

16 April 2024

**Davide Racco**

**ETH** zürich



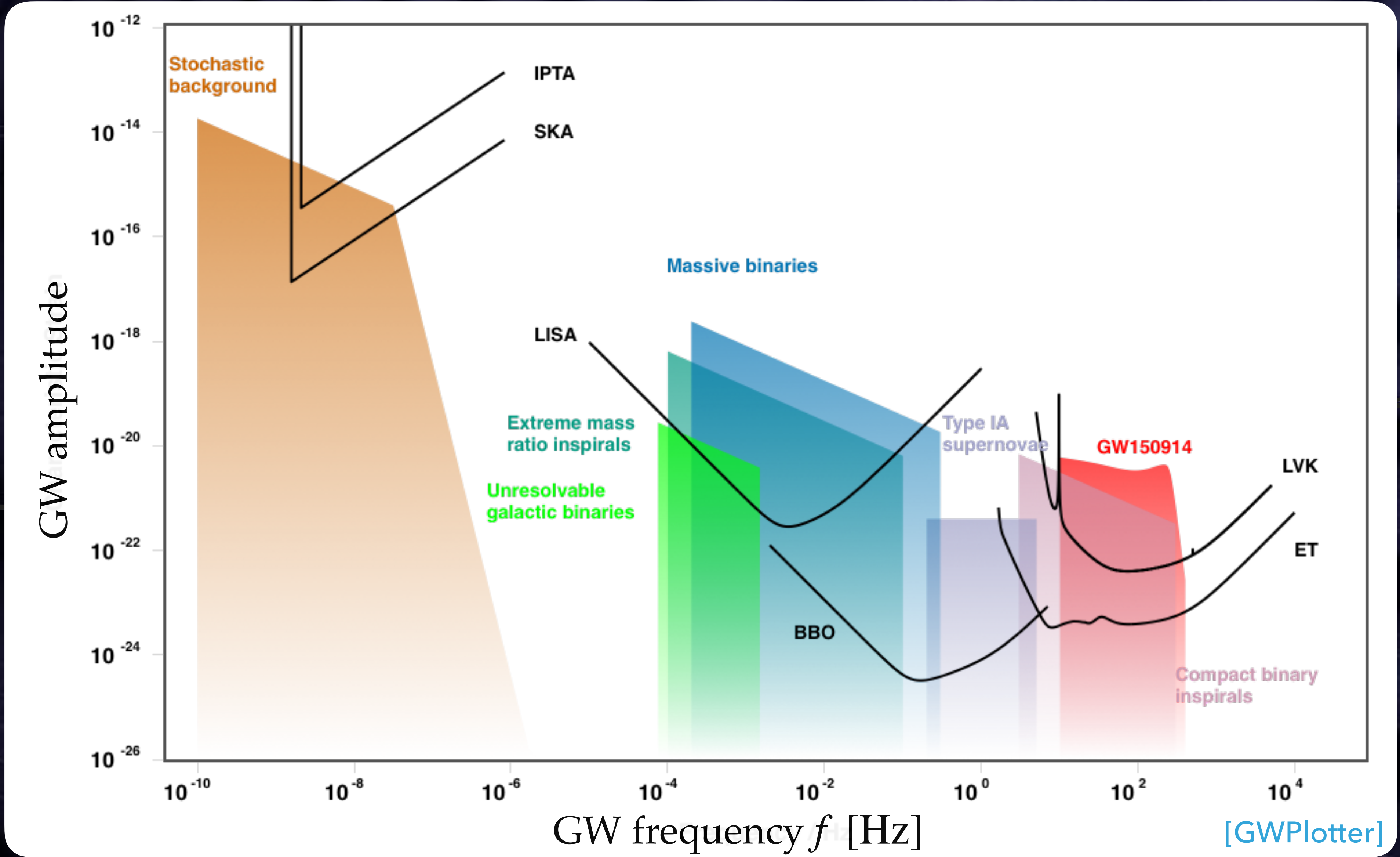
Universität  
Zürich <sup>UZH</sup>

LAFPTA

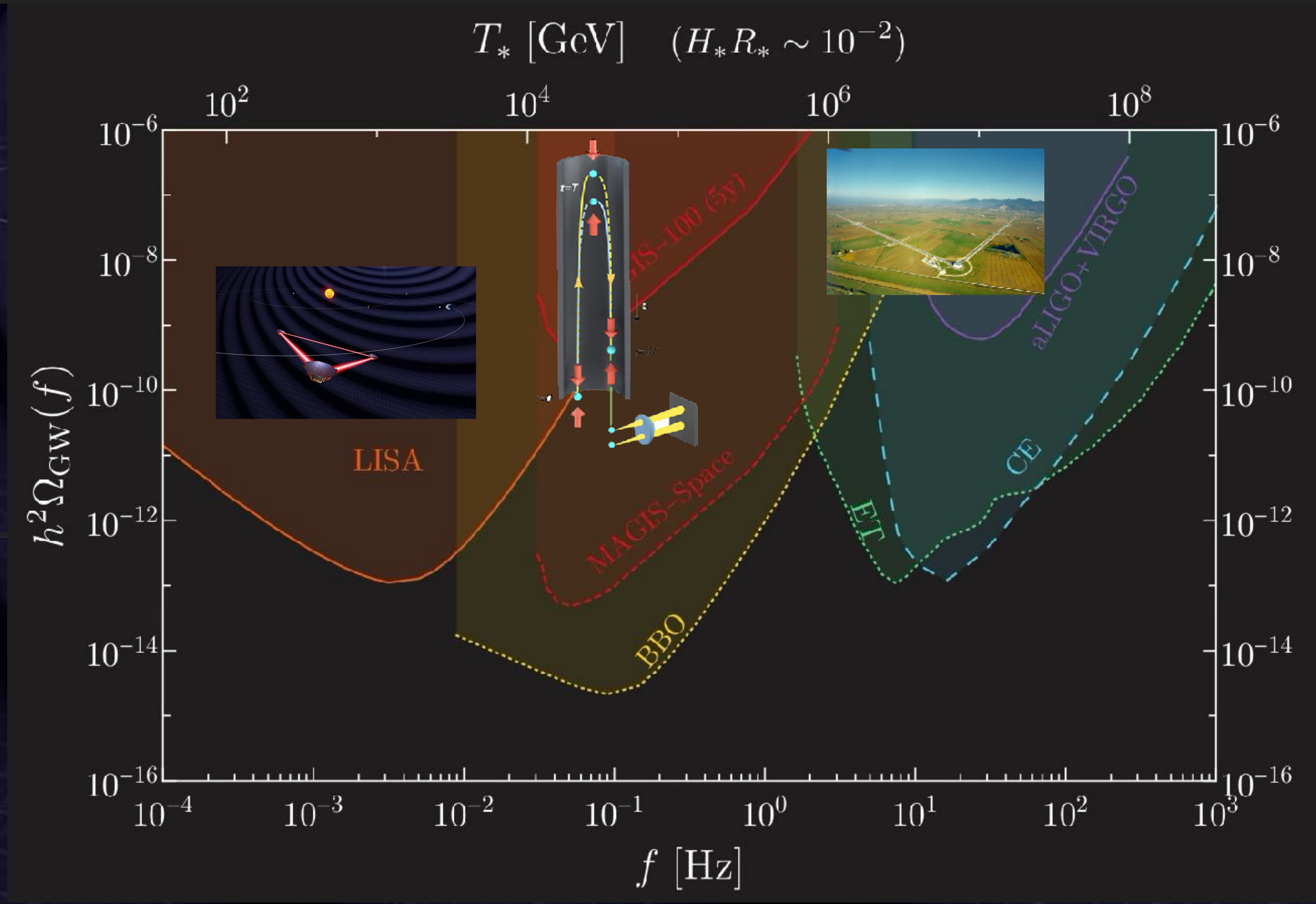
Laboratoire de  
Physique Théorique  
Annecy

Utilising the Causal  
Spectrum of GW  
Backgrounds: Implications  
for Pulsar Timing Arrays





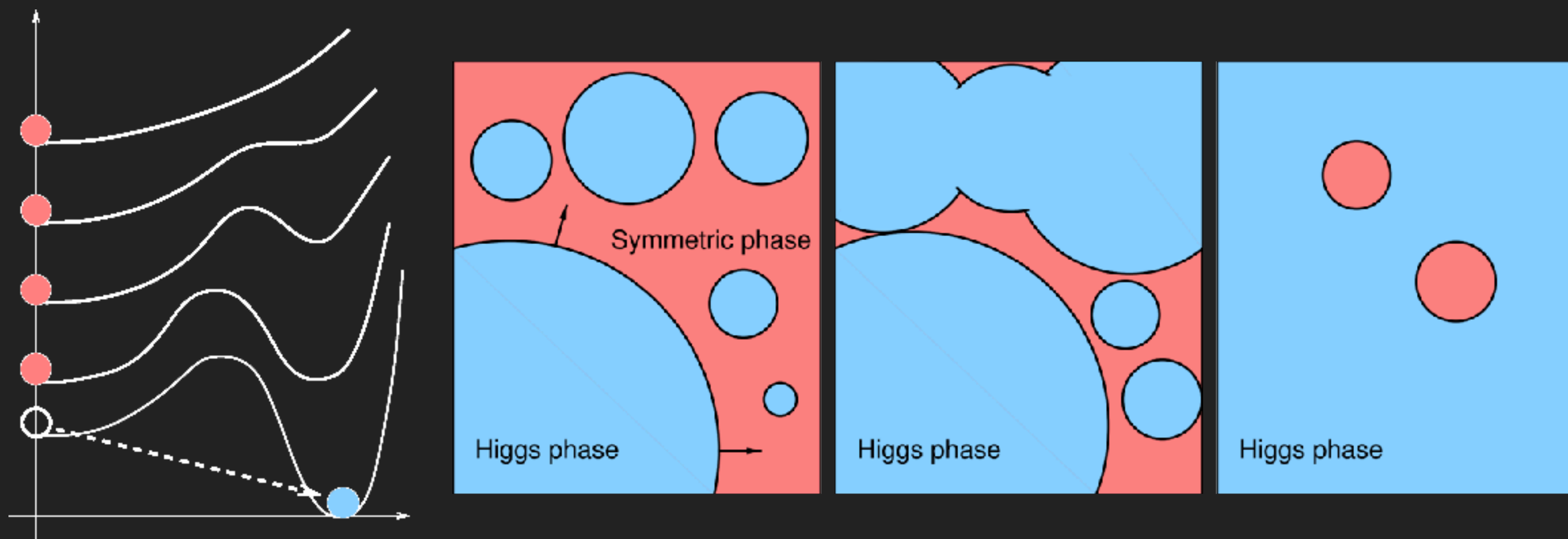




- ▶ GWs from primordial source **local** in *space* and *time*
- ▶ Phase transitions, preheating, peak in  $\mathcal{P}_\zeta$  ...

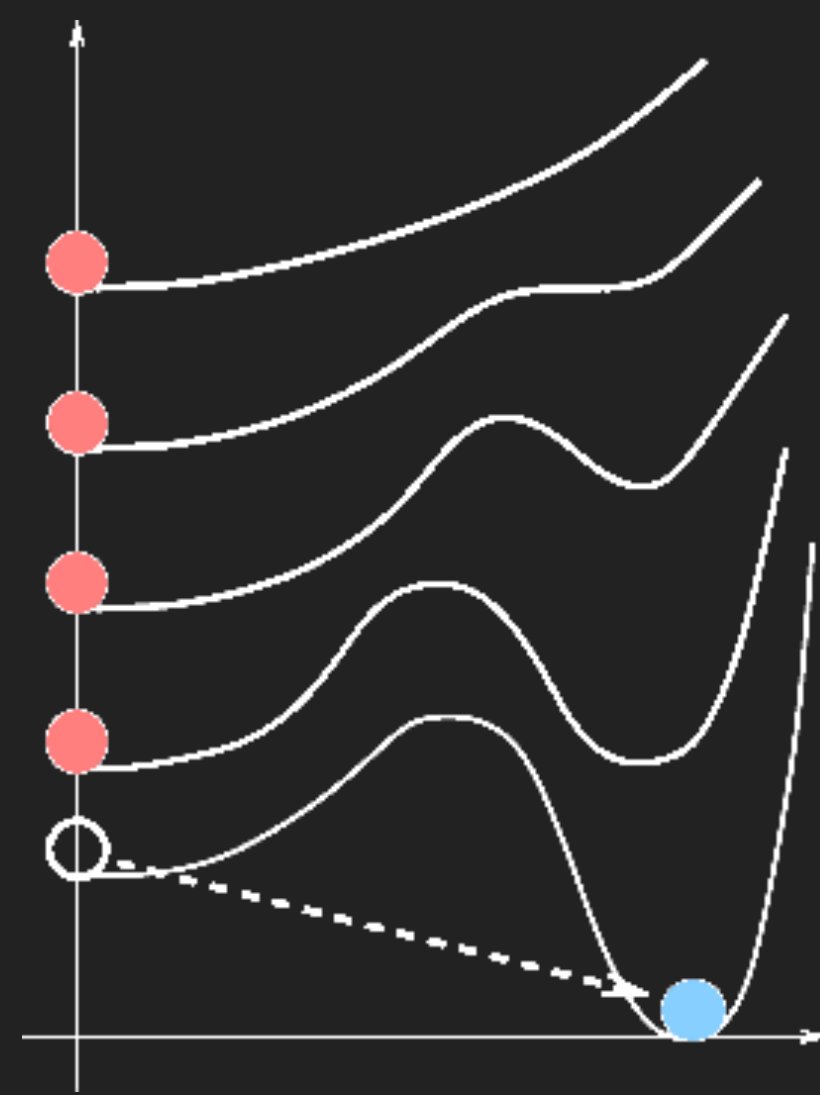


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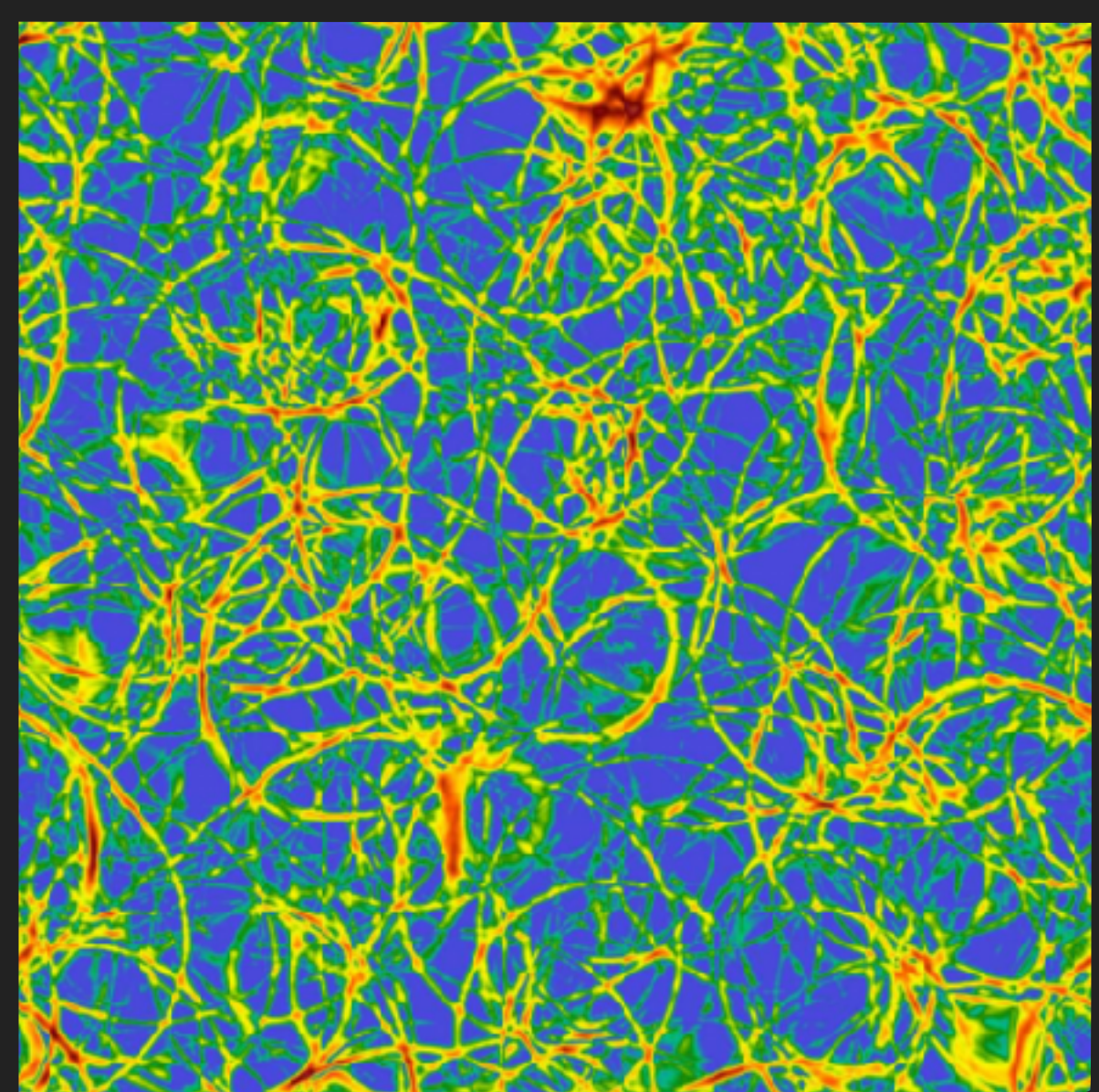
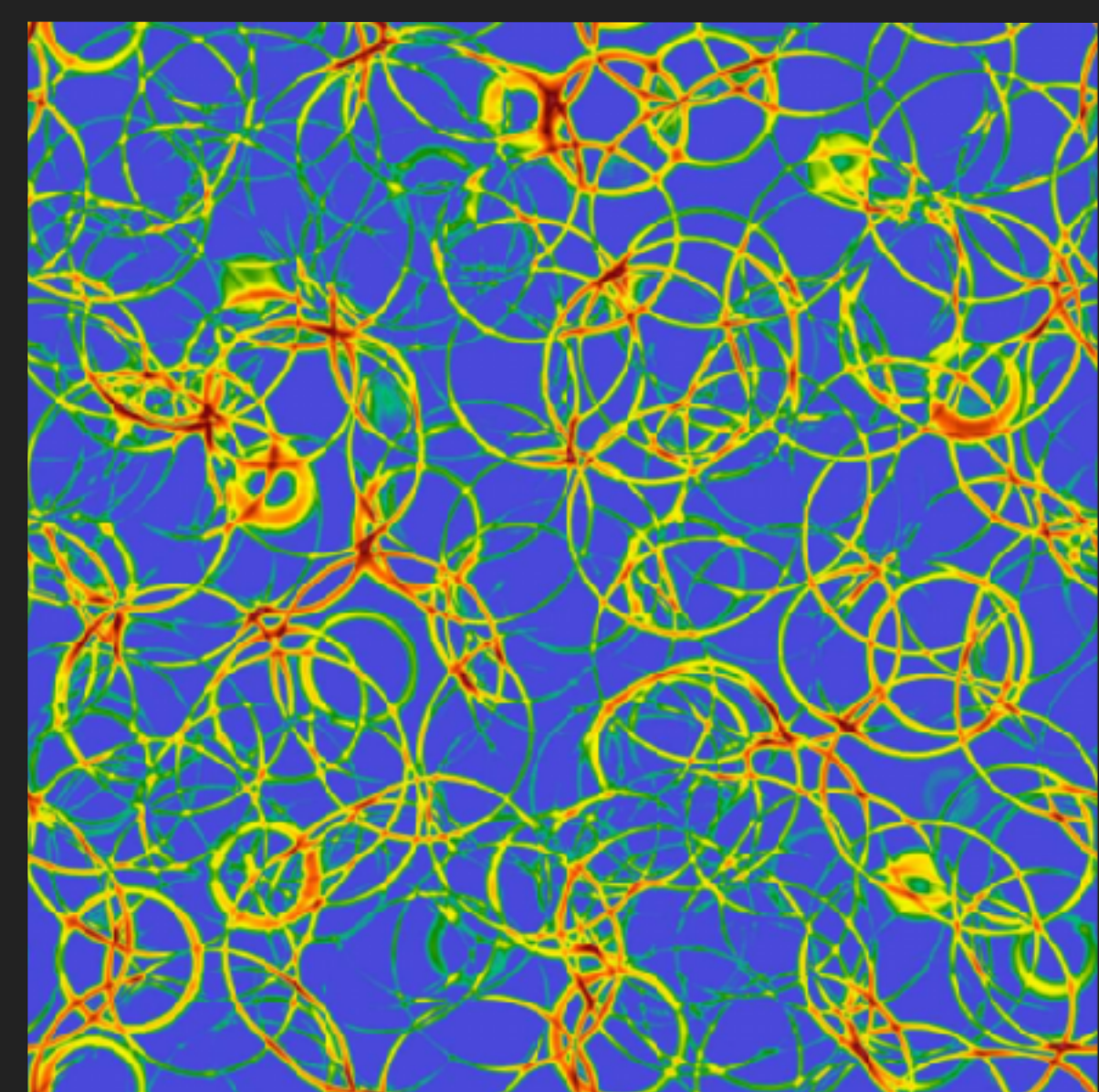
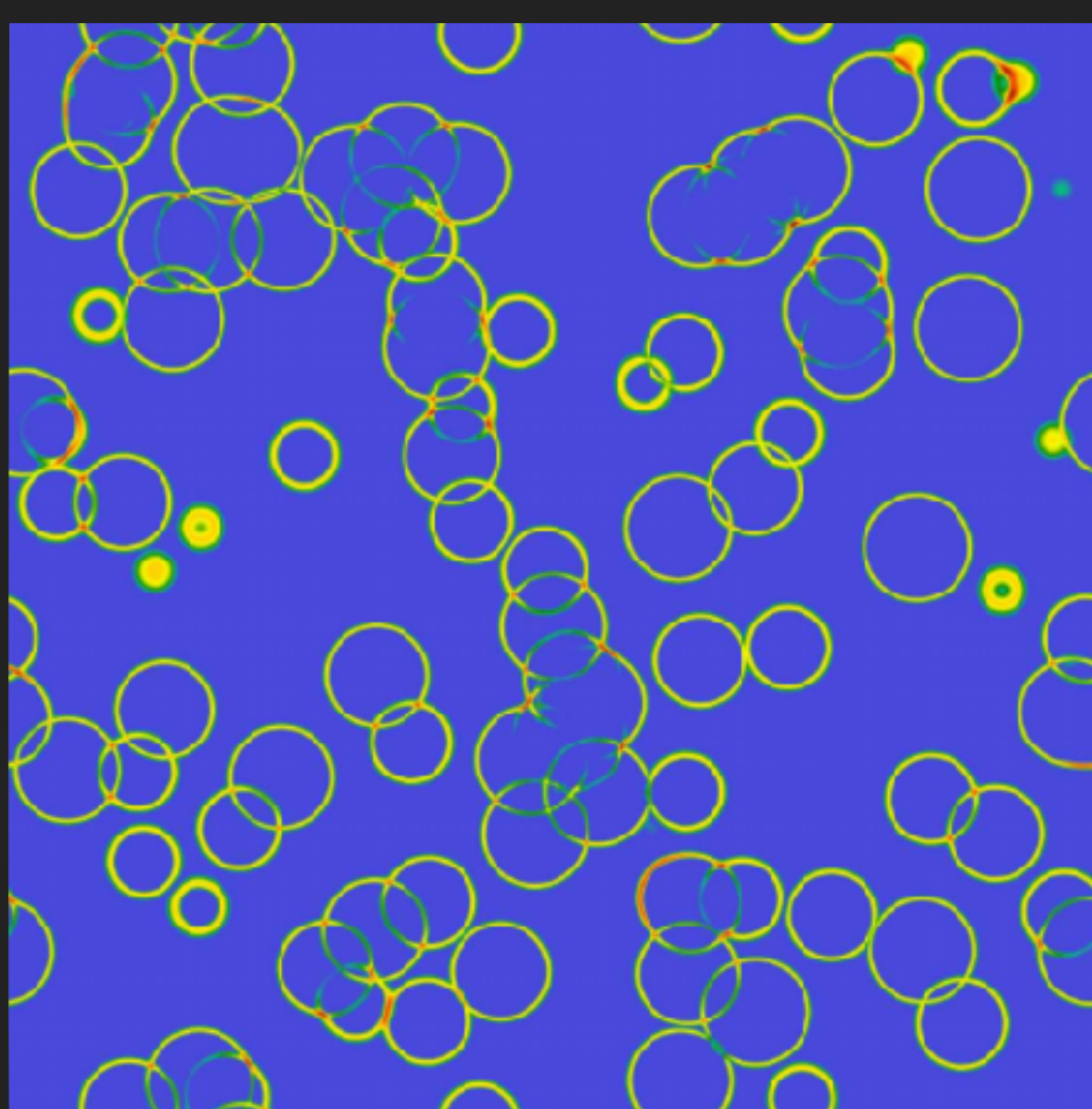
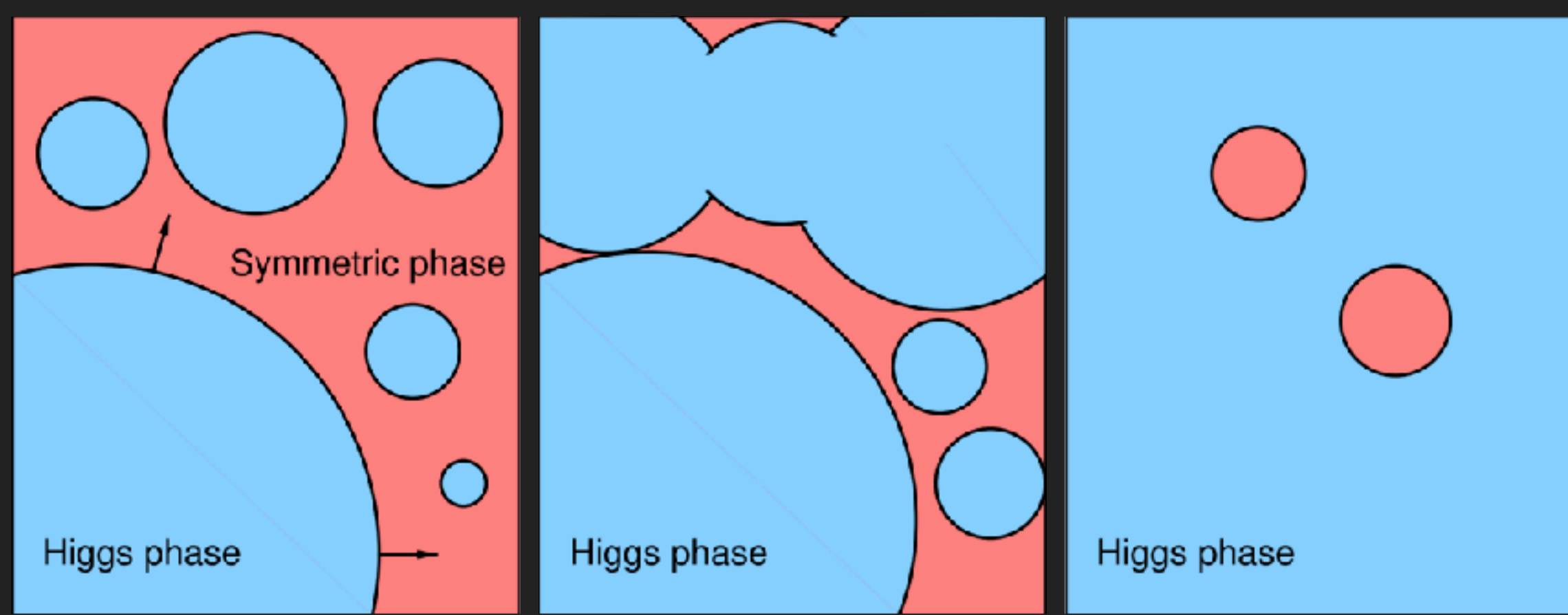




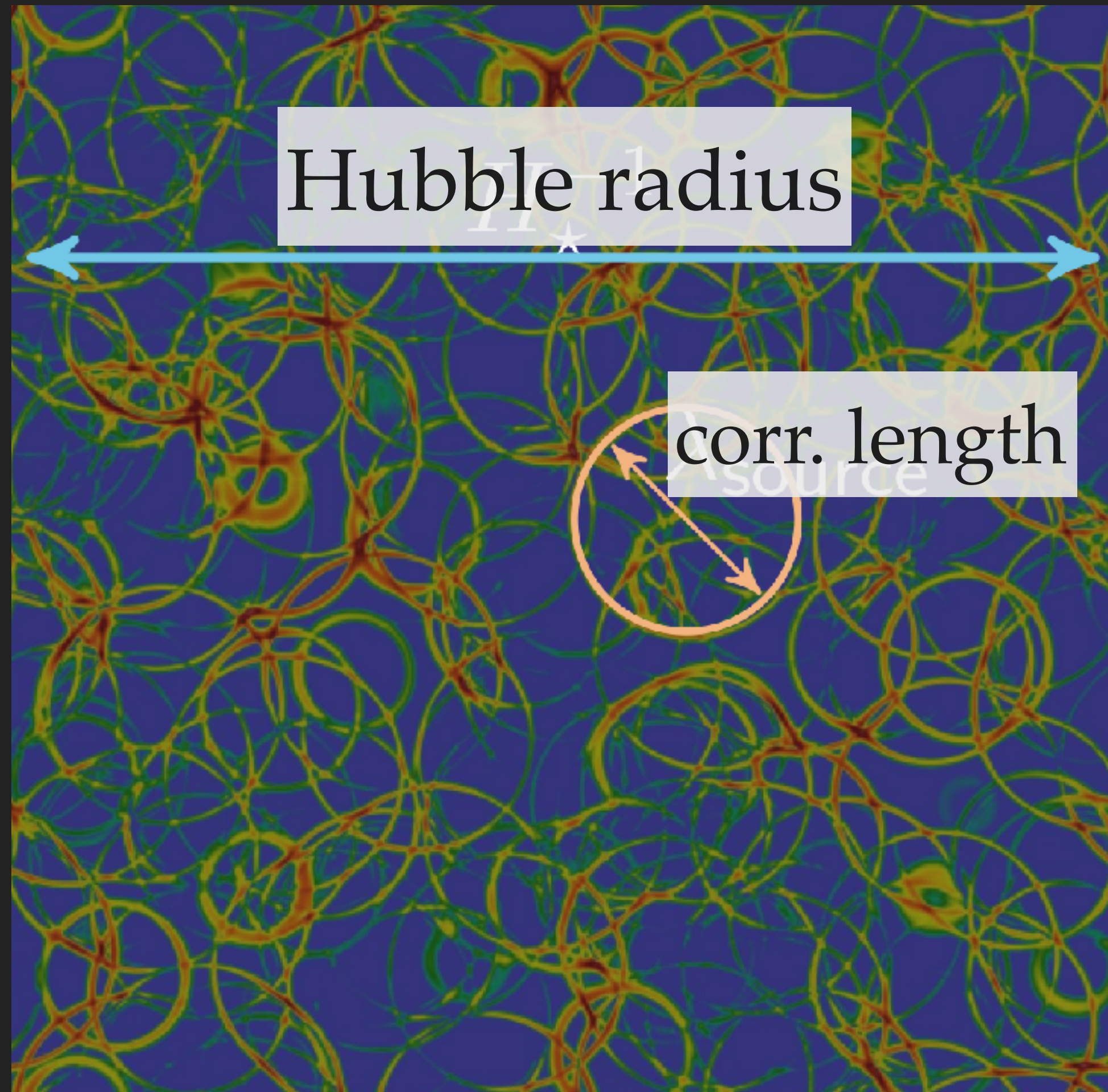
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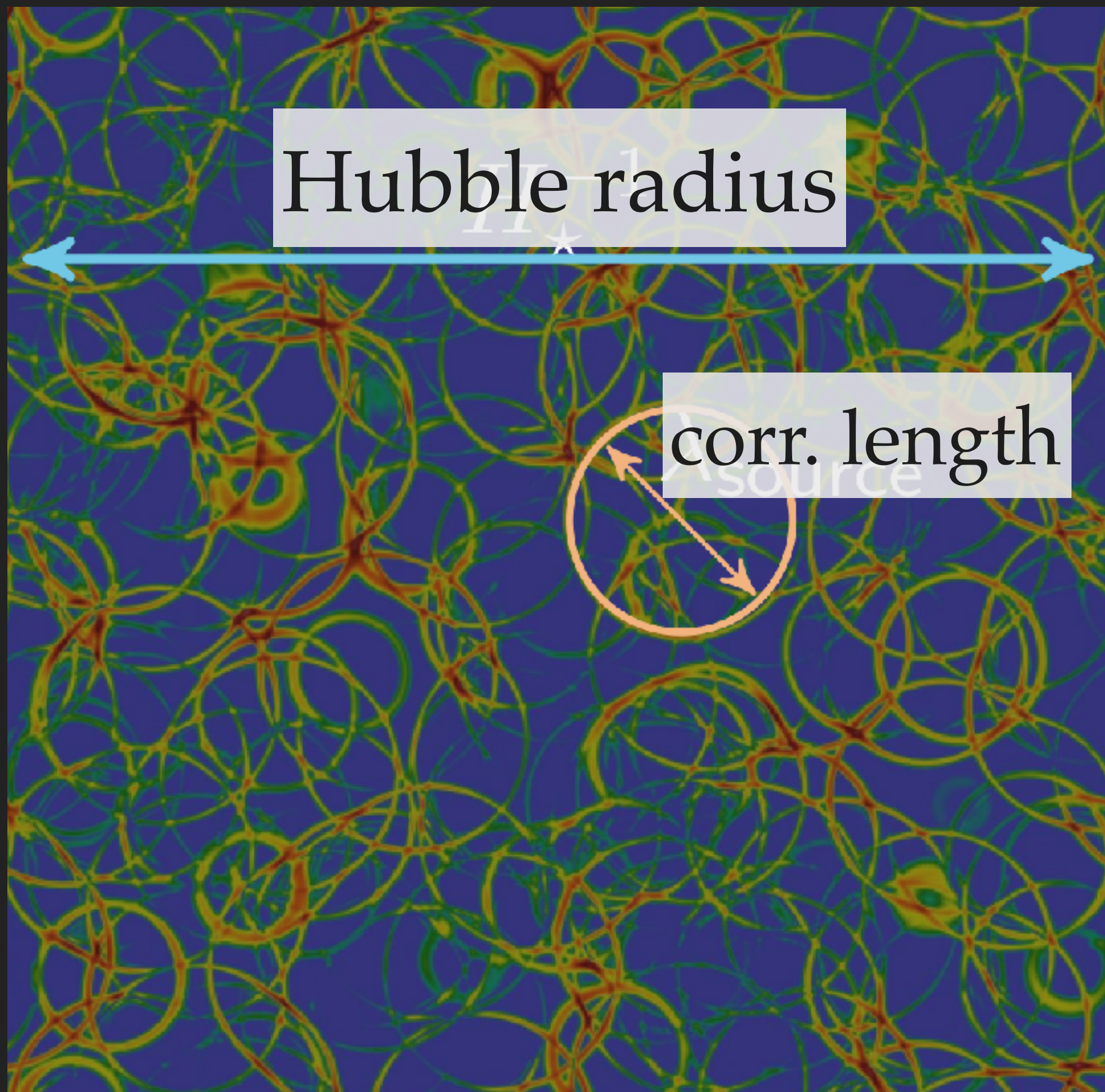
['15 Hindmarsh+]







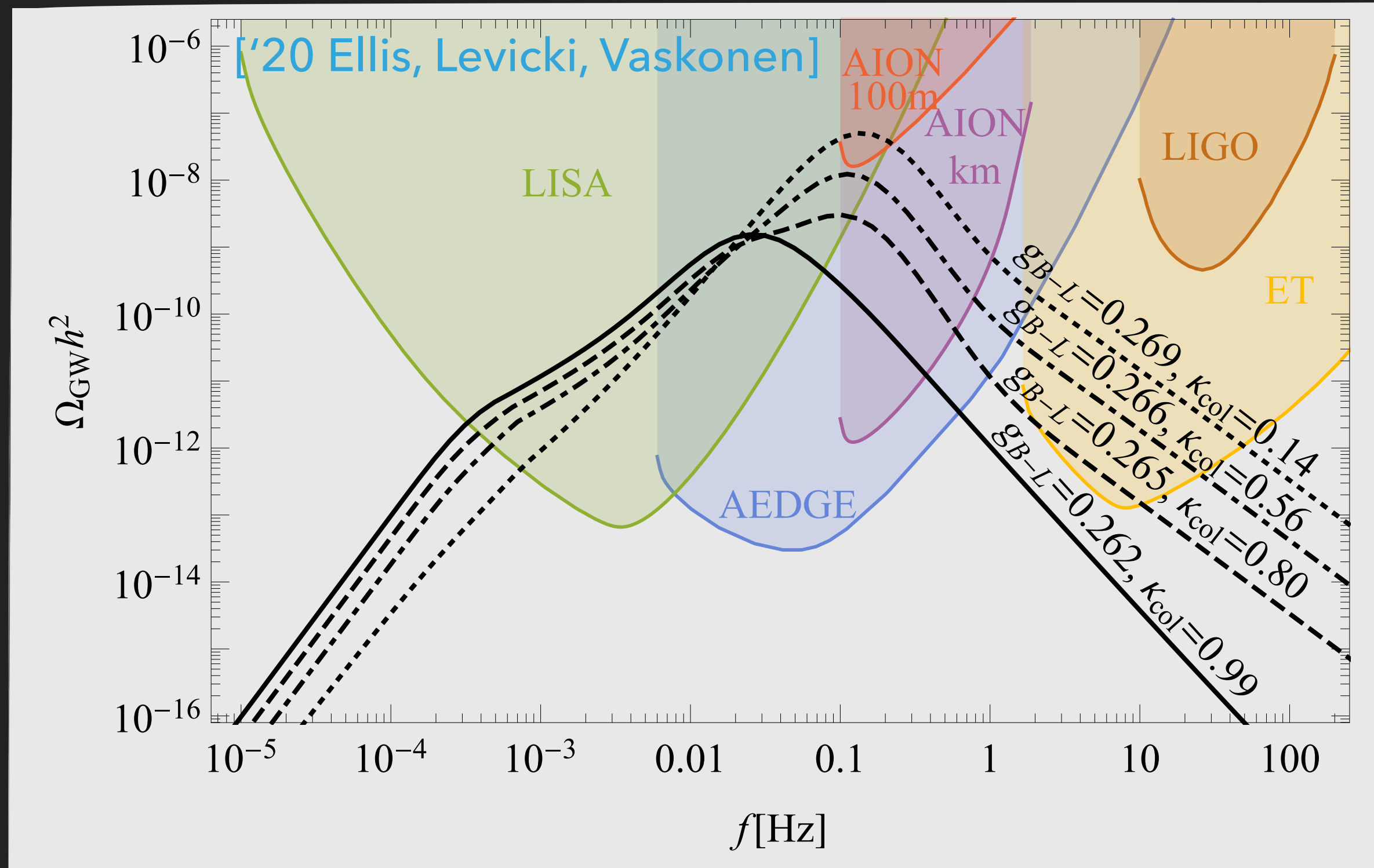




$$\langle \Pi(0) \Pi(d \gg \lambda_{\text{source}}) \rangle = 0 \implies$$

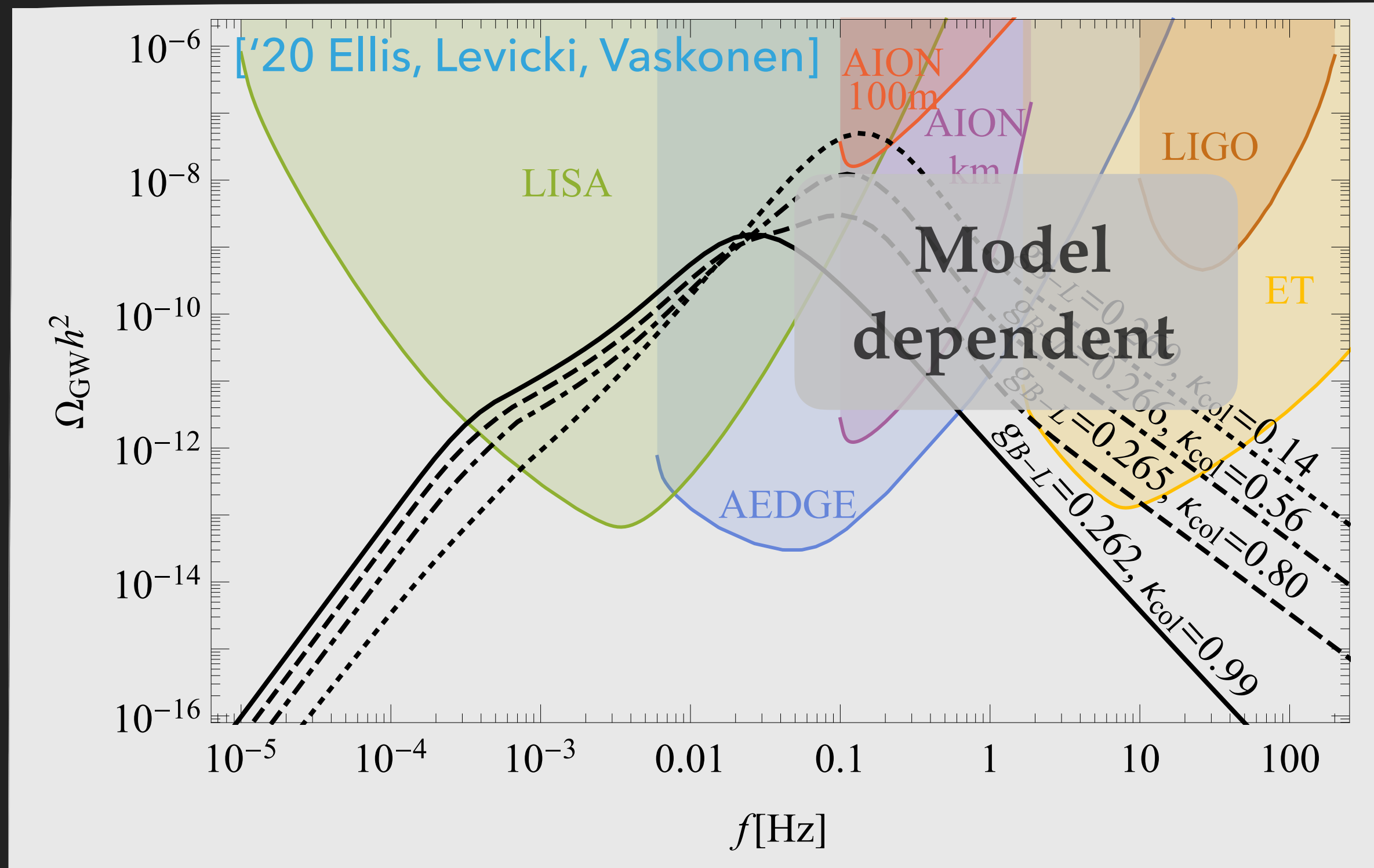
$$\langle \tilde{\Pi}(k) \tilde{\Pi}(-k) \rangle \xrightarrow{k \ll \lambda_{\text{source}}^{-1}} \text{constant}$$





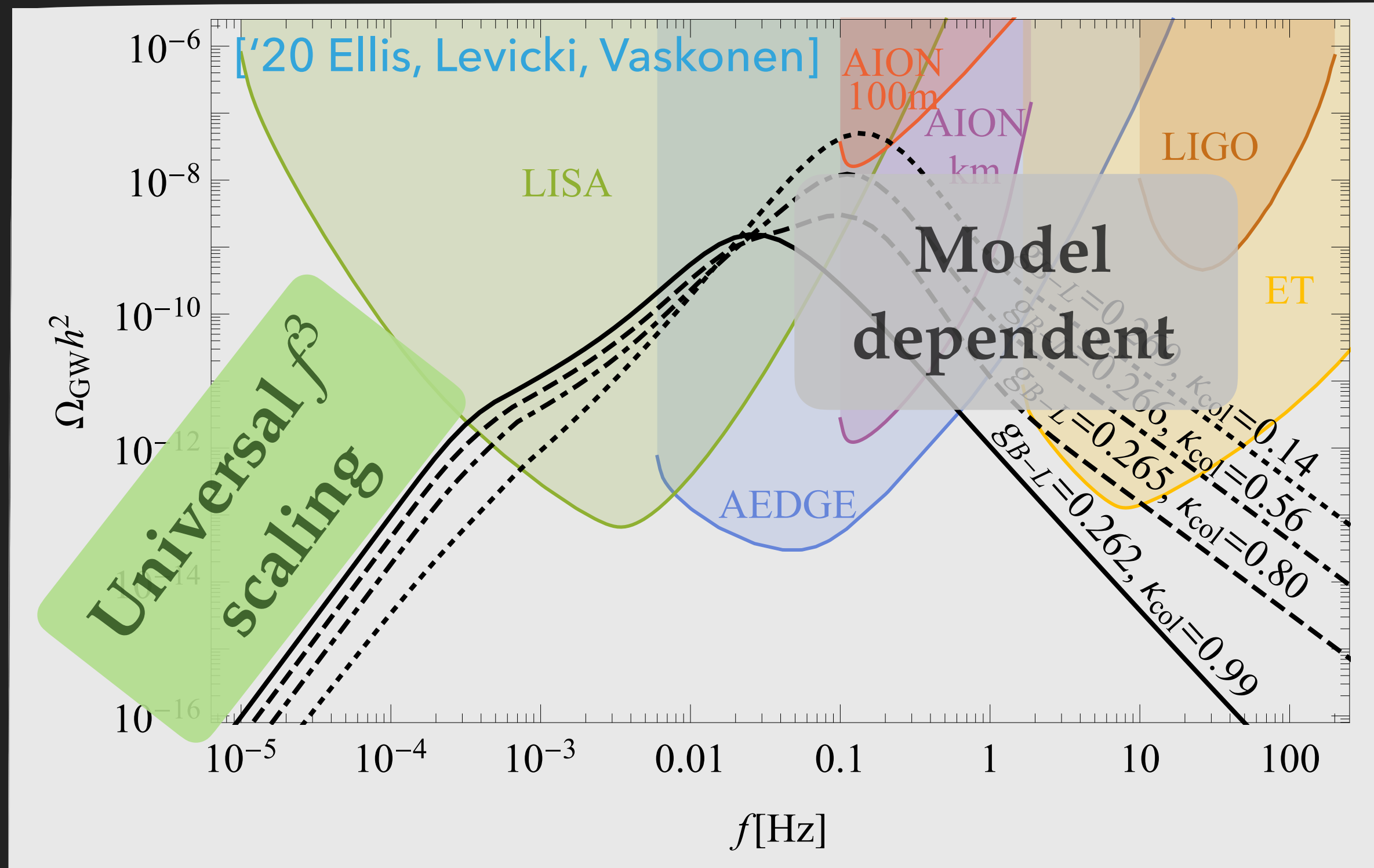
- [ '03 Seto, Yokoyama;
- '05 Boyle, Steinhardt;
- '06 Watanabe, Komatsu;
- '09 Caprini, Durrer, Konstandin, Servant;
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- '18 Cui, Lewicki, Morrissey, Wells;
- '19 D'Eramo, Schmitz;
- ...]





