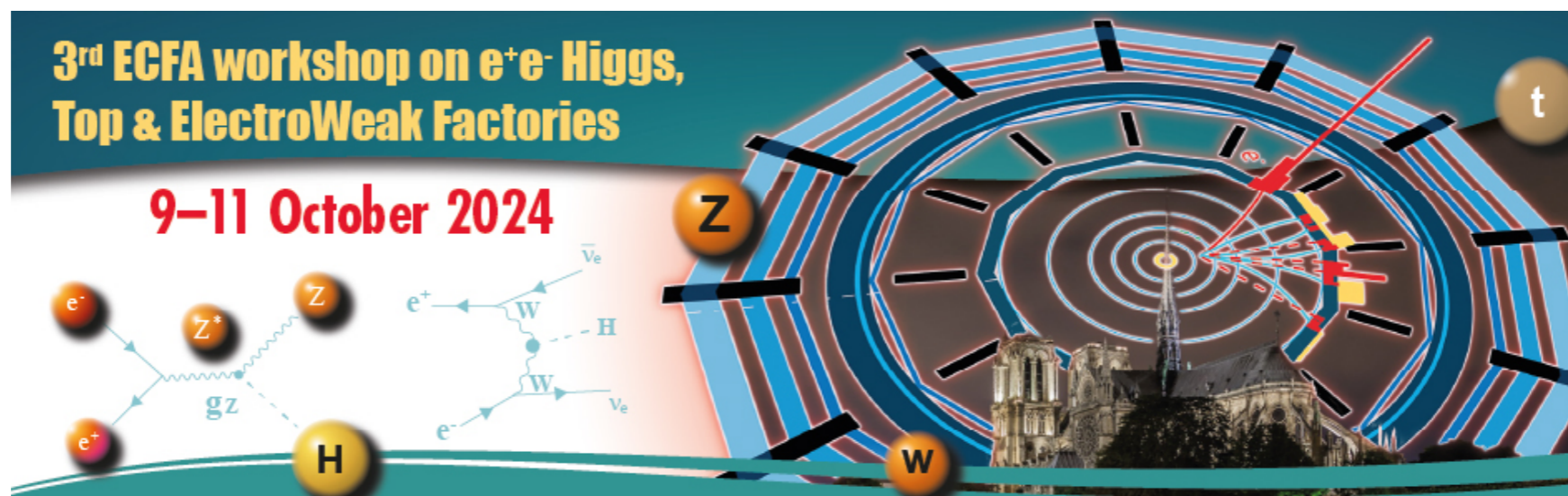


# Updates from the SMEFiT collaboration

## towards the next ESPPU

Based on JHEP 09 (2024) 091

with E. Celada, T. Giani, L. Mantani, J. Rojo, A. N. Rossia, M. Thomas, E. Vryonidou



Jaco ter Hoeve

09/10/24

# The high energy landscape

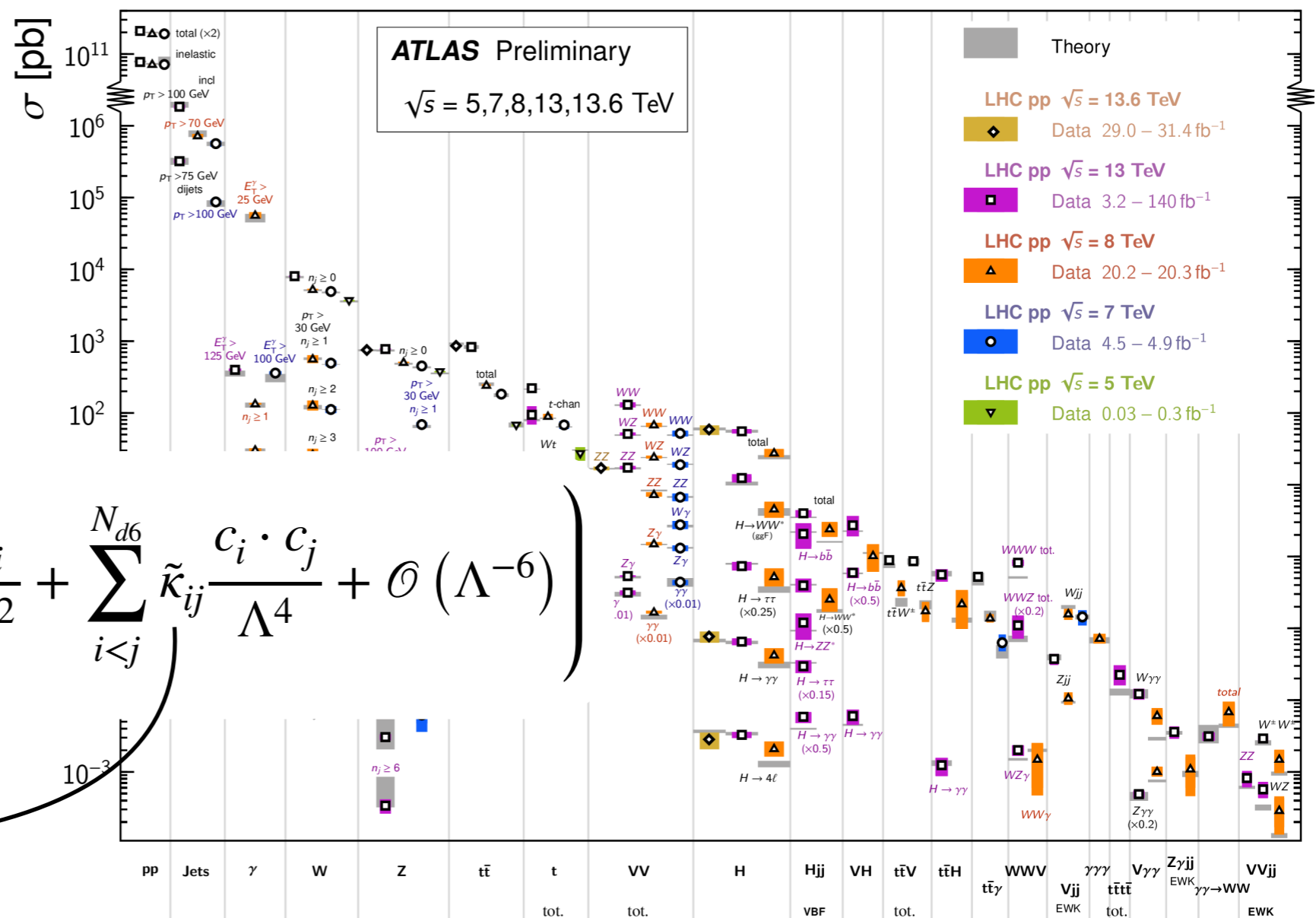
Lots of impressive cross-section measurements, but no clear deviation from the SM (yet) ...

... so we study their overall pattern!

[ATL-PHYS-PUB-2023-039]

Status: October 2023

Standard Model Production Cross Section Measurements



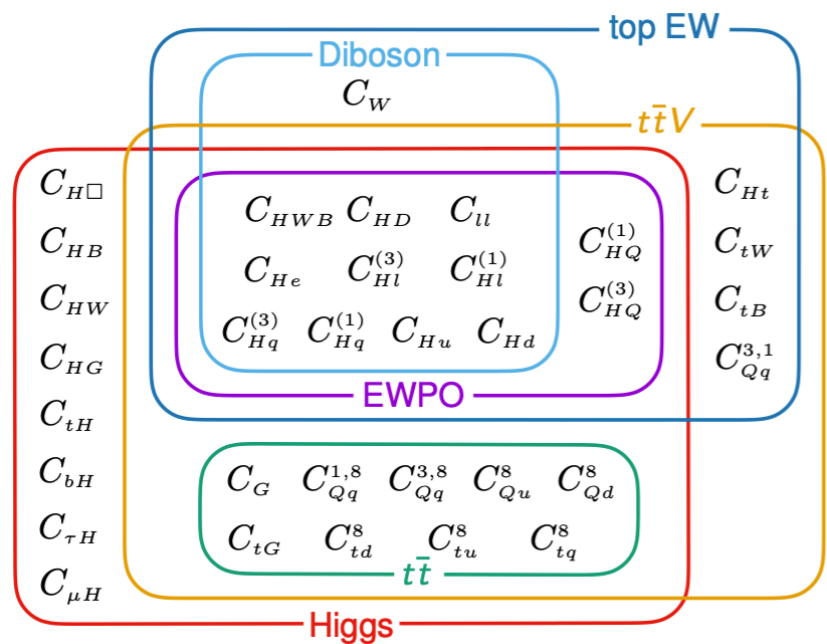
Linear EFT corrections:  
interference SM-EFT<sub>d6</sub>  
@NLO QCD

$$\sigma(c, \Lambda) = \sigma_{\text{SM}} \times \left( 1 + \sum_i^{N_{d6}} \kappa_i \frac{c_i}{\Lambda^2} + \sum_{i<j}^{N_{d6}} \tilde{\kappa}_{ij} \frac{c_i \cdot c_j}{\Lambda^4} + \mathcal{O}(\Lambda^{-6}) \right)$$

Quadratic EFT  
corrections:  
EFT<sub>d6</sub>-EFT<sub>d6</sub>  
@NLO QCD

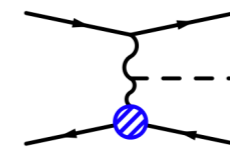
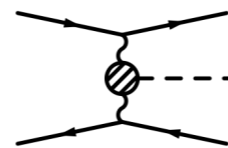
# Why global SMEFT fits?

- ▶ The SMEFT is our **universal** tool to search for BSM physics above the EW scale, with **minimal assumptions** on what it may look like
- ▶ Given the **cross-talk** between Higgs, top, diboson and EWPO (and flavour and low energy observables), a simultaneous fit is our only way forward
- ▶ **Challenge:** a large number of operators, with many datasets needed to break degeneracies

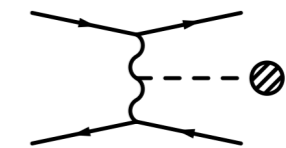


[2012.02779] Fitmaker collaboration

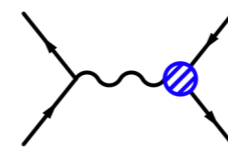
One observable can be influenced by many operators



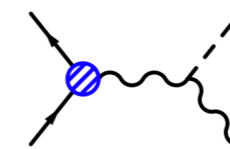
Higgs decay



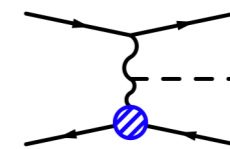
One operator can contribute to many different observables



$e^+e^- \rightarrow f\bar{f}$



Zh production



Weak boson fusion  
Higgs production

Anke Biekötter - HET seminar Brookhaven

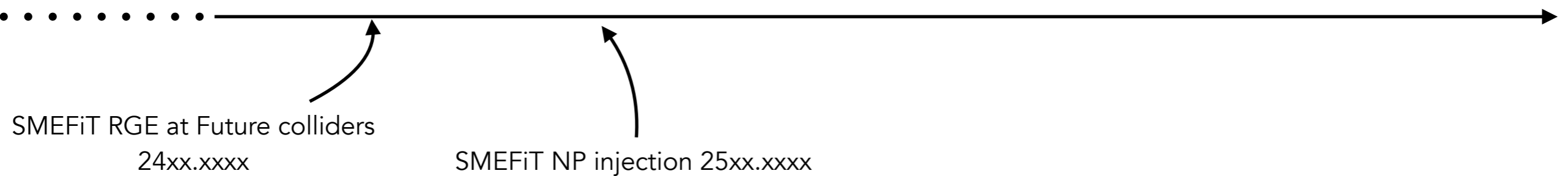
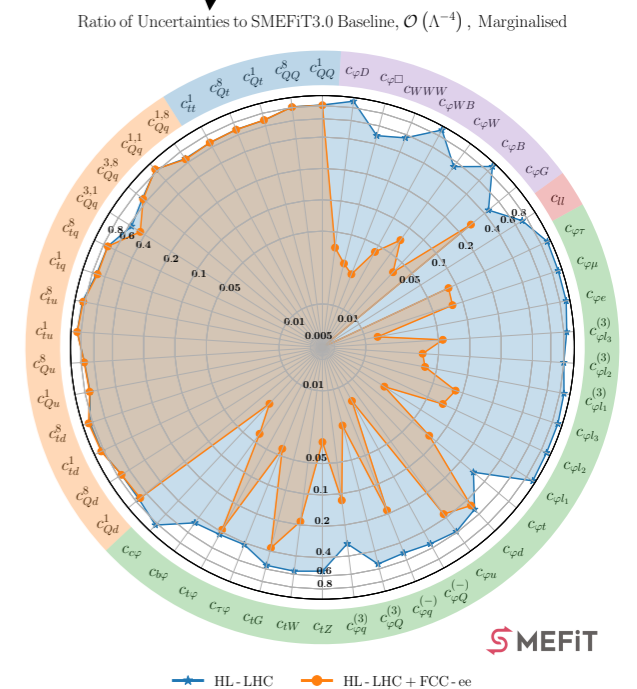
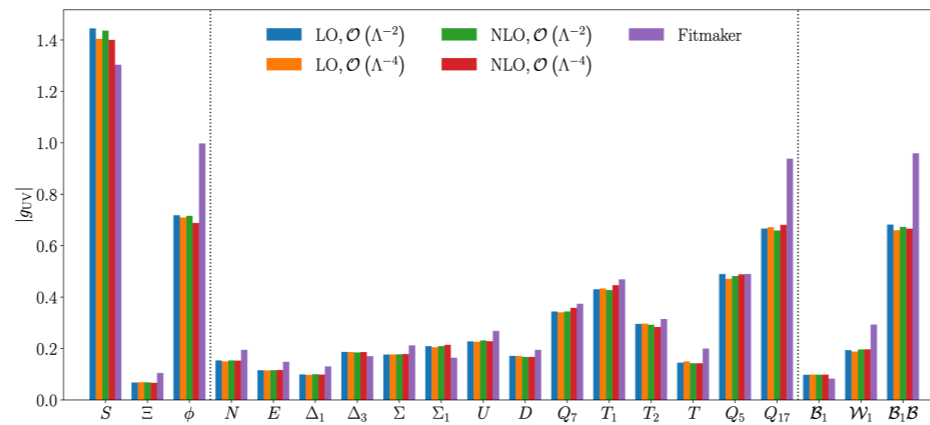
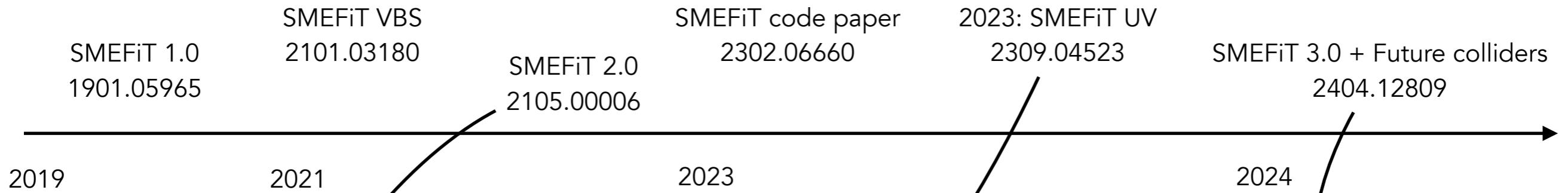
# The SMEFiT3.0 framework

E. Celada, T. Giani, L. Mantani, J. Rojo, A. Rossia, M. Thomas, E. Vryonidou , JtH

*JHEP 09 (2024) 091 [2404.12809]*



# The SMEFiT timeline



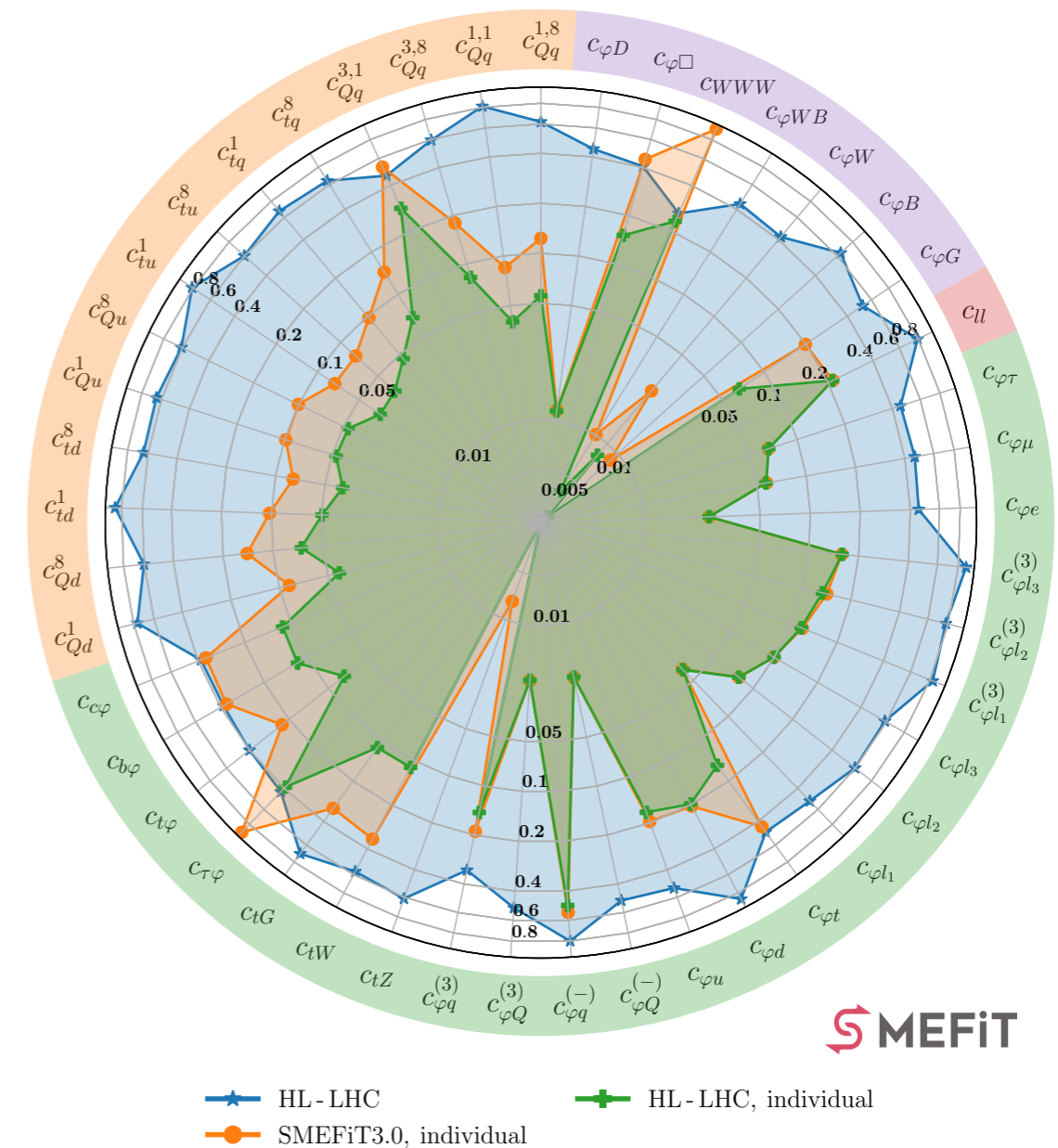
# SMEFiT3.0 in a nutshell

- ✓ SMEFiT2.0 extended with recent datasets in **top, diboson and Higgs production** based on the full Run II luminosity
- ✓ Full independent treatment of the **EWPOs** from LEP and SLD
- ✓ Dedicated **projection module** to extrapolate Run II data to HL-LHC
- ✓ **FCC-ee and CEPC pseudodata** from Snowmass predictions, updated to 4 IPs as per the FCC feasibility midterm report [2206.08326], CERN/3789/RA
- ✓ Both results in terms of Wilson coefficients and **UV-complete models**
- ✓ **Public code, data and theory:** results are fully reproducible



