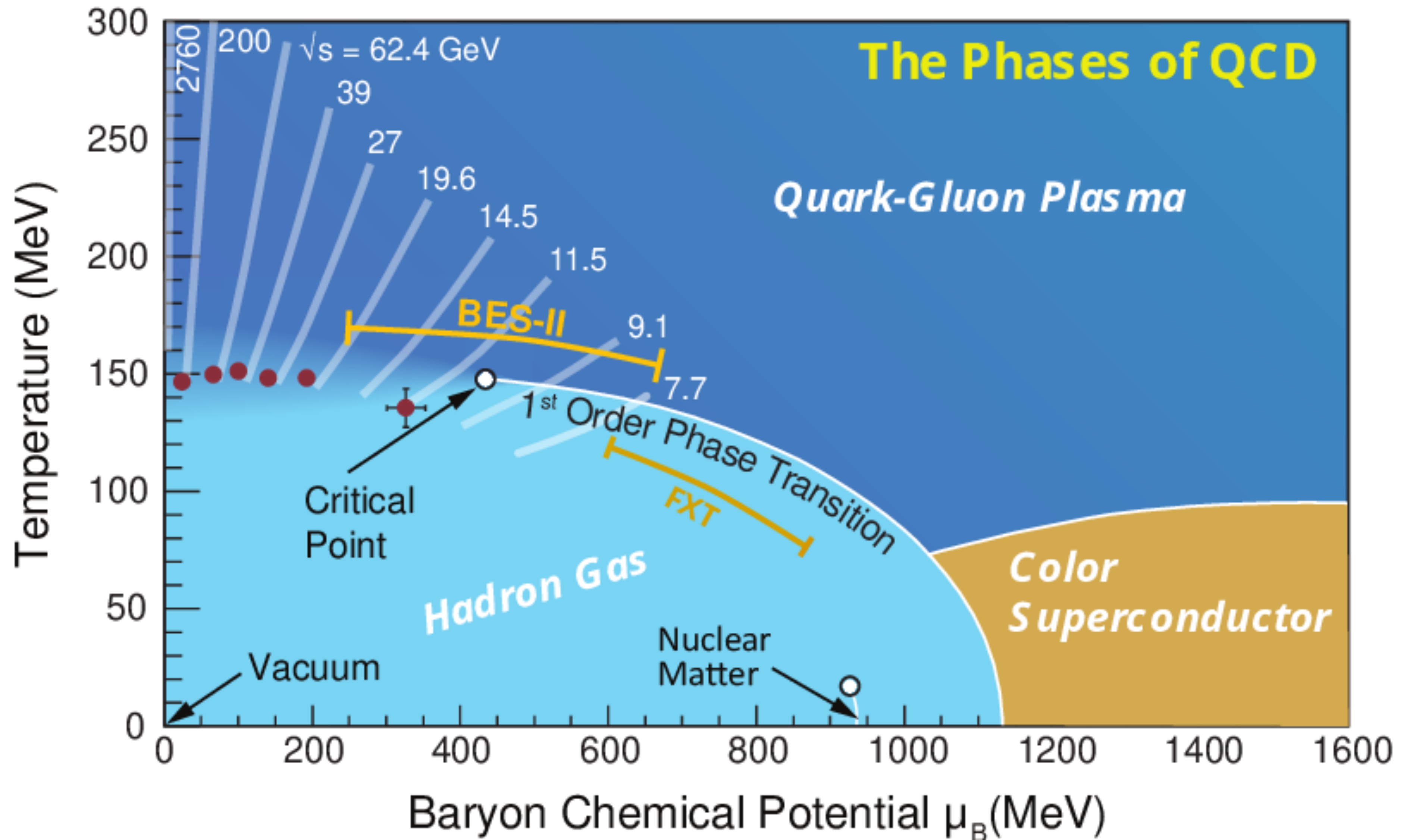


Bulk evolution of linearized fluctuations

Jakub Štěřba, CTU in Prague, Subatech

Introduction



Fluctuations in general

- In viscous relativistic fluid dynamics
- Quantum fluctuations
 - Initial fluctuations
- Thermal fluctuations
 - related to susceptibilities and EoS \rightarrow phase structure of QCD
 - fluctuation-dissipation relations