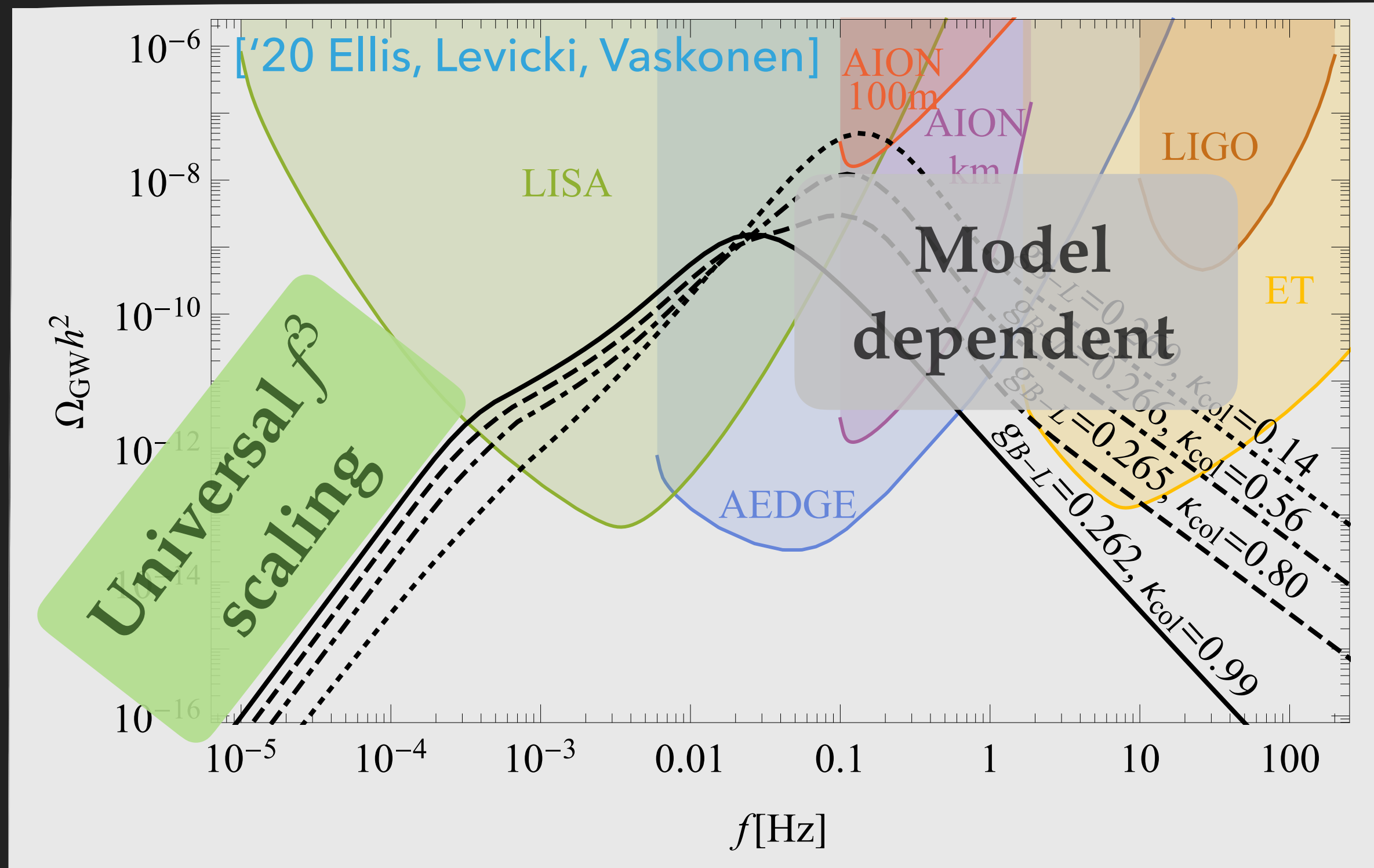
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**Causality tail**

▶ Low- $f$  spectrum GWs  $\Leftrightarrow$  universal (model-independent) probe of the Universe



wavelength  $k^{-1} \gg$  corr. length  $\lambda_{\text{source}}$

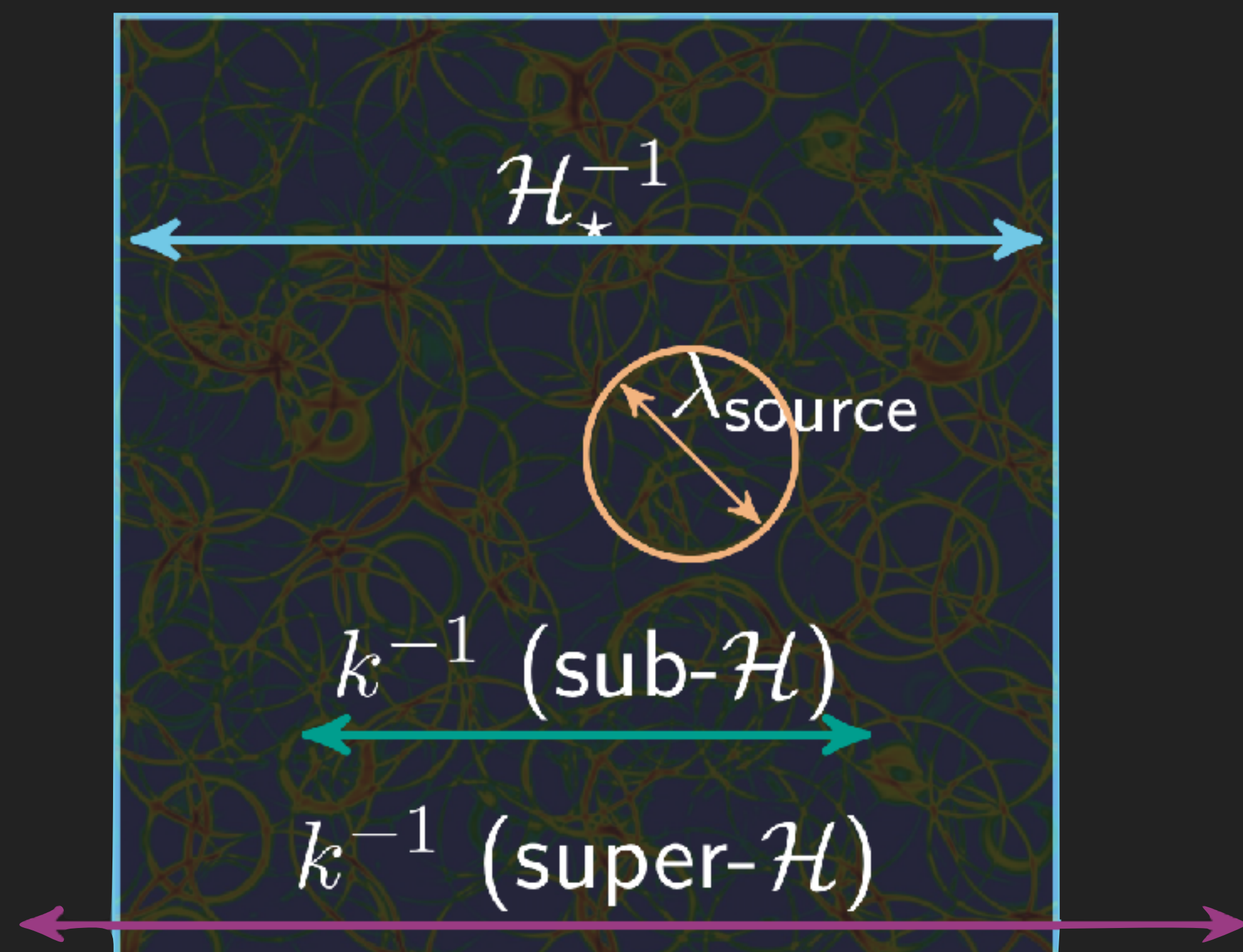
period  $f^{-1} \gg$  duration of phase transition  $\beta^{-1}$

$$\square h = J_{\star} \delta(\tau - \tau_{\star}) \implies \begin{cases} h = 0 \\ h' = J_{\star} \end{cases}$$

$$|h(\tau)| \sim \underbrace{\frac{a_{\star}}{a}}_{\text{redshift}} \cdot \underbrace{\frac{1}{k}}_{\text{for RD}} \cdot J_{\star}$$

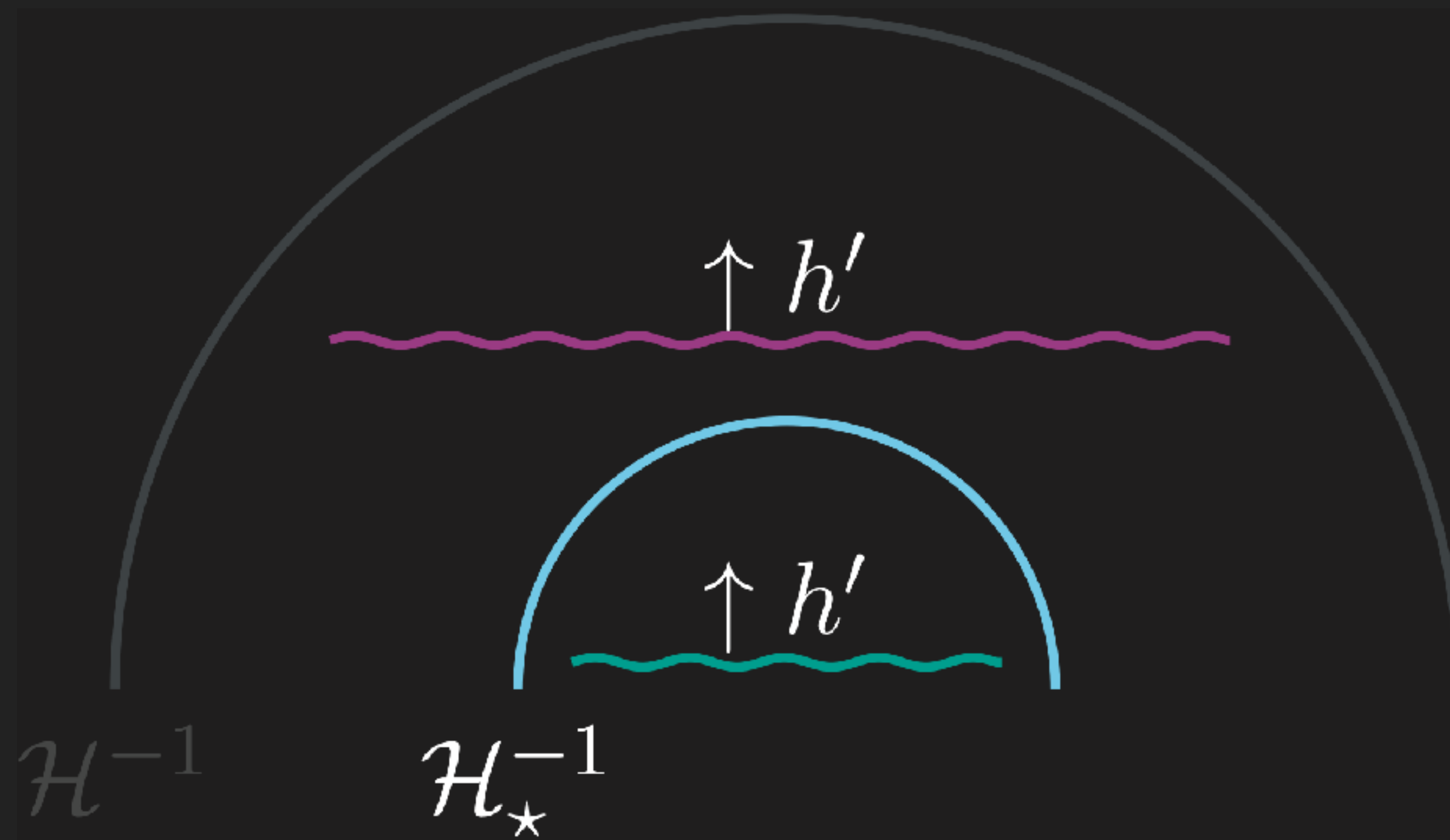
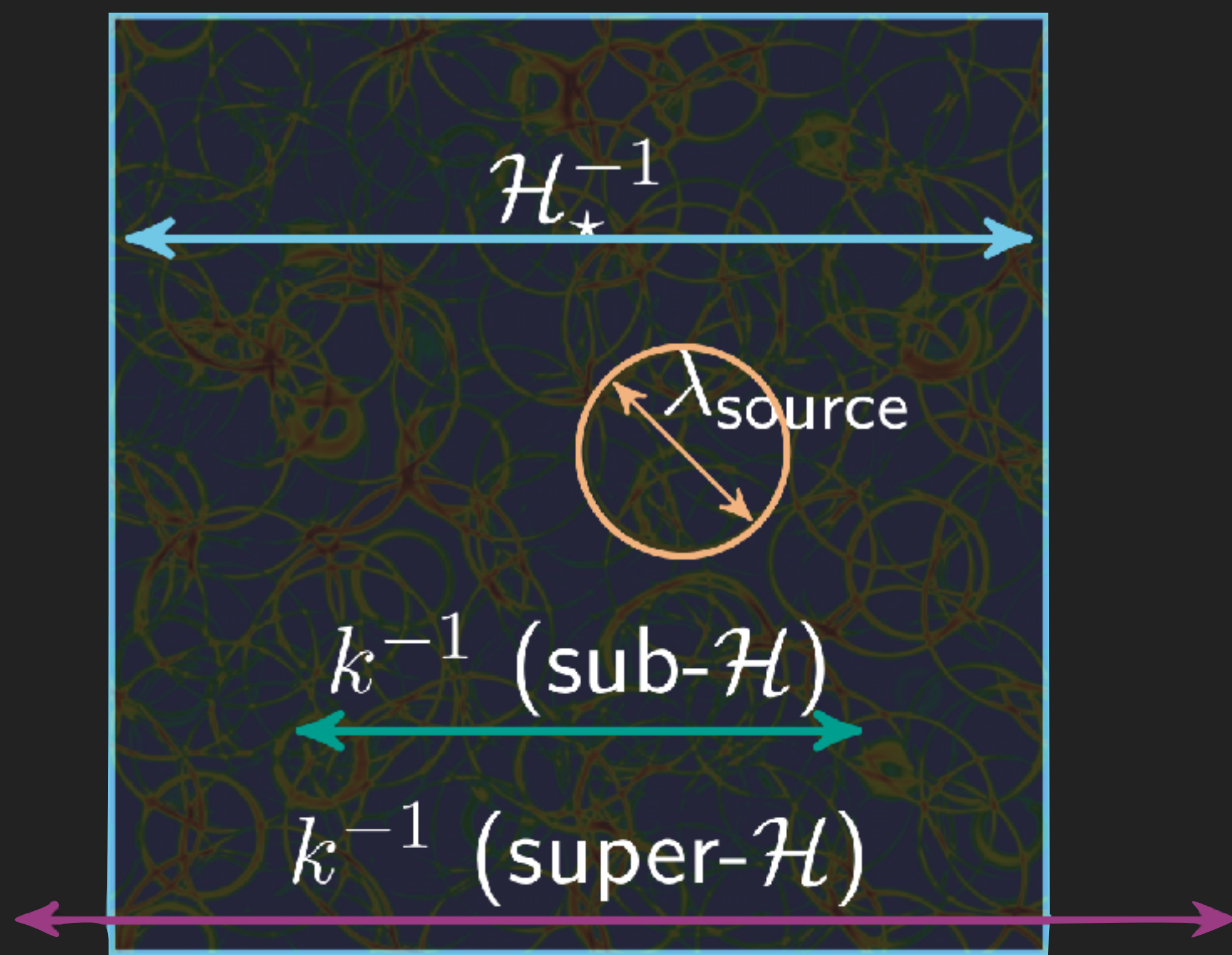
$$\frac{d\Omega_{\text{GW}}}{d \ln k} \sim \underbrace{k^3}_{\text{phase space}} \cdot \underbrace{k^2}_{\rho_{\text{GW}} \sim h'^2} \cdot \underbrace{\frac{1}{k^2}}_{\text{for RD}} \cdot \underbrace{P_{\Pi}(k)}_{k \text{ ind. from causality}} \sim k^3.$$





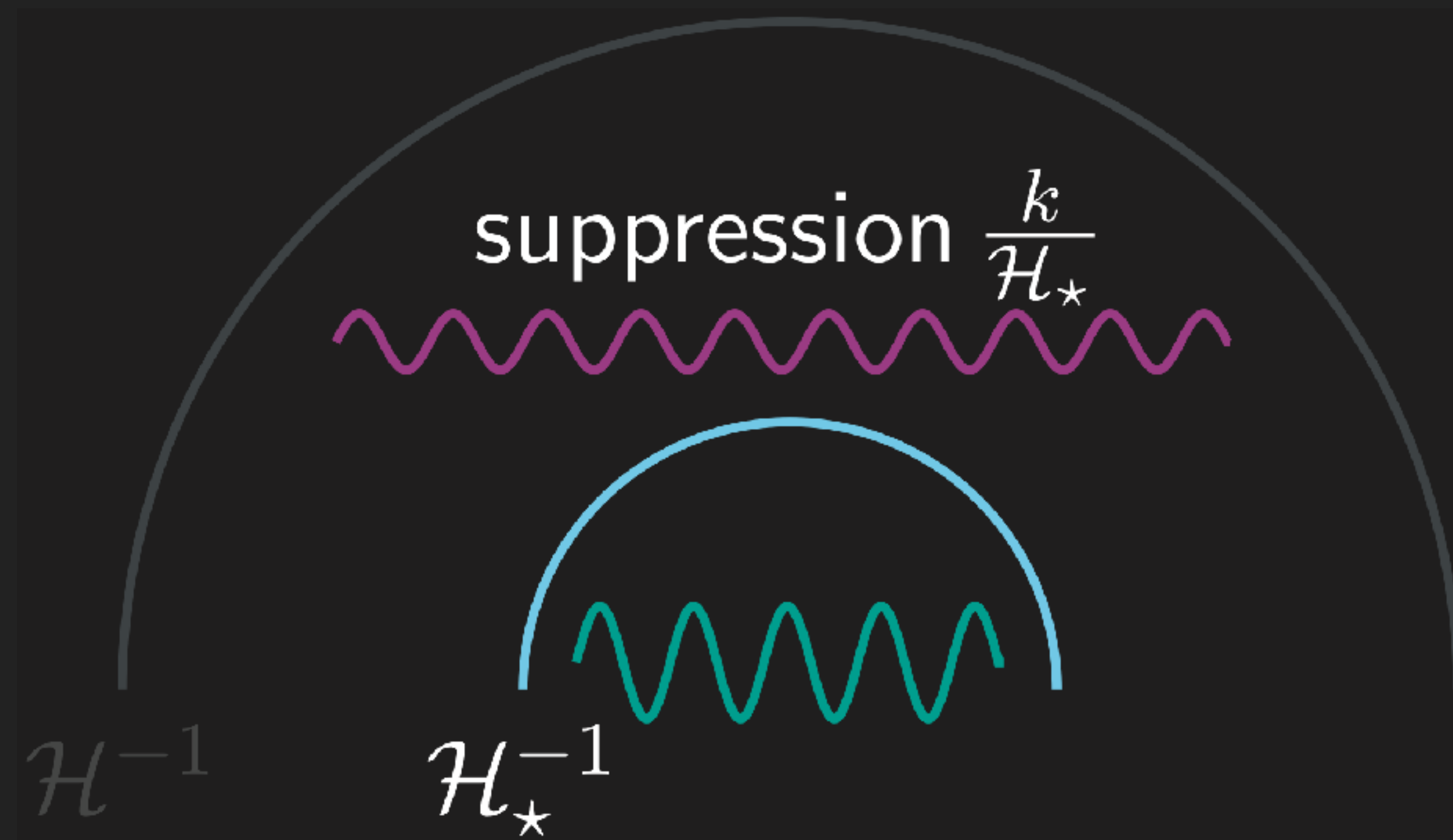
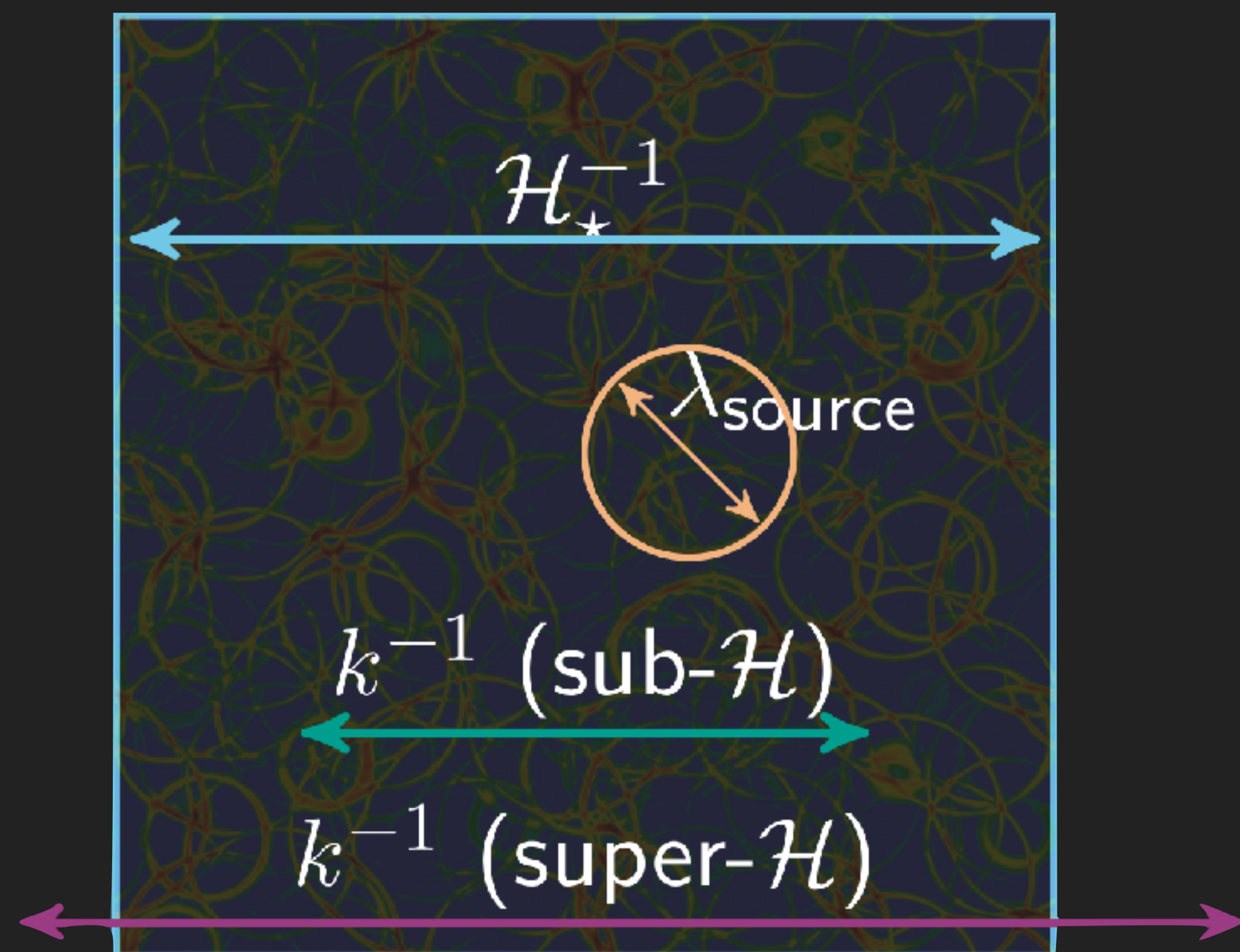


$$h'' + 2\mathcal{H}h' + k^2h = 0$$



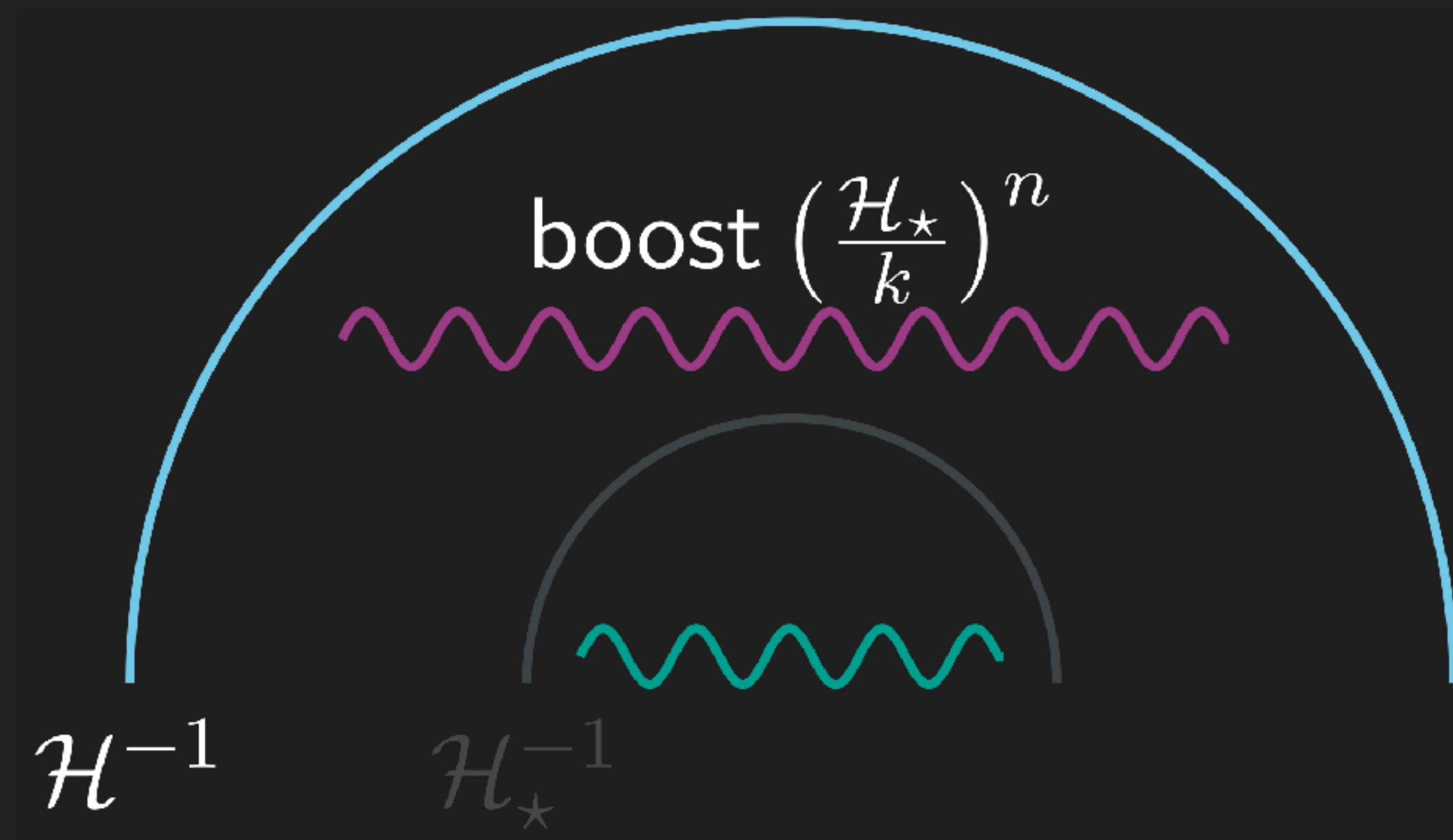
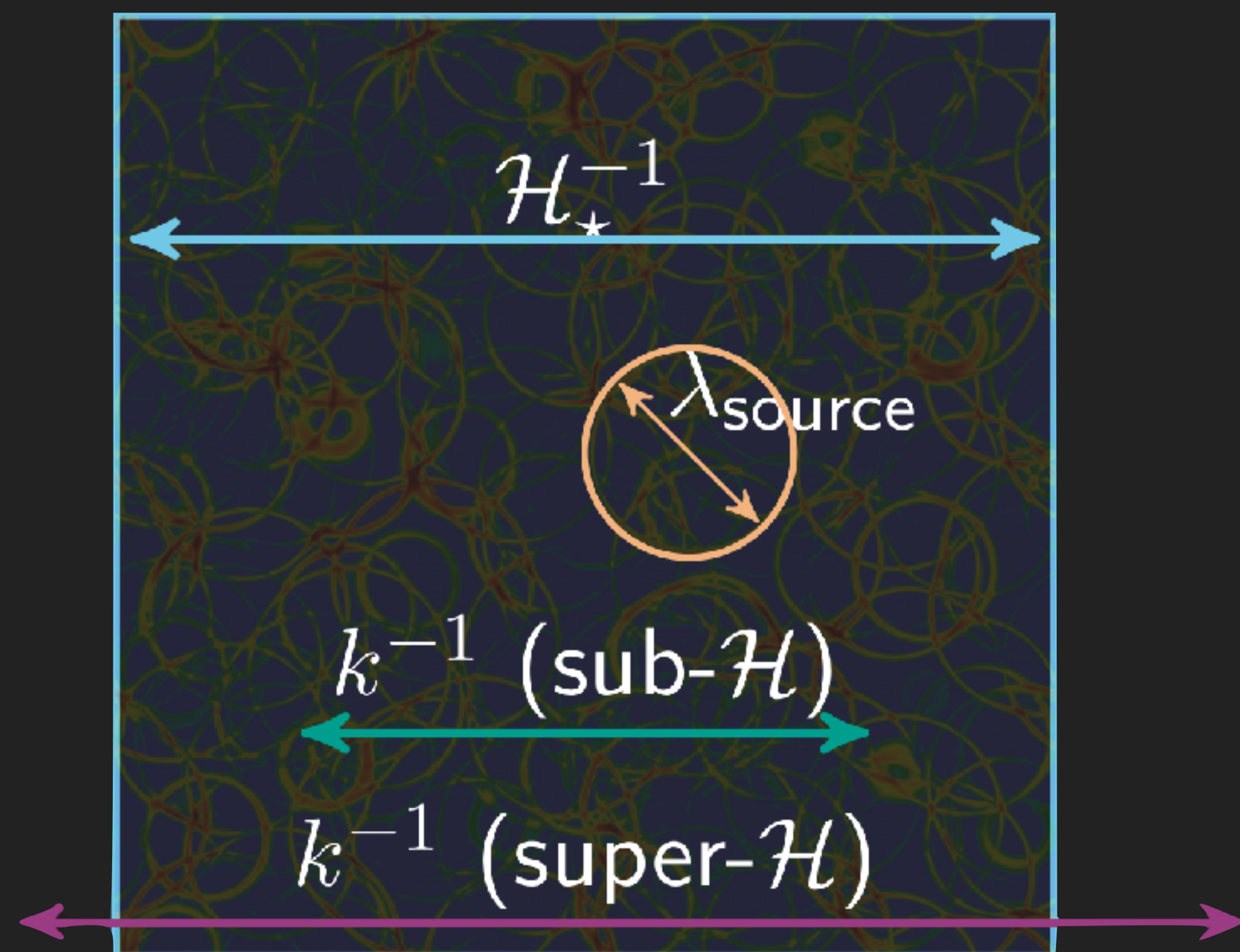


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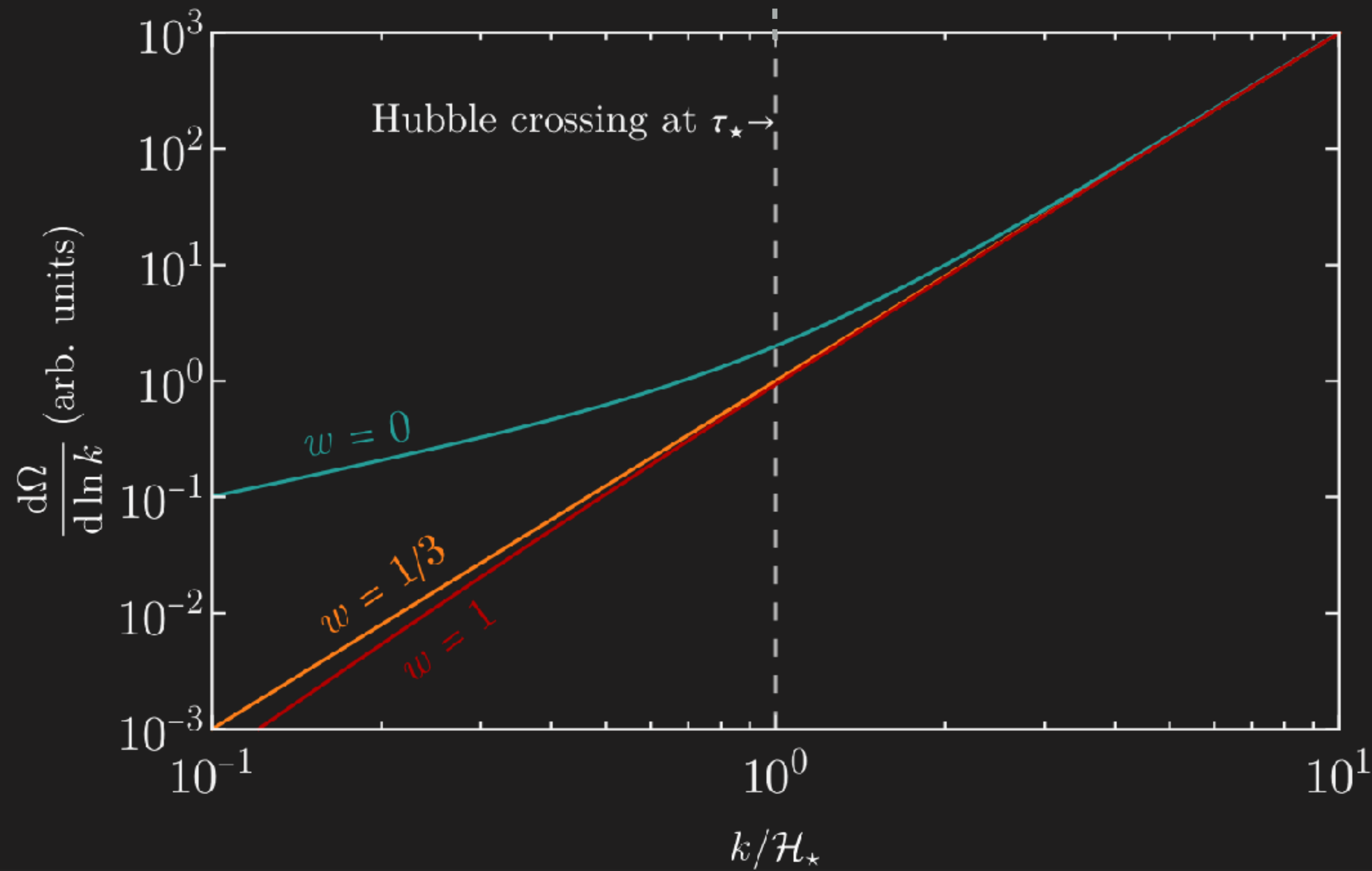
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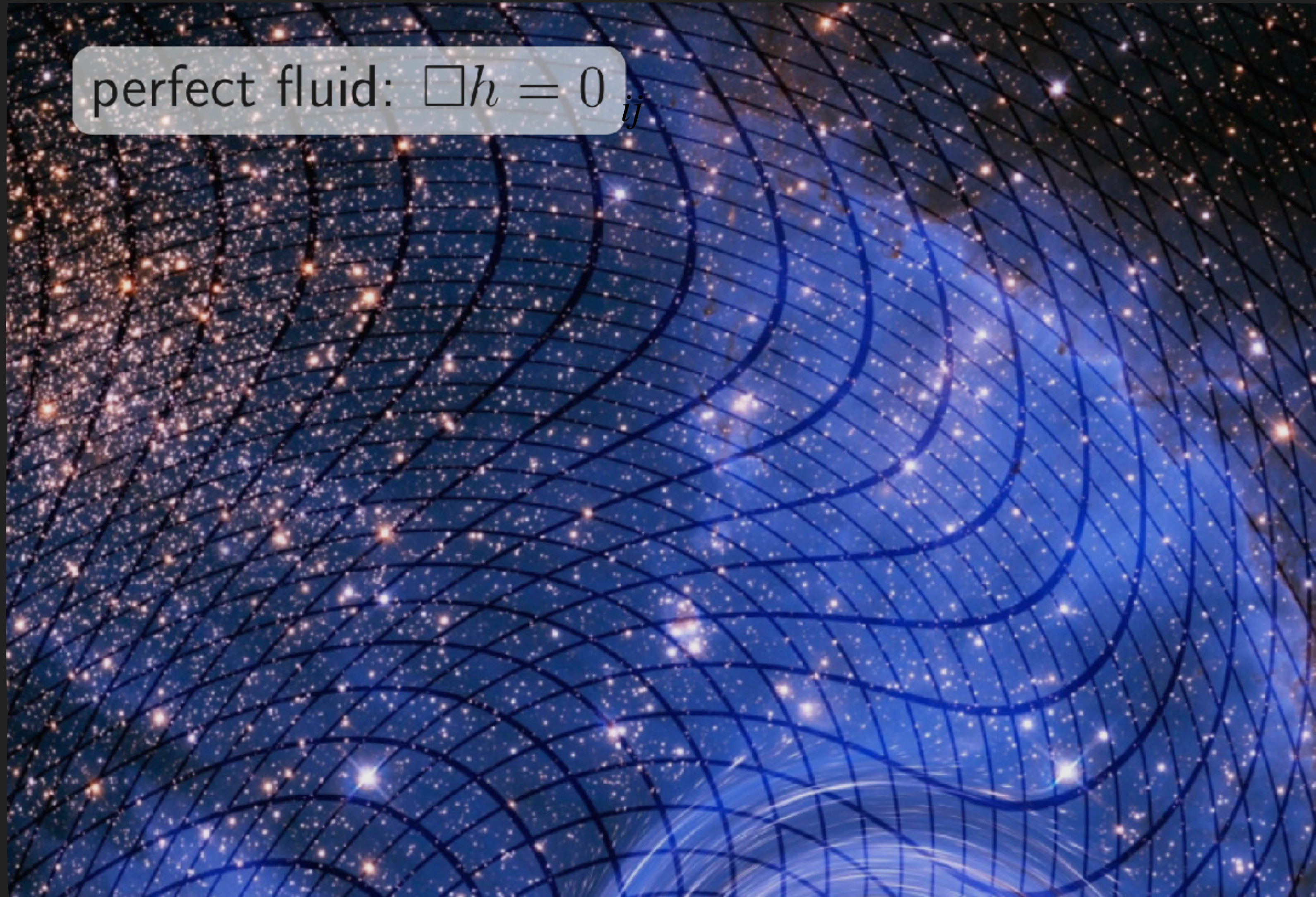
$$\Omega_{\text{GW}} \sim \begin{cases} k^3 & \text{RD} \\ k & \text{MD} \end{cases} \text{ (super-horizon)}$$

$$\Omega_{\text{GW}} \sim k^3 \text{ (sub-horizon)}$$



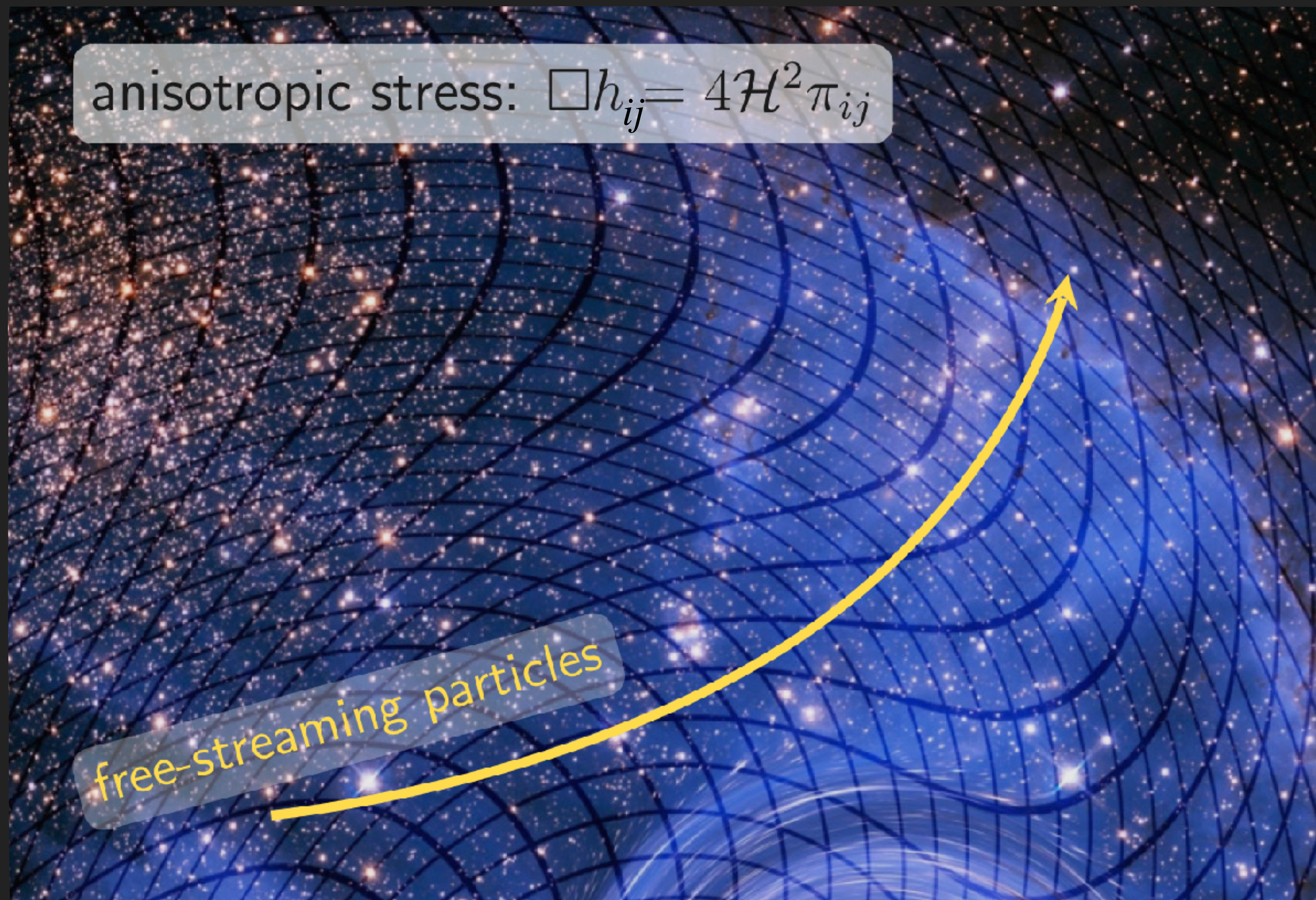
► Causality tail  $\Leftrightarrow$  probe equation-of-state  $w$





