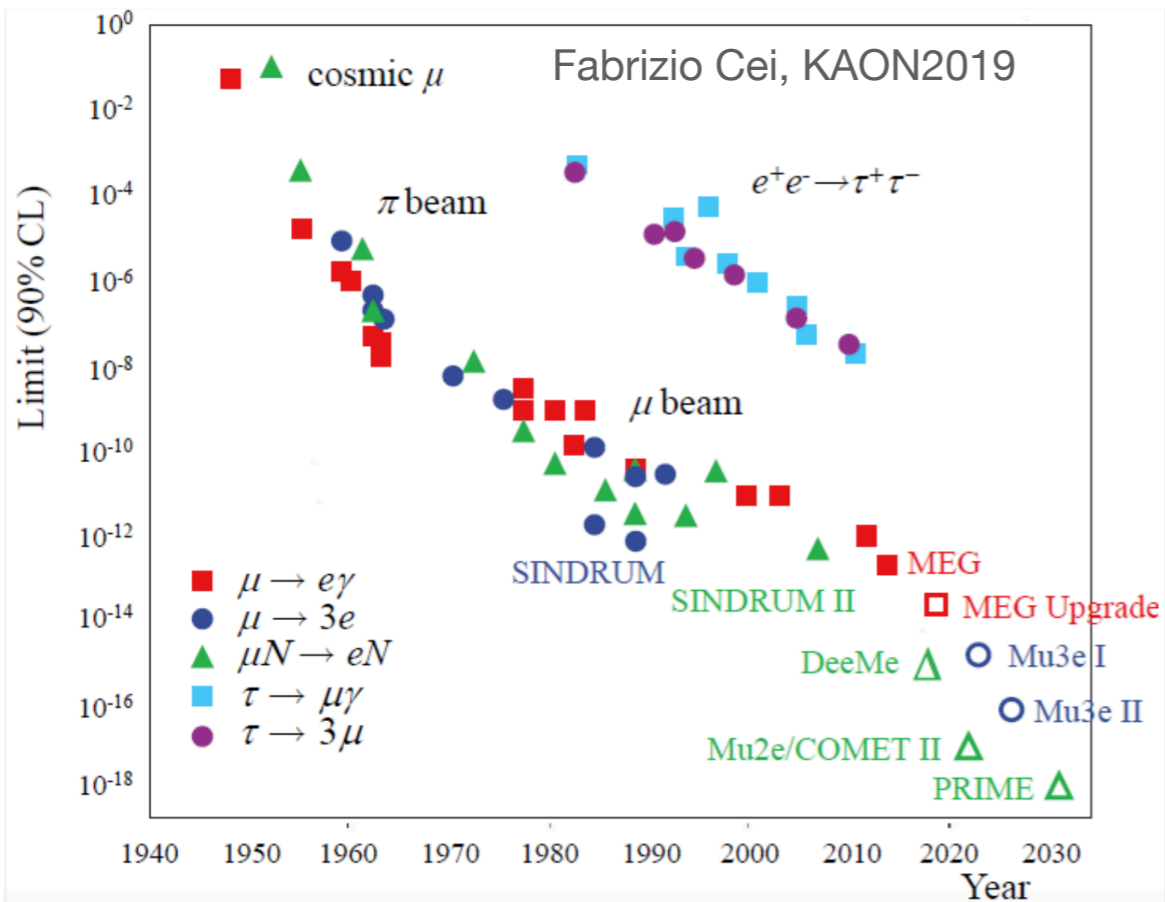
$\mu \rightarrow eee$, $\mu - e$ conversion

Promising potential of TeraZ factory for

$Z \rightarrow \tau \ell$ decays

(competitive with low-energy cLFV)

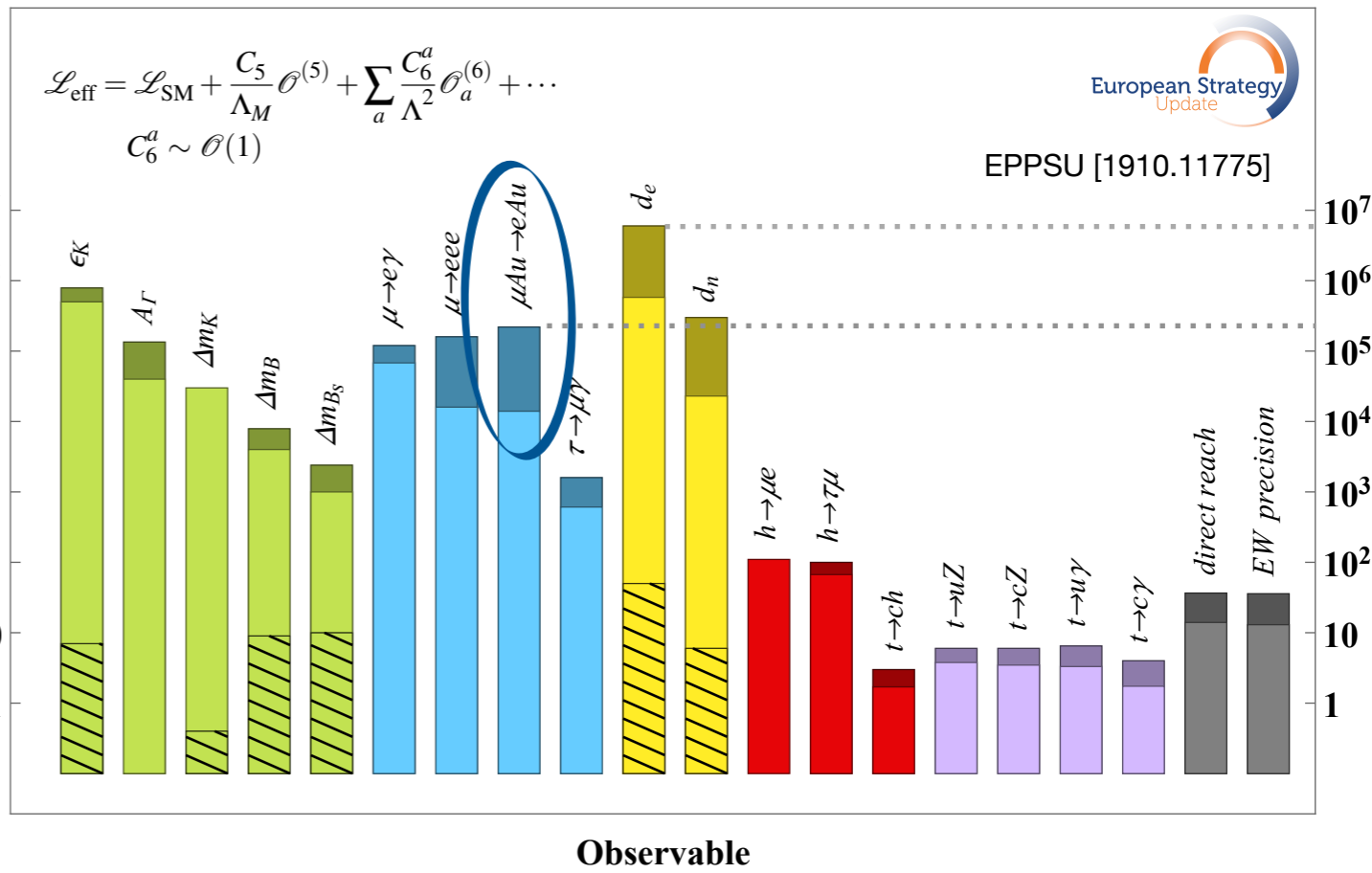
Muon cLFV prospects



Muons: *lightest "unstabiles"* - clean objects, ideal & versatile probes for NP searches
At the centre of a world-wide comprehensive programme - *experiments and theory*



cLFV $\mu - e$ conversion in nuclei:
most sensitive **flavoured probe** of NP!
Unprecedented **sensitivity $\mathcal{O}(10^{-17})$**
Unique opportunities for NP discovery



$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{C_5}{\Lambda_M} \mathcal{O}^{(5)} + \sum_a \frac{C_6^a}{\Lambda^2} \mathcal{O}_a^{(6)} + \dots$$

$C_6^a \sim \mathcal{O}(1)$

EPPSU [1910.11775]

Electric dipole moments (EDM)

$$i\langle p' | J^\mu(0) | p \rangle = (-ie) \bar{\Psi}(p') \left[\gamma^\mu F_1(k^2) + \frac{i\sigma^{\mu\nu} k_\nu}{2m} F_2(k^2) \right. \\ \left. + \gamma_5 \frac{i\sigma^{\mu\nu} k_\nu}{2m} F_3(k^2) + \gamma_5 (k^2 \gamma^\mu - \not{k}^\mu) F_4(k^2) \right] \Psi(p)$$

$F_1(0) = 1$ (charge renormalisation)
 $\mu_e = \frac{e}{2m} (F_1(0) + F_2(0))$ (magnetic dipole moment)
 $a_e = F_2(0)$ (anomalous magnetic moment)
 $d_e = -\frac{e}{2m} F_3(0)$ (electric dipole moment, $T\&P$ violating)
 $F_4(0) = 0$ (anapole moment, P violating)



Electric dipole moments (EDM)



Flavour & CPV: the "usual graveyard of BSM electroweak theories"

Neutron EDM: observable (likely) responsible for falsifying the largest number of BSM...



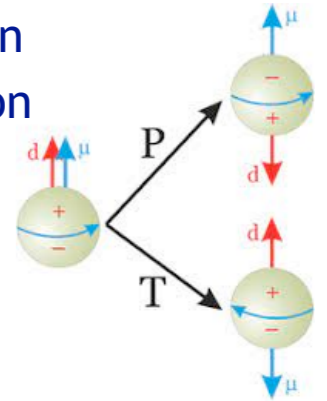
Electric dipole moments (EDM)

Electric Dipole Moments - observables sensitive to **CP violation**

non-relativistic approach: $\mathcal{H} \propto -(\mu_f \vec{\sigma} \cdot \vec{B} + d_f \vec{\sigma} \cdot \vec{E})$

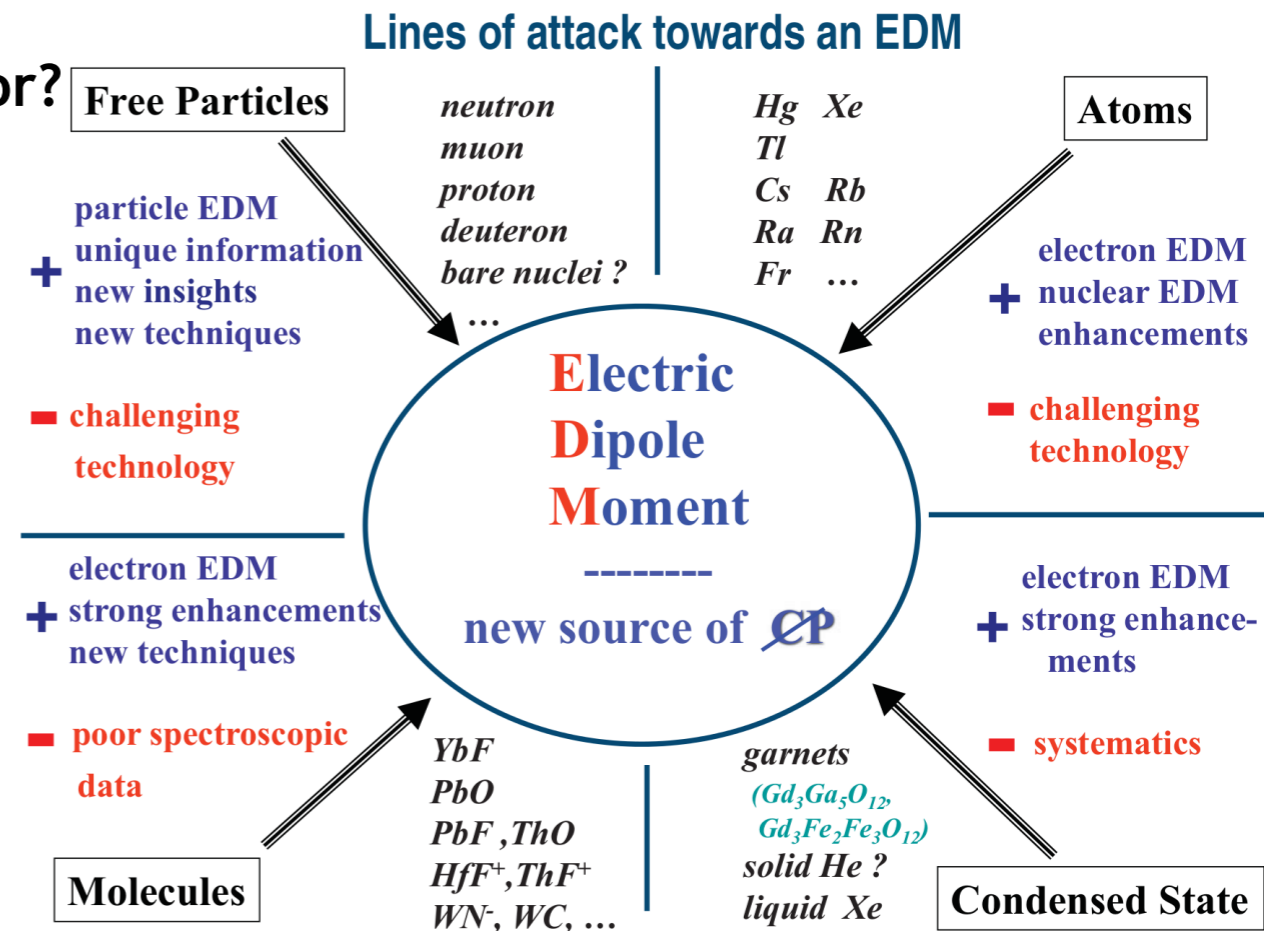
relativistic generalisation $\leadsto \mathcal{L}_{\text{CP-odd}} = -\frac{i}{2} \mathbf{d} \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$

P & T violation
 \Rightarrow CP violation



EDMs: sensitive to **SM sources of CP violation** (weak δ_{CKM} and strong $\bar{\theta}$),
and **NP CPV interactions** - required to **explain the BAU** (baryo- or leptogenesis)
(especially "flavour blind" new phases)

► Which **EDM observables** are being "hunted" for?
As many as possible!



[adapted from Jungmann, 2013]

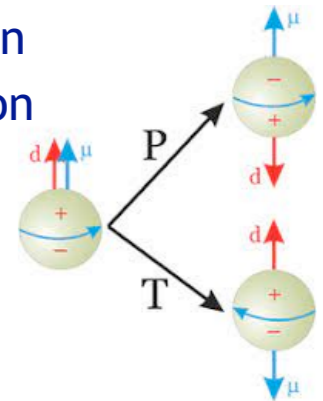
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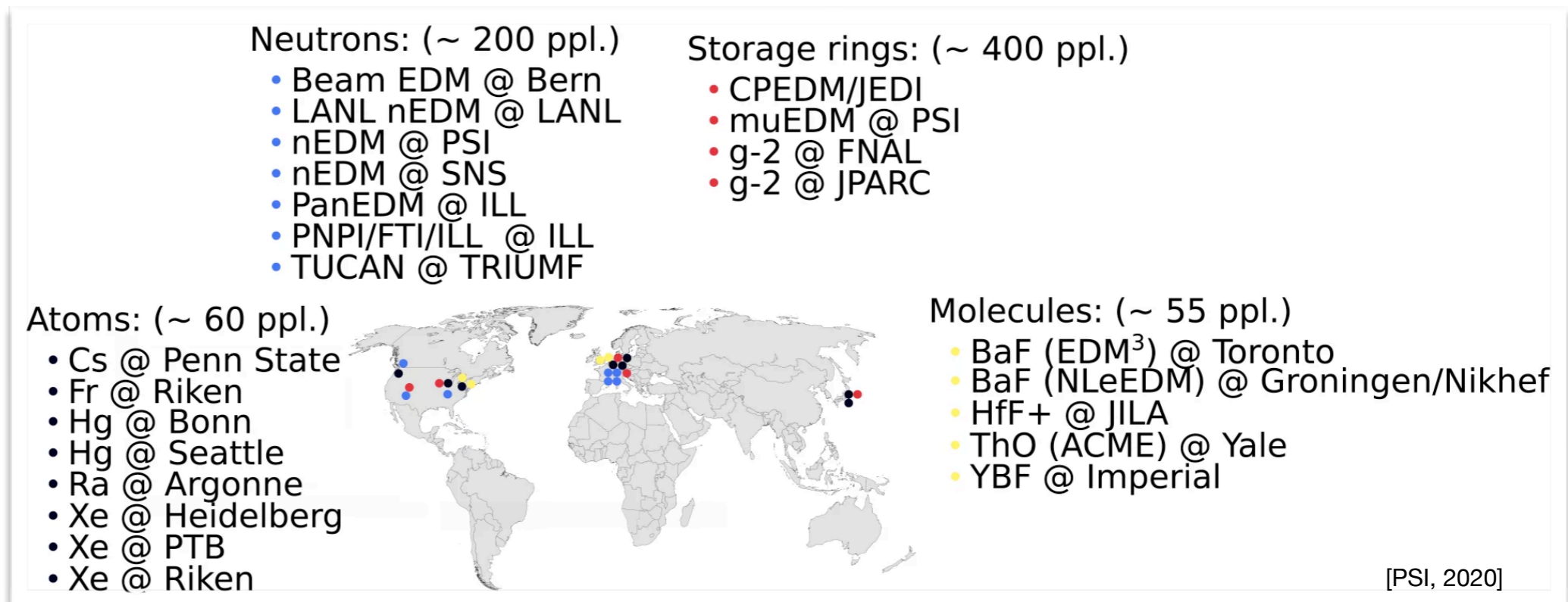


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► Where?

Worldwide!



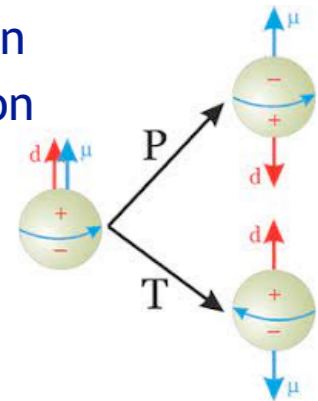
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- ▶ Which **EDM observables** are being "hunted" for? *As many as possible!*
- ▶ Where? *Worldwide!*
- ▶ Bounds obtained? *Impressive!*

	Result	95% u.l.
Paramagnetic systems		
Xe ^m	$d_A = (0.7 \pm 1.4) \times 10^{-22}$	$3.1 \times 10^{-22} \text{ e cm}$
Cs	$d_A = (-1.8 \pm 6.9) \times 10^{-24}$	$1.4 \times 10^{-23} \text{ e cm}$
	$d_e = (-1.5 \pm 5.7) \times 10^{-26}$	$1.2 \times 10^{-25} \text{ e cm}$
	$C_S = (2.5 \pm 9.8) \times 10^{-6}$	2×10^{-5}
	$Q_m = (3 \pm 13) \times 10^{-8}$	$2.6 \times 10^{-7} \mu_N R_{Cs}$
Tl	$d_A = (-4.0 \pm 4.3) \times 10^{-25}$	$1.1 \times 10^{-24} \text{ e cm}$
	$d_e = (6.9 \pm 7.4) \times 10^{-28}$	$1.9 \times 10^{-27} \text{ e cm}$
YbF	$d_e = (-2.4 \pm 5.9) \times 10^{-28}$	$1.2 \times 10^{-27} \text{ e cm}$
ThO	$d_e = (-2.1 \pm 4.5) \times 10^{-29}$	$9.7 \times 10^{-29} \text{ e cm}$
	$C_S = (-1.3 \pm 3.0) \times 10^{-9}$	6.4×10^{-9}
HfF ⁺	$d_e = (0.9 \pm 7.9) \times 10^{-29}$	$1.6 \times 10^{-28} \text{ e cm}$

	Result	95% u.l.
Diamagnetic systems		
¹⁹⁹ Hg	$d_A = (2.2 \pm 3.1) \times 10^{-30}$	$7.4 \times 10^{-30} \text{ e cm}$
¹²⁹ Xe	$d_A = (0.7 \pm 3.3) \times 10^{-27}$	$6.6 \times 10^{-27} \text{ e cm}$
²²⁵ Ra	$d_A = (4 \pm 6) \times 10^{-24}$	$1.4 \times 10^{-23} \text{ e cm}$
TlF	$d = (-1.7 \pm 2.9) \times 10^{-23}$	$6.5 \times 10^{-23} \text{ e cm}$

ACME (ThO)
 $|d_e| < 1.1 \times 10^{-29} \text{ e cm}$

$d_n = (0.0 \pm 1.2) \times 10^{-26} \text{ e cm}$
 $|d_n| < 1.8 \times 10^{-26} \text{ e cm}$

	Result	95% u.l.
Particle systems		
μ	$d_\mu = (0.0 \pm 0.9) \times 10^{-19}$	$1.8 \times 10^{-19} \text{ e cm}$
τ	$Re(d_\tau) = (1.15 \pm 1.70) \times 10^{-17}$	$3.9 \times 10^{-17} \text{ e cm}$
Λ	$d_\Lambda = (-3.0 \pm 7.4) \times 10^{-17}$	$1.6 \times 10^{-16} \text{ e cm}$

[nEDM Coll - Abel et al, 2001.11966]

[Chupp et al, 1710.02504]

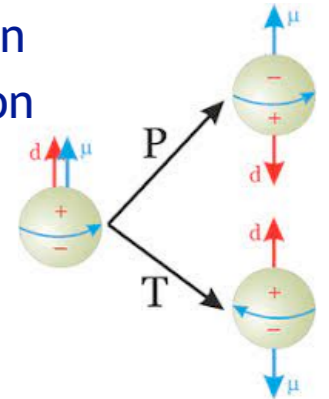
Electric dipole moments (EDM)

Electric Dipole Moments - observables sensitive to **CP violation**

non-relativistic approach: $\mathcal{H} \propto -(\mu_f \vec{\sigma} \cdot \vec{B} + d_f \vec{\sigma} \cdot \vec{E})$

relativistic generalisation $\rightsquigarrow \mathcal{L}_{\text{CP-odd}} = -\frac{i}{2} \mathbf{d} \bar{\Psi} \sigma^{\mu\nu} \gamma_5 \Psi F_{\mu\nu}$

P & T violation
 \Rightarrow CP violation



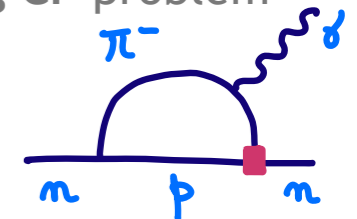
Sources of **CP violation** in the Standard Model

Strong interactions: $\mathcal{L}_{\text{CP}}^{\text{QCD}} = \theta \frac{g^2}{32\pi^2} G^{\mu\nu} \tilde{G}_{\mu\nu} - i \bar{q} \text{Im}(M_q) \gamma_5 q$ $\bar{\theta} = \theta + \text{Arg}[\det(M_q)]$

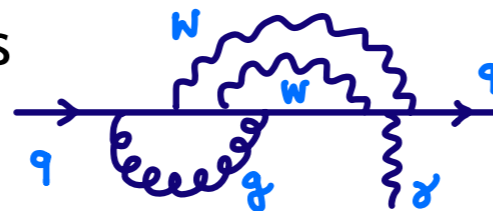
$\bar{\theta} < 10^{-10} \rightsquigarrow$ Strong CP problem

Electroweak CPV: $Y^f \rightsquigarrow \delta_{\text{CKM}}$

$$(J_{\text{CP}} = \mathcal{J}[V_{ts}^* V_{td} V_{ud}^* V_{us}] \approx 3 \times 10^{-5})$$

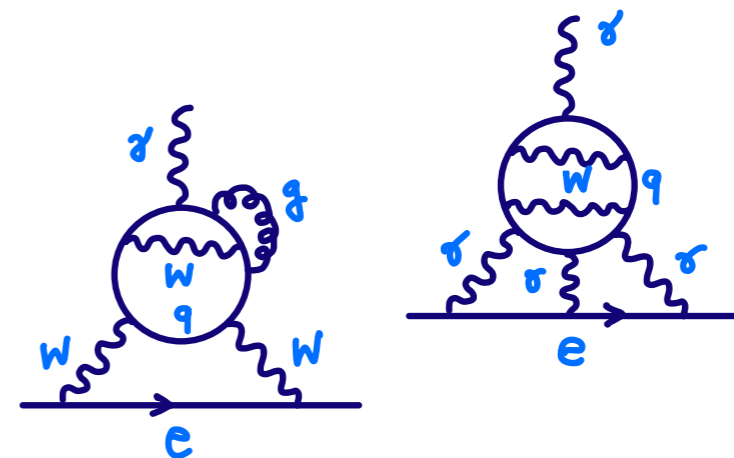


"quark" EDMs @ 3 loops



lepton EDMs @ 4 loops

(no leptonic sources of CPV in the SM...)



