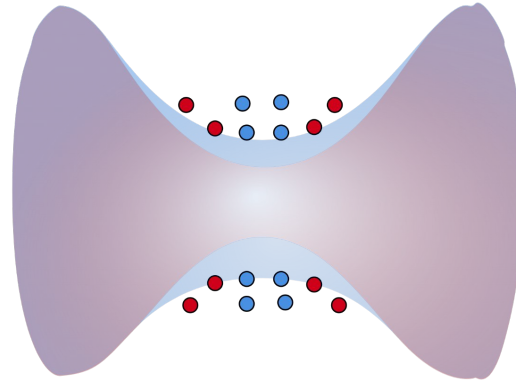
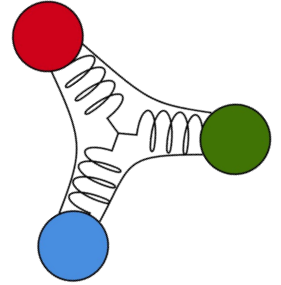
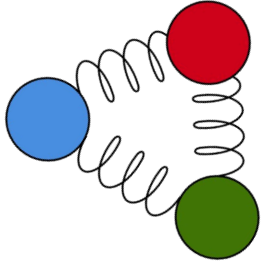


Tracing the conserved charges in high energy collisions



Hydrodynamics and related observables in heavy-ion collisions
Oct. 31, 2024, Nantes, France

Grégoire Pihan¹, Akihiko Monnai², Bjoern Schenke³, Chun Shen^{1,3}



WAYNE STATE
UNIVERSITY

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³Brookhaven National Lab, Upton, USA

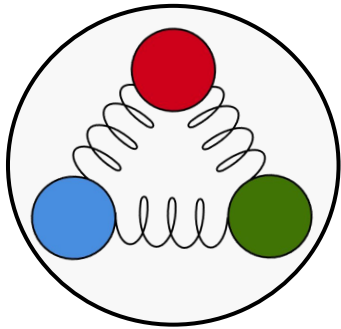


U.S. DEPARTMENT OF
ENERGY

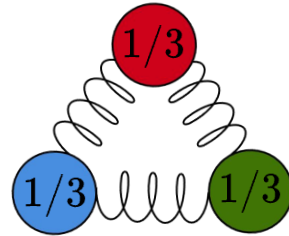
What carries the baryon number?

How do we build a baryon from quarks?

$$B = 1$$



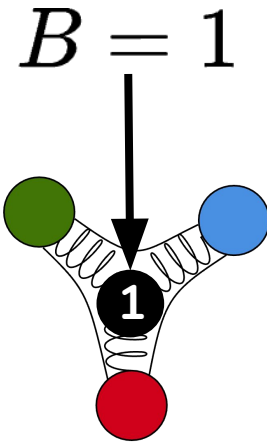
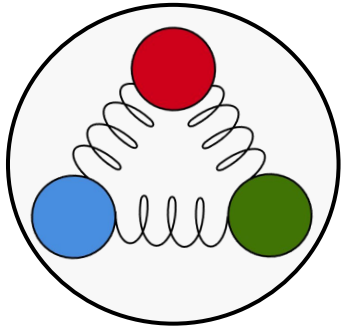
$$B = \frac{1}{3} (n_q - n_{\bar{q}})$$



This is an assumption!

String junction?

$$B = 1$$

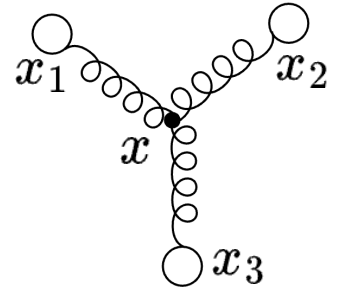


What carries the baryon number?

String junction: The most simple way to build a baryon from quarks

Non-perturbative configuration of gluons represented by a locally **gauge-invariant** state vector.

[G.C Rossi and G.Veneziano PHYSICS REPORTS 63, No. 3 \(1980\)](#)



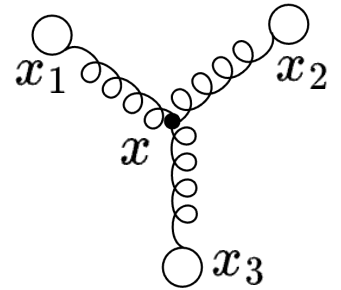
$$B = \epsilon^{ijk} \left[P \exp \left(ig \int_{x_1}^x A_\mu dx^\mu \right) q(x_1) \right]_i \left[P \exp \left(ig \int_{x_2}^x A_\mu dx^\mu \right) q(x_2) \right]_j \left[P \exp \left(ig \int_{x_3}^x A_\mu dx^\mu \right) q(x_3) \right]_k$$

What carries the baryon number?

String junction: The most simple way to build a baryon from quarks

Non-perturbative configuration of gluons represented by a locally **gauge-invariant** state vector.

G.C Rossi and G.Veneziano PHYSICS REPORTS 63, No. 3 (1980)

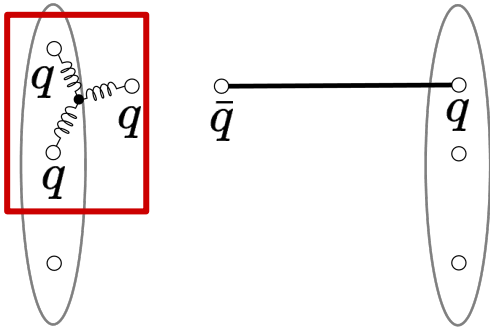


$$B = \epsilon^{ijk} \left[P \exp \left(ig \int_{x_1}^{\bullet x} A_\mu dx^\mu \right) q(x_1) \right]_i \left[P \exp \left(ig \int_{x_2}^{\bullet x} A_\mu dx^\mu \right) q(x_2) \right]_j \left[P \exp \left(ig \int_{x_3}^{\bullet x} A_\mu dx^\mu \right) q(x_3) \right]_k$$

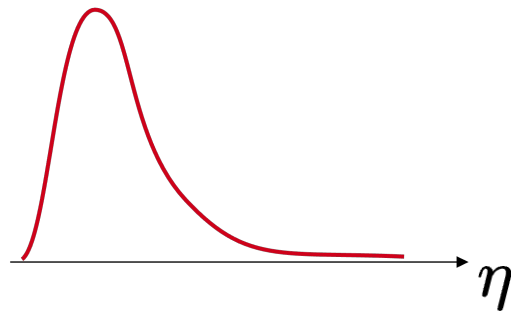
The string junction x carries the baryon number inside the baryon

Baryon stopping in central pp and AA collisions

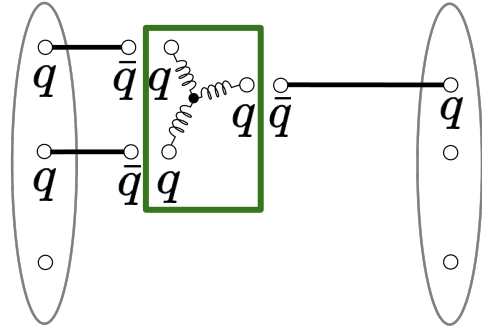
$B = 1$



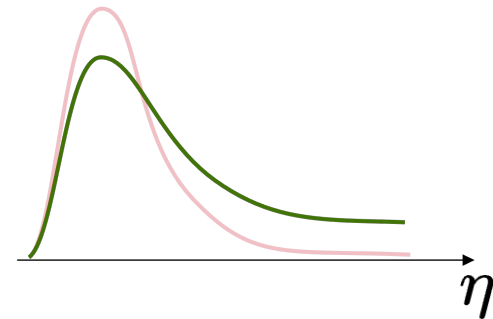
The baryon number remain attached to the nucleon

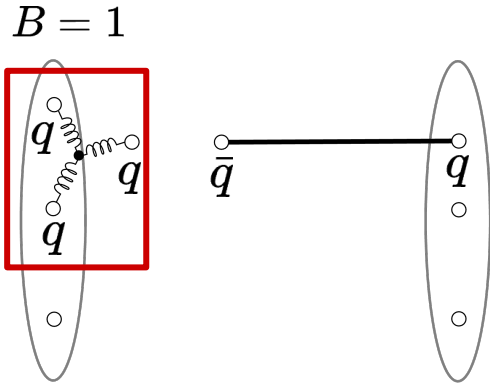


$B = 1$

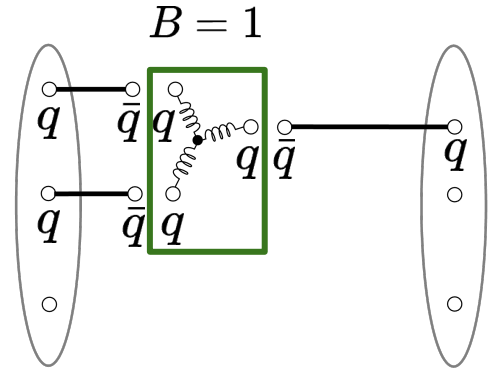
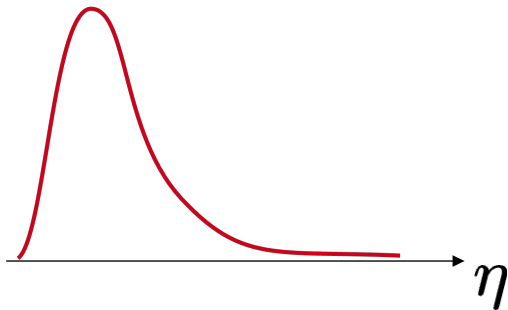


The baryon number to fluctuate towards mid-rapidity

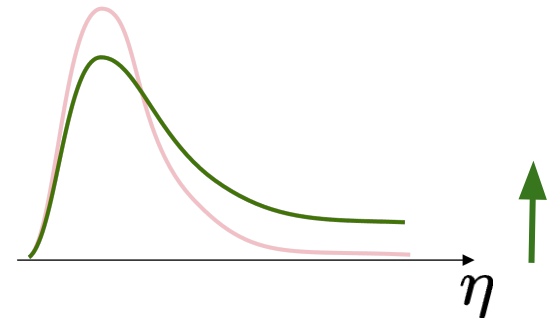




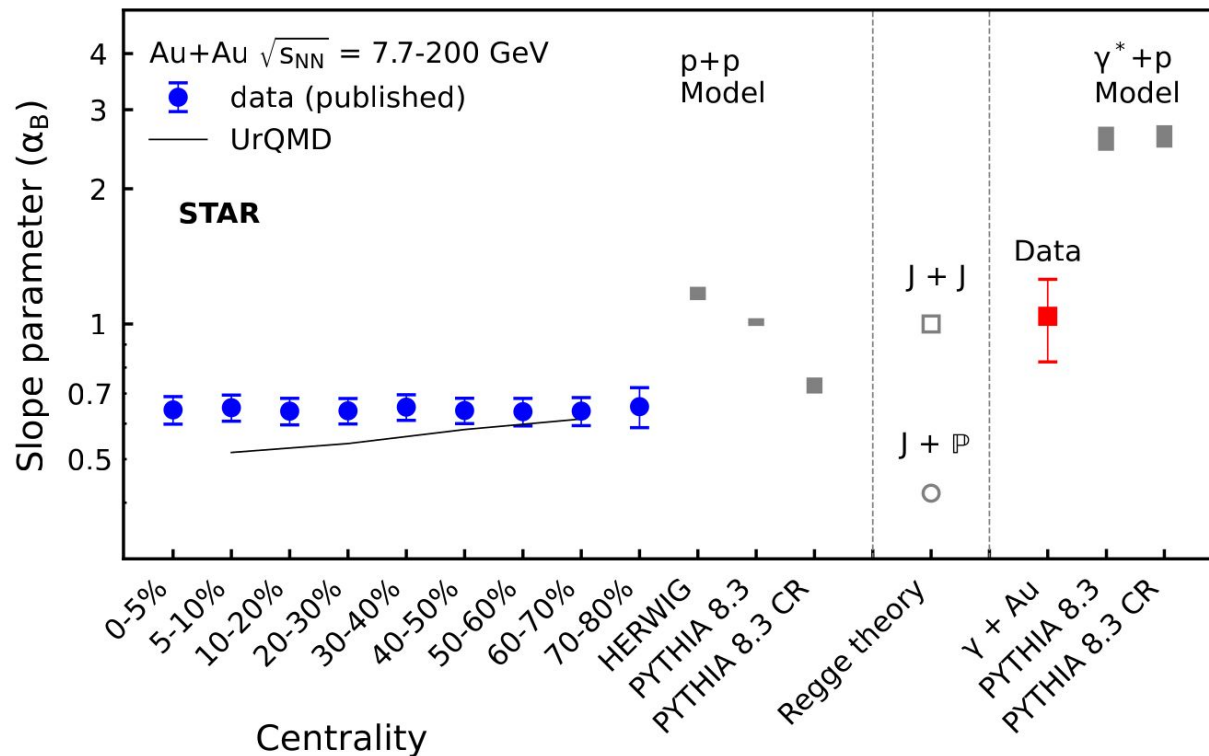
The baryon number remain attached to the nucleon



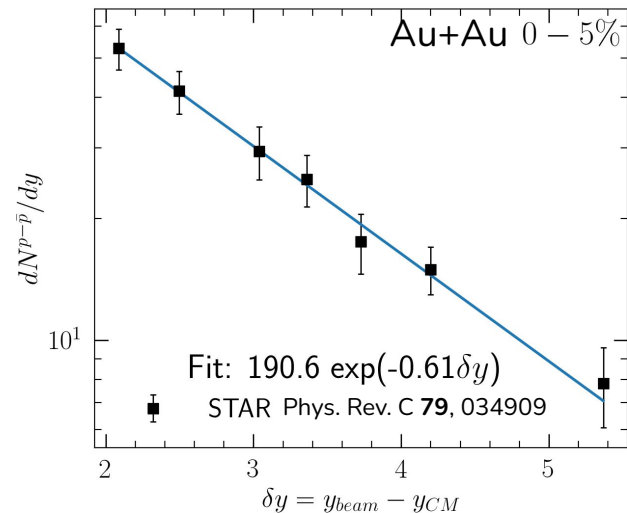
The baryon number to fluctuate towards mid-rapidity