Dynamical fluctuations

- Stochastic
 - discretized noise sampled event-by-event
 - observables calculated after statistical averaging
- Hydro-kinetics
 - deterministic kinetic equations for the two-point functions of fluid dynamical fields
 - linearization of stochastic fluid dynamics
- Critical fluctuations - inclusion of non-linearities

Stochastic fluctuations

$$\begin{aligned} \partial_\mu T^{\mu\nu} &= 0, & T^{\mu\nu} &= T^{\mu\nu}_{ideal} + T^{\mu\nu}_{viscous} + \Xi^{\mu\nu} \\ \partial_\mu J^\mu &= 0, & J^\mu &= J^\mu_{ideal} + J^\mu_{viscous} + I^\mu_{noise} \end{aligned} \quad \langle u^\gamma \partial_{;\gamma} \pi^{\mu\nu} \rangle = -\frac{\pi^{\mu\nu} - \pi_{NS}^{\mu\nu}}{\tau_\pi} - \frac{4}{3} \pi^{\mu\nu} \partial_{;\gamma} u^\gamma$$

- From fluctuation-dissipation relation

$$\partial_t \Xi^{ij} = -\frac{1}{\tau_\pi} (\Xi^{ij} - \xi^{ij})$$

- Correlator of the noise

$$\langle \xi^{\mu\nu}(x) \xi^{\alpha\beta}(x') \rangle = \left[2\eta T (\Delta^{\mu\alpha} \Delta^{\nu\beta} + \Delta^{\mu\beta} \Delta^{\nu\alpha}) + 2 \left(\zeta - \frac{2}{3} \eta \right) T \Delta^{\mu\nu} \Delta^{\alpha\beta} \right] \delta^4(x - x')$$

- Discretization of delta function leads to

- Lattice spacing dependence
- Large noise contributions can locally lead to negative densities

Linearized equations

- Introducing a perturbation to hydro equations

$$\partial_\nu \left(T_0^{\mu\nu} + \delta T^{\mu\nu} \right) = 0$$

- Rewriting the equations as

$$\partial_0 Q^\mu + \partial_i \vec{F}^\mu = 0, \text{ where } Q^\mu = T^{0\mu} \text{ and } F^{i\mu} = T^{i\mu}$$

- Decoupling for background and perturbations

$$\begin{aligned} \partial_0 Q_0^\mu + \partial_i \vec{F}_0^\mu &= 0 \\ \partial_0 \delta Q^\mu + \partial_i \delta \vec{F}^\mu &= 0 \end{aligned}$$

- perturbations have zero mean over the ensemble average

Linearized equations

- Introducing the perturbation to primitive variables

$$\varepsilon = \varepsilon_0 + \delta\varepsilon$$

$$p = p_0 + \delta p \quad \text{and} \quad \pi^{\mu\nu} = \pi_0^{\mu\nu} + \delta\pi^{\mu\nu}$$

$$u^\mu = u_0^\mu + \delta u^\mu$$

- We arrive at the set of equations

$$Q_0^\mu = (\varepsilon_0 + p_0)u_0^0 u_0^\mu - p_0 g^{\mu 0} + \pi_0^{\mu 0}$$

$$F_0^{\mu i} = (\varepsilon_0 + p_0)u_0^\mu u_0^i - p_0 g^{\mu i} + \pi_0^{\mu i}$$

$$\delta Q^\mu = (\varepsilon_0 + p_0)(u_0^\mu \delta u^0 + u_0^0 \delta u^\mu) + (\delta\varepsilon + \delta p)u_0^0 u_0^\mu - \delta p g^{\mu 0} + \delta\pi^{\mu 0}$$

$$\delta F^{\mu i} = (\varepsilon_0 + p_0)(u_0^\mu \delta u^i + u_0^i \delta u^\mu) + (\delta\varepsilon + \delta p)u_0^\mu u_0^i - \delta p g^{\mu i} + \delta\pi^{\mu i}$$

