Integrable Quantum Field Theories, Irrelevant Perturbations and Minimal Form Factors

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My talk today is based on the following preprints/papers:

1. Olalla Castro-Alvaredo, Stefano Negro and Fabio Sailis, Completing the Bootstrap Program for $T\bar{T}$ -Deformed Massive Integrable Quantum Field

2. Olalla Castro-Alvaredo, Stefano Negro and Fabio Sailis, Form Factors and Correlation Functions of $T\bar{T}$ -Deformed Integrable Quantum Field Theories,

3. Olalla Castro-Alvaredo, Stefano Negro and Fabio Sailis, Entanglement Entropy from Form Factors in $T\bar{T}$ -Deformed Integrable Quantum Field Theories, JHEP

- Theories, J. Phys. **A57** 265401 (2024); ArXiv:2305.17068
- JHEP **09** 2023 048; ArXiv:2306.01640
- **11** 2023 129; ArXiv:2306.11064.
- Nucl. Phys. B1000 (2024) 116459. ArXiv: 2311.16955

4. Olalla Castro-Alvaredo, Stefano Negro and István M. Szécsényi, On the Representation of Minimal Form Factors in Integrable Quantum Field Theory,

As you can see, this is work with Stefano Negro, Fabio Sailis and István M. Szécsényi

PhD Student at City, University of London (UK)

Lecturer, University of York (UK)

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C provides a nice way to interpret some of our results (see later)

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[Cardy & Doyon'10]

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F_n^{\mathcal{O}}(\theta_1, ..., \theta_i, \theta_{i+1}, ..., \theta_n; \boldsymbol{\alpha}) = S_{\boldsymbol{\alpha}}(\theta_i - \theta_{i+1}) F_n^{\mathcal{O}}(\theta_1, ..., \theta_{i+1}, \theta_i, ..., \theta_n; \boldsymbol{\alpha})
$$

 $F_n^{\mathcal{O}}(\theta_1 + 2\pi i, \theta_2 \dots, \theta_n; \boldsymbol{\alpha}) = \gamma_{\mathcal{O}} F_n^{\mathcal{O}}$

 $\lim_{\bar{\theta}\to 0} (\bar{\theta}-\theta) F_{n+2}^{\mathcal{O}}(\bar{\theta}+i\pi,\theta,\theta_1,\ldots,\theta_n;\boldsymbol{\alpha})=i\left[1-\gamma_{\mathcal{O}}\right]$ $\bar{\theta} \rightarrow \theta$

$$
\ldots, \theta_n \rangle
$$

$$
=\gamma_{\mathcal{O}}F_n^{\mathcal{O}}(\theta_2,\ldots,\theta_n,\theta_1;\boldsymbol{\alpha})
$$

$$
= i \left(1 - \gamma_{\odot} \prod_{j=1}^{n} S_{\alpha}(\theta - \theta_{j}) \right) F_{n}^{\odot}(\theta_{1}, ..., \theta_{n}; \alpha)
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• The most general solution to this equation is:

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	- 2. The solution depends on additional parameters not present in the S-matrix. In a sense, it contains its own CDD factor!
	- 3. Understanding what the different choices of β s actually mean is still an open question (more on this later).

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• For some theories and fields, it seems that higher particle form factors also factorise. We have shown this to be the case quite generally, and worked out the details for the

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F_n^{\mathcal{O}}(\theta_1, ..., \theta_n; \boldsymbol{\alpha}) = F_n^{\mathcal{O}}(\theta_1, ..., \theta_n; \mathbf{0}) G_n^{\mathcal{O}}(\theta_1, ..., \theta_n; \boldsymbol{\alpha})
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	- 1. This is quite general for symmetry fields (only constraint is that $\gamma^{0} = \pm 1$)
	- 2. The function $G_n^{\mathscr{O}}(\theta_1,...,\theta_n;\bm{x})$ can be written totally explicitly. It is typically an oscillatory function of the S-matrix and of $\gamma^{\mathscr{O}}$ times a product of functions $\varphi(\theta; \boldsymbol{\alpha}).$ $\theta_n^0(\theta_1, ..., \theta_n; \boldsymbol{\alpha})$

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	- 3. For local fields there can be an additional non-trivial multiplicative factor (like Θ).