\rightsquigarrow tiny theoretical predictions ($d_e, d_N, d_{\text{Hg}}, \dots$)

EDMs: CPV (SM and beyond)

- A *non-trivial theory problem*: numerous scales and approaches (elementary, QCD, nuclear & atomic physics, and effective description...)
- **Paramagnetic & diamagnetic observables**

SM pioneering results for EDMs (J_{CP}):

$$d_N \propto C_{qq}(J_{CP}) \propto J_{CP} G_F^2 \sim \mathcal{O}(10^{-32})$$

[Khriplovich, Zhitnitsky '82; McKellar et al '87; Mannel, Uraltsev '12]

$$d_{Hg} \propto C_{qq}(J_{CP}) \propto J_{CP} G_F^2 \sim \mathcal{O}(10^{-36})$$

[Flambaum et al '84; Donoghue et al '87]

$$d_e^{equiv} \propto r C_S(J_{CP}) \propto r J_{CP} G_F^2 \sim \mathcal{O}(10^{-38})$$

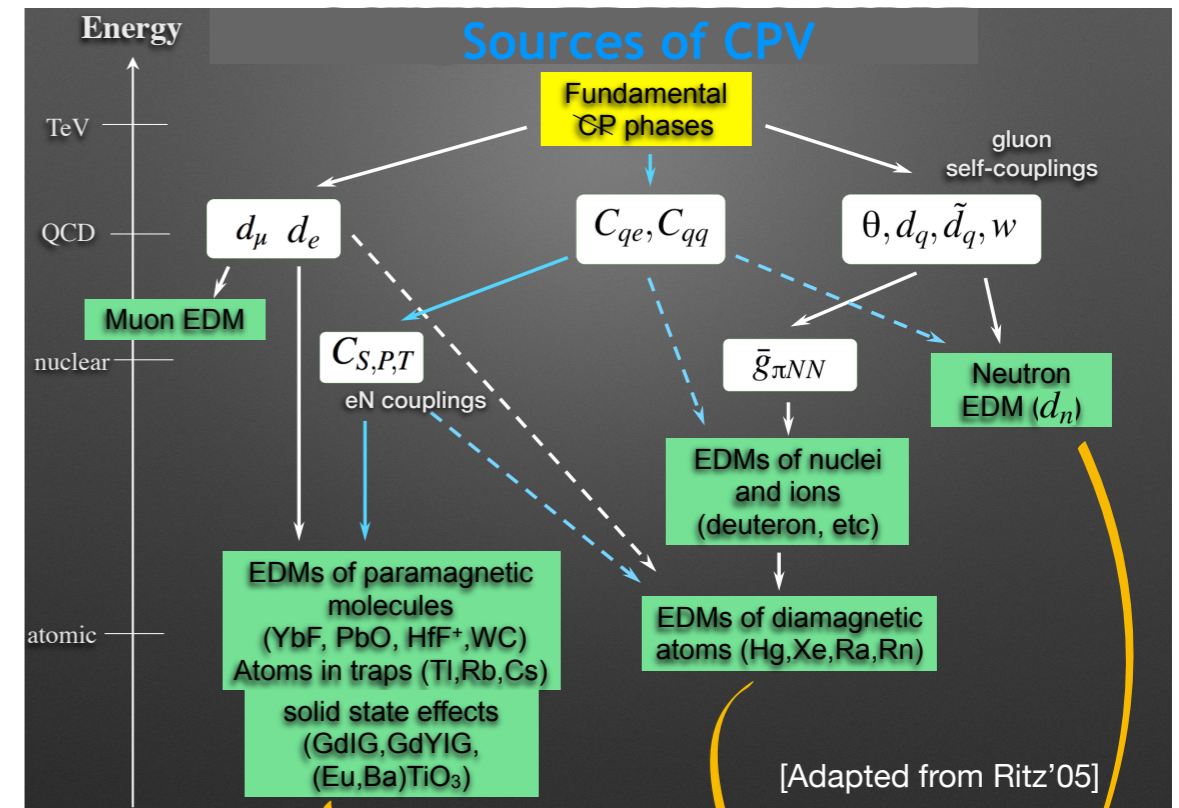
[Pospelov, Ritz '13]

⇒ Recent developments in d_e^{equiv} : **larger CP-odd**

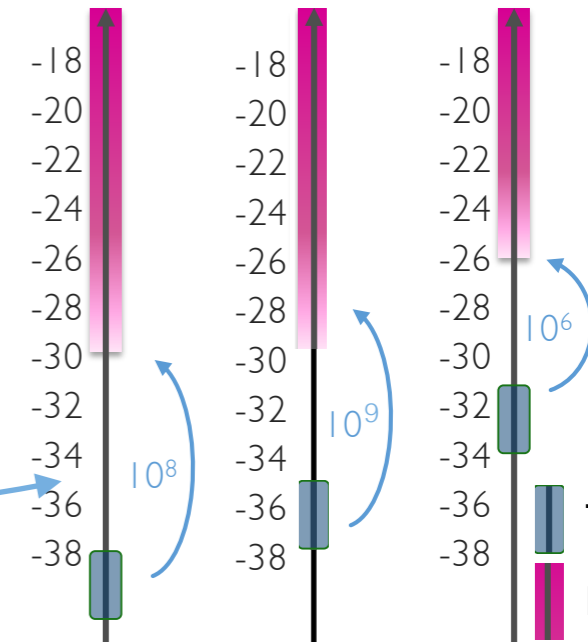
electron-nucleon interaction! $d_e^{equiv} \sim \mathcal{O}(10^{-35})$

[Ema, Gao, Pospelov '22]

⇒ **Still - room for New Physics sources of CPV!**



Electron log(d) / ecm Mercury Hg log(d) / ecm Neutron log(d) / ecm



Theory
Experiment

New physics contributions to EDMs

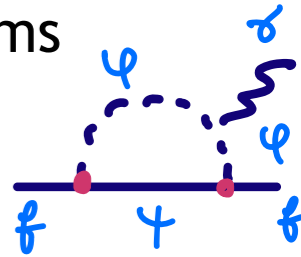
New CPV sources must necessarily be present in Nature - not enough CPV for BAU
 From the mechanism of neutrino mass generation (Dirac & Majorana phases) !?
 And from generic sources present in SM extensions...

► EDM constraints on New Physics sources of CP violation

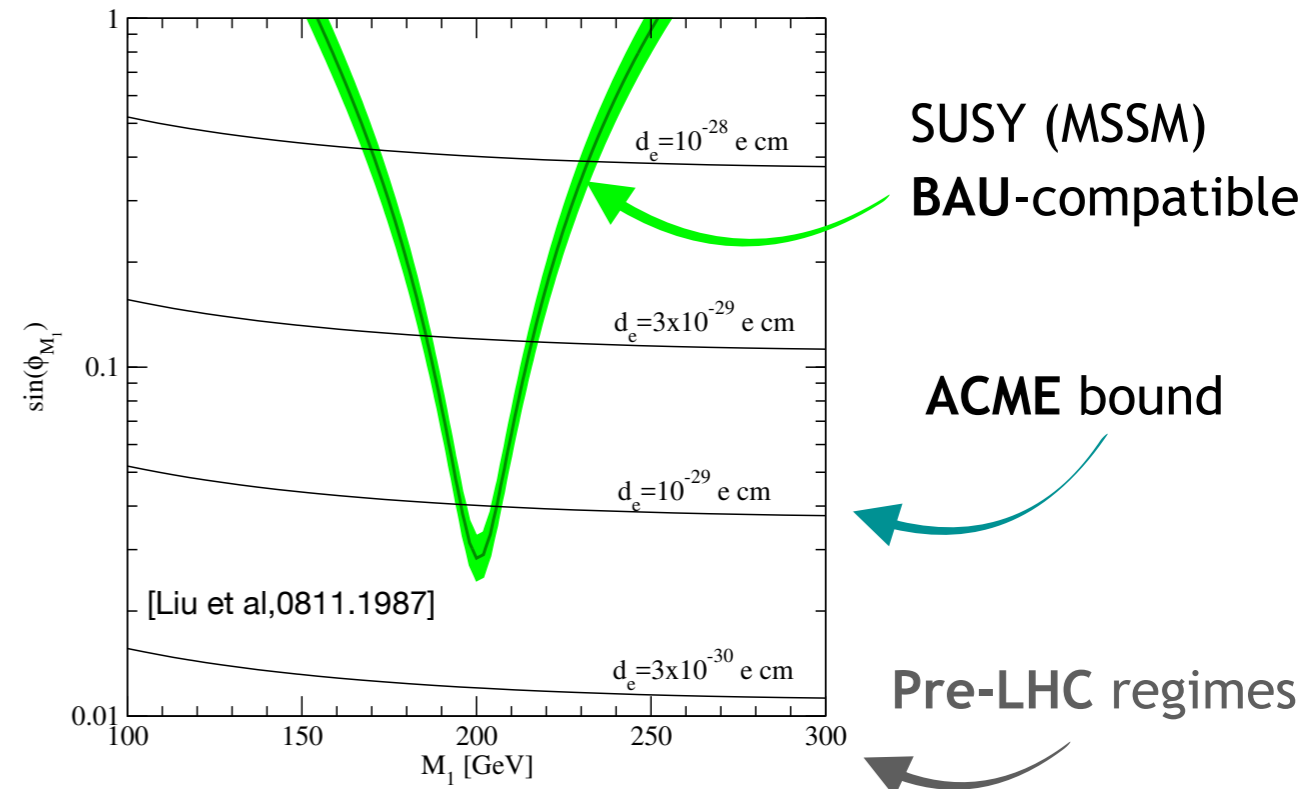
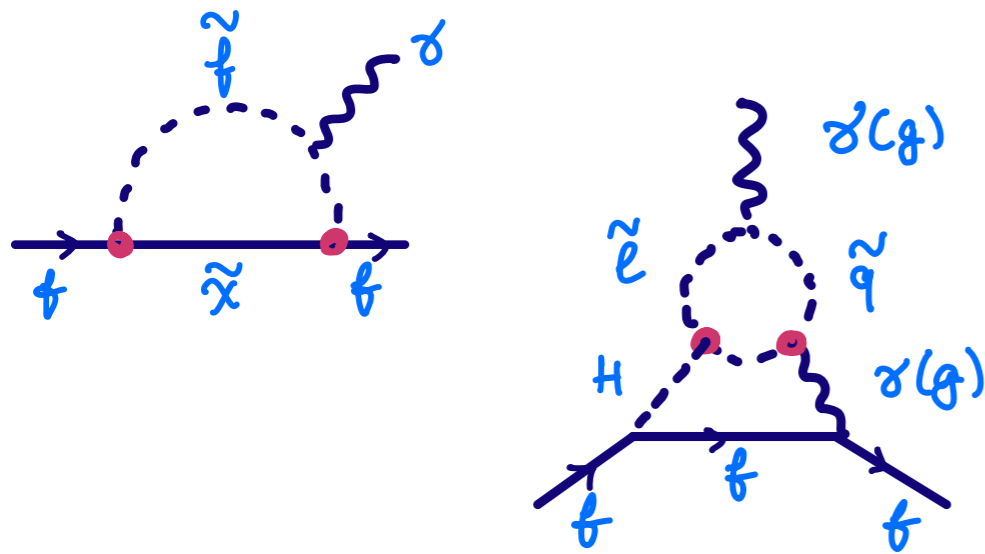
A very naive example - Supersymmetry

Generically - numerous flavour-blind CP phases in the soft SUSY-breaking terms

"Light" scalar states, extra CPV \Rightarrow SUSY EW baryogenesis!



But... unless very heavy s-particles, EDMs constrain phases to be *unnaturally small*
SUSY CP (and flavour) problem



New physics contributions to EDMs (EFT)

New CPV sources must necessarily be present in Nature - not enough CPV for BAU

From the mechanism of neutrino mass generation (Dirac & Majorana phases) !?

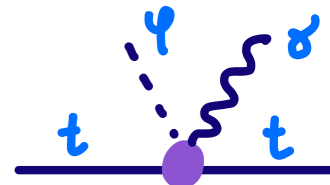
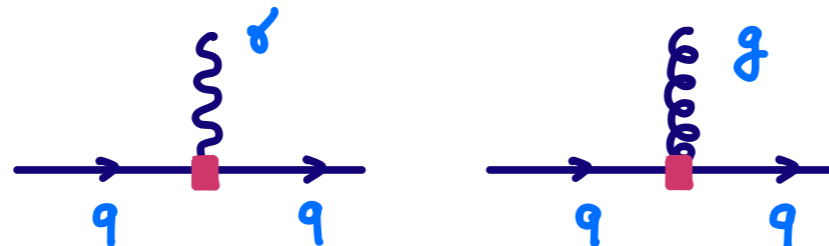
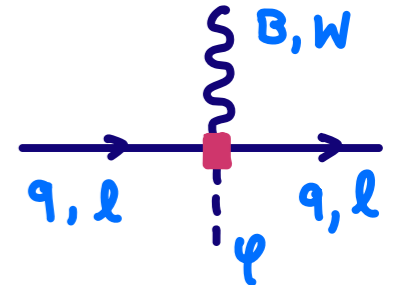
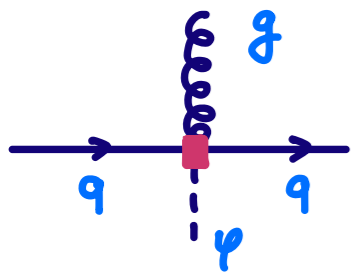
And from generic sources present in SM extensions...

EFT approach (dim 6 operators, flavour conserving but CP violating)

$$\mathcal{L}^{\text{dim } 6} \supset (\mathcal{C}_{eB}^{6\alpha\beta} / \Lambda^2) (\bar{L}_\alpha H \bar{\sigma}^{\mu\nu} \bar{E}_\beta^c) B_{\mu\nu} + (\mathcal{C}_{eW}^{6\alpha\beta} / \Lambda^2) (\bar{L}_\alpha \sigma^k H \bar{\sigma}^{\mu\nu} \bar{E}_\beta^c) W_{\mu\nu}^k +$$

$$+ (\mathcal{C}_{uG}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha \tilde{H} T^a \bar{\sigma}^{\mu\nu} \bar{U}_\beta^c) G_{\mu\nu}^a + (\mathcal{C}_{dG}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha H T^a \bar{\sigma}^{\mu\nu} \bar{D}_\beta^c) G_{\mu\nu}^a + \dots + H.c.$$

$$+ (\mathcal{C}_{uH}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha \tilde{H} \bar{U}_\beta^c) H^\dagger H + \dots$$

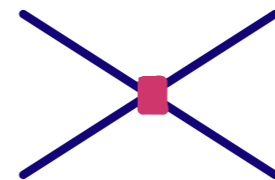
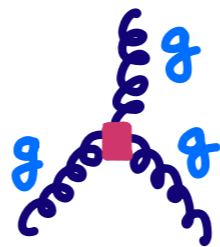


At low energies:

"diamagnetic" EDMs

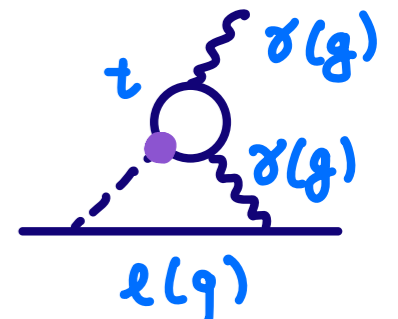
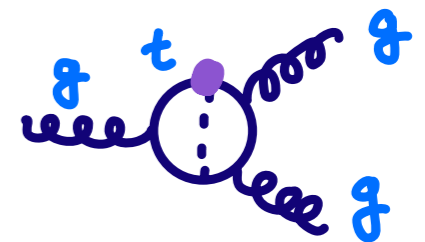
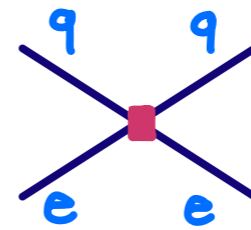
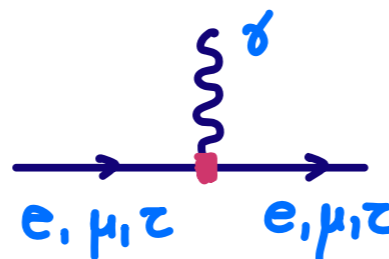
(neutron, ¹⁹⁹Hg, ²²⁵Ra, ...)

$\bar{\theta}$



"paramagnetic" EDMs

(muon, ²⁰⁵Ta, ThO, HfF, ...)



New CPV sources must necessarily **be present in Nature** - not enough CPV for BAU
 From the mechanism of neutrino mass generation (Dirac & Majorana phases) !?
 And from **generic sources** present in **SM extensions...**

EFT approach (dim 6 operators, *flavour conserving but CP violating*)

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$$+ (\mathcal{C}_{uG}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha H T^a \bar{\sigma}^{\mu\nu} \bar{U}_\beta^c) G_{\mu\nu}^a + (\mathcal{C}_{dG}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha H T^a \bar{\sigma}^{\mu\nu} \bar{D}_\beta^c) G_{\mu\nu}^a + \dots + H.c.$$

$$+ (\mathcal{C}_{uH}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha \tilde{H} \bar{U}_\beta^c) H^\dagger H + \dots$$

$\Rightarrow d_e = -2\sqrt{2} \cos \theta_w v \mathcal{F}([\mathcal{C}_{eB}^6]_{11} / \Lambda^2)$

ACME (2018): $|d_e| < 1.1 \times 10^{-29}$ e.cm $\Rightarrow \frac{|\mathcal{F}[\mathcal{C}_{eB}^6]_{11}|}{\Lambda^2} \leq \frac{1}{(1.9\text{EeV})^2}$

Sensitivity to **New Physics**: $\mathcal{F}(\mathcal{C}_{eB}^6) \sim \mathcal{O}(1) \rightarrow \Lambda \geq 10^6$ TeV

$$\mathcal{F}(\mathcal{C}_{eB}^6) \sim \mathcal{O}(Y^e) \rightarrow \Lambda \geq 10^3 \text{ TeV}$$

well above that of **direct LHC (or any collider) discovery!**

Clear impact on CPV models of new physics...