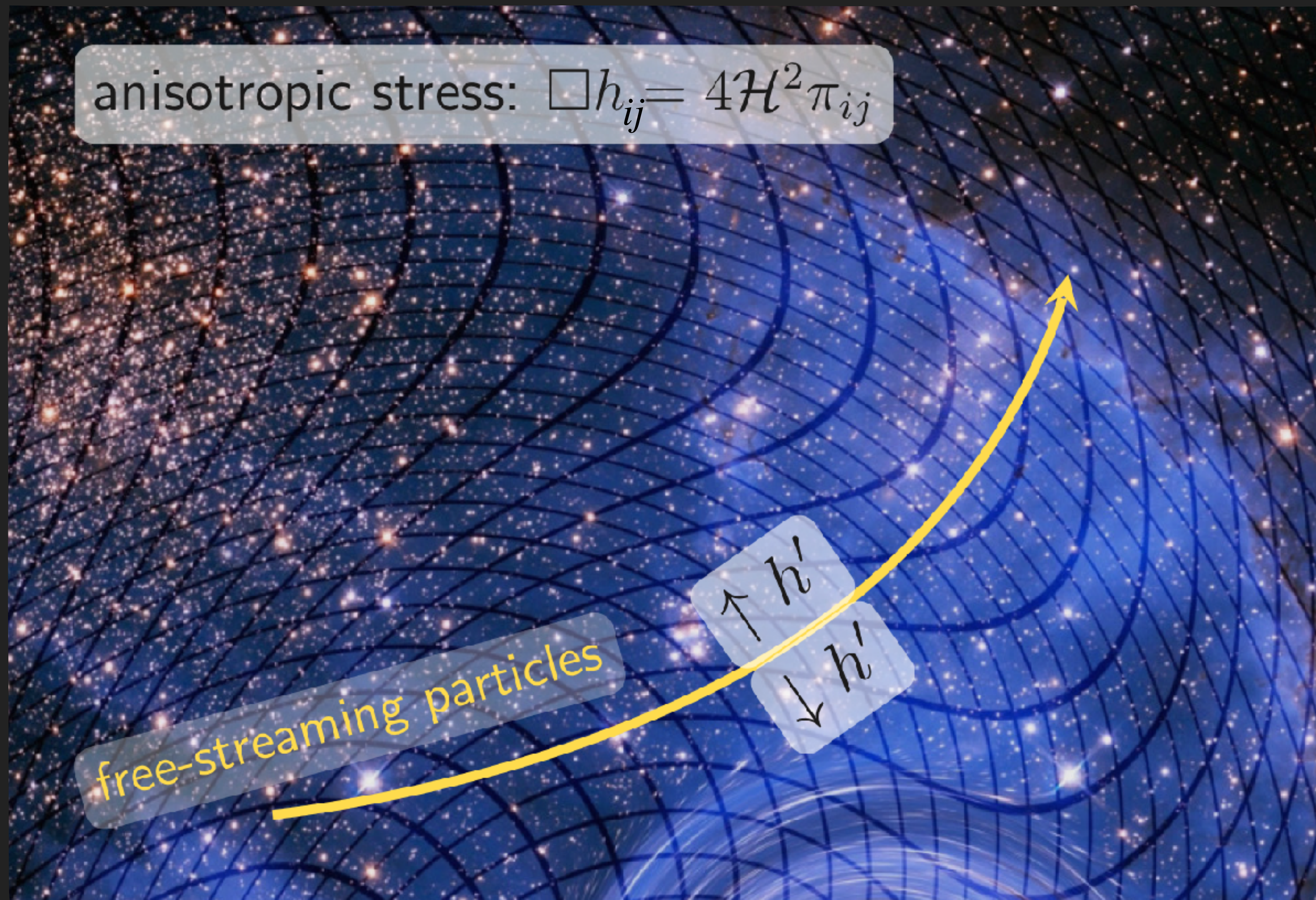
[ '86 Bond]  
[ '04 Weinberg]  
[ '06 Watanabe,  
Komatsu]





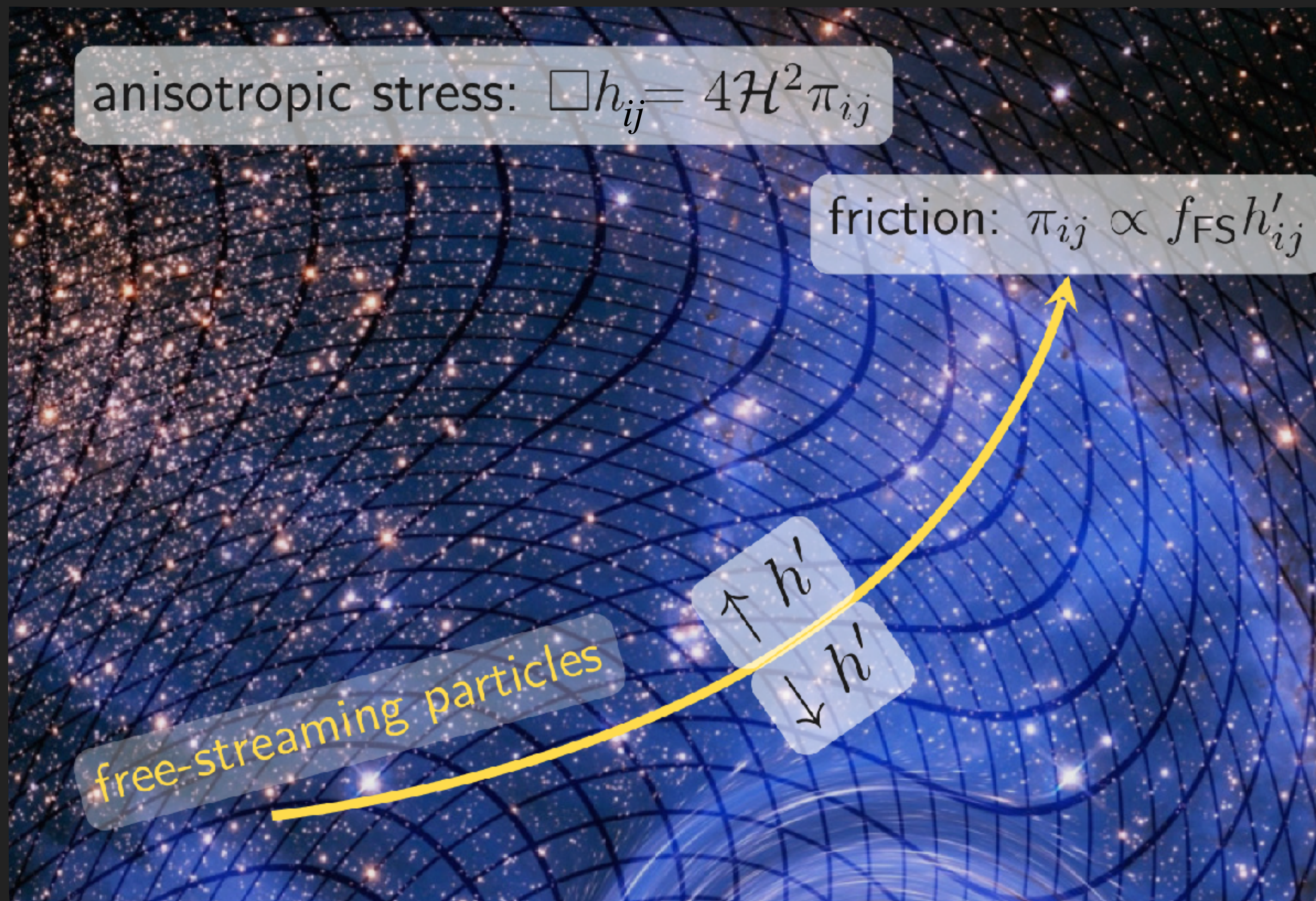
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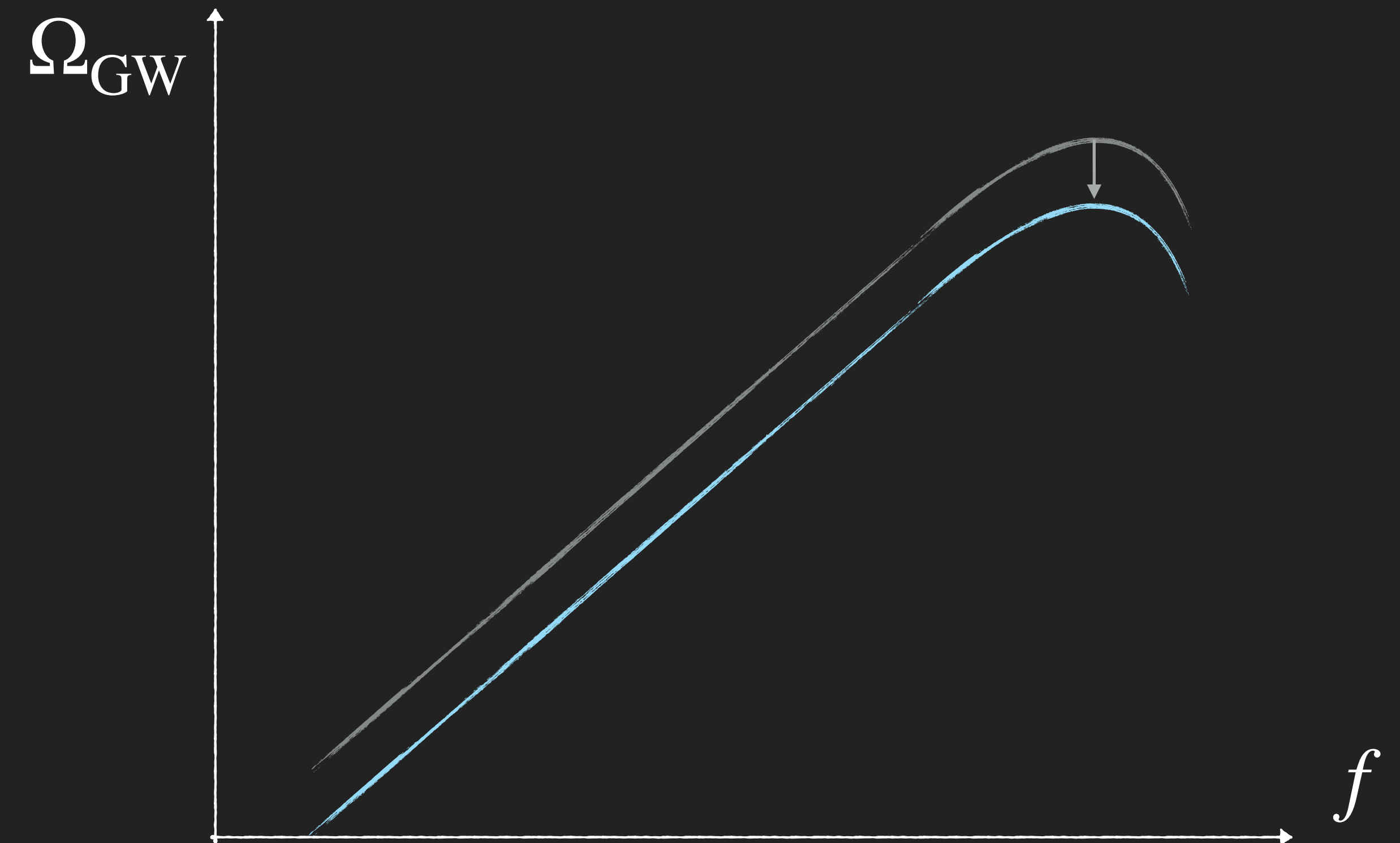




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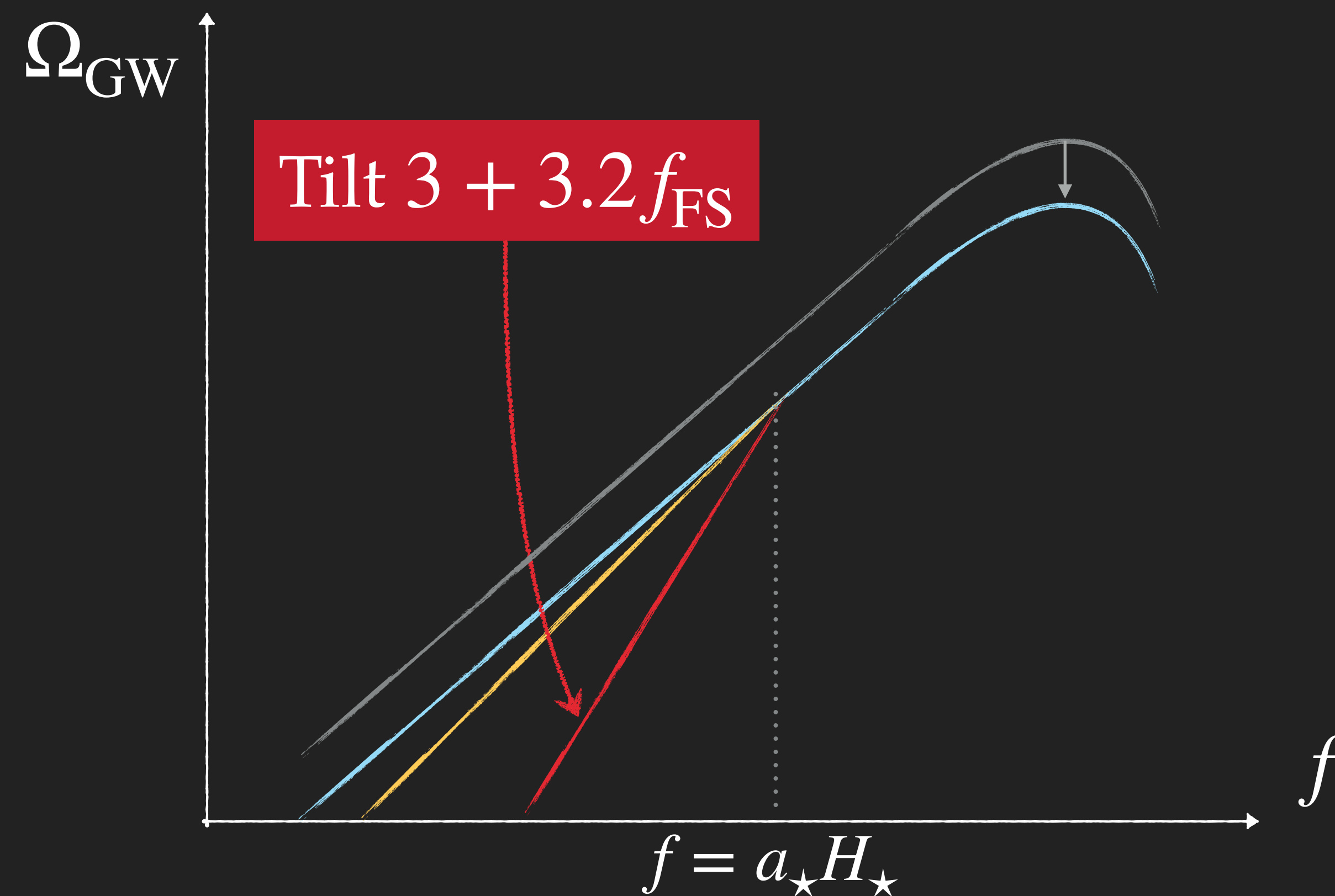
- ▶ Weinberg damping:  $h' \neq 0$  at horizon crossing
- ▶ Frequency-independent [‘04 Weinberg]





- ▶ Weinberg damping:  $h' \neq 0$  at horizon crossing
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- ▶ Causality-limited sources: **also** affected at GW production
- ▶ Frequency-dependent

[‘20 Hook, Marques-Tavares, DR]

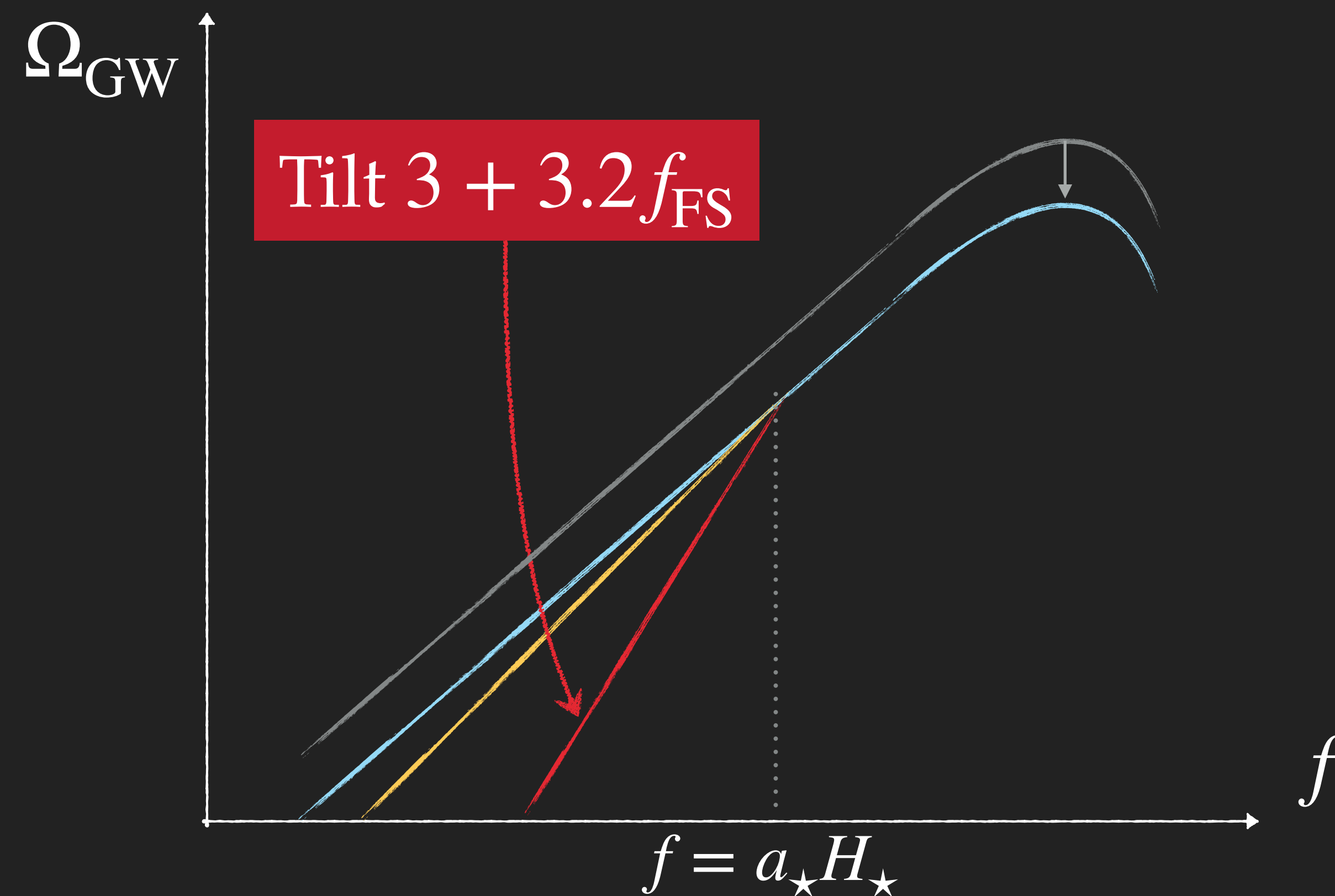




- ▶ Weinberg damping:  $h' \neq 0$  at horizon crossing
- ▶ Frequency-independent [‘04 Weinberg]
- ▶ Causality-limited sources: **also** affected at GW production
- ▶ Frequency-dependent

[‘20 Hook, Marques-Tavares, DR]

- ▶ Causality tail  $\Leftrightarrow$  probe new free-streaming particles before BBN





- ▶ SM particles are massless  $\Rightarrow w = \frac{1}{3}$

---

$\longleftarrow T$   
time  $\longrightarrow$



- ▶ SM particles are massless  $\Rightarrow w = \frac{1}{3}$
- ▶ Heavy particle decays  $\Rightarrow \Omega_{\text{GW}} \sim g_{\star}(T)^{1/3}$

 $v_{\text{EW}}$ 

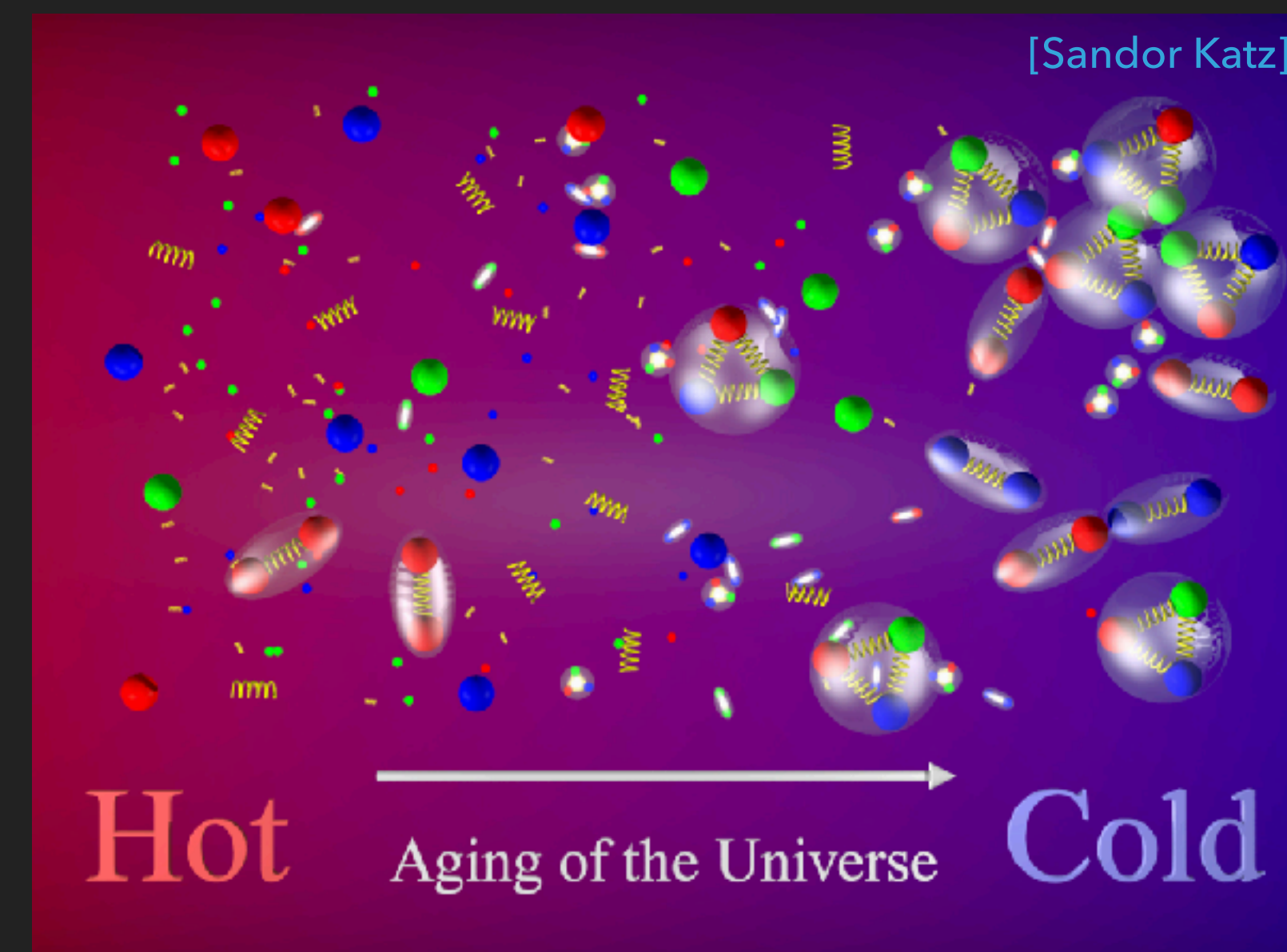
$\leftarrow T$   
time  $\rightarrow$



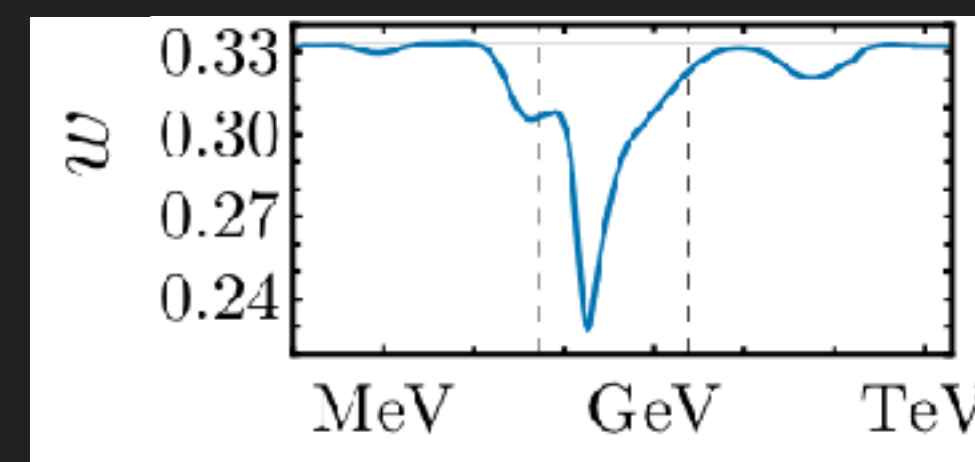
- ▶ SM particles are massless  $\Rightarrow w = \frac{1}{3}$
- ▶ Heavy particle decays  $\Rightarrow \Omega_{\text{GW}} \sim g_{\star}(T)^{1/3}$
- ▶ QCD confines
- ▶  $w$  deviates from  $\frac{1}{3}$

 $\nu_{\text{EW}}$  $\Lambda_{\text{QCD}}$ 

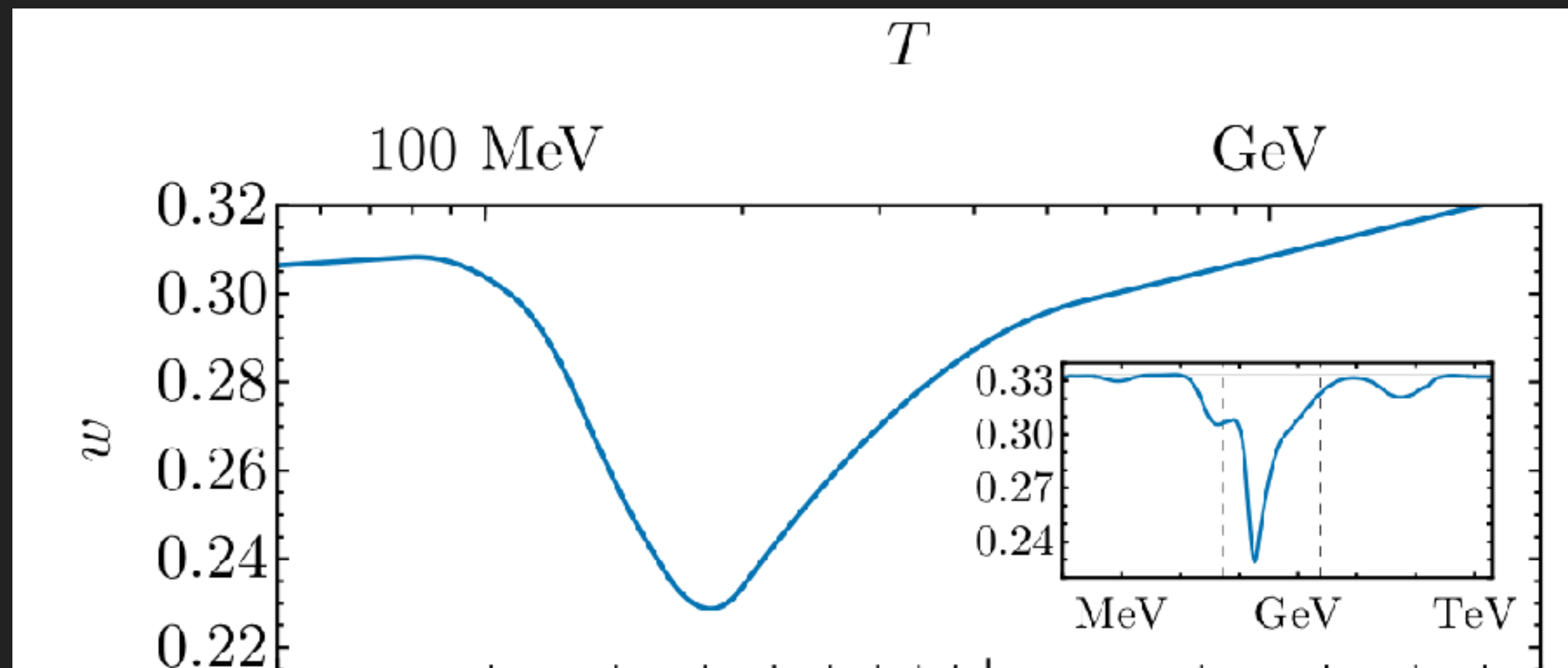
$\leftarrow T$   
time  $\rightarrow$



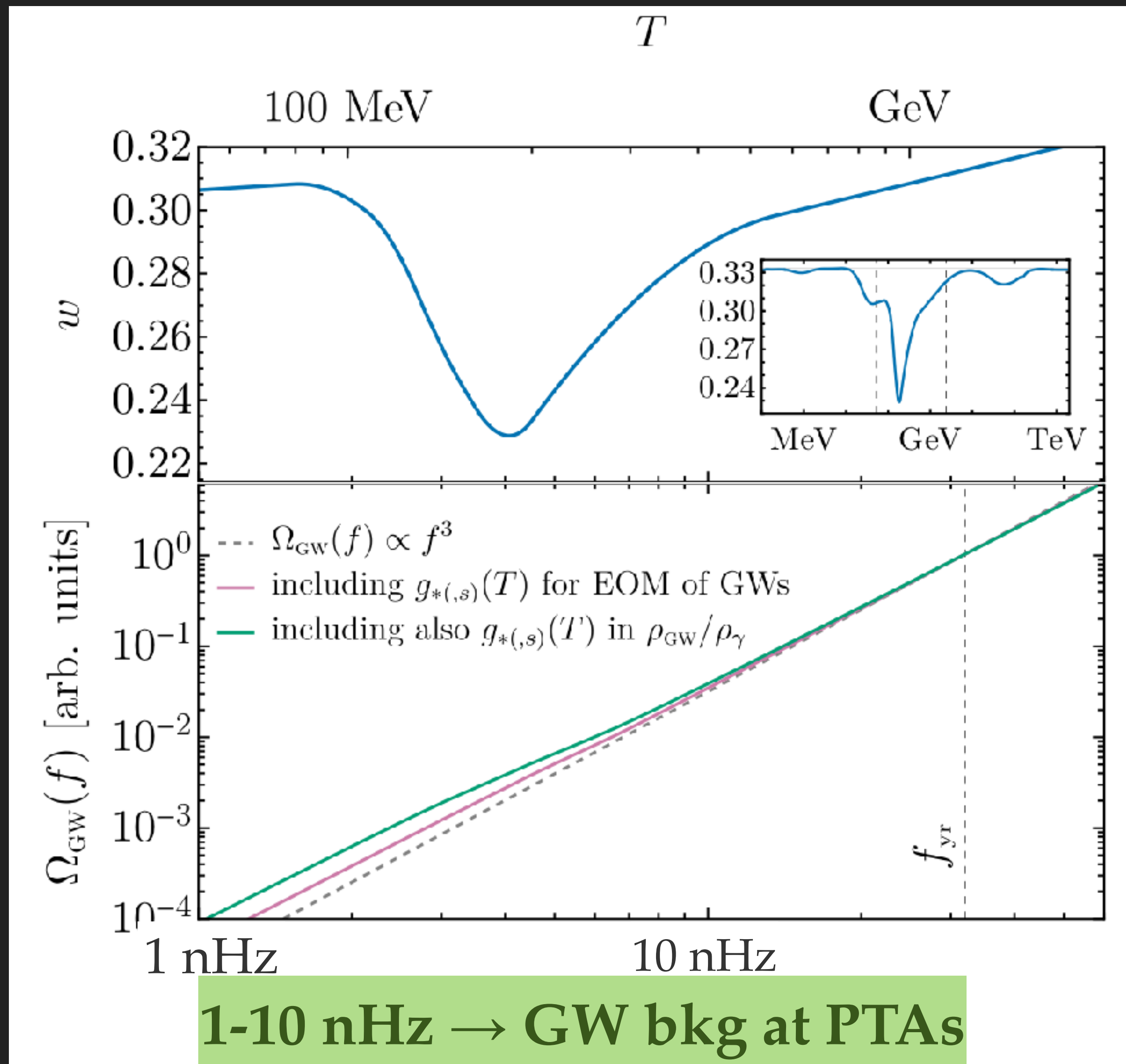




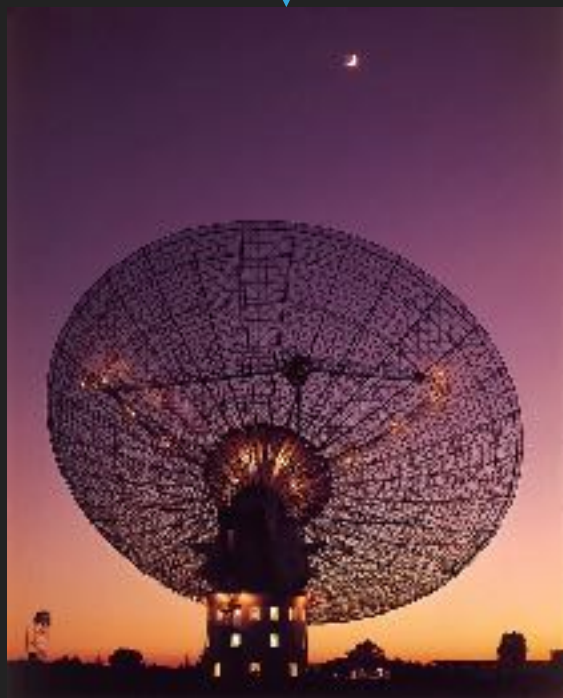
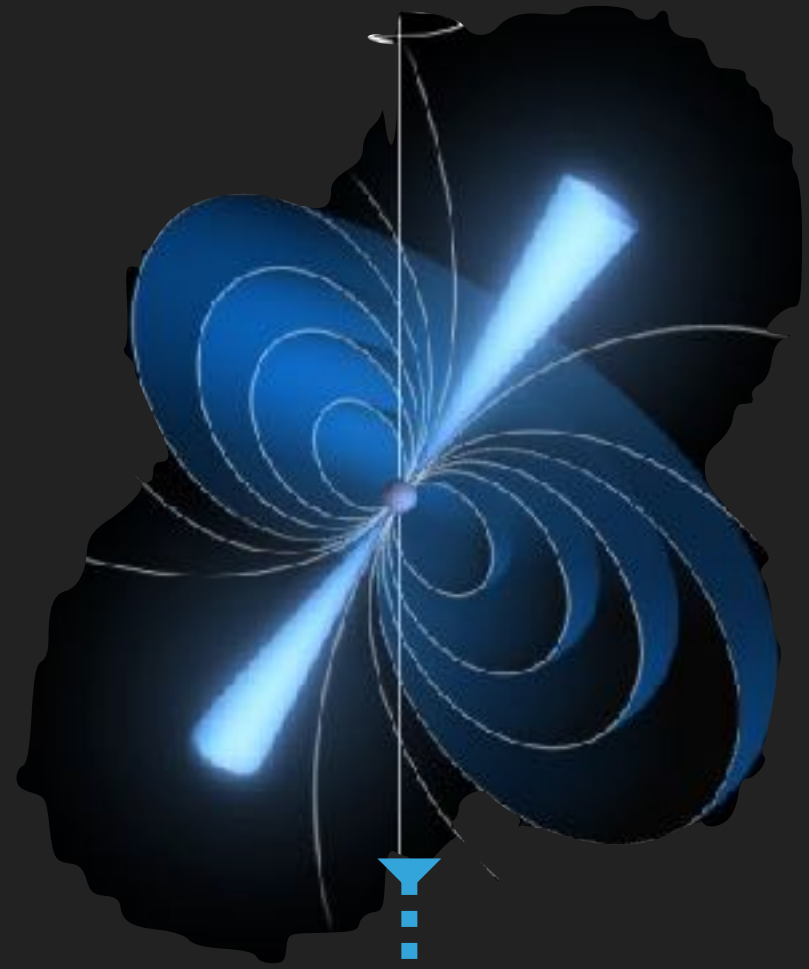