Baryon stopping and string junction



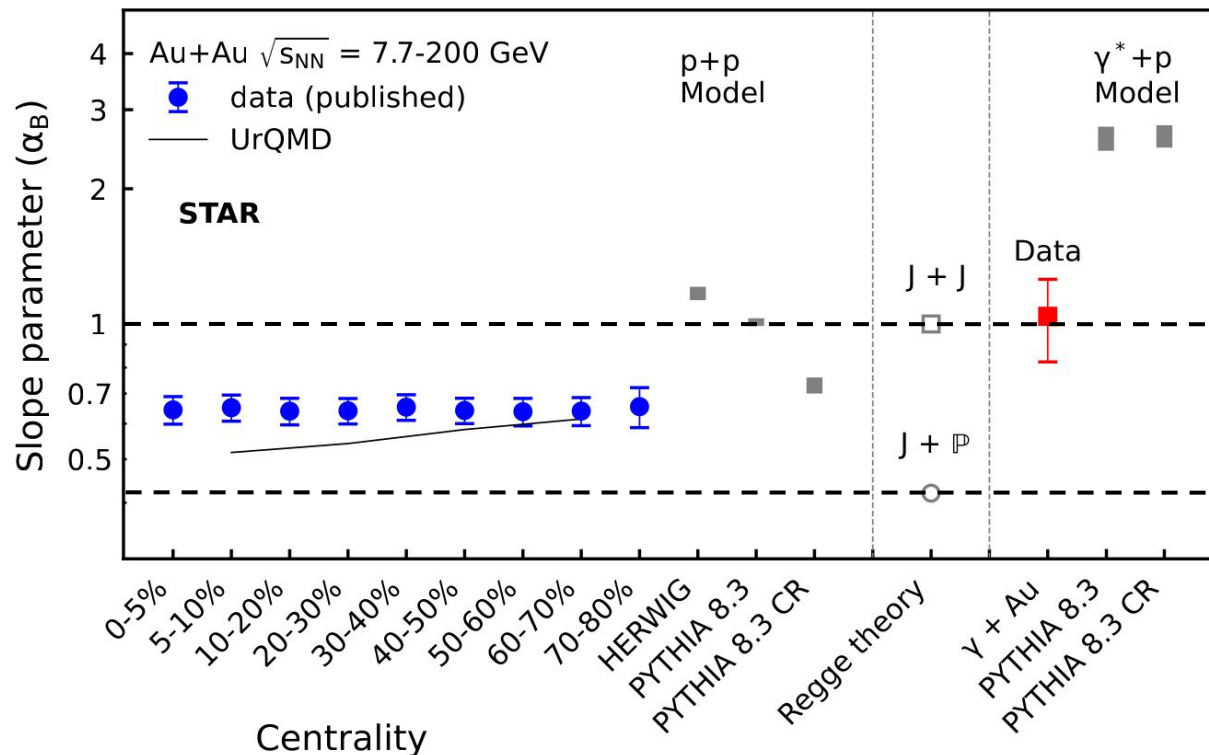
STAR Collab, arxiv:2408.15441



$$\frac{dN^{p-\bar{p}}}{d\delta y} \propto e^{-\alpha_B \delta y}$$

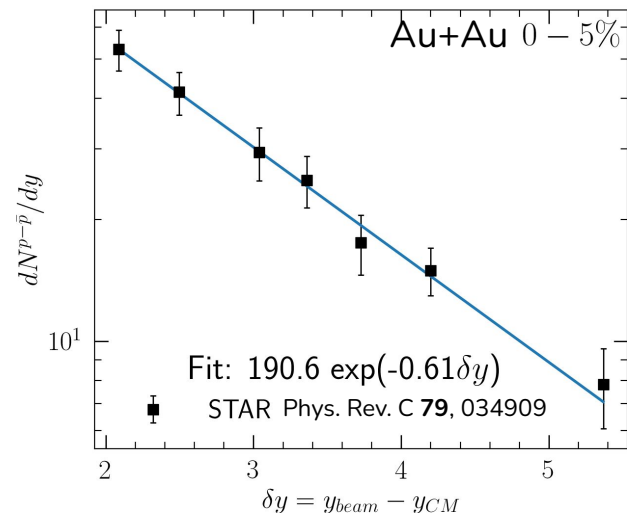
$$\alpha_B = 0.64 \pm 0.05$$

Baryon stopping and string junction



STAR Collab, [arxiv:2408.15441](https://arxiv.org/abs/2408.15441)

Compatible with
Baryon junction theory predictions!



$$\frac{dN^{p-\bar{p}}}{d\delta y} \propto e^{-\alpha_B \delta y}$$

$$\alpha_B = 0.64 \pm 0.05$$

Insight from the isobar collisions at RHIC

Baryon Junction: carries no electric charge!



Decorrelation of B and Q!
if B remains attached to the junction

Isobar Runs: Same number of nucleons A, different number of protons Z



Baryon stopping compared to **electric charge stopping!**

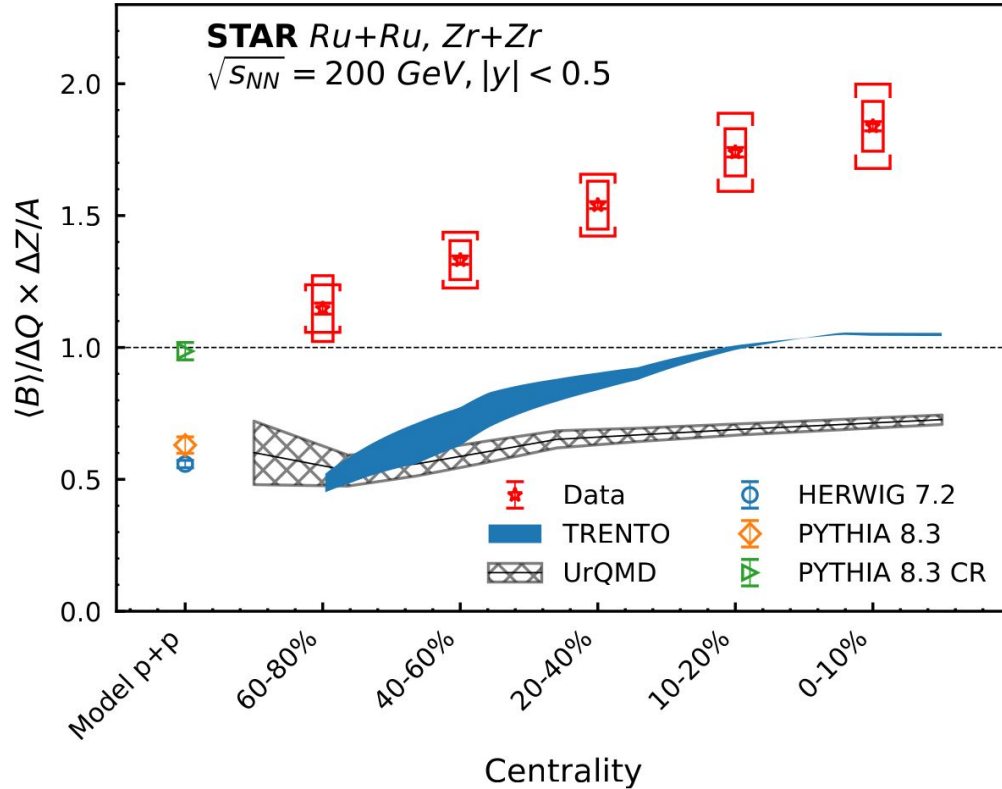
“Equal stoppings”

No longitudinal
decorrelation from
junction:
**B and Q carried by
valence quarks!**

“Different stoppings”

B and Q are less
correlated in the
longitudinal direction:
**B is carried by the
junction!**

Insight from the isobar collisions at RHIC



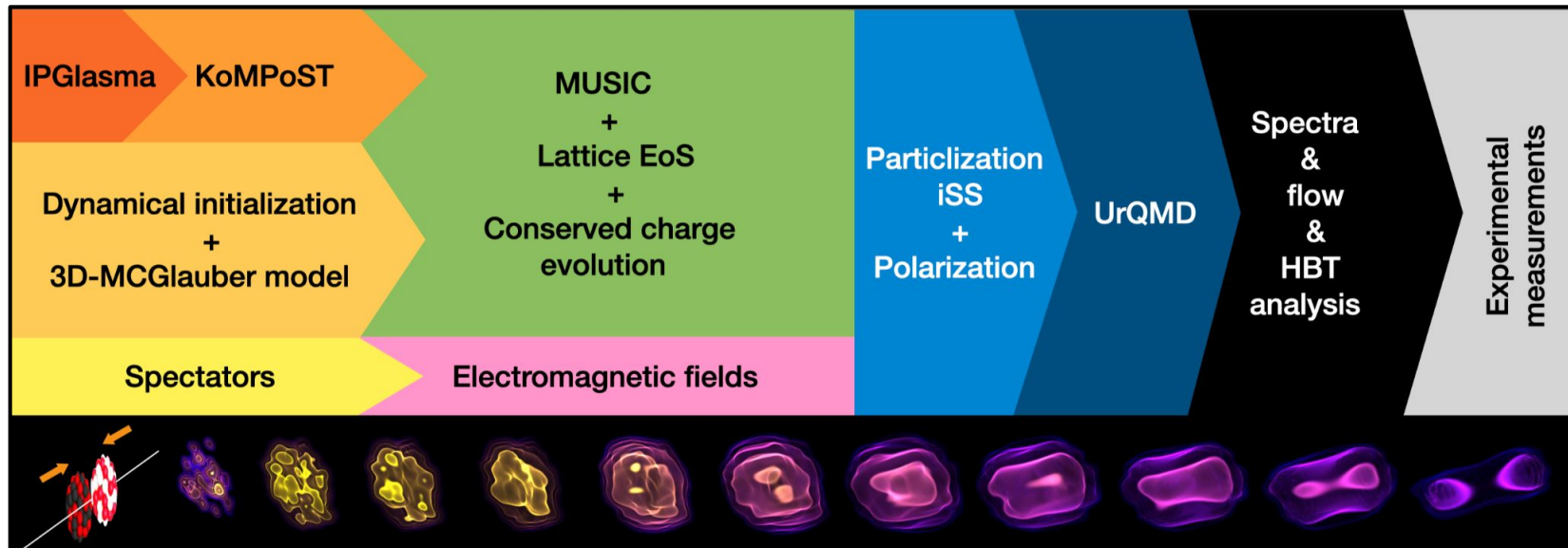
STAR measurement of stopping ratio deviates from unity

More baryon stopping than electric charge stopping:

Different stoppings!

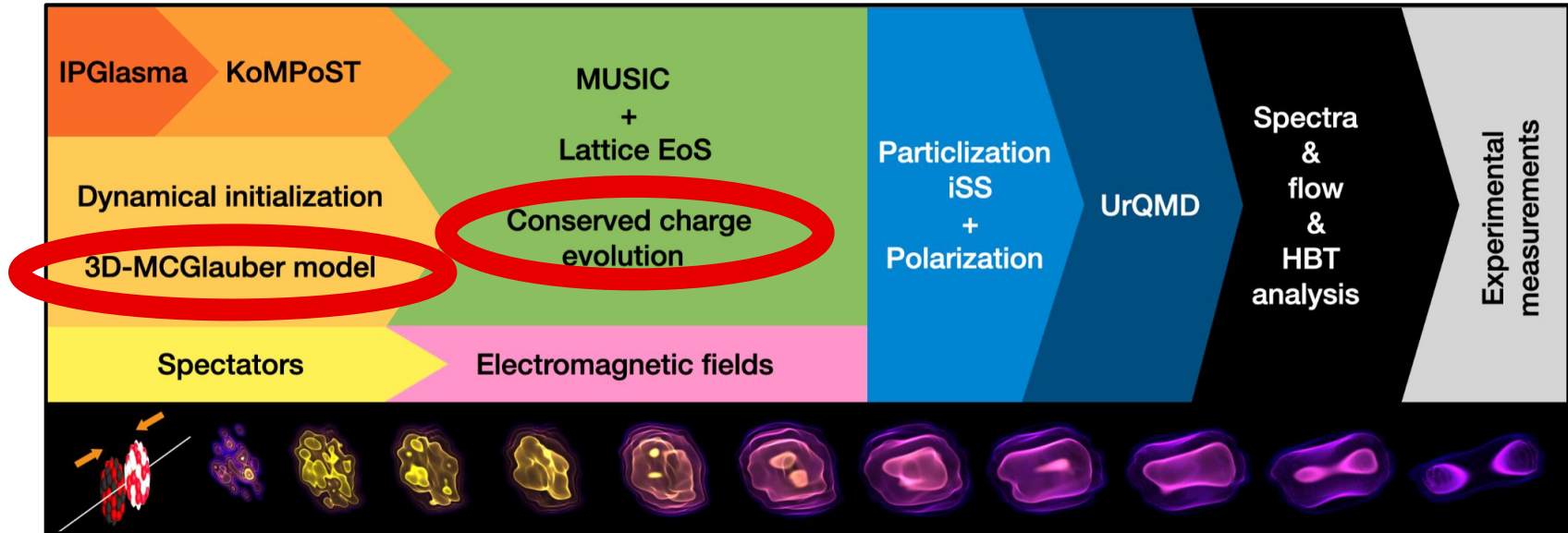
Is this a sign of the existence of the baryon junction?