New physics contributions to EDMs (EFT)

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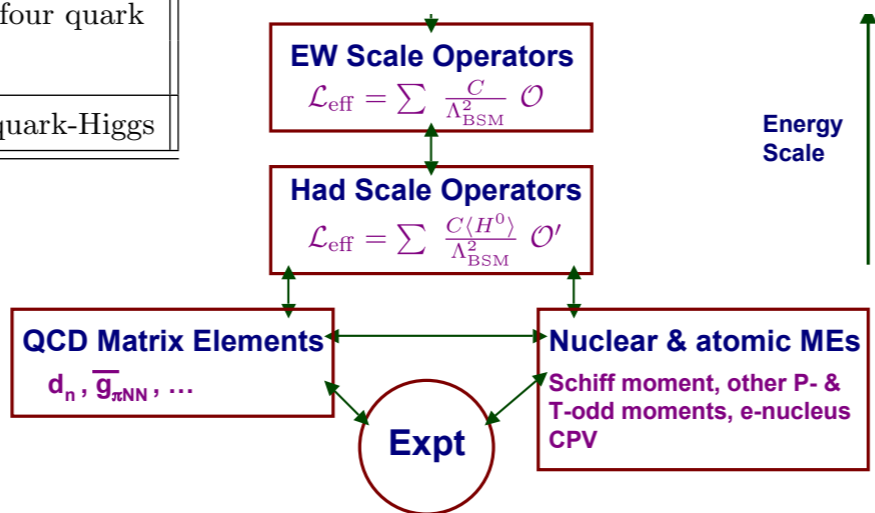
$$+ (\mathcal{C}_{uG}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha \tilde{H} T^a \bar{\sigma}^{\mu\nu} \bar{U}_\beta^c) G_{\mu\nu}^a + (\mathcal{C}_{dG}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha H T^a \bar{\sigma}^{\mu\nu} \bar{D}_\beta^c) G_{\mu\nu}^a + \dots + H.c.$$

$$+ (\mathcal{C}_{uH}^{6\alpha\beta} / \Lambda^2) (\bar{Q}_\alpha \tilde{H} \bar{U}_\beta^c) H^\dagger H + \dots$$

[Engel et al, 1303.2371]

\mathcal{O}_{fW}	$(\bar{F} \sigma^{\mu\nu} f_R) \tau^I \Phi W_{\mu\nu}^I$	fermion $SU(2)_L$ EDM
\mathcal{O}_{fB}	$(\bar{F} \sigma^{\mu\nu} f_R) \Phi B_{\mu\nu}$	fermion $U(1)_Y$ EDM
\mathcal{O}_{uG}	$(\bar{Q} \sigma^{\mu\nu} T^A u_R) \tilde{\varphi} G_{\mu\nu}^A$	u-quark Chromo EDM
\mathcal{O}_{dG}	$(\bar{Q} \sigma^{\mu\nu} T^A d_R) \varphi G_{\mu\nu}^A$	d-quark Chromo EDM
Q_{ledq}	$(\bar{L}^j e_R) (\bar{d}_R Q^j)$	CP-violating semi-leptonic
$Q_{lequ}^{(1)}$	$(\bar{L}^j e_R) \epsilon_{jk} (\bar{Q}^k u_R)$	
$Q_{lequ}^{(3)}$	$(\bar{L}^j \sigma_{\mu\nu} e_R) \epsilon_{jk} (\bar{Q}^k \sigma^{\mu\nu} u_R)$	
$\mathcal{O}_{\tilde{G}}$	$f^{ABC} \tilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	CP-violating 3 gluon
$Q_{quqd}^{(1)}$	$(\bar{Q}^j u_R) \epsilon_{jk} (\bar{Q}^k d_R)$	CP-violating four quark
$Q_{quqd}^{(8)}$	$(\bar{Q}^j T^A u_R) \epsilon_{jk} (\bar{Q}^k T^A d_R)$	
$Q_{\varphi ud}$	$i (\tilde{\varphi}^\dagger D_\mu \varphi) \bar{u}_R \gamma^\mu d_R$	CP-violating quark-Higgs

EFTs across scales & systems

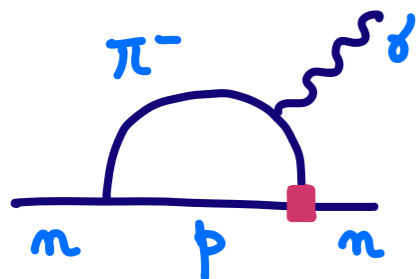


Nucleon level:

$$\mathcal{L}_N^S = -\frac{G_F}{\sqrt{2}} (\bar{e} i \gamma_5 e) \bar{N} \left[C_S^{(0)} + C_S^{(1)} \tau_3 \right] N$$

$$\mathcal{L}_N^T = \frac{8G_F}{\sqrt{2}} (\bar{e} \sigma_{\mu\nu} e) v_\nu \bar{N} \left[C_T^{(0)} + C_T^{(1)} \tau_3 \right] S_\mu N$$

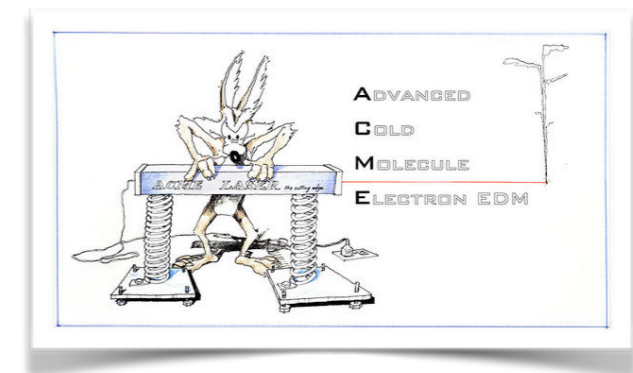
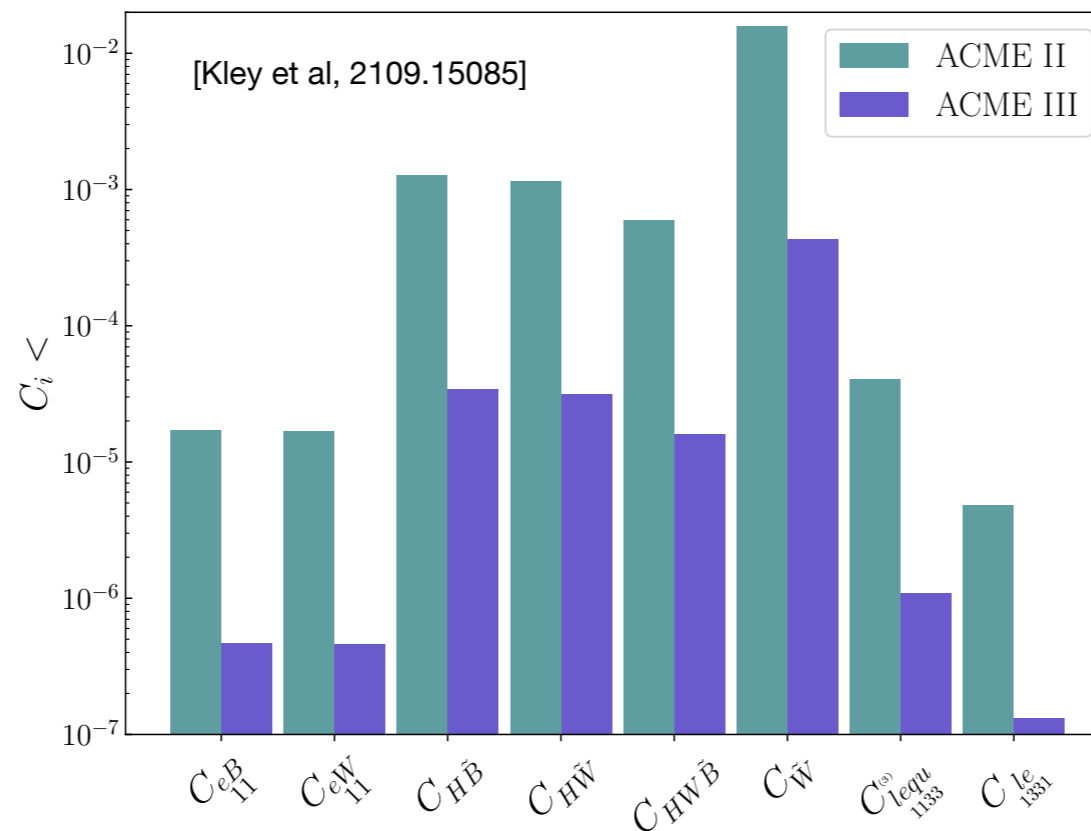
$$C_{S,T} = C_{S,T}(g_{S,T}, \mathfrak{S} C_{\ell equ(d)}, (v/\Lambda_{\text{NP}})^2, \dots)$$



New physics contributions to eEDM

New CPV sources must necessarily **be present in Nature** - not enough CPV for BAU
 From the mechanism of neutrino mass generation (Dirac & Majorana phases) !?
 And from **generic sources** present in **SM extensions...**

► **EDM constraints** on New Physics sources of **CP violation** - Wilson coefficients \mathcal{C}_{ij}
 (assuming $\Lambda_{NP} \approx 5 \text{ TeV}$)

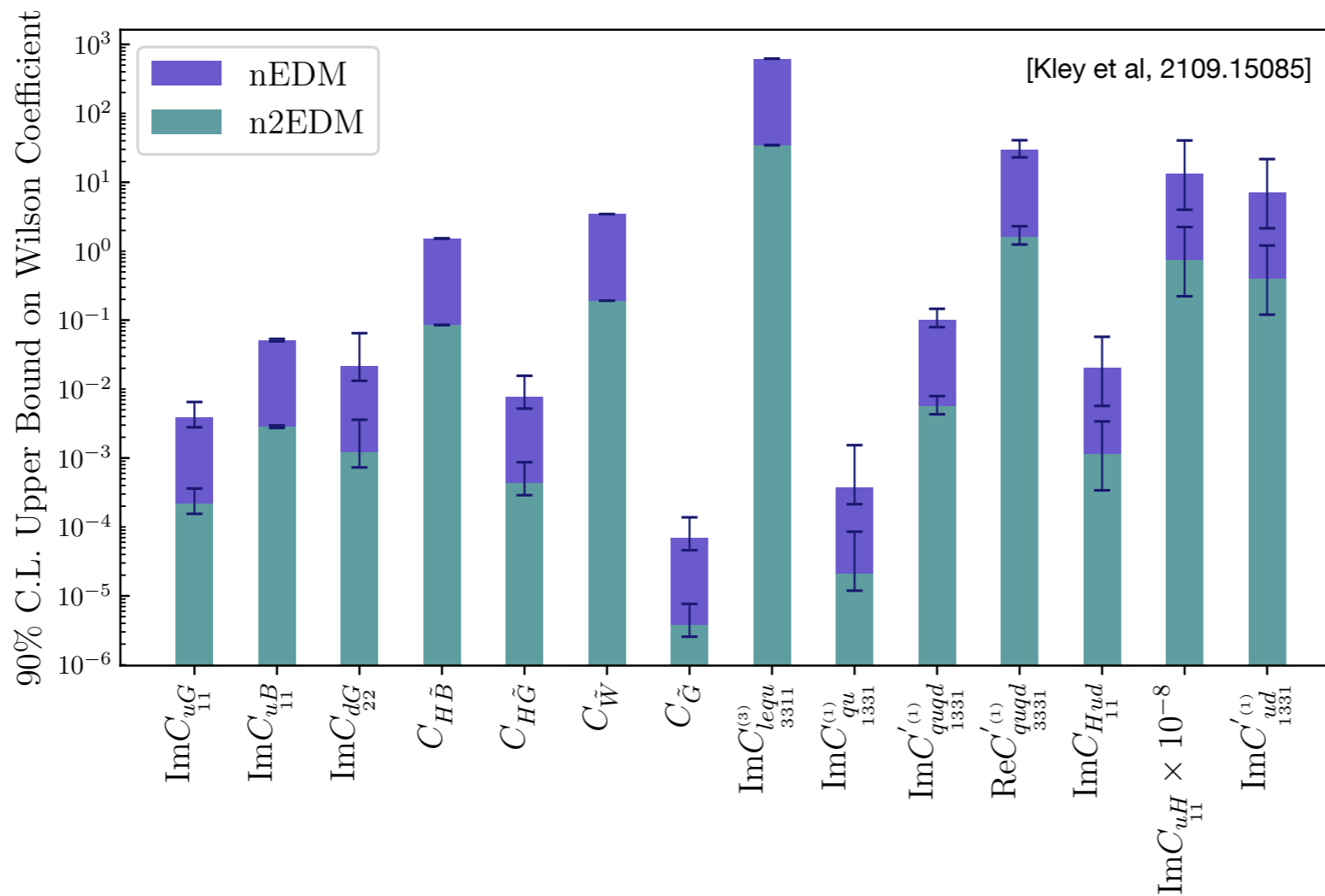


⇒ Impact of **ACME III** expected sensitivity: over **one order of magnitude** in \mathcal{C}_{ij} bounds

New physics contributions to nEDM

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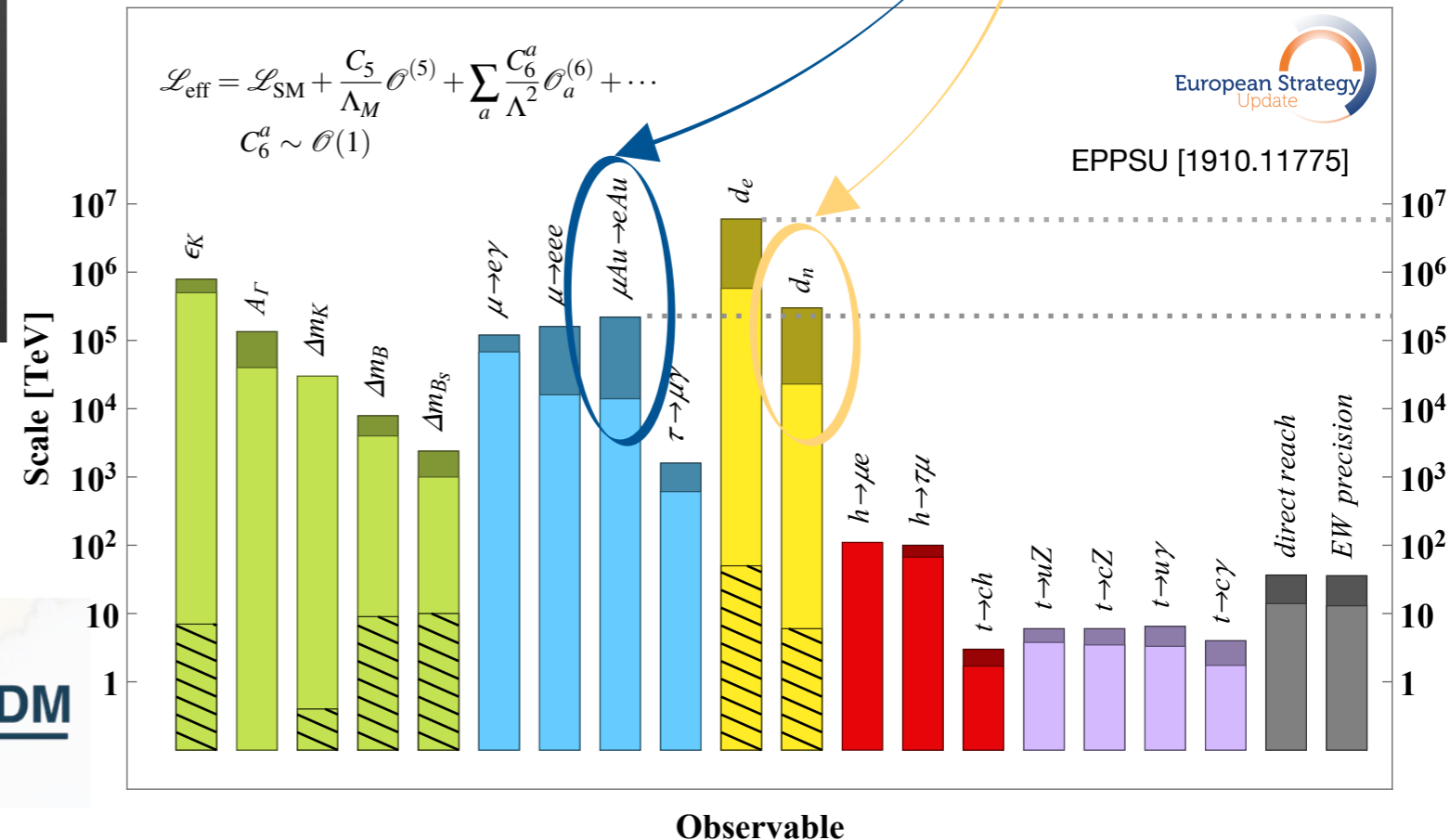
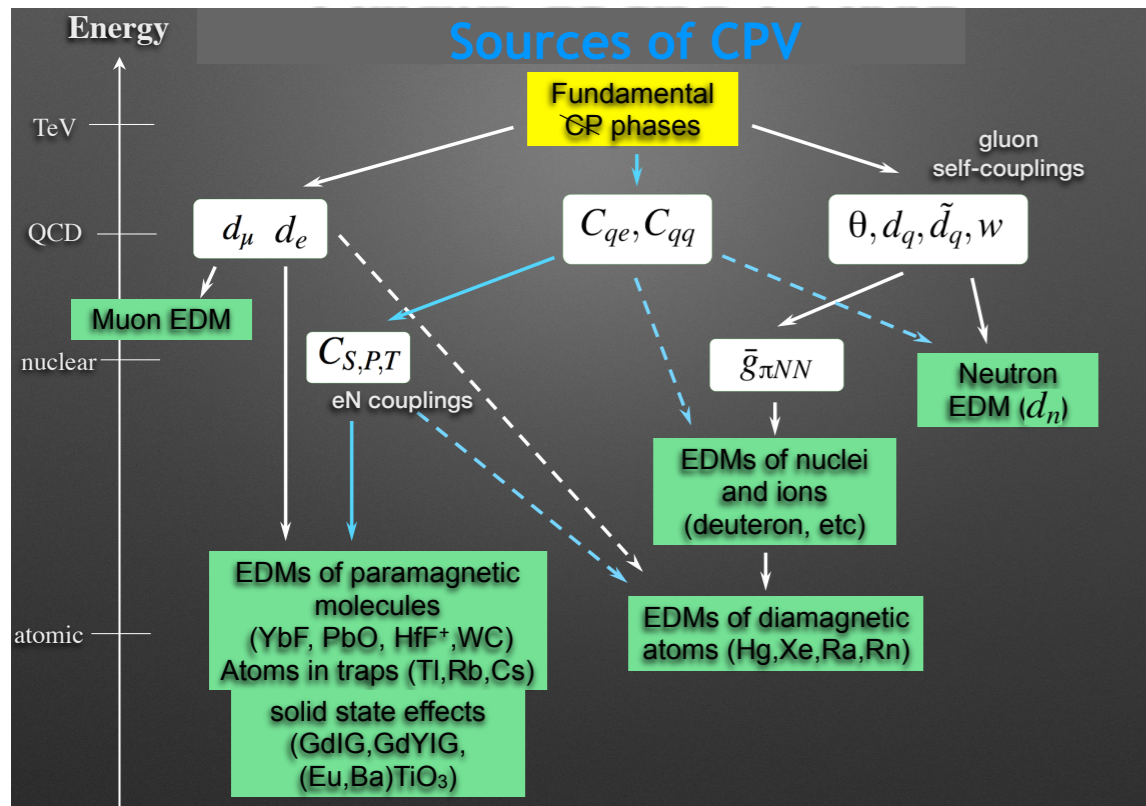


⇒ **Impact of n2EDM expected sensitivity:** over **one order of magnitude** in \mathcal{C}_{ij} bounds as well!

New CPV sources must be present- not enough CPV for BAU

EDM searches: powerful probes for New Physics sources of CP violation

$\mu - e$ conversion & d_n :
reaching unreachable scales

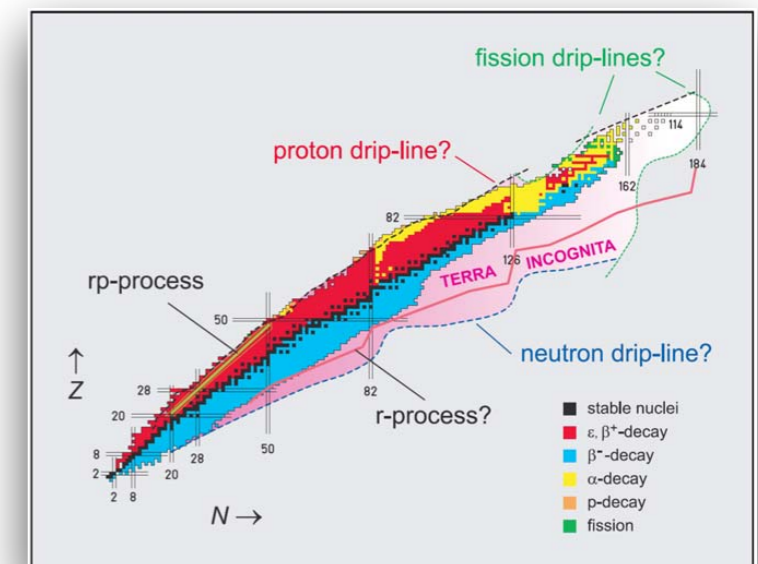


Neutron EDM:
among the most sensitive probes of NP!

Instrumental in the path to
understanding the BAU



New Physics searches: electroweak precision tests in nuclear beta decays



- **Nuclear β decays:** instrumental in **determining the structure of weak interactions** and **establishing the SM** (hadrons, charged leptons and neutrinos!)

SM semileptonic charged-current (weak) processes

- (i) **dominant $V - A$** component

$V + A, S, P, T$ only at **higher orders** in radiative corrections (or recoil momentum)

- (ii) **effective Fermi constants** (extracted from β decays) obey

lepton and quark-lepton (Cabibbo) **universality** \Rightarrow **unitarity of CKM matrix**

Low-energy charged-current interaction **Lagrangian** sensitive to many **SM extensions**

Theoretical and experimental progress \Rightarrow **0.1% level precision**

powerful constraints and **hints** on **BSM realisations at the TeV scale**

- ▶ **Nuclear β decays:** instrumental in determining the structure of **weak interactions** and establishing the **SM** (hadrons, charged leptons and neutrinos!)

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- ▶ **"Broad band" probes of New Physics**

Superaligned $0^+ \rightarrow 0^+$ transitions (& theoretical progress on radiative corrections...)

$\Rightarrow \delta V_{ud} \sim 3 \times 10^{-4}$ (sensitivity to $\Lambda_{NP} \sim 10$ TeV)

Measurement of τ_n and **β -asymmetry** (& LQCD precise determination of nucleon g_A)

\Rightarrow **probe right-handed currents** @subpercent level...

Superaligned transitions & neutron decay & mirror β decays

\Rightarrow **strong limits on strength** of $V + A, S, P, T$ interactions ($\leq 0.001 g_w$)

Angular correlations of β decay products

\Rightarrow search for **non-standard sources of CP violation!**

- ▶ **Nuclear β decays:** instrumental in determining the structure of **weak interactions** and establishing the **SM** (hadrons, charged leptons and neutrinos!)

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lepton and quark-lepton (Cabibbo) **universality** \Rightarrow **unitarity of CKM matrix**

- ▶ **"Broad band" probes of New Physics**

- ▶ In the **LHC era**, β decays remain uniquely **precise probes of New Physics**, highly **competitive** on their own, and **complementary** to searches at high-energies

Charged- and neutral current **Drell-Yann production at LHC** - directly access the **TeV scale**;

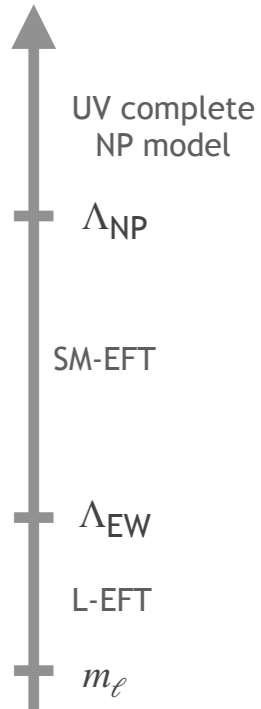
$$pp \rightarrow e + \nu + X$$

\Rightarrow **synergy of constrains places strong bounds on EFT couplings!**

Nuclear (effective) interactions

- At low-energies ($\mu \sim 2$ GeV), relevant information on underlying physics (**SM, NP**) for β decays - involving CC transitions in first family (u, d and e)

$$\begin{aligned} \mathcal{L}_{\text{CC}}^{\text{eff}} = & -\frac{G_F^{(0)} V_{ud}}{\sqrt{2}} \left[(1 + \delta_\beta) \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \bar{u} \gamma^\mu (1 - \gamma_5) d \right. \\ & + \epsilon_L \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \bar{u} \gamma^\mu (1 - \gamma_5) d + \tilde{\epsilon}_L \bar{e} \gamma_\mu (1 + \gamma_5) \nu_e \bar{u} \gamma^\mu (1 - \gamma_5) d \\ & + \epsilon_R \bar{e} \gamma_\mu (1 - \gamma_5) \nu_e \bar{u} \gamma^\mu (1 + \gamma_5) d + \tilde{\epsilon}_R \bar{e} \gamma_\mu (1 + \gamma_5) \nu_e \bar{u} \gamma^\mu (1 + \gamma_5) d \\ & + \epsilon_S \bar{e} (1 - \gamma_5) \nu_e \bar{u} d + \tilde{\epsilon}_S \bar{e} (1 + \gamma_5) \nu_e \bar{u} d \\ & - \epsilon_P \bar{e} (1 - \gamma_5) \nu_e \bar{u} \gamma_5 d - \tilde{\epsilon}_P \bar{e} (1 + \gamma_5) \nu_e \bar{u} \gamma_5 d \\ & \left. + \epsilon_T \bar{e} \sigma_{\mu\nu} (1 - \gamma_5) \nu_e \bar{u} \sigma^{\mu\nu} (1 - \gamma_5) d + \tilde{\epsilon}_T \bar{e} \sigma_{\mu\nu} (1 + \gamma_5) \nu_e \bar{u} \sigma^{\mu\nu} \gamma_\mu (1 + \gamma_5) d \right] \end{aligned}$$



A few remarks:

$$G_F^{(0)} = \sqrt{2} \frac{g^2}{8M_W^2} \text{ (tree-level Fermi constant); } G_F^{(0)} = G_\mu^{\text{exp}} (1 - \delta_\mu - \epsilon_e)$$

$\delta_{\mu(\beta)}$ \rightsquigarrow **SM EW corrections** in purely (semi)leptonic decays

ϵ (left-handed **neutrino** currents); $\tilde{\epsilon}$ (right-handed **neutrino** currents)

$$\Rightarrow \Gamma \propto f(\epsilon), g(\tilde{\epsilon}^2), h(\tilde{\epsilon} \times m_\nu / E_\nu)$$

"All" associated operators can produce collider signatures!