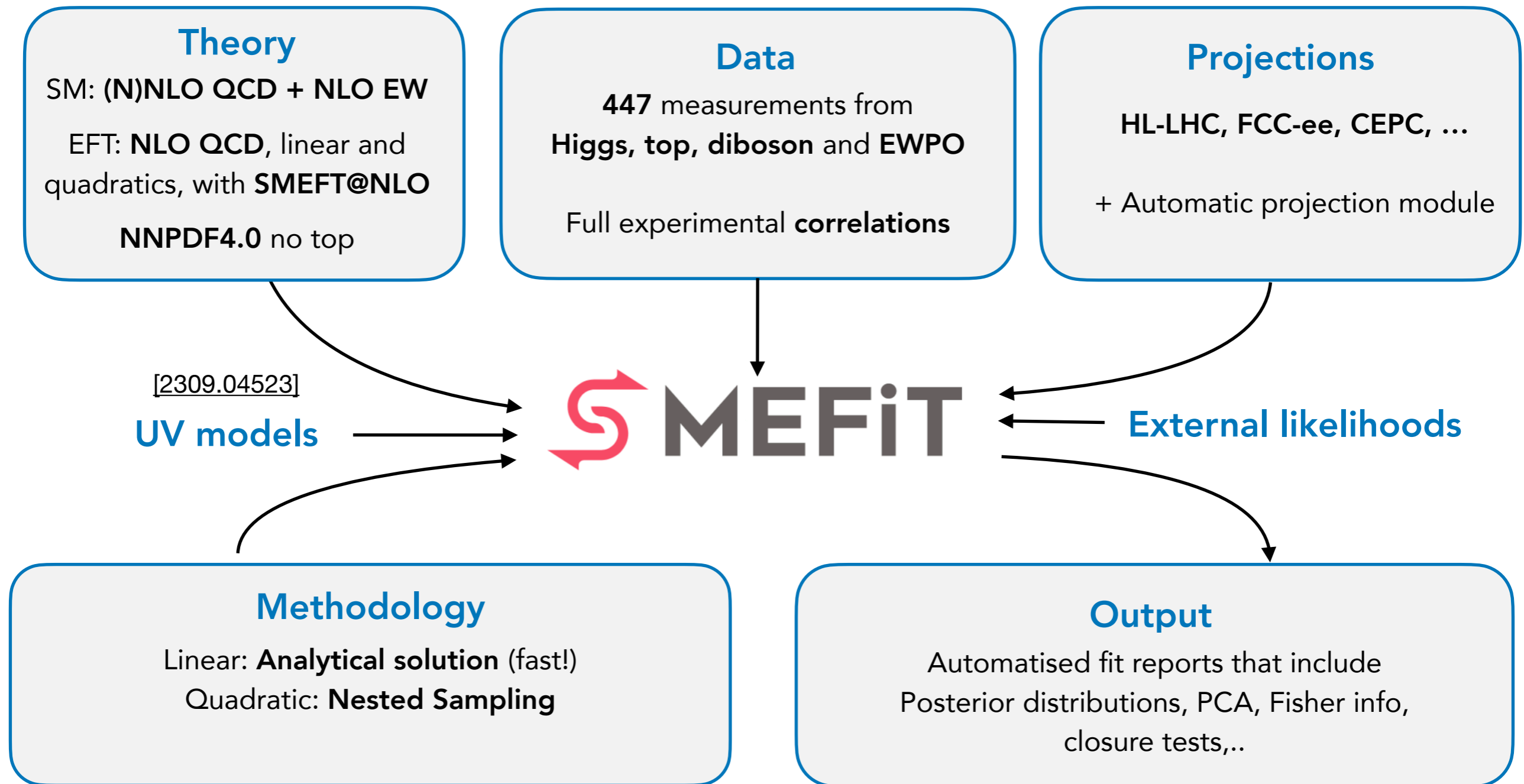
[lhcfitnikhef.github.io/smefit\\_release](https://lhcfitnikhef.github.io/smefit_release)

Ratio of Uncertainties to SMEFiT3.0 Baseline,  $\mathcal{O}(\Lambda^{-2})$ , Marginalised



"Spider plots / Antarctica plots"

# SMEFiT under the hood



# Full treatment of EWPOs

- ▶ In the SMEFT, the SM couplings receive corrections from dim-6 operators

$$\begin{aligned}
 \delta g_V^{l_i} &= \delta \bar{g}_Z \bar{g}_V^{l_i} + Q^{l_i} \delta s_\theta^2 + \Delta_V^{l_i} = 0, \quad i = 1, 2, 3, \\
 \delta g_A^{l_i} &= \delta \bar{g}_Z \bar{g}_A^{l_i} + \Delta_A^{l_i} = 0, \quad i = 1, 2, 3, \\
 \delta g_V^u &= \delta \bar{g}_Z \bar{g}_V^u + Q^u \delta s_\theta^2 + \Delta_V^u = 0, \\
 \delta g_A^u &= \delta \bar{g}_Z \bar{g}_A^u + \Delta_A^u = 0, \\
 \delta g_V^d &= \delta \bar{g}_Z \bar{g}_V^d + Q^d \delta s_\theta^2 + \Delta_V^d = 0, \\
 \delta g_A^d &= \delta \bar{g}_Z \bar{g}_A^d + \Delta_A^d = 0, \\
 \delta g_V^{W,l_i} &= \frac{c_{ll} + 2c_{\varphi l_i}^{(3)} - c_{\varphi l_1}^{(3)} - c_{\varphi l_2}^{(3)}}{4\sqrt{2}G_F} = 0, \quad i = 1, 2, 3, \\
 \delta g_V^{W,q} &= \frac{c_{ll} + c_{\varphi q}^{(3)} - c_{\varphi l_1}^{(3)} - c_{\varphi l_2}^{(3)}}{4\sqrt{2}G_F} = 0,
 \end{aligned}$$

$$\begin{pmatrix} c_{\varphi l_i}^{(3)} \\ c_{\varphi l_i}^{(1)} \\ c_{\varphi e/\mu/\tau} \\ c_{\varphi q}^{(-)} \\ c_{\varphi q}^{(3)} \\ c_{\varphi u} \\ c_{\varphi d} \\ c_{ll} \end{pmatrix} = \begin{pmatrix} -\frac{1}{t_W} & -\frac{1}{4t_W^2} \\ 0 & -\frac{1}{4} \\ 0 & -\frac{1}{2} \\ \frac{1}{t_W} & \frac{1}{4s_W^2} - \frac{1}{6} \\ -\frac{1}{t_W} & -\frac{1}{4t_W^2} \\ 0 & \frac{1}{3} \\ 0 & -\frac{1}{6} \\ 0 & 0 \end{pmatrix} \begin{pmatrix} c_{\varphi WB} \\ c_{\varphi D} \end{pmatrix}$$

- ▶ **SMEFiT2.0:** assumed measurements at LEP were precise enough to set the coupling shifts to zero: 14 constraints, 16 d.o.f
- ▶ **SMEFiT3.0:** hardwired constraints get no longer imposed, EWPOs are treated on the same footing as (existing) LHC data: 14 **extra** d.o.f



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SMEFiT3.0 is simultaneously sensitive to **45 (50) Wilson coefficients** at the linear (quadratic) level!

# Dataset upgrade

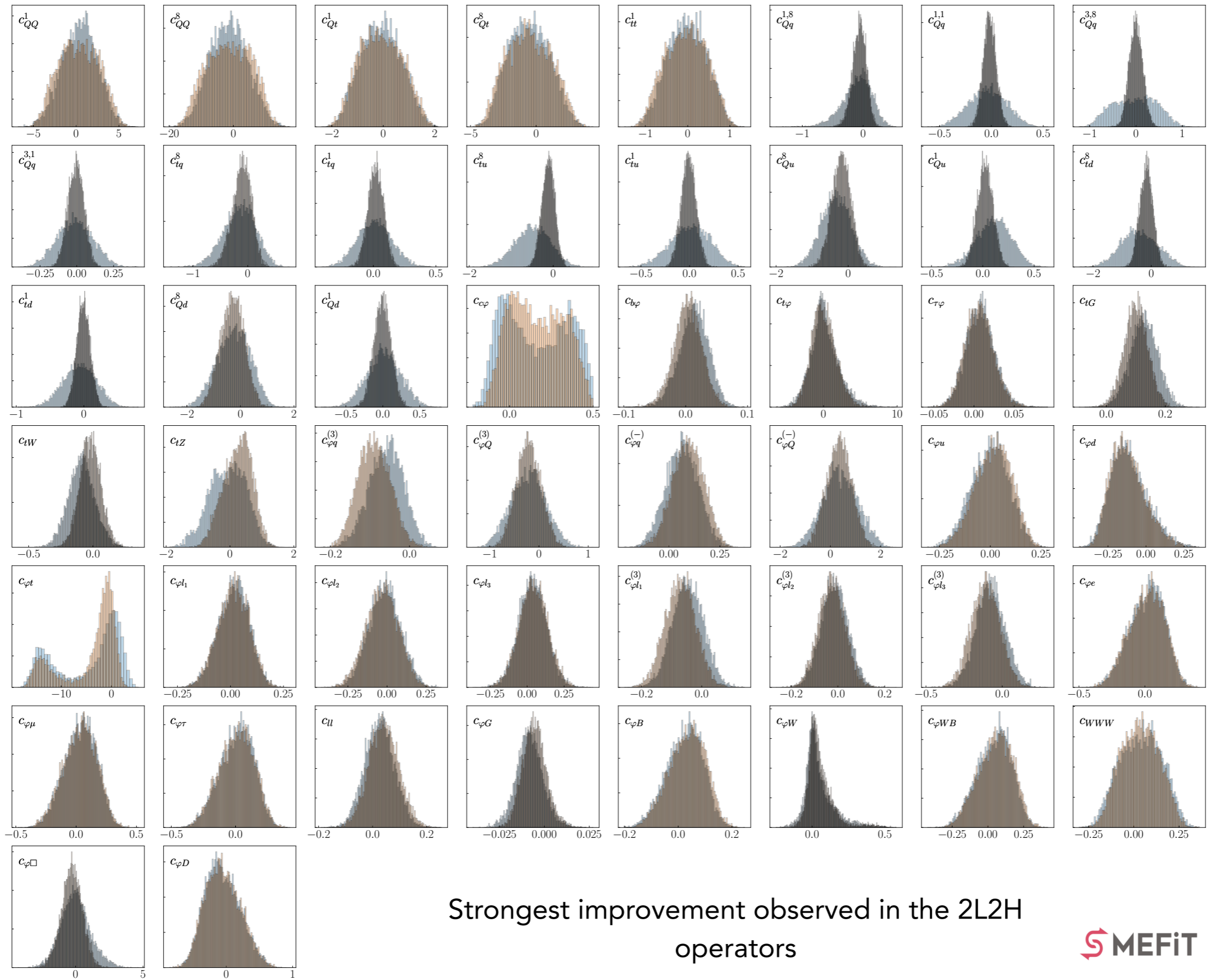
We extended SMEFiT2.0 with recent Run II datasets from top, diboson and Higgs production

Category	Processes	$n_{\text{dat}}$	
		SMEFiT2.0	SMEFiT3.0
Top quark production	$t\bar{t} + X$	94	115
	$t\bar{t}Z, t\bar{t}W$	14	21
	$t\bar{t}\gamma$	-	2
	single top (inclusive)	27	28
	$tZ, tW$	9	13
	$t\bar{t}t\bar{t}, t\bar{t}b\bar{b}$	6	12
	<b>Total</b>	<b>150</b>	<b>189</b>
Higgs production and decay	Run I signal strengths	22	22
	Run II signal strengths	40	40
	Run II, differential distributions & STXS	35	71
	<b>Total</b>	<b>97</b>	<b>133</b>
Diboson production	LEP-2	40	40
	LHC	30	41
	<b>Total</b>	<b>70</b>	<b>81</b>
Z-pole EWPOs	LEP-2	-	44
Baseline dataset	<b>Total</b>	<b>317</b>	<b>449</b>

Flavour assumption:  $U(2)_q \times U(3)_d \times U(2)_u \times (U(1)_\ell \times U(1)_e)^3$

■ SMEFiT2.0 Dataset, NLO  $\mathcal{O}(\Lambda^{-4})$

■ SMEFiT3.0 Dataset, NLO  $\mathcal{O}(\Lambda^{-4})$

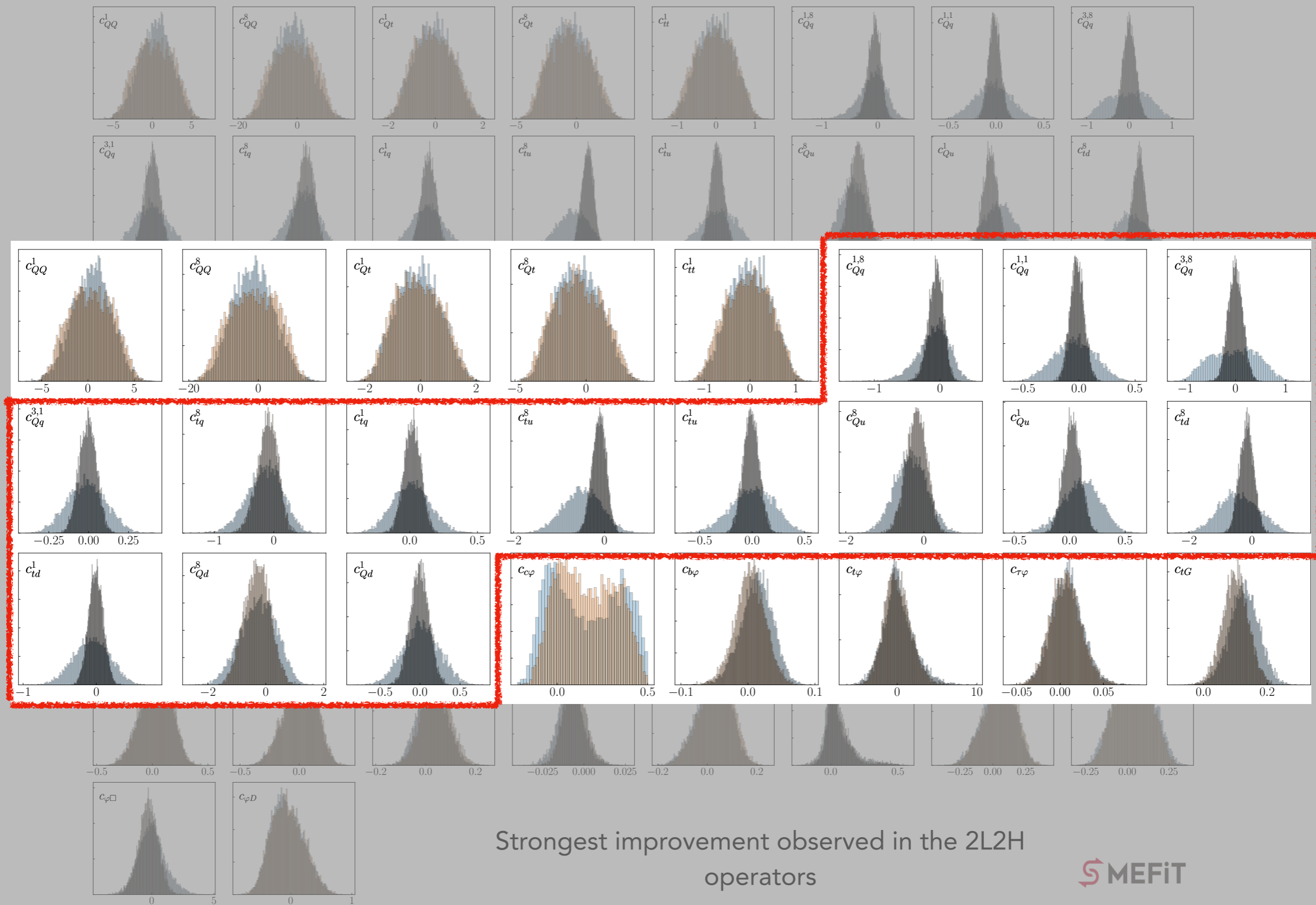


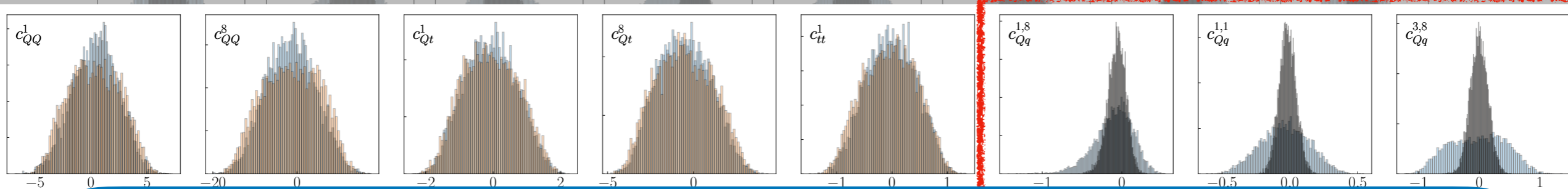
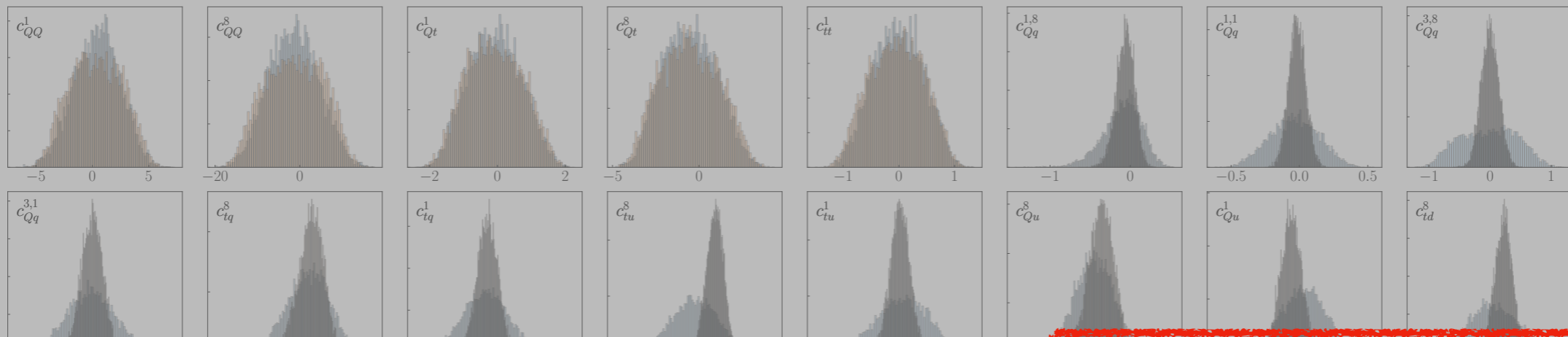
Strongest improvement observed in the 2L2H operators