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[Yurov & Zamolodchikov'91]

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F_{2n}^{\mu}(\theta_1, ..., \theta_{2n}; \alpha) = i^{n} \langle \mu \rangle_{\alpha} \sqrt{\prod_{i=1}^{2n} \cos \left(\sum_{s \in \mathcal{S}} \frac{\alpha_s}{2} \sum_{j=1}^{2n} \sinh(\theta_j) \right)}
$$

 $\sinh(s\theta_{ij})$ | \prod *i*<*j* tanh *θij* $\frac{J}{2}\varphi(\theta_{ij})$

; *^α*) [Yurov & Zamolodchikov'91]

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$$

 $F_{2n+1}^{\sigma}(\theta_1, ..., \theta_{2n+1}; \alpha) = i$

$$
= i^{n} F_{1}^{\sigma}(\boldsymbol{\alpha}) \sqrt{\prod_{i=1}^{2n+1} \cos \left(\sum_{s \in \mathcal{S}} \frac{\alpha_{s}}{2} \sum_{j=1}^{2n+1} \sinh(s\theta_{ij})\right) \prod_{i < j} \tanh \frac{\theta_{ij}}{2} \varphi(\boldsymbol{\alpha})}
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[Yurov & Zamolodchikov'91]

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[Yurov & Zamolodchikov'91]

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\Theta_n^{\emptyset}(\{\theta_i\}_n; \alpha) = \prod_{i=1}^n \sqrt{\prod_{j=1}^n S_{\alpha}(\theta_{ij})^{1/2} - \gamma_{\emptyset} \prod_{j=1}^n S_{\alpha}(\theta_{ij})^{-1/2}} = \prod_{i=1}^n \sqrt{\frac{\sin \left(\frac{1}{2} \sum_{j=1}^n \left[\delta(\theta_{ij}) - i \log \Phi_{\alpha}(\theta_{ij}) \right] - \frac{\omega_{\emptyset}}{2} \right)}{\sin \left(\frac{1}{2} \sum_{j=1}^n \delta(\theta_{ij}) - \frac{\omega_{\emptyset}}{2} \right)}}
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$$

• Here $S_0(\theta) = e^{i\delta(\theta)}$ and $\gamma^{\circ} = e^{i\omega_{\theta}}$.

• Form factors are building blocks for correlation functions in the ground state so we can now study those, particularly their long and short-distance behaviours.

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• These statements apply both to mr large and small, although for $\alpha^{\star} > 0$ and mr sufficiently large there is a rapidity cut-off that can make the series convergent.

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- An intuitive picture emerges in conjunction with [Cardy & Doyon'20]:
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	- 2. When $\alpha^{\star} < 0$ the fundamental excitations acquire a "negative" length so both the UV and IR can probed. In a sense "extra space" is created at short distances.

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•For $\alpha < 0$ correlation functions exhibit power-law scaling at short distances, just like in CFT. The coefficients however, are now functions of α and there is no underlying CFT.

How do these Powers Behave?

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Here $z_2^\mathscr{O}(\alpha) = 4\Delta^\mathscr{O}(\alpha)$ in the two-particle approximation. The field ε is the "energy" field in the Ising model which is proportional to $\Theta.$

How it All fits Together: Minimal Form Factors

• We have just learned about a particular type of CDD factor:

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 $S_{\alpha}(\theta) = S_0(\theta) \Phi_{\alpha}(\theta), \ \alpha = {\alpha_1, \alpha_2, \dots}$ $\Phi_{\alpha}(\theta) = \exp[-i\theta]$ ∑ *s*∈ $\alpha_s m^{2s} \sinh(s\theta)$]

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• We know however that an S-matrix such as that of the sinh-Gordon model is also a CDD factor.

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S(\theta) = \frac{\tanh\frac{1}{2}\left(\theta - \frac{i\pi B}{2}\right)}{\tanh\frac{1}{2}\left(\theta + \frac{i\pi B}{2}\right)}
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$$
-i\frac{d}{d\theta}\log S(\theta) = 4\sum_{k=0}^{\infty} (-1)^k \cos\frac{(2k\theta)^k}{k!}
$$

$$
S(\theta) = -\exp\left[-4i\theta\right]
$$

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with
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S_0(\theta) = -1
$$
 and
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• The sinh-Gordon model can be seen as the Ising model perturbed by an infinite number of irrelevant perturbations for carefully chosen couplings [LeClair'21; Ahn & LeClair'22]

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 $= e^{-\frac{\vartheta}{2a}}$ 2*π i* log Φ*α*(*θ*)

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\varphi(\theta; \alpha) = e^{-\frac{i\pi - \theta}{2\pi} \sum_{s \in \mathcal{S}} \alpha_s m^{2s} \sinh(s\theta) + \sum_{s \in \mathcal{S}'} \beta_s m^{2s} \cosh(s\theta)}
$$

$$
\log \Phi_{\alpha}(\theta) C_{\beta}(\theta) \quad \text{with} \quad \theta = i\pi - \theta \text{ and } C_{\beta}(\theta) = e^{\sum_{s} \beta_s m^{2s} \cosh(s\theta)}
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	-