- Studying **nuclear and hadronic transition amplitudes** - far more involved!
short-distance couplings evolved to appropriate matching scale
hadronic & nuclear matrix elements

Inferring information on the **Wilson coefficients** (sensitive to the presence of **NP**) from nuclear and hadronic observables \Rightarrow knowledge of the **matrix elements!**

A few examples:

- (Semi)leptonic **decays of mesons** $M \rightarrow \ell \nu$ and $M \rightarrow M' \ell \nu$:
decay constants and form factors from LQCD!
- At the nucleon level, **matrix elements** of **neutron to proton decays** (dim-3 quark bilinears)
- And the ultimate **challenge**, from **nucleon level to nuclear beta decays...**
numerous spin sequences, unstable daughter nucleus (under em or strong interactions)
large Q -value (e^- and e^+)...

Nuclear (effective) interactions

► **SM precision tests** in nuclear decays: **identify NP contributions** in ϵ and $\tilde{\epsilon}$ from **low-energy observables!**

⇒ connect the **quark-level Lagrangian** to the **nucleon-level** formulation

$$\begin{aligned} \mathcal{L}_{L-Y}^{\text{eff}} = & -\bar{p}\gamma^\mu n (C_V \bar{e}\gamma_\mu \nu - C'_V \bar{e}\gamma_\mu \gamma_5 \nu) + \bar{p}\gamma^\mu \gamma_5 n (C_A \bar{e}\gamma_\mu \gamma_5 \nu - C'_A \bar{e}\gamma_\mu \nu) \\ & -\bar{p}n (C_S \bar{e}\nu - C'_S \bar{e}\gamma_5 \nu) - \frac{1}{2} \bar{p}\sigma^{\mu\nu} n (C_T \bar{e}\sigma^{\mu\nu} \nu - C'_T \bar{e}\sigma_{\mu\nu} \gamma_5 \nu) \\ & -\bar{p}\gamma_5 n (C_P \bar{e}\gamma_5 \nu - C'_P \bar{e}\nu) + \text{H.c.} \end{aligned}$$

Relating coefficients: $C^{(\prime)} = \bar{C}^{(\prime)} V_{ud} G_F^{(0)} / \sqrt{2}$

$$\bar{C}_V^{(\prime)} = g_V (1 + \delta_\beta + \epsilon_L + \epsilon_R \pm \tilde{\epsilon}_L \pm \tilde{\epsilon}_R)$$

$$\bar{C}_A^{(\prime)} = -g_A (1 + \delta_\beta + \epsilon_L - \epsilon_R \mp \tilde{\epsilon}_L \pm \tilde{\epsilon}_R)$$

$$\bar{C}_S^{(\prime)} = g_S (\epsilon_S \pm \tilde{\epsilon}_S)$$

$$\bar{C}_P^{(\prime)} = g_P (\epsilon_P \mp \tilde{\epsilon}_P)$$

$$\bar{C}_T^{(\prime)} = 4g_T (\epsilon_T \pm \tilde{\epsilon}_T)$$

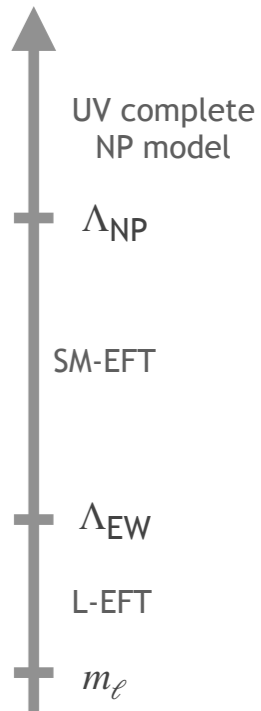
g_V, g_A, g_S, g_P, g_T

vector, axial, (pseudo)scalar and tensor *nuclear* charges

⇒ determined from LQCD or other theoretical methods....

[Gonzalez-Alonso et al, 1803.732]

Charge	Value	Ref.
g_A	1.278(33)	[35]
g_T	0.987(55)	[34]
g_S	1.02(11)	[24]
g_P	349(9)	[24]



⇒ Understand the role of **short-distance couplings** ϵ and $\tilde{\epsilon}$ (and constrain them!) from **β decays** and other **low-energy observables**

- ▶ **SM precision tests** in nuclear decays: **identify NP contributions** in ϵ and $\tilde{\epsilon}$ from **low-energy observables!** Which hadronic and nuclear observables?

⇒ Ideally, looking for **experimentally feasible set-ups**, investigating **systems theoretically "under control"** (small degree of uncertainties), **thoroughly exploring decays** (rates, angular correlations, spectrum shape, recoil, ...)

- ▶ **"Standard" experimental systems:** allowed **Fermi** and/or **Gamow-Teller** decays
Neutron decays - simplest baryon!

Theoretical description mastered to high precision (progress in LQCD, ...)

(Super)allowed decays - small number of matrix elements

Theoretical precision challenged...

Mirror nuclei - mixed Fermi and Gamow-Teller

Excellent field for CPV searches and right-handed neutrino currents

$T = 1$ (isospin triplet) **pure Gamow-Teller decays**

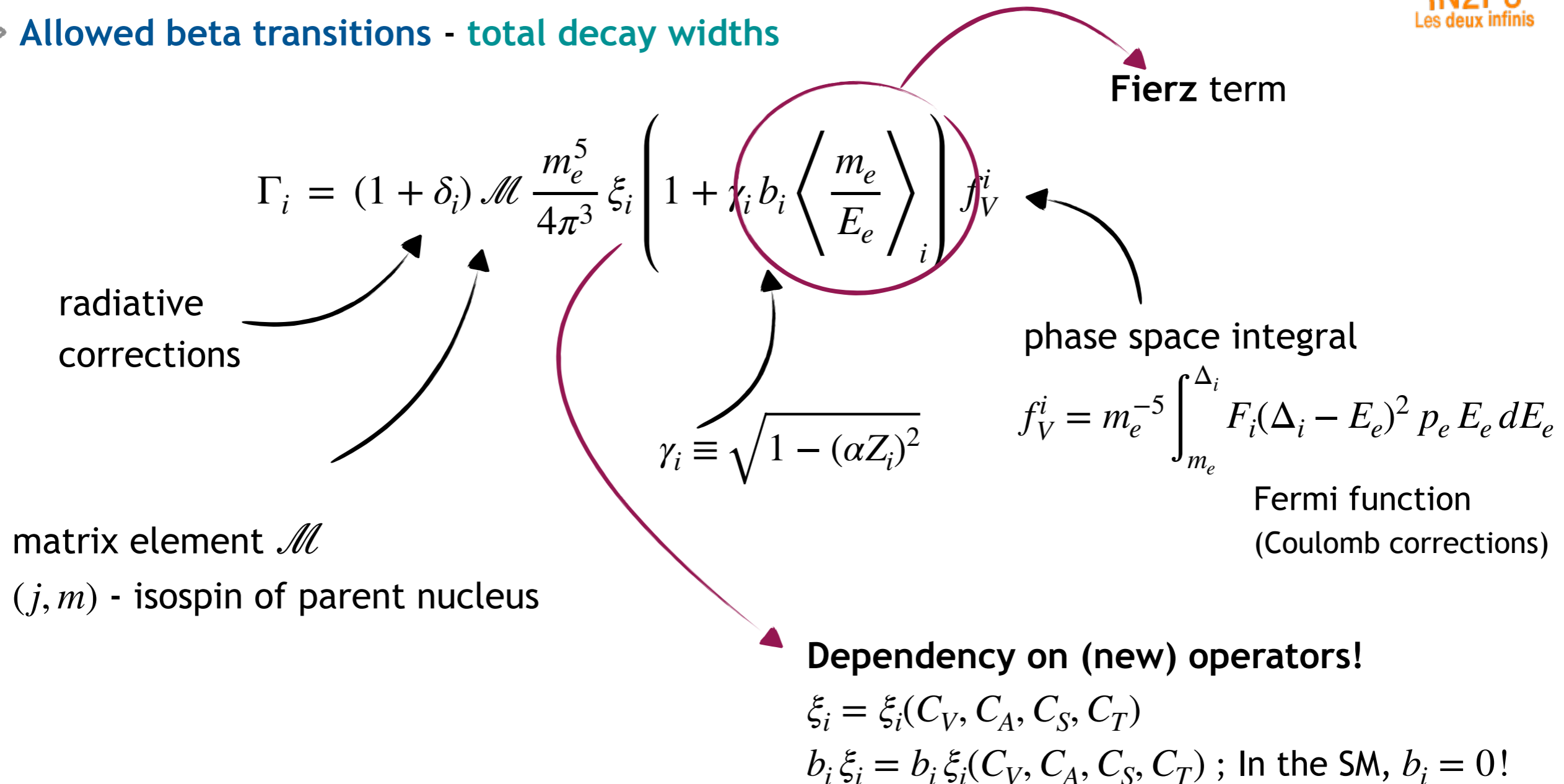
Small theoretical uncertainties, access to exotic tensor currents

- ▶ **"New" opportunities?**

Pseudoscalar decays? Unique forbidden beta decays?

Total decay widths (integrated rates)

► Allowed beta transitions - total decay widths



$$\xi = |M_F|^2 \left(|C_V|^2 + |C_V'|^2 + |C_S|^2 + |C_S'|^2 \right) + |M_{GT}|^2 \left(|C_A|^2 + |C_A'|^2 + |C_T|^2 + |C_T'|^2 \right)$$

$$b \xi = \pm \sqrt{1 + \alpha^2 Z^2} \Re \left[|M_F|^2 \left(C_V C_S^* + C_V' C_S'^* \right) + |M_{GT}|^2 \left(C_A C_T^* + C_A' C_T'^* \right) \right]$$

Total decay widths (integrated rates)

► Allowed beta transitions - total decay widths

radiative corrections

matrix element \mathcal{M}
(j, m) - isospin of parent nucleus

$$\Gamma_i = (1 + \delta_i) \mathcal{M} \frac{m_e^5}{4\pi^3} \xi_i \left(1 + \gamma_i b_i \left\langle \frac{m_e}{E_e} \right\rangle_i \right) f_V^i$$

$\gamma_i \equiv \sqrt{1 - (\alpha Z_i)^2}$

Fierz term

phase space integral

$$f_V^i = m_e^{-5} \int_{m_e}^{\Delta_i} F_i(\Delta_i - E_e)^2 p_e E_e dE_e$$

Fermi function
(Coulomb corrections)

Dependency on (new) operators!
 $\xi_i = \xi_i(C_V, C_A, C_S, C_T)$
 $b_i \xi_i = b_i \xi_i(C_V, C_A, C_S, C_T)$; In the SM, $b_i = 0!$

Corrected half-life: $\mathcal{F}t_i \equiv \frac{f_V^i(1 + \delta_i)\log 2}{\Gamma_i} = \frac{4\pi^3 \log 2}{M_F^2 m_e^5} \left(\xi_i + \gamma_i b_i \xi_i \left\langle \frac{m_e}{E_e} \right\rangle_i \right)^{-1}$

Total decay widths (integrated rates)

▶ Allowed beta transitions - total decay widths

Extensively studied in the literature for numerous nuclei (and the neutron!)

▶ A class of nuclear beta decays emerges as a uniquely powerful probe of the SM description of quark flavour violation (CKM paradigm)

⇒ **superallowed $0^+ \rightarrow 0^+$ decays**: depend uniquely on **vector part** of interaction

~ decay width directly related to C_V

[Hardy and Towner, '15 - '21]

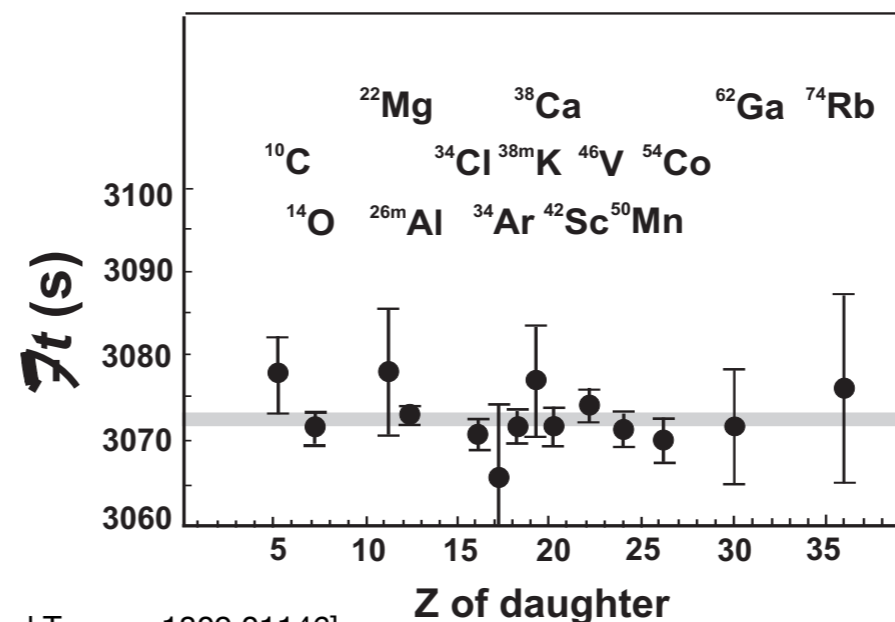
Parent nucleus	$ft(s)$	δ'_R (%)	$\delta_C - \delta_{NS}$ (%)	$\mathcal{F}t(s)$
$T_z = -1$:				
^{10}C	3042.4 ± 4.1	1.679	0.575 ± 0.039	3075.7 ± 4.4^a
^{14}O	3042.2 ± 0.8	1.543	0.613 ± 0.056	3070.2 ± 1.9^a
^{18}Ne	2912 ± 79	1.506	0.886 ± 0.052	2930 ± 80
^{22}Mg	3051.1 ± 6.9	1.466	0.635 ± 0.026	3076.2 ± 7.0^a
^{26}Si	3052.2 ± 5.6	1.439	0.669 ± 0.033	3075.4 ± 5.7^a
^{30}S	3015 ± 41	1.423	1.001 ± 0.049	3027 ± 41
^{34}Ar	3058.0 ± 2.8	1.412	0.840 ± 0.043	3075.1 ± 3.1^a
^{38}Ca	3062.8 ± 6.0	1.414	0.912 ± 0.049	3077.8 ± 6.2^a
^{42}Ti	3090 ± 88	1.424	1.193 ± 0.066	3097 ± 88
^{46}Cr	3126 ± 100	1.426	0.924 ± 0.089	3141 ± 100
^{50}Fe	3099 ± 71	1.426	0.800 ± 0.053	3118 ± 72
^{54}Ni	3062 ± 50	1.423	0.933 ± 0.070	3077 ± 50
$T_z = 0$:				
^{26m}Al	3037.61 ± 0.67	1.478	0.329 ± 0.026	3072.4 ± 1.1^a
^{34}Cl	$3049.43^{+0.95}_{-0.88}$	1.443	0.706 ± 0.051	3071.6 ± 1.8^a
^{38m}K	3051.45 ± 0.92	1.440	0.726 ± 0.056	3072.9 ± 2.0^a
^{42}Sc	3047.7 ± 1.2	1.453	0.657 ± 0.050	3071.7 ± 2.0^a
^{46}V	$3050.33^{+0.54}_{-0.44}$	1.445	0.651 ± 0.063	3074.3 ± 2.0^a
^{50}Mn	3048.4 ± 1.2	1.444	0.689 ± 0.033	3071.1 ± 1.6^a
^{54}Co	$3050.8^{+1.4}_{-1.1}$	1.443	0.787 ± 0.068	$3070.4^{+2.5a}_{-2.4}$
^{62}Ga	3074.1 ± 1.5	1.459	1.49 ± 0.21	3072.4 ± 6.7^a
^{66}As		1.468	1.58 ± 0.40	
^{70}Br		1.486	1.74 ± 0.25	
^{74}Rb	3082.8 ± 6.5	1.499	1.65 ± 0.27	3077 ± 11^a
Average (best 15), $\overline{\mathcal{F}t}$				3072.24 ± 0.57
				χ^2/ν 0.47

$$\mathcal{F}t = \frac{K}{2 G_F^2 \bar{V}_{ud}^2 (1 + \Delta_R^V)}$$

$$K = 2\pi^3 \hbar \log 2 (\hbar c)^6 / (m_e c^2)^5$$

Δ_R^V - transition independent RC

$\mathcal{F}t$ values almost "constant" for superallowed $0^+ \rightarrow 0^+$ decays



Extremely precise determination!

[Hardy and Towner, 1809.01146]

Total decay widths (integrated rates)

▶ **Allowed beta transitions - total decay widths**

Extensively studied in the literature for numerous nuclei (and the neutron!)

▶ **A class of nuclear beta decays emerges as a uniquely powerful probe of the SM description of quark flavour violation (CKM paradigm)**

⇒ **superallowed $0^+ \rightarrow 0^+$ decays**: depend uniquely on **vector part** of interaction

→ decay width directly related to C_V

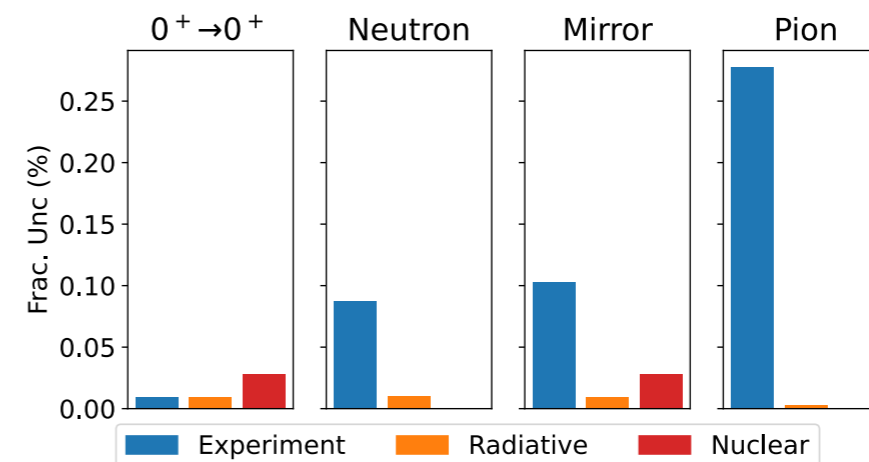
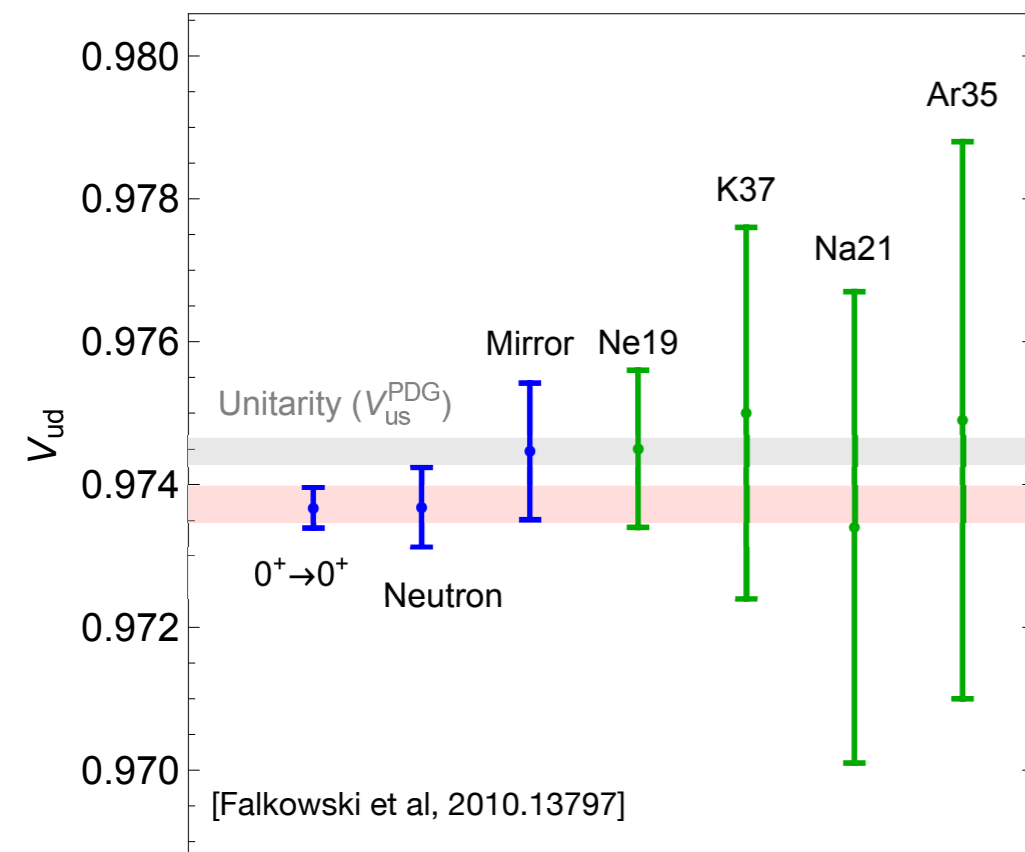
$$Ft = \frac{K}{2 G_F^2 \bar{V}_{ud}^2 (1 + \Delta_R^V)}$$

$$K = 2\pi^3 \hbar \log 2 (\hbar c)^6 / (m_e c^2)^5$$

Δ_R^V - transition independent RC

Determination of \bar{V}_{ud} dominated by

superallowed $0^+ \rightarrow 0^+$ decays (very precise! $\delta \sim 0.03\%$)



Best determination of \bar{V}_{us} from **Kaon decays** ($K_{\ell 2}, K_{\ell 3}$)

Testing the CKM paradigm

CKM paradigm of flavour mixing: **FV** encoded in *strongly hierarchical unitary* matrix
Mostly **successful description** of **hadron flavour dynamics!**

EW fit and V_{CKM} fit appear to be in good agreement with SM hypotheses!