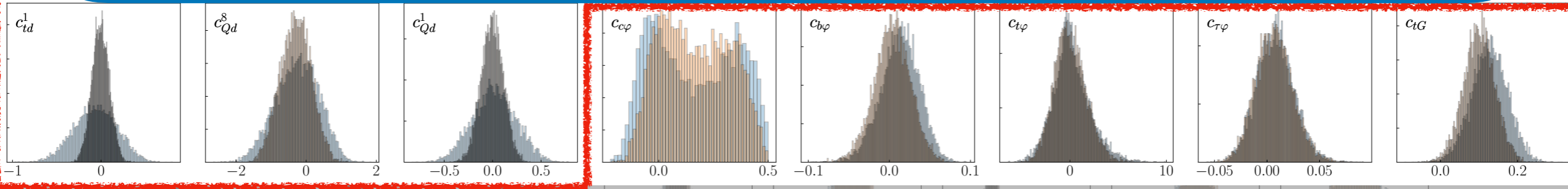
SMEFiT2.0 Dataset, NLO  $\mathcal{O}(\Lambda^{-4})$

SMEFiT3.0 Dataset, NLO  $\mathcal{O}(\Lambda^{-4})$





How will this pattern change at HL-LHC and future colliders?



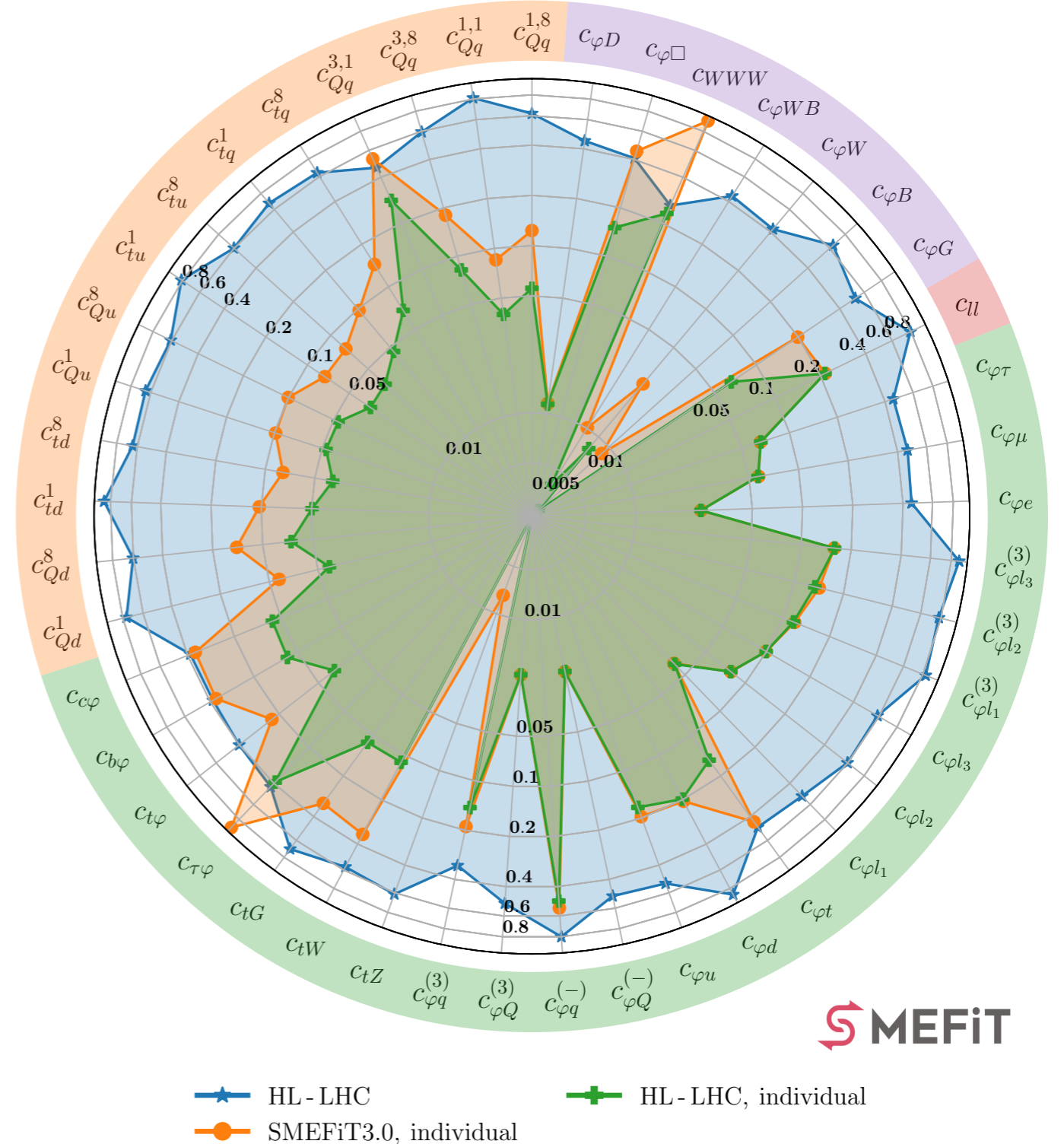
Strongest improvement observed in the 2L2H operators



# Result: HL-LHC

- ▶ We project all RunII datasets from the SMEFiT 3.0 baseline: one for each process and final state [see backup for details](#)
- ▶ We see an improvement ranging from 20 to 70 % in the marginalised fit
- ▶ The EW operators only improve in the marginalised fit because of correlations

Ratio of Uncertainties to SMEFiT3.0 Baseline,  $\mathcal{O}(\Lambda^{-2})$ , Marginalised



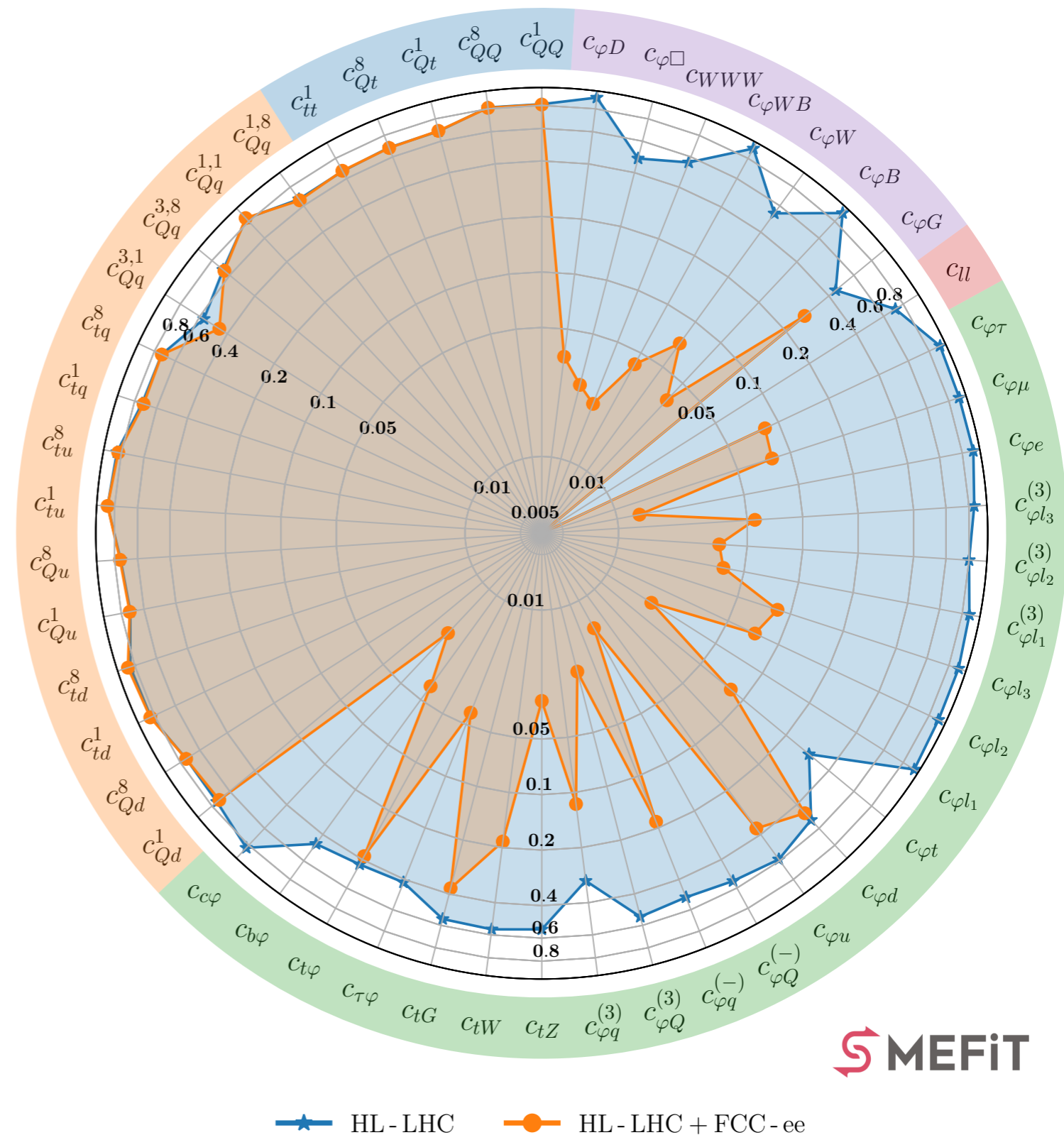
# Result: FCC-ee

Ratio of Uncertainties to SMEFiT3.0 Baseline,  $\mathcal{O}(\Lambda^{-4})$ , Marginalised

## Dataset input

- ▶ EWPOs at the Z-pole
- ▶ Light fermion pair prediction
- ▶ Higgstrahlung and VBF
- ▶ Gauge boson pair production
- ▶ Top-quark pair production
- ▶ Optimal Observables

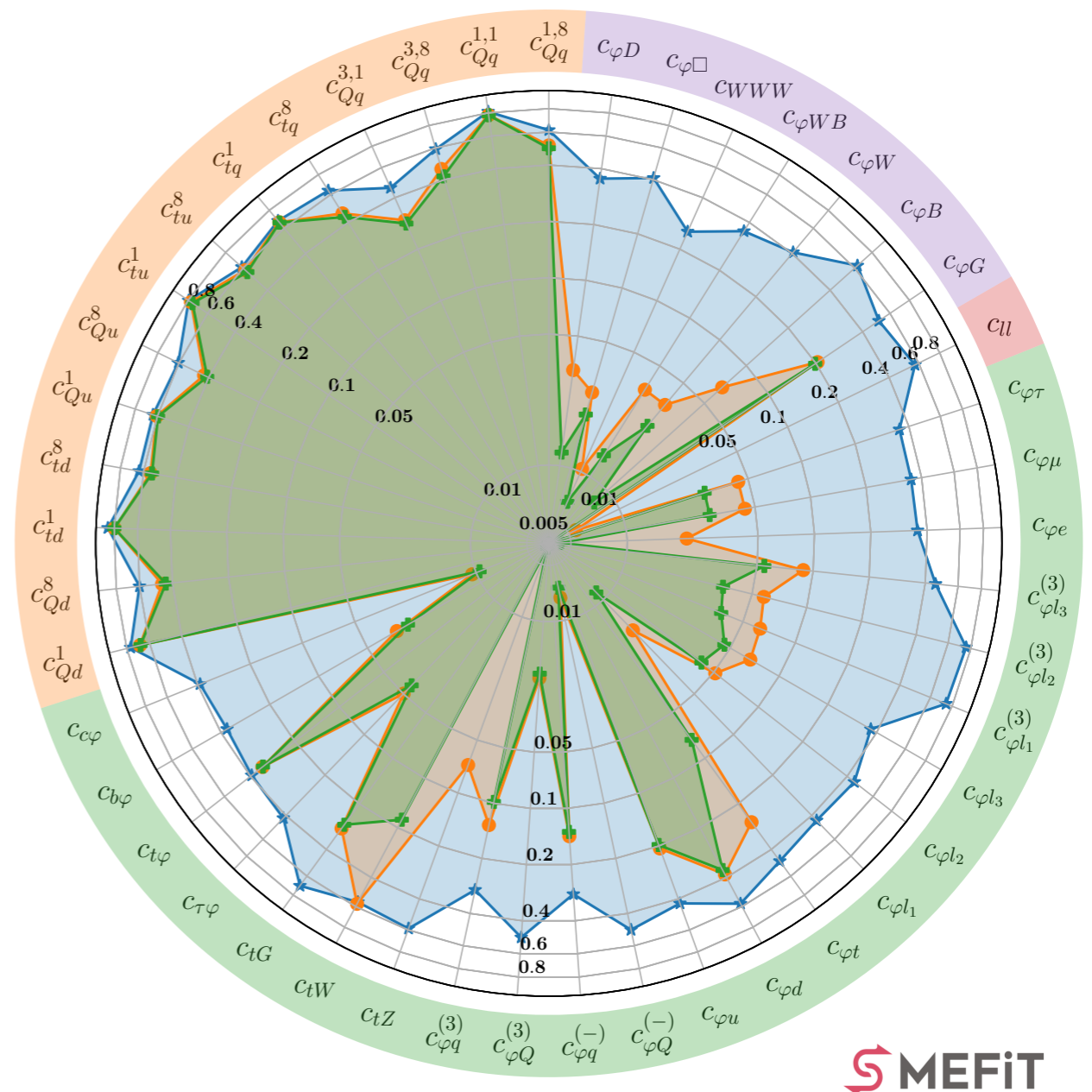
Energy ( $\sqrt{s}$ )	$\mathcal{L}_{\text{int}}$ (Run time)	
	FCC-ee	CEPC
91 GeV (Z-pole)	300 $\text{ab}^{-1}$ (4 years)	100 $\text{ab}^{-1}$ (2 years)
161 GeV ( $2m_W$ )	20 $\text{ab}^{-1}$ (2 years)	6 $\text{ab}^{-1}$ (1 year)
240 GeV	10 $\text{ab}^{-1}$ (3 years)	20 $\text{ab}^{-1}$ (10 years)
350 GeV	0.4 $\text{ab}^{-1}$ (1 years)	-
365 GeV ( $2m_t$ )	3 $\text{ab}^{-1}$ (4 years)	1 $\text{ab}^{-1}$ (5 years)



# Result: FCC-ee energy breakdown

Ratio of Uncertainties to SMEFiT3.0 Baseline,  $\mathcal{O}(\Lambda^{-2})$ , Marginalised

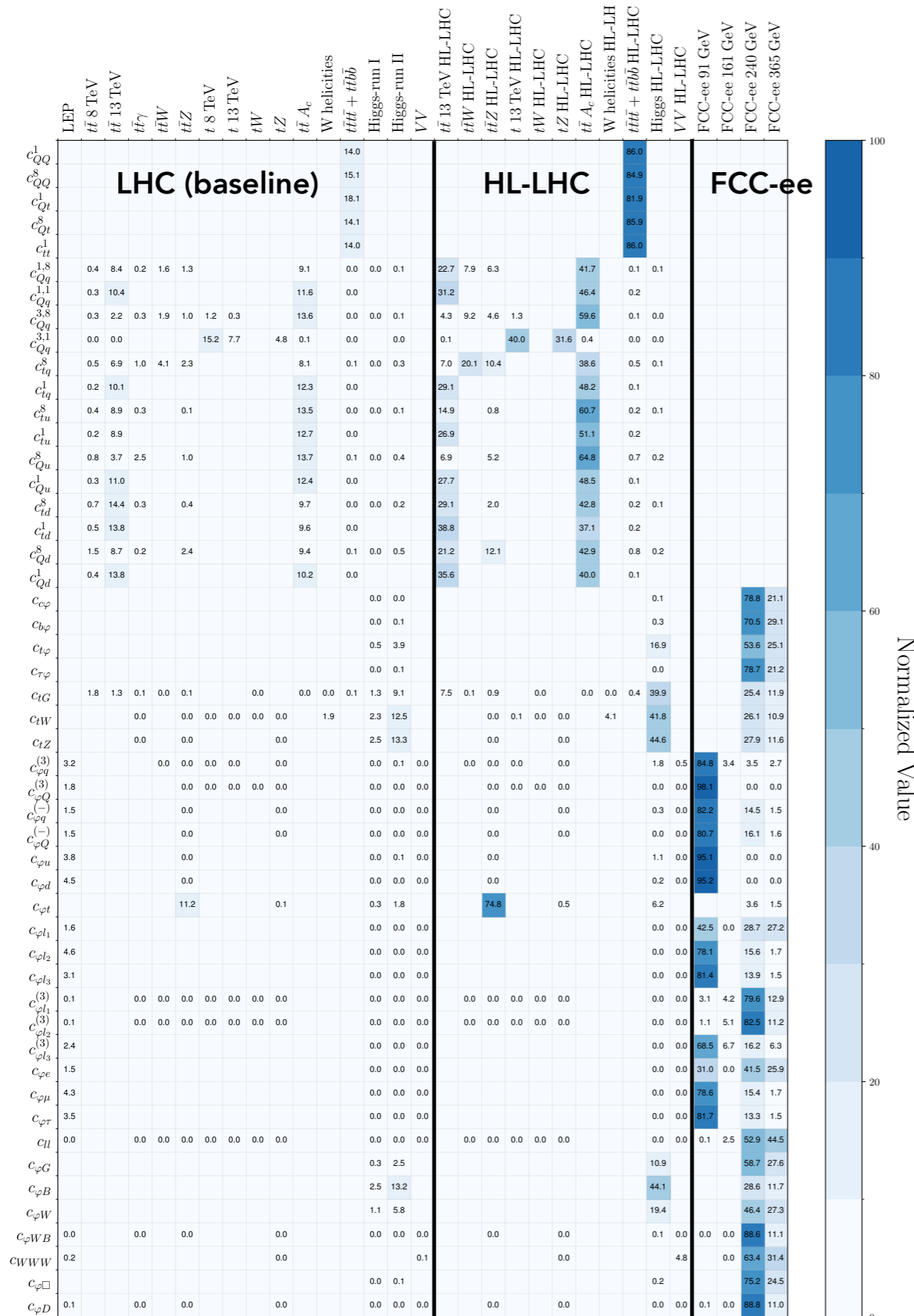
- ▶ The FCC-ee plans to operate **sequentially**, hence we need to study the impact at the various energies
- ▶ Largest impact for Z-pole at 91 GeV plus the Higgs factory run at 240 GeV
- ▶ We can try other combinations too in order to find the most optimal run order for the SMEFT.