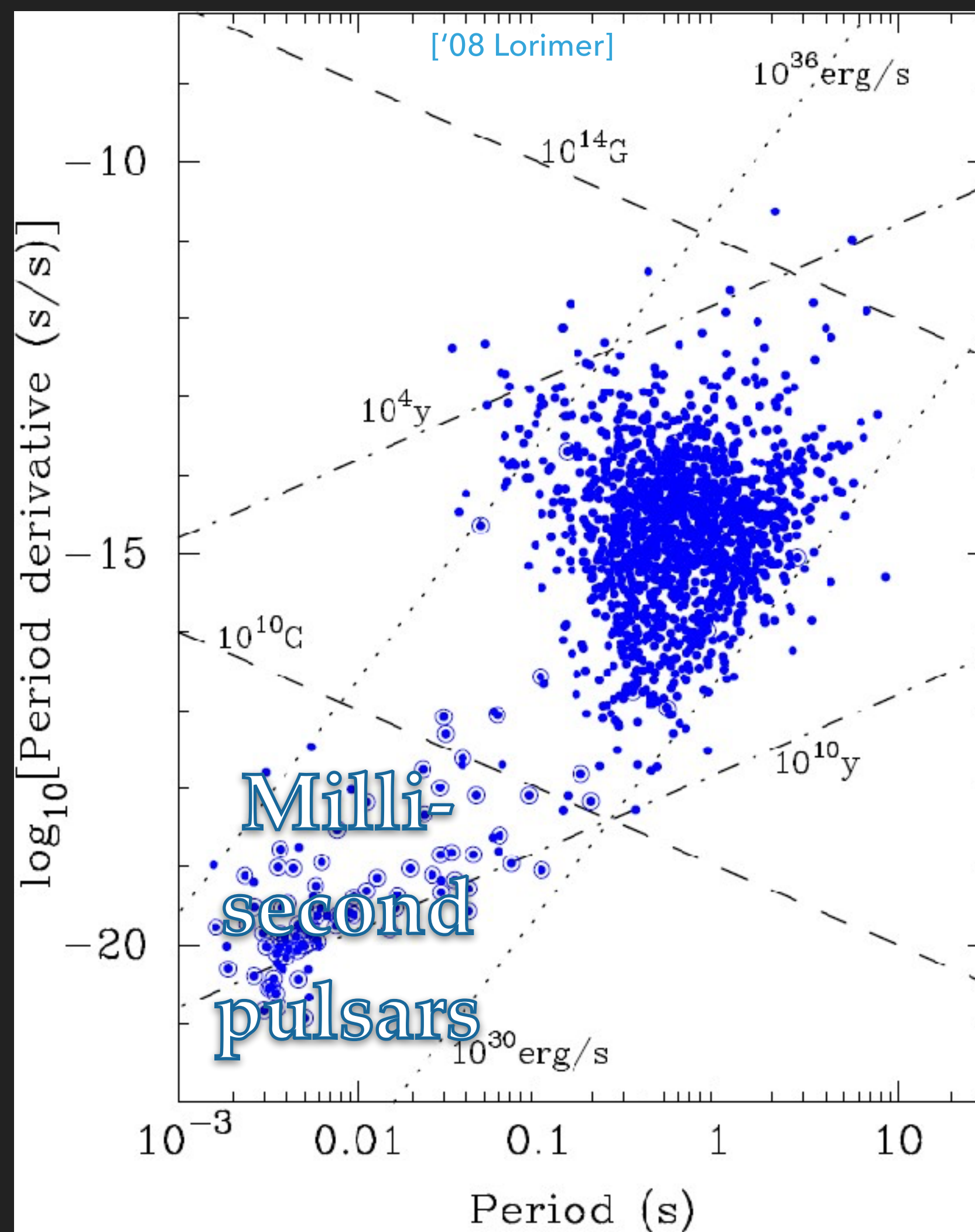
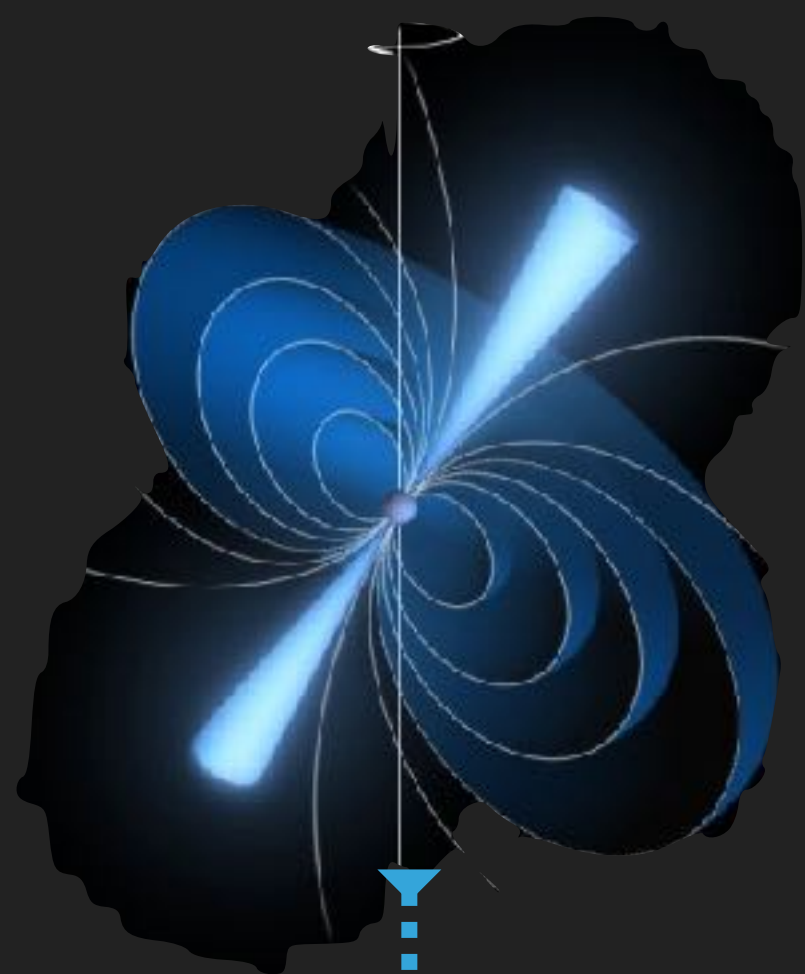






# PULSAR TIMING ARRAY MEASUREMENTS

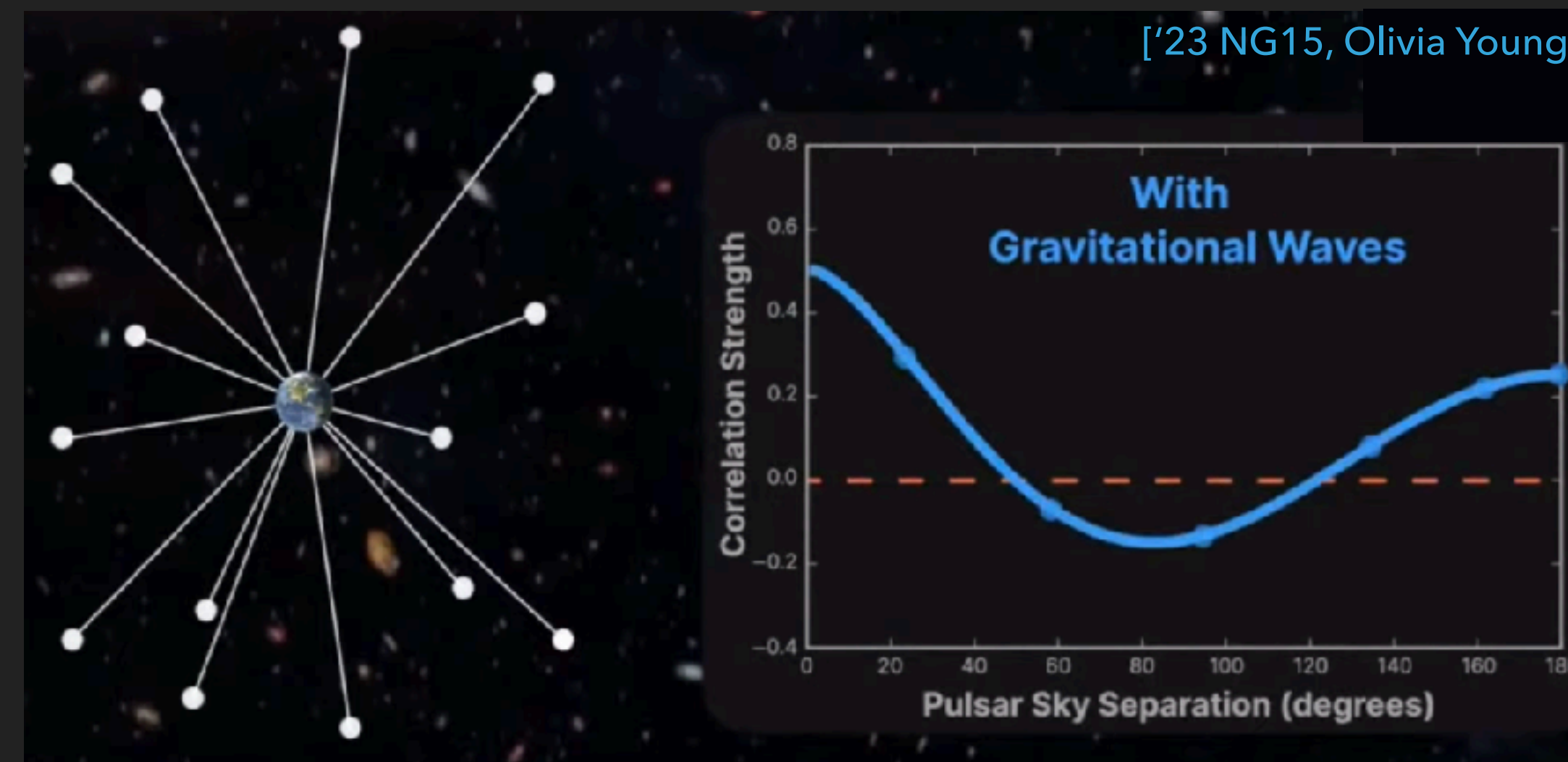
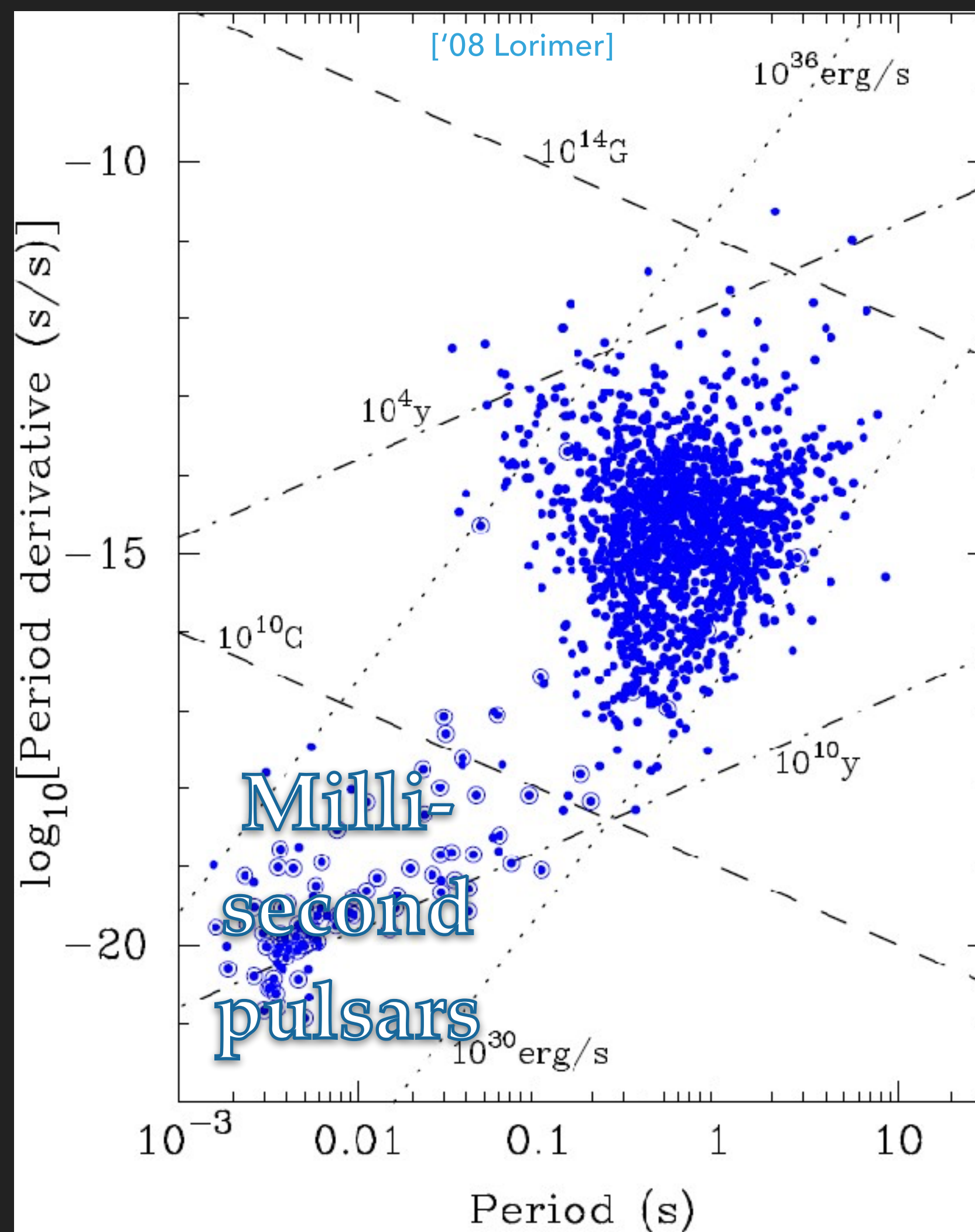
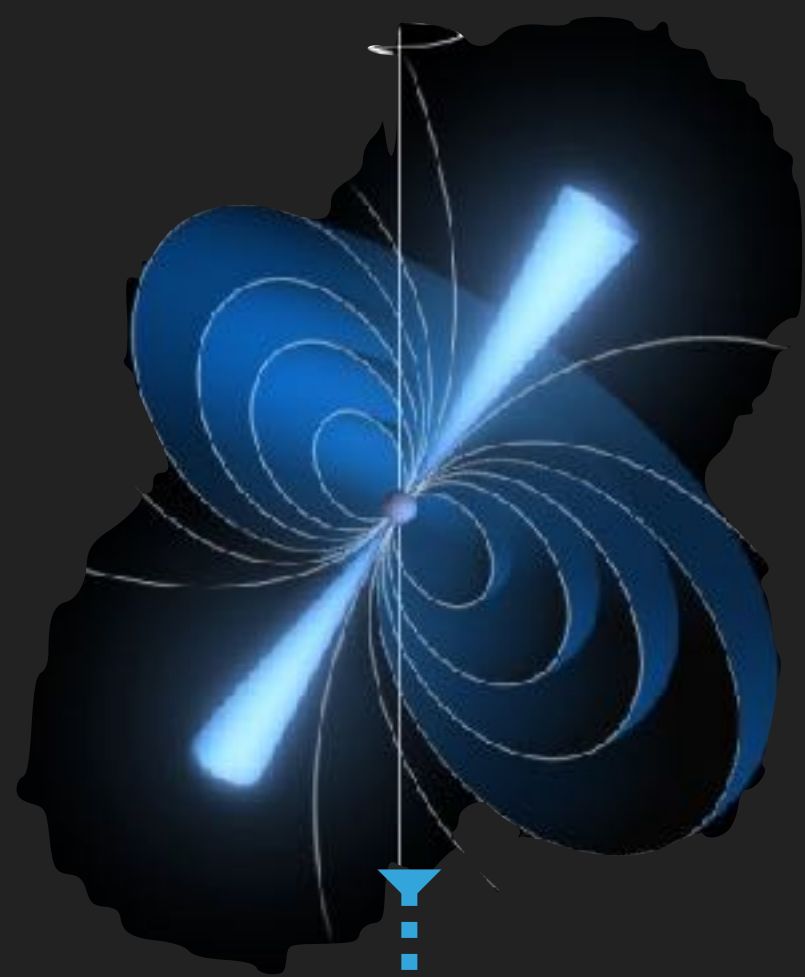
[ '78 Sazhin; '79 Detweiler;  
'83 Hellings, Downs] 14





# PULSAR TIMING ARRAY MEASUREMENTS

[ '78 Sazhin; '79 Detweiler; '83 Hellings, Downs] 14

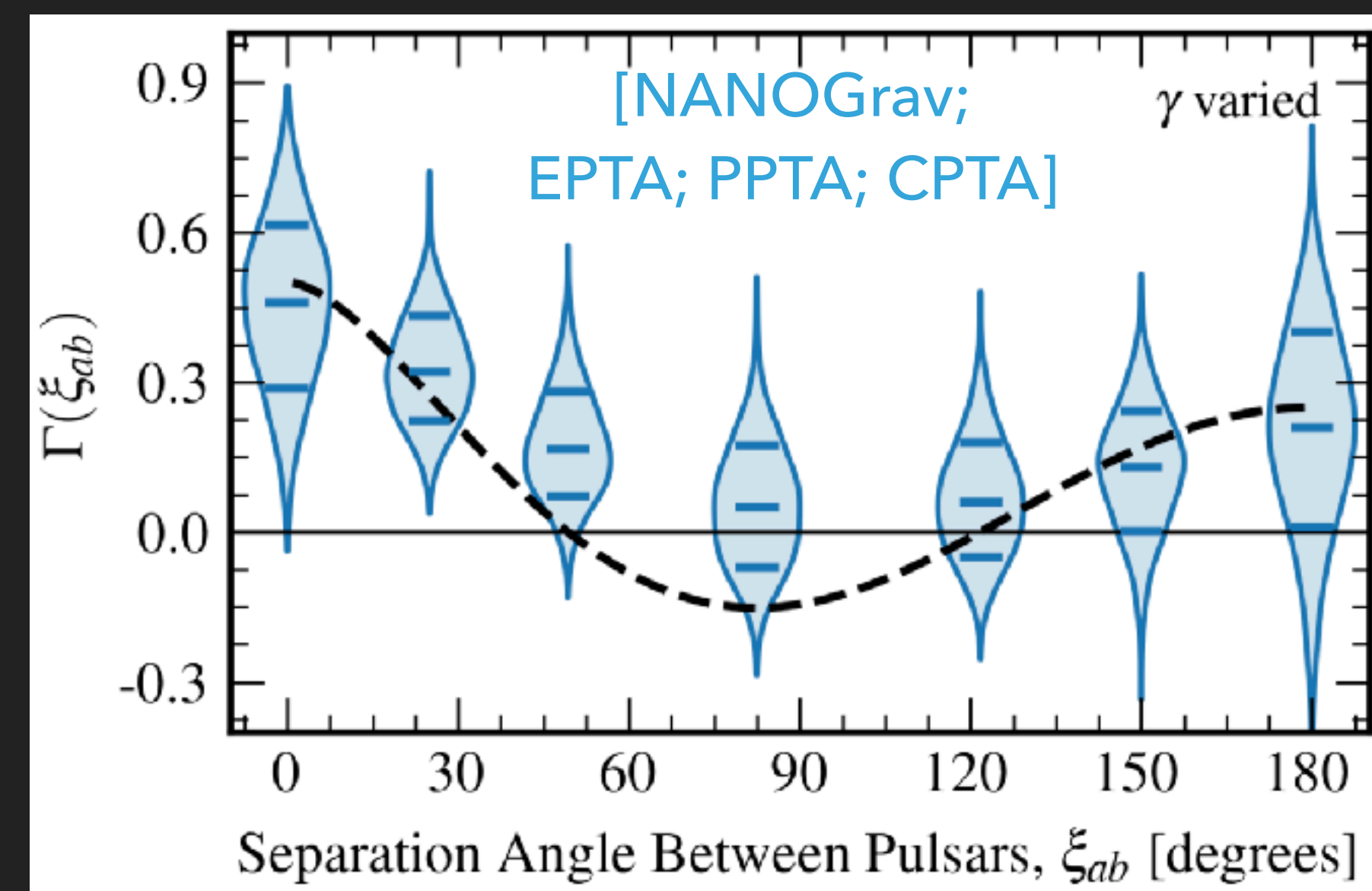
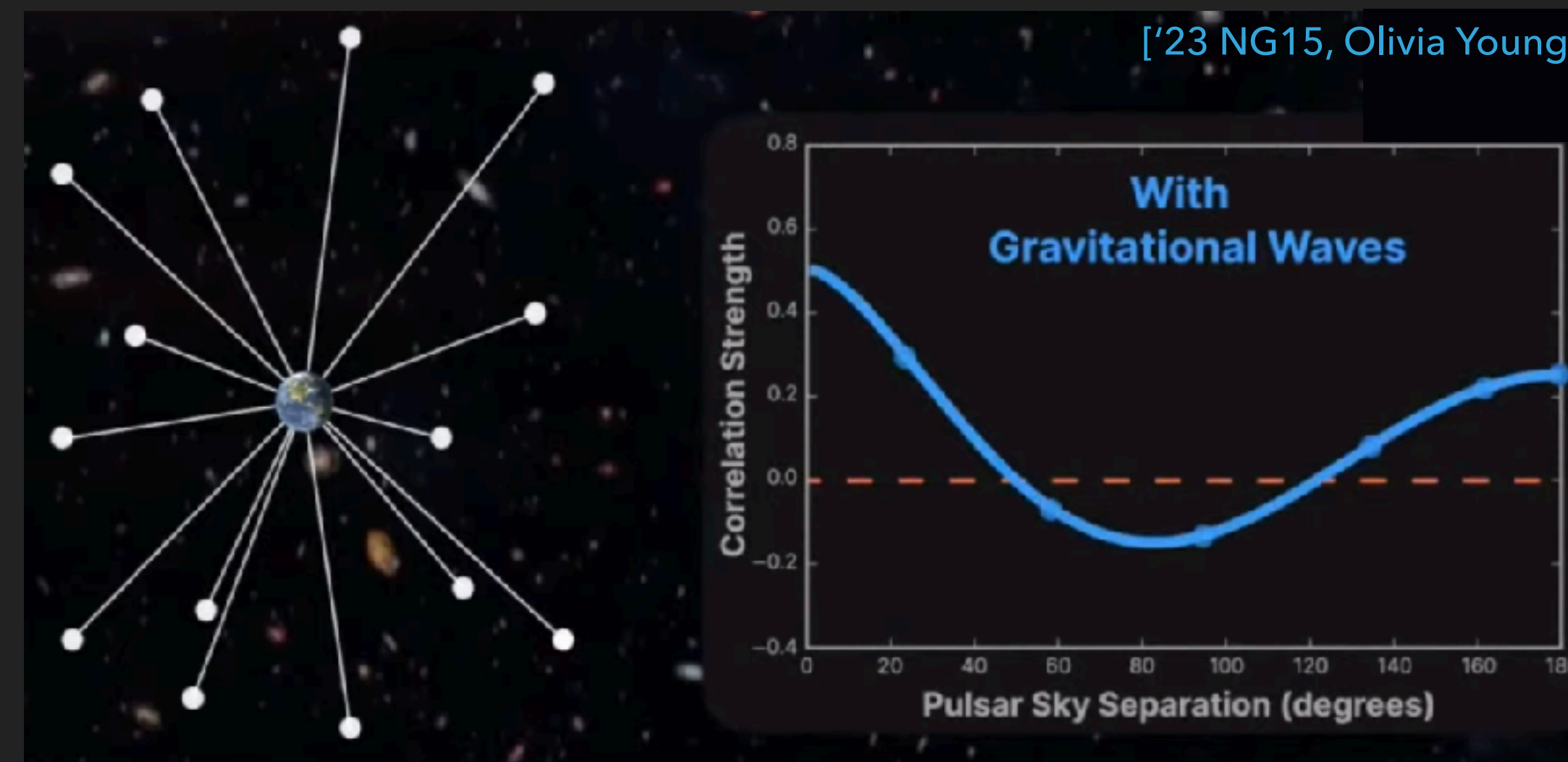
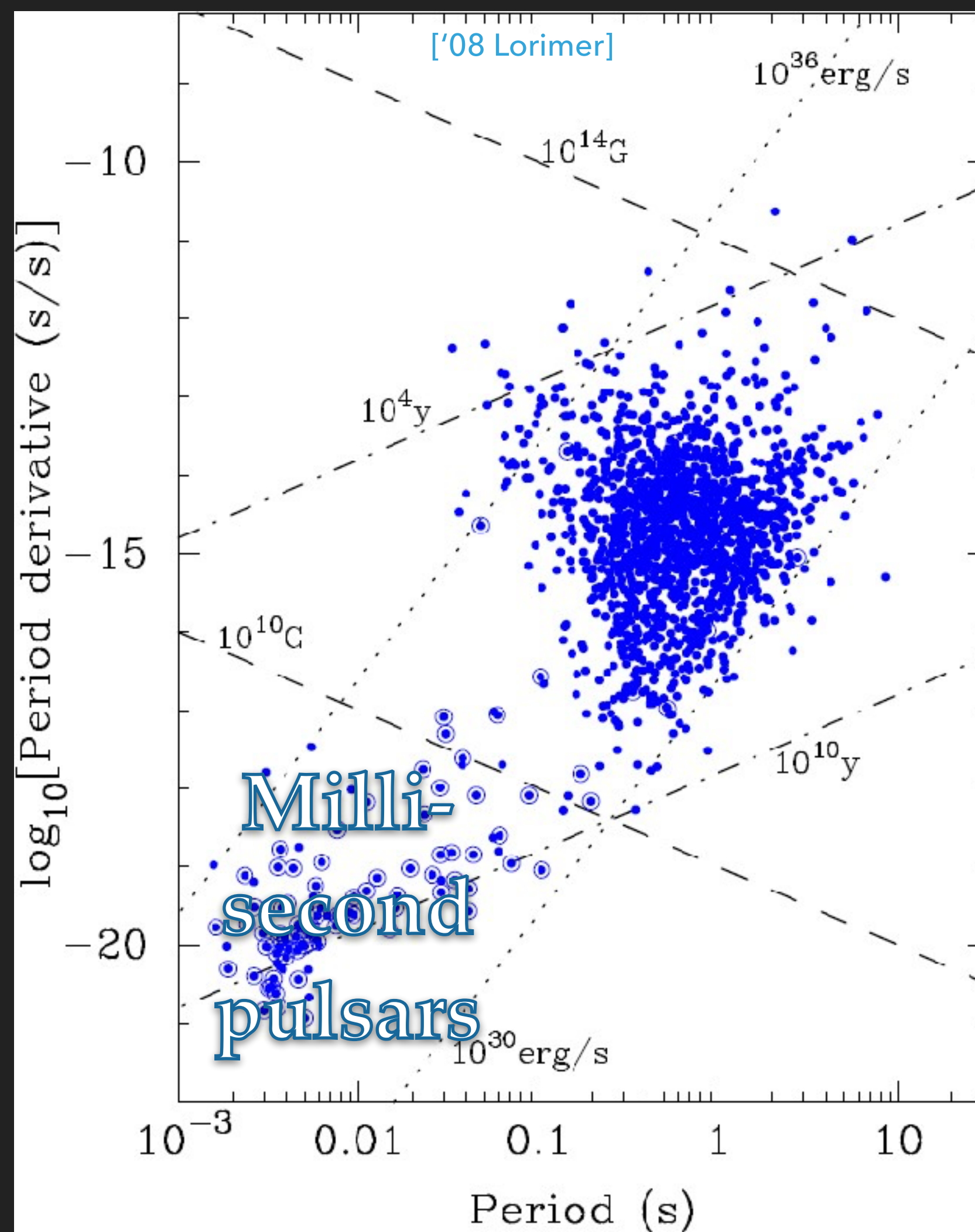
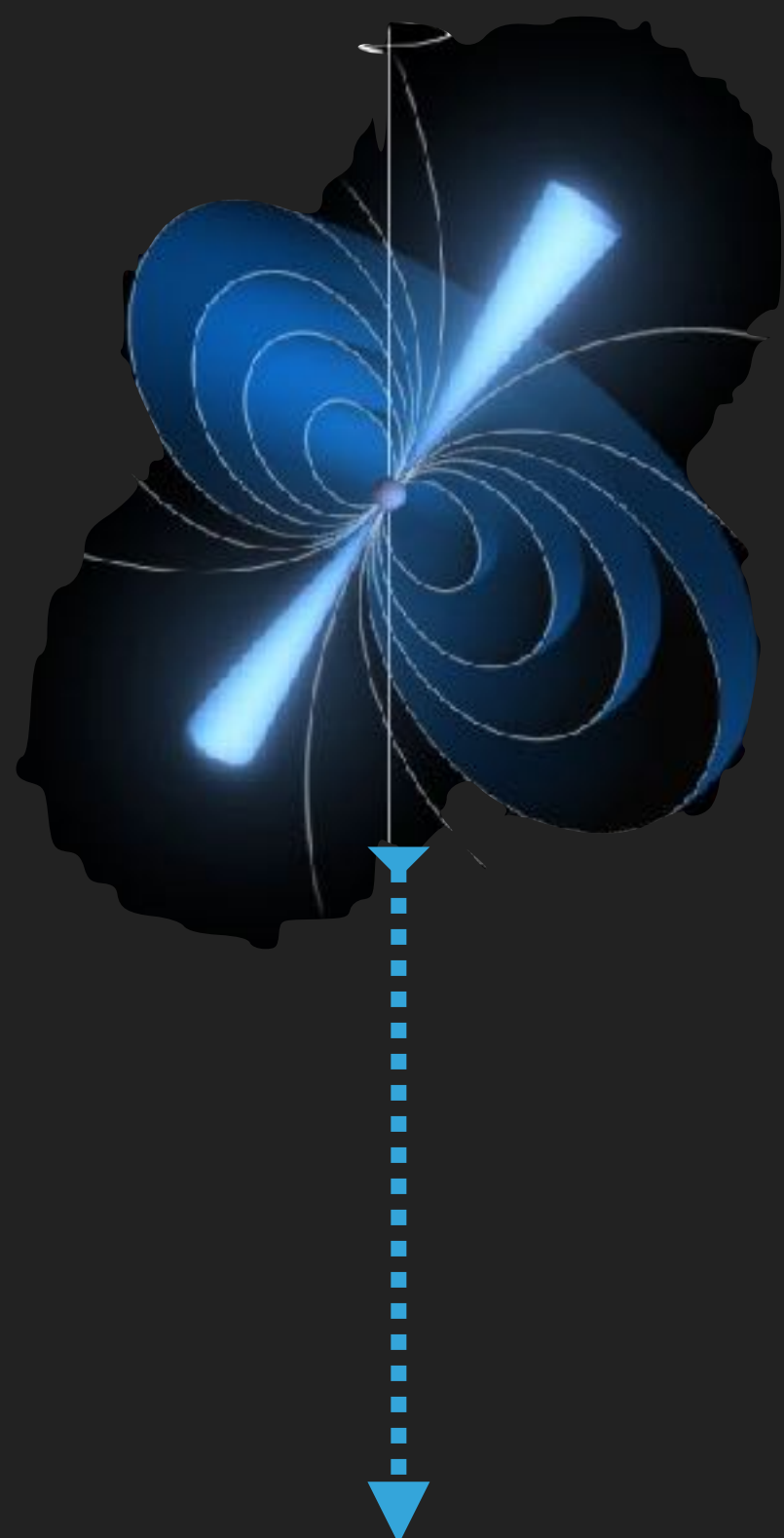


[ '23 NG15, Olivia Young]

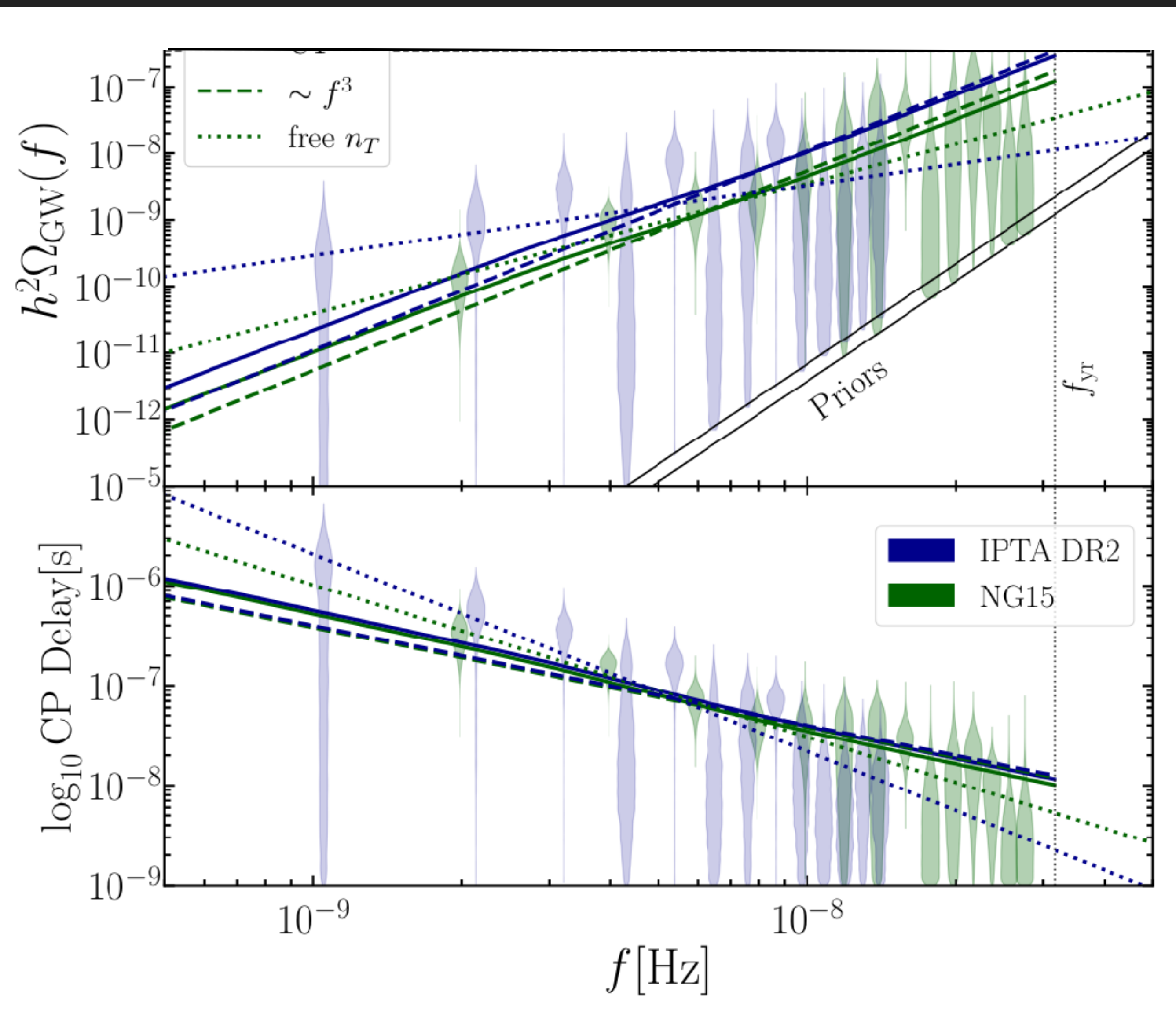


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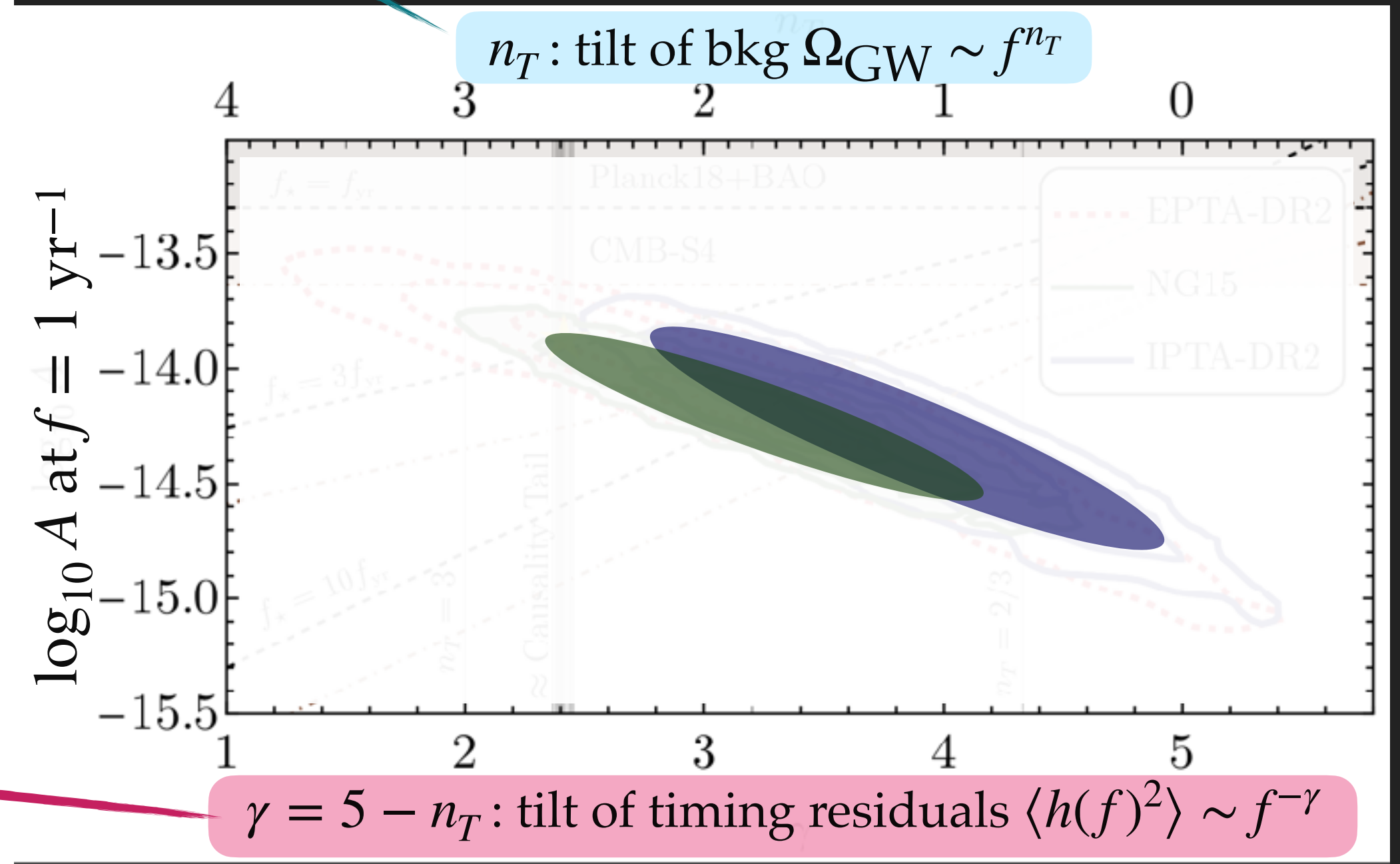
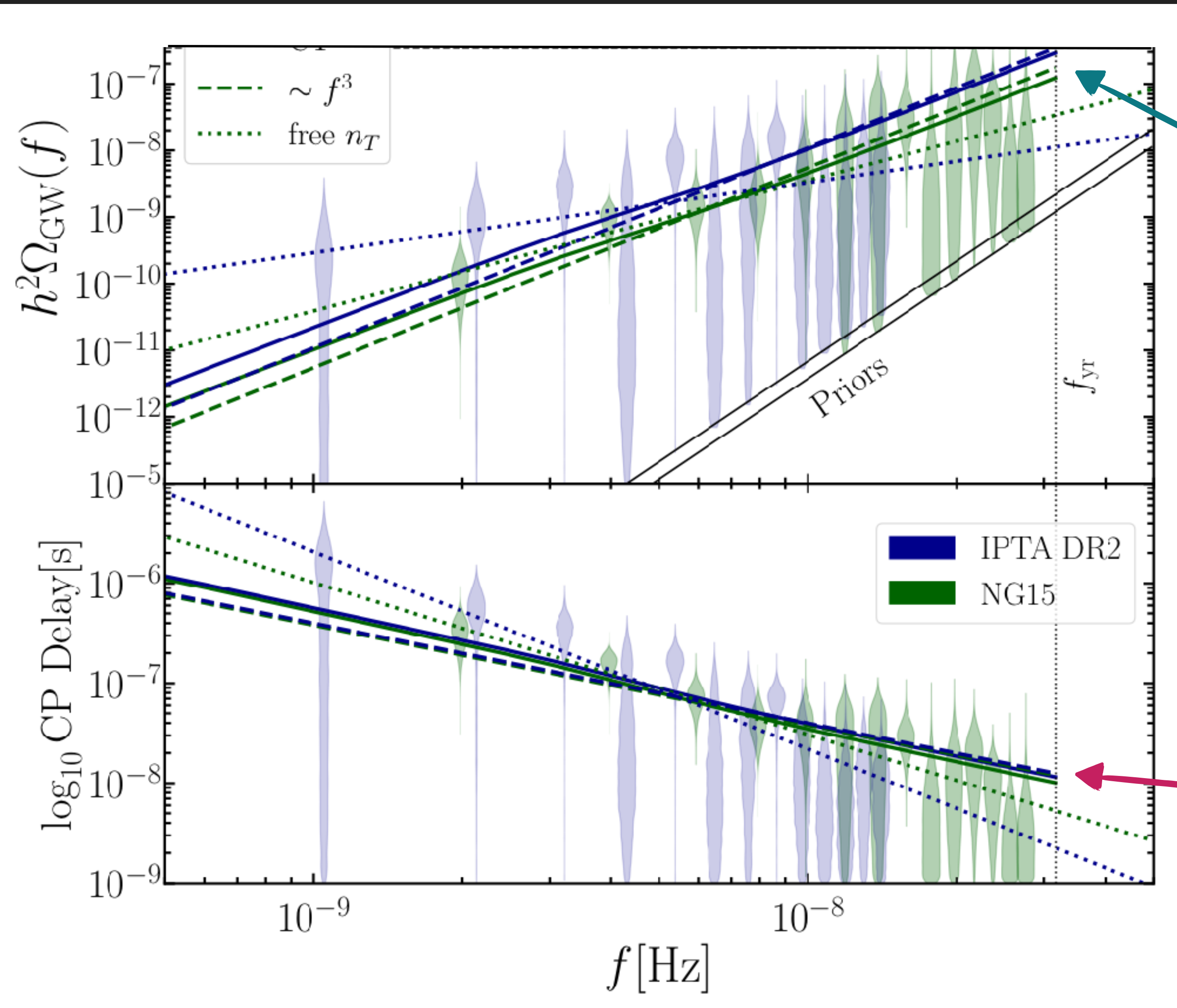
[ '78 Sazhin; '79 Detweiler; '83 Hellings, Downs] 14