But recent **tensions** in the determination of the "**Cabibbo angle**" (and of V_{ud} and V_{us})

$$V_{ud}^2 + V_{us}^2 + V_{ub}^2 = 1 \Rightarrow \sin \theta_C = V_{ud}, \cos \theta_C = V_{us}$$

$$\Rightarrow \Delta_{CKM} = |\bar{V}_{ud}|^2 + |\bar{V}_{us}|^2 - 1 = 0$$

Phenomenological determination of

\bar{V}_{ud} and \bar{V}_{us} (recall that $|\bar{V}_{ub}| \approx \mathcal{O}(10^{-3})$)

\Rightarrow test unitarity of **1st row of CKM**

Overview of \bar{V}_{ud} and \bar{V}_{us} constraints:

nuclear, nucleon, meson & τ decays (1σ bands for V_{ij})

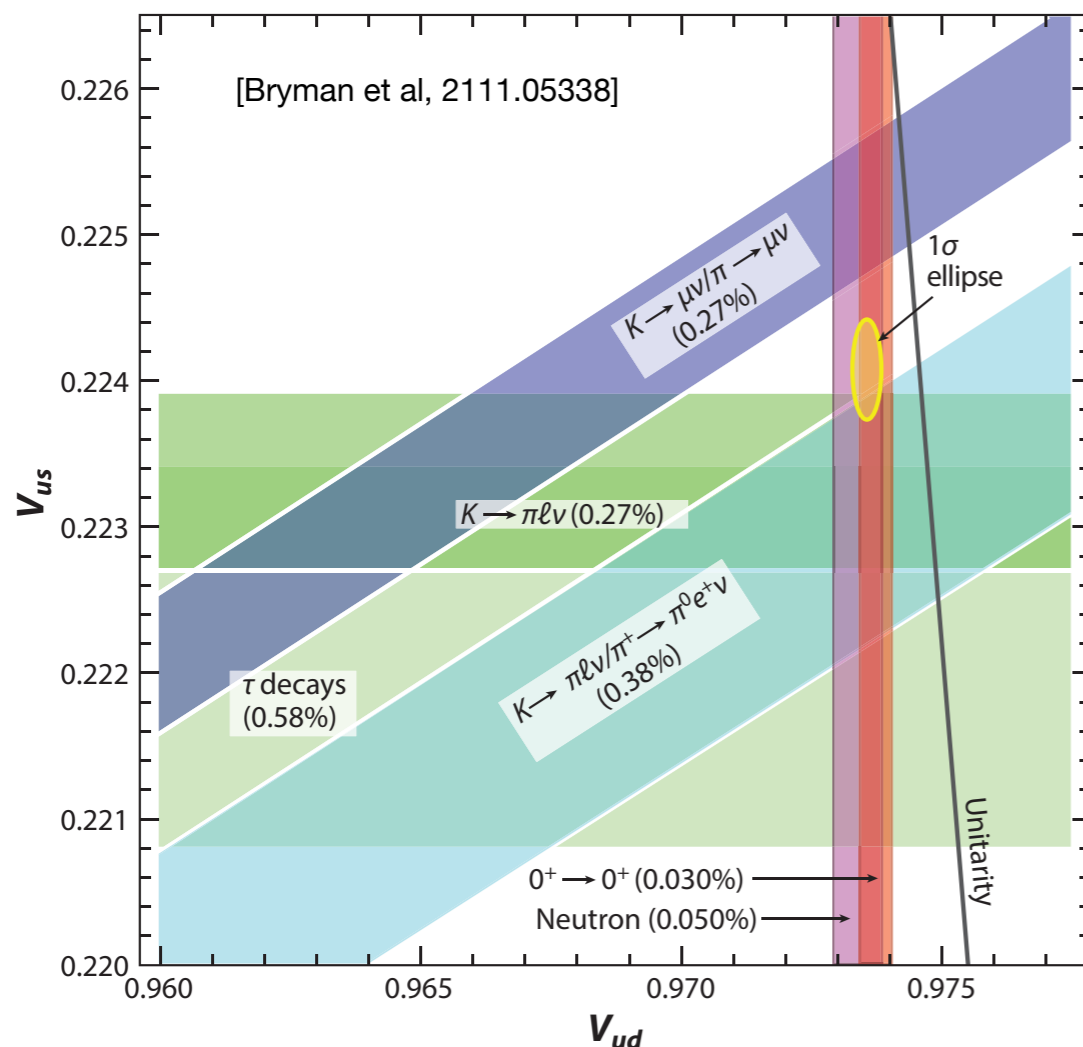
Global fit: $V_{ud} = 0.97379 \pm 0.00025$

$V_{us} = 0.22405 \pm 0.00035$

$\Rightarrow \Delta_{CKM} = (-19.5 \pm 5.3) \times 10^{-4}$

[Crivellin et al, 2212.06862]

\Rightarrow **Deviation from SM unitarity @ 2.8σ**



Differential decay distributions

► **SM precision tests in nuclear β decays:** a vast array of additional observables from the exploration of **angular correlations** between **decay products!**

► **Differential decay distributions** - very sensitive to the **underlying Lorentz structure**

⇒ searches for **non (V - A)** components in weak interactions!

$$\frac{d^3\Gamma}{dE_e d\Omega_e d\Omega_\nu} = \frac{1}{(2\pi)^5} p_e E_e (E_0 - E_e)^2 \xi \left\{ 1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \left\langle \frac{\vec{J}}{J} \right\rangle \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} + \dots \right] \right\}$$

Parent nucleus polarisation \vec{J}

Fierz interference term

$e - \bar{\nu}_e$ correlation

beta asymmetry

$\bar{\nu}_e$ asymmetry

T-odd term (CPV!)

Asymmetries and angular correlations

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T-odd term (CPV!)

Some remarks:

Overall factor ξ , and **correlation coefficients** depend on M_F , M_{GT} , on $C_i^{(\prime)}$ (and possibly E_e)

$$\xi = |M_F|^2 (|C_V|^2 + |C'_V|^2 + |C_S|^2 + |C'_S|^2) + |M_{GT}|^2 (|C_A|^2 + |C'_A|^2 + |C_T|^2 + |C'_T|^2)$$

$$a \xi = |M_F|^2 (|C_V|^2 + |C'_V|^2 - |C_S|^2 - |C'_S|^2) - \frac{1}{3} |M_{GT}|^2 (|C_A|^2 + |C'_A|^2 - |C_T|^2 - |C'_T|^2)$$

$$b \xi = \pm \sqrt{1 + \alpha^2 Z^2} \Re \left[|M_F|^2 (C_V C_S^* + C'_V C_S'^*) + |M_{GT}|^2 (C_A C_T^* + C'_A C_T'^*) \right]$$

[Gonzalez-Alonso et al, 1803.08732]

Asymmetries and angular correlations

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Fierz interference term
 $e - \bar{\nu}_e$ correlation
beta asymmetry
 $\bar{\nu}_e$ asymmetry
T-odd term (CPV!)

Some remarks:

Overall factor ξ , and **correlation coefficients** depend on M_F , M_{GT} , on $C_i^{(0)}$ (and possibly E_e)

$$\xi = |M_F|^2 \left(|C_V|^2 + |C_V'|^2 + |C_S|^2 + |C_S'|^2 \right) + |M_{GT}|^2 \left(|C_A|^2 + |C_A'|^2 + |C_T|^2 + |C_T'|^2 \right)$$

If **pure Fermi** transitions, dependence only on $C_{V,S}^{(0)}$

If **pure Gamow-Teller**, dependence only on $C_{A,T}^{(0)}$

If **mixed** (e.g. neutron decay), **combination** of $C_{V,S}^{(0)}$ & $C_{A,T}^{(0)}$

► **SM precision tests** in **nuclear β decays**: a vast array of additional observables from the exploration of **angular correlations** between **decay products**!

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⇒ searches for **non (V – A)** components in weak interactions!

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And further remarks:

To a good *first* approximation, and for pure F(GT) transitions

$$b_F \approx \pm \Re \left(\frac{C_S + C'_S}{C_V} \right) \text{ and } b_{GT} \approx \pm \Re \left(\frac{C_T + C'_T}{C_A} \right); \quad a_F \approx 1 - \frac{|C_S|^2 + |C'_S|^2}{|C_V|^2} \text{ and } a_{GT} \approx -\frac{1}{3} + \frac{1}{3} \frac{|C_T|^2 + |C'_T|^2}{|C_A|^2}$$

Measurement of the **Fierz** term, **$e - \bar{\nu}_e$ asymmetry** and **correlation parameters**

⇒ probe (combinations) of **non-standard** (particle-level) **coefficients**

in particular **ϵ_S** and **ϵ_T**

If measured, **a, A, B** include contributions from **Fierz** term

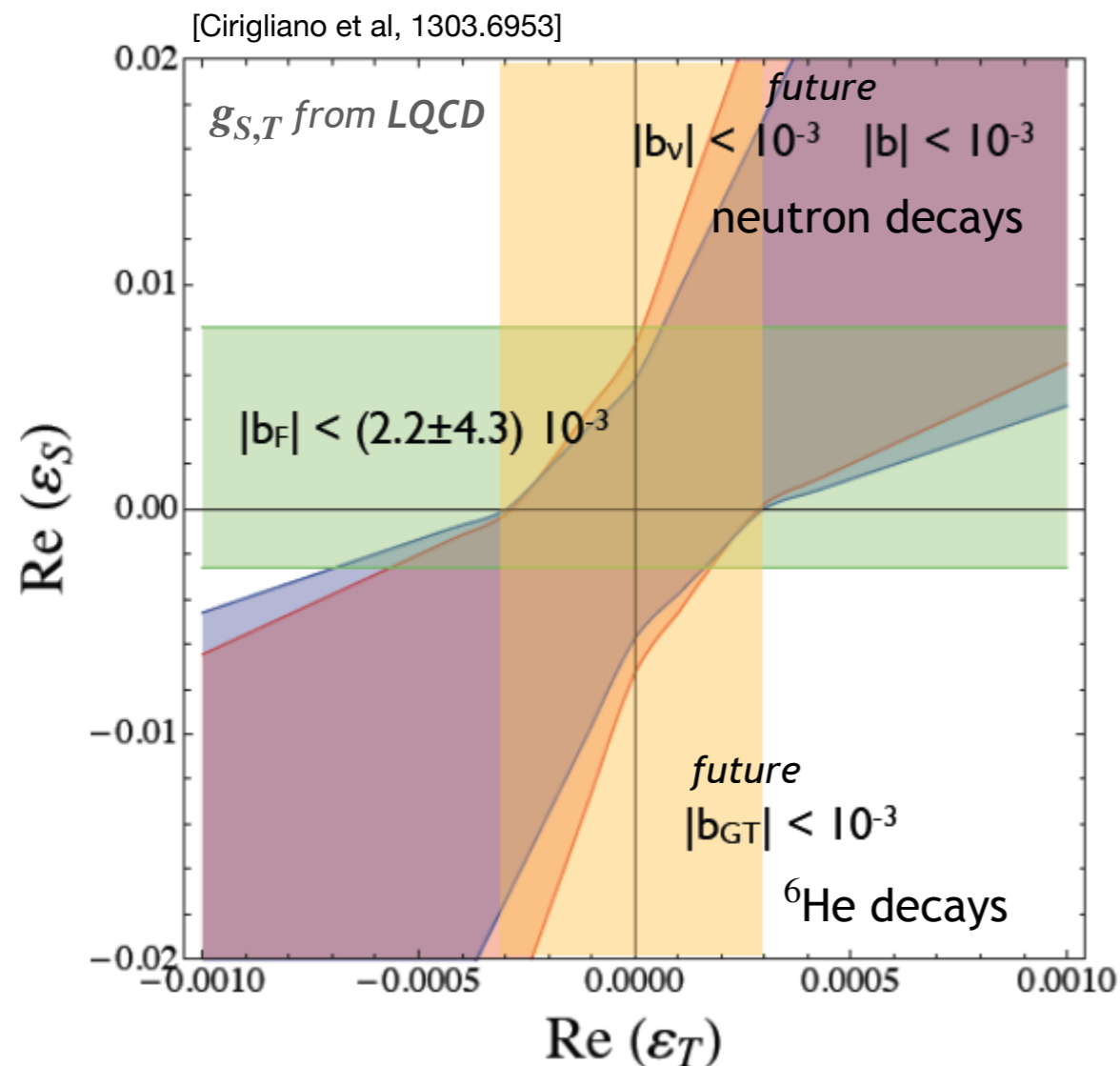
$$\tilde{a} = \frac{a}{1 + \alpha(E_e) b}$$

Asymmetries and angular correlations

► **SM precision tests in nuclear β decays:** a vast array of additional observables from the exploration of **angular correlations** between decay products!

► **Differential decay distributions** - very sensitive to the underlying Lorentz structure

⇒ searches for **non $(V - A)$** components in weak interactions!



A **summary** of important constraints:

ϵ_S - **superallowed decays**

ϵ_T - Dalitz plot study of $\pi^+ \rightarrow e^+ \nu_e \gamma$

Bounds from **other observables** (ϵ_S, ϵ_T)

^{60}Co measurements of A_{GT}

$$|g_T \Re \epsilon_T| < 10^{-2} \text{ (similar from } ^{114}\text{In decays)}$$

Long. polarisation of photon P_F/P_{GT}

$$|g_S \Re \epsilon_S + 4g_A/g_V g_T \Re \epsilon_T| < 10^{-2}$$

Long. polarisation of e^+ from polarised ^{107}In

$$|g_T \Re \epsilon_T| < 3 \times 10^{-3}$$

... ..

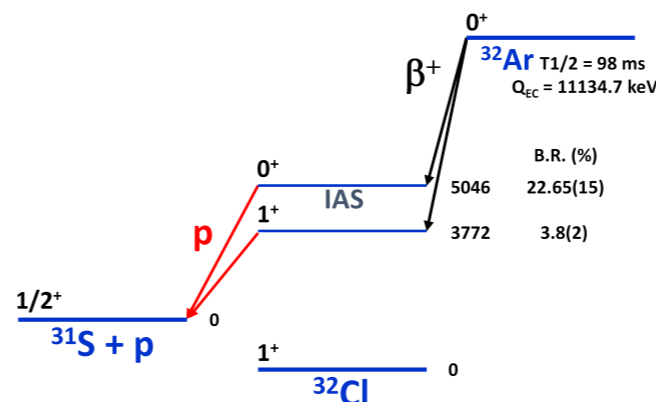
Recent comprehensive "global fit" [Falkowski et al, 2010.13797]

- ▶ **SM precision tests** in **nuclear β decays**: a vast array of additional observables from the exploration of **angular correlations** between **decay products**!
- ▶ **Recoil spectroscopy** offers many interesting and (**powerful**) features - access to both **a** and **b**
 - ⇒ **direct measurement** of daughter nucleus **recoil**
 - ⇒ **kinematic shifts** in energy spectrum of secondary (energetic) emitted particles
 - determination** of (unstable) **daughter momentum**
 - ⇒ simultaneous **study of multiple decay transitions** ;
 - precise determination** of $\tilde{a} \rightsquigarrow \delta\tilde{a}_{F(GT)} \sim 5 \times 10^{-3} (3 \times 10^{-3})$

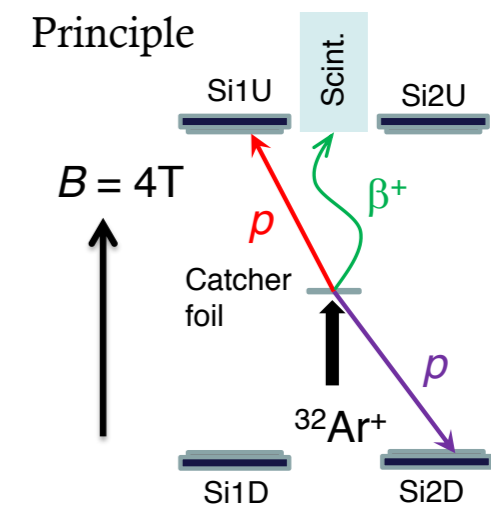
$$\tilde{a} = \frac{a}{1 + \alpha(E_e)b}$$

WISArD: Fermi (and Gamow-Teller) transitions in ^{32}Ar
simultaneous determination of \tilde{a}_F and \tilde{a}_{GT}

$$(\delta\tilde{a}_F < 10^{-3})$$



[Araujo-Escalona et al, 1906.05135]



Angular correlations & recoil spectroscopy

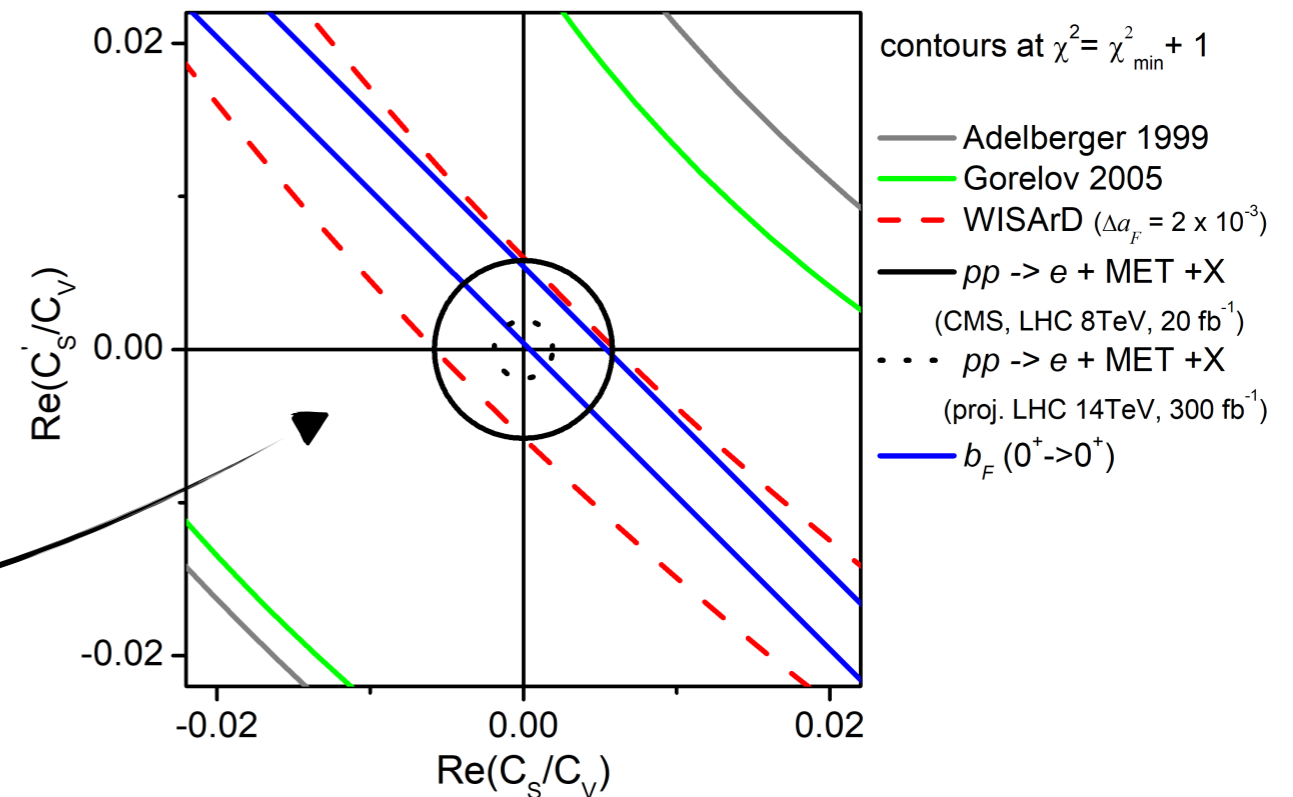
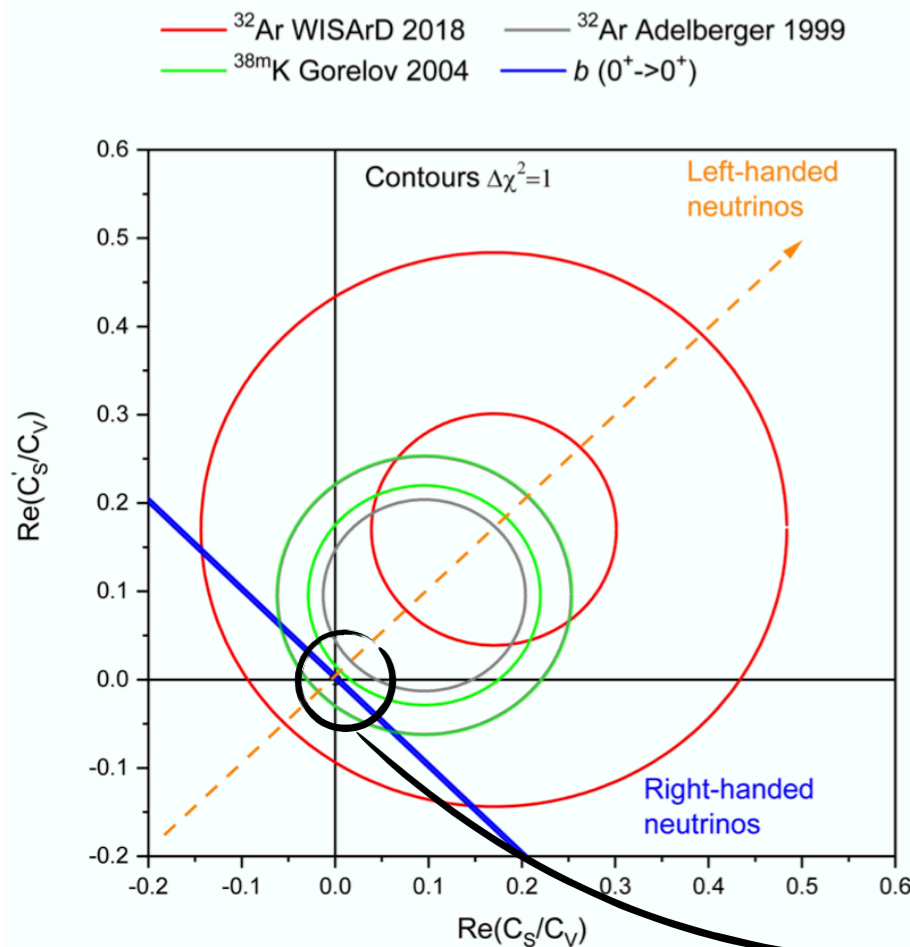
► **SM precision tests in nuclear β decays:** a vast array of additional observables from the exploration of **angular correlations** between decay products!

► **Differential decay distributions** - very sensitive to the underlying Lorentz structure

⇒ searches for **non $(V - A)$** components in weak interactions!