Open source hydrodynamics + hadronic transport hybrid framework

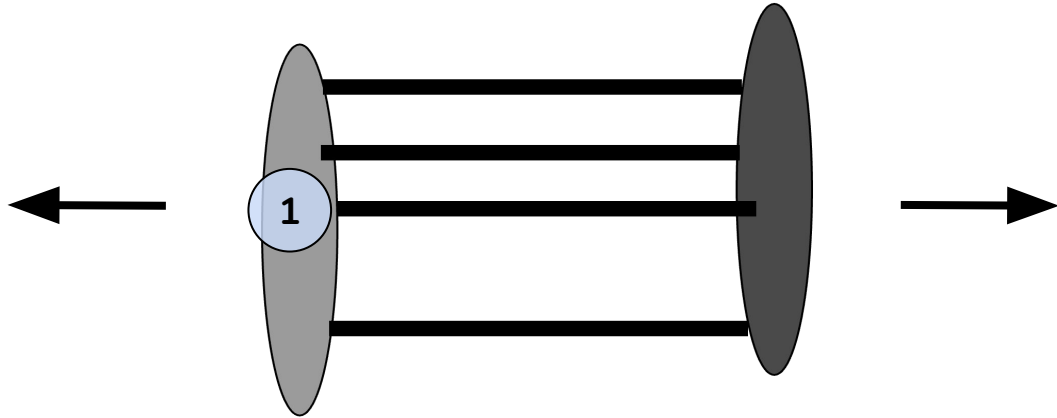


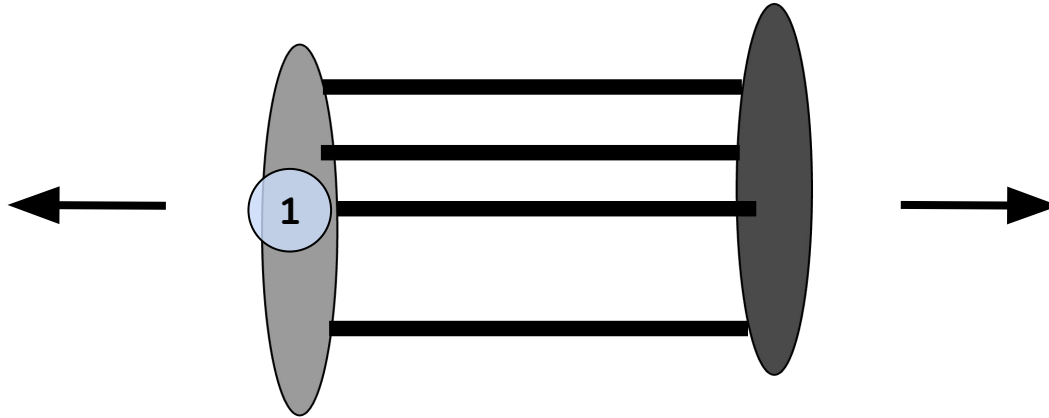
<https://github.com/chunshen1987/iEBE-MUSIC>

Open source hydrodynamics + hadronic transport hybrid framework

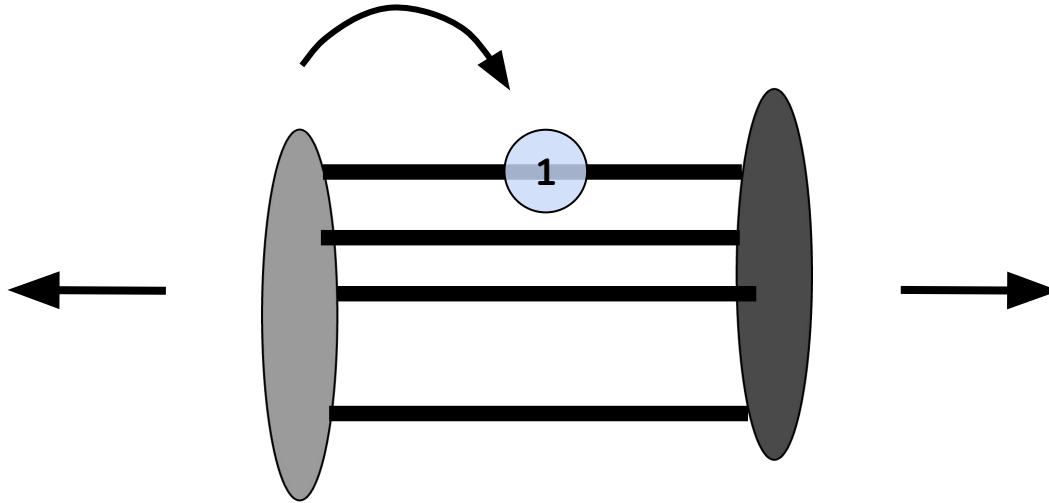


<https://github.com/chunshen1987/iEBE-MUSIC>

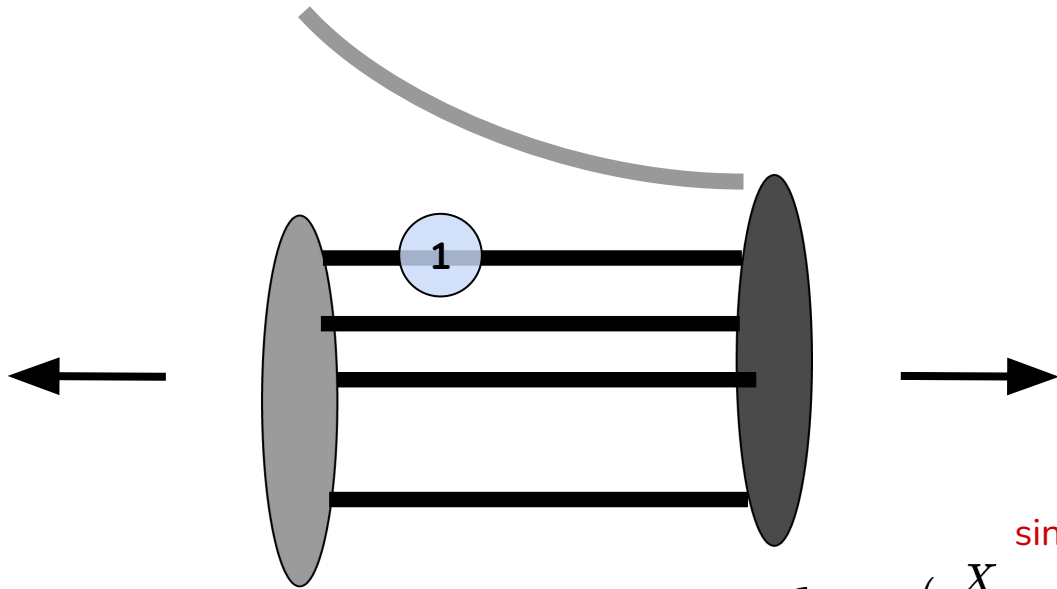




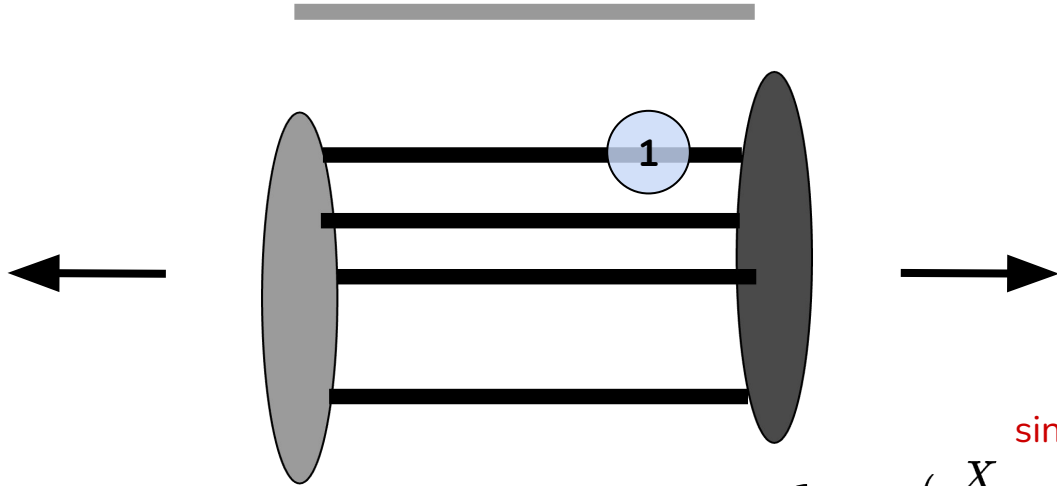
$$P(y_{P/T}^X) = (1 - \lambda_X) y_{P/T}$$



$$P(y_{P/T}^X) = (1 - \lambda_X)y_{P/T} + \lambda_X$$



$$P(y_{P/T}^X) = (1 - \lambda_X)y_{P/T} + \lambda_X \left\{ \begin{array}{l} \text{single junction} \\ e^{(y_{P/T}^X - (y_P + y_T)/2)/2} \\ 4 \sinh((y_P - y_T)/4) \end{array} \right.$$

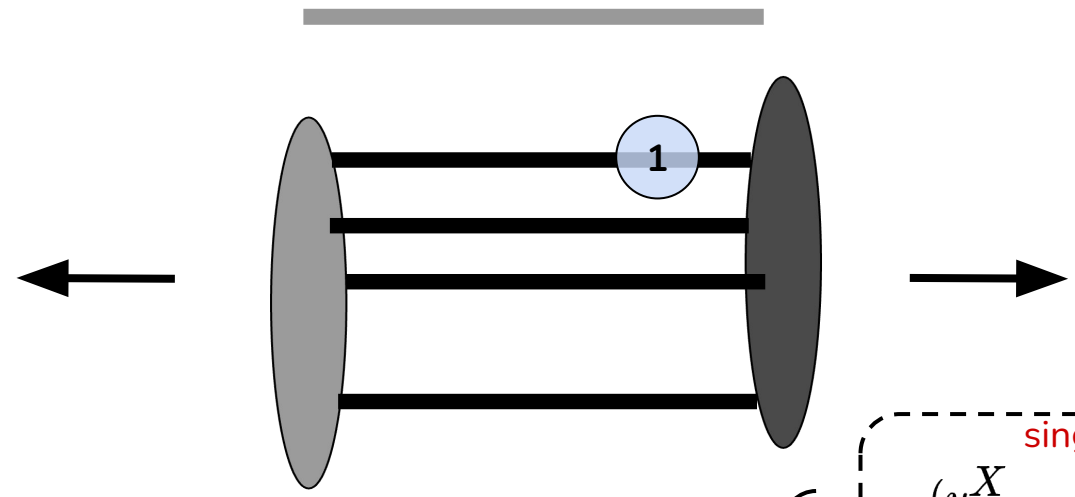


$$P(y_{P/T}^X) = (1 - \lambda_X)y_{P/T} + \lambda_X$$

single junction

$$\frac{e^{(y_{P/T}^X - (y_P + y_T)/2)/2}}{4 \sinh((y_P - y_T)/4)}$$

double junction constant



$$P(y_{P/T}^X) = (1 - \lambda_X)y_{P/T} + \lambda_X$$

single junction

$$\frac{e^{(y_{P/T}^X - (y_P + y_T)/2)/2}}{4 \sinh((y_P - y_T)/4)}$$

~~double junction constant~~

Initial conditions from the string junction 1/3

$$P(y_{P/T}^X) = (1 - \lambda_X)y_{P/T} + \lambda_X \frac{e^{(y_{P/T}^X - (y_P + y_T)/2)/2}}{4 \sinh((y_P - y_T)/4)}$$

Net-baryon number:

$$\lambda_B = 0.2$$

Net-electric charge:

free parameter

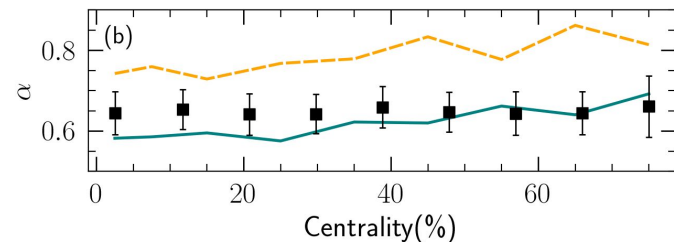
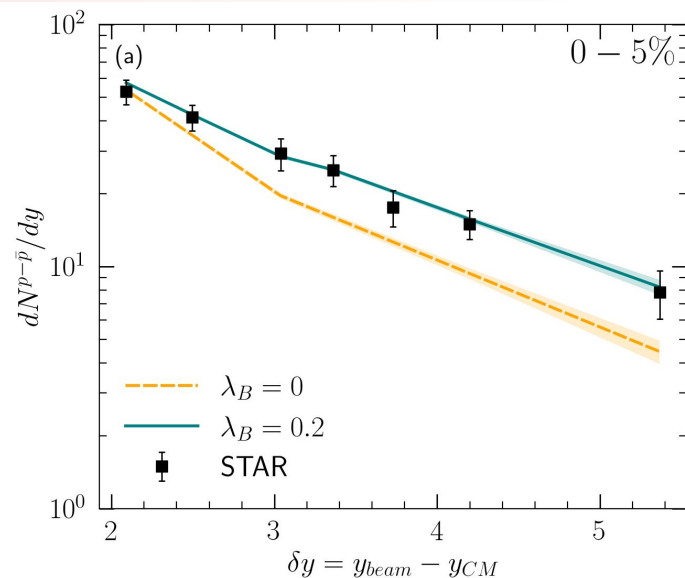
$$\lambda_Q$$

$$\lambda_Q = \lambda_B$$

$$\lambda_Q \neq \lambda_B$$

Equal stoppings

Different stoppings



C. Shen and B. Schenke Phys. Rev. C **105**, 064905 (2022)

GP, A. Monnai, B. Schenke, C. Shen Phys. Rev. Lett. **133**, 182301

STAR collaboration, Phys. Rev. C **79**, 034909, arxiv:2408.15441

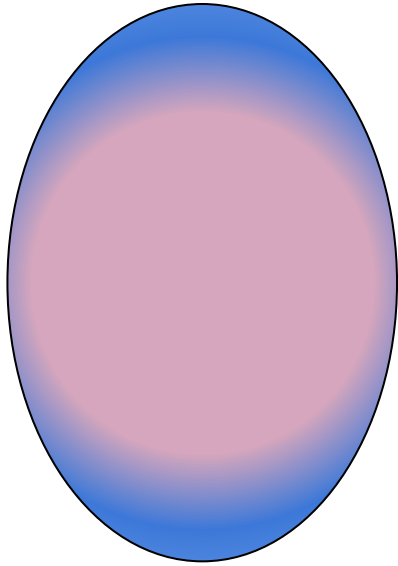
Wood-Saxon potential: **nuclear structure** and **neutron skin**

$$R_{p,n}(\theta, \phi) = R_{p,n}(1 + \beta_2 Y_2^0(\theta, \phi) + \beta_3 Y_3^0(\theta, \phi))$$