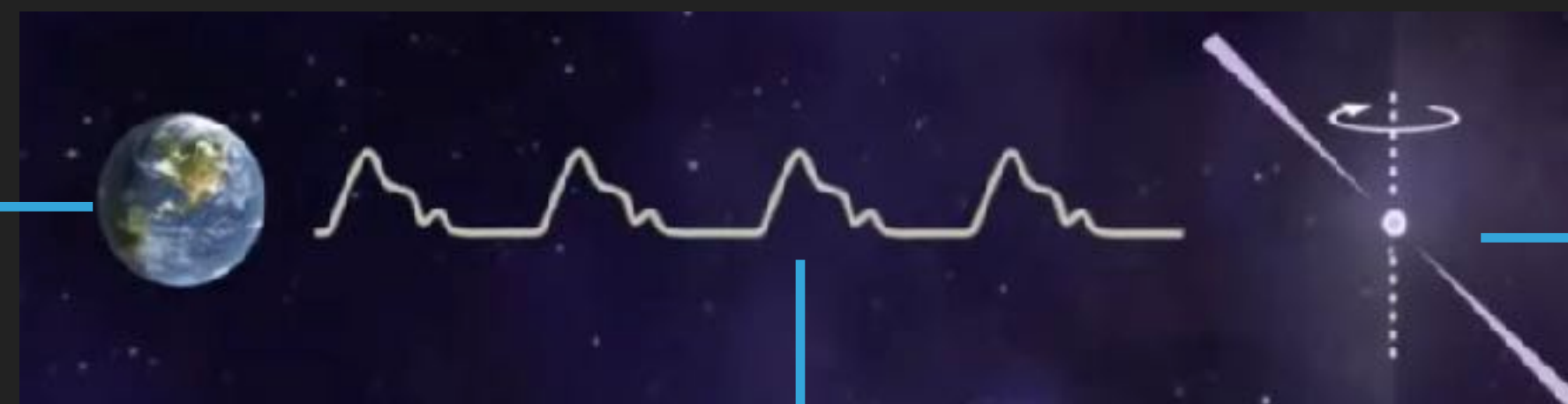








▶ Earth motion

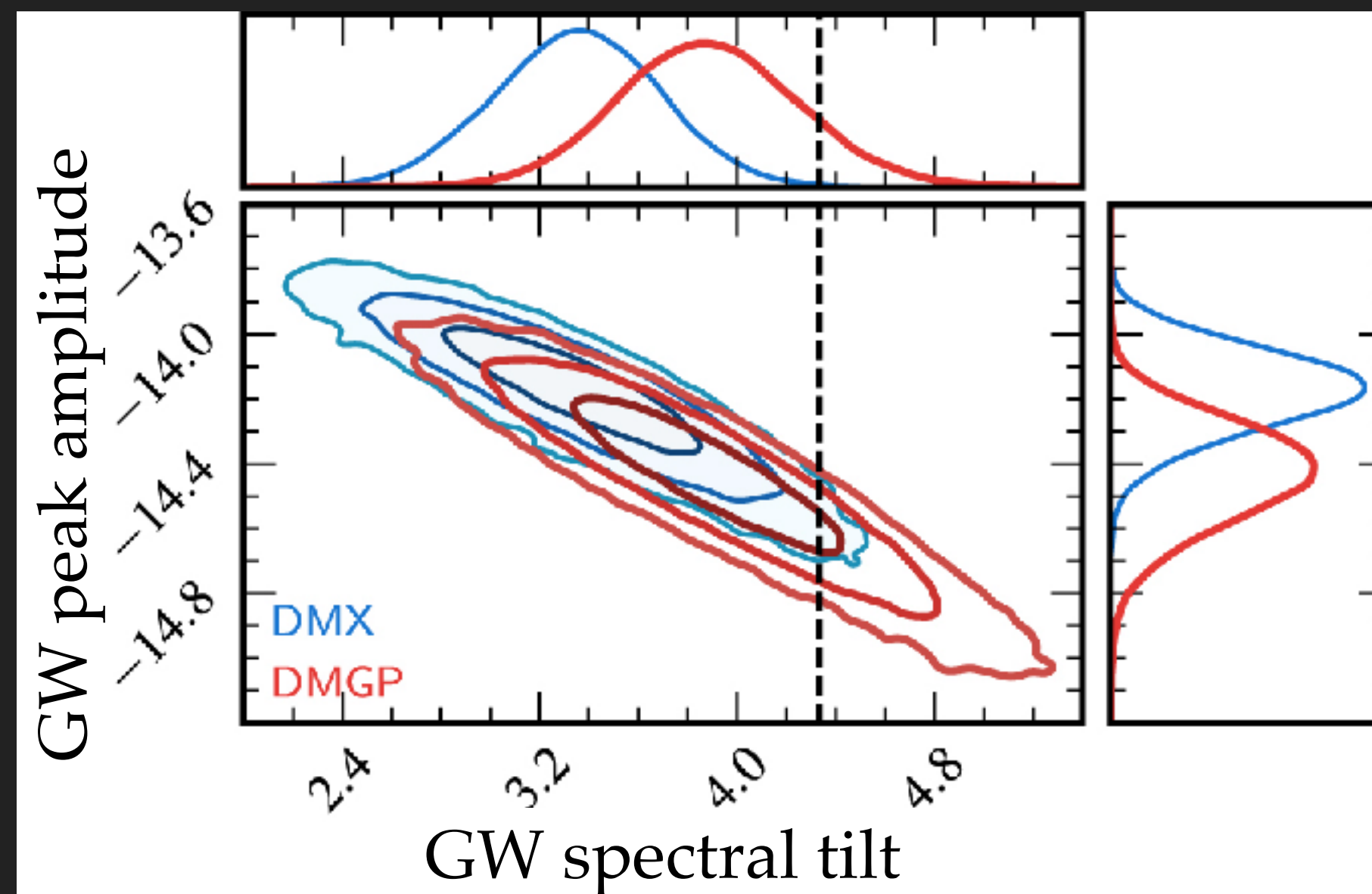


▶ Pulsar evolution: timing model

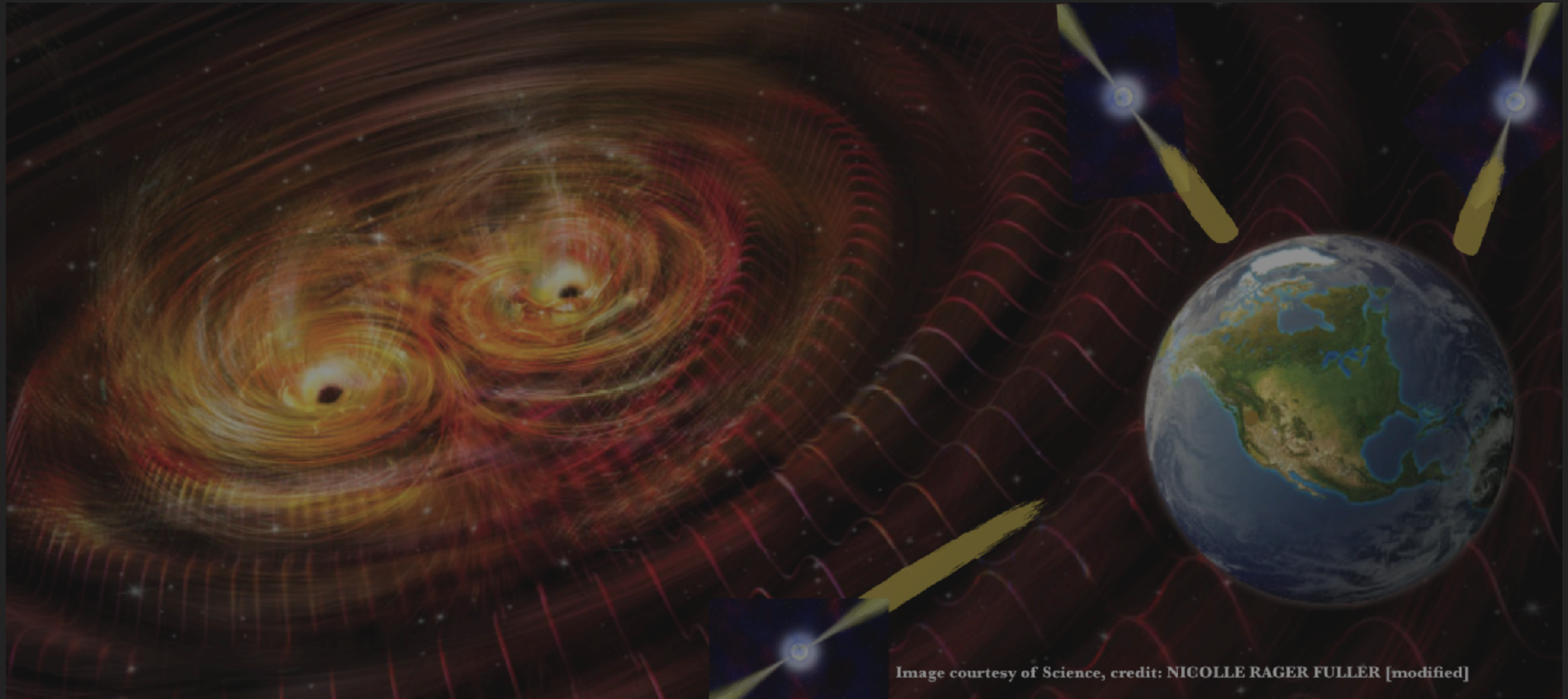
▶ Dispersion measure: ionised gas clouds



[ '23 NANOGrav-15 ]









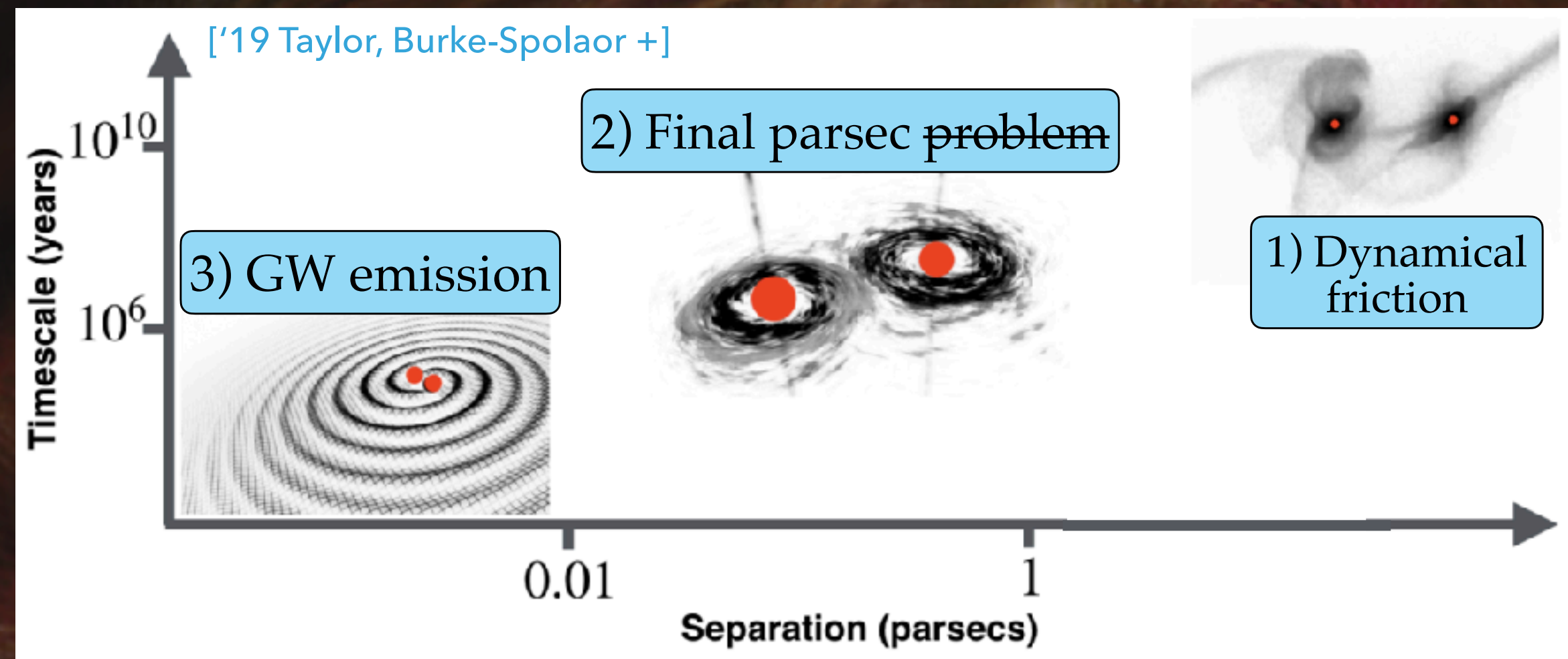


Image courtesy of Science, credit: NICOLLE RAGER FULLER [modified]



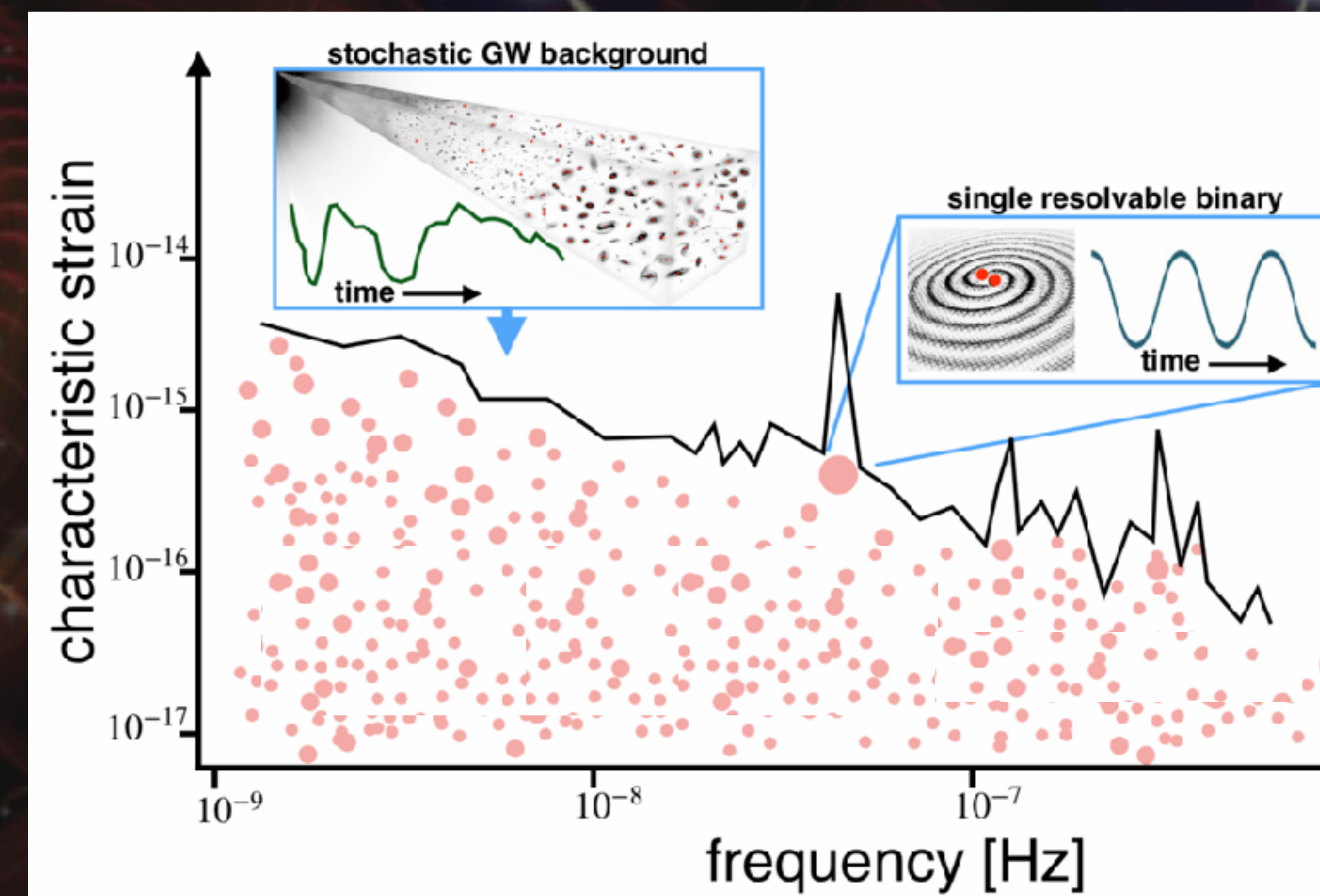
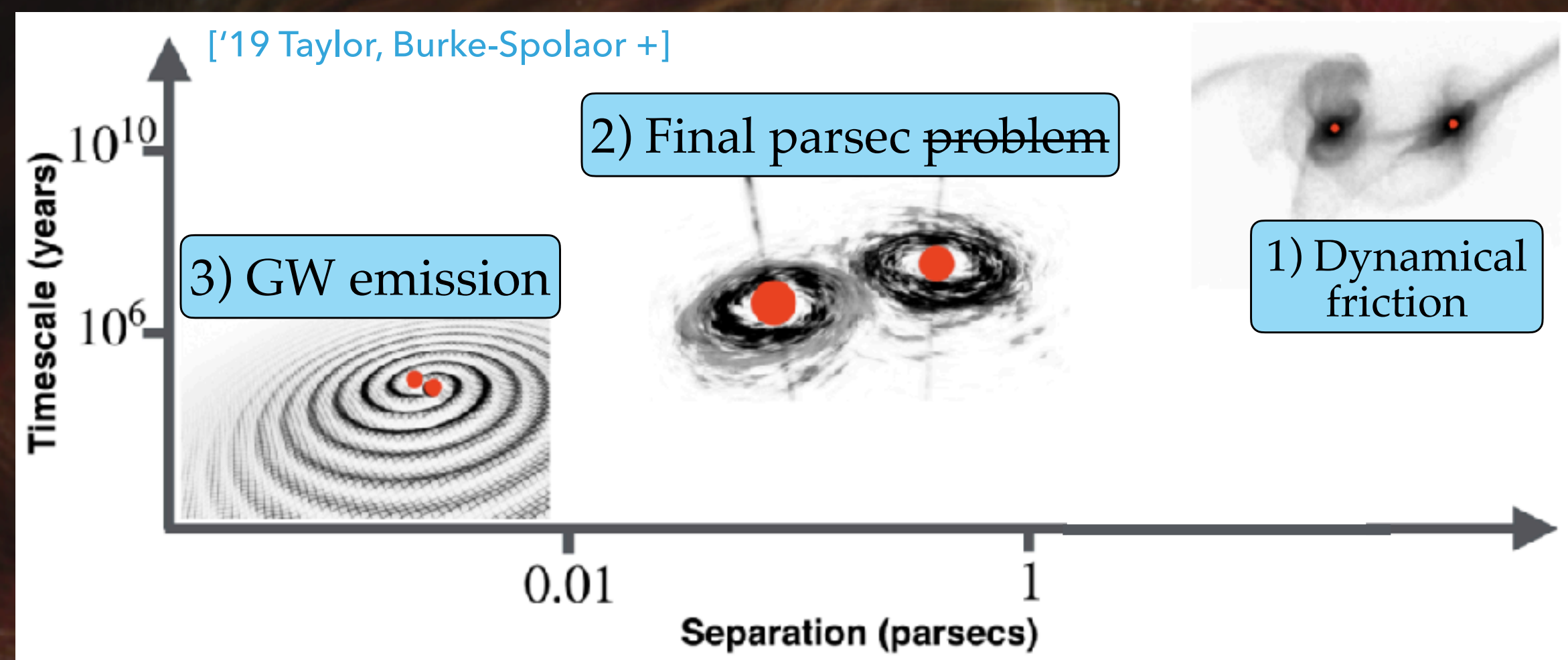
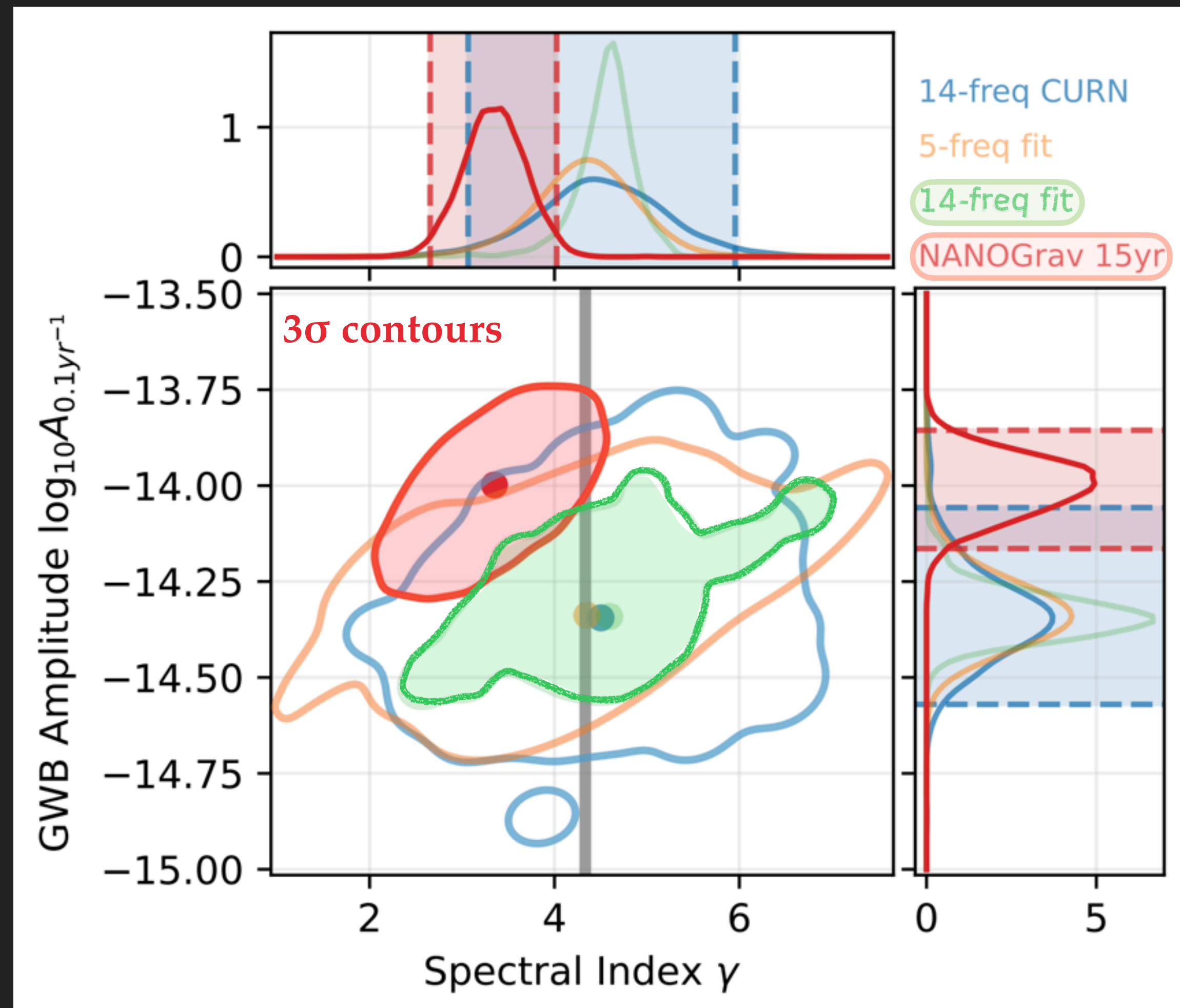


Image courtesy of Science, credit: NICOLLE RAGER FULLER [modified]

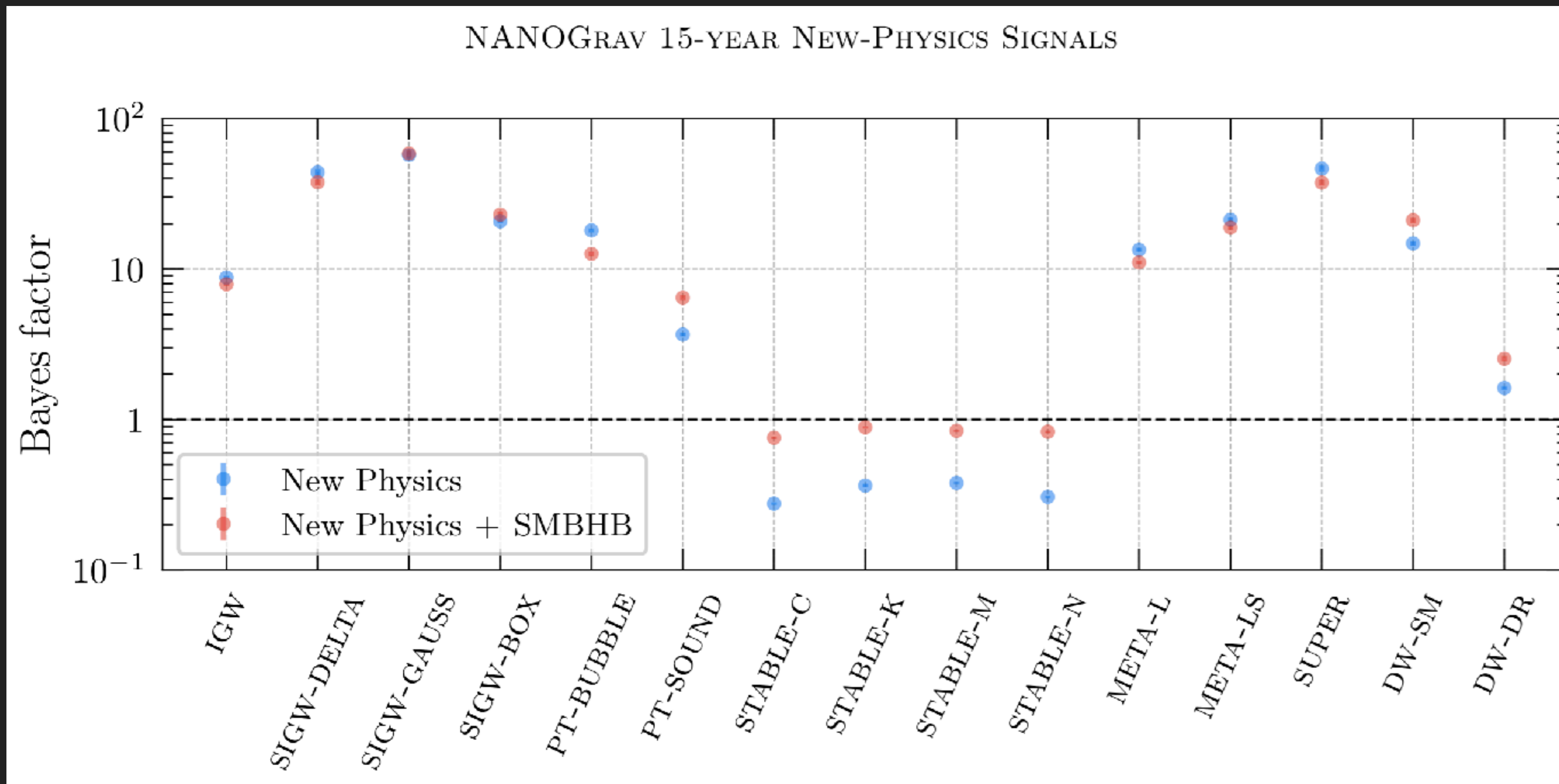


[ '23 Becsy+ (NG15) ]

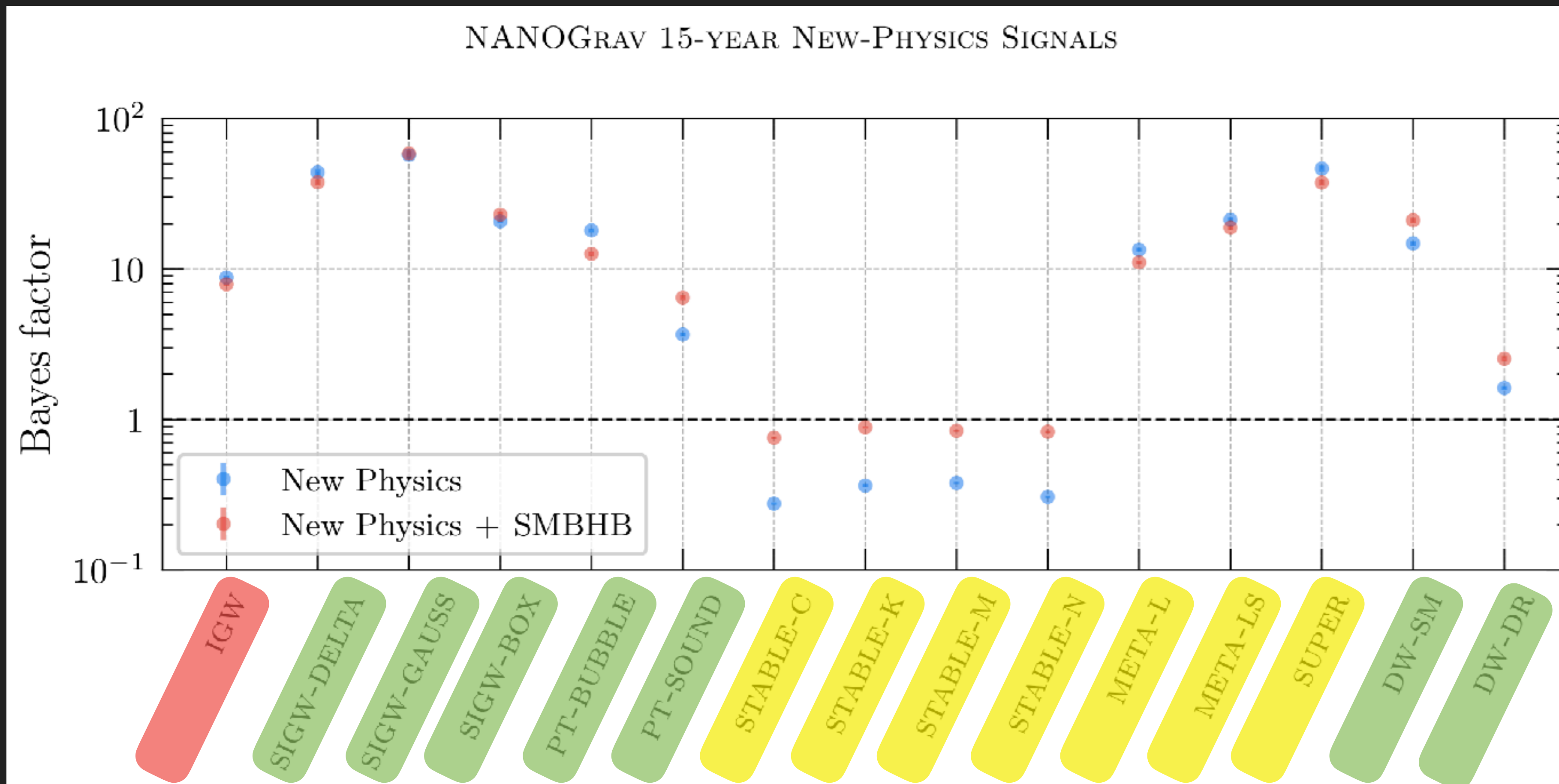


- Handle to discriminate different origins: identify **robust** spectral features







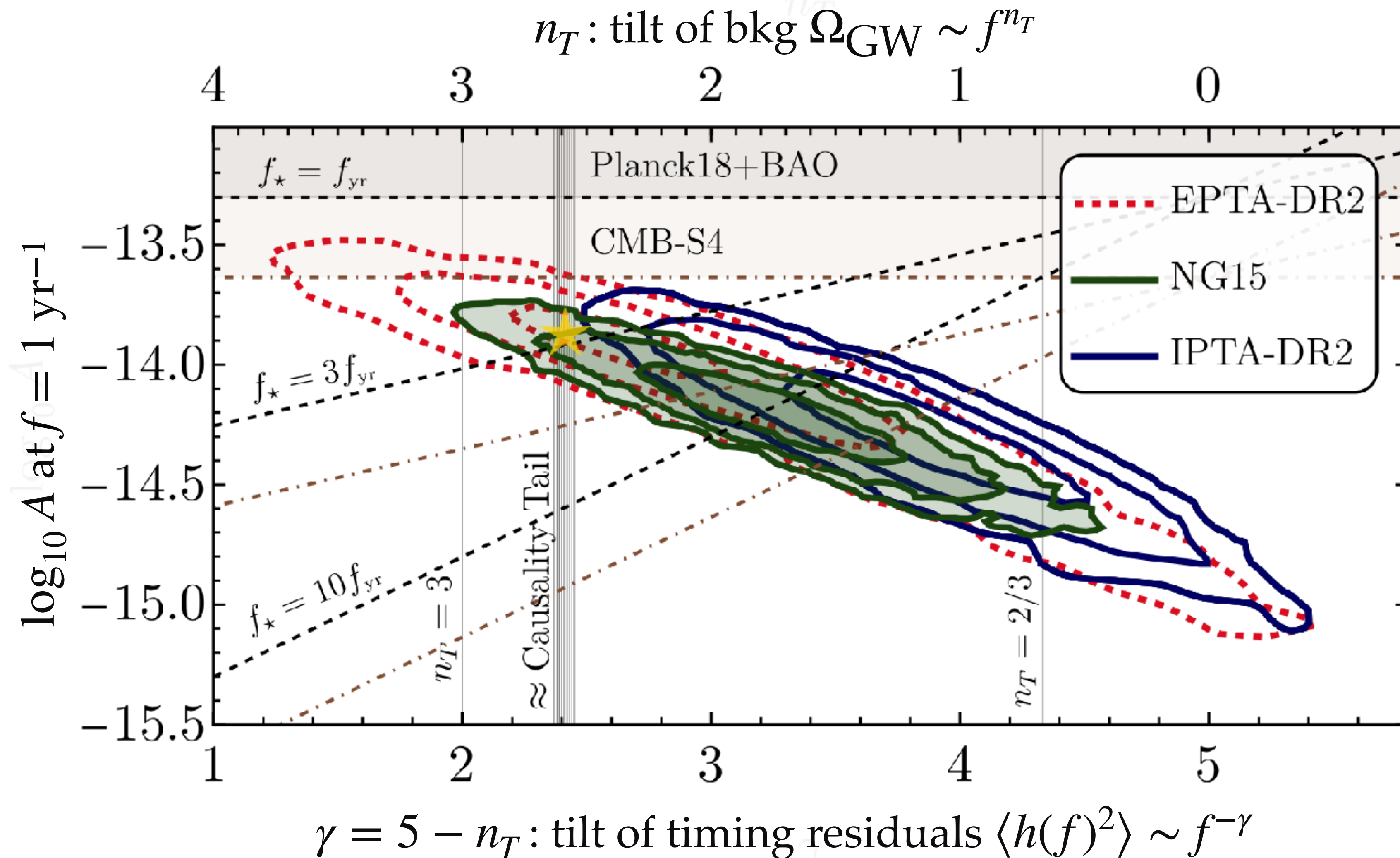
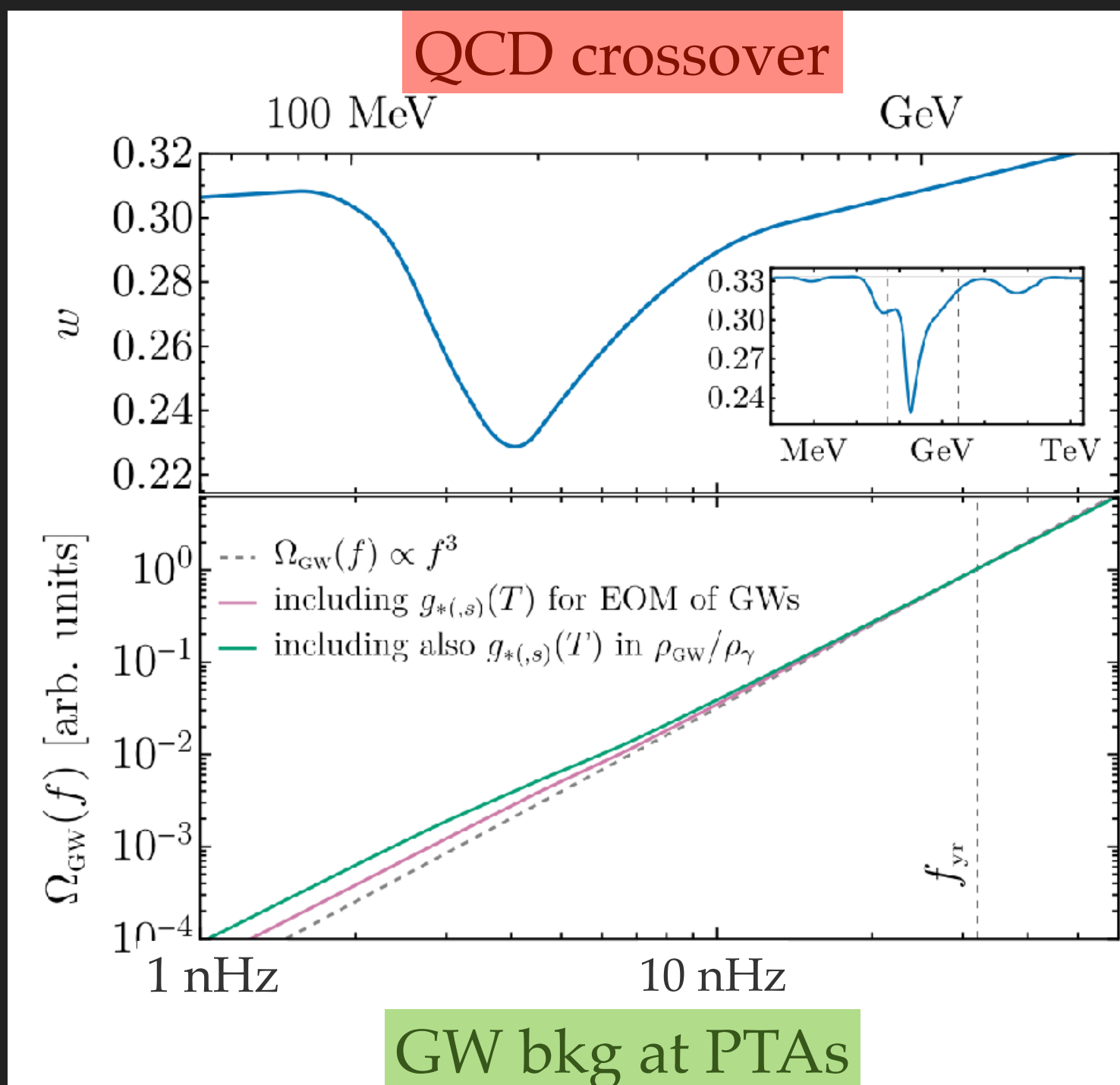