- ★ HL - LHC + FCC - ee (91 GeV)
- HL - LHC + FCC - ee (91 + 240 GeV)
- ✚ HL - LHC + FCC - ee (91 + 161 + 240 + 365 GeV)



# Fisher information matrix

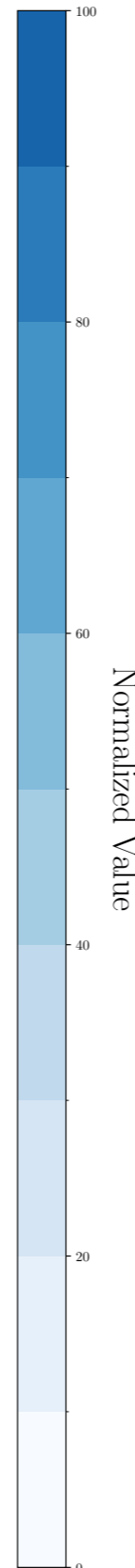
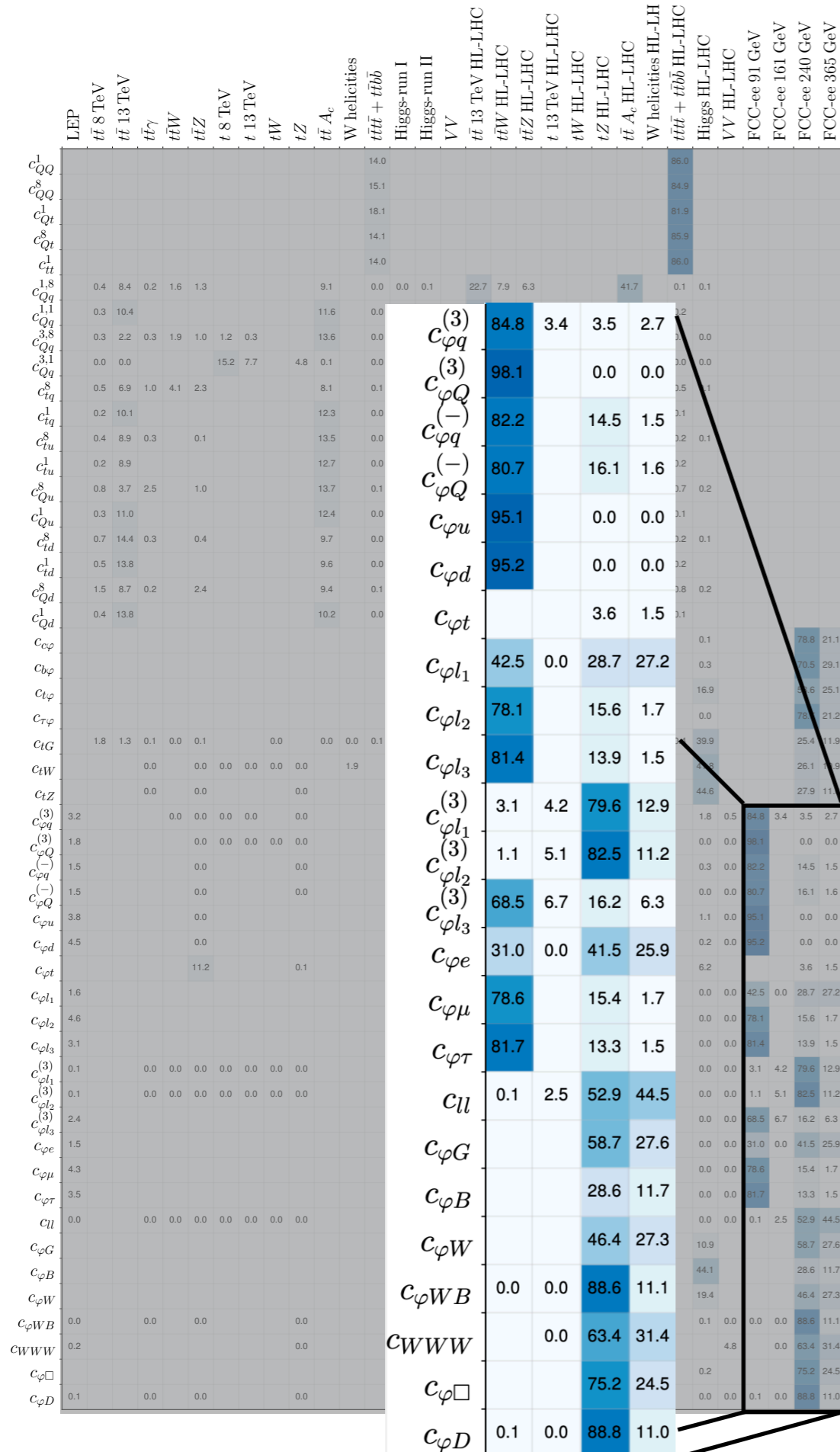


- ▶ The sensitivity of the EFT parameters to the Run II, HL-LHC and FCC-ee datasets is quantified by the **fisher information**

$$I_{ij} = \sum_{m=1}^{n_{\text{dat}}} \frac{\sigma_{m,i}^{(\text{eft})} \sigma_{m,j}^{(\text{eft})}}{\delta_{\text{exp},m}^2}$$

- ▶ The highest sensitivity in the 2FB and bosonic sectors comes in via the FCC-ee
- ▶ The FCC-ee run at 161 GeV is the least sensitive for the SMEFT

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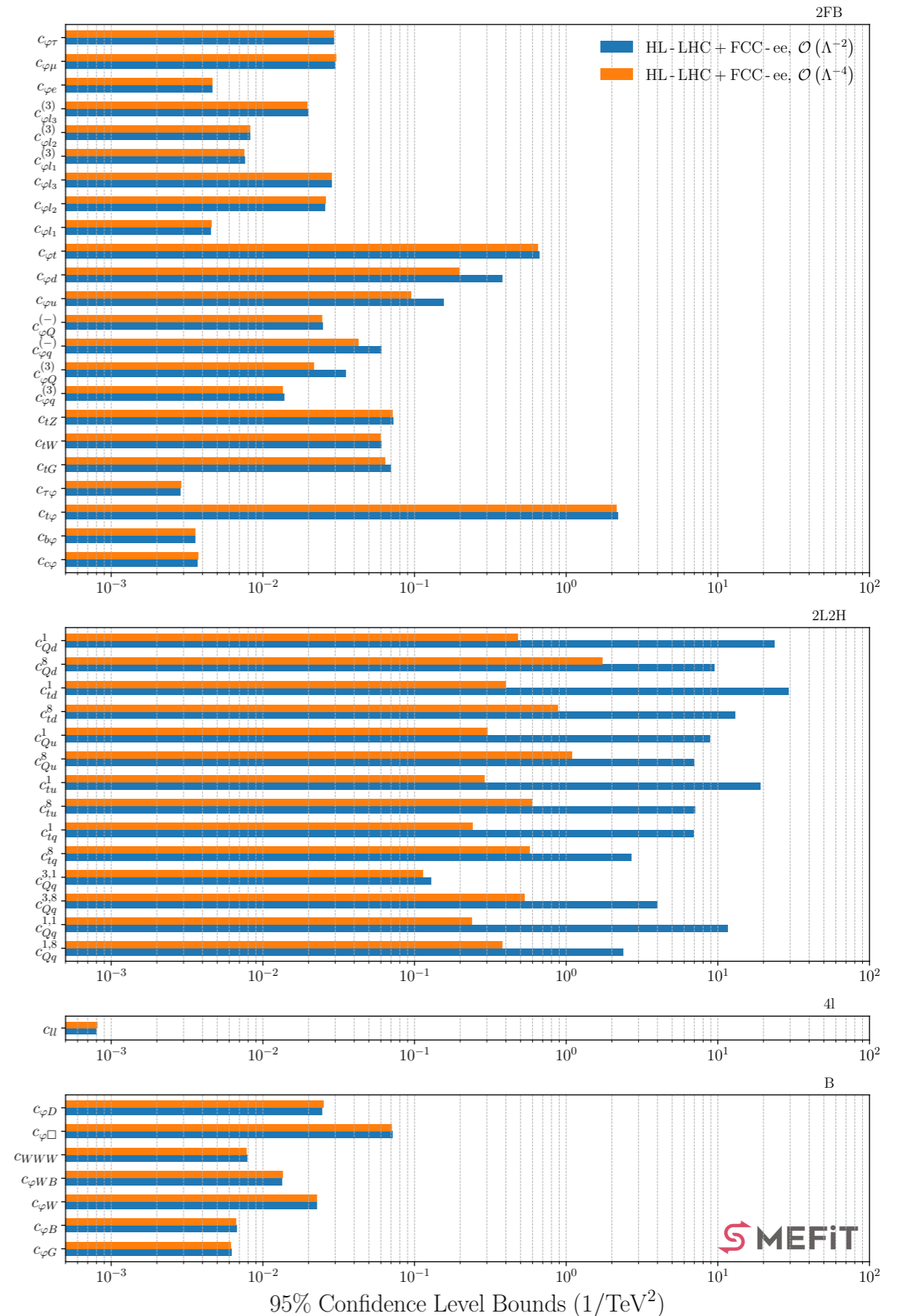
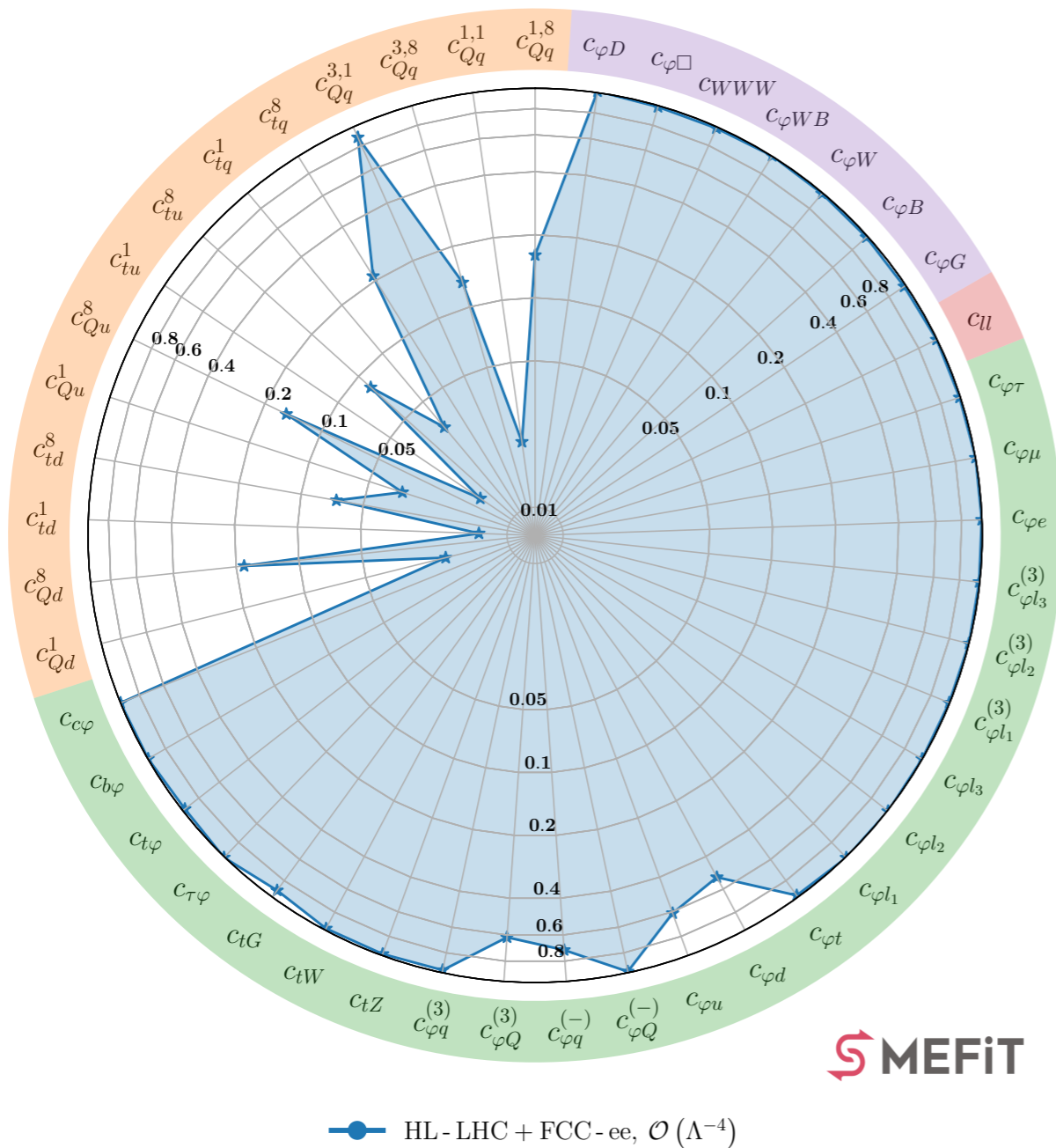
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- ▶ The highest sensitivity in the 2FB and bosonic sectors comes in via the FCC-ee
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# Impact of quadratics

The uncertainty due to the EFT truncation is small except for 4F

Ratio of Uncertainties to HL - LHC + FCC - ee,  $\mathcal{O}(\Lambda^{-2})$ , Marginalised

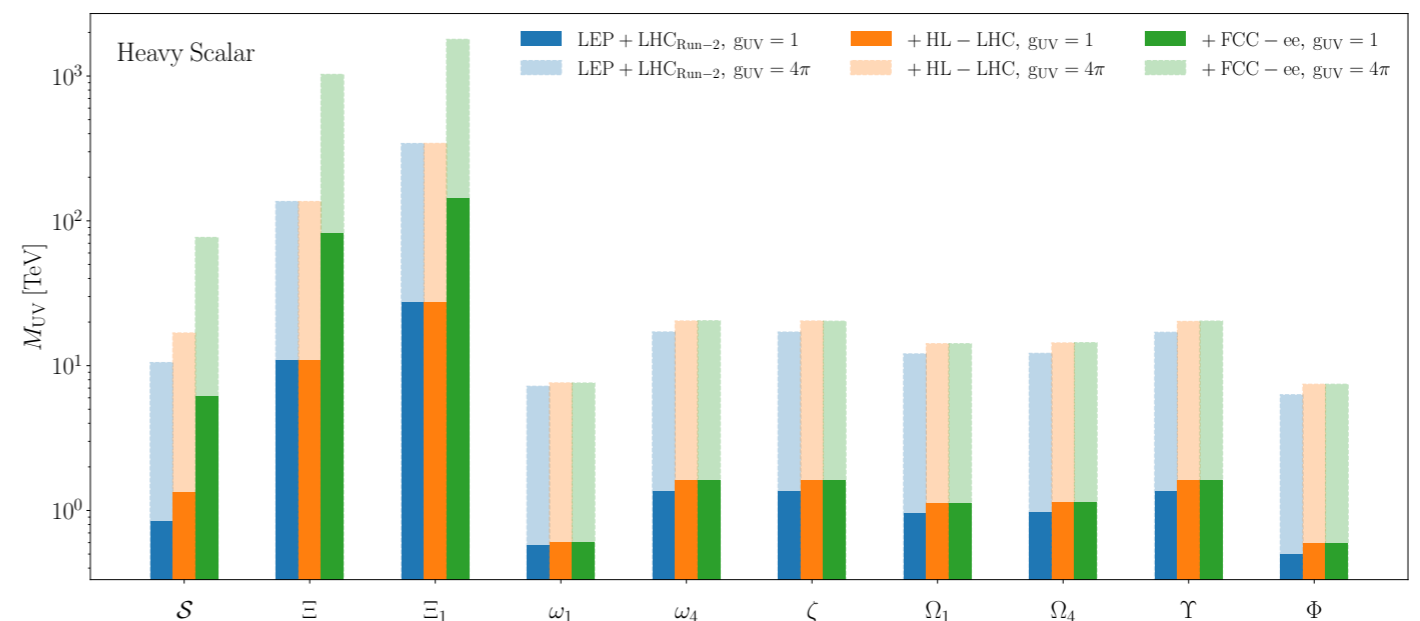


# RGE effects

- ▶ Experimental input to global fits spans a **wide range of different energy scales**, from  $m_Z$  at LEP to  $m_{t\bar{t}} \sim 3$  TeV in tails at LHC
- ▶ At FCC-ee this gap becomes **even more sizeable** due to an unprecedented indirect mass reach up  $\mathcal{O}(100)$  TeV
- ▶ A consistent treatment that connects measurements across different scales thus becomes necessary and is provided by the **Renormalisation Group Equations (RGEs)**

$$\frac{dc_i(\mu)}{d \ln \mu} = \sum_{j=1}^{n_{\text{op}}} \gamma_{ij} c_j(\mu)$$