► **Recoil spectroscopy** offers access to both **a** and **b**

⇒ expect **significant (near) future progress**
precise determination of \tilde{a}



[WISArD@PANIC 2021]

[Araujo-Escalona et al, 1906.05135]

- ▶ **SM precision tests in nuclear β decays:** a vast array of additional observables from the exploration of **angular correlations** between decay products!
- ▶ **Shape of β energy spectrum** (upon integration of angular distributions)

$$W(E_e) dE_e = \frac{F(\pm Z, E_e)}{2\pi^3} p_e E_e (E_0 - E_e)^2 \xi \left(1 + b \frac{m_e}{E_e} \right) dE_e$$

β spectrum shape modified by **Fierz term!**

$$\text{For transparency } b_{\text{Fierz}} \approx \pm \frac{1}{1 + \rho^2} \left[\Re \left(\frac{C_S + C'_S}{C_V} \right) + \rho^2 \Re \left(\frac{C_T + C'_T}{C_A} \right) \right] \quad \rho = \frac{C_V}{C_A} \frac{M_{\text{GT}}}{M_F}$$

⇒ sensitivity to **BSM scalar** and **tensor** operators

$$b_{\text{Fierz}} \approx \pm \frac{2\gamma}{1 + \rho^2} \left[\Re \frac{g_S \epsilon_S}{g_V (1 + \epsilon_L + \epsilon'_R)} + \rho^2 \frac{4g_T \epsilon_T}{-g_A (1 + \epsilon_L - \epsilon'_R)} \right]$$

⇒ fully explore the **experimental precision** (and maximise sensitivity to NP)

include SM QCD form factors: dominant "weak electromagnetism" term

$$b \frac{m_e}{E_e} \rightarrow b \frac{m_e}{E_e} + \mathcal{O}(b_{\text{wm}} E_e / m_e)$$

[See, e.g., Severijns and Naviliat-Cuncic, '13;
Fenker et al, '16;]

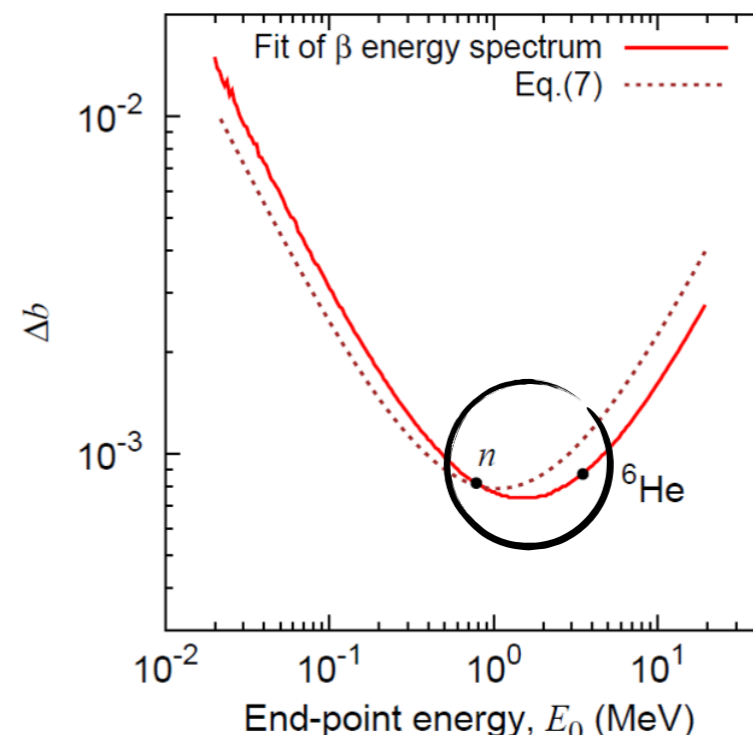
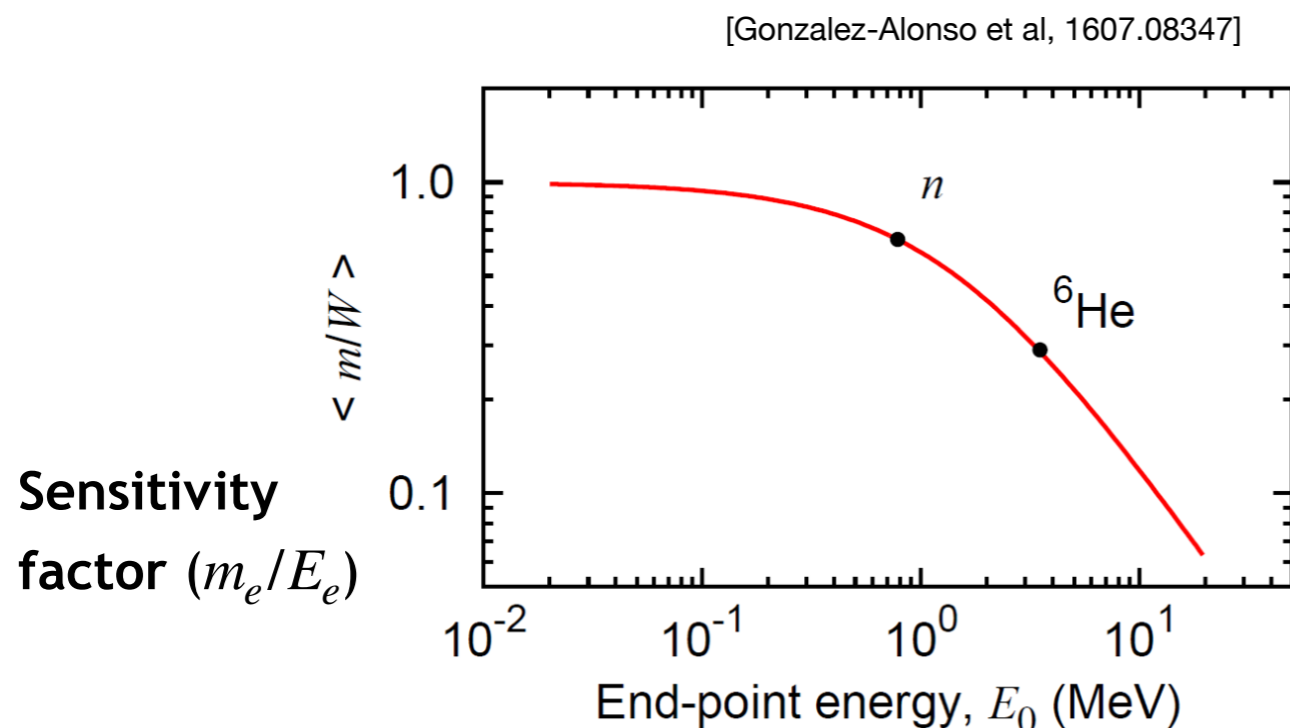
Shape of β spectrum

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β spectrum shape modified by **Fierz term!**

- ▶ **Maximal effects** (sensitivity to New Physics effects) for **endpoint energies** $\sim 1 - 2$ MeV (rapid decrease for smaller/larger values)



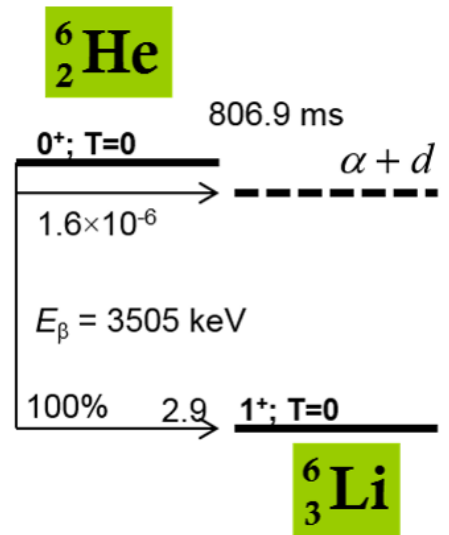
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► **Shape of β energy spectrum** (upon integration of angular distributions)

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β spectrum shape modified by **Fierz term!**



► **Excellent candidate ${}^6\text{He}$: pure Gamow-Teller transition to ground state**

Energy endpoint ~ 3.5 MeV

$$b_{GT} \propto g_T \mathcal{R} \epsilon_T$$

b-STILED:

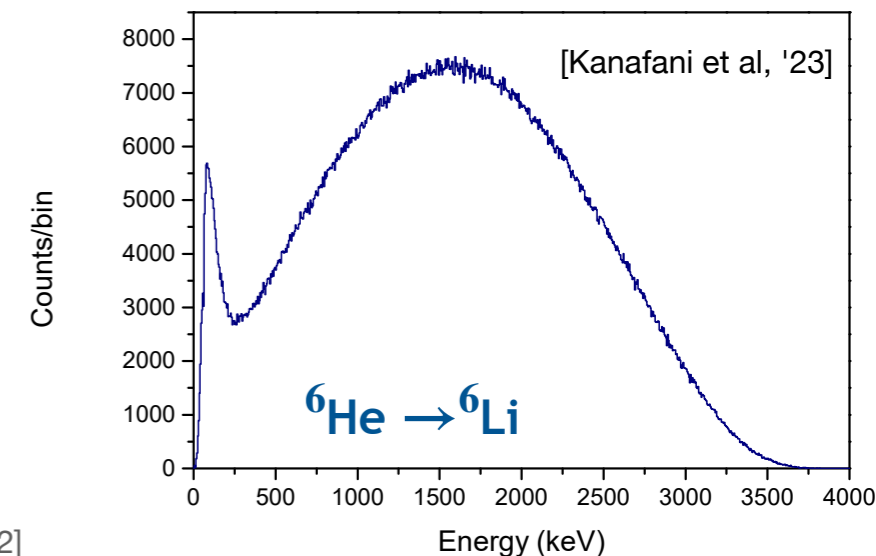
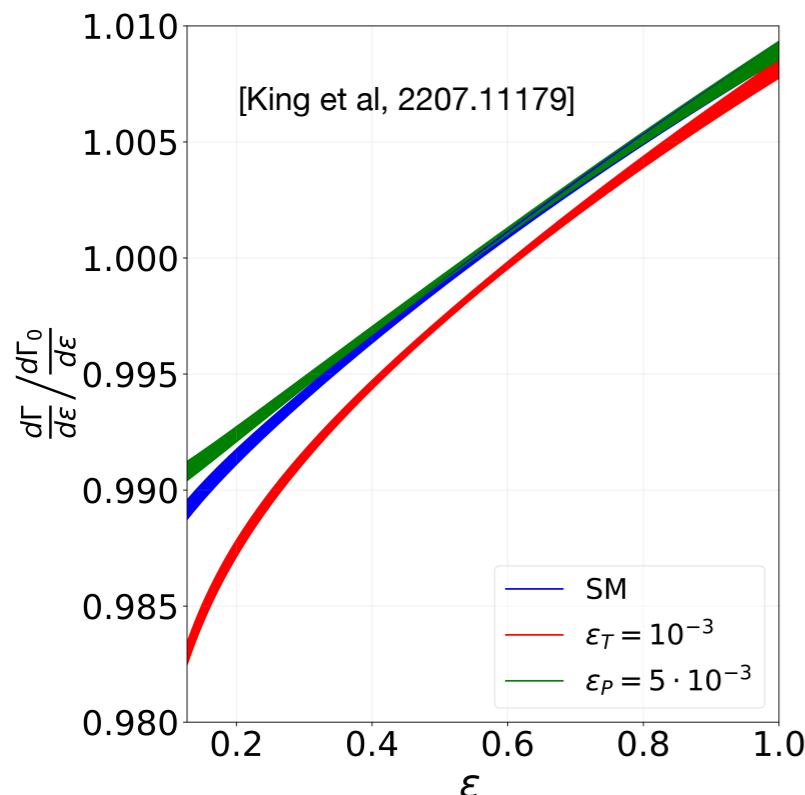
pure Gamow-Teller ${}^6\text{He} \rightarrow {}^6\text{Li}$

Expected **precision** on b_{GT}

$$\delta b_{GT} \sim 10^{-3}$$

"Ab-initio" calculations of ${}^6\text{He}$ beta decays for BSM !

[Glick-Magid et al, 2107.10212]



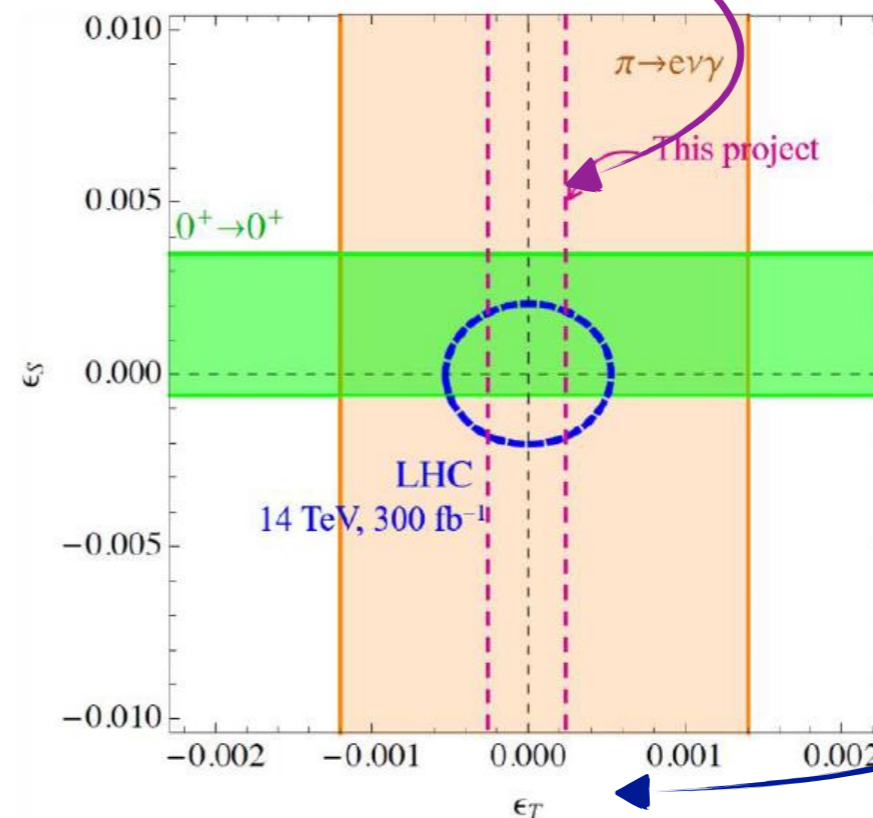
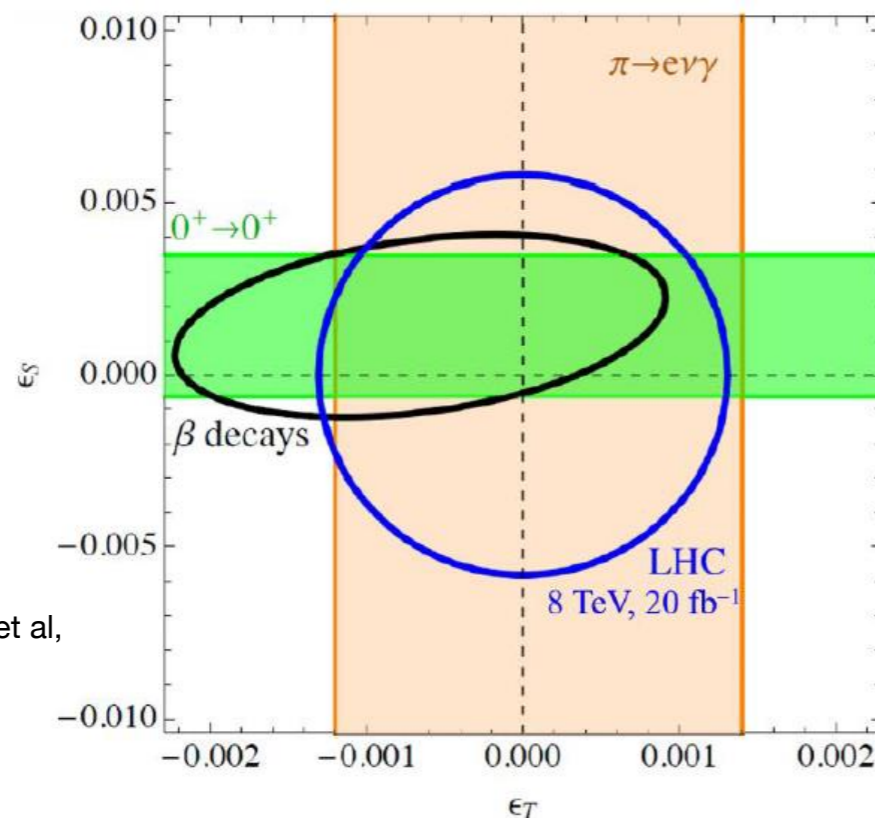
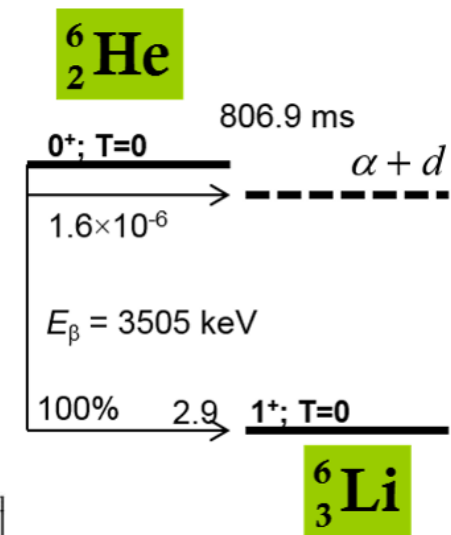
Shape of β spectrum

- ▶ **SM precision tests in nuclear β decays:** a vast array of additional observables from the exploration of **angular correlations** between decay products!
- ▶ **Shape of β energy spectrum** (upon integration of angular distributions)

$$W(E_e) dE_e = \frac{F(\pm Z, E_e)}{2\pi^3} p_e E_e (E_0 - E_e)^2 \xi \left(1 + b \frac{m_e}{E_e} \right) dE_e$$

β spectrum shape modified by Fierz term!

b-STILED: pure Gamow-Teller ${}^6\text{He} \rightarrow {}^6\text{Li}$ precision $\delta b_{\text{GT}} \sim 10^{-3}$



Constraints
on ϵ_T !

[Adapted from
Gonzalez-Alonso et al,
1803.08732]

Angular correlations and T-odd terms

- ▶ **SM precision tests** in **nuclear β decays**: a vast array of additional observables from the exploration of **angular correlations** between **decay products**!
- ▶ **Differential decay distributions** - very sensitive to the **underlying Lorentz structure** and to **new sources** of T-violation (if CPT \Rightarrow **CP violation**)

A correlation of **odd number of spins & momenta**: among other possibilities **D-correlation**

$$D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \quad \text{probed in mixed Fermi/Gamow-Teller transitions}$$

Contributions to **D**: **T-violating interactions** and **final state effects**

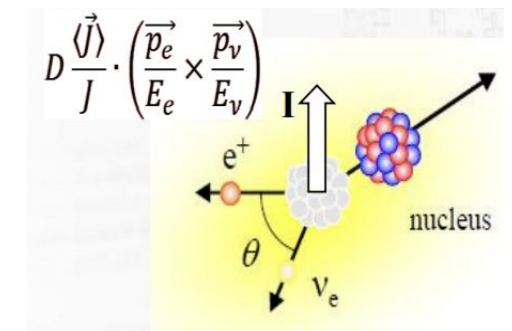
$$D = D_{TV} + D_{FSI}$$

$$D_{FSI} \approx D_1 \frac{p_e}{p_e^{\max}} + D_2 \frac{p_e^{\max}}{p_e} \quad \begin{aligned} D_1 &\sim \mathcal{O}(10^{-5 \div -4}) \\ D_2 &\sim \mathcal{O}(10^{-6 \div -5}) \end{aligned}$$

$$\text{Neutron: } D_{FSI} \approx 1.2 \times 10^{-5}$$

$$\delta \lesssim 1 \%$$

[Jackson et al, '57; Callan and S.B. Treiman, '67; Ando et al, 2009]



- ▶ **SM precision tests** in **nuclear β decays**: a vast array of additional observables from the exploration of **angular correlations** between **decay products**!
- ▶ **Differential decay distributions** - very sensitive to the **underlying Lorentz structure** and to **new sources** of T-violation (if CPT \Rightarrow **CP violation**)

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Contributions to **D**: **T-violating interactions** and final state effects

$$D = D_{TV} + D_{FSI}$$

$$D_{TV} \approx \frac{1}{1 + 3|\lambda|^2} \times \left[-2 \frac{\Im(C_V C_A^*)}{|C_V|^2} + \frac{\Im(C_S C_T^* + C_S' C_T'^*)}{|C_V|^2} + \frac{\alpha m_e}{p_e} \Re \left(\lambda^* \frac{C_T + C_T^*}{C_A^*} - \lambda^* \frac{C_S + C_S^*}{C_V^*} \right) \right]$$

$$D \approx \frac{4r g_V g_A}{g_V^2 + r^2 g_A^2} \sqrt{\frac{J}{J+1}} \Im \left[\epsilon_R (1 + \epsilon_L^*) + \frac{g_S g_T}{2g_V g_A} (\epsilon_S \epsilon_T^* + \tilde{\epsilon}_S \tilde{\epsilon}_T^*) - \tilde{\epsilon}_R \tilde{\epsilon}_L^* \right]$$

Angular correlations and T-odd terms

- ▶ **SM precision tests** in **nuclear β decays**: a vast array of additional observables from the exploration of **angular correlations** between **decay products**!
- ▶ **Differential decay distributions** - very sensitive to the **underlying Lorentz structure** and to **new sources** of T-violation (if CPT \Rightarrow **CP violation**)

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Contributions to **D**: **T-violating interactions** and final state effects

$$D = D_{TV} + D_{FSI}$$

**\Rightarrow Can nucleons and nuclei unveil new sources of CP violation?
Very much needed to overcome SM observational problems!**

$$D \approx \frac{4r g_V g_A}{g_V^2 + r^2 g_A^2} \sqrt{\frac{J}{J+1}} \Im \left[\epsilon_R (1 + \epsilon_L^*) + \frac{g_S g_T}{2g_V g_A} (\epsilon_S \epsilon_T^* + \tilde{\epsilon}_S \tilde{\epsilon}_T^*) - \tilde{\epsilon}_R \tilde{\epsilon}_L^* \right]$$

- ▶ **SM precision tests in nuclear β decays:** a vast array of additional observables from the exploration of **angular correlations** between **decay products!**
- ▶ **Differential decay distributions** - very sensitive to the **underlying Lorentz structure** and to **new sources** of T-violation (if CPT \Rightarrow **CP violation**)
And we **do need new CPV sources!**

So far, experimental searches in **only two systems** - neutron and ^{19}Ne ($J = 1/2$)

[Calaprice et al, '85] $^{19}\text{Ne} \rightsquigarrow D = 0.0001(6)$

[PDG, '20] **neutron** $\rightsquigarrow D_n = -0.00012(20)$ (world average)

... consistent with **absence** of **new sources of CPV** in **exotic scalar and tensor interactions**

Under the assumptions of **no new CPV**, constrain $\phi_{AV} = 180.012^\circ \pm 0.028^\circ$

[Chupp et al, 1205.6588]

$$\phi_{AV} = \arg \lambda \equiv \arg C_A/C_V$$

- ▶ **New experimental directions & sensitivity to NP:** β decays have different sensitivities to ϕ_{AV}

$$\text{Sensitivity} \rightsquigarrow D_{n,N} = F(n,N) \sin \phi_{AV}$$

\Rightarrow **Maximise D_N and polarisation degree !**

Angular correlations and T-odd terms

- ▶ **SM precision tests in nuclear β decays:** a vast array of additional observables from the exploration of **angular correlations** between **decay products!**
- ▶ **Differential decay distributions** - very sensitive to the **underlying Lorentz structure** and to **new sources** of T-violation (if CPT \Rightarrow **CP violation**)
Only two systems explored: neutron and ^{19}Ne ($J = 1/2$)
- ▶ **New experimental directions & sensitivity to NP \Rightarrow Maximise D_N and polarisation degree !**

	n	^{19}Ne	^{23}Mg	^{35}Ar	^{39}Ca
Sensitivity $F(X)$	0,43	-0,52	-0,65	0,41	0,71
D_1 ($\times 10^{-4}$)	0,108	2,326	1,904	0,386	-0,489
D_2 ($\times 10^{-4}$)	0,023	0,169	0,099	0,010	-0,024

[Delahaye et al - MORA project
1802.02970]



First focus: ^{23}Mg (easier to polarise...) \rightsquigarrow future ^{39}Ca

Expected precision $\delta D_{\text{Mg}} \sim \text{few} \times 10^{-5}$

Sensitivity to ϕ_{AV} : $\phi_{AV} = (1.6 \pm 6.3) \times 10^{-4} \rightarrow \mathcal{O}(10^{-5})$

Synergy with other searches of New Physics sources of CP?