Linearized equations

- Linearizing the transport coefficients

$$\begin{aligned}\eta &= \eta_0 + \delta\eta \\ \zeta &= \zeta_0 + \delta\zeta \\ \tau_\pi &= \tau_{\pi 0} + \delta\tau_\pi \\ \tau_\Pi &= \tau_{\Pi 0} + \delta\tau_\Pi\end{aligned}$$

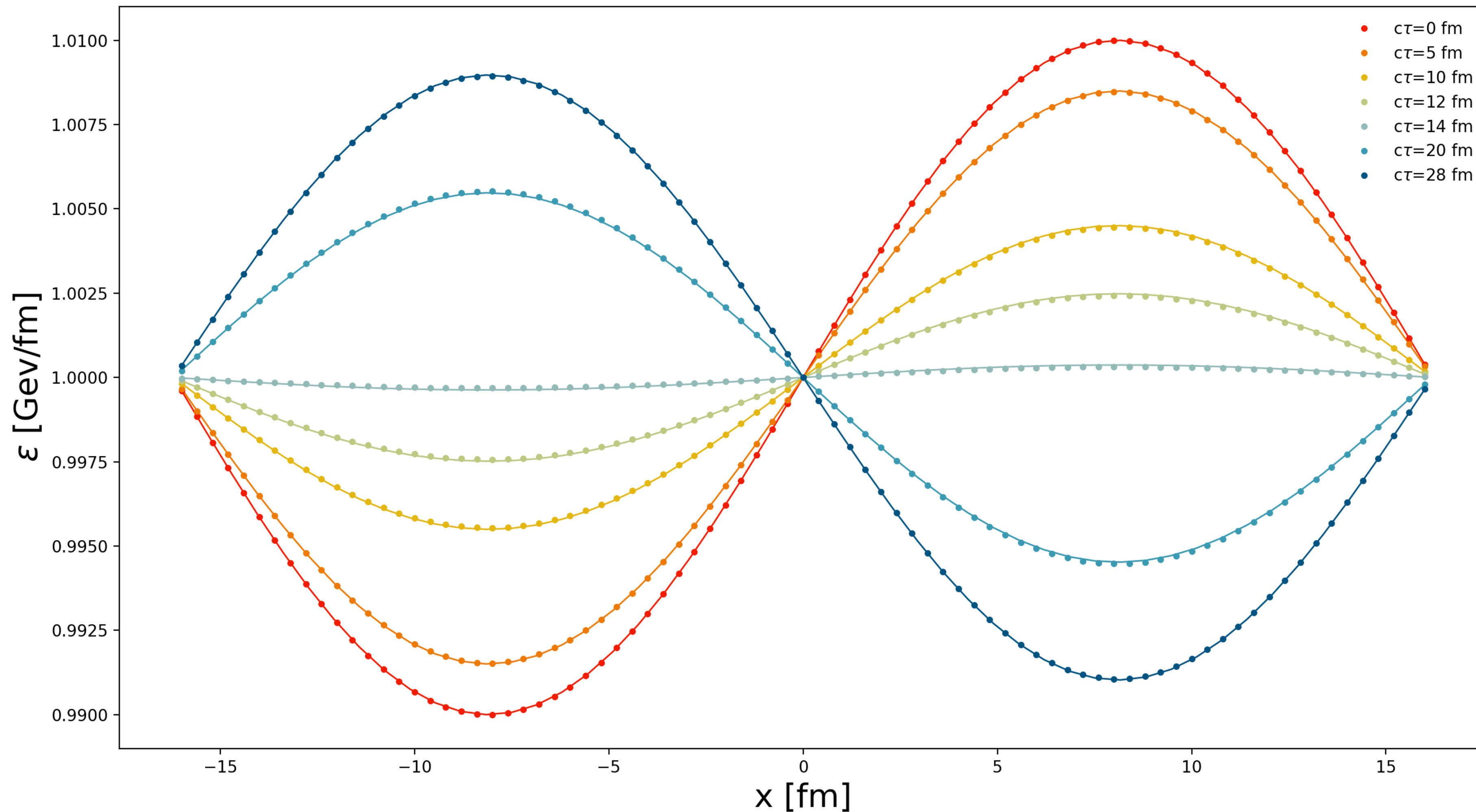
$$\langle \rangle = \frac{1}{2} \Delta_\alpha^\mu \Delta_\beta^\nu + \frac{1}{2} \Delta_\alpha^\nu \Delta_\beta^\mu - \frac{1}{3} \Delta^{\mu\nu} \Delta_{\alpha\beta} = \langle \rangle_0 + \langle \rangle_\delta$$