inflationary GWB: no CT

must go to CT at low f

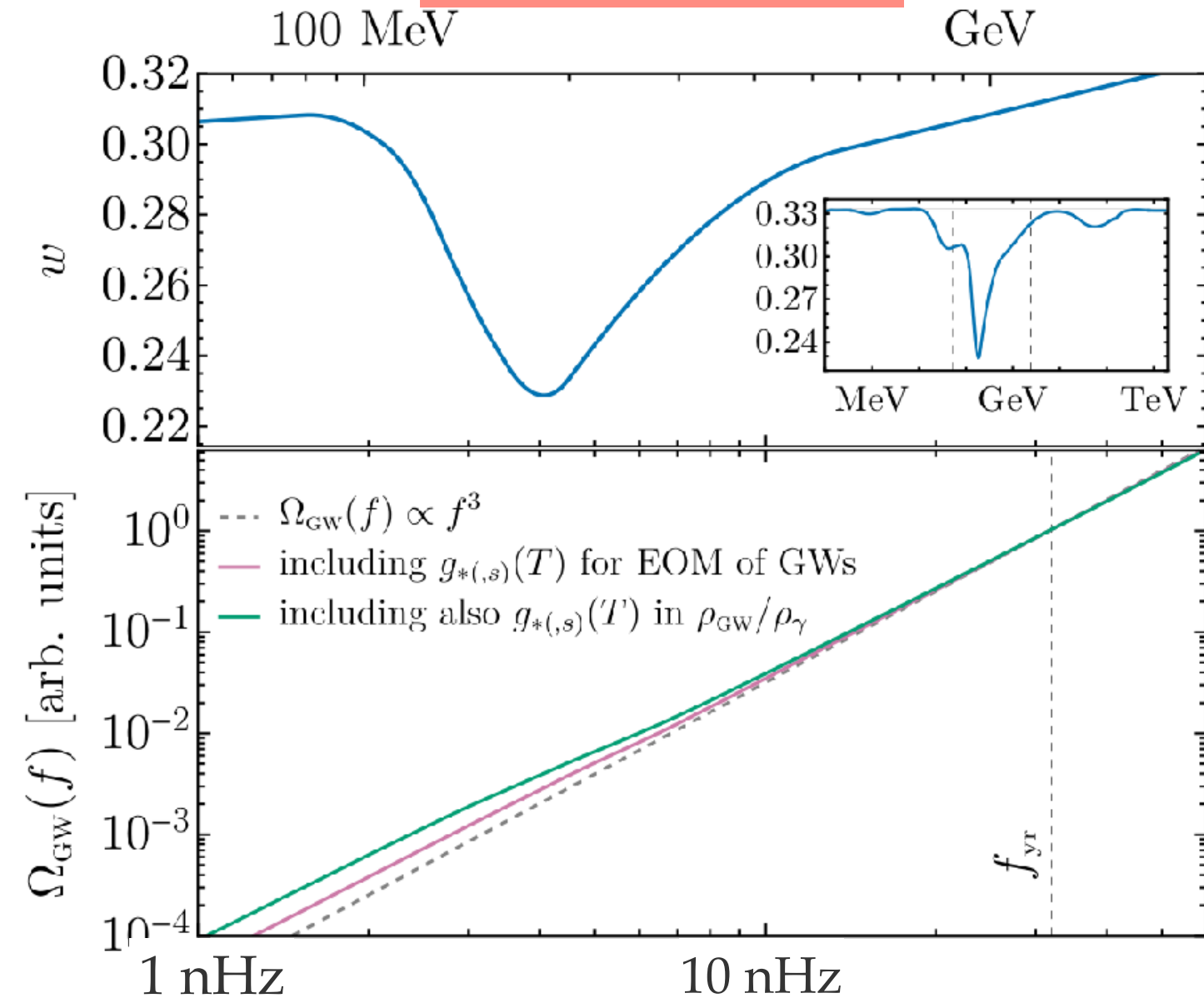
can go to CT at low f





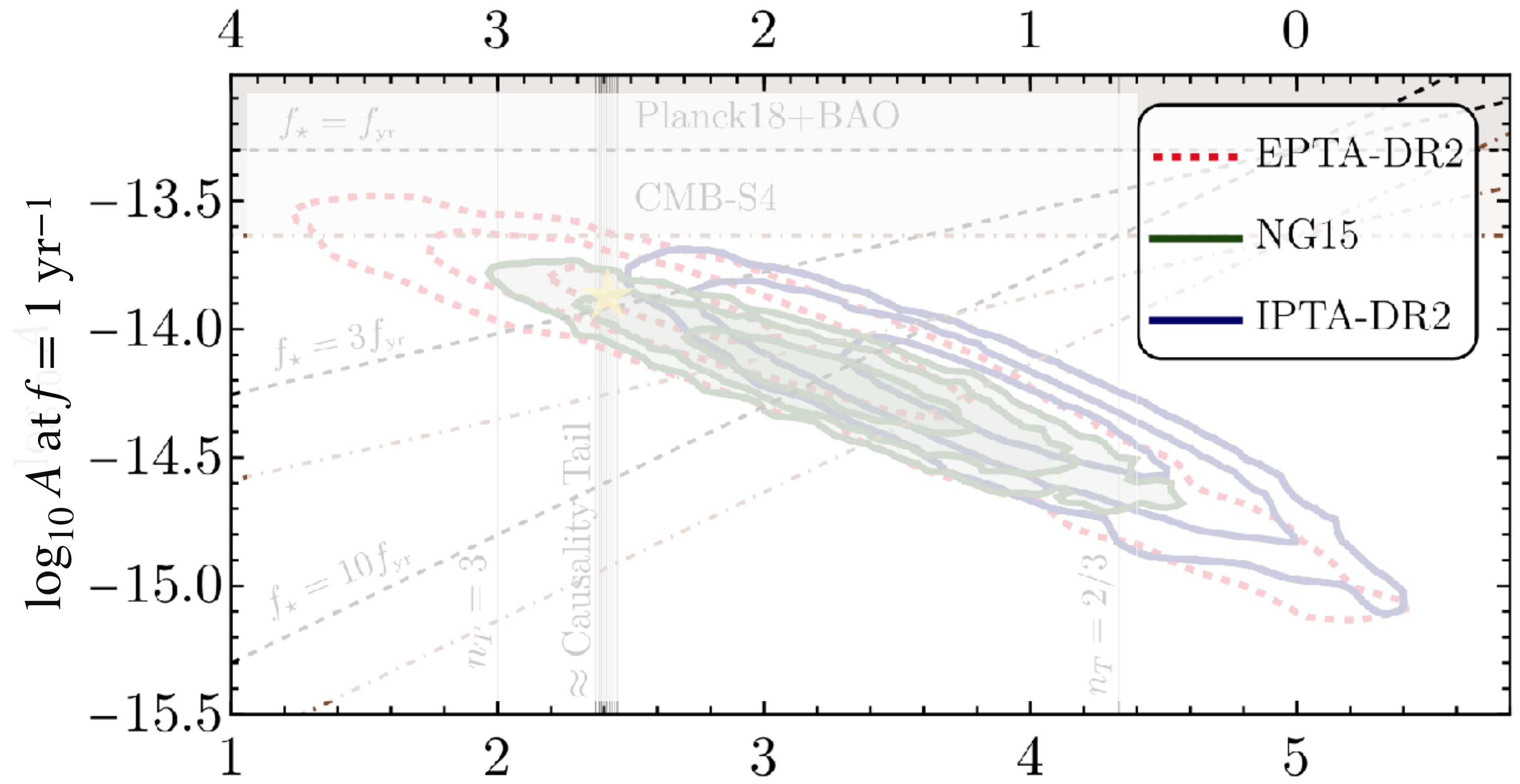


QCD crossover



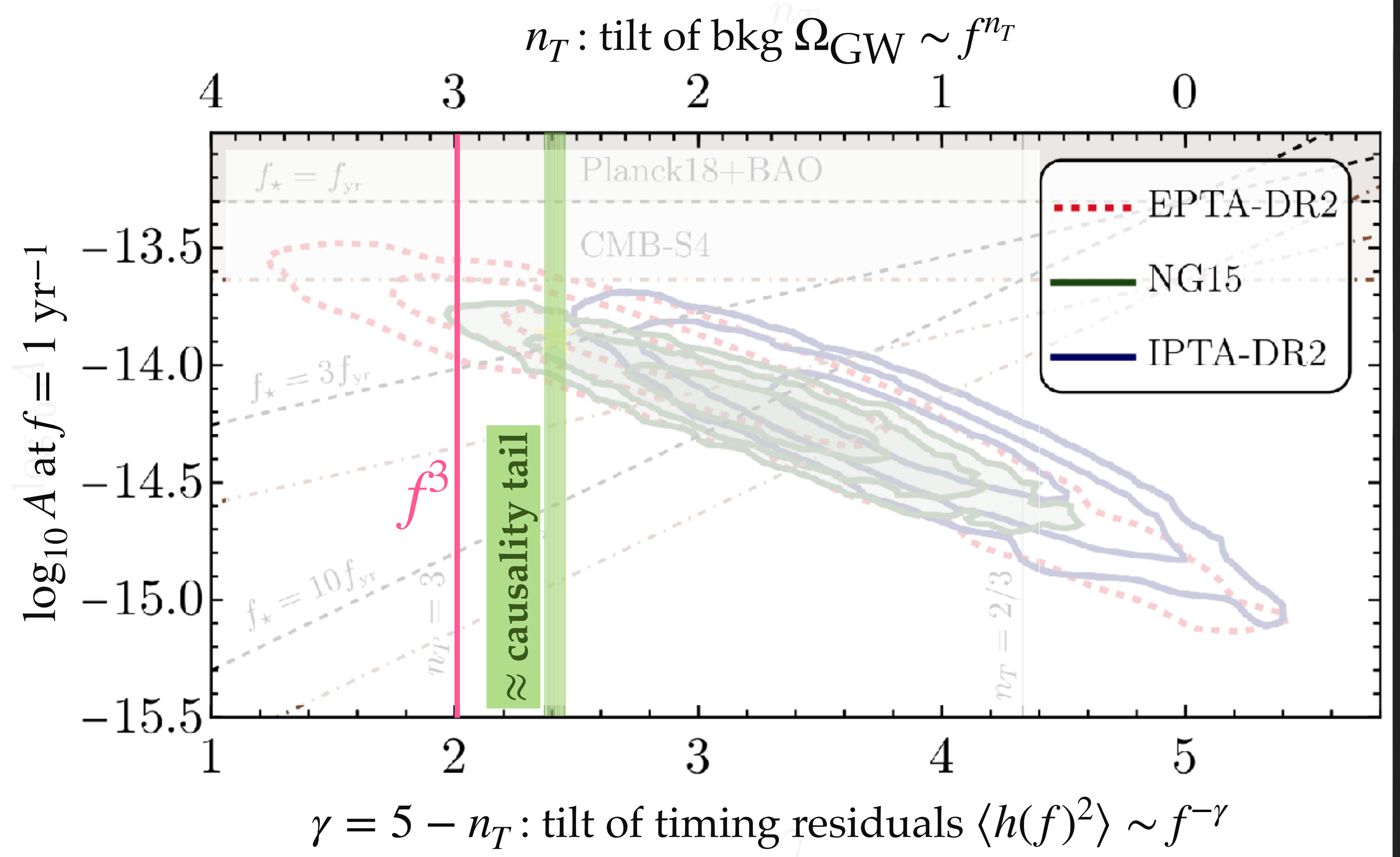
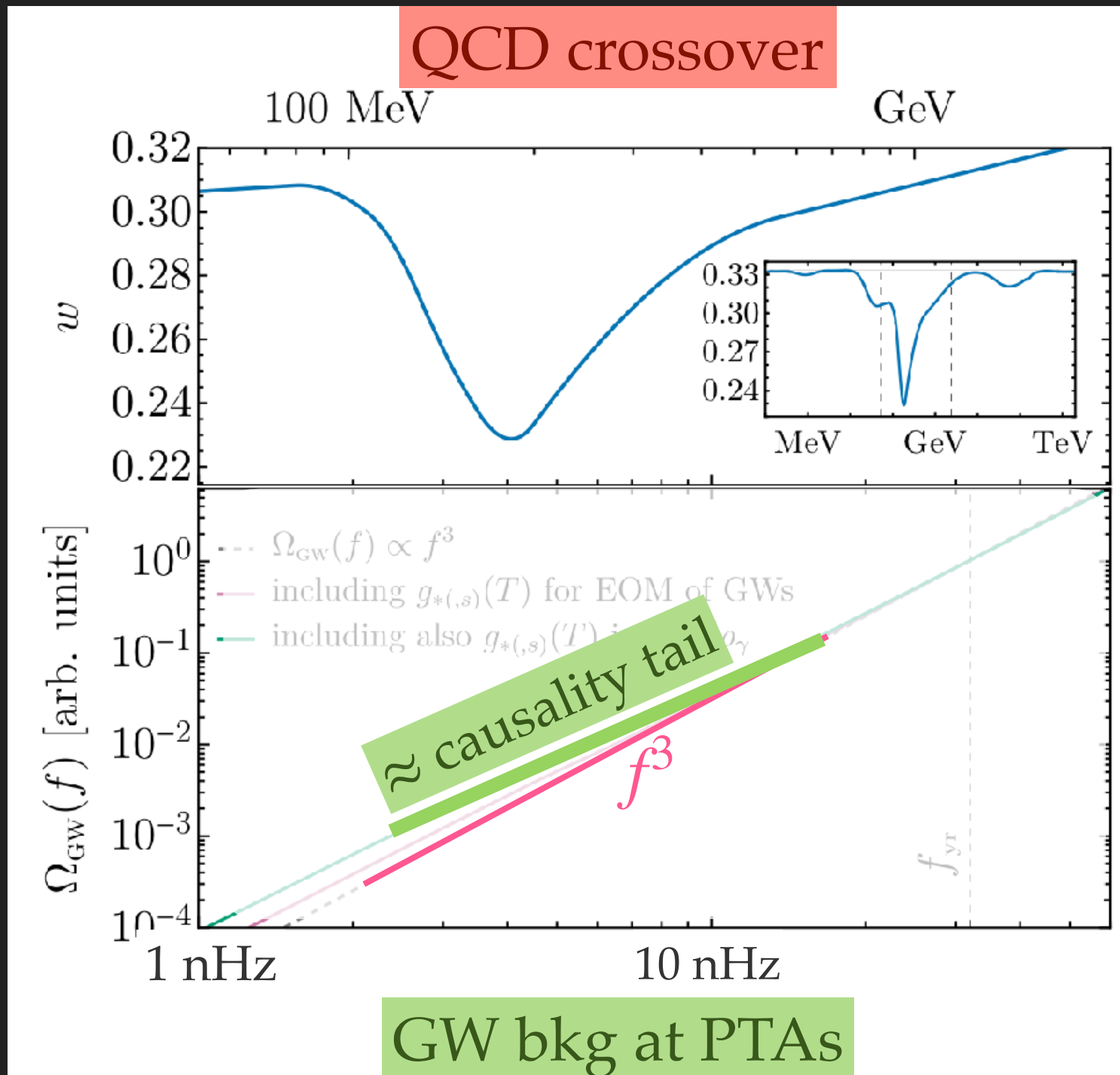
GW bkg at PTAs

$n_T$ : tilt of bkg  $\Omega_{\text{GW}} \sim f^{n_T}$

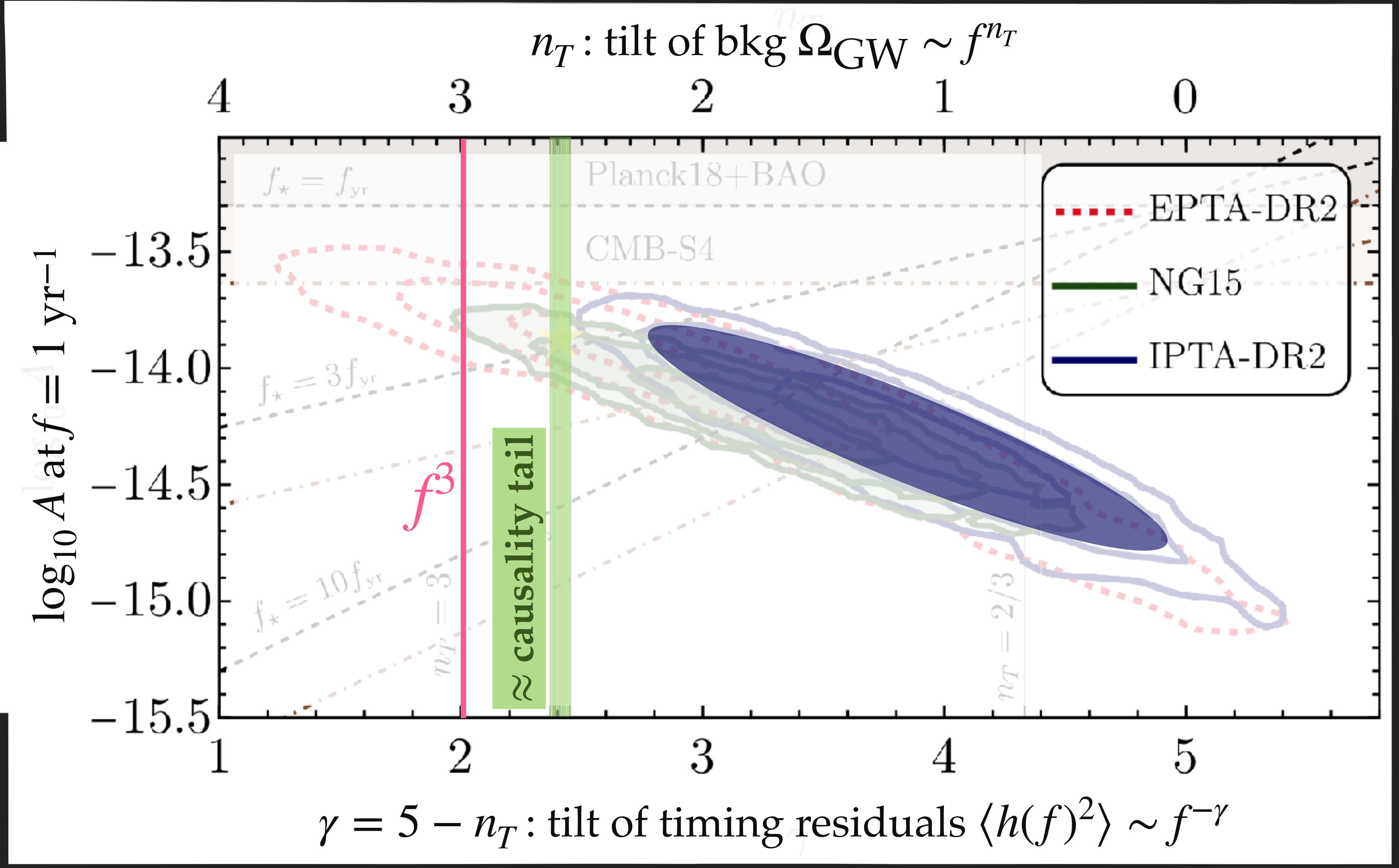
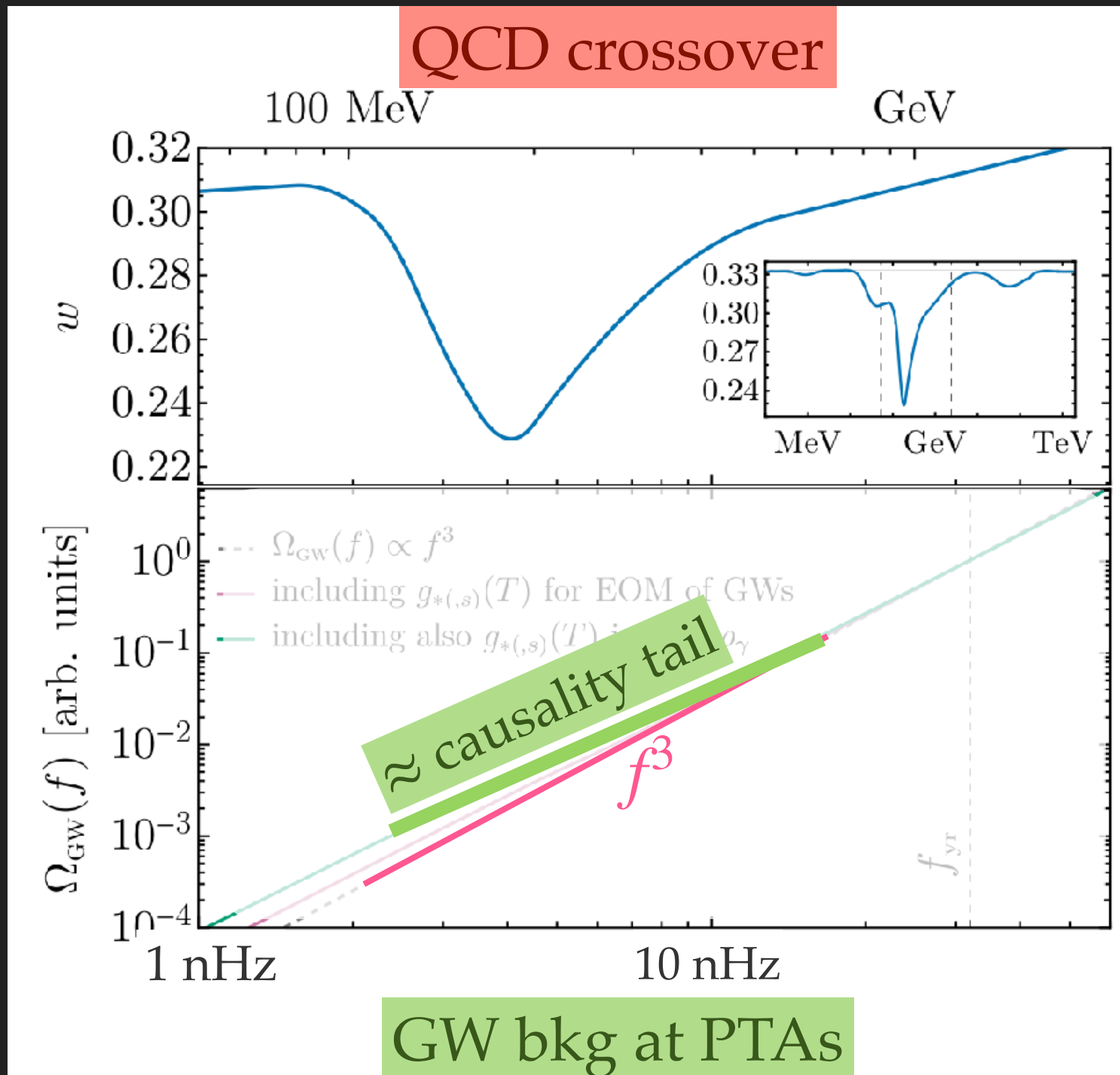


$\gamma = 5 - n_T$ : tilt of timing residuals  $\langle h(f)^2 \rangle \sim f^{-\gamma}$

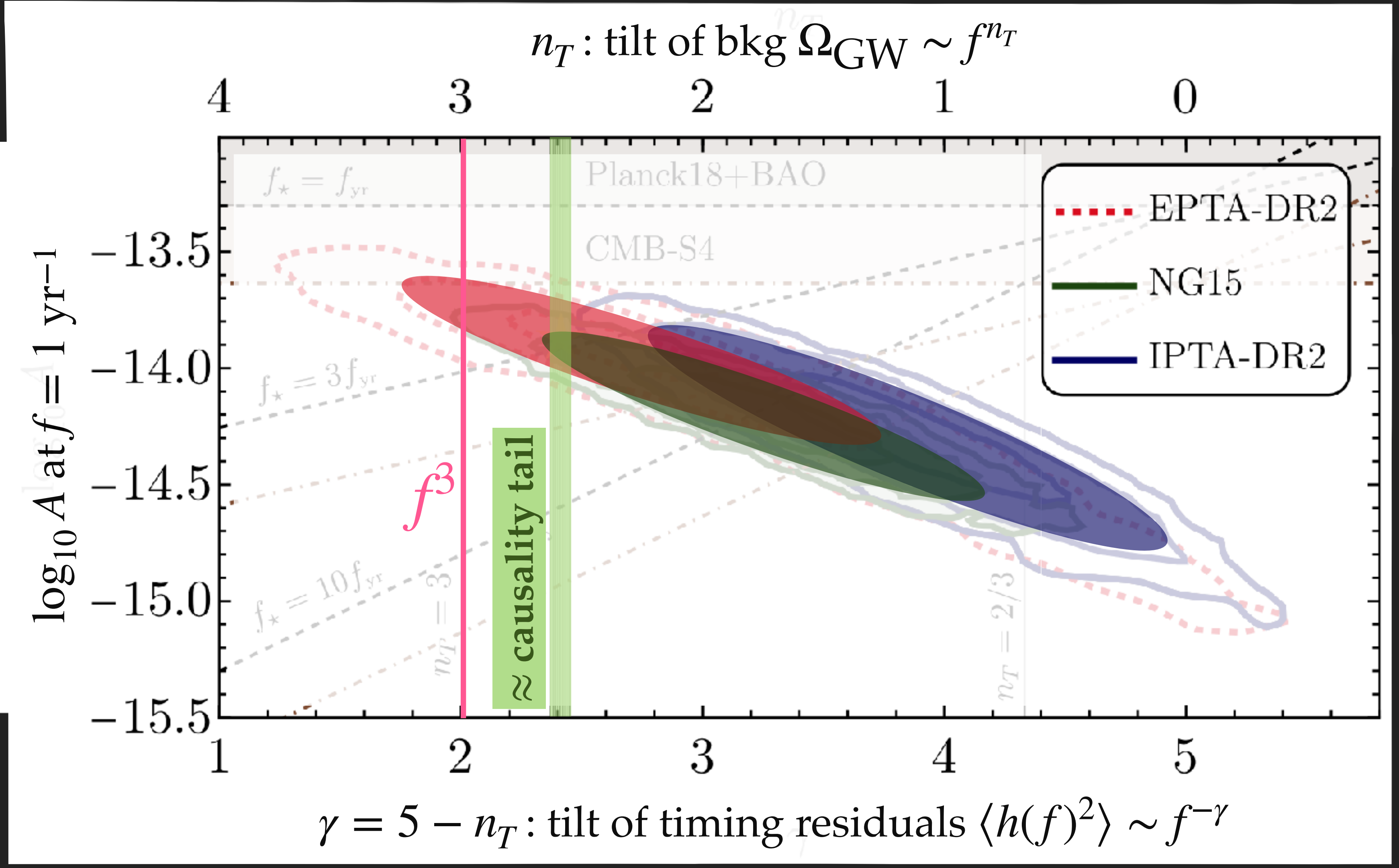
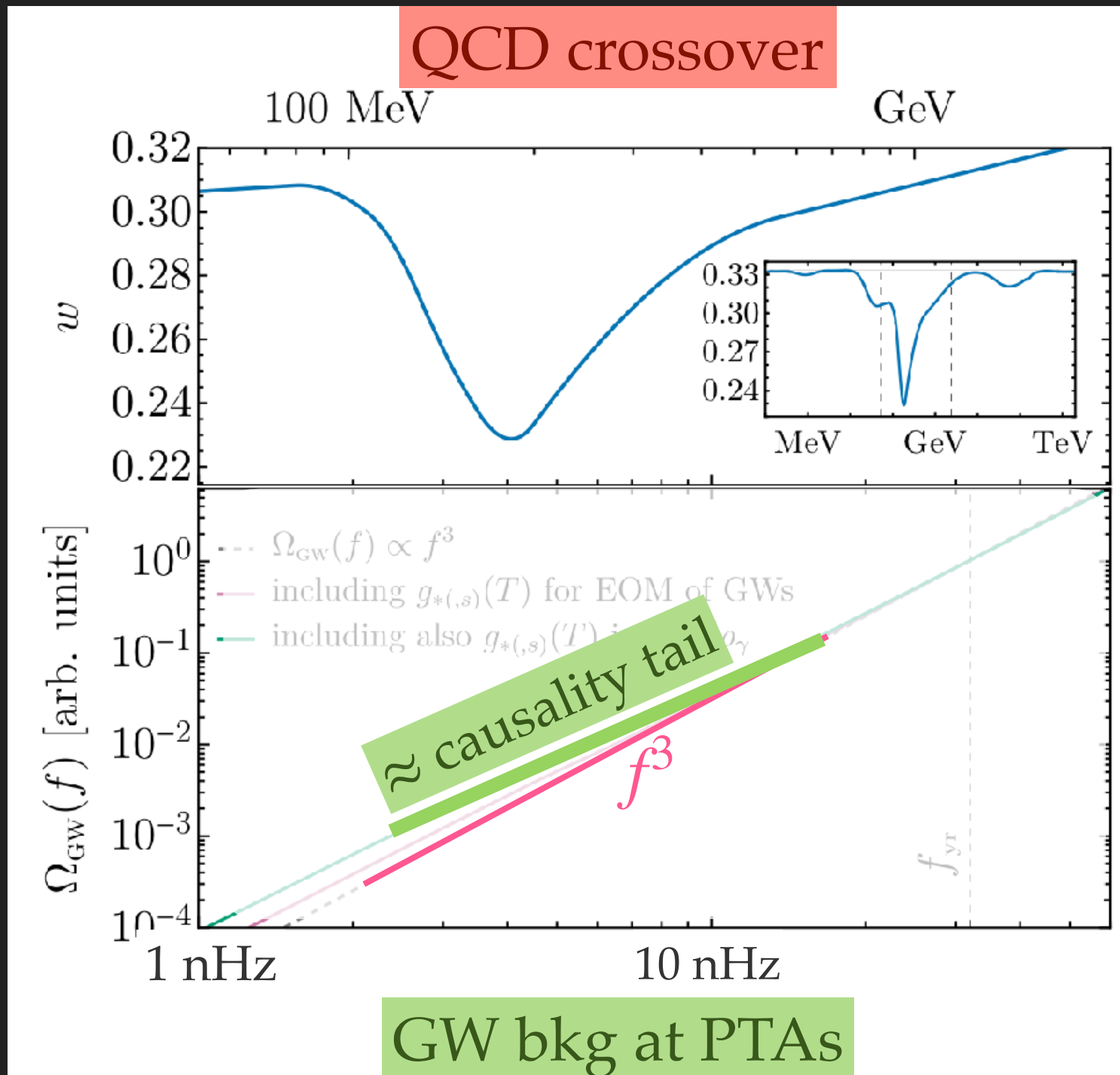












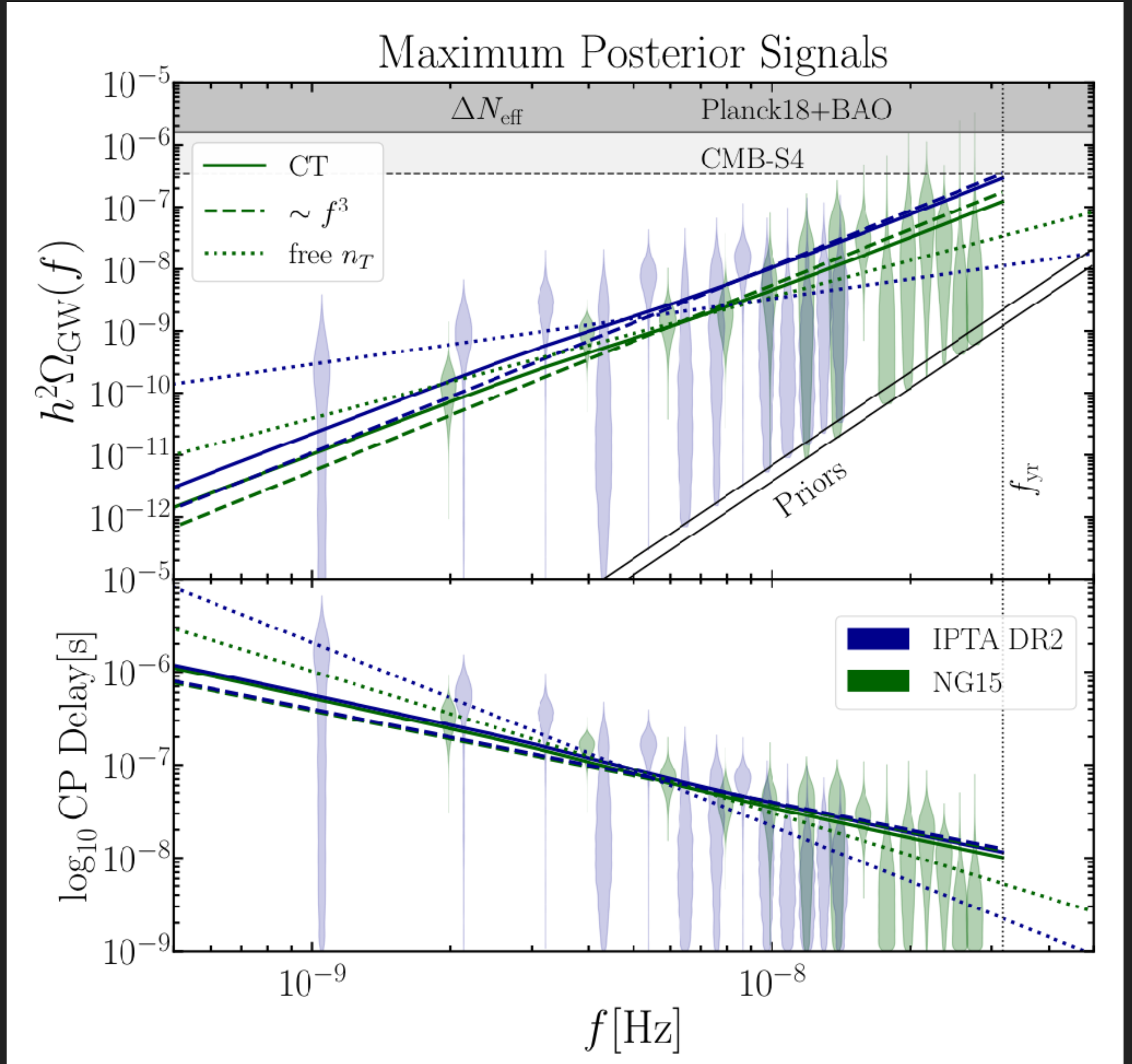


- ▶ Bayesian model comparison:  
causality tail vs.  $f^3$

$$\log_{10} \mathcal{B} = \begin{cases} 1.6 & (\text{NANOGrav-15}) \\ 1.3 & (\text{IPTA-DR2}) \end{cases}$$

(strong evidence)

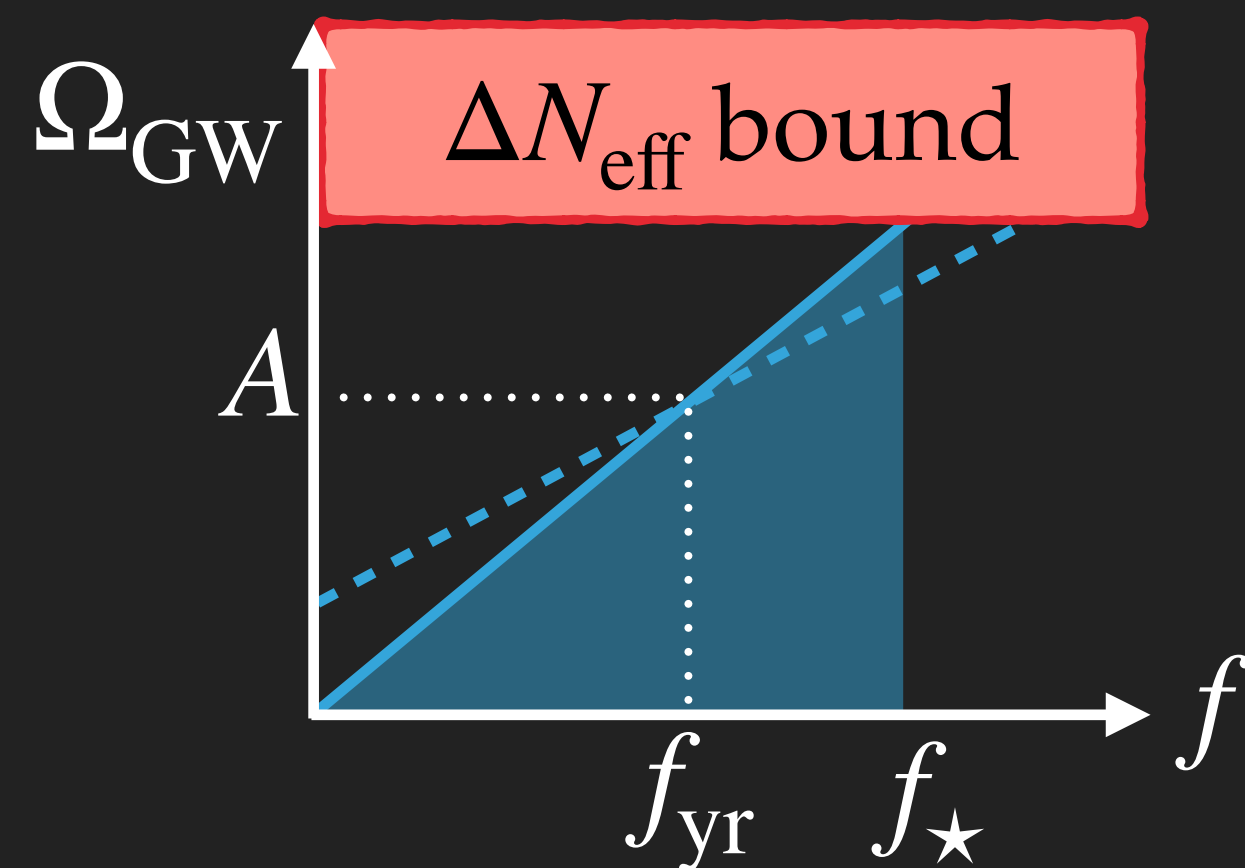
- ▶ Must account for causality tail  
in model comparison for PTAs





- ▶ Primordial GWs are relativistic energy not coupled to SM bath
- ▶  $\Rightarrow \Delta N_{\text{eff}}$  constrained by BBN, CMB

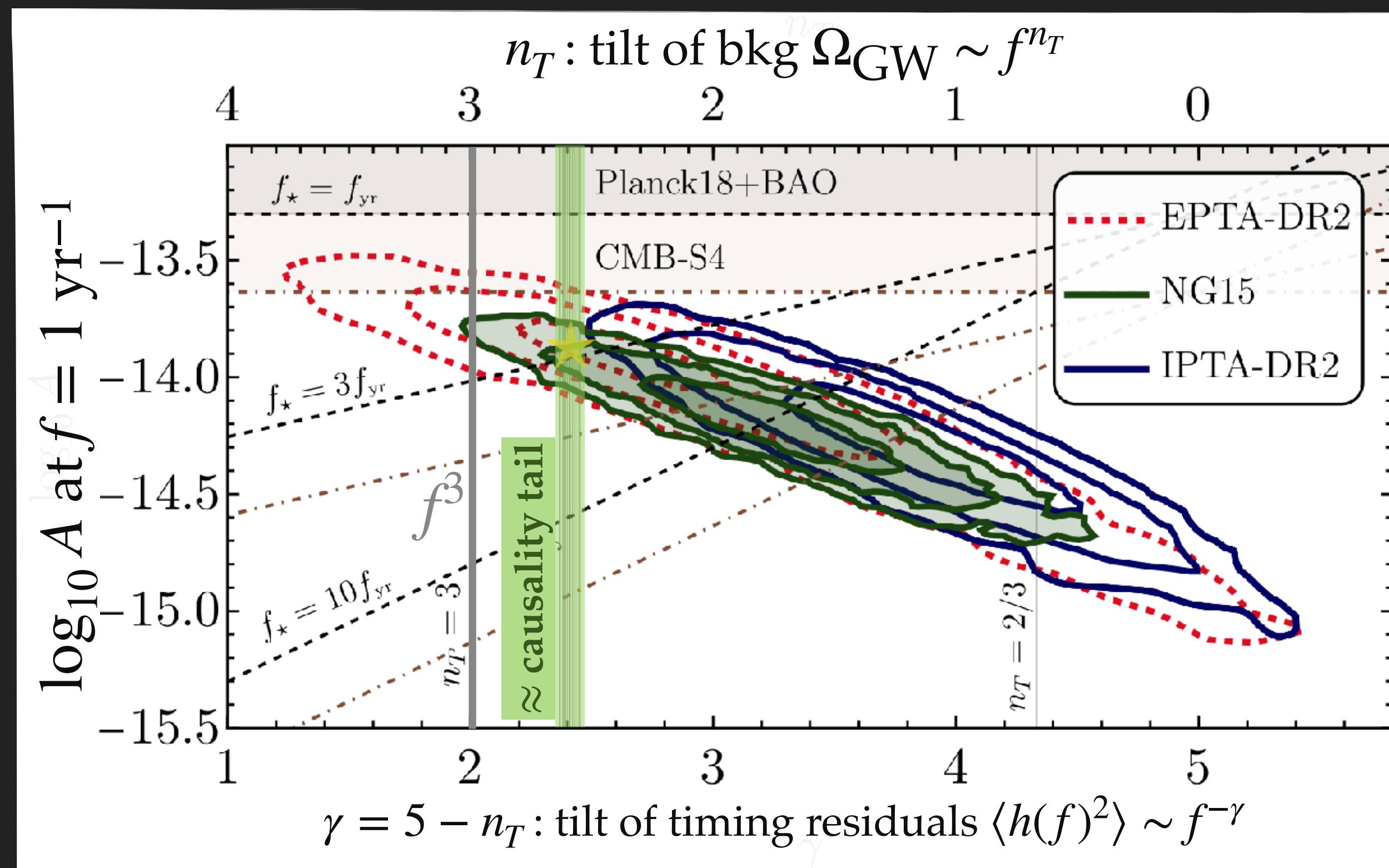
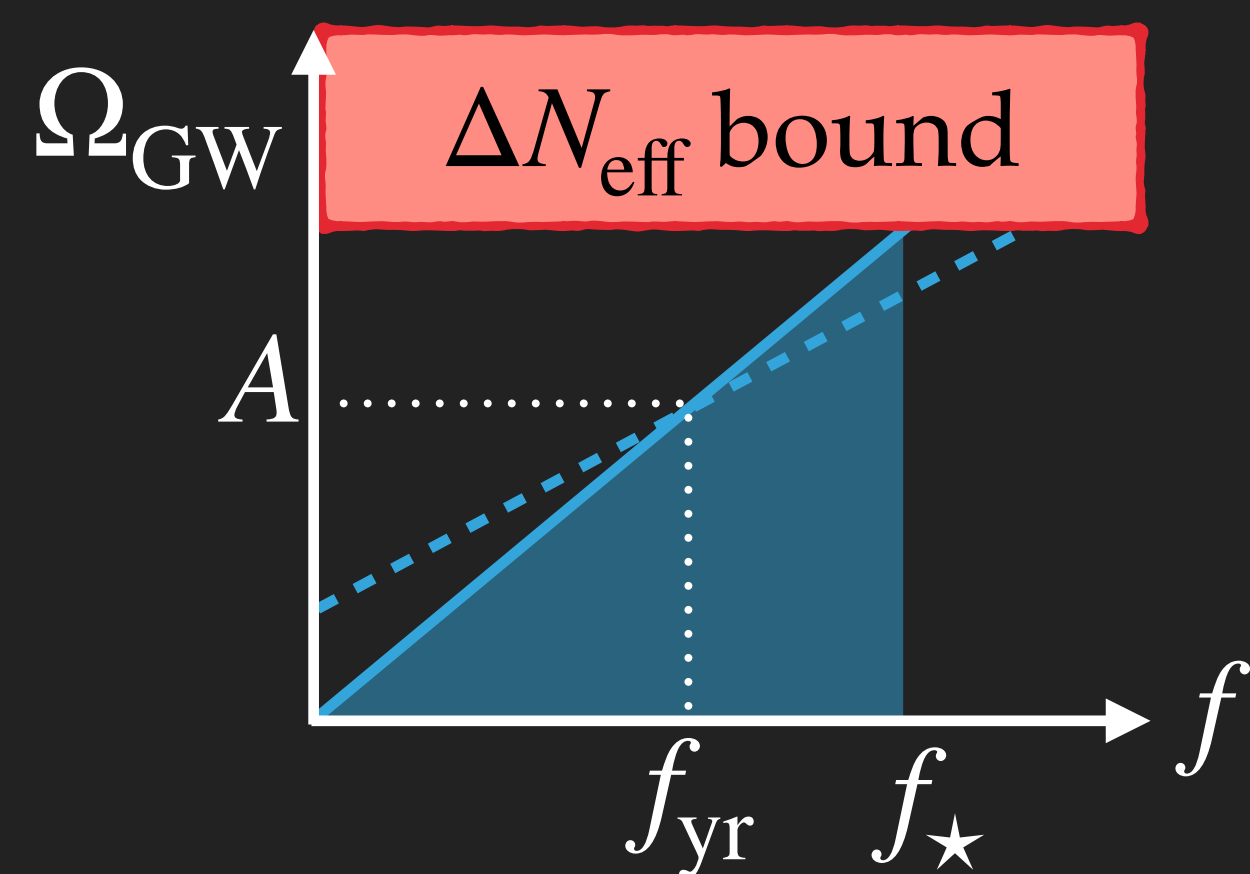
$$\Omega_{\text{CGW}} h^2 = 1.6 \cdot 10^{-6} \cdot \frac{\Delta N_{\text{eff}}^{\text{CGW}}}{0.28}$$





- ▶ Primordial GWs are relativistic energy not coupled to SM bath
- ▶  $\Rightarrow \Delta N_{\text{eff}}$  constrained by BBN, CMB

$$\Omega_{\text{CGW}} h^2 = 1.6 \cdot 10^{-6} \cdot \frac{\Delta N_{\text{eff}}^{\text{CGW}}}{0.28}$$



[ '23 Franciolini, DR, Rompineve ]

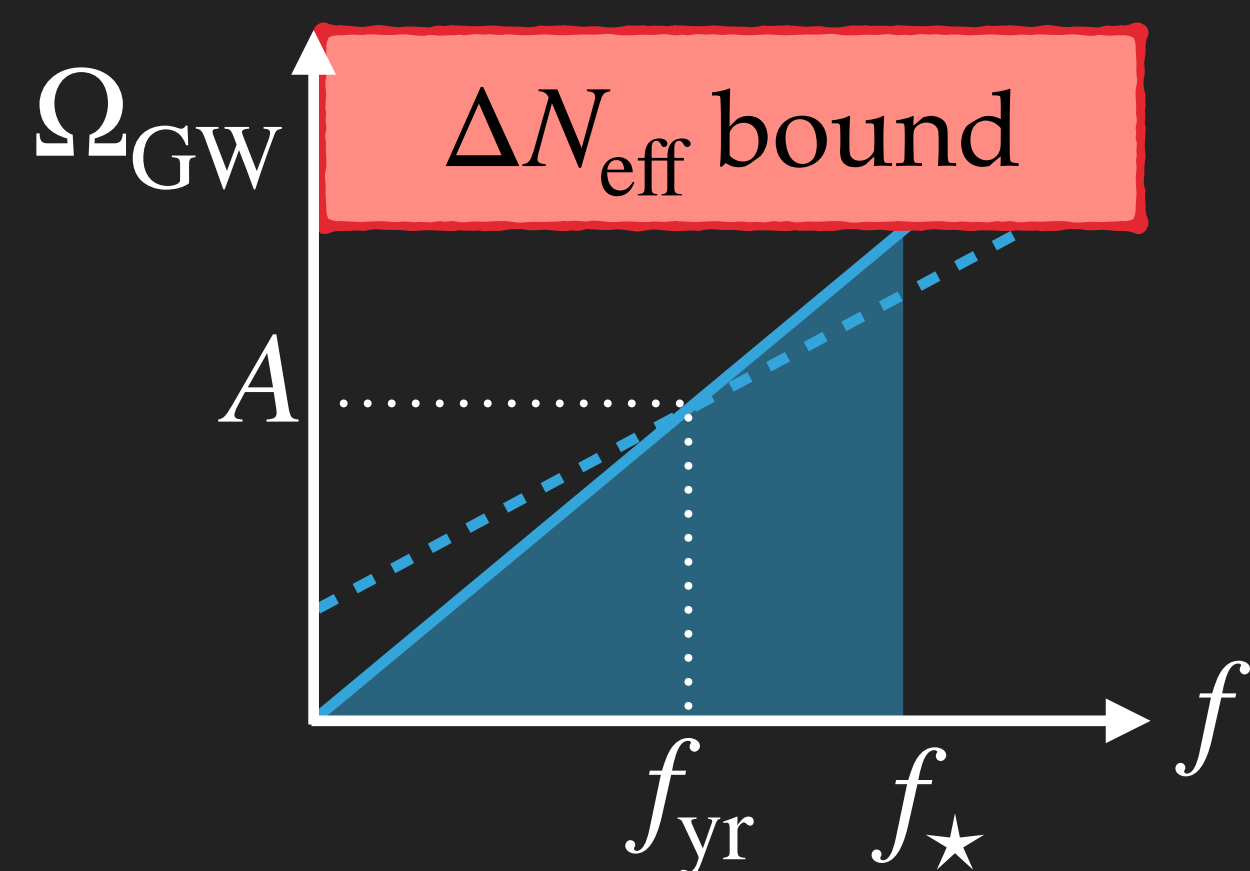




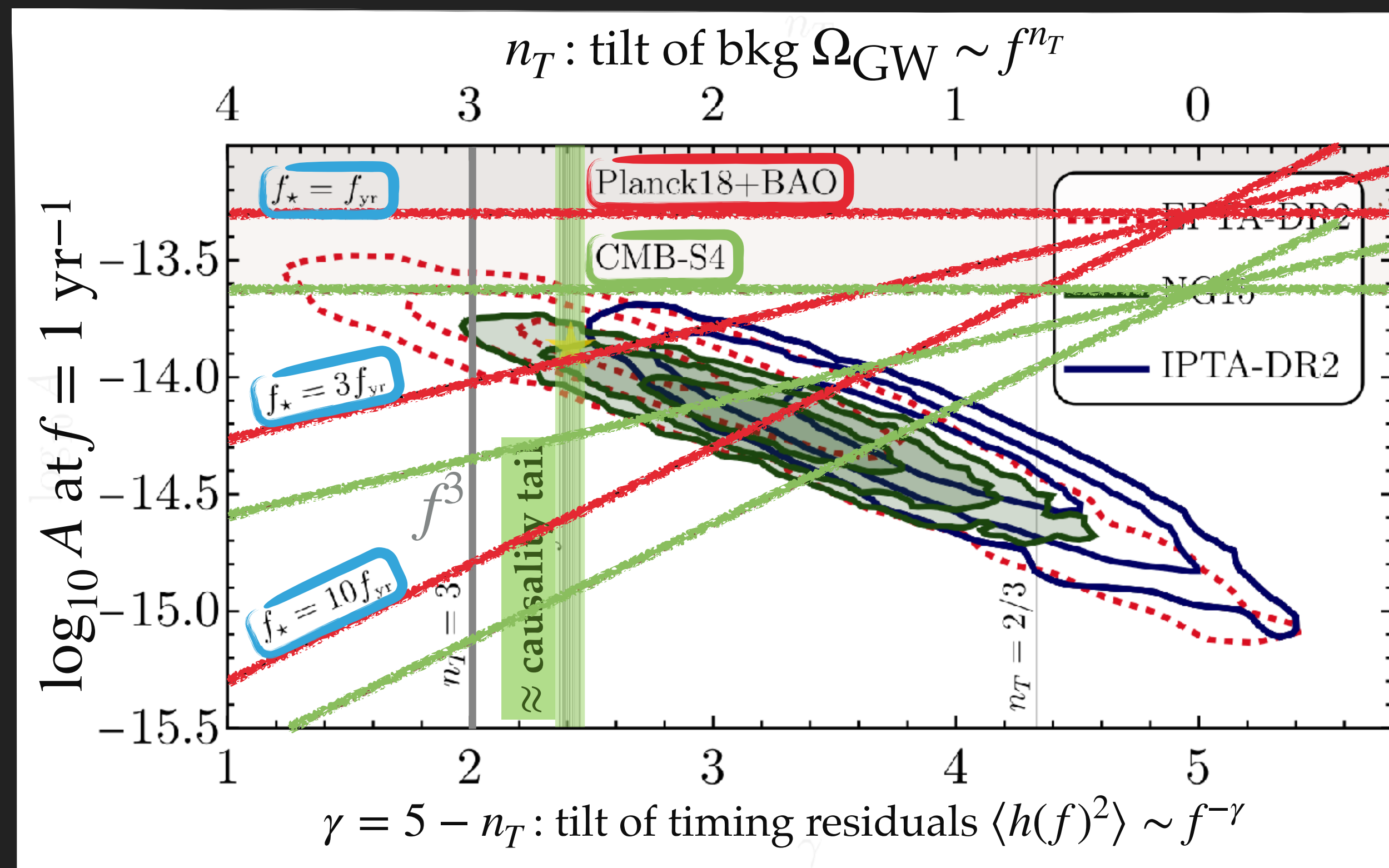


- ▶ Primordial GWs are relativistic energy not coupled to SM bath
- ▶  $\Rightarrow \Delta N_{\text{eff}}$  constrained by BBN, CMB

$$\Omega_{\text{CGW}} h^2 = 1.6 \cdot 10^{-6} \cdot \frac{\Delta N_{\text{eff}}^{\text{CGW}}}{0.28}$$



- ▶ Primordial GWB can't grow until too large  $f_\star$
- ▶ Signature in CMB if GWB is primordial!