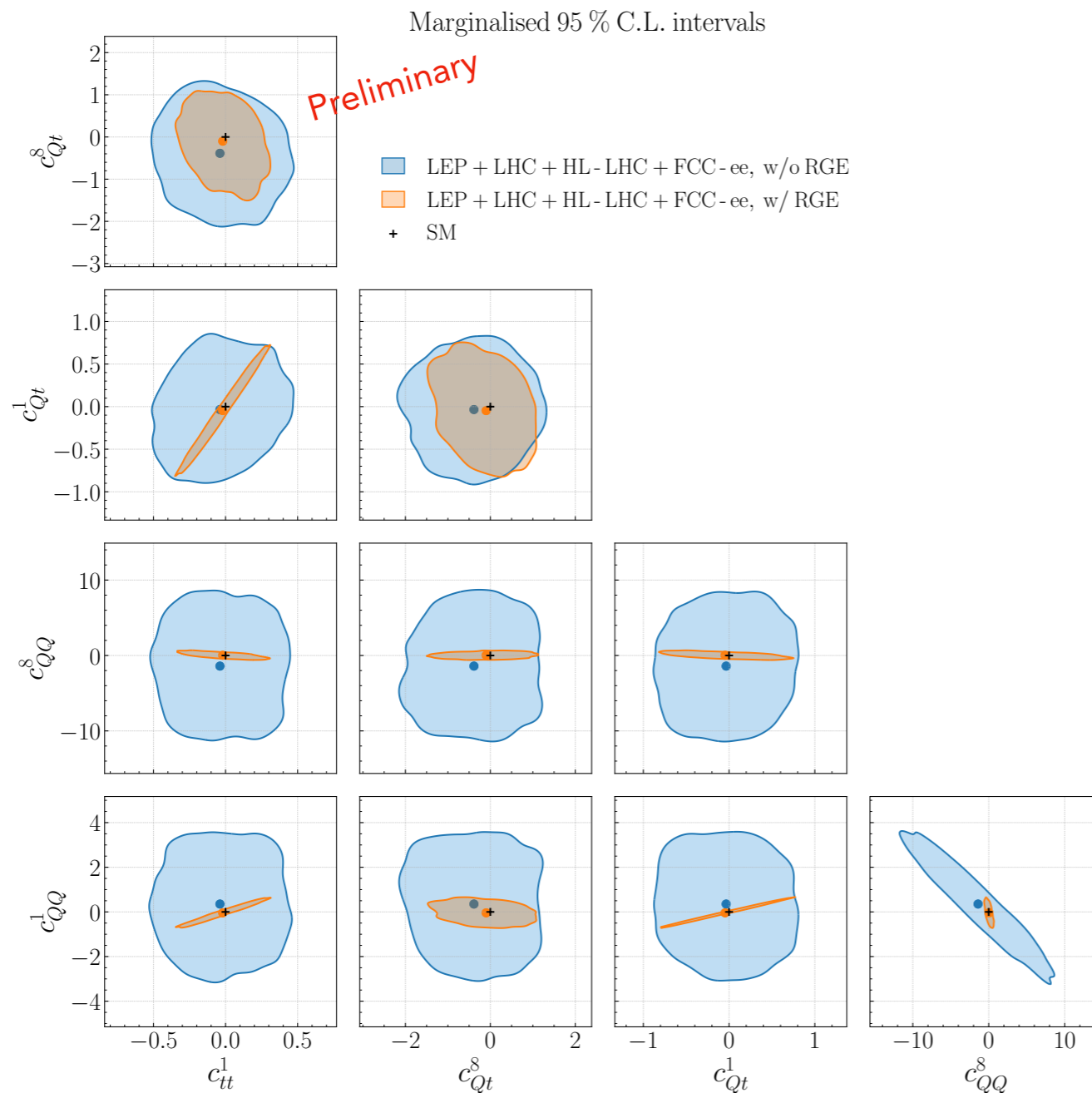
Anomalous dimension



ter Hoeve et al [2309.04523]

# RGE effects

Wilson Coefficients run and mix with energy under the RGEs, possibly resulting in higher sensitivity



Two competing effects:

- More operators enter the same observable
- An ill constrained operator could flow into a precisely determined observable

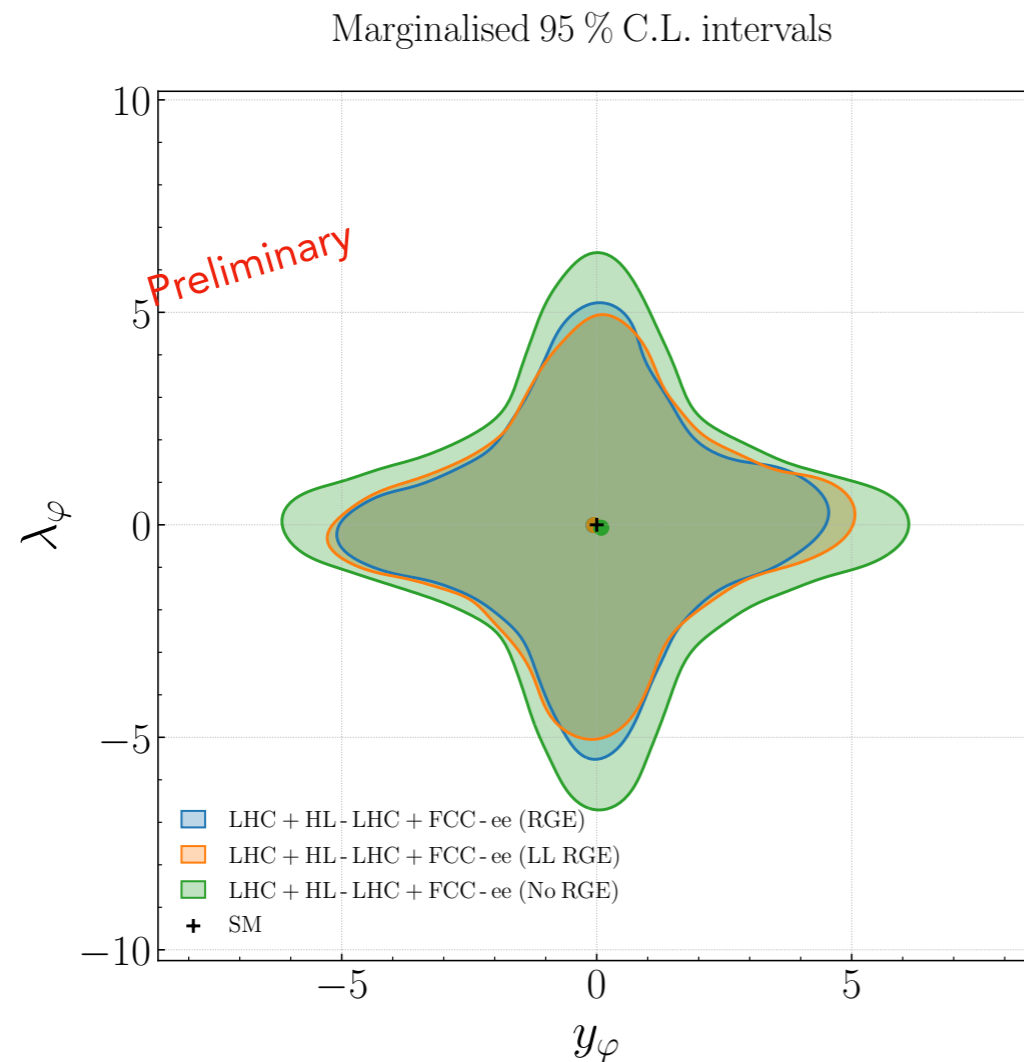


[1804.05033] Aebischer, Kumar, Straub

# RGE effects in the UV

The 2HDM\* is seen to benefit from including RG effects, with bounds improving by around 20%

\* In the decoupling limit



Two competing effects:

- More operators enter the same observable
- An ill constrained operator could flow into a precisely determined observable



[1804.05033] Aebischer, Kumar, Straub

$$\begin{aligned}
 \mathcal{L}_{UV} = & \mathcal{L}_{SM} + |D_\mu \phi|^2 - m_\phi^2 \phi^\dagger \phi - \left( (y_\phi^e)_{ij} \phi^\dagger \bar{e}_R^i \ell_L^j + (y_\phi^d)_{ij} \phi^\dagger \bar{d}_R^i q_L^j \right. \\
 & \left. + (y_\phi^u)_{ij} \phi^\dagger i \sigma_2 \bar{q}_L^{T,i} u_R^j + \lambda_\phi \phi^\dagger \phi |\phi|^2 + \text{h.c.} \right) - \text{scalar potential}
 \end{aligned}$$

# Conclusion and outlook

- Presented **SMEFiT3.0**, a global fit of 50 Wilson coefficients to Higgs, top, diboson and EWPOs, including **quadratic corrections**
- We are becoming **increasingly sensitive** to possible new physics effects, both through the still expanding LHC datasets, as well as through future colliders
- The FCC-ee offers an **unprecedented indirect mass reach** on new heavy particles
- RGE effects are relevant to include to connect experiments at widely separated scales
- The inclusion of other proposed future colliders are WIP for the ESPPU

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Contact: [jaco.ter.hoeve@ed.ac.uk](mailto:jaco.ter.hoeve@ed.ac.uk)

**Thanks for your attention!**



# Backup

# Theory input

We fit a total of **50 Wilson coefficients** simultaneously at quadratic order in the EFT

## 2 fermion bosonic operators

Operator	Coefficient	Definition	Operator	Coefficient	Definition
3rd generation quarks					
$\mathcal{O}_{\varphi Q}^{(1)}$	$c_{\varphi Q}^{(1)}$ (*)	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{Q} \gamma^\mu Q)$	$\mathcal{O}_{tW}$	$c_{tW}$	$i(\bar{Q} \tau^{\mu\nu} \tau_I t) \tilde{\varphi} W_{\mu\nu}^I + \text{h.c.}$
$\mathcal{O}_{\varphi Q}^{(3)}$	$c_{\varphi Q}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{Q} \gamma^\mu \tau^I Q)$	$\mathcal{O}_{tB}$	$c_{tB}$ (*)	$i(\bar{Q} \tau^{\mu\nu} t) \tilde{\varphi} B_{\mu\nu} + \text{h.c.}$
$\mathcal{O}_{\varphi t}$	$c_{\varphi t}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{t} \gamma^\mu t)$	$\mathcal{O}_{tG}$	$c_{tG}$	$ig_S (\bar{Q} \tau^{\mu\nu} T_A t) \tilde{\varphi} G_{\mu\nu}^A + \text{h.c.}$
$\mathcal{O}_{t\varphi}$	$c_{t\varphi}$	$(\varphi^\dagger \varphi) \bar{Q} t \tilde{\varphi} + \text{h.c.}$	$\mathcal{O}_{b\varphi}$	$c_{b\varphi}$	$(\varphi^\dagger \varphi) \bar{Q} b \varphi + \text{h.c.}$
1st, 2nd generation quarks					
$\mathcal{O}_{\varphi q}^{(1)}$	$c_{\varphi q}^{(1)}$ (*)	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}_i \gamma^\mu q_i)$	$\mathcal{O}_{\varphi d}$	$c_{\varphi d}$	$\sum_{i=1,2,3} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{d}_i \gamma^\mu d_i)$
$\mathcal{O}_{\varphi q}^{(3)}$	$c_{\varphi q}^{(3)}$	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{q}_i \gamma^\mu \tau^I q_i)$	$\mathcal{O}_{c\varphi}$	$c_{c\varphi}$	$(\varphi^\dagger \varphi) \bar{q}_2 c \tilde{\varphi} + \text{h.c.}$
$\mathcal{O}_{\varphi u}$	$c_{\varphi u}$	$\sum_{i=1,2} i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}_i \gamma^\mu u_i)$			
two-leptons					
$\mathcal{O}_{\varphi l_i}$	$c_{\varphi l_i}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{l}_i \gamma^\mu l_i)$	$\mathcal{O}_{\varphi \mu}$	$c_{\varphi \mu}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{\mu} \gamma^\mu \mu)$
$\mathcal{O}_{\varphi l_i}^{(3)}$	$c_{\varphi l_i}^{(3)}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi)(\bar{l}_i \gamma^\mu \tau^I l_i)$	$\mathcal{O}_{\varphi \tau}$	$c_{\varphi \tau}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{\tau} \gamma^\mu \tau)$
$\mathcal{O}_{\varphi e}$	$c_{\varphi e}$	$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{e} \gamma^\mu e)$	$\mathcal{O}_{\tau\varphi}$	$c_{\tau\varphi}$	$(\varphi^\dagger \varphi) \bar{l}_3 \tau \varphi + \text{h.c.}$

## Four lepton operators

four-leptons					
$\mathcal{O}_{\ell\ell}$	$c_{\ell\ell}$	$(\bar{l}_1 \gamma_\mu l_2)(\bar{l}_2 \gamma^\mu l_1)$			

## 4 fermion operators

DoF	Definition (in Warsaw basis notation)	DoF	Definition (in Warsaw basis notation)
$c_{QQ}^1$	$2c_{qq}^{1(3333)} - \frac{2}{3}c_{qq}^{3(3333)}$	$c_{QQ}^8$	$8c_{qq}^{3(3333)}$
$c_{Qt}^1$	$c_{qu}^{1(3333)}$	$c_{Qt}^8$	$c_{qu}^{8(3333)}$
$c_{Qq}^{1,8}$	$c_{qq}^{1(i33i)} + 3c_{qq}^{3(i33i)}$	$c_{Qq}^{1,1}$	$c_{qq}^{1(i33i)} + \frac{1}{6}c_{qq}^{1(i33i)} + \frac{1}{2}c_{qq}^{3(i33i)}$
$c_{Qq}^{3,8}$	$c_{qq}^{1(i33i)} - c_{qq}^{3(i33i)}$	$c_{Qq}^{3,1}$	$c_{qq}^{3(i33i)} + \frac{1}{6}(c_{qq}^{1(i33i)} - c_{qq}^{3(i33i)})$
$c_{tq}^8$	$c_{qu}^{8(i33i)}$	$c_{tq}^1$	$c_{qu}^{1(i33i)}$
$c_{tu}^8$	$2c_{uu}^{(i33i)}$	$c_{tu}^1$	$c_{uu}^{(i33i)} + \frac{1}{3}c_{uu}^{(i33i)}$
$c_{Qu}^8$	$c_{qu}^{8(33ii)}$	$c_{Qu}^1$	$c_{qu}^{1(33ii)}$
$c_{td}^8$	$c_{ud}^{8(33jj)}$	$c_{td}^1$	$c_{ud}^{1(33jj)}$
$c_{Qd}^8$	$c_{qd}^{8(33jj)}$	$c_{Qd}^1$	$c_{qd}^{1(33jj)}$

## Purely bosonic

Operator	Coefficient	Definition	Operator	Coefficient	Definition
$\mathcal{O}_{\varphi G}$	$c_{\varphi G}$	$(\varphi^\dagger \varphi) G_A^{\mu\nu} G_{\mu\nu}^A$	$\mathcal{O}_{\varphi \square}$	$c_{\varphi \square}$	$\partial_\mu(\varphi^\dagger \varphi) \partial^\mu(\varphi^\dagger \varphi)$
$\mathcal{O}_{\varphi B}$	$c_{\varphi B}$	$(\varphi^\dagger \varphi) B^{\mu\nu} B_{\mu\nu}$	$\mathcal{O}_{\varphi D}$	$c_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\dagger (\varphi^\dagger D_\mu \varphi)$
$\mathcal{O}_{\varphi W}$	$c_{\varphi W}$	$(\varphi^\dagger \varphi) W_I^{\mu\nu} W_{\mu\nu}^I$	$\mathcal{O}_W$	$c_{WWW}$	$\epsilon_{IJK} W_{\mu\nu}^I W^{\nu\rho J} W_\rho^{K,\mu}$
$\mathcal{O}_{\varphi WB}$	$c_{\varphi WB}$	$(\varphi^\dagger \tau_I \varphi) B^{\mu\nu} W_{\mu\nu}^I$			