- ▶ If **new sources of CP violation** are present - as required to explain the BAU generically expect **contributions to a vast array of CP-odd observables** from LHC, to **meson-decay observables** (asymmetries), ..., **EDMs** (elementary, nucleon, atomic...) and **nuclear decays!**
- ▶ Synergy between **EDMs** and D_{TV} - a very naïve first approach

Consider **T-violating dim-6 term**

$$\mathcal{L}_{\text{SMEFT}}^{\text{eff}} \supset i\mathcal{C}_{Hud}^6 \tilde{H}^\dagger D_\mu H (u^c \sigma^\mu \bar{d}^c) + \text{H.c.}$$

$$\mathcal{L}_{\text{WEFT}}^{\text{eff}} \supset -\frac{2V_{ud}}{v^2} \left[(\bar{e} \bar{\sigma}_\mu \nu) (\bar{u} \bar{\sigma}_\mu d) + \frac{v^2}{2V_{ud}} \mathcal{C}_{Hud}^{11} (\bar{e} \bar{\sigma}_\mu \nu) (u^c \sigma_\mu \bar{d}^c) \right]$$

At nucleon level $\Rightarrow D_n \approx \frac{4g_V g_A}{g_V^2 + g_A^2} \Im \epsilon_R \approx 0.4v^2 \Im \mathcal{C}_{Hud}^{11}$

But **nEDM** constrains $|v^2 \Im \mathcal{C}_{Hud}^{11}| \leq 6 \times 10^{-6}$

$\Rightarrow D_n \lesssim 2 \times 10^{-6} ??$

(Recall however that cancellations might occur and final implications are model-dependent)

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Likewise one also finds $|d_n| \approx 1 \times 10^{-19} e \cdot \text{cm} |D_{TV}/\kappa|$

$$|d_{\text{Hg}}| \approx 7 \times 10^{-24} e \cdot \text{cm} |D_{TV}/\kappa|$$

[Ng and Tulin, 1111.0649]

$$|d_D| \approx 4.5 \times 10^{-20} e \cdot \text{cm} |D_{TV}/\kappa|$$

$$\kappa = \frac{4g_V g_A M_F M_{GT}}{g_V^2 M_F^2 + g_A^2 M_{GT}^2} \sqrt{\frac{J}{J+1}}$$

Bounds generically apply to **classes of UV-models** leading to **distinct operators in \mathcal{L}^{eff}**

Scenario	ν WEFT	ν SMEFT	order D	max $ D $
Ia	ϵ_R	$HD_\mu H u^c \sigma^\mu \bar{d}^c$	Λ^{-2}	$\mathcal{O}(10^{-6})$
Ib	ϵ_R	$(\bar{l} H \bar{\sigma}_\mu H l)(u^c \sigma^\mu \bar{d}^c)$	Λ^{-4}	$\mathcal{O}(10^{-4}) \frac{v^2}{\Lambda^2}$
II	ϵ_S, ϵ_T	$(\bar{l} \bar{\sigma}_{\mu\nu} \bar{e}^c)(\bar{q} \bar{\sigma}^{\mu\nu} \bar{u}^c), (\bar{l} \bar{e}^c)(\bar{q} \bar{u}^c), (\bar{l} \bar{e}^c)(\bar{d}^c q)$	Λ^{-4}	$\mathcal{O}(10^{-14})$
III	$\tilde{\epsilon}_S, \tilde{\epsilon}_T$	$(\bar{l} \bar{\sigma}^{\mu\nu} \bar{\nu}^c)(\bar{q} \bar{\sigma}_{\mu\nu} \bar{d}^c), (\bar{l} \bar{\nu}^c)(\bar{q} \bar{d}^c), (\bar{l} \bar{\nu}^c)(u^c q)$	Λ^{-4}	$\mathcal{O}(10^{-6})$
IVa	$\tilde{\epsilon}_L, \tilde{\epsilon}_R$	$H^\dagger D_\mu H^\dagger e^c \sigma^\mu \bar{\nu}^c, (e^c \sigma^\mu \bar{\nu}^c)(u^c \sigma_\mu \bar{d}^c)$	Λ^{-4}	$\mathcal{O}(10^{-4}) \frac{v^2}{\Lambda^2}$
IVb	$\tilde{\epsilon}_L, \tilde{\epsilon}_R$	$e^c \sigma^\mu \bar{\nu}^c \bar{q} H^\dagger \sigma_\mu H^\dagger q, (e^c \sigma^\mu \bar{\nu}^c)(u^c \sigma_\mu \bar{d}^c)$	Λ^{-6}	$\mathcal{O}(10^{-4}) \frac{v^2}{\Lambda^2}$

[Falkowski and Rodriguez-Sanchez, 2207.02161]

(Recall however that cancellations might occur and final implications are model-dependent)

- ▶ If **new sources of CP violation** are present - as required to explain the BAU generically expect **contributions to a vast array of CP-odd observables** from LHC, to **meson-decay observables** (asymmetries), ..., **EDMs** (elementary, nucleon, atomic...) and **nuclear decays!**
- ▶ **Synergy** between **EDMs** and D_{TV} - a very naïve first approach \Rightarrow **strong constraints on D_{TV}** (*Recall however that cancellations might occur and final implications are model-dependent*)
- ▶ A **strong case** for the **precise determination** of D_{TV} in nuclear beta decays:
 - EDM** measurements offer stronger bounds on new **CP violation** sources
 - However, a *single EDM measurement has little discriminating power*
 - Especially in the advent of **EDM** observation, more **independent observables required!**
 - Sensitivity of D_{TV} to exotic currents might help **untangling nature of CPV!**



► SM precision tests in nuclear β decays: LHC probing the same parton level processes

Contributions of common set of EFT operators to $pp \rightarrow e + \text{MET}(+X)$

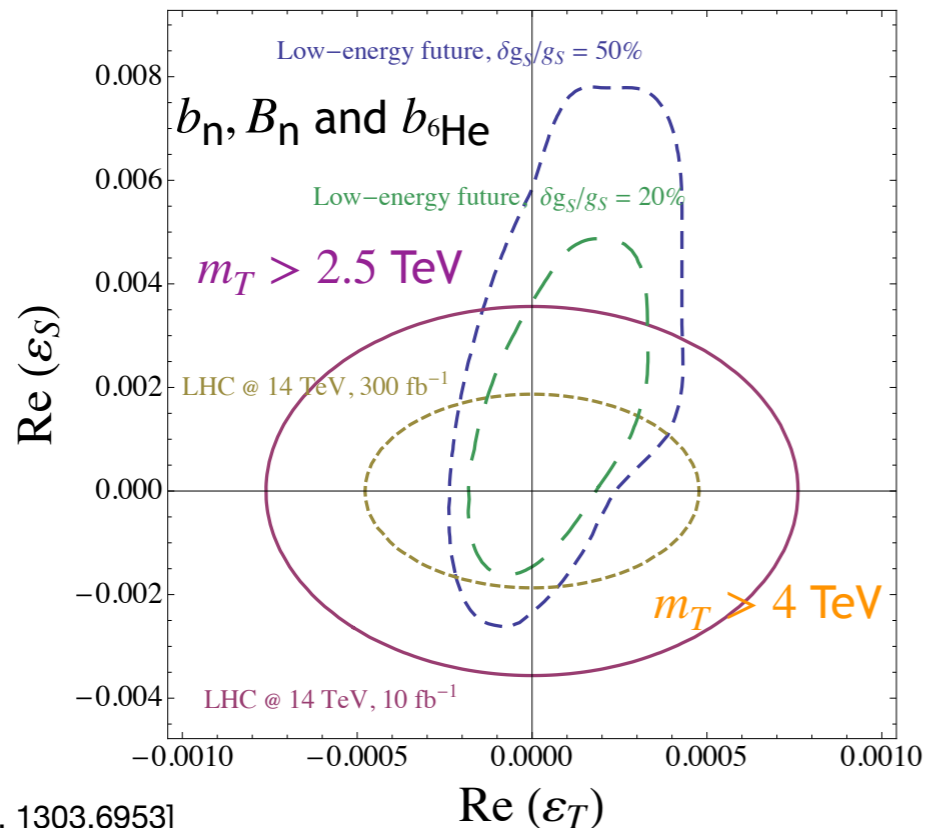
Unsuppressed effects at collider energies $\Rightarrow \mathcal{O}(s^2/v^4)$ enhancement

If $m_T \equiv \sqrt{2E_T^e E_T^\nu(1 - \cos \Delta\phi_{e\nu})}$ larger than threshold production

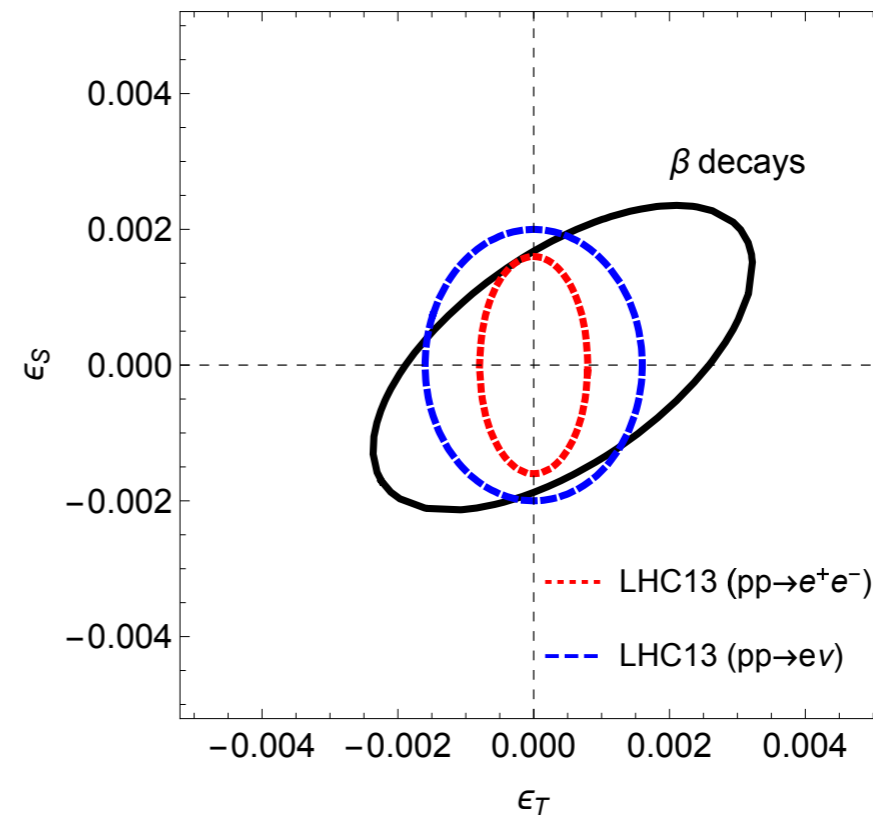
$$\sigma(m_T > m_{\bar{T}}) = \sigma_W \left[(1 + \epsilon_L^{(v)})^2 + |\tilde{\epsilon}_L|^2 + |\epsilon_R|^2 \right] - 2\sigma_{WL} \Re(\epsilon_L^{(c)} + \epsilon_L^{(c)} \epsilon_L^{(v)*})$$

[See, e.g., Cirigliano et al, 1210.4553]

$$+ \sigma_R \left[|\tilde{\epsilon}_R|^2 + |\epsilon_L^{(c)}|^2 \right] + \sigma_S \left[|\epsilon_S|^2 + |\tilde{\epsilon}_S|^2 + |\epsilon_P|^2 + |\tilde{\epsilon}_P|^2 \right] + \sigma_T \left[|\epsilon_T|^2 + |\tilde{\epsilon}_T|^2 \right]$$



[Cirigliano et al, 1303.6953]



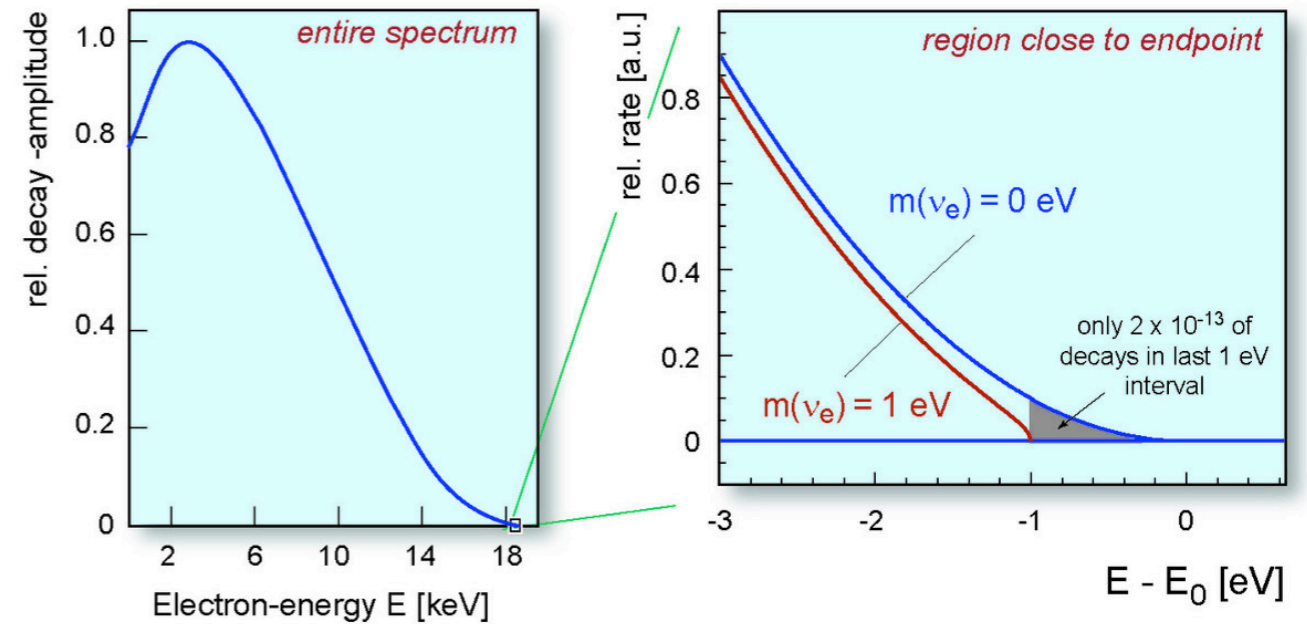
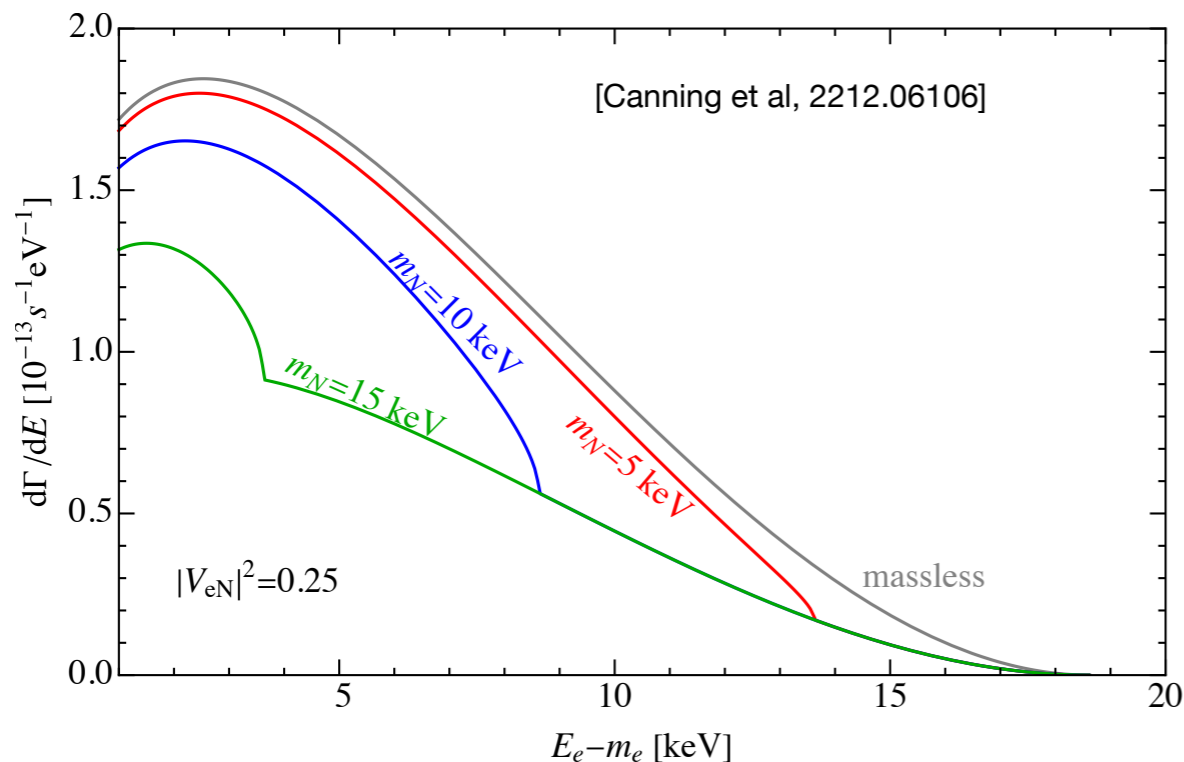
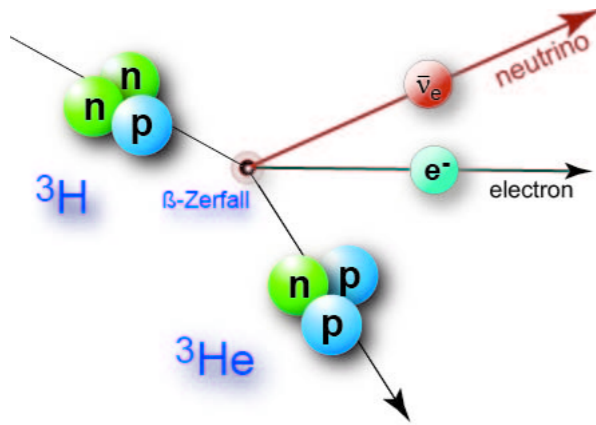
[Falkowski et al, 2010.13797]

Nuclear decays and neutrino physics

► Historically, a treasure trove for **neutrinos**: from **discovery**, to the characterisation of **fundamental properties** (mass, nature, ...) & searches for **new (sterile) states**

► Highlights: **Tritium decays** ${}^3\text{H} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_e$
measure the **shape of beta spectrum**, at its endpoint

⇒ determine **neutrino mass** $m_\beta < 1.1 \text{ eV}$; $m_\beta \sim 0.2 \text{ eV}$
(little sensitivity to exotic currents)



⇒ **search for light sterile states**
(non-negligible mixings)
Kurie plots, differential decay rates

▶ Historically, a treasure trove for **neutrinos**: from **discovery**, to the characterisation of **fundamental properties** (mass, nature, ...) & searches for **new (sterile) states**

▶ **Highlights**: massive neutrinos and the role of **sterile neutrinos**

Modified charged current interactions (PMNS matrix) $\nu_\alpha = U_{\alpha i}^{\text{PMNS}} \nu_i$

$${}^3\text{H} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_e \Rightarrow m_\beta \rightarrow \left(\sum_k |U_{ek}|^2 m_{\nu_k}^2 \right)^{1/2} \quad \& \text{ kinks on the beta spectrum!}$$

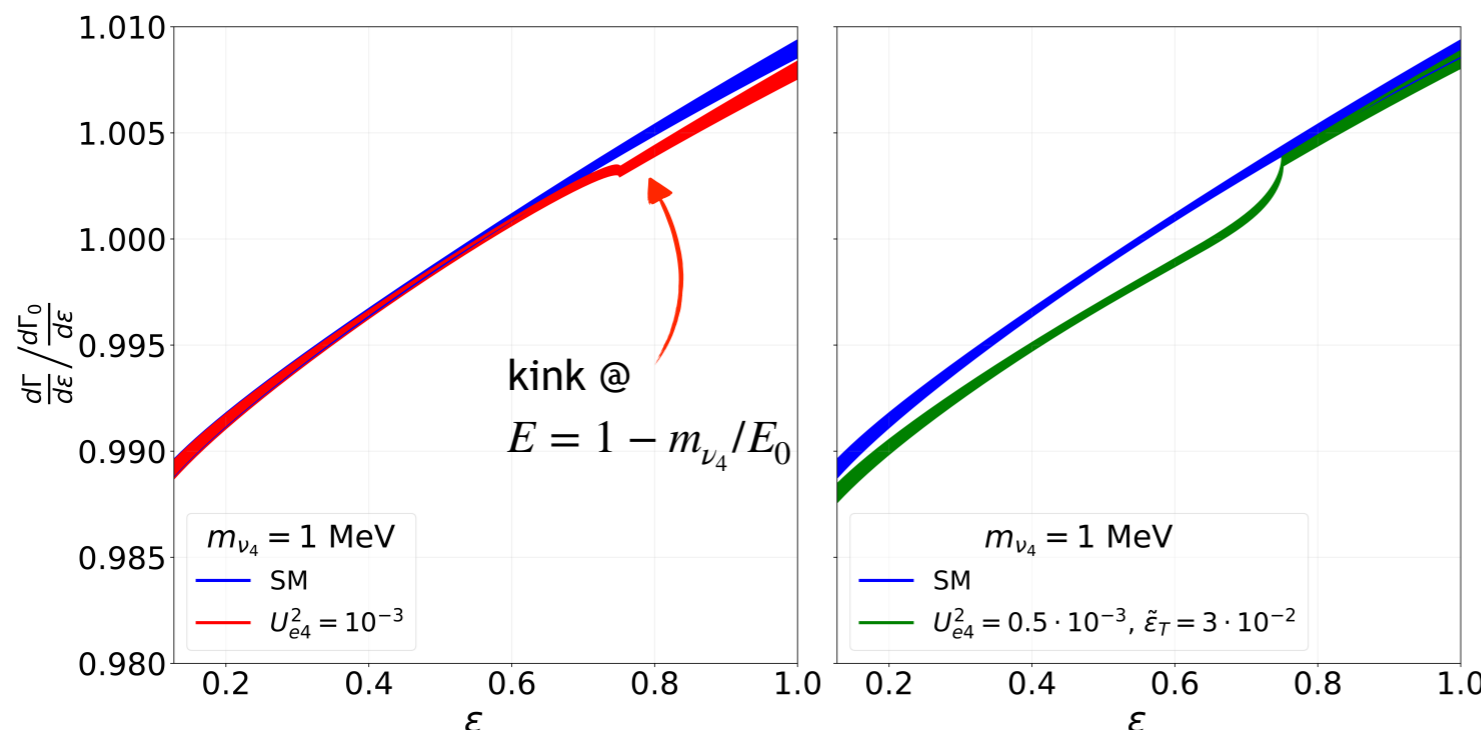
If **sterile states decoupled**, or if **negligibly light** ($m_{\nu_s} \ll M_N$) **neglect ν_s operators**

If $m_{\nu_s} \approx \mathcal{O}(m_e)$ **new terms** in \mathcal{L}^{eff} , e.g. $\propto V_{ud} U_{e4} (1 + \epsilon_L) (\bar{e}_L \gamma_\mu \nu_4) (\bar{u}_L \gamma^\mu d_L) +$

$$V_{ud} U_{e4} \epsilon_T (\bar{u}_R \sigma_{\mu\nu} d_L) (\bar{e}_R \sigma^{\mu\nu} \nu_4) + \dots$$

⇒ **Corrections to the β spectrum**

[King et al, 2207.11197]



- ▶ **SM precision tests in nuclear β decays:** a possible path to NP discovery!