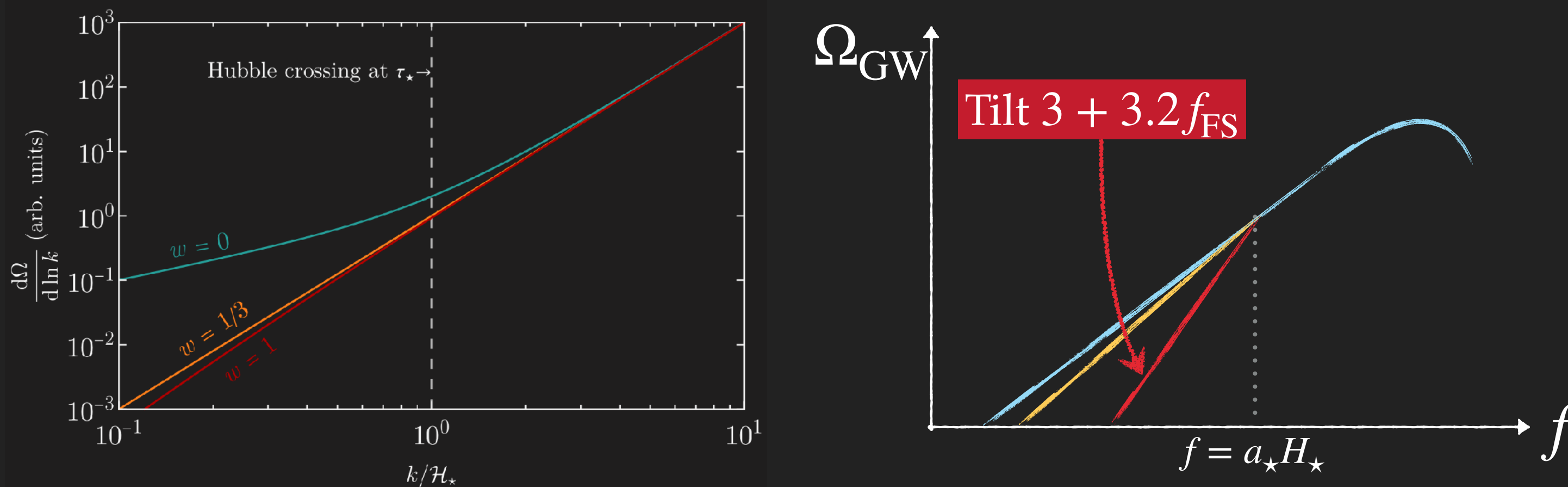


[ '23 Franciolini, DR, Rompineve ]

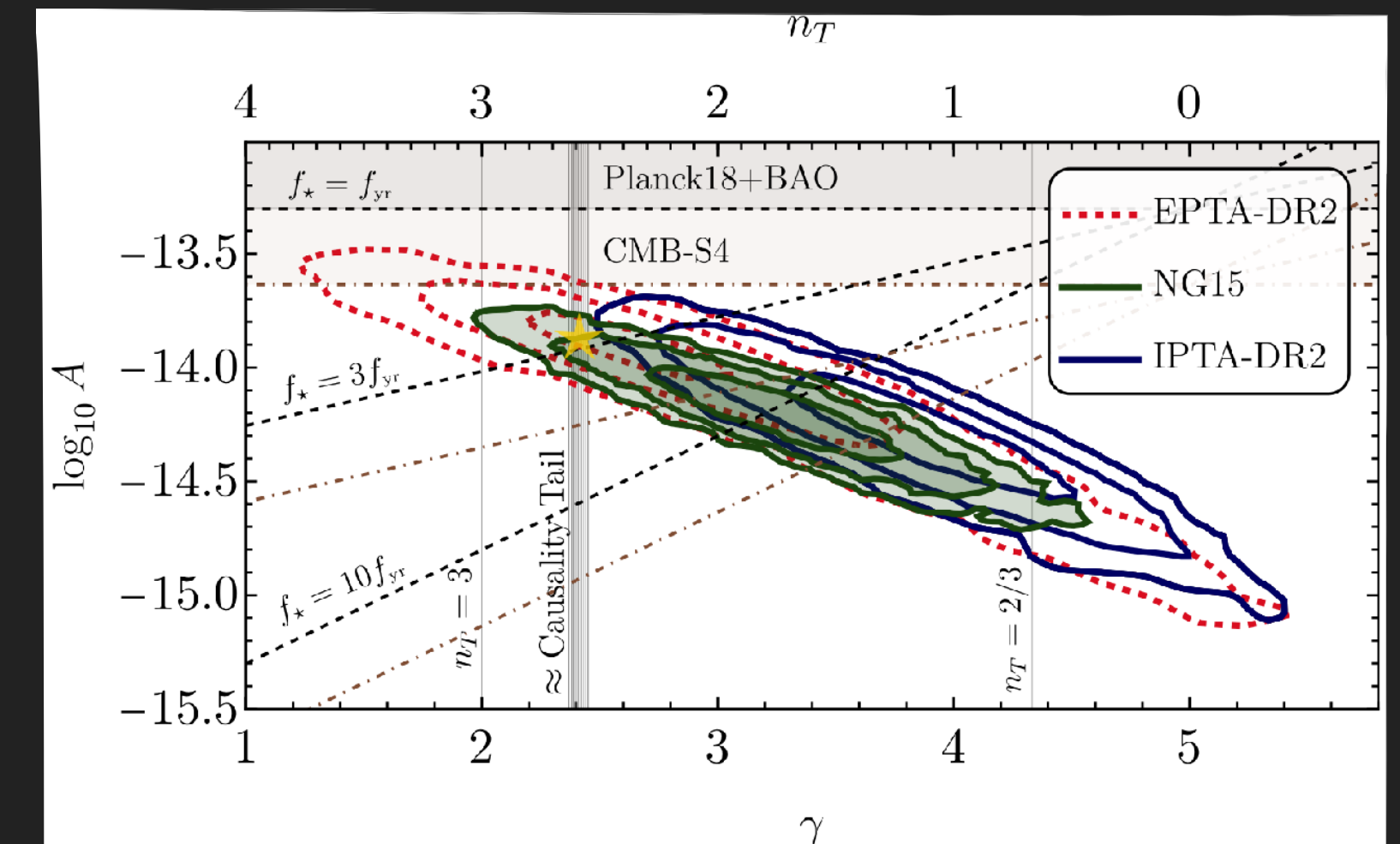


- ▶ Causality tail  $\Leftrightarrow$  probe  $w$  and  $f_{\text{FS}}$  independently of the primordial GW source

- ▶ SM effects on the causality tail must be included at PTAs



Hook, Marques-Tavares, **DR**  
JHEP (2021)



Franciolini, **DR**, Rompineve  
Phys. Rev. Lett. (2024)

Buchalter Cosmology Prize 2023



9 — 20 September 2024



POLLICA  
PHYSICS  
CENTRE

# Fundamental Physics and Gravitational Wave Detectors

Gravitational waves (GWs) from compact binaries have been detected in the kHz regime, and Pulsar Timing Arrays are revealing a GW background in the nHz band. Experimental ideas spanning from nHz to MHz are being explored to cover the gravitational spectrum. This is an exciting time to explore what GWs can uncover about fundamental physics and cosmology by detecting astrophysical and primordial GW sources, as well as exploring effects that could be induced by various Dark Matter candidates.

This workshop will unite astrophysicists, cosmologists, and particle physicists to discuss open questions and define theoretical targets, guiding the field in experimental strategies across the gravitational spectrum to learn about our Universe. All of this is to be set against the backdrop of the medieval town of Pollica in southwestern Italy (Cilento region).



## Participants will include:

Bruce Allen  
(AEI Potsdam)

Daniel Holz  
(U. of Chicago)

Nataliya Porayko  
(MPIFR Bonn)

Chiara Caprini  
(CERN & U. Geneva)

Andrea Mitridate  
(DESY)

Nicholas Rodd  
(LBNL Berkeley)

Reed Essick  
(CITA Toronto)

Samaya Nissanke  
(UvA)

Jorinde Van De Vis  
(Leiden U.)

José María Ezquiaga  
(NBI Copenhagen)

Chris Overstreet  
(Johns Hopkins)

Sarah Vigeland  
(UW Milwaukee)

Daniel Figueroa  
(IFIC Valencia)

Marco Peloso  
(Padova U.)

Matias Zaldarriaga  
(IAS Princeton)

Maya Fishbach  
(CITA Toronto)

Antoine Petiteau  
(CEA)

## ORGANIZING COMMITTEE

Sebastian Baum (RWTH Aachen)

Djuna Croon (Durham University)

Paolo Pani (Sapienza University &  
INFN Roma)

Davide Racco (ETH & University of  
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Géraldine Servant (DESY & Universität  
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Sponsored by



# APPLICATIONS ARE OPEN!

Davide Racco (ETH - U. Zürich)

LAPTh Annecy





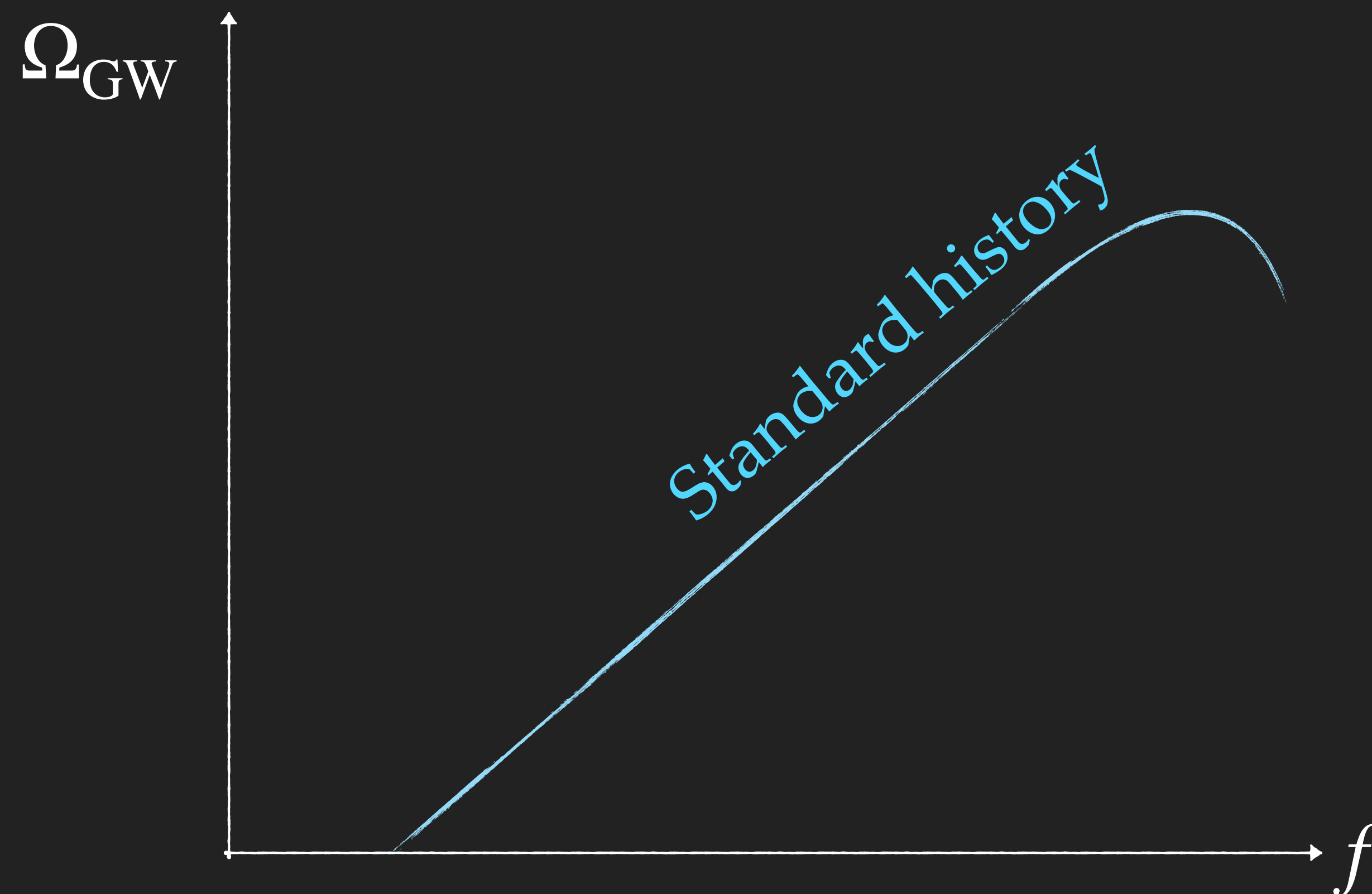
**Thank you for your attention!**



BACKUP SLIDES

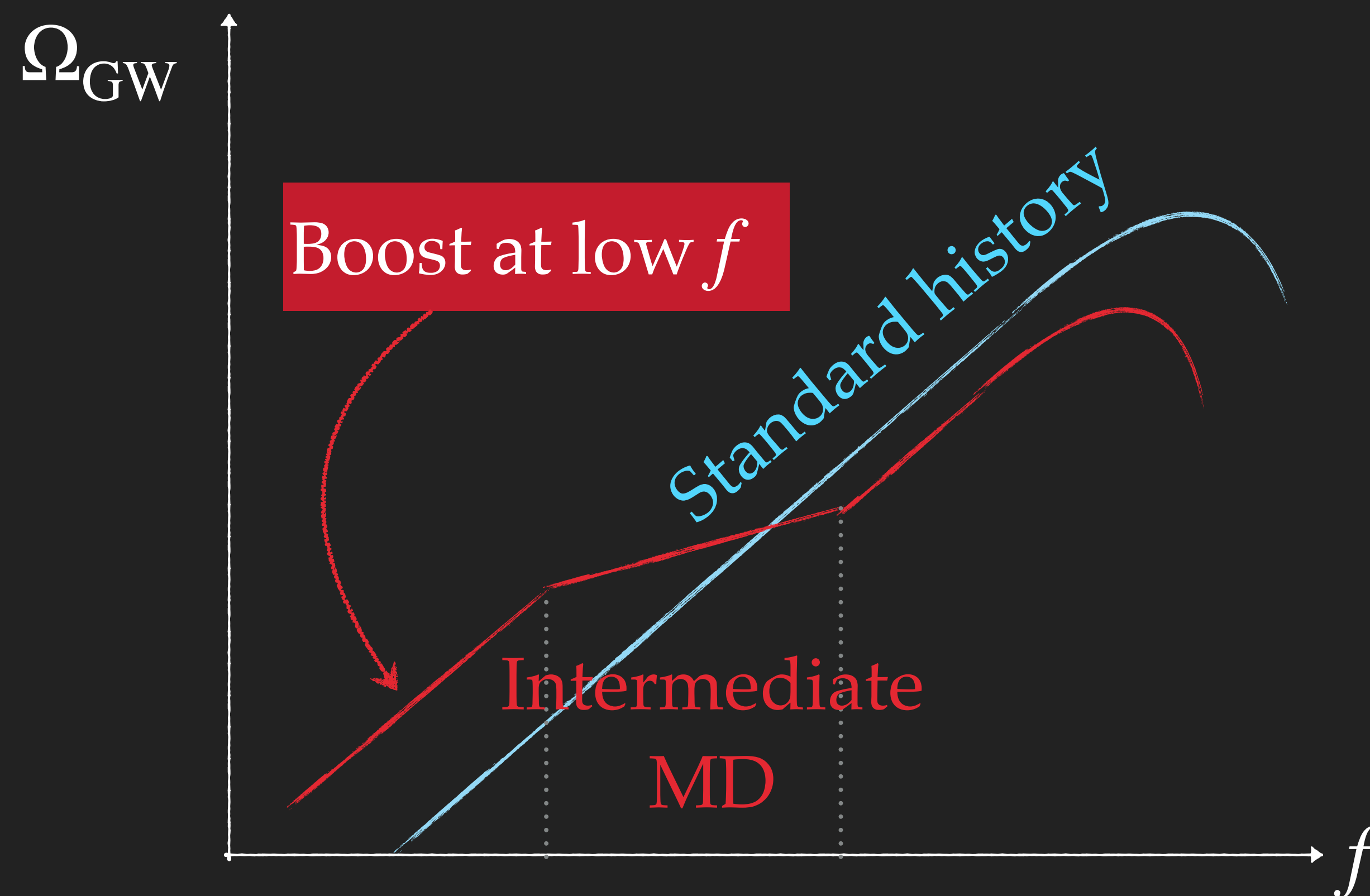


- ▶ Intermediate MD phase  $\rightarrow$  amplify low- $f$  spectrum





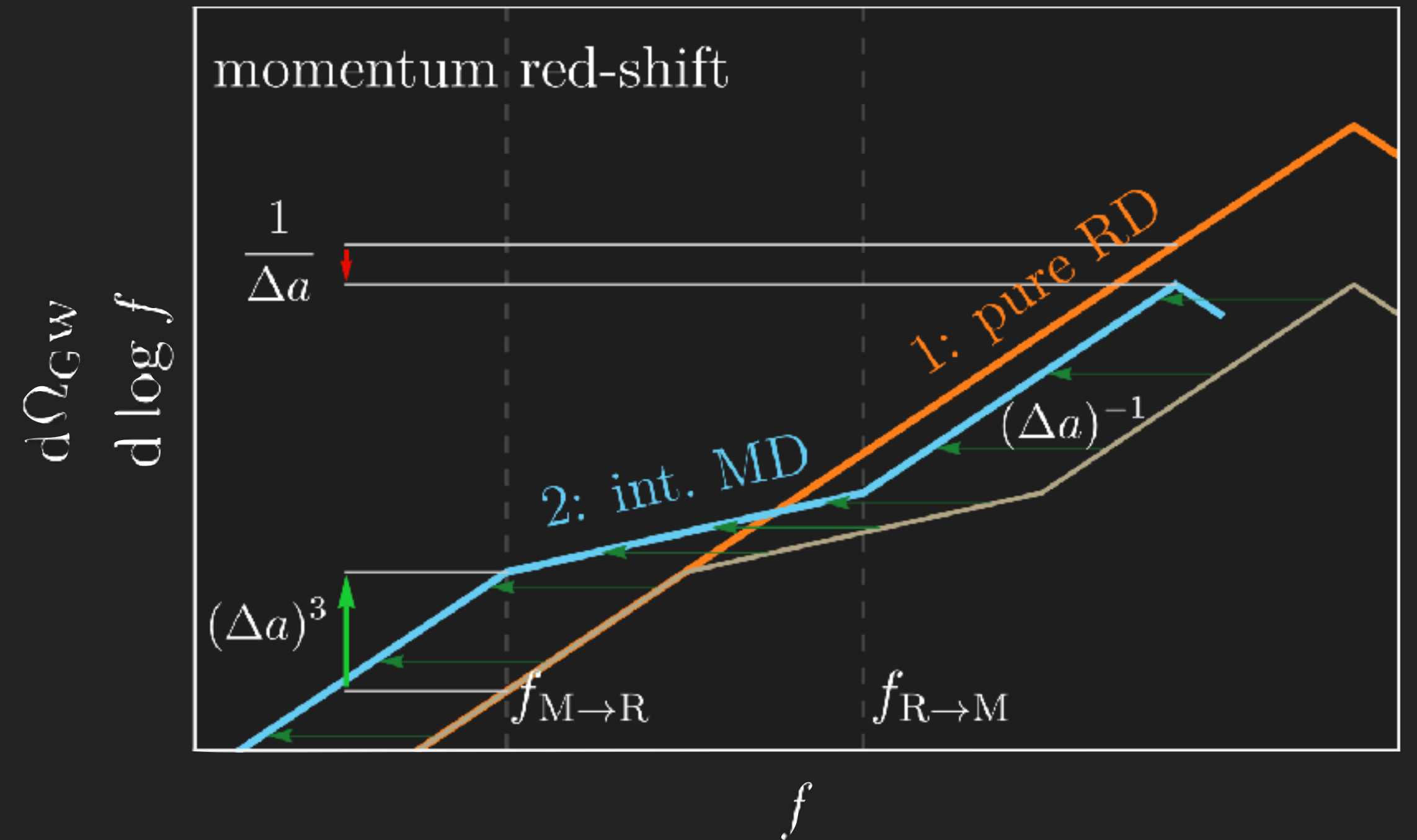
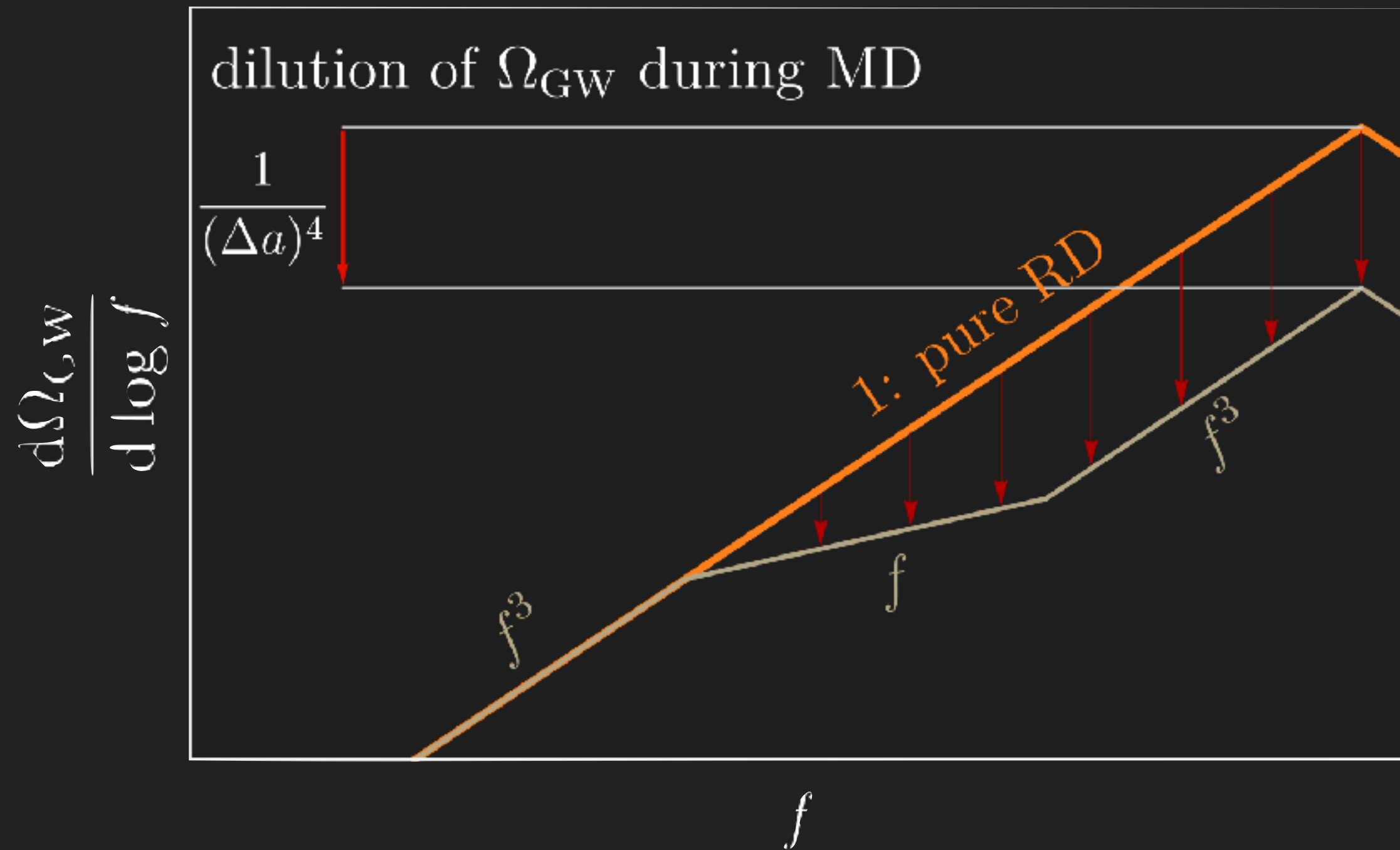
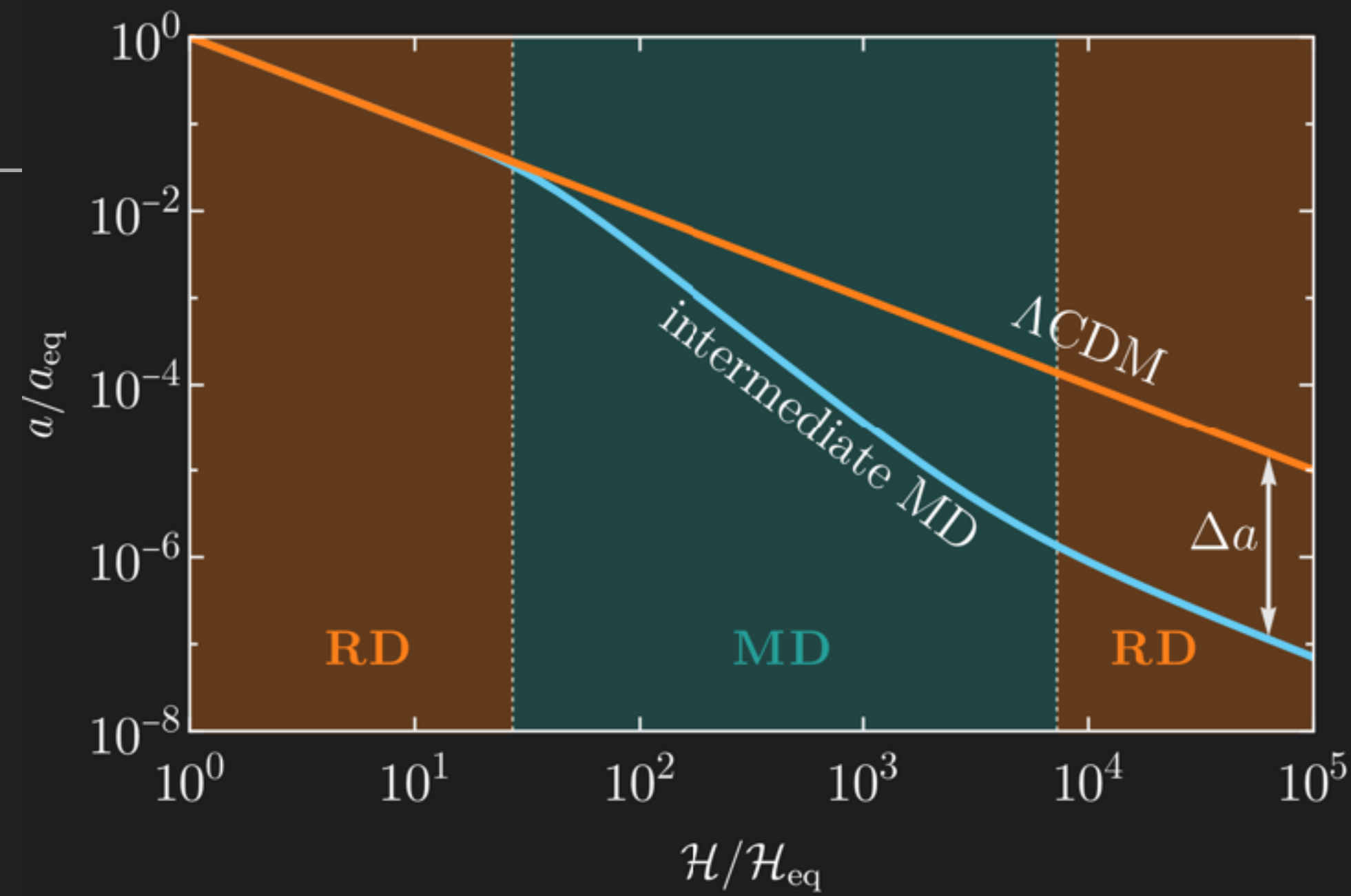
- ▶ Intermediate MD phase  $\rightarrow$  amplify low- $f$  spectrum





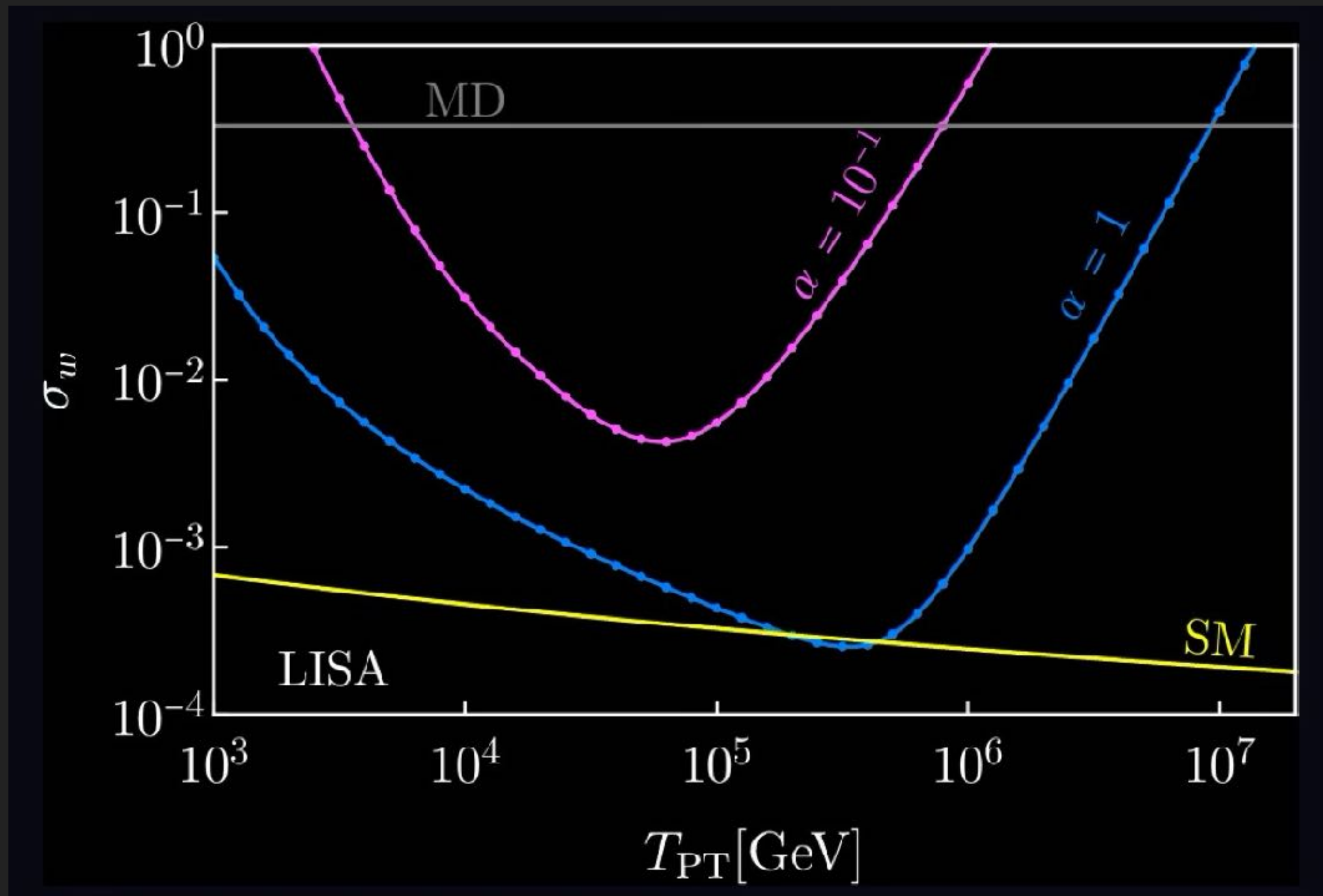
# INTERMEDIATE MD

[‘20 Hook, Marques-Tavares, DR] 28

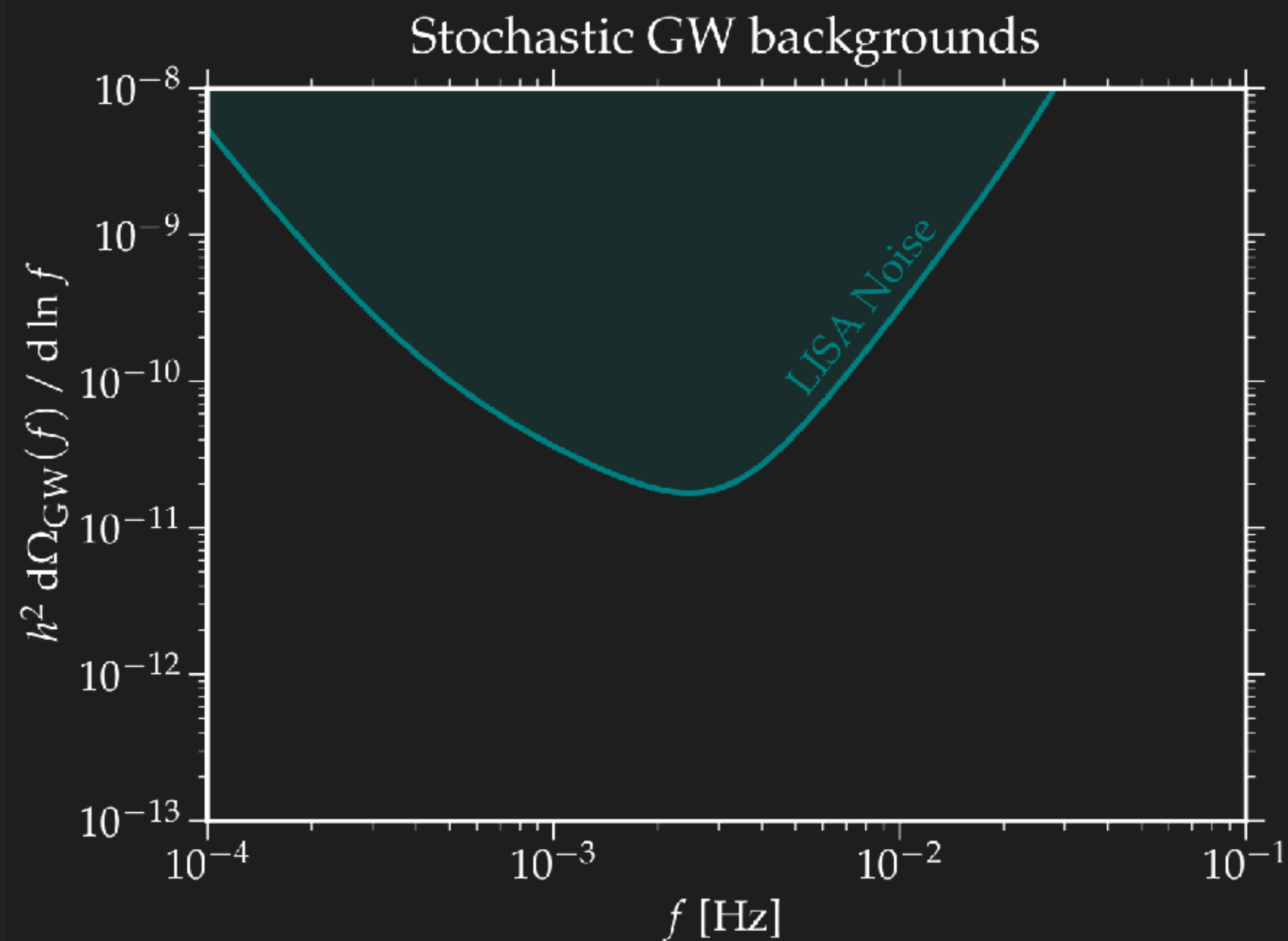




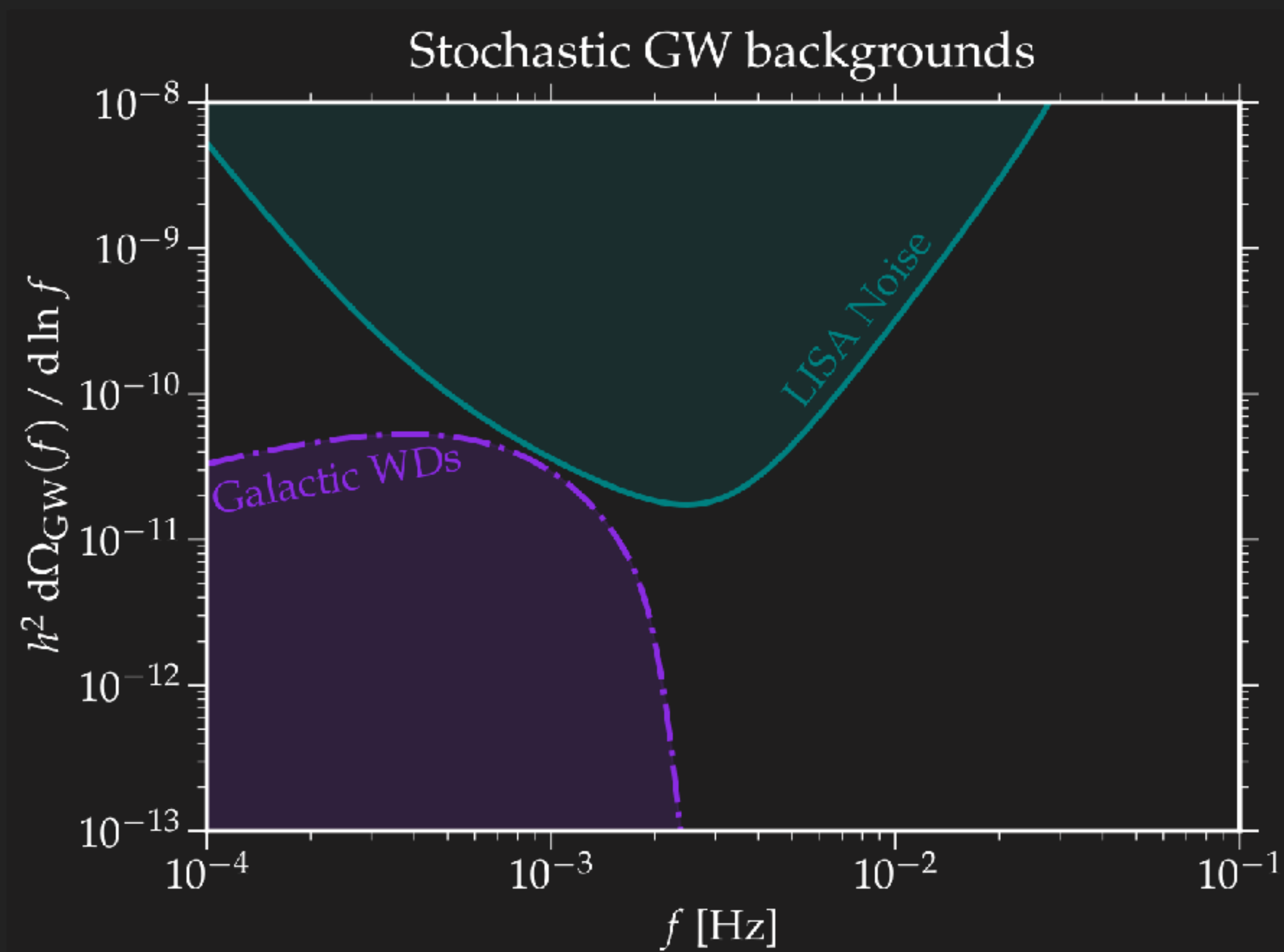
[ '22 Brzeminski, Hook, Marques-Tavares ]