Flavour assumption:

$$U(2)_q \times U(3)_d \times U(2)_u \times (U(1)_\ell \times U(1)_e)^3$$

# HL-LHC projections

- ▶ The central values of the pseudo data are fluctuated around the SM

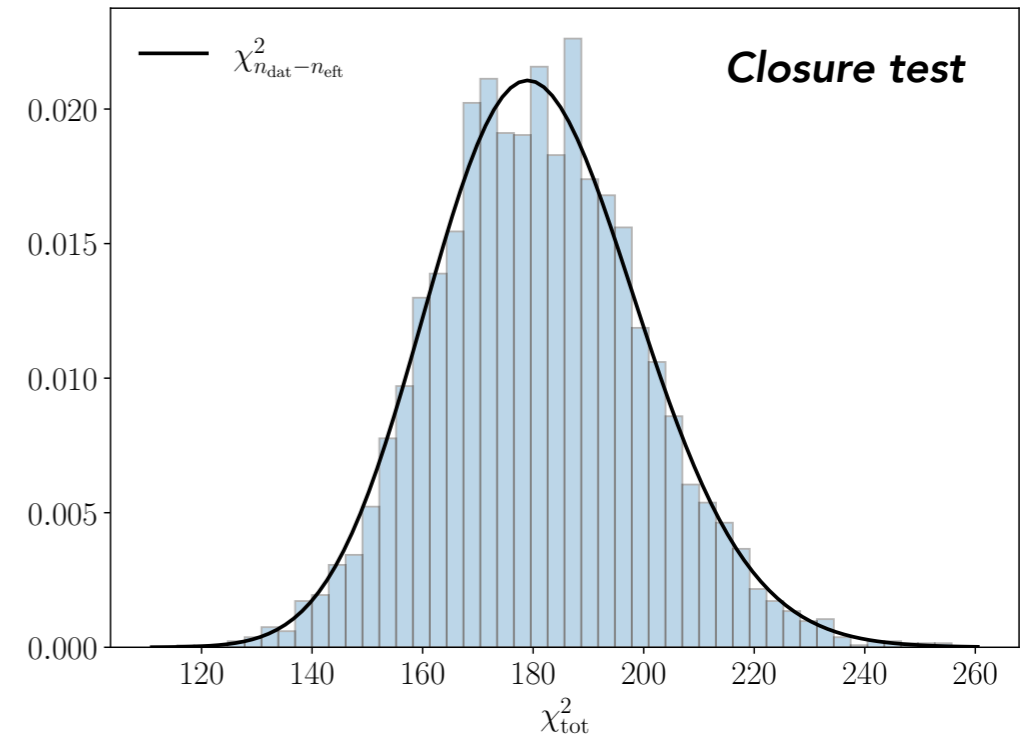
$$\mathcal{O}_i^{(\text{exp})} = \mathcal{O}_i^{(\text{th})} \left( 1 + r_i \delta_i^{(\text{stat})} + \sum_{k=1}^{n_{\text{sys}}} r_{k,i} \delta_{k,i}^{(\text{sys})} \right)$$

- ▶ Statistical uncertainties we rescale according to the improved luminosity

$$\delta_i^{(\text{stat})} = \tilde{\delta}_i^{(\text{stat})} \sqrt{\frac{\mathcal{L}_{\text{Run2}}}{\mathcal{L}_{\text{HLLHC}}}}$$

- ▶ While systematics are rescaled by an overall factor, namely 1/2 for all datasets

$$\delta_{k,i}^{(\text{sys})} = \tilde{\delta}_{k,i}^{(\text{sys})} \times f_{\text{red}}^{(k)} \quad k = 1, \dots, n_{\text{sys}}$$



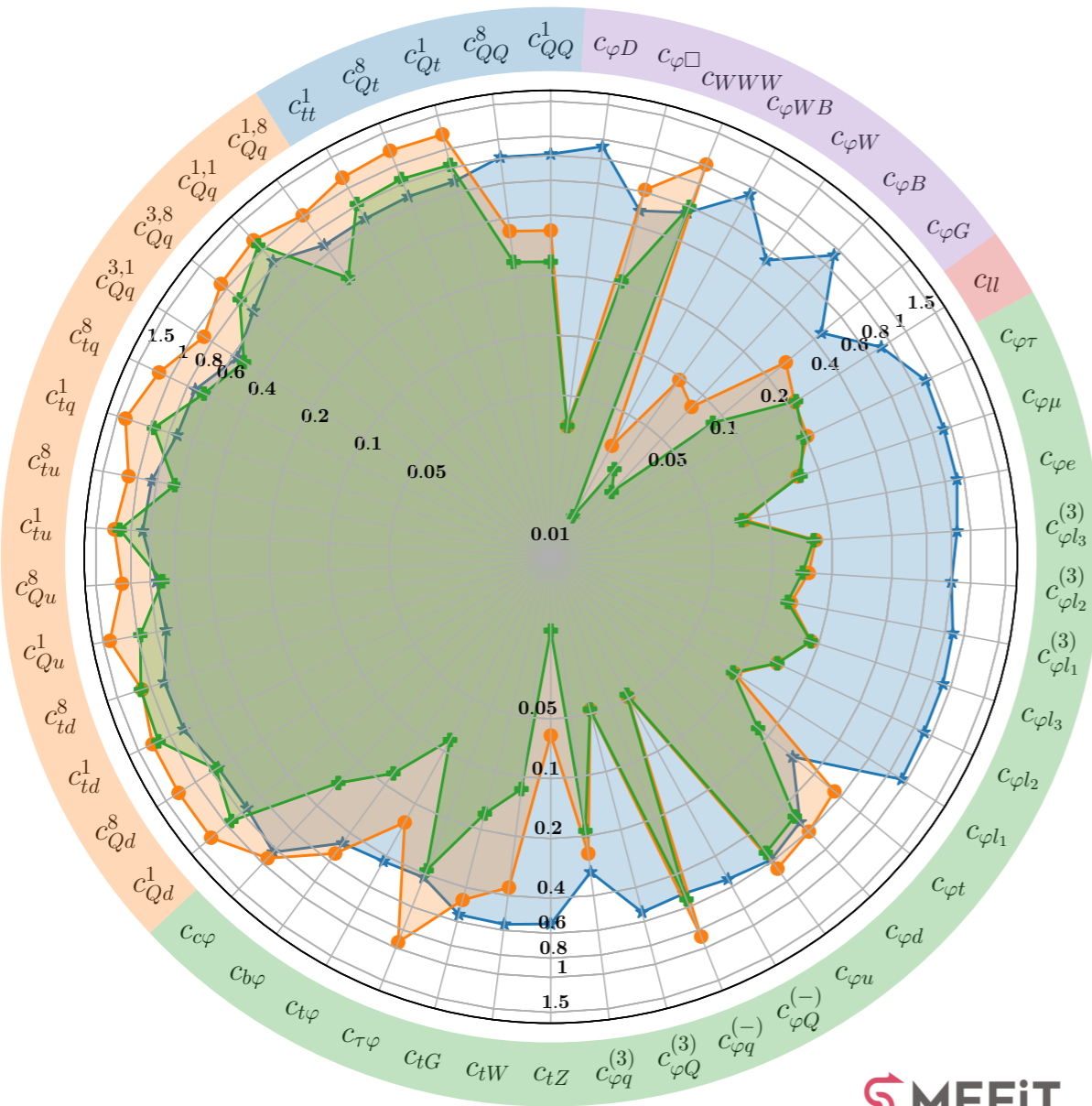
+ flexible framework that can project any Run II dataset

+ SMEFT predictions can be recycled

- No additional bins in the tails

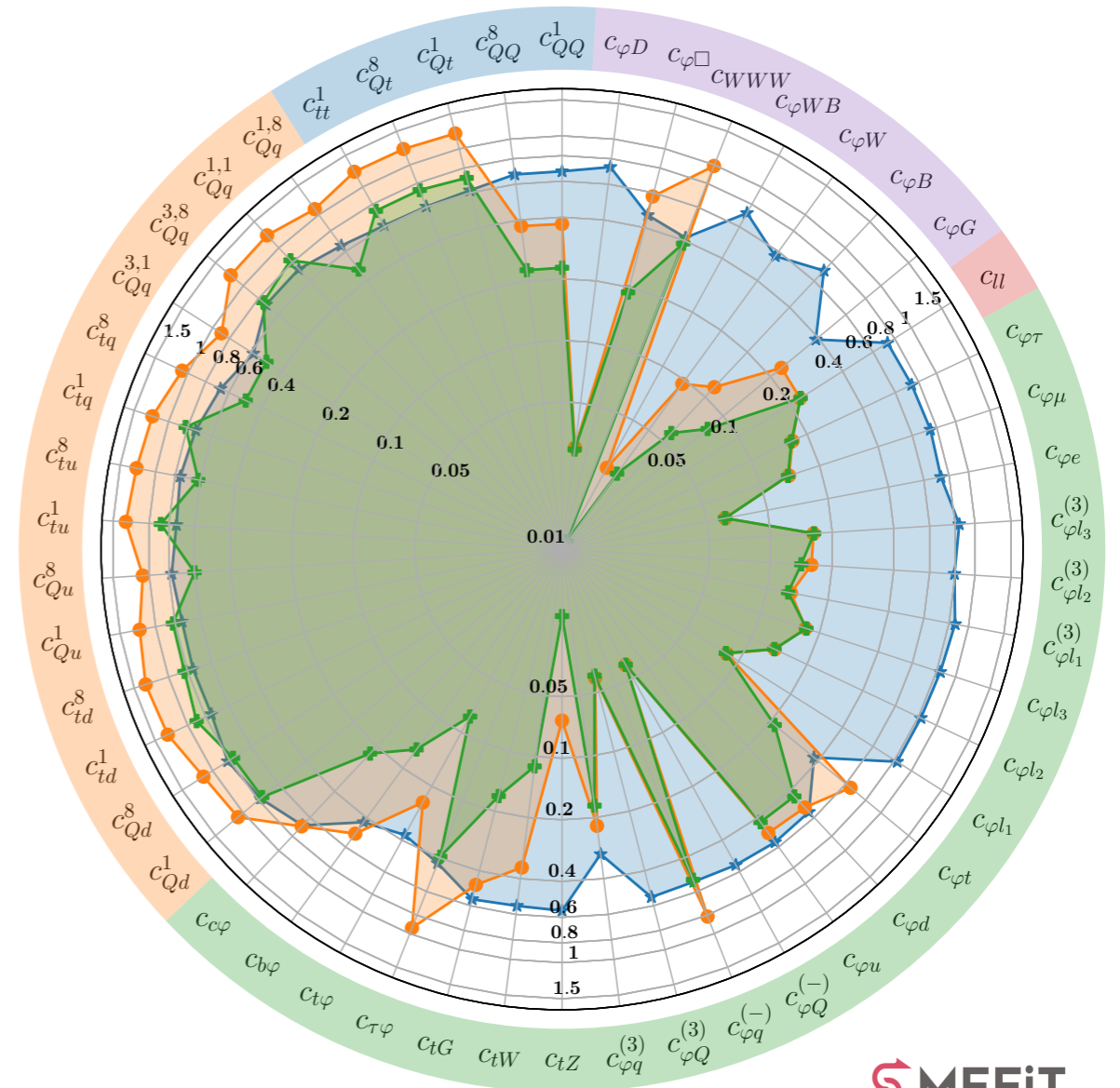
### With statistical noise = L1

Ratio of Uncertainties to SMEFiT3.0 Baseline,  $\mathcal{O}(\Lambda^{-4})$ , Marginalised

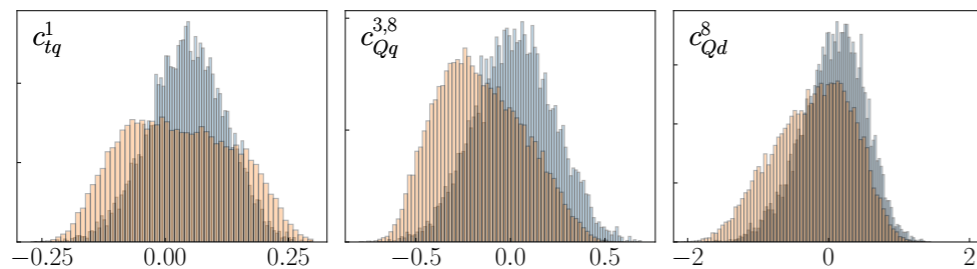


### Without statistical noise = L0

Ratio of Uncertainties to SMEFiT3.0 Baseline,  $\mathcal{O}(\Lambda^{-4})$ , Marginalised



■ SMEFiT3.0, NLO  $\mathcal{O}(\Lambda^{-4})$ , Marginalised
 ■ SMEFiT3.0, NLO  $\mathcal{O}(\Lambda^{-4})$ , Individual