- We get Israel-Stewart equations

$$\begin{aligned}\langle u_0^\gamma \partial_{\delta;\gamma} \delta\pi^{\mu\nu} \rangle_0 + \langle \delta u^\gamma \partial_{0;\gamma} \pi_0^{\mu\nu} \rangle_0 + \langle u_0^\gamma \partial_{0;\gamma} \pi_0^{\mu\nu} \rangle_\delta = \\ = - \frac{\delta\pi^{\mu\nu} - \delta\pi_{NS}^{\mu\nu}}{\tau_{\pi 0}} - \frac{\pi_0^{\mu\nu} - \pi_{0NS}^{\mu\nu}}{\tau_{\pi 0}^2} \delta\tau_\pi - \frac{4}{3} (\pi_0^{\mu\nu} \partial_{\delta;\gamma} \delta u^\gamma + \delta\pi^{\mu\nu} \partial_{0;\gamma} u_0^\gamma)\end{aligned}$$

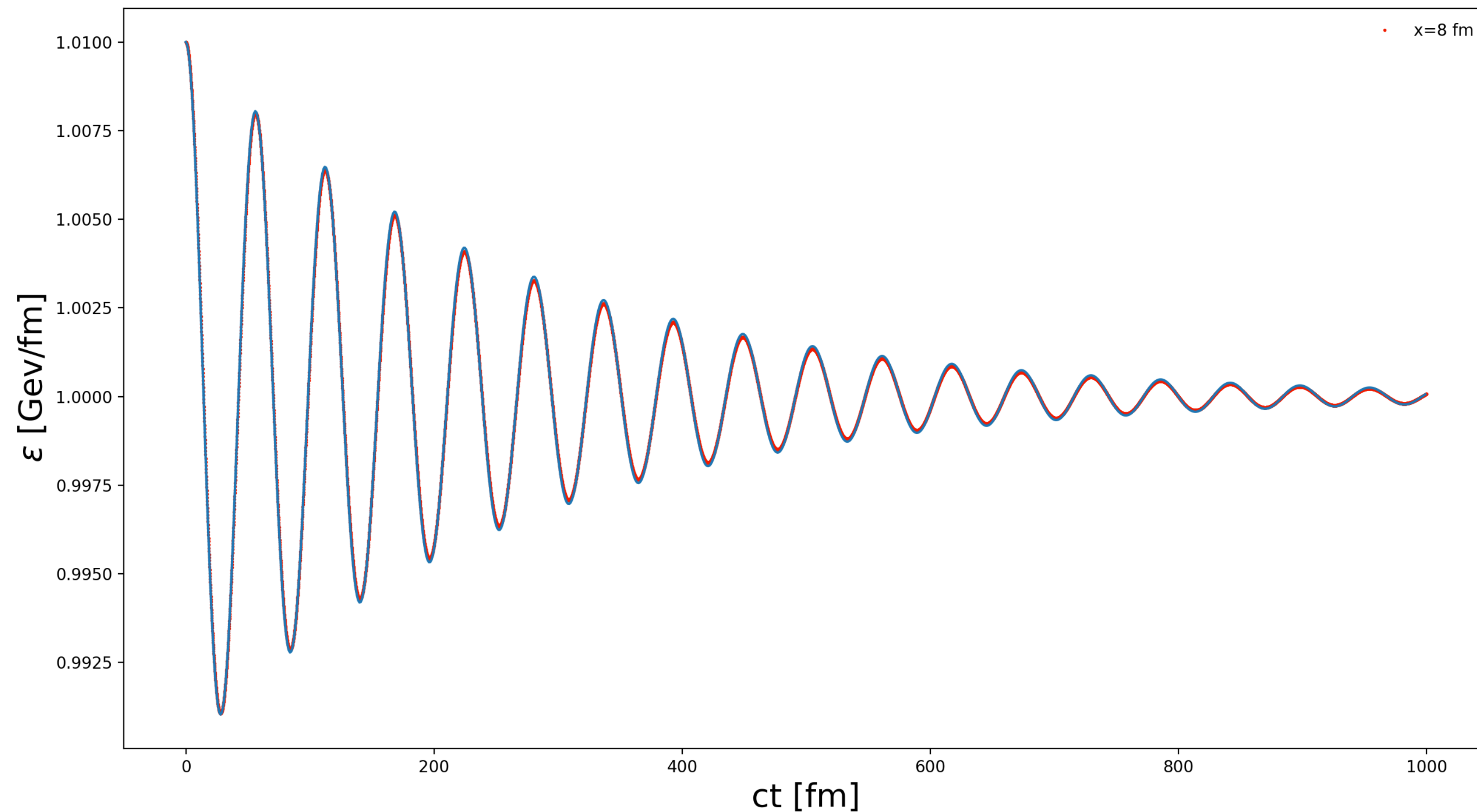
Test of linearized equations in box mode