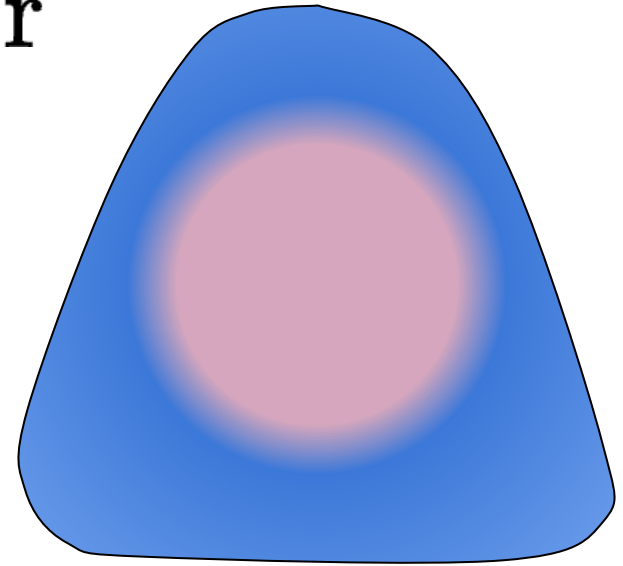
	R_p	a_p	R_n	a_n	β_2	β_3
Ru	5.09	0.46	5.105	0.47	0.16	0.0
Zr	5.02	0.52	5.12	0.57	0.06	0.2

p: protons
n: neutrons

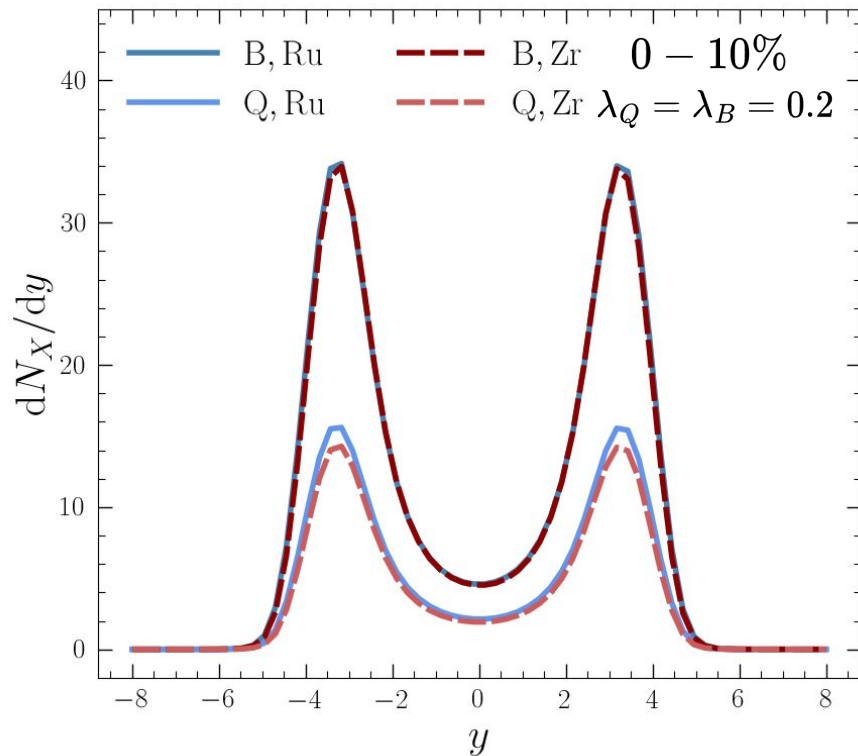
R_u



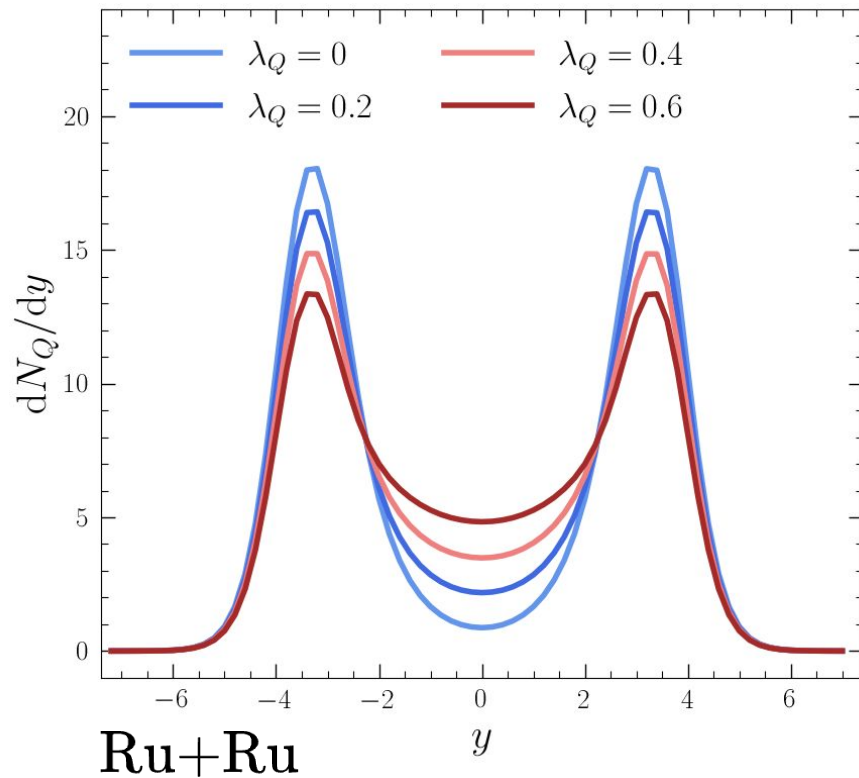
Z_r



Initial baryon and electric charge density rapidity distributions for isobar runs at $\sqrt{s_{NN}} = 200$ GeV



Initial electric charge density rapidity distributions for different values of λ_Q



MUSIC with 4D equation of state

NEOS 4D equation of state

Taylor expansion at finite chemical potentials

$$\frac{P_{\text{Latt}}}{T^4} = \frac{P_0}{T^4} + \sum_{l,n,m} \chi_{l,n,m}^{B,Q,S} \left(\frac{\mu_B}{T}\right)^l \left(\frac{\mu_Q}{T}\right)^n \left(\frac{\mu_S}{T}\right)^m$$

Hadron Resonance Gas

$$P_{\text{HRG}} = \pm T \sum_i g_i \int \frac{d^3k}{(2\pi)^3} \ln [1 \pm e^{(E_i(k) - \mu_i)/T}]$$

i : hadronic species $\mu_i = B_i \mu_B + Q_i \mu_Q + S_i \mu_S$

Matching

$$\frac{P}{T^4} = \frac{1}{2} [1 - f(T, \mu_X)] \frac{P_{\text{HRG}}}{T^4} + \frac{1}{2} [1 + f(T, \mu_X)] \frac{P_{\text{Latt}}}{T^4}$$

No assumptions on the relation between conserved charge densities!

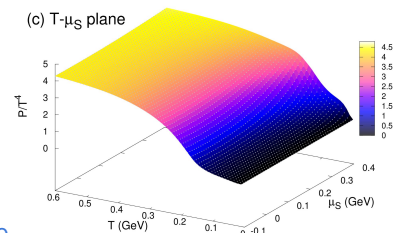
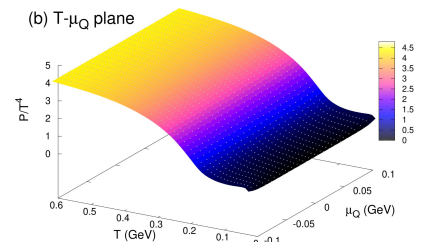
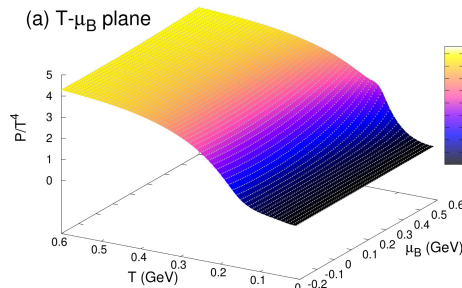
~~$$\rho_Q \approx 0.4 \rho_B$$~~

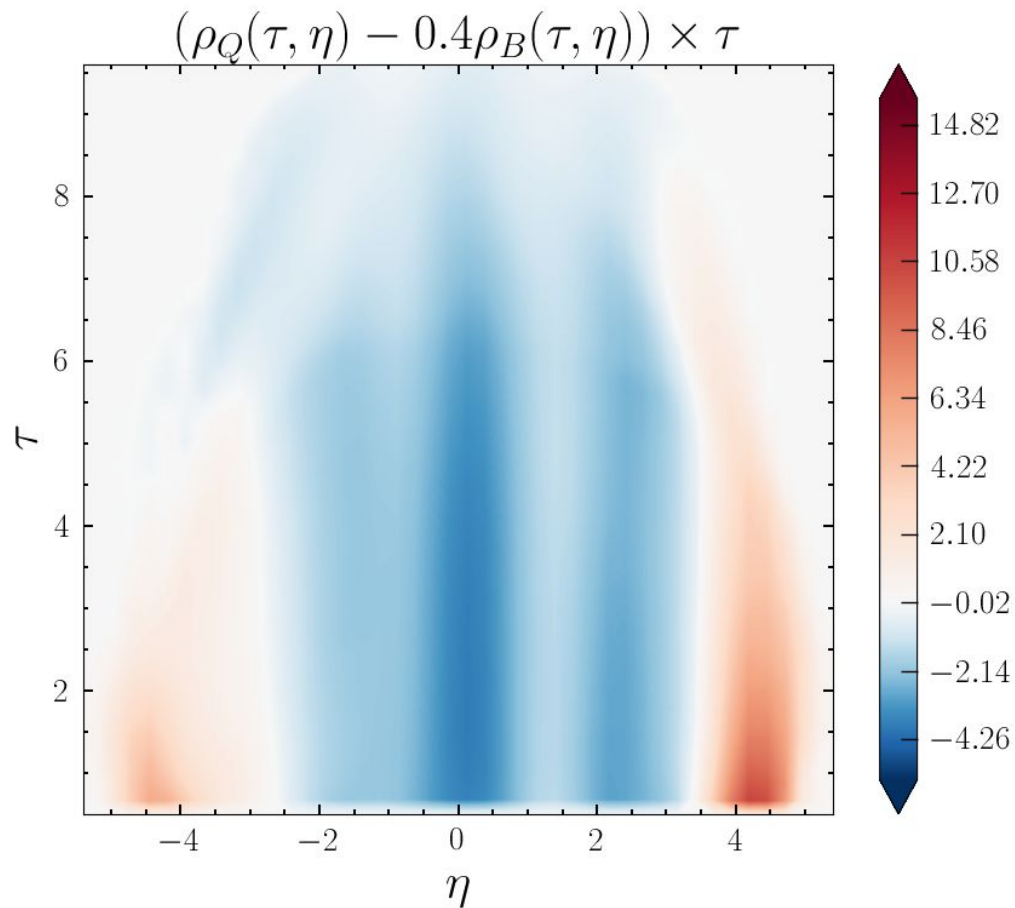
MUSIC with BQS conserved charges

$$\begin{cases} \partial_\mu T^{\mu\nu} = 0 \\ \partial_\mu N_X^\mu = 0 \end{cases} \quad N_X^\mu = \rho_X u^\mu$$