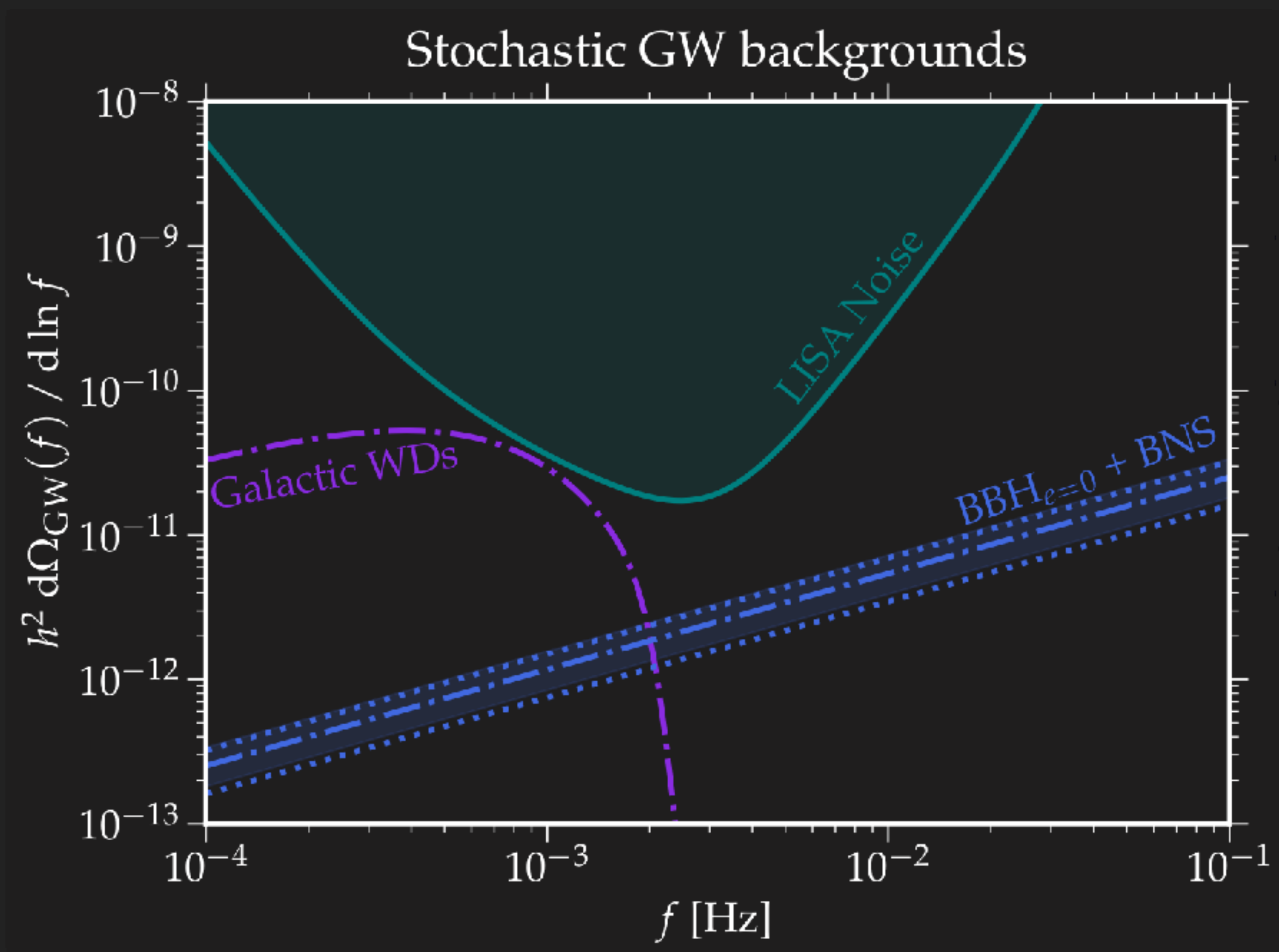




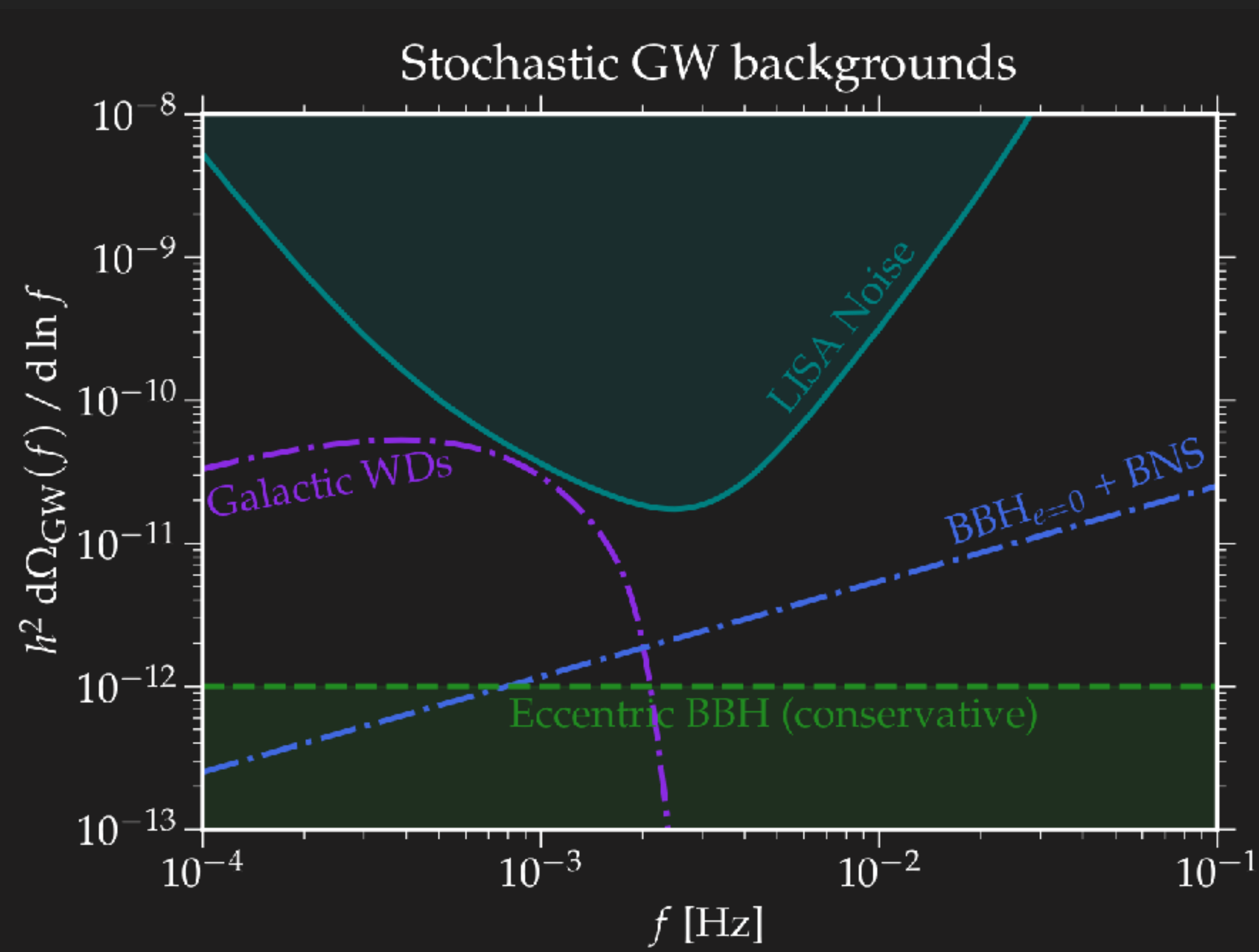




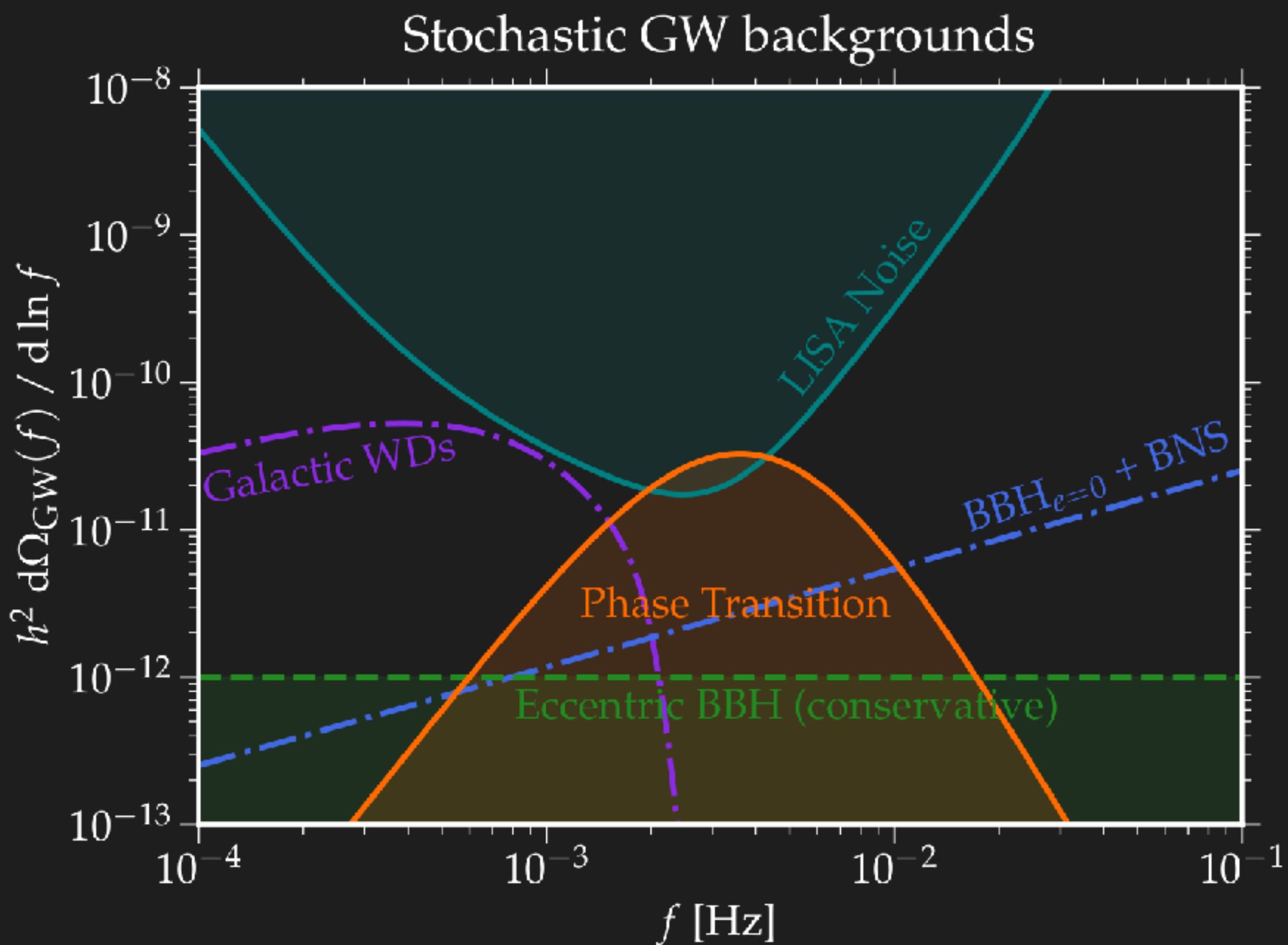






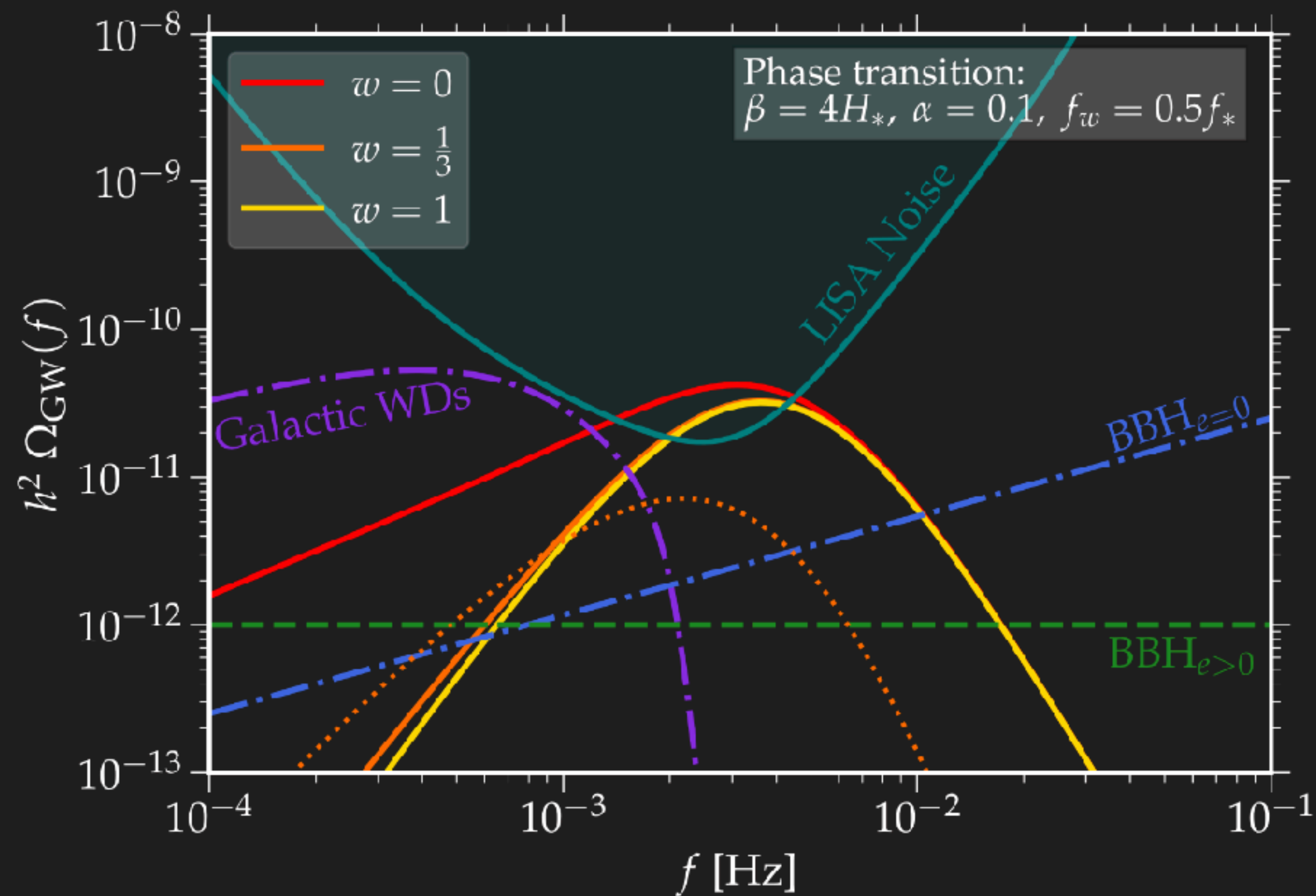






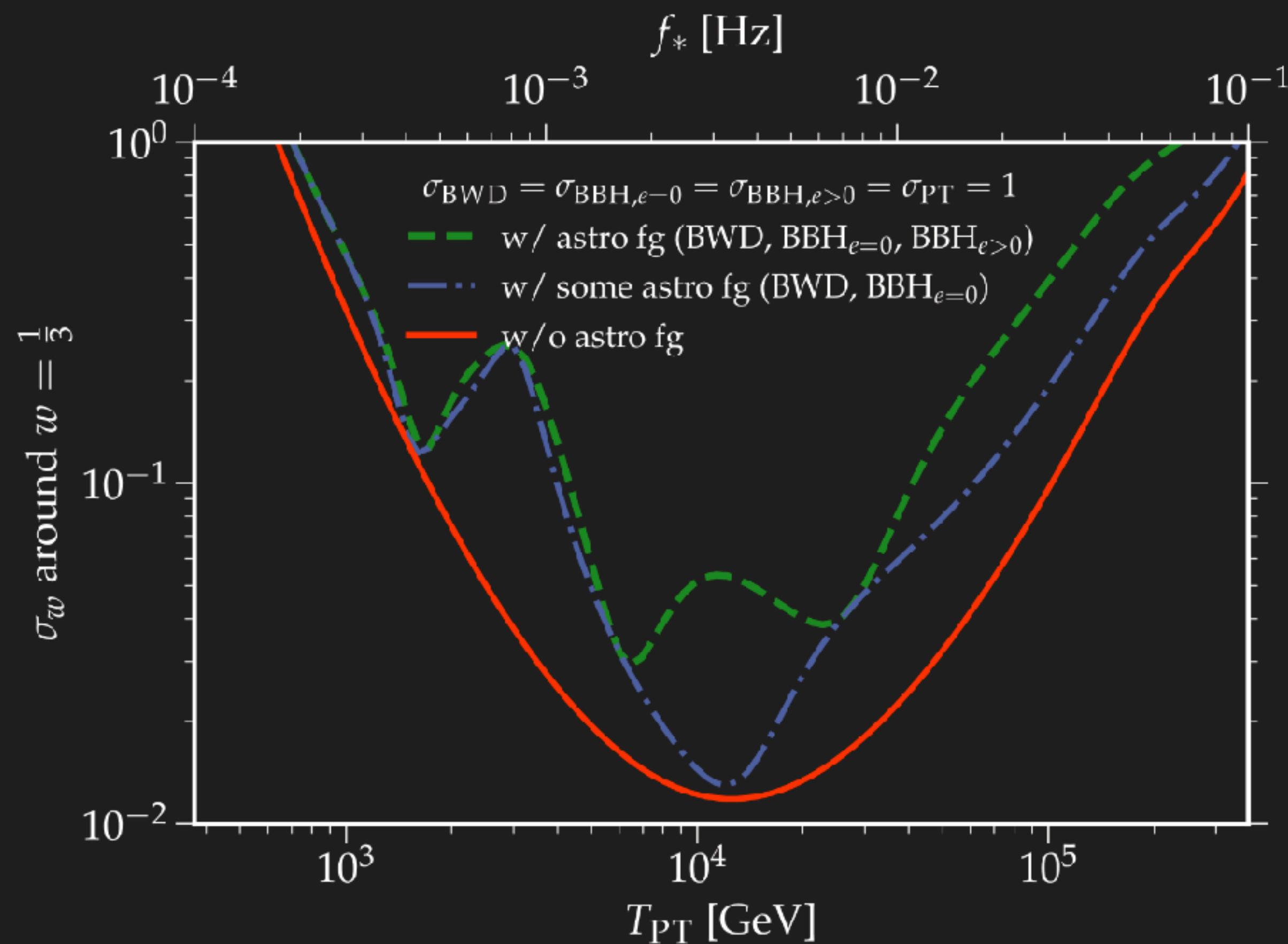
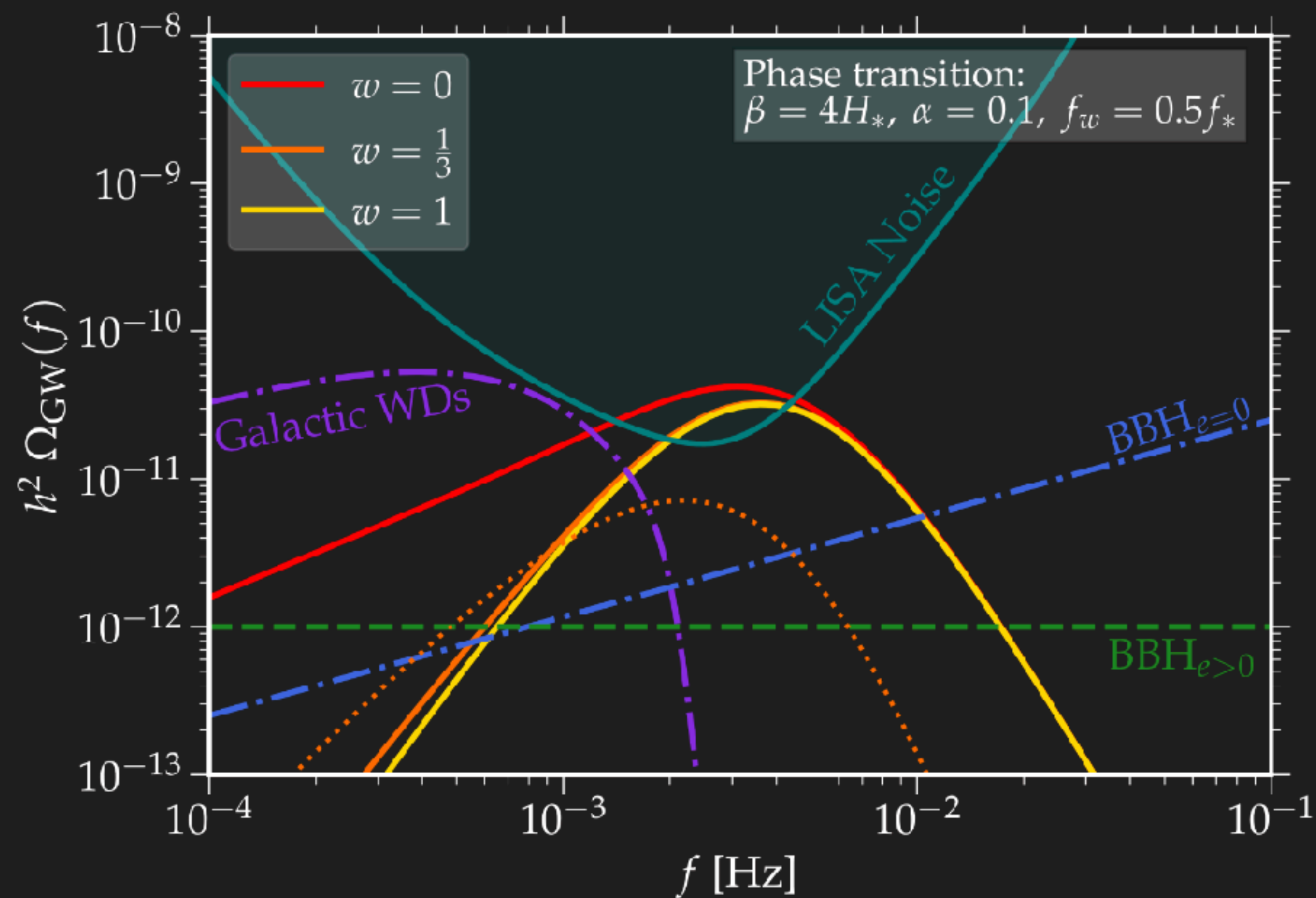


Variation of  $w$  after PT



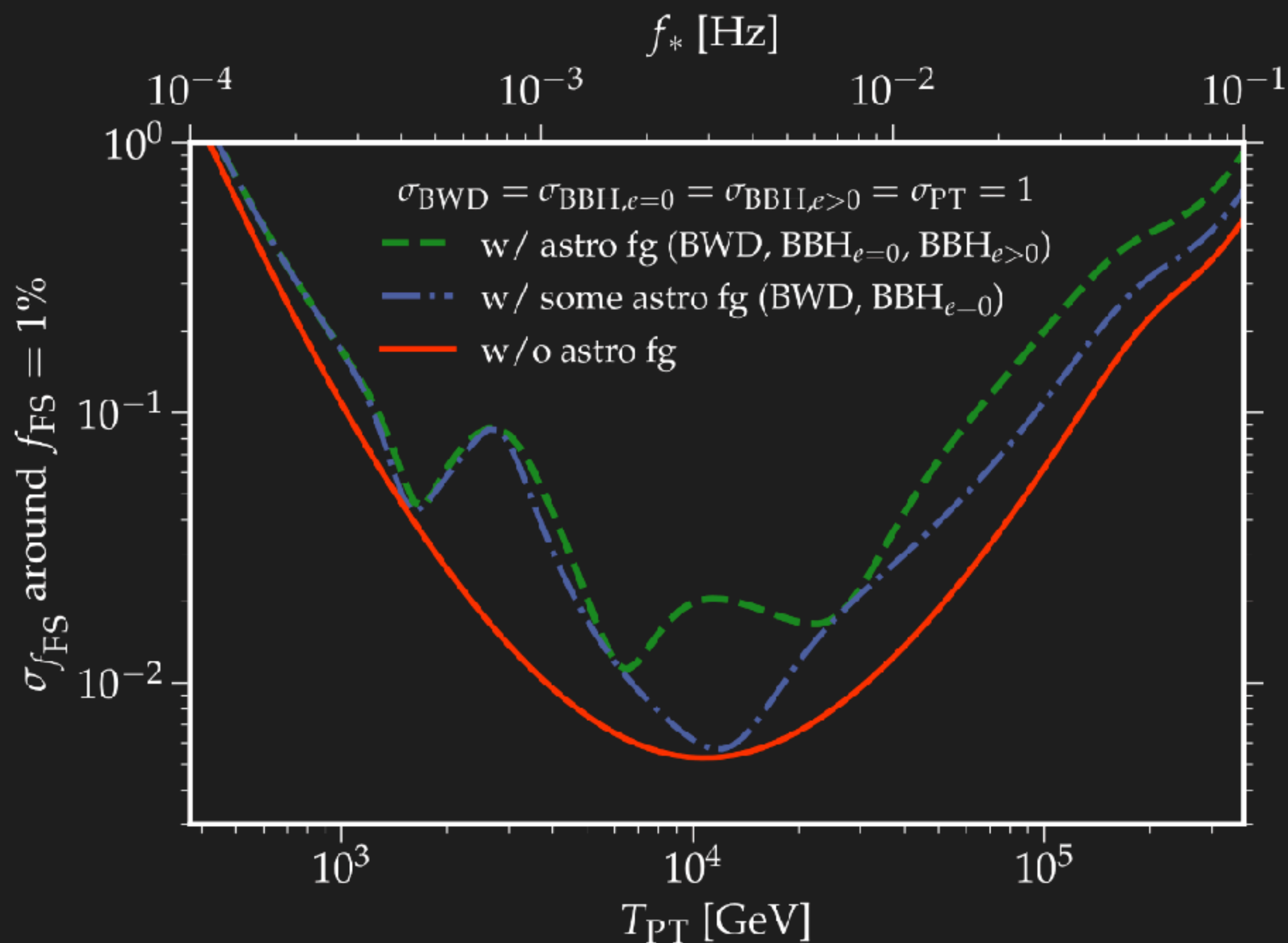
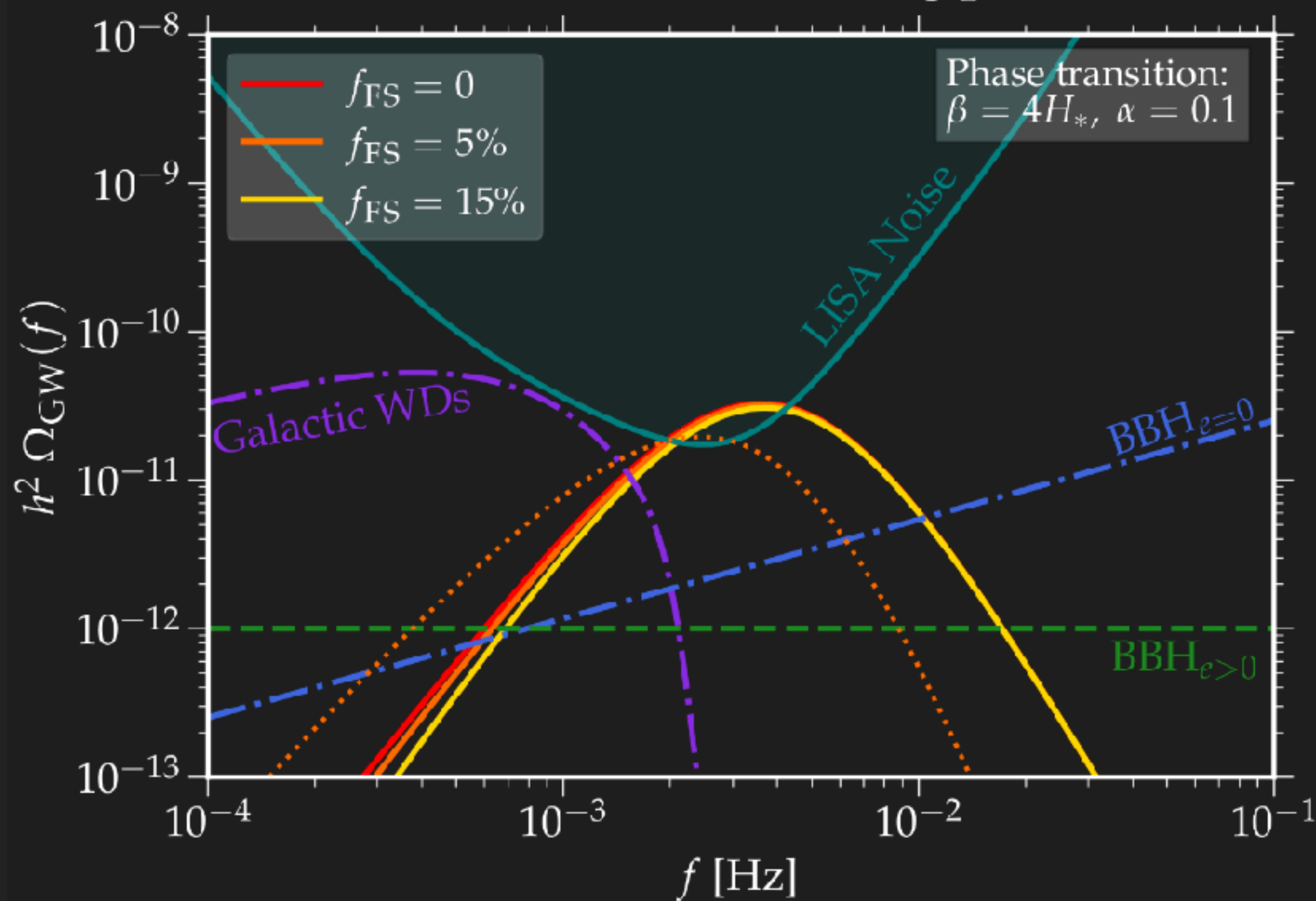


Variation of  $w$  after PT





Presence of free-streaming particles





- **Standard Model:** free-streaming *neutrinos* ( $T \lesssim 2$  MeV).  
 Primordial GWs:  $h' = 0$  until horizon-entry  $\Rightarrow$  frequency-ind. suppression.

$$h'' + 2\mathcal{H}h' + k^2h = -24f_\nu\mathcal{H}^2 \int_{\tau_0}^{\tau} d\tau' \frac{j_2[k(\tau - \tau')]}{k^2(\tau - \tau')^2} h'(\tau')$$

- **Phase transitions:** with extra FS species, also suppression at *generation*:  
 $h'(\tau_\star) \neq 0$ . [**'20 Hook, Marques-Tavares, DR**]

sub- $\mathcal{H}$

no effect

super- $\mathcal{H}$

$$h'' + \underbrace{2\mathcal{H}h'}_{\text{friction}} + \underbrace{\left( k^2 + \frac{8f_{\text{FS}}}{5}\mathcal{H}^2 \right)}_{\text{mass term}} h = 0$$



## Schematic derivation of Weinberg damping

[Weinberg '04; Watanabe, Komatsu '06]

- 

$$h''_{ij} + 2\mathcal{H}h'_{ij} + k^2 h_{ij} = 4\mathcal{H}^2 \pi_{ij}$$

$$T_{ij} = p g_{ij} + a^2 \pi_{ij}, \quad T_{ij}^{(\nu)} = \frac{1}{\sqrt{-g}} \int \frac{d^3 q}{q^0} q_i q_j F^{(\nu)}(q)$$

- Boltzmann eq.:

$$0 = \frac{dF}{dt} = \frac{\partial F}{\partial \tau} + \frac{dx^i}{dt} \frac{\partial F}{\partial x^i} + \frac{dp^0}{dt} \frac{\partial F}{\partial p^0}$$

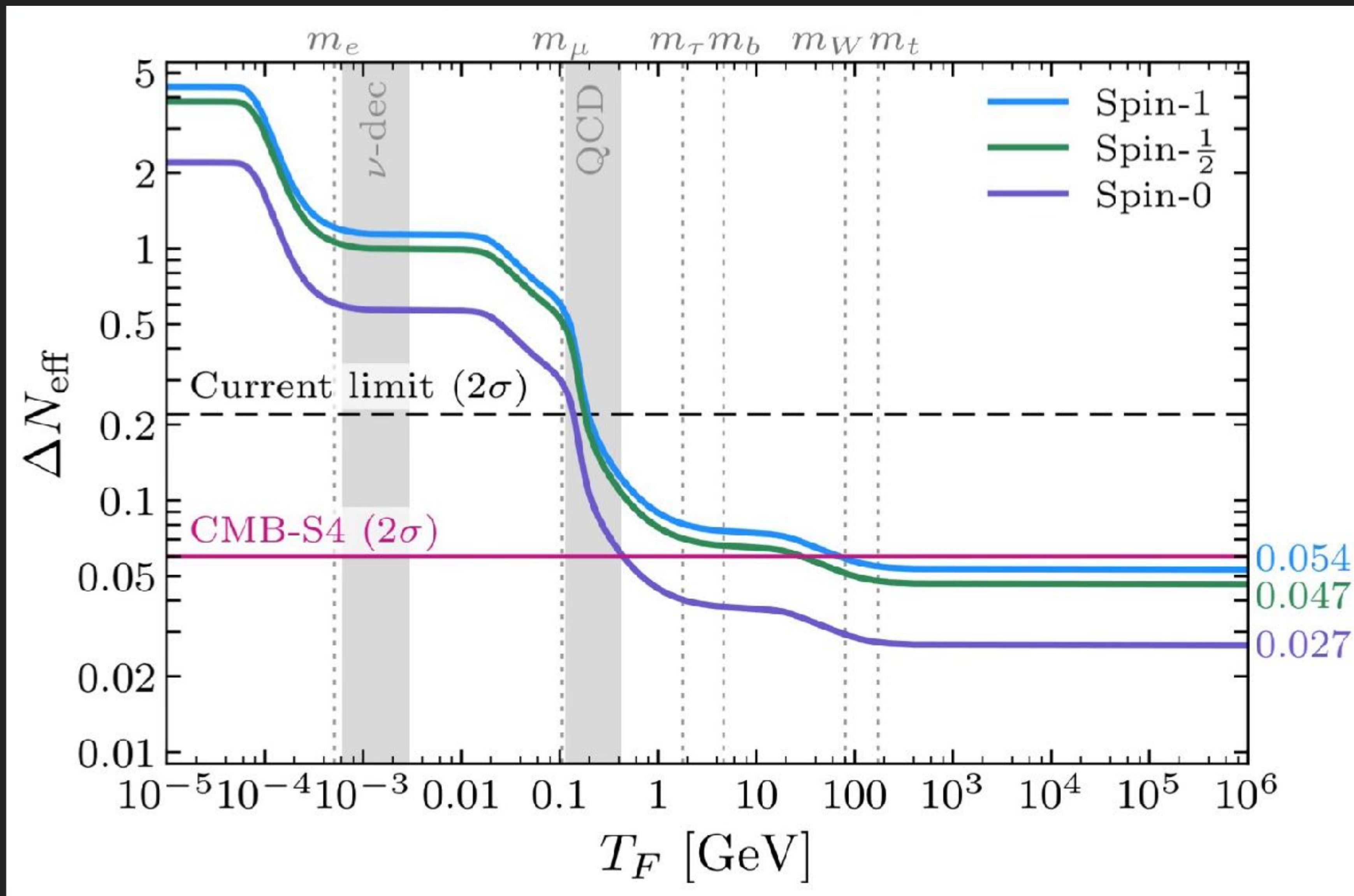
- Geodesic eq.:

$$\frac{dp^\mu}{d\lambda} = -\Gamma^\mu_{\alpha\beta} p^\alpha p^\beta \implies \frac{1}{p^0} \frac{dp^0}{dt} = -H - \frac{1}{2} \frac{\partial h_{ij}}{\partial t} \frac{p^i p^j}{(p^0)^2}$$

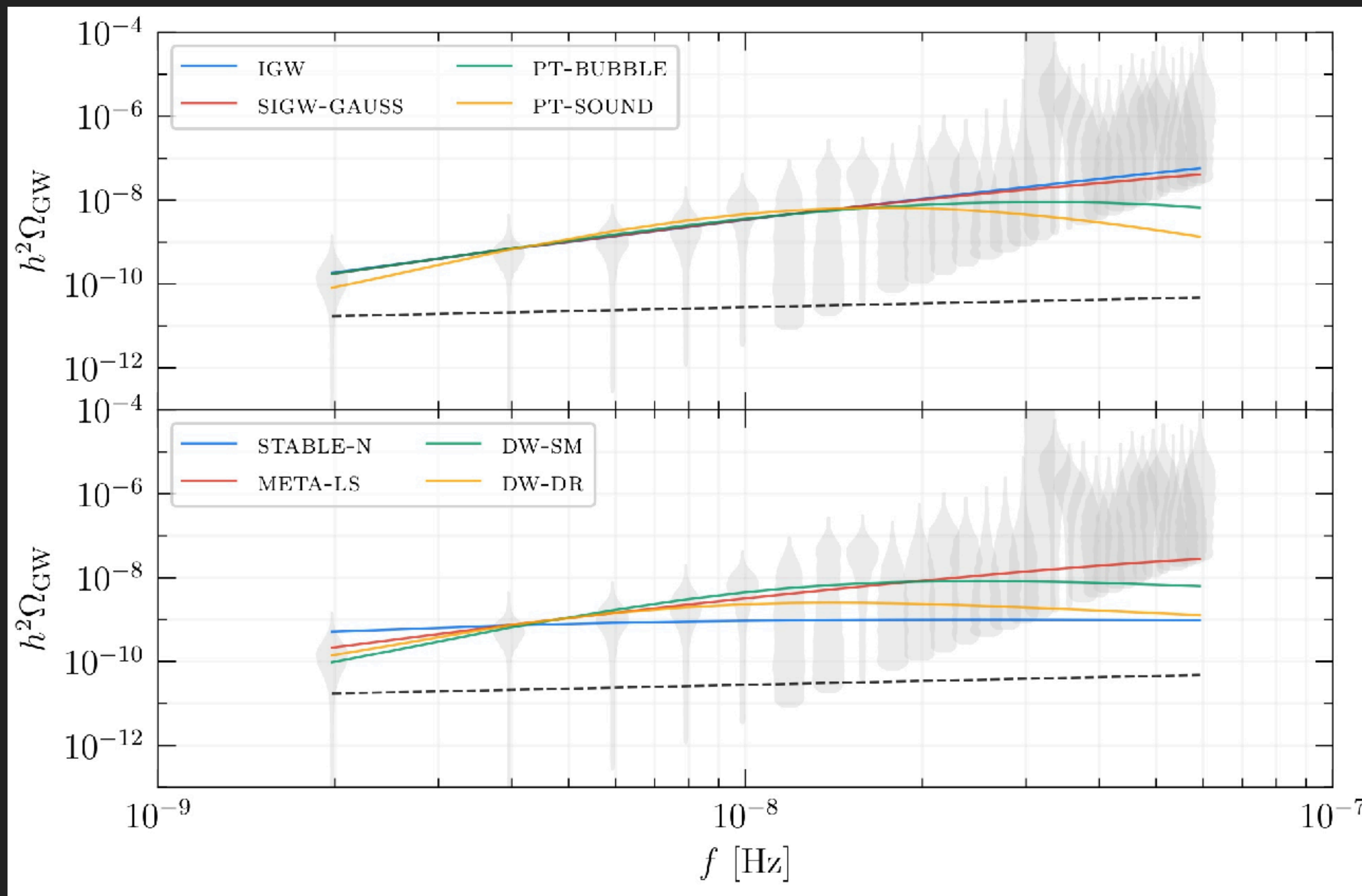
- $\nu$ 's lose (or gain) energy depending on the sign of  $h'$ .

$$h'' + 2\mathcal{H}h' + k^2 h = -24f_\nu \mathcal{H}^2 \int_{\tau_0}^{\tau} d\tau' \frac{j_2[k(\tau - \tau')]}{k^2(\tau - \tau')^2} h'(\tau')$$











[ '18 Espinosa, **DR**, Riotto ]

[ '18 Bartolo, De Luca, Franciolini, Peloso, **DR**, Riotto ]