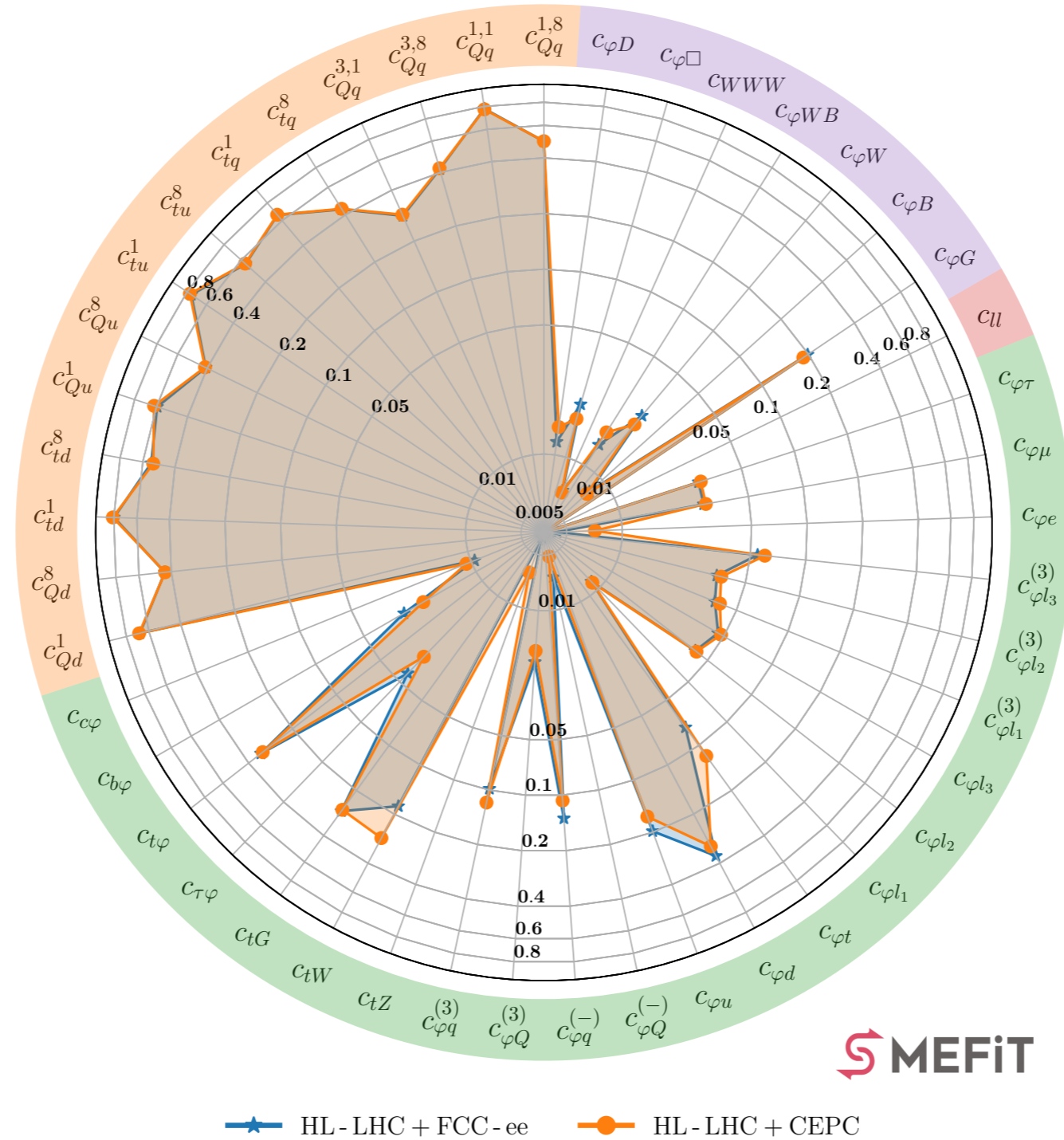


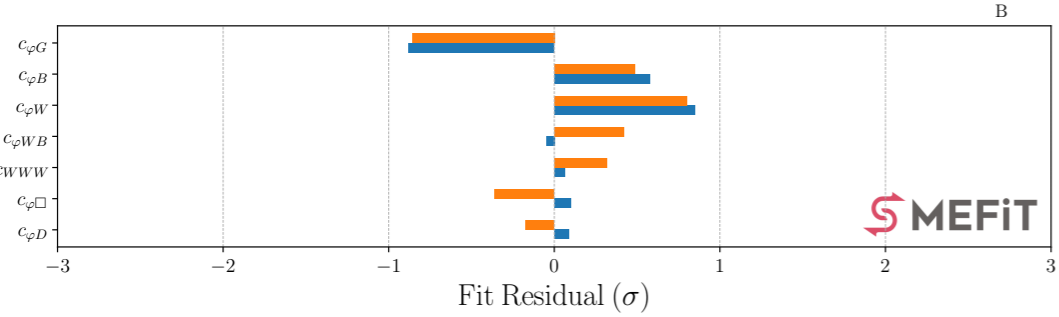
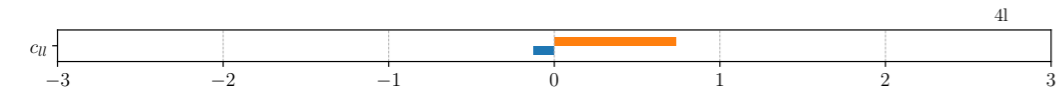
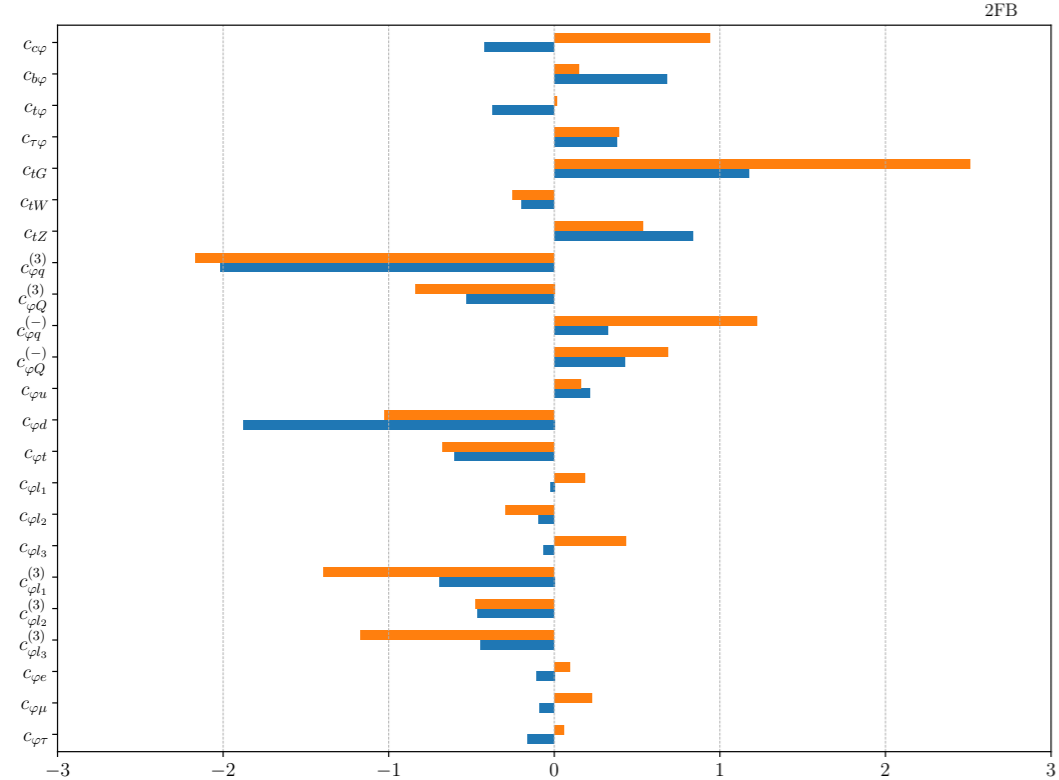
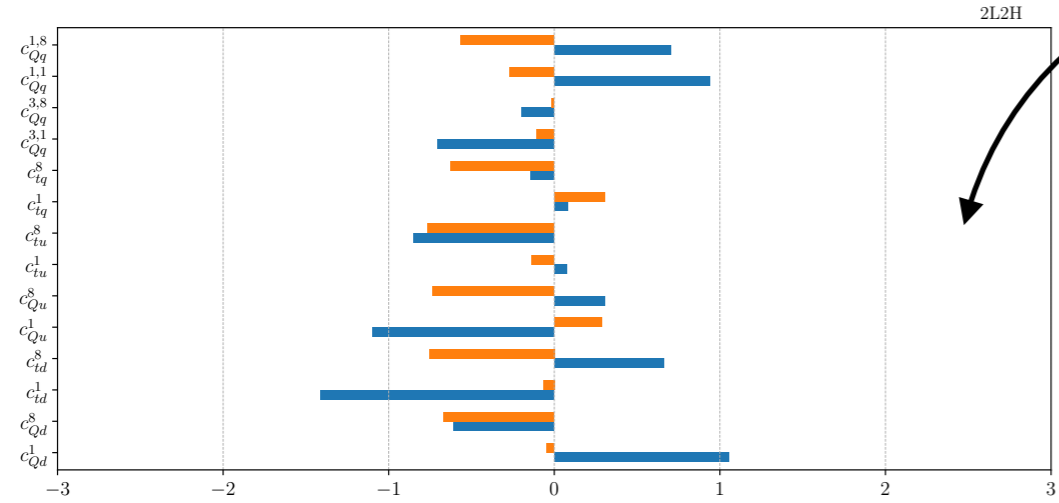
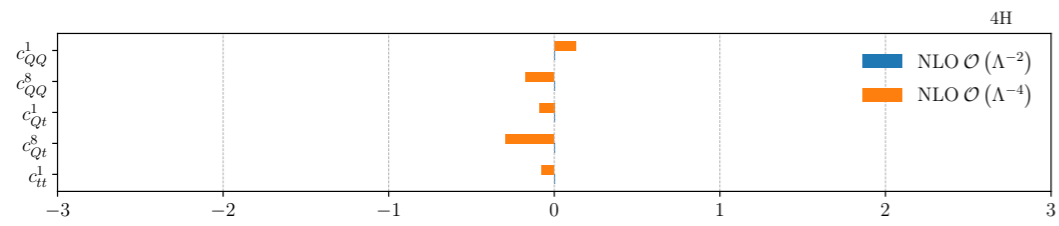
★ HL-LHC
 ★ HL-LHC, individual

● SMEFiT3.0, individual

# FCC-ee and CEPC

Ratio of Uncertainties to SMEFiT3.0 Baseline,  $\mathcal{O}(\Lambda^{-2})$ , Marginalised

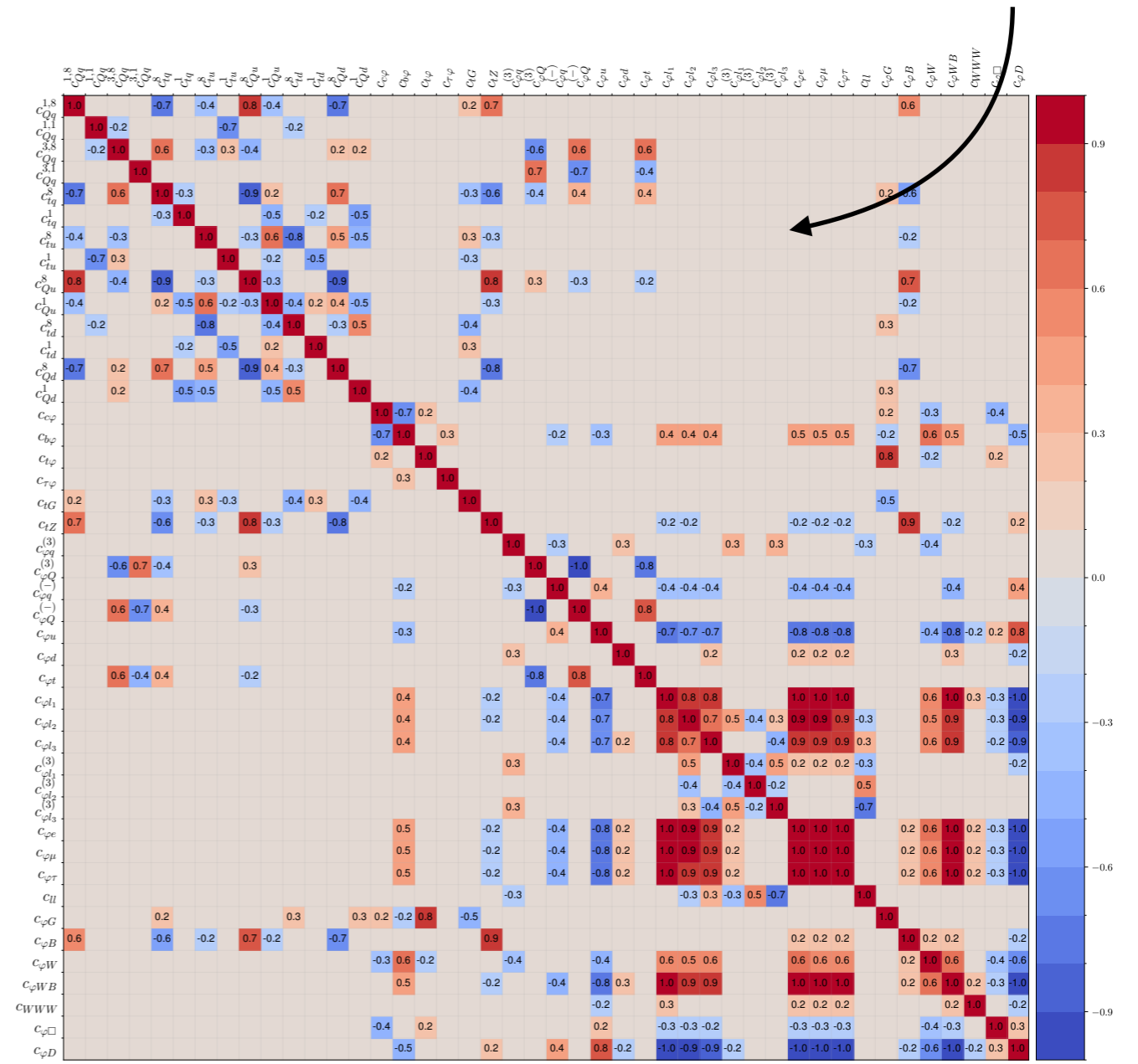




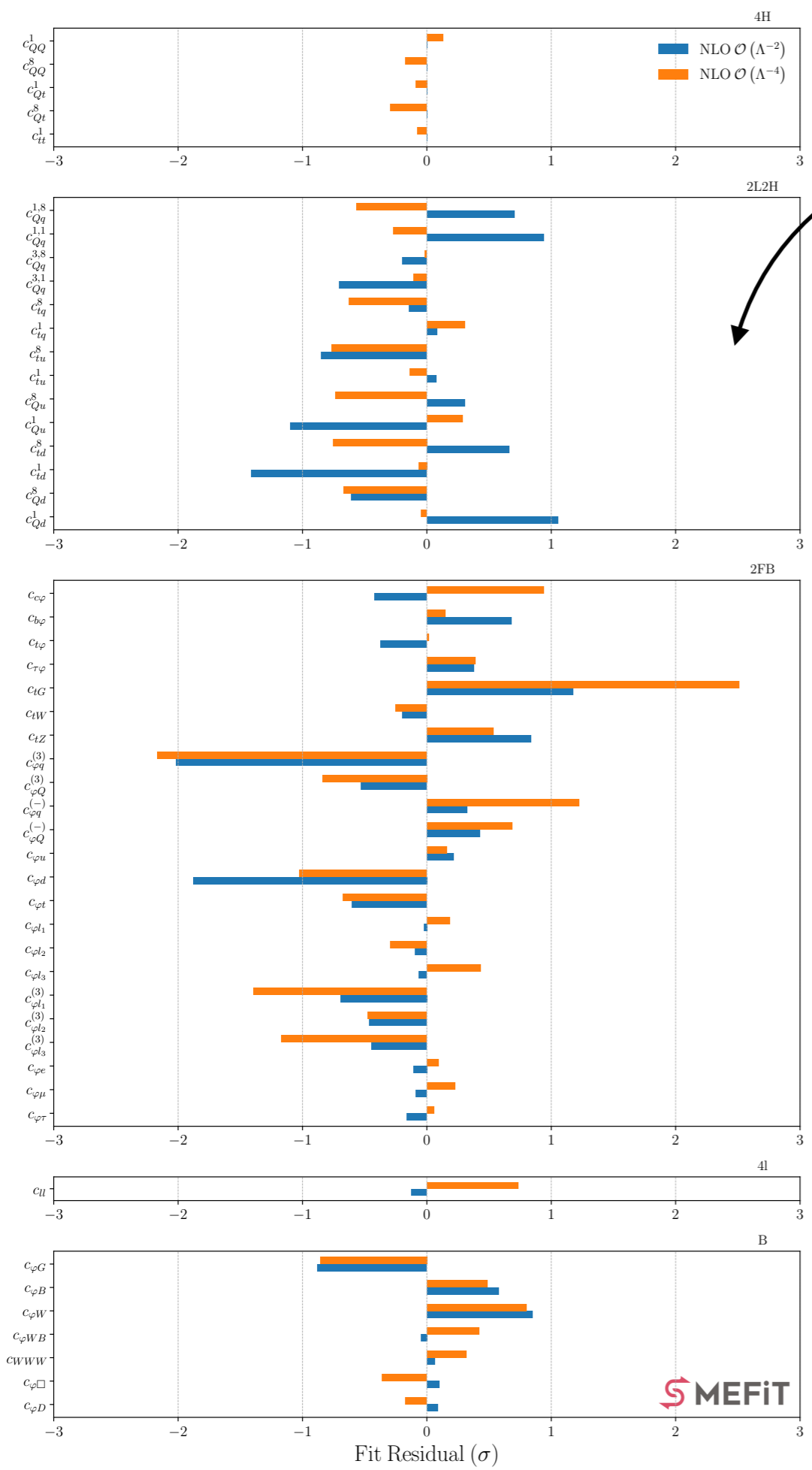
Fit residuals (pulls) are largely **consistent** with the SM

$$P_i \equiv \frac{\langle c_i \rangle - c_i^{(\text{SM})}}{[c_i^{\text{min}}, c_i^{\text{max}}]^{68\% \text{ CL}}}$$

Large correlations in linear fit get lifted in the quadratic fit



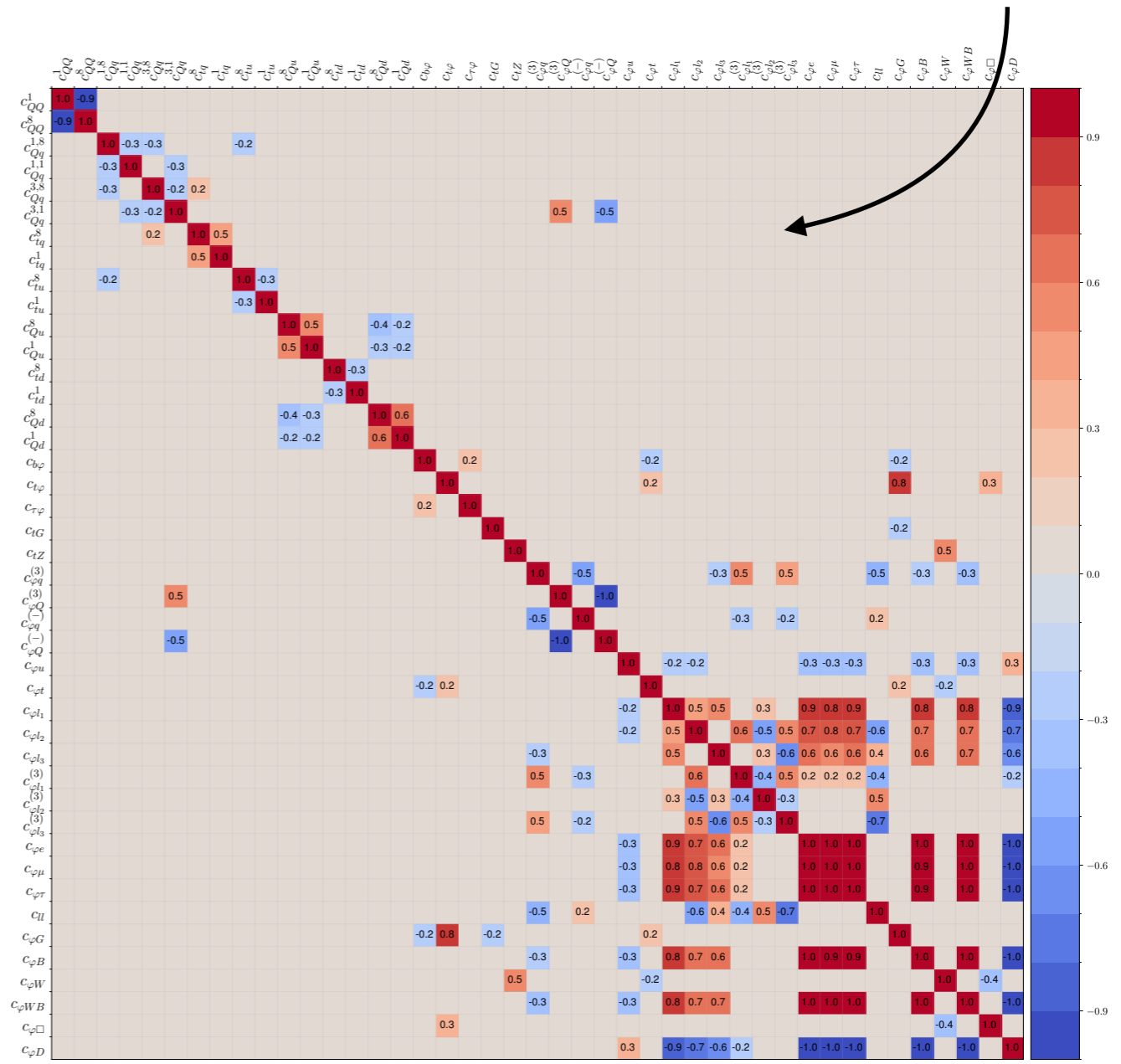
Correlation: NLO  $\mathcal{O}(\Lambda^{-2})$



Fit residuals (pulls) are largely **consistent** with the SM

$$P_i \equiv \frac{\langle c_i \rangle - c_i^{(SM)}}{[c_i^{\min}, c_i^{\max}]^{68\% \text{ CL}}}$$

Large correlations in linear fit get lifted in the quadratic fit



# Building the likelihood

From (differential) cross sections ...

$$\sigma_{\text{SMEFT}}(c, \Lambda) = \sigma_{\text{SM}} \times \left( 1 + \sum_i^{N_{d6}} \kappa_i \frac{c_i}{\Lambda^2} + \sum_{i<j}^{N_{d6}} \tilde{\kappa}_{ij} \frac{c_i \cdot c_j}{\Lambda^4} + \mathcal{O}(\Lambda^{-6}) \right)$$