- vHLE - periodic boundaries, Cartesian coordinates, static background, NS limit
- Perturbation of sinus wave $\varepsilon = 0.01 \sin\left(\frac{2\pi}{L}x\right)$



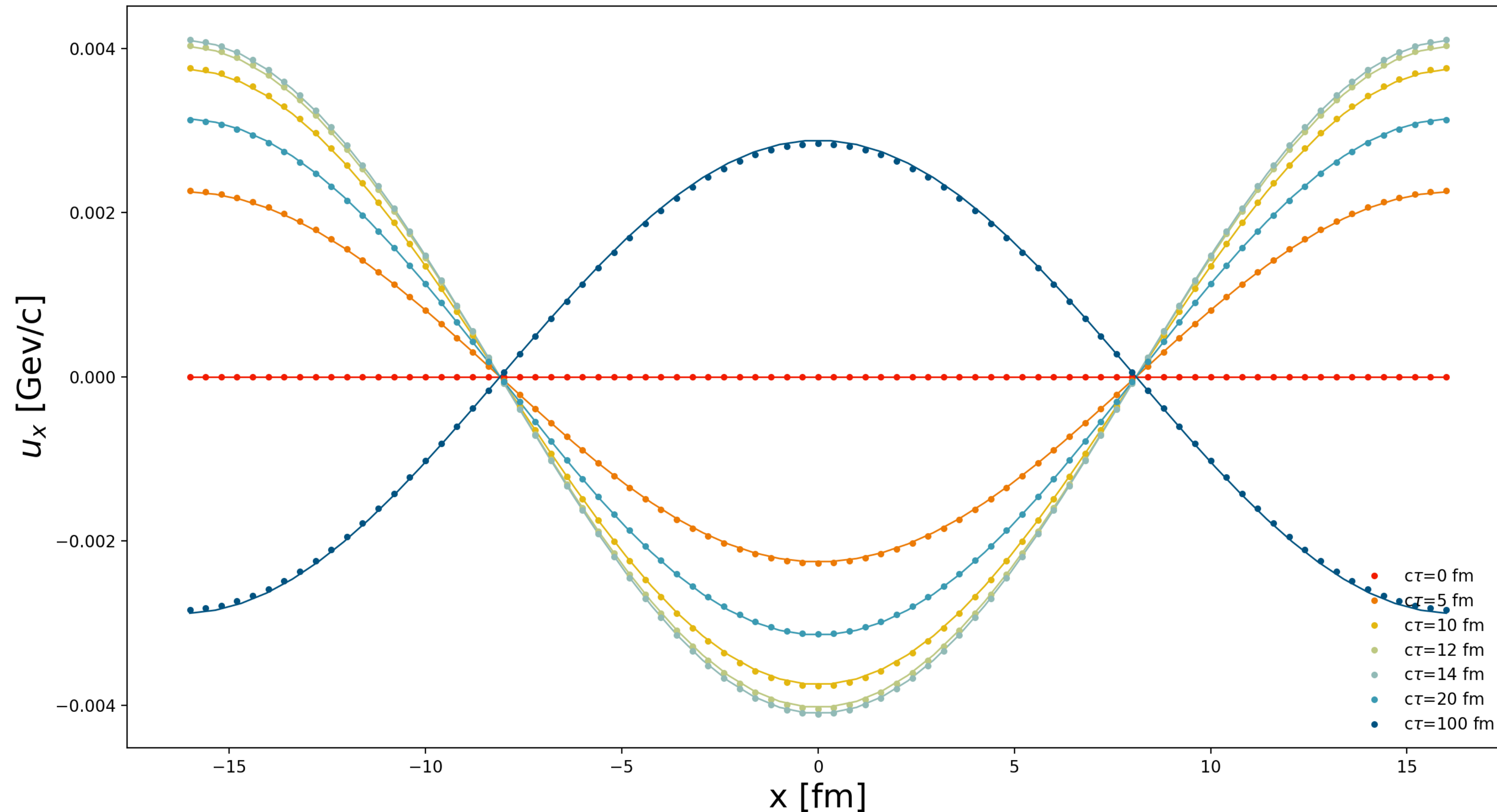
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Introducing of stochastic noise to linearized equations