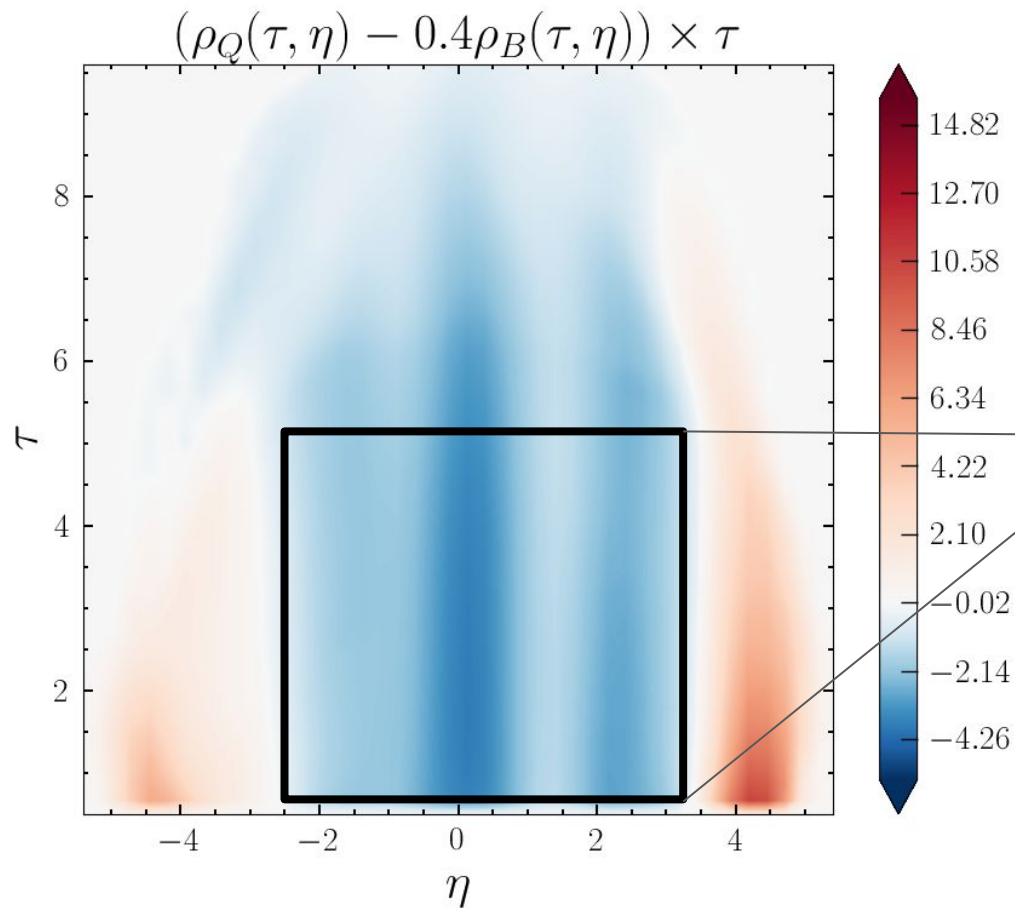
B, Q and S currents evolve independently!

$$X = B, Q, S$$





~~$\rho_Q \approx 0.4\rho_B$~~



~~$\rho_Q \approx 0.4\rho_B$~~

Local excess baryon density!

Case 1
"Equal stopping"

$$\lambda_Q = \lambda_B = 0.2$$

Naive expectation from
charge conservation

$$r \sim 1$$

Selection

Case 2
"No extra Q stopping"

$$\lambda_Q = 0, \lambda_B = 0.2$$

No extra stopping mechanism
for electric charge:

**Maximal difference between
B and Q stopping at fixed λ_B**

$$r > 1$$

In each centrality class

$$N_{\text{ch,Ru}} = N_{\text{ch,Zr}}$$

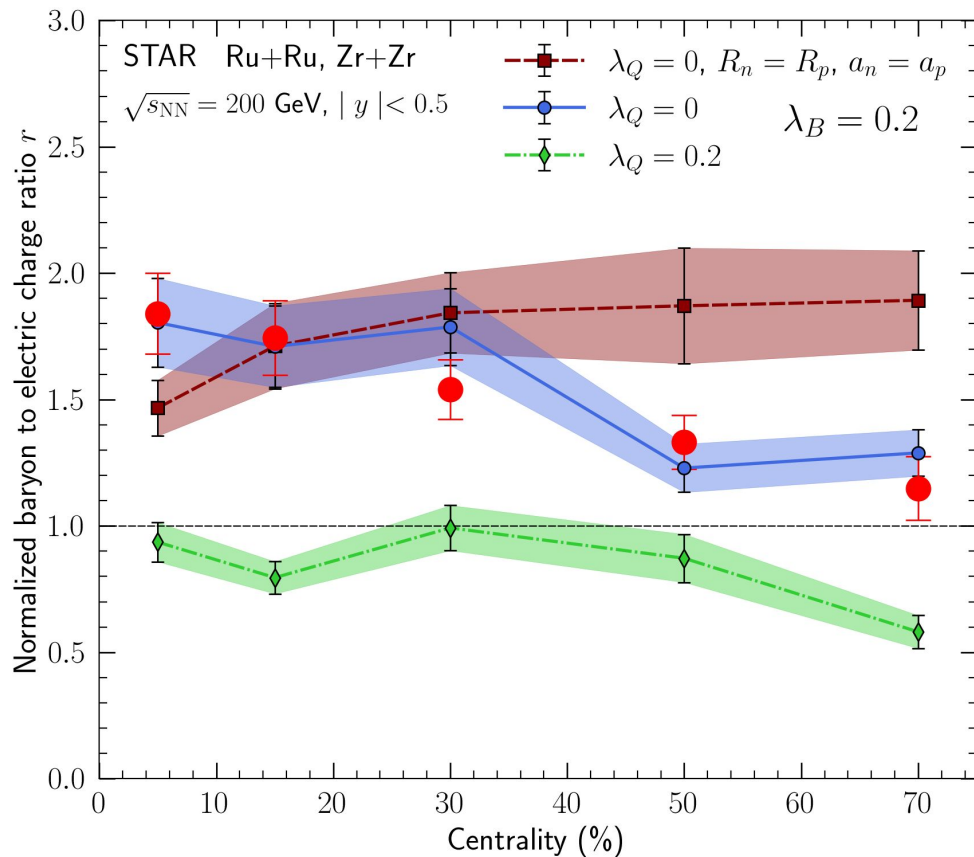
Case 3
Case 1 + "No neutron skin"

$$\lambda_Q = 0, \lambda_B = 0.2$$

$$R_p = R_n, a_p = a_n$$

No extra stopping for electric
charge and no structure as a
function of centrality

$$r > 1 \quad \frac{dr}{dN_{\text{part}}} \sim 0$$



GP, A. Monnai, B. Schenke, C. Shen Phys. Rev. Lett. **133**, 182301

Case 2:

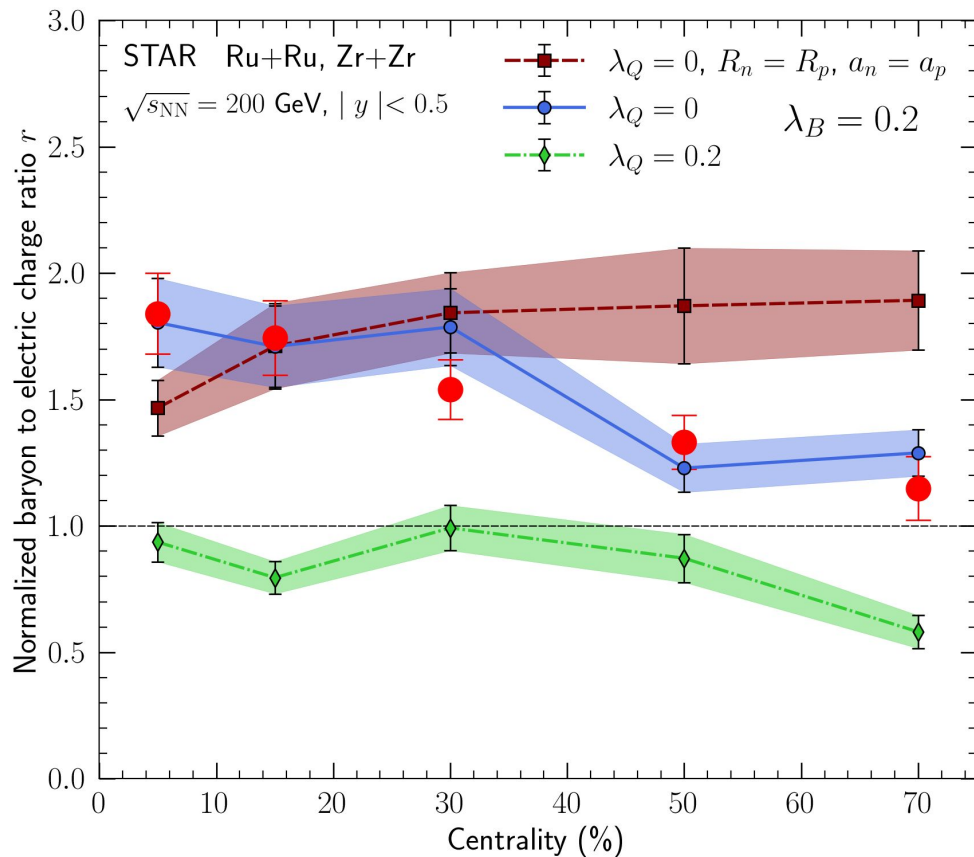
Unequivocally favored by experimental data!

Case 1:

The naive expectation indeed falls to unity but has decreasing behavior at large centrality

Case 3:

The structure is due to the neutron skin of the Zirconium!



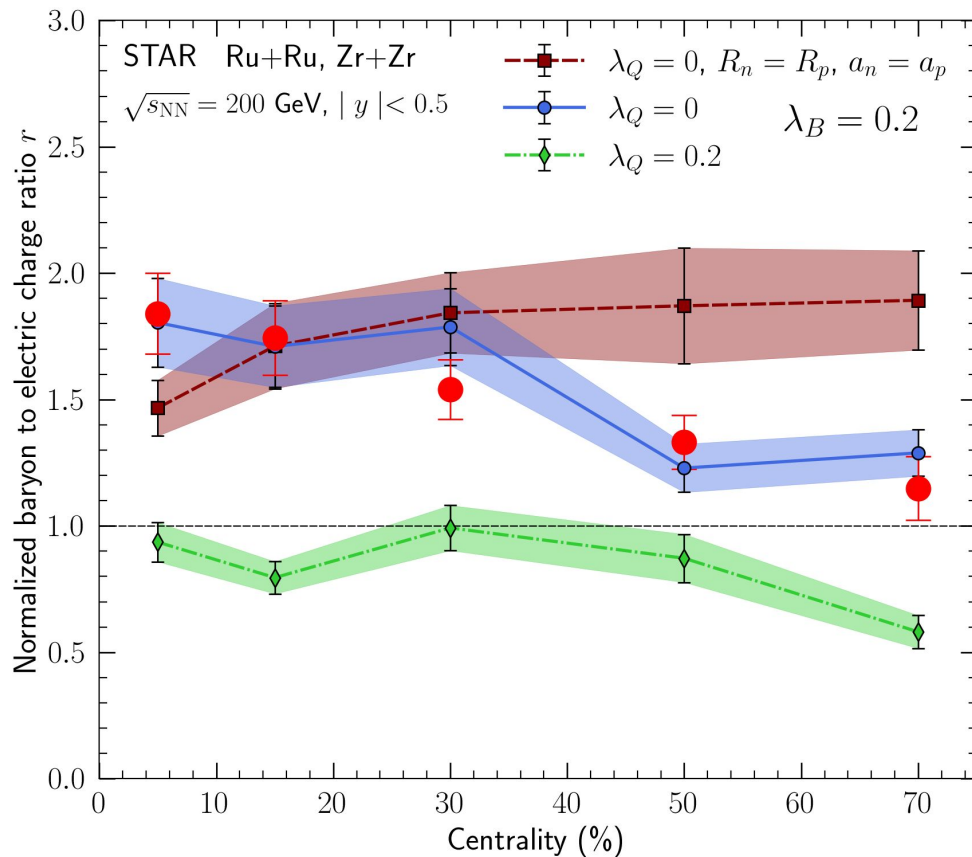
GP, A. Monnai, B. Schenke, C. Shen Phys. Rev. Lett. **133**, 182301

Case 2:

Unequivocally favored by experimental data!

This study:

Any stopping models that predicts an extra baryon stopping mechanism with the same longitudinal probability distribution as the single baryon junction is largely favored by the experimental measurement in the isobar runs at 200 GeV!