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- There are several properties of interest:
	- 1. The solution is very simple and general. The minimal form factors gets modified by a universal multiplicative factor, just like the S-matrix!
	- 2. The solution depends on additional parameters not present in the S-matrix. In a sense, it contains its own CDD factor

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$$
F_{\min}^{\text{SG}}(\theta) = F_{\min}^{\text{Ising}}(\theta)e^{-\frac{\vartheta}{2\pi}i\log(\theta)}
$$

 $i \log \Phi_{\alpha}^{\text{SG}}(\theta)$ *C*sG $\beta^{\text{SU}}(\theta)$ with $\theta = i\pi - \theta$

- representation as well.
- This is indeed the case:

• Let's have a closer look at all these functions…..

$$
F^{\rm SG}_{\rm min}(\theta) = F^{\rm Ising}_{\rm min}(\theta)e^{-\frac{\theta}{2\pi}i\log\Phi_{\alpha}^{\rm SG}(\theta)}C_{\beta}^{\rm SG}(\theta) \quad \text{with} \quad \theta = i\pi - \theta
$$

The Formula

$$
\omega(\vartheta) = \frac{1}{2}\log 2 + \log \cosh \frac{\vartheta}{2} - \frac{1+b}{4}\log \left[\cosh \vartheta + \sin \frac{\pi b}{2}\right] - \frac{1-b}{4}\log \left[\cosh \vartheta - \sin \frac{\pi b}{2}\right]
$$

$$
-\frac{i\vartheta}{2\pi}\log \left[\frac{i\cos \frac{\pi b}{2} - \sinh \vartheta}{i\cos \frac{\pi b}{2} + \sinh \vartheta}\right] - \frac{i}{4\pi}\left[\text{Li}_2\left(-ie^{\vartheta - i\frac{\pi}{2}b}\right) - \text{Li}_2\left(ie^{\vartheta - i\frac{\pi}{2}b}\right) + \right]
$$

$$
+\text{Li}_2\left(-ie^{\vartheta + i\frac{\pi}{2}b}\right) - \text{Li}_2\left(ie^{\vartheta + i\frac{\pi}{2}b}\right) + \left(\vartheta \to -\vartheta\right)\right] \quad \text{with} \quad \vartheta = i\pi - \theta
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The Formula

$log F_{min}^{\text{SG}}(\theta)$

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$$
\n
$$
-\frac{i\vartheta}{2\pi} i \log \Phi_{\alpha}^{\text{SG}}(\vartheta)
$$

Formula	"Integration
constant" fixed by asymptotics	
$\cosh \vartheta + \sin \frac{\pi b}{2} - \frac{1-b}{4} \log \left[\cosh \vartheta - \sin \vartheta \right]$	
$\text{Li}_2 \left(-ie^{\vartheta - i\frac{\pi}{2}b} \right) - \text{Li}_2 \left(ie^{\vartheta - i\frac{\pi}{2}b} \right) +$	
$+(\vartheta \to -\vartheta) \left \right $ with $\vartheta = i\pi - \theta$	
$\log(C_{\beta}^{SG}(\vartheta))$	

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\beta = b-1
$$

$$
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$$

 $\times \exp \left[8 \int_0^\infty \frac{dx}{x} \frac{\sinh\left(\frac{xB}{4}\right) \sinh\left(\frac{x}{2}(1 - \frac{B}{2})\right) \sinh\frac{x}{2}}{\sinh^2 x} (N + 1 - N e^{-2x}) e^{-2Nx} \sin^2\left(\frac{x\hat{\beta}}{2\pi}\right)\right]$

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\omega(\vartheta) = \frac{1}{2}\log 2 + \log \cosh \frac{\vartheta}{2} - \frac{1+b}{4}\log \left[\frac{i\cos \frac{\pi b}{2} - \sinh \vartheta}{i\cos \frac{\pi b}{2} + \sinh \vartheta} - \frac{i}{2\pi}\right]
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How does it Compare?

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• A byproduct of our investigation has been finding a new representation for the MFFs of standard IQFTs. This representation is totally explicit, convergent and numerically efficient, given in terms of elementary and (a finite number of) special functions.

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• Having this representation for sinh-Gordon means that we effectively have it for every diagonal IQFT since the S-matrix and MFF of sinh-Gordon is a "standard block" for more complicated theories [Dorey, Exact S-Matrices'98; Mussardo, Book'10]

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- asymptotics. Understanding its role for $T\bar{T}$ -models remains an open question.

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