- Given

$$\partial_t \Xi^{ij} = -\frac{1}{\tau_\pi} (\Xi^{ij} - \xi^{ij})$$

$$\langle \xi^{\mu\nu}(x) \xi^{\alpha\beta}(x') \rangle = \left[2\eta T_0 (\Delta_0^{\mu\alpha} \Delta_0^{\nu\beta} + \Delta_0^{\mu\beta} \Delta_0^{\nu\alpha}) + 2 \left(\zeta - \frac{2}{3}\eta \right) T_0 \Delta_0^{\mu\nu} \Delta_0^{\alpha\beta} \right] \delta^4(x - x')$$

- $\xi^{\mu\nu}$ has the same structure as $\pi^{\mu\nu}$

$$T^{\mu\nu} = T_{id}^{\mu\nu} + T_{visc}^{\mu\nu} + \Xi^{\mu\nu} = T_{id}^{\mu\nu} + T_{visc}'^{\mu\nu} \text{ and } \delta\pi'^{\mu\nu} = \delta\pi^{\mu\nu} + \xi^{\mu\nu}$$

$$\begin{aligned} & \langle u_0^\gamma \partial_{\delta;\gamma} \delta\pi'^{\mu\nu} \rangle_0 + \langle \delta u^\gamma \partial_{0;\gamma} \pi_0^{\mu\nu} \rangle_0 + \langle u_0^\gamma \partial_{0;\gamma} \pi_0^{\mu\nu} \rangle_\delta = \\ & = -\frac{\delta\pi'^{\mu\nu} - \delta\pi_{NS}^{\mu\nu} - \xi^{\mu\nu}}{\tau_{\pi 0}} - \frac{\pi_0^{\mu\nu} - \pi_{0NS}^{\mu\nu}}{\tau_{\pi 0}^2} \delta\tau_\pi - \frac{4}{3} (\pi_0^{\mu\nu} \partial_{\delta;\gamma} \delta u^\gamma + \delta\pi'^{\mu\nu} \partial_{0;\gamma} u_0^\gamma) \end{aligned}$$