GP, A. Monnai, B. Schenke, C. Shen Phys. Rev. Lett. **133**, 182301

Case 2:

**Unequivocally favored by
experimental data!**

In all likelihood:

**This is strongly in favor of the
baryon junction picture!**

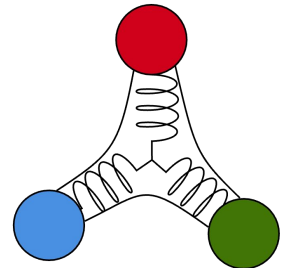
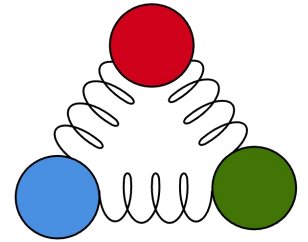
Baryon junction: take home

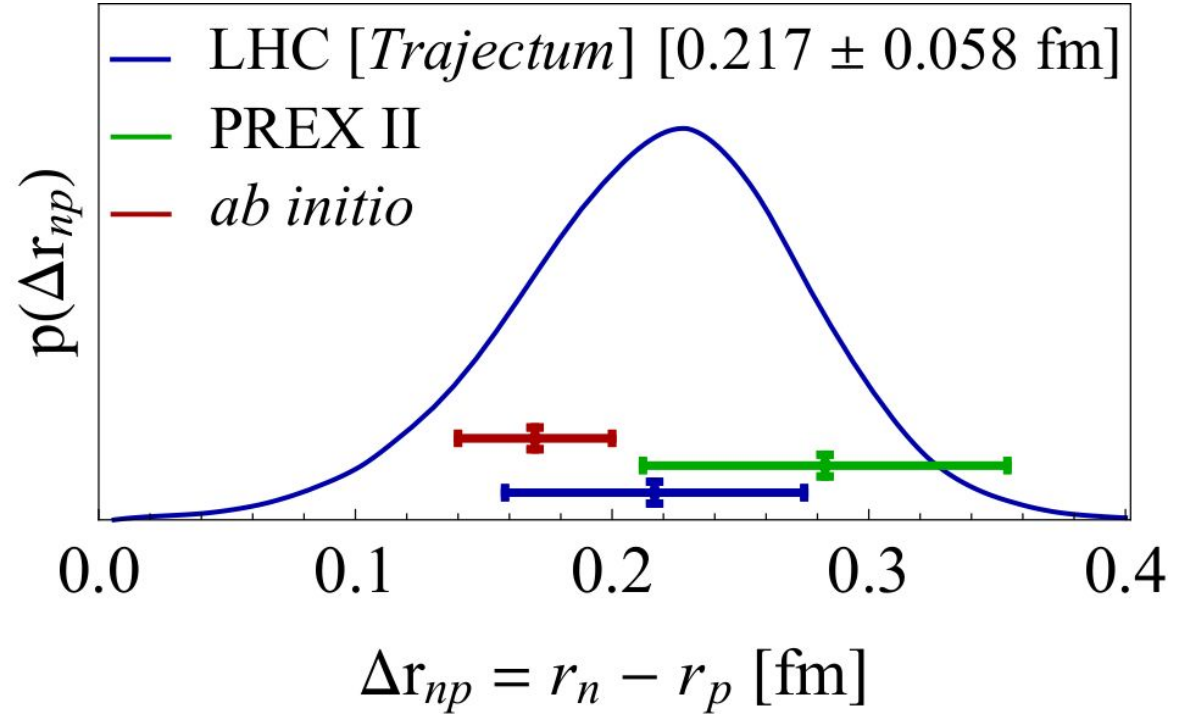
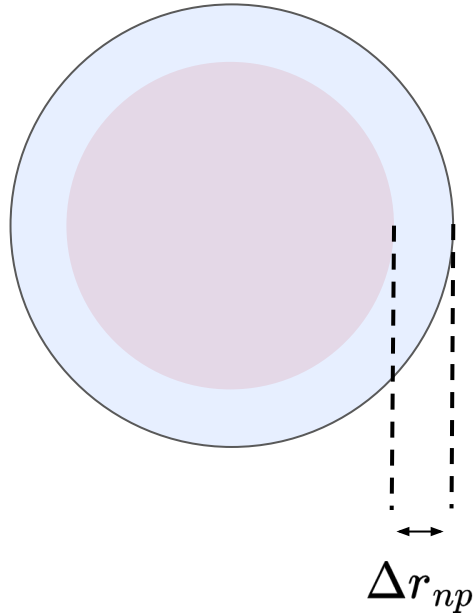
“Can gluon junction trace the baryon number?”

The isobar simulations are in **quantitative agreement** with the STAR measurement.

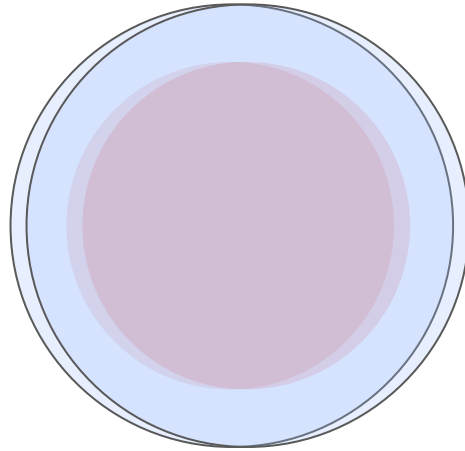
Strong evidence for the existence of the junction!

The ratio is sensitive to the Zr neutron skin
Study of the nuclear structure in Heavy-ion collisions?





Ultra-central collisions in ALICE Pb-Pb 5.02 TeV



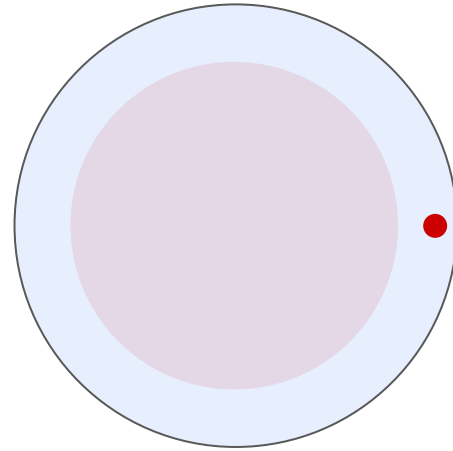
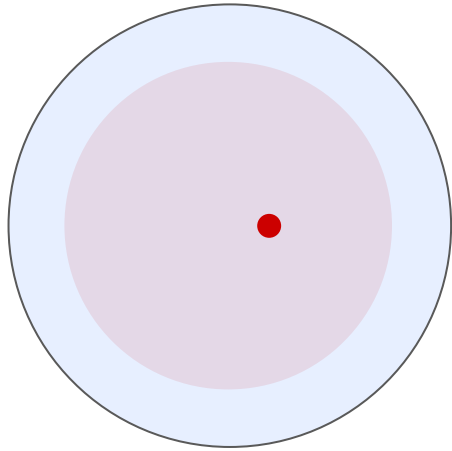
A drop of $\langle n_p \rangle / \langle n_n \rangle$ in the ZDC!

Isobars:

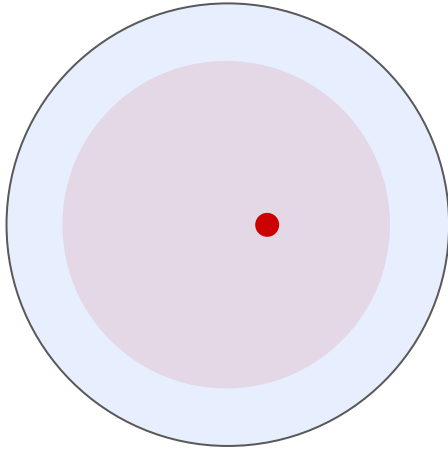
Ratios involving dN_B/dy and dN_Q/dy :

Sensitive to **protons** and **neutron** distribution

Nuclear structure and neutron skin!



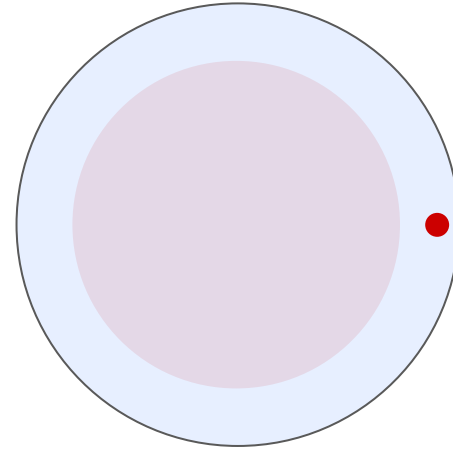
Central collisions



more p-p collisions

More Q/B than baseline

Peripheral collisions



more p-n collisions

Less Q/B than baseline

$$R_{c_1, c_2}(y) = r_{c_1}(y) / r_{c_2}(y)$$

$c_1 > c_2$

$$r_{(c)}(y) = \frac{dN_Q}{dy} / \frac{dN_B}{dy}(c) \times A/Z$$

$$R_{c_1, c_2}(y) \sim 1$$

$$R_{c_1, c_2}(y) < 1$$

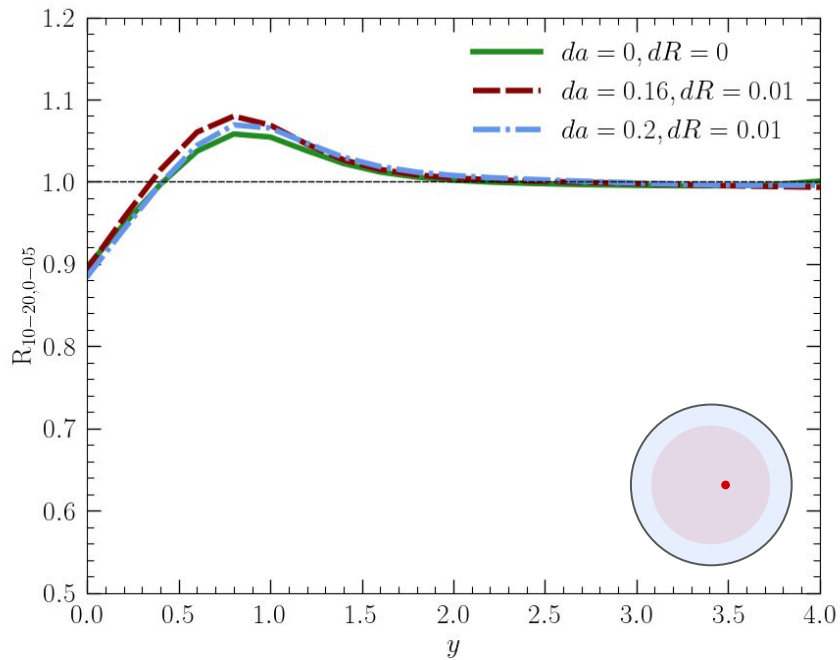
No neutron skin Q depletion effect

Neutron skin Q depletion effect!

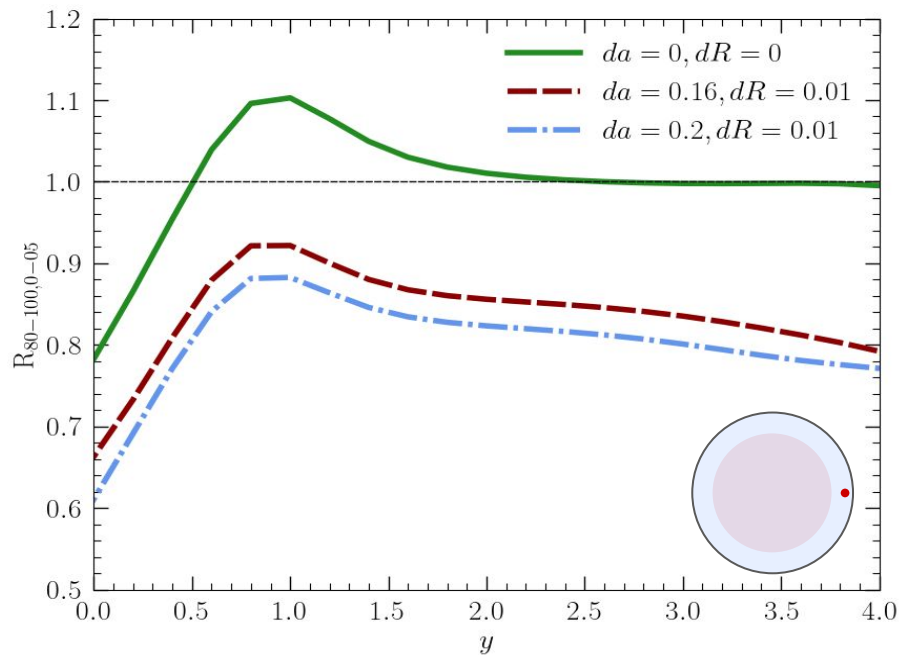
Pb neutron skin in p-Pb at LHC

Initial stage pPb 72 GeV

10-20%



80-100%



Pb neutron skin in p-Pb at LHC

Initial stage pPb 72 GeV

Constraining the Pb neutron skin

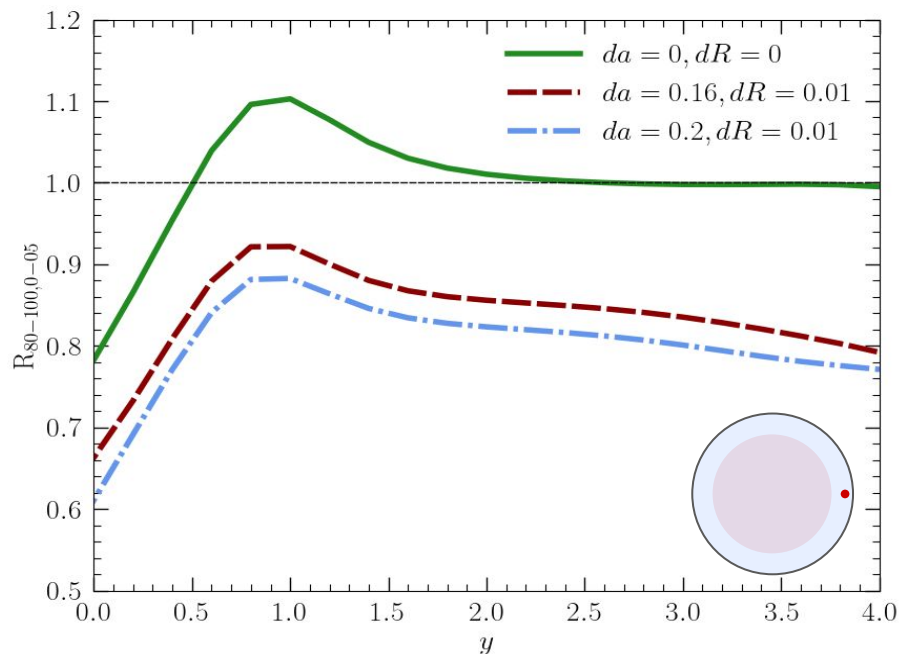
SMOG 2 at LHCb!

The ratio R_{c_1, c_2} could be defined from proxies

net-pion/net-proton

Weaker signal but easier access experimentally!