Improvement in numerous experimental fronts \Rightarrow and Λ_{NP} sensitivities

Complementary to LHC direct searches and to other high-intensity probes

- ▶ **Identifying tensions between theoretical prediction and expectation**

excellent experimental precision, and reduction of theoretical uncertainties

On the **theory** side, *many mountains to climb*:

computation of **nuclear charges**, reduce **approximations** in $V_{us}(\mathcal{F}t)$,

proper inclusion of **small effects** (critical for SM-like observables) -

radiative corrections, induced hadronic form factors, ...

\Rightarrow Excellent news from **first principle approaches**: **LQCD** and **ab-initio nuclear computations**

- ▶ *Discrepancy between nuclear observable and SM prediction ...*

\Rightarrow Identifying the **UV particle physics model at work!**

\Rightarrow Relate **nucleonic** form factors with **quark level (NP) charges!**

Understanding the nature of antimatter & ultimate tests of the weak equivalence principle



- ▶ "Immediate" **Universe** is composed of **matter**: electrons, protons, neutrons, nuclei and atoms...
antimatter only in cosmic rays, (certain) radioactive decays or laboratory produced
- ▶ Matter and antimatter in the **SM**: **identical elementary particles**, with **opposite CP charges** (or strictly identical in the case of *Majorana neutral fermions*!)
- ▶ **SM gauge interactions** only distinguish matter and antimatter via "**sign**" of charges ($\pm |\lambda|$)
What about **mass**? **No kinematical difference** between electrons and positrons...
 p and \bar{p} charge-to-mass ratio: identical (precision of 16 parts in a trillion!)
[BASE Collaboration, 2022]
- ▶ What about **gravitation**?
Observation suggests that **gravity effects** on the motion of neutral antimatter (\bar{H})
$$\bar{g} = (0.75 \pm 0.13_{\text{sys+stat}} \pm 0.16_{\text{sim}}) g$$
[ALPHA-g Collaboration, 2023]
⇒ **consistent with a downward gravitational acceleration (1 g) for antihydrogen**

Is there still room for **non-standard gravitational interactions**? What would be the impact?
(for gravity theories and particle physics!)

Weak Equivalence Principle

► One principle to challenge (or confirm!):

"In a uniform gravitational field, **all objects** (regardless of nature and composition) free-fall with precisely the **same acceleration**"

Newtonian interpretation \Rightarrow identity of **inertial** and **gravitational masses** ($m^g \equiv m^I$)

Asymmetry in gravitational interactions: challenge universality of free-fall acceleration

$$g_{(i)} = m^g / m_{(i)}^I$$

$$\eta_{ij}^\oplus = \frac{\Delta g^\oplus}{\langle g^\oplus \rangle} = \frac{g_i^\oplus - g_j^\oplus}{(g_i^\oplus + g_j^\oplus)/2} \quad (\text{Eötvös parameter})$$

Extensive **tests of WEP violation** for bodies of different compositions falling in the Earth's field: variations on free torsion pendulum experiments to constrain η^\oplus between **pairs of matter elements**



Experiment	Test bodies	Measurement
Eöt-Wash	Be - Ti	$\eta_{\oplus, \text{Be-Ti}} = (0.3 \pm 1.8) \times 10^{-13}$
Eöt-Wash	Be - Al	$\eta_{\oplus, \text{Be-Al}} = (-1.5 \pm 1.5) \times 10^{-13}$
Eöt-Wash	Be - Cu	$\eta_{\oplus, \text{Be-Cu}} = (-1.9 \pm 2.5) \times 10^{-12}$

$$\eta_{\text{E-W}}^\oplus \lesssim \mathcal{O}(10^{-12})$$

[See, e.g., Wagner et al, 1207.2442]

MICROSCOPE (satellite experiment): $\eta_{\text{Ti-Pl}}^\oplus \lesssim \mathcal{O}(10^{-15})$

[MICROSCOPE Collaboration, 2209.15487]

[Torsion pendulum used in Eöt-Wash experiments]

Having **antimatter** gravitating in a distinct way corresponds to the more general possibility that **different forms of energy** gravitate differently...

Two broad classes of theoretical possibilities

(i) **modified General Relativity**

(ii) **new forces**, mediated by vectors and/or scalars, (sub)gravitational strength

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(i) **Modified gravitation** (subatomic scales) \leadsto *impact fundamental principles of GR and QFT*
(**CPT** invariance, **Lorentz** invariance, ...)

CPT: invariance of **local Lorentz-invariant QFTs** (point-like particles) - e.g. QED, SM, ...

Does this hold for more fundamental theories, upon combining the SM and gravity?

In **string theory** one can have **spontaneous breaking of CPT and Lorentz...**

Can one look for **laboratory signals of CPT & Lorentz violation** (at the Planck scale)?

Exceptionally sensitive experiments required...

\Rightarrow Hydrogen and antihydrogen spectroscopy!

[See, for instance, Charlton et al, 2002.09348]

Antimatter and modified gravity

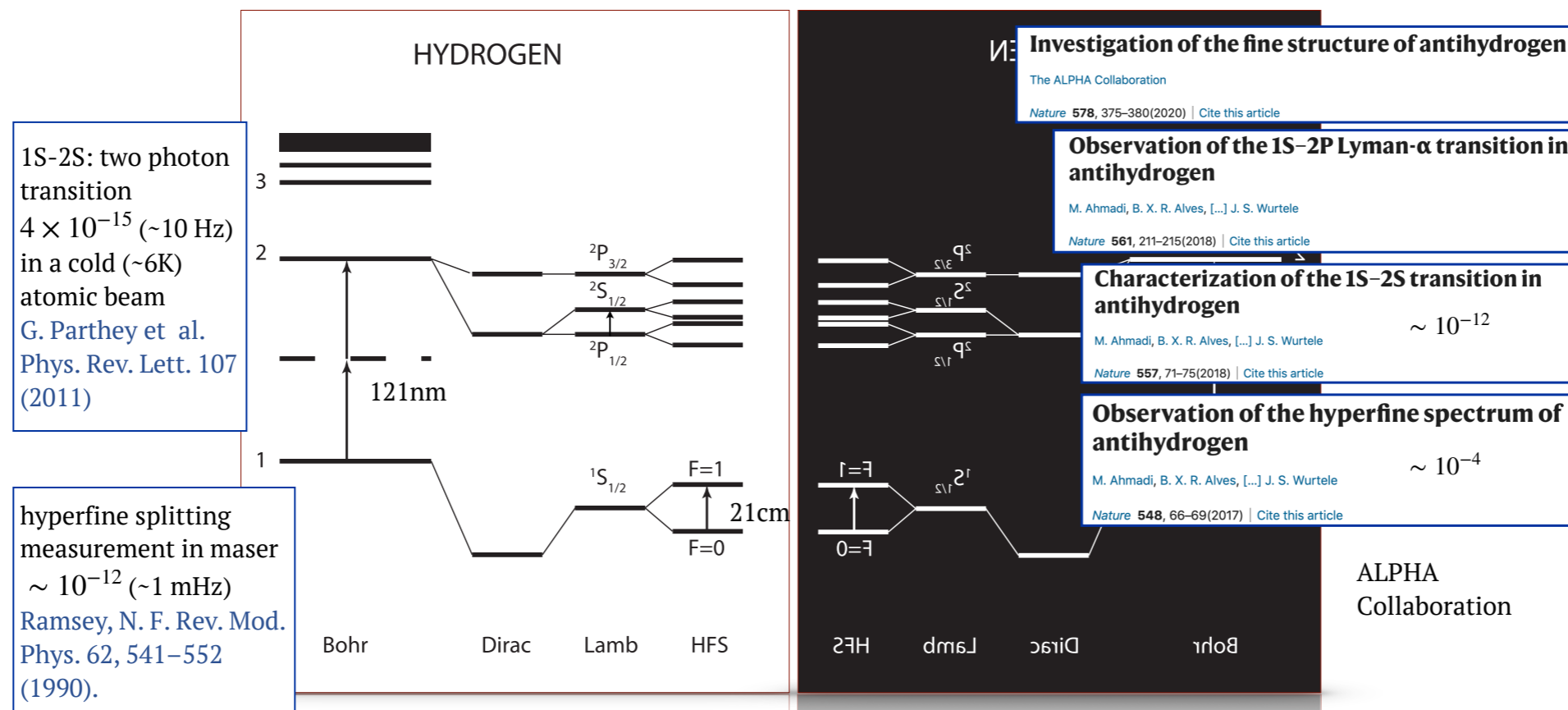
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Extraordinary progress in recent years!



[From Malbrunot, CERN'21]

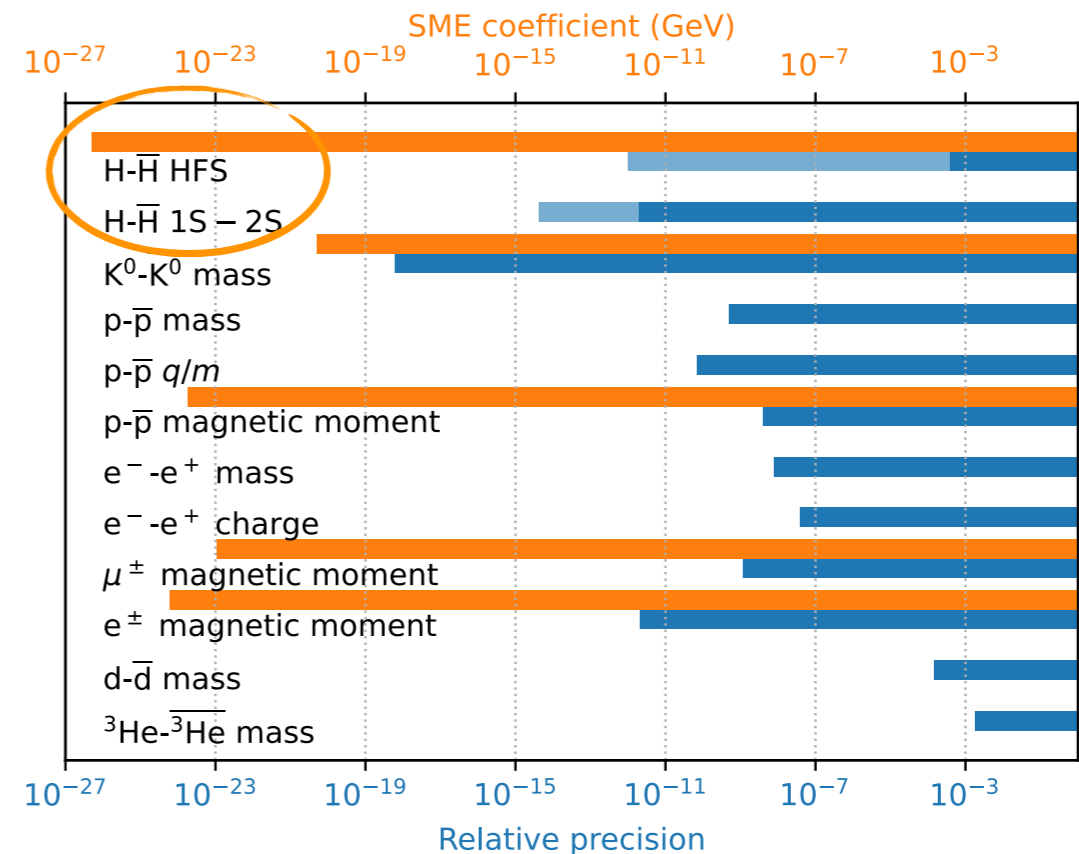
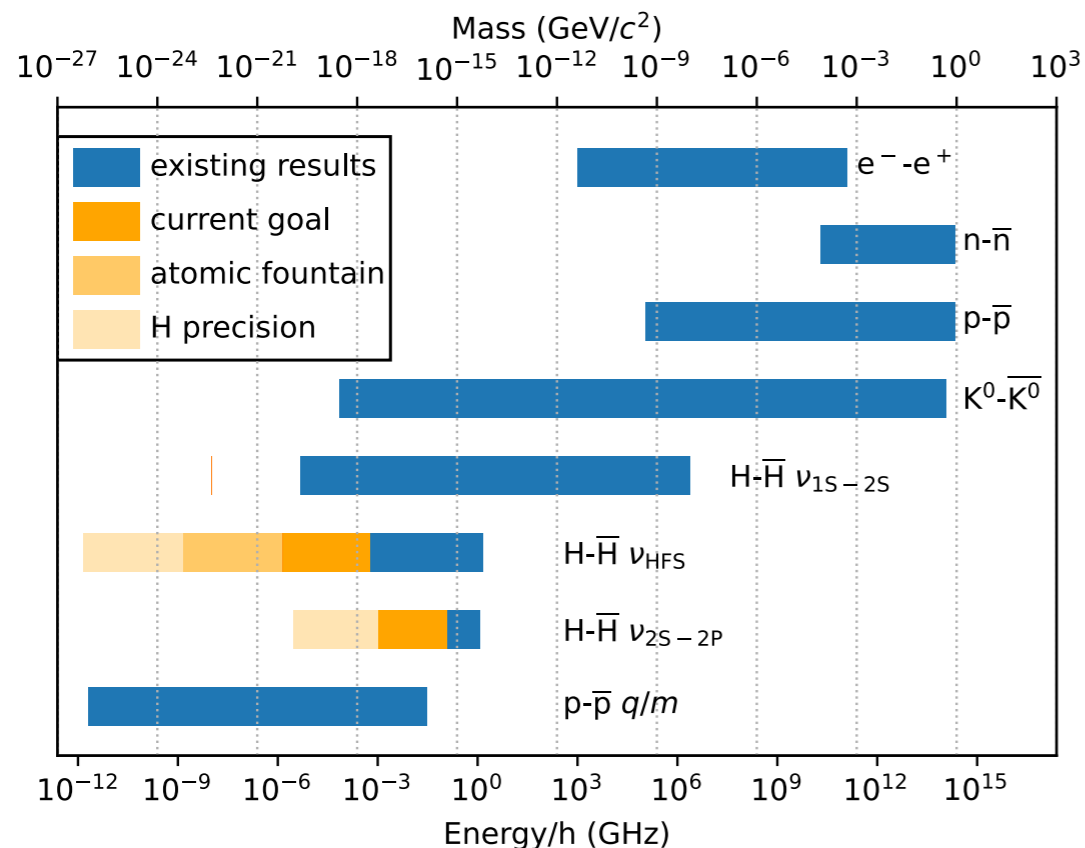
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Extraordinary progress in recent years! Impressive results for **CPT violation tests!**



[Widmann et al, 2111.04056]

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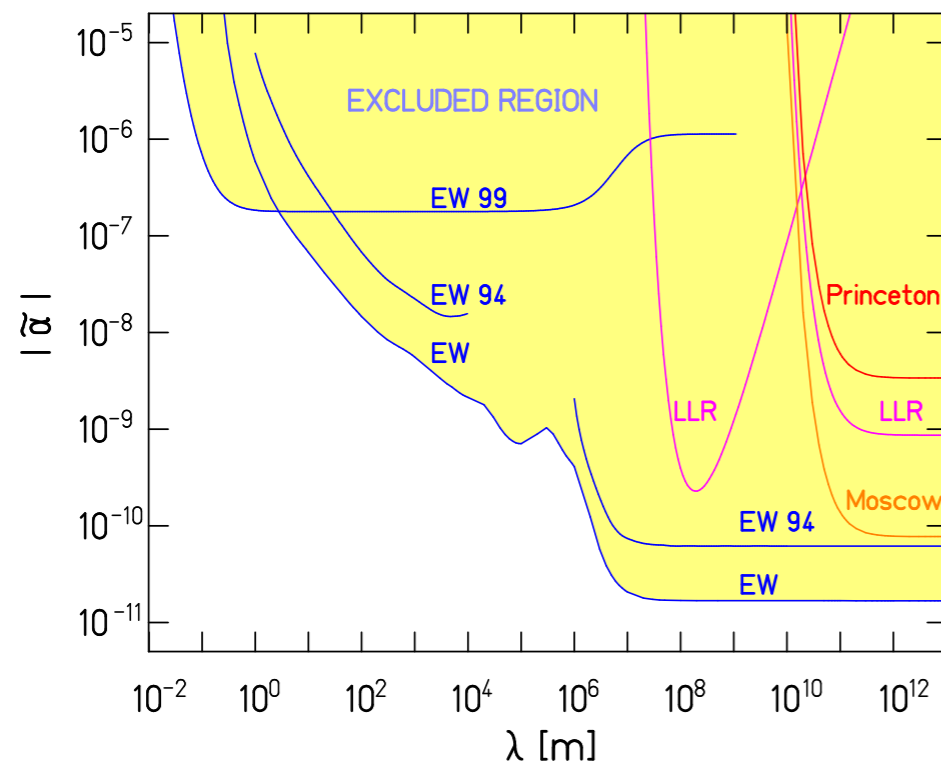
Two broad classes of theoretical possibilities

- (i) modified General Relativity
- (ii) **new forces**, mediated by vectors and/or scalars, **(sub)gravitational strength**

A new (**5th force**) interaction: potentially much weaker than gravity

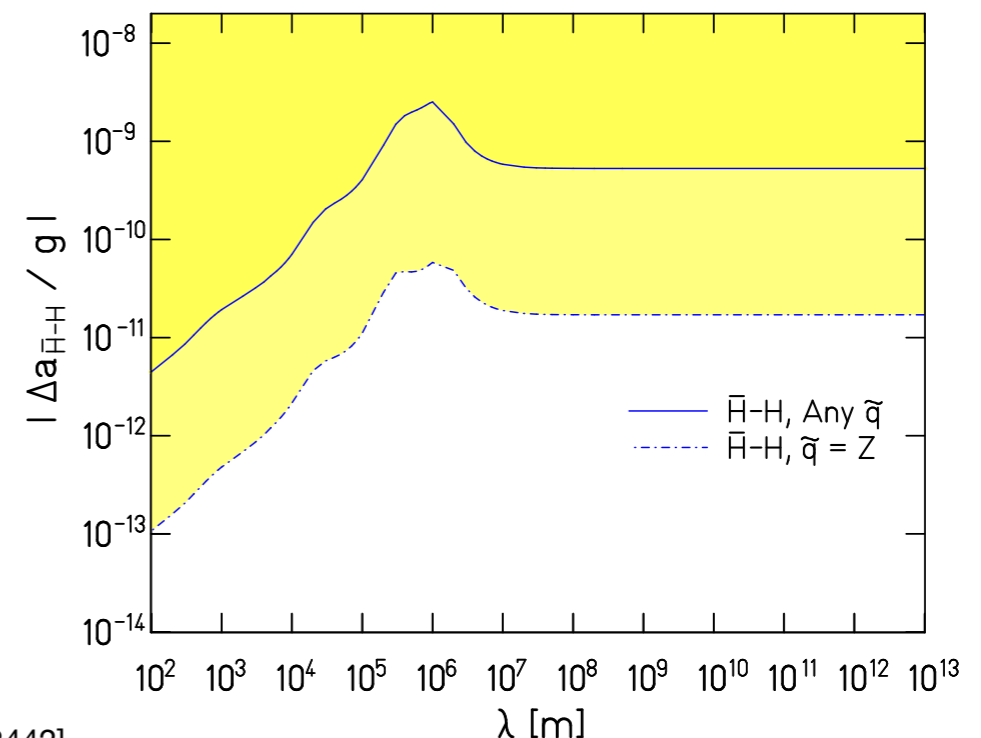
mediated by **bosons** (M_B), **strength** \tilde{g} , coupling to **fermions** (\tilde{q})

bounds on $\tilde{\alpha}$ for new vector interaction
(with $\tilde{q} = B - L$) vs. range



[Wagner et al, 1207.2442]

bounds on $(g_H - g_{\bar{H}})/g_H$



Thorough characterisation of antimatter



IN2P3
Les deux infinis

► Explaining a pressing **observational problem** of the SM:

the matter-antimatter asymmetry of the Universe

The SM offers a *strikingly simple description of antimatter*: **Charge-Parity** transformation

Can **antimatter** couple differently to **gravity**? Free-fall in Earth's gravitational field

How to study matter made of **anti-constituents**? **Anti-hydrogen** is the best object to consider

⇒ thorough studies (spectroscopy, ...) with **the best available precision!**

⇒ **\bar{H} free-fall** in Earth's gravitational field and direct measure of **\bar{g} !**

Recent results: $\bar{g} = (0.75 \pm 0.13_{\text{sys+stat}} \pm 0.16_{\text{sim}})g$

[ALPHA Collaboration, 2023]

Improve precision to ascertain $g = \bar{g} \Rightarrow$ **GBAR: below 1%!**



A first step in **understanding gravity effects in nuclear and elementary interactions**:

Couplings to virtual pairs? Couplings to binding energies? Flavour content of valence q ?

⇒ Aim at **precision well below 10^{-2} !**

From **antihydrogen** to antimatter - a long road ahead! **Muonium** as a next stop?

Concluding remarks and theory prospects



► **New Physics** paths to discovery at three frontiers

Precision tests of the SM offer uniquely **promising prospects**

► Explored here several fronts:

cLFV transitions in the **muon sector** \Rightarrow **muon-electron conversion** offers an amazing **probing power to NP**

EDMs in the quest for **new sources of CPV** \Rightarrow **neutron EDM remarkably competitive**

EW precision tests in beta-decays \Rightarrow so many observables to explore, offering a joint **probing power of NP sources of CPV and of NP interactions**
strong synergy with direct LHC searches and EDMs!

Precision tests of WEK - antihydrogen and **new gravitational interactions!**

► **Very strong experimental prospects!**

Theory must reduce its uncertainties to be on par with experimental precision!

EFT is an extremely powerful tool \leadsto explore (UV) **models of New Physics!**

Outlook and perspectives



IN2P3
Les deux infinis

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Thank you for the attention!

Additional material