Discretization and sampling of noise

- Discretizing the delta function

$$\langle \xi^{\mu\nu}(x) \xi^{\alpha\beta}(x') \rangle = \left[2\eta_0 T_0 (\Delta_0^{\mu\alpha} \Delta_0^{\nu\beta} + \Delta_0^{\mu\beta} \Delta_0^{\nu\alpha}) + 2 \left(\zeta_0 - \frac{2}{3} \eta_0 \right) T_0 \Delta_0^{\mu\nu} \Delta_0^{\alpha\beta} \right] \frac{1}{\Delta t \Delta V}$$

- Sampling from Gaussian with covariance

$$\langle \xi^{\mu\nu}(x) \xi^{\alpha\beta}(x') \rangle = 2\eta_0 T_0 (\Delta_0^{\mu\alpha} \Delta_0^{\nu\beta} + \Delta_0^{\mu\beta} \Delta_0^{\nu\alpha}) \frac{1}{\Delta t \Delta V}$$

- Symmetric tensor
- Subtracting 1/3 of trace from spatial elements

[C. Young, *Phys.Rev.C* 89 (2014) 2]

Structure factor

- Static constant background
- Structure factor - correlation of fields - power spectrum

$$S(\omega, \vec{k}) = A \cdot \langle \delta \hat{U}(\omega, \vec{k}) \delta \hat{U}(\omega', -\vec{k}) \rangle$$

- where A is normalization
- Related to susceptibilities via fluctuation-dissipation relation
- Equal time correlation - static structure factor

$$S(\vec{k}) = A \cdot \langle \delta \hat{U}(\vec{k}) \delta \hat{U}(-\vec{k}) \rangle$$

Normalization of structure factor

$$\partial_t U = LU + KW$$

- in NS limit in 1D [A. Donev et al., CAMCOS (2009)]

$$\begin{bmatrix} \partial_t \delta \varepsilon \\ \partial_t \delta u^x \end{bmatrix} = -\partial_x \begin{bmatrix} (\varepsilon_0 + p_0) \delta u^x \\ \frac{c_s^2}{\varepsilon_0 + p_0} \delta \varepsilon \end{bmatrix} + \partial_x \begin{bmatrix} 0 \\ \frac{4}{3} \frac{\eta}{\varepsilon_0 + p_0} \partial_x \delta u^x \end{bmatrix} + \frac{1}{\varepsilon_0 + p_0} \partial_x \begin{bmatrix} 0 \\ \xi \end{bmatrix}$$

- identifying L and K matrices

$$\hat{L} = -ik \begin{pmatrix} 0 & \varepsilon_0 + p_0 \\ \frac{c_s^2}{\varepsilon_0 + p_0} & -ik \frac{4}{3} \frac{\eta}{\varepsilon_0 + p_0} \end{pmatrix}, \hat{K} = ik \begin{pmatrix} 1 & 0 \\ 0 & \frac{1}{\varepsilon_0 + p_0} \end{pmatrix} = -\hat{K}^* \text{ and } C_W = \begin{pmatrix} 0 & 0 \\ 0 & \xi \end{pmatrix} \text{ where } \xi = \frac{8}{3} \eta_0 T_0$$