80-100%



Tracing the conserved charge in heavy-ion collisions:

- Distributions of the conserved charges B and Q

Give access to protons and neutrons difference in
stopping mechanisms

The baryon junction certainly carries the baryon number!

spatial distribution

The nuclear shape and neutron skin can be studied in
heavy-ion collisions in this way

- Future related work:

Predict the pPb at final stage

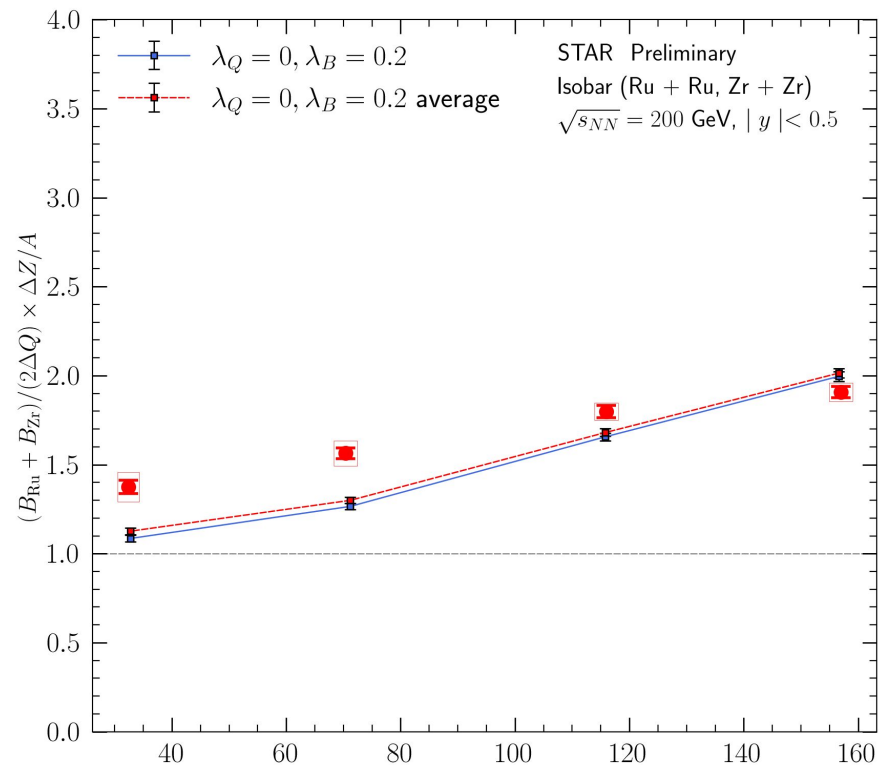
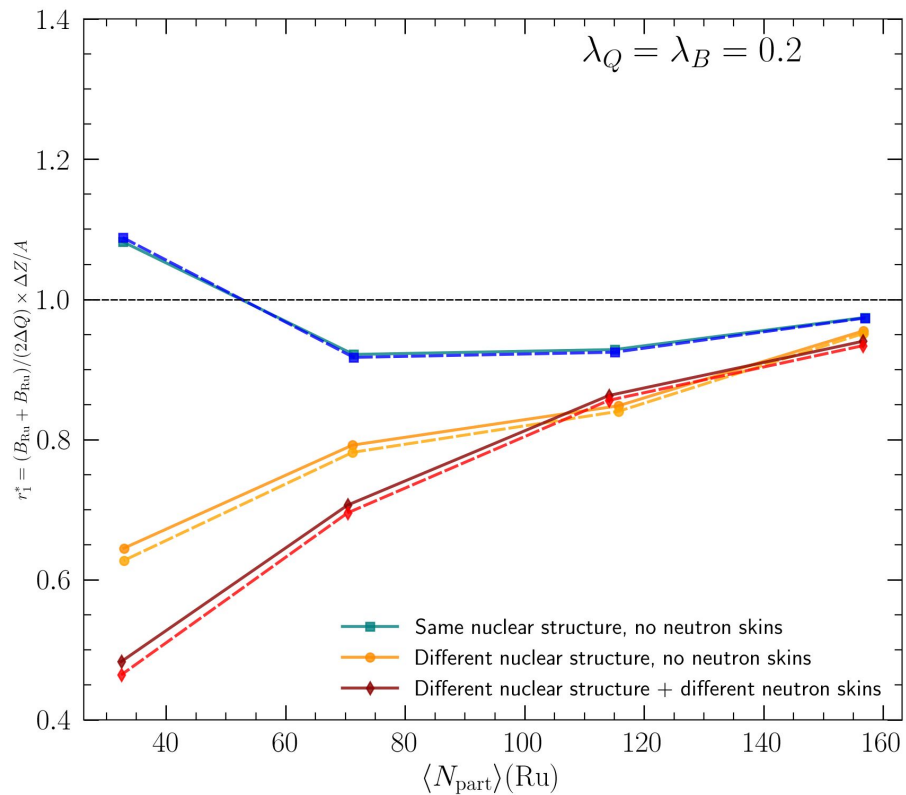
Enhance statistics via machine learning

Study the baryon junction in pPb net-proton number slope!

backup

Comparison with STAR ratio

The ratio is defined as: $r_1^* = (B_{\text{Ru}} + B_{\text{Zr}})/(2\Delta Q) \times \Delta Z/A$



Experimental RuB ratio

$$B/\Delta Q \times \Delta Z/A$$

Net-baryon number:

STAR does not measure neutrons,

Evaluation of neutrons from deuterons yields via HRG model

$$N_B = (N_p - N_{\bar{p}}) + (N_n - N_{\bar{n}}) \approx (N_p - N_{\bar{p}}) + \bar{p} \sqrt{\frac{d}{\bar{d}}} - p \sqrt{\frac{\bar{d}}{d}}$$

STAR Collaboration, Phys Rev.99.064905

Net-charge difference:

The electric charge is a non-trivial measurement at mid-rapidity (small yields!).

Making use of the convenient double ratios to cancel uncertainties accessible in isobar collisions.

$$\Delta Q = [(N_{\pi}^+ + N_K^+ + N_p) - (N_{\pi}^- + N_K^- + N_{\bar{p}})]_{\text{Ru}} - []_{\text{Zr}}$$

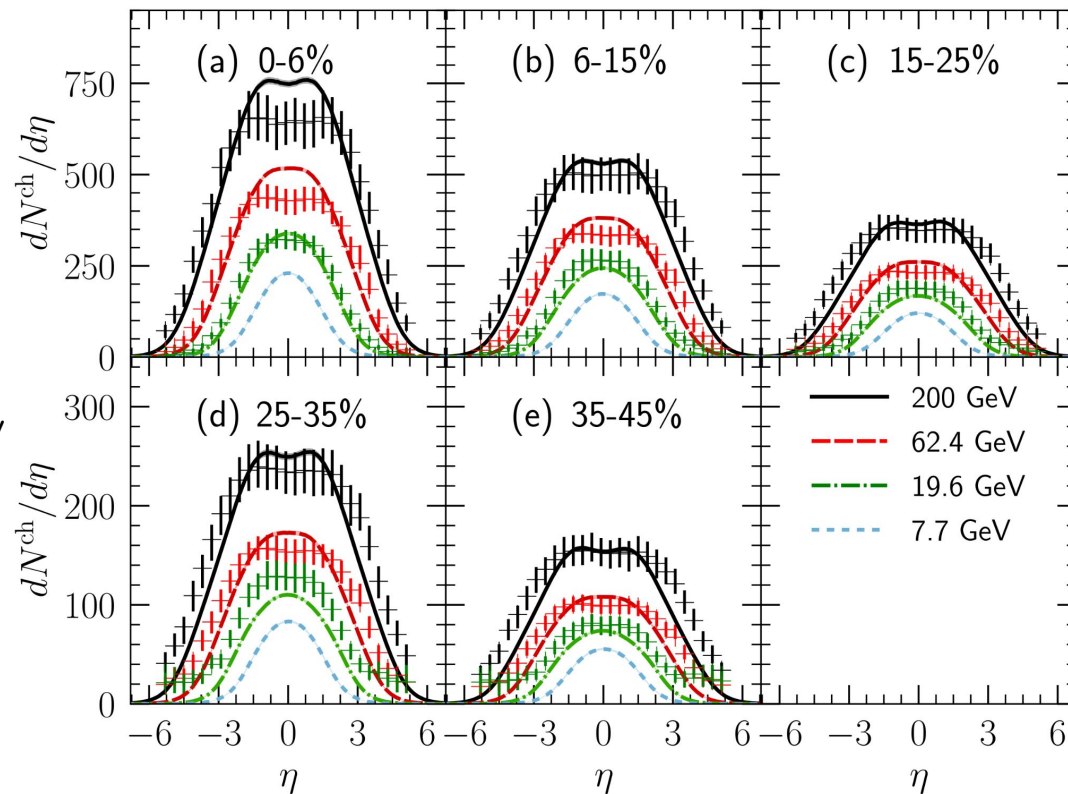
$$R2_{\pi} = \frac{(N_{\pi}^+/N_{\pi}^-)_{\text{Ru}}}{(N_{\pi}^+/N_{\pi}^-)_{\text{Zr}}} \approx 1 + (N_{\pi}^+ - N_{\pi}^-)_{\text{Ru}} - (N_{\pi}^+ - N_{\pi}^-)_{\text{Zr}}$$

$$\Delta Q = N_{\pi}(R2_{\pi} - 1) + N_K(R2_K - 1) + N_p(R2_p - 1)$$

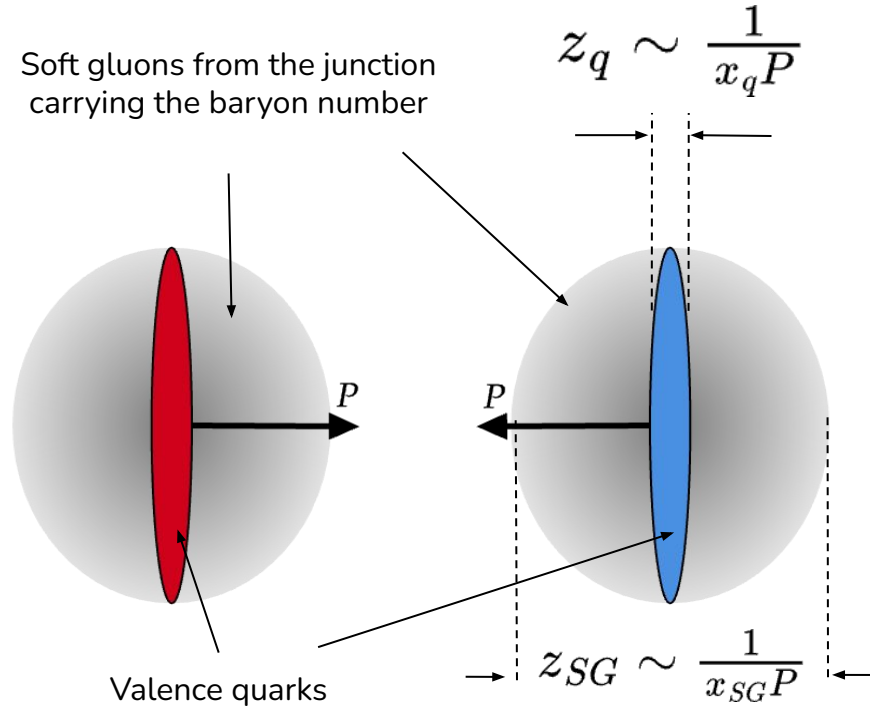
MUSIC tuning on PHOBOS Au+Au data

Current version:

- Tuned on charged particle rapidity distributions for Au+Au collisions at RHIC PHOBOS
- Overestimate yields at mid-rapidity for most central collision
- Overall good agreement ✓



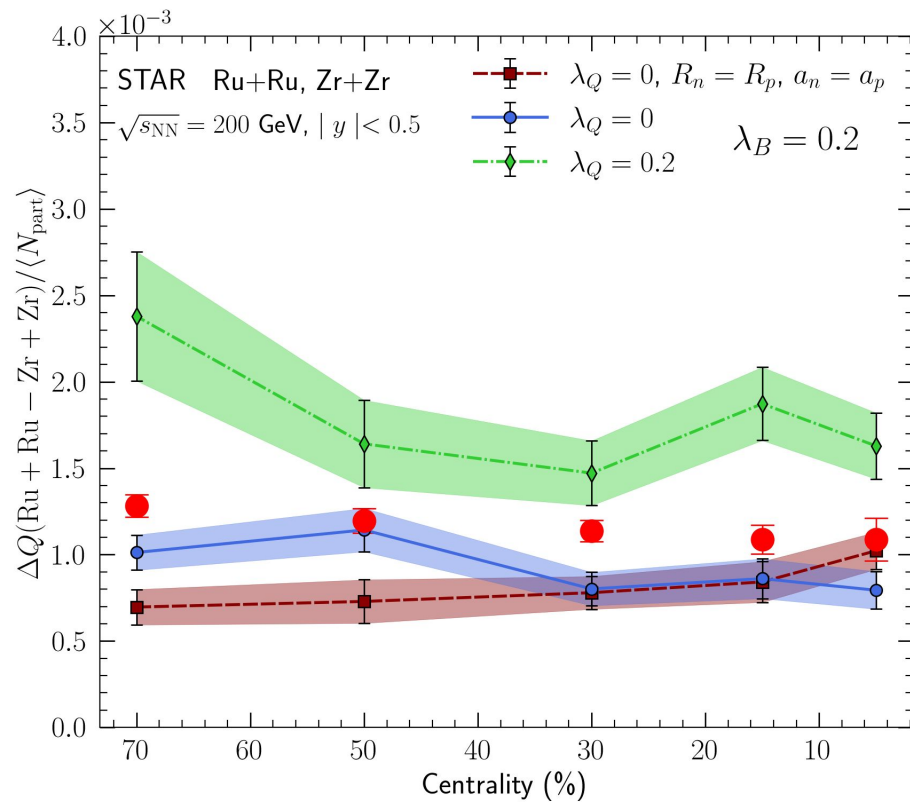
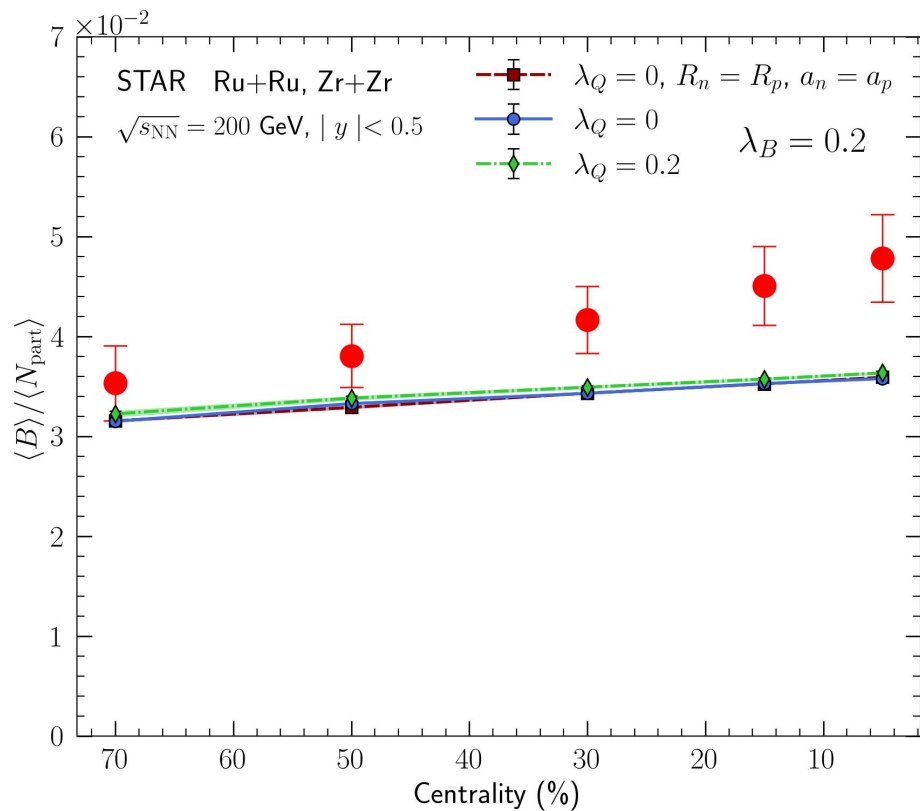
Backup: Gluon cloud interpretation