Linear EFT corrections:  
interference SM-EFT<sub>d6</sub>  
@NLO QCD

Quadratic EFT  
corrections:  
EFT<sub>d6</sub>-EFT<sub>d6</sub>  
@NLO QCD

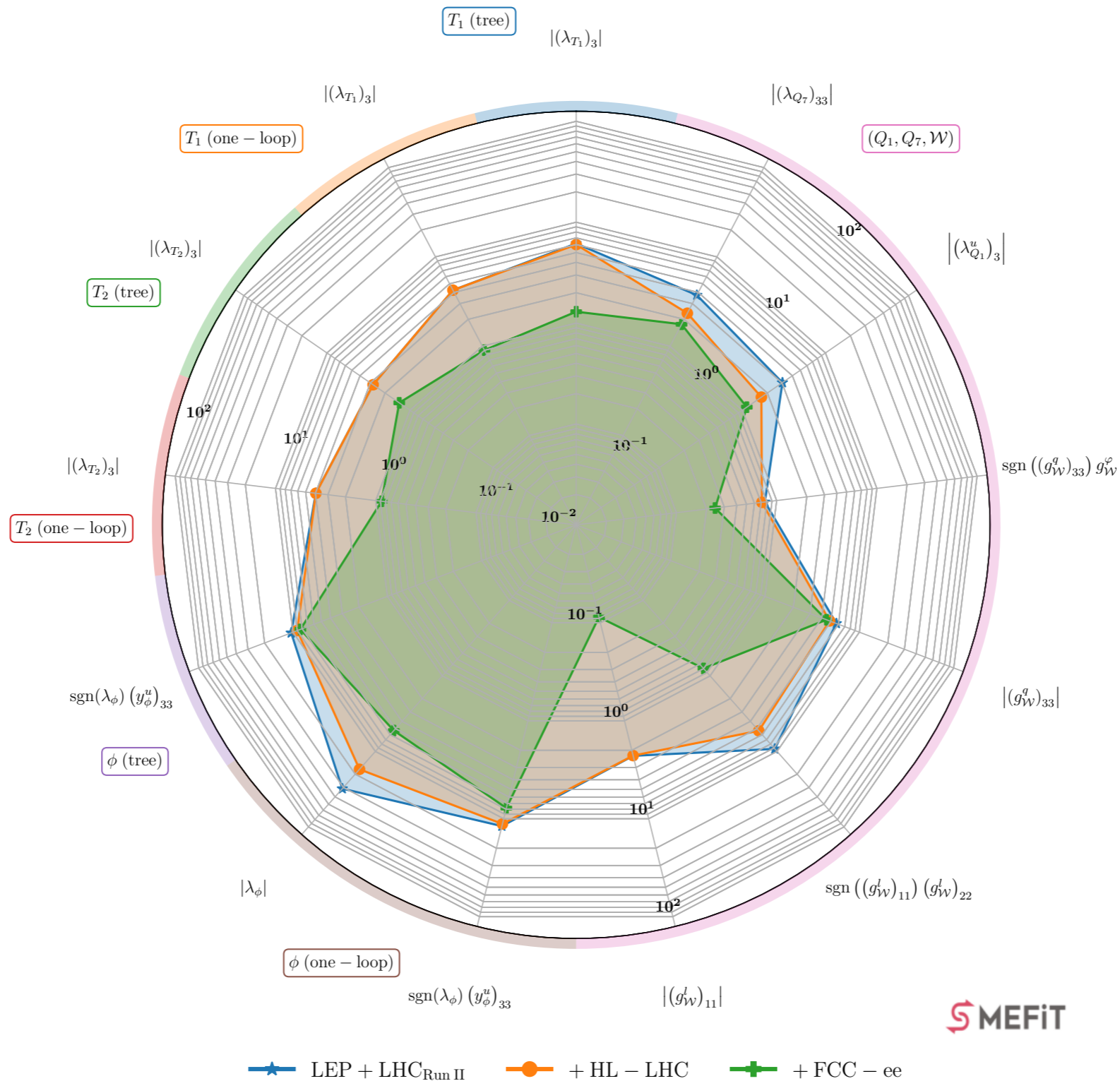
To a combined likelihood ready for optimisation ...

$$-2 \log \mathcal{L} = \frac{1}{n_{\text{dat}}} \sum_{i,j=1}^{n_{\text{dat}}} \left( \sigma_{i,\text{SMEFT}}(c) - \sigma_{i,\text{exp}} \right) (\text{cov}^{-1})_{ij} \left( \sigma_{j,\text{SMEFT}}(c) - \sigma_{j,\text{exp}} \right)$$

Theory (pdf + scale) and experimental uncertainties (stat + systematics):  $\text{cov}^{(\text{tot})}_{ij} = \text{cov}^{(\text{th})}_{ij} + \text{cov}^{(\text{exp})}_{ij}$



# 1-loop & multi-particle matching



# SM predictions

Category	Process	SM	Code/Ref	SMEFT
Top quark production	$t\bar{t}$ (incl)	NNLO QCD	MG5_aMC NLO + NNLO $K$ -fact	NLO QCD
	$t\bar{t} + V$	NLO QCD	MG5_aMC NLO	LO QCD + NLO SM $K$ -fact
	single- $t$ (incl)	NNLO QCD	MG5_aMC NLO + NNLO $K$ -fact	NLO QCD
	$t + V$	NLO QCD	MG5_aMC NLO	LO QCD + NLO SM $K$ -fact
	$t\bar{t}t\bar{t}, t\bar{t}b\bar{b}$	NLO QCD	MG5_aMC NLO	LO QCD + NLO SM $K$ -fact
Higgs production and decay	$gg \rightarrow h$	NNLO QCD + NLO EW	HXSWG	NLO QCD
	VBF	NNLO QCD + NLO EW	HXSWG	LO QCD
	$h + V$	NNLO QCD + NLO EW	HXSWG	NLO QCD
	$ht\bar{t}$	NNLO QCD + NLO EW	HXSWG	NLO QCD
	$h \rightarrow X$	NNLO QCD + NLO EW	HXSWG	NLO QCD ( $X = b\bar{b}$ ) LO QCD ( $X \neq b\bar{b}$ )
Diboson production	$e^+e^- \rightarrow W^+W^-$	NNLO QCD + NLO EW	LEP EWWG	LO QCD
	$pp \rightarrow VV'$	NNLO QCD	MATRIX	NLO QCD

# HL-LHC projected datasets

Dataset	$\mathcal{L}$ (fb <sup>-1</sup> )	Info	Observables	$n_{\text{dat}}$	Ref.
ATLAS_STXS_RunII_13TeV_2022	139	$ggF, \text{VBF}, Vh, t\bar{t}h, th$	$d\sigma/dp_T^h$ $d\sigma/dm_{jj}$ $d\sigma/dp_T^V$	36	[55]
CMS_ggF_aa_13TeV	77.4	$ggF, h \rightarrow \gamma\gamma$	$\sigma_{ggF}(p_T^h, N_{\text{jets}})$	6	[83]
ATLAS_ggF_ZZ_13TeV	79.8	$ggF, h \rightarrow ZZ$	$\sigma_{ggF}(p_T^h, N_{\text{jets}})$	6	[84]
ATLAS_ggF_13TeV_2015	36.1	$ggF, h \rightarrow ZZ, h \rightarrow \gamma\gamma$	$d\sigma(ggF)/dp_T^h$	9	[85]
ATLAS_WH_Hbb_13TeV	79.8	$Wh, h \rightarrow b\bar{b}$	$d\sigma^{(\text{fid})}/dp_T^W$ (stage 1 STXS)	2	[86]
ATLAS_ZH_Hbb_13TeV	79.8	$Zh, h \rightarrow b\bar{b}$	$d\sigma^{(\text{fid})}/dp_T^Z$ (stage 1 STXS)	2	[86]
CMS_H_13TeV_2015_pTH	35.9	$h \rightarrow b\bar{b}, h \rightarrow \gamma\gamma, h \rightarrow ZZ$	$d\sigma/dp_T^h$	9	[87]
ATLAS_WW_13TeV_2016_memu	36.1	fully leptonic	$d\sigma^{(\text{fid})}/dm_{e\mu}$	13	[88]
ATLAS_WZ_13TeV_2016_mTWZ	36.1	fully leptonic	$d\sigma^{(\text{fid})}/dm_T^{WZ}$	6	[89]
CMS_WZ_13TeV_2016_pTZ	35.9	fully leptonic	$d\sigma^{(\text{fid})}/dp_T^Z$	11	[90]
CMS_WZ_13TeV_2022_pTZ	137	fully leptonic	$d\sigma/dp_T^Z$	11	[56]

Dataset	$\mathcal{L}$ (fb $^{-1}$ )	Info	Observables	$n_{\text{dat}}$	Ref.
ATLAS_tt_13TeV_ljets_2016.Mtt	36.1	$\ell$ +jets	$d\sigma/dm_{t\bar{t}}$	7	[91]
CMS_tt_13TeV_dilep_2016.Mtt	35.9	dilepton	$d\sigma/dm_{t\bar{t}}$	7	[92]
CMS_tt_13TeV.Mtt	137	$\ell$ +jets	$1/\sigma d\sigma/dm_{t\bar{t}}$	14	[57]
CMS_tt_13TeV_ljets_inc	137	$\ell$ +jets	$\sigma(t\bar{t})$	1	[57]
ATLAS_tt_13TeV_asy_2022	139	$\ell$ + jets	$A_C$	5	[59]
CMS_tt_13TeV_asy	138	$\ell$ + jets	$A_C$	3	[58]
ATLAS_Whel_13TeV	139	W-helicity fraction	$F_0, F_L$	2	[60]
ATLAS_ttbb_13TeV_2016	36.1	lepton + jets	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[93]
CMS_ttbb_13TeV_2016	35.9	all-jets	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[94]
CMS_ttbb_13TeV_dilepton_inc	35.9	dilepton	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[68]
CMS_ttbb_13TeV_ljets_inc	35.9	lepton + jets	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[68]
ATLAS_tttt_13TeV_run2	139	multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[95]
CMS_tttt_13TeV_run2	137	same-sign or multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[96]
ATLAS_tttt_13TeV_slep_inc	139	single-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[64]
CMS_tttt_13TeV_slep_inc	35.8	single-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[65]
ATLAS_tttt_13TeV_2023	139	multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[66]
CMS_tttt_13TeV_2023	139	same-sign or multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[67]
CMS_ttZ_13TeV_pTZ	77.5	$t\bar{t}Z$	$d\sigma(t\bar{t}Z)/dp_T^Z$	4	[97]
ATLAS_ttZ_13TeV_pTZ	139	$t\bar{t}Z$	$d\sigma(t\bar{t}Z)/dp_T^Z$	7	[61]
ATLAS_ttW_13TeV_2016	36.1	$t\bar{t}W$	$\sigma_{\text{tot}}(t\bar{t}W)$	1	[98]
CMS_ttW_13TeV	35.9	$t\bar{t}W$	$\sigma_{\text{tot}}(t\bar{t}W)$	1	[99]
ATLAS_t_tch_13TeV_inc	3.2	$t$ -channel	$\sigma_{\text{tot}}(tq), \sigma_{\text{tot}}(\bar{t}q)$	2	[100]
CMS_t_tch_13TeV_2019_diff_Yt	35.9	$t$ -channel	$d\sigma/d y_t $	5	[101]
ATLAS_t_sch_13TeV_inc	139	$s$ -channel	$\sigma(t + \bar{t})$	1	[69]
ATLAS_tW_13TeV_inc	3.2	multi-lepton	$\sigma_{\text{tot}}(tW)$	1	[102]
CMS_tW_13TeV_inc	35.9	multi-lepton	$\sigma_{\text{tot}}(tW)$	1	[103]
CMS_tW_13TeV_slep_inc	36	single-lepton	$\sigma_{\text{tot}}(tW)$	1	[71]
ATLAS_tZ_13TeV_run2_inc	139	multi-lepton + jets	$\sigma_{\text{fid}}(t\ell^+\ell^-q)$	1	[104]
CMS_tZ_13TeV_pTt	138	multi-lepton + jets	$d\sigma_{\text{fid}}(tZj)/dp_T^t$	3	[70]

# FCC-ee and CEPC datasets

$Zh$  and VBF ( $h\nu\nu$ )

EWPOs

Z-pole EWPOs ( $\sqrt{s} = 91.2$ GeV)		
$\mathcal{O}_i$	$\delta/\Delta \mathcal{O}_i$	
	FCC-ee	CEPC
$\alpha(m_Z)^{-1} (\times 10^3)$	$\Delta = 2.7$ (1.2)	$\Delta = 17.8$
$\Gamma_W$ (MeV)	$\Delta = 0.85$ (0.3)	$\Delta = 1.8$ (0.9)
$\Gamma_Z$ (MeV)	$\Delta = 0.0028$ (0.025)	$\Delta = 0.005$ (0.025)
$A_e (\times 10^5)$	$\Delta = 0.5$ (2)	$\Delta = 1.5$
$A_\mu (\times 10^5)$	$\Delta = 1.6$ (2.2)	$\Delta = 3.0$ (1.8)
$A_\tau (\times 10^5)$	$\Delta = 0.35$ (20)	$\Delta = 1.2$ (6.9)
$A_b (\times 10^5)$	$\Delta = 1.7$ (21)	$\Delta = 3$ (21)
$A_c (\times 10^5)$	$\Delta = 14$ (15)	$\Delta = 6$ (30)
$\sigma_{\text{had}}^0$ (pb)	$\Delta = 0.025$ (4)	$\Delta = 0.05$ (2)
$R_e (\times 10^3)$	$\delta = 0.0028$ (0.3)	$\delta = 0.003$ (0.2)
$R_\mu (\times 10^3)$	$\delta = 0.0021$ (0.05)	$\delta = 0.003$ (0.1)
$R_\tau (\times 10^3)$	$\delta = 0.0021$ (0.1)	$\delta = 0.003$ (0.1)
$R_b (\times 10^3)$	$\delta = 0.001$ (0.3)	$\delta = 0.005$ (0.2)
$R_c (\times 10^3)$	$\delta = 0.011$ (1.5)	$\delta = 0.02$ (1)

$e^+e^- \rightarrow Zh$				
	$\sqrt{s} = 240$ GeV		$\sqrt{s} = 365$ GeV	
$\mathcal{O}_i$	$\delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\delta_{\text{exp}} \mathcal{O}_i$ (CEPC)	$\delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\delta_{\text{exp}} \mathcal{O}_i$ (CEPC)
$\sigma_{Zh}$	0.0035	0.0026	0.0064	0.014
$\sigma_{Zh} \times \text{BR}_{b\bar{b}}$	0.0021	0.0014	0.0035	0.009
$\sigma_{Zh} \times \text{BR}_{c\bar{c}}$	0.0156	0.0202	0.046	0.088
$\sigma_{Zh} \times \text{BR}_{gg}$	0.0134	0.0081	0.0247	0.034
$\sigma_{Zh} \times \text{BR}_{ZZ}$	0.0311	0.0417	0.0849	0.2
$\sigma_{Zh} \times \text{BR}_{WW}$	0.0085	0.0053	0.0184	0.028
$\sigma_{Zh} \times \text{BR}_{\tau^+\tau^-}$	0.0064	0.0042	0.0127	0.021
$\sigma_{Zh} \times \text{BR}_{\gamma\gamma}$	0.0636	0.0302	0.127	0.11
$\sigma_{Zh} \times \text{BR}_{\gamma Z}$	0.12	0.085	-	-
$e^+e^- \rightarrow h\nu\nu$				
	$\sqrt{s} = 240$ GeV		$\sqrt{s} = 365$ GeV	
$\mathcal{O}_i$	$\delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\delta_{\text{exp}} \mathcal{O}_i$ (CEPC)	$\delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\delta_{\text{exp}} \mathcal{O}_i$ (CEPC)
$\sigma_{h\nu\nu} \times \text{BR}_{b\bar{b}}$	0.0219	0.0159	0.0064	0.011
$\sigma_{h\nu\nu} \times \text{BR}_{c\bar{c}}$	-	-	0.0707	0.16
$\sigma_{h\nu\nu} \times \text{BR}_{gg}$	-	-	0.0318	0.045
$\sigma_{h\nu\nu} \times \text{BR}_{ZZ}$	-	-	0.0707	0.21
$\sigma_{h\nu\nu} \times \text{BR}_{WW}$	-	-	0.0255	0.044
$\sigma_{h\nu\nu} \times \text{BR}_{\tau^+\tau^-}$	-	-	0.0566	0.042
$\sigma_{h\nu\nu} \times \text{BR}_{\gamma\gamma}$	-	-	0.156	0.16

# FCC-ee and CEPC datasets

Light fermion production

$e^+e^- \rightarrow f\bar{f}$				
	$\sqrt{s} = 240 \text{ GeV}$		$\sqrt{s} = 365 \text{ GeV}$	
$\mathcal{O}_i$	$\Delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\Delta_{\text{exp}} \mathcal{O}_i$ (CEPC)	$\Delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\Delta_{\text{exp}} \mathcal{O}_i$ (CEPC)
$\sigma_{\text{tot}}(e^+e^-)$ [fb]	2.29	1.62	2.74	4.68
$A_{\text{FB}}(e^+e^-)$	$9.79 \cdot 10^{-6}$	$6.92 \cdot 10^{-6}$	$2.83 \cdot 10^{-5}$	$4.83 \cdot 10^{-5}$
$\sigma_{\text{tot}}(\mu^+\mu^-)$ [fb]	0.405	0.287	0.48	0.82
$A_{\text{FB}}(\mu^+\mu^-)$	$1.98 \cdot 10^{-4}$	$1.397 \cdot 10^{-4}$	$5.69 \cdot 10^{-4}$	$9.7 \cdot 10^{-4}$
$\sigma_{\text{tot}}(\tau^+\tau^-)$ [fb]	0.374	0.264	0.443	0.756
$A_{\text{FB}}(\tau^+\tau^-)$	$2.17 \cdot 10^{-4}$	$1.53 \cdot 10^{-4}$	$6.24 \cdot 10^{-4}$	0.00106
$\sigma_{\text{tot}}(c\bar{c})$ [fb]	0.088	0.062	0.102	0.175
$A_{\text{FB}}(c\bar{c})$	0.000813	$5.74 \cdot 10^{-4}$	0.00238	0.00405
$\sigma_{\text{tot}}(b\bar{b})$ [fb]	0.151	0.107	0.171	0.29
$A_{\text{FB}}(b\bar{b})$	$4.86 \cdot 10^{-4}$	$3.44 \cdot 10^{-4}$	0.00142	0.00243

$e^+e^- \rightarrow W^+W^-$						
$\mathcal{O}_i$	$\sqrt{s} = 161 \text{ GeV}$		$\sqrt{s} = 240 \text{ GeV}$		$\sqrt{s} = 365 \text{ GeV}$	
	$\delta_{\text{exp}}$ (FCC-ee)	$\delta_{\text{exp}}$ (CEPC)	$\delta_{\text{exp}}$ (FCC-ee)	$\delta_{\text{exp}}$ (CEPC)	$\delta_{\text{exp}}$ (FCC-ee)	$\delta_{\text{exp}}$ (CEPC)
$\sigma_{WW}$	$1.36 \cdot 10^{-4}$	$2.48 \cdot 10^{-4}$	$1.22 \cdot 10^{-4}$	$8.63 \cdot 10^{-5}$	$2.81 \cdot 10^{-4}$	$4.87 \cdot 10^{-4}$
$\text{BR}_{W \rightarrow \ell_i \nu_i}$	$2.72 \cdot 10^{-4}$	$4.95 \cdot 10^{-4}$	$2.44 \cdot 10^{-4}$	$1.73 \cdot 10^{-4}$	$5.63 \cdot 10^{-4}$	$9.75 \cdot 10^{-4}$