- Using the equation

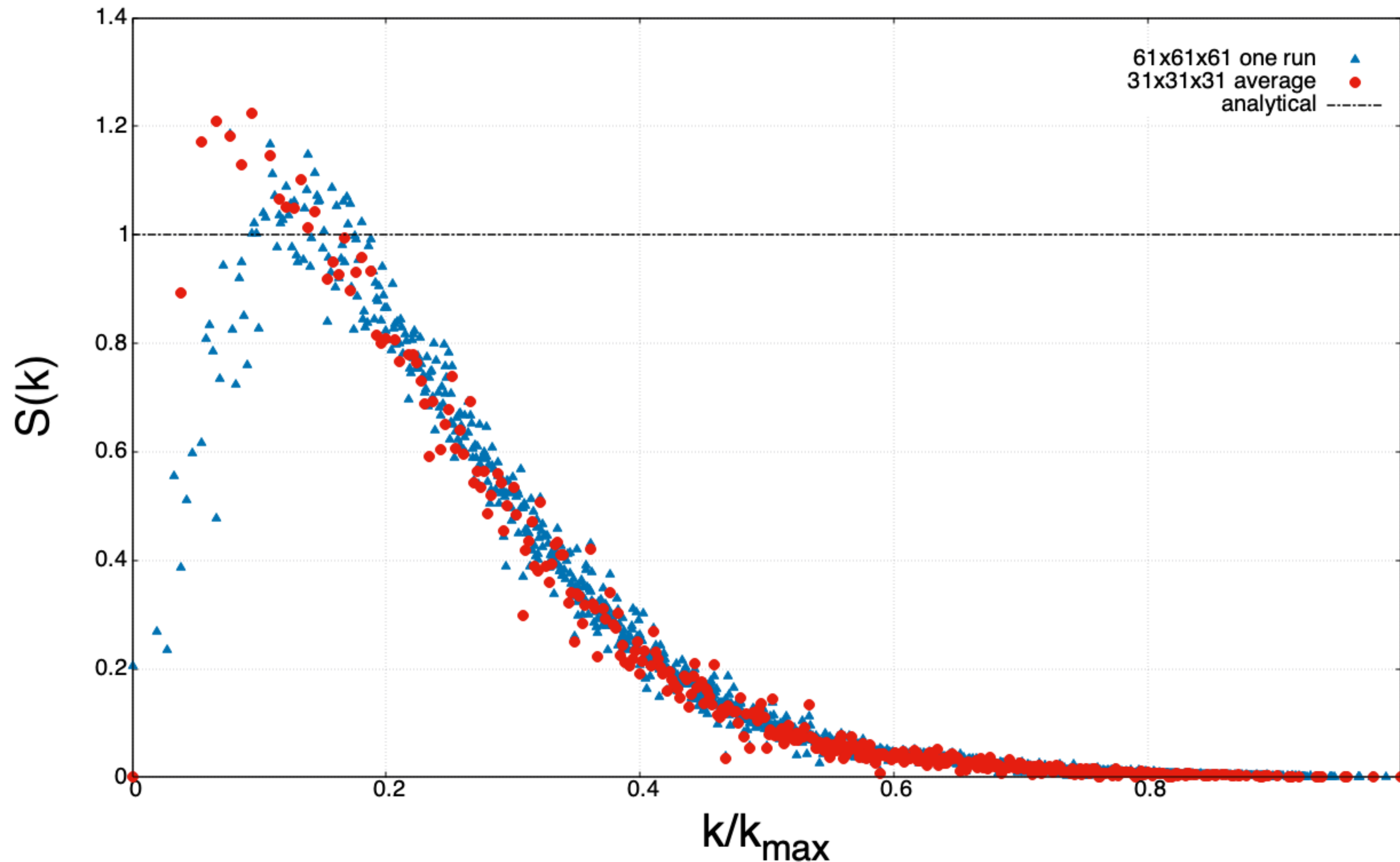
$$\hat{L}S + S\hat{L}^* = -\hat{K}C_W\hat{K}^*$$

The structure factor matrix $S(k) = \begin{pmatrix} c_s^{-2}(\varepsilon_0 + p_0)T_0 & 0 \\ 0 & (\varepsilon_0 + p_0)^{-1}T_0 \end{pmatrix}$ independent of k

Current status

- Using KISS FFT to transform fields to Fourier space

$$S(k) = \frac{V}{c_s^2(\epsilon_0 + p_0)T_0} \langle \delta\epsilon(k)\delta\epsilon(-k) \rangle$$



Conclusion and further steps

- Thermal fluctuations should be included - Fluctuation dissipation theorem
- Fluctuations provide good basis for studying phase diagram
 - Critical fluctuation for studying critical point
- Stochastic fluid dynamics can include non-linear terms
 - But it has some difficulties - fluctuation larger than background, discretization dependence
- Further steps
 - Calculating static structure factor for finer grid
 - Dynamic structure factor
 - Renormalization of the grid dependence