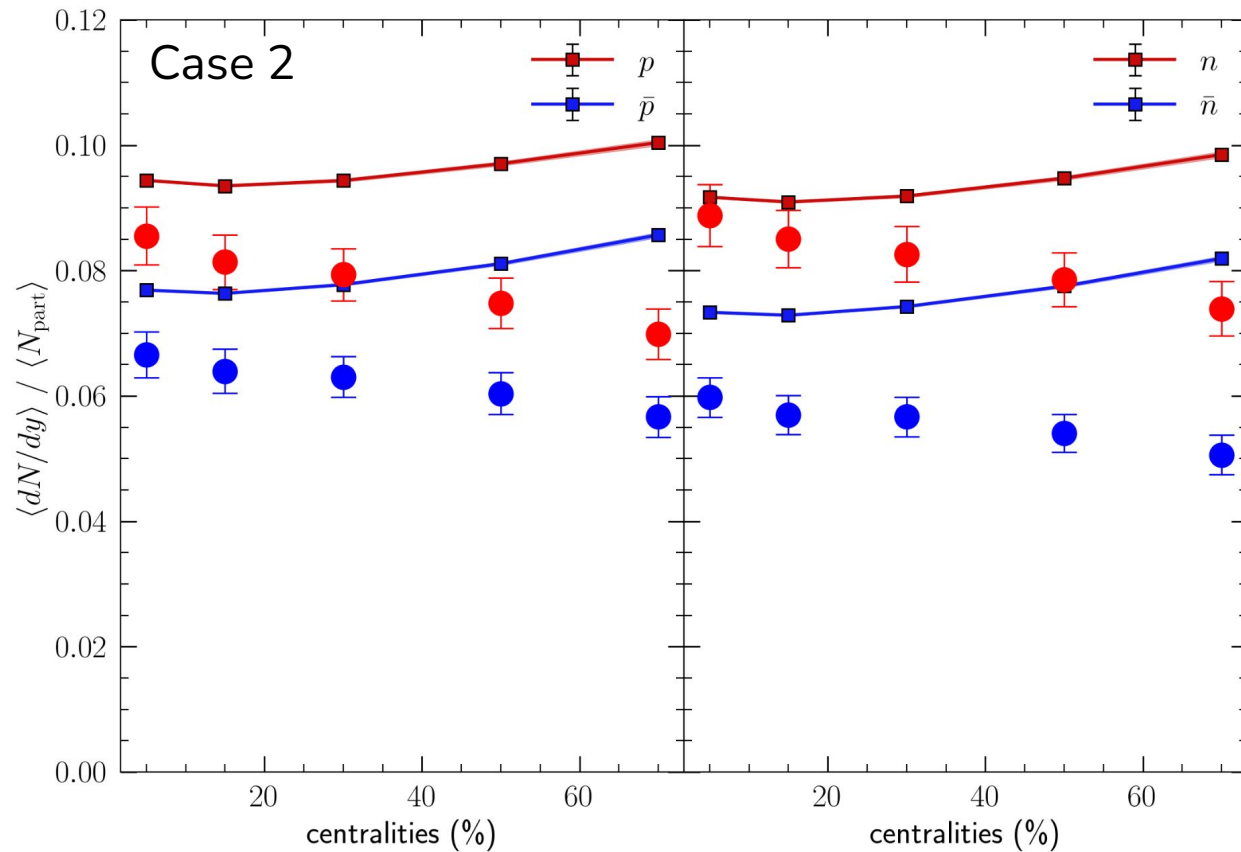
Quarks: $x_q \sim 1/3$
interaction time is short
 $\sim 6 \cdot 10^{-3}$ fm

Soft gluons: $x_{SG} \ll x_q$
interaction time is large
 ~ 1 fm

**The baryon number
is stopped!**



Proton, neutron yield vs STAR measurement



A fine tuning of λ_B and λ_Q

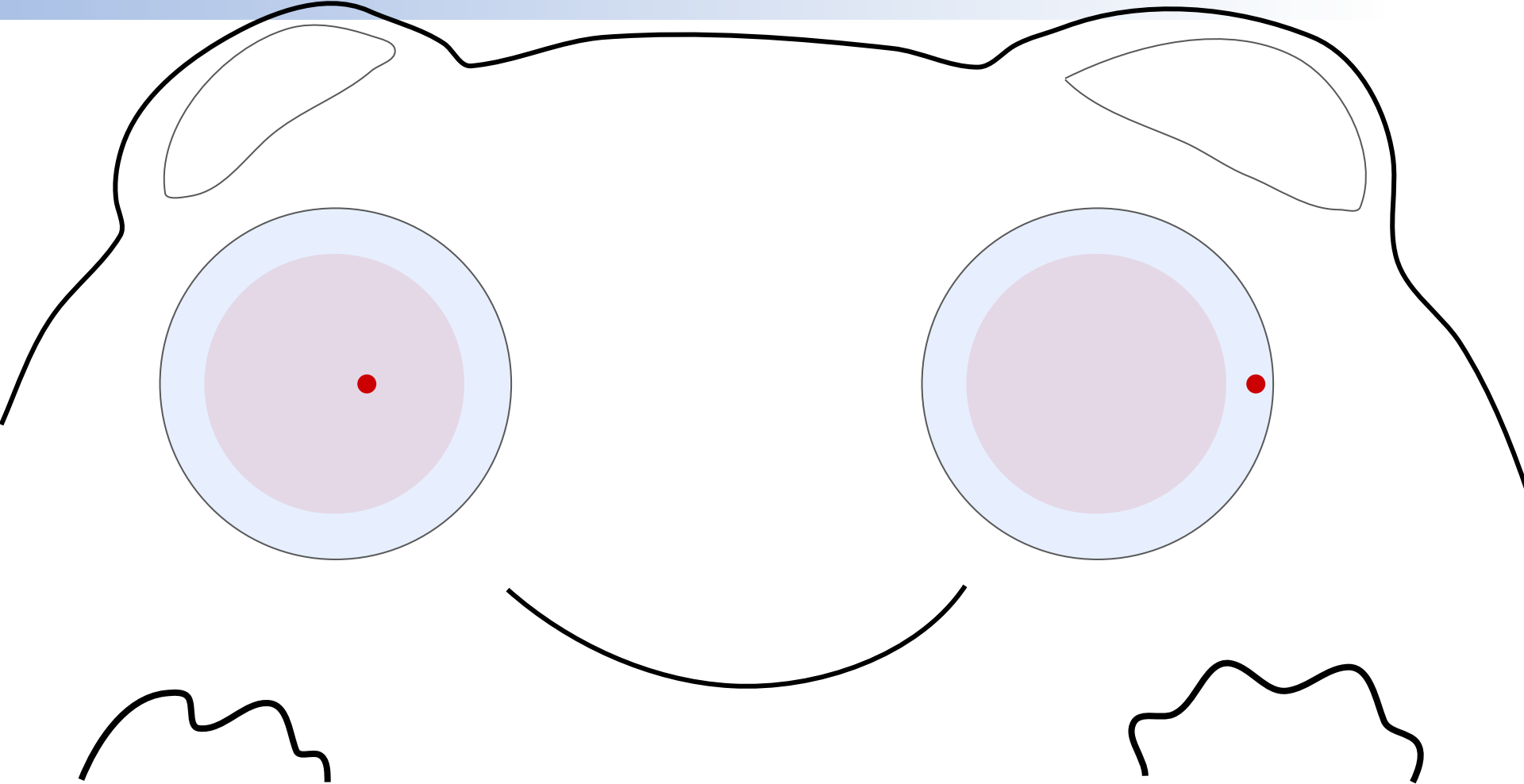
Overall amplitude can be reduced

A fine tuning of freeze out energy density

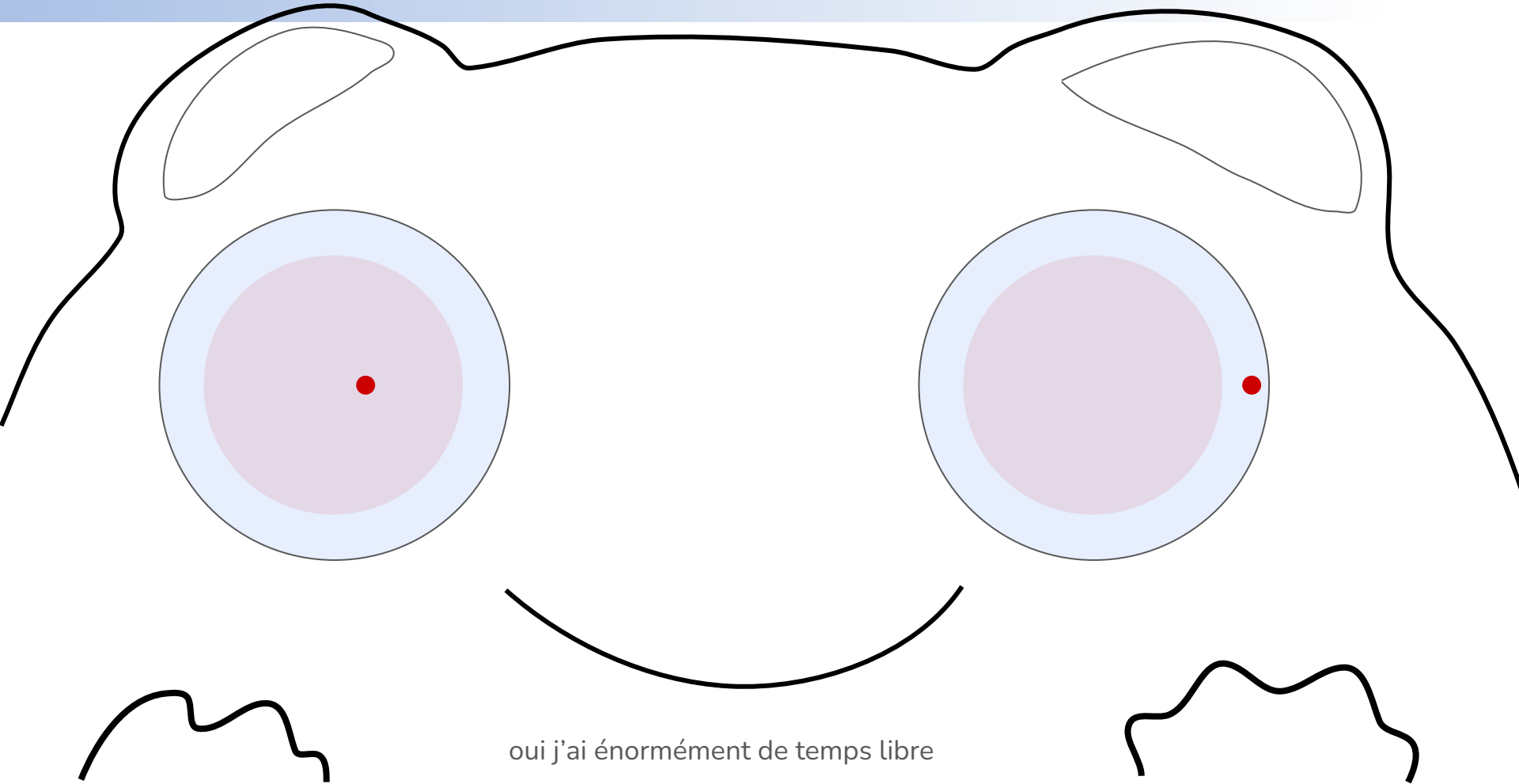
Difference in particles and antiparticles can be fitted

Left for future study

Pilède the kind Monster



Pilède the kind Monster



oui j'ai énormément de temps libre