

# Higgsed Millicharge particles

## Status report

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# Project overview

- Happy to be a new scientific project board at DMLab with Felix !
- Our general common research axis will be Feebly Interacting Particles
- Ongoing project, started with Patrick Foldenauer on millicharge particles
  - Today is more of a status report than a full presentation ...

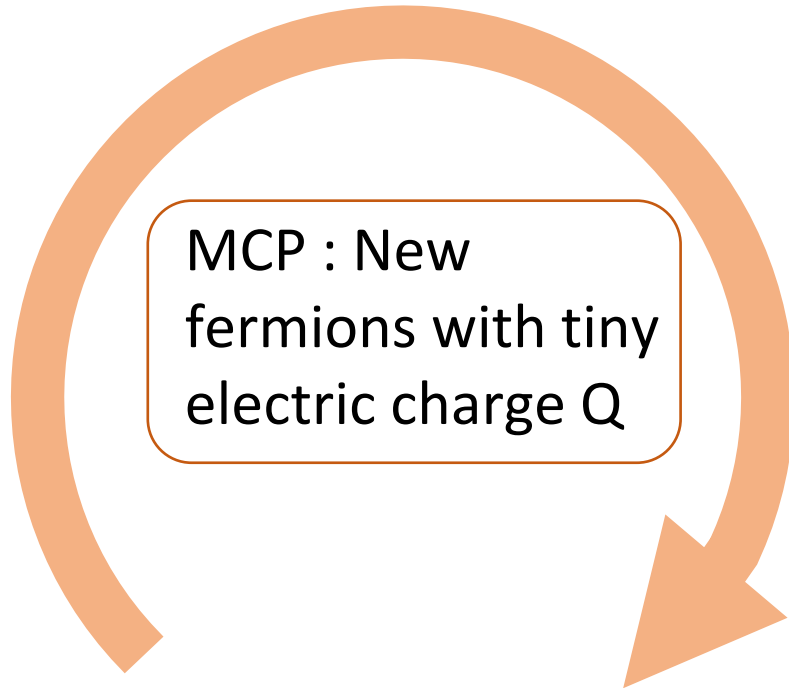
# All millicharge particles are born equal, but...

## Pure Millicharge

$\chi$ , massive Dirac fermion with tiny hypercharge  $Y_\chi$



MCP : New fermions with tiny electric charge  $Q$



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## Kinetic mixing Millicharge

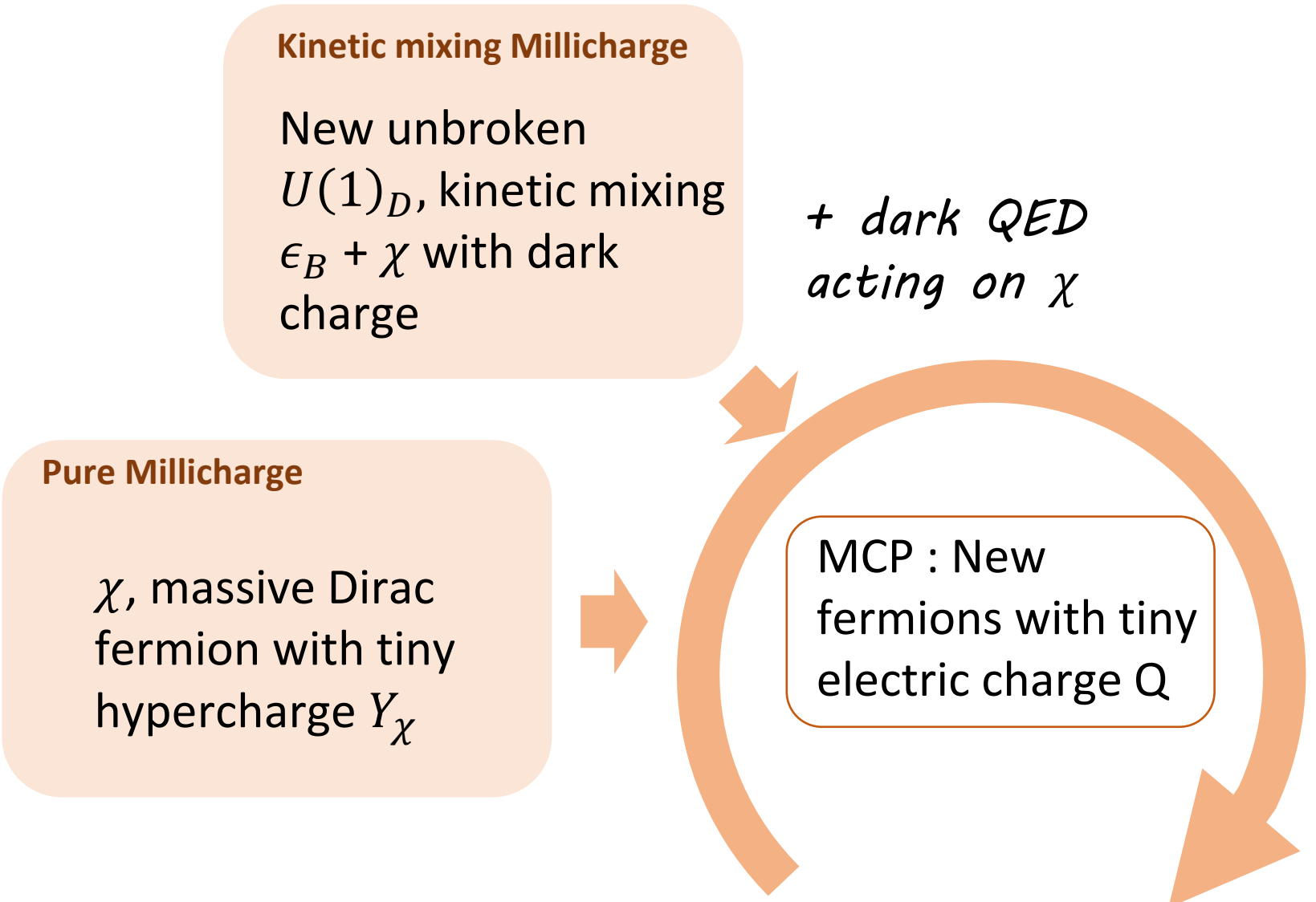
New unbroken  
 $U(1)_D$ , kinetic mixing  
 $\epsilon_B + \chi$  with dark  
charge

*+ dark QED  
acting on  $\chi$*

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## Higgsed Millicharge

New broken  
 $U(1)_D$ , dark Higgs  
with tiny  
hypercharge  $Y_S + \chi$   
with dark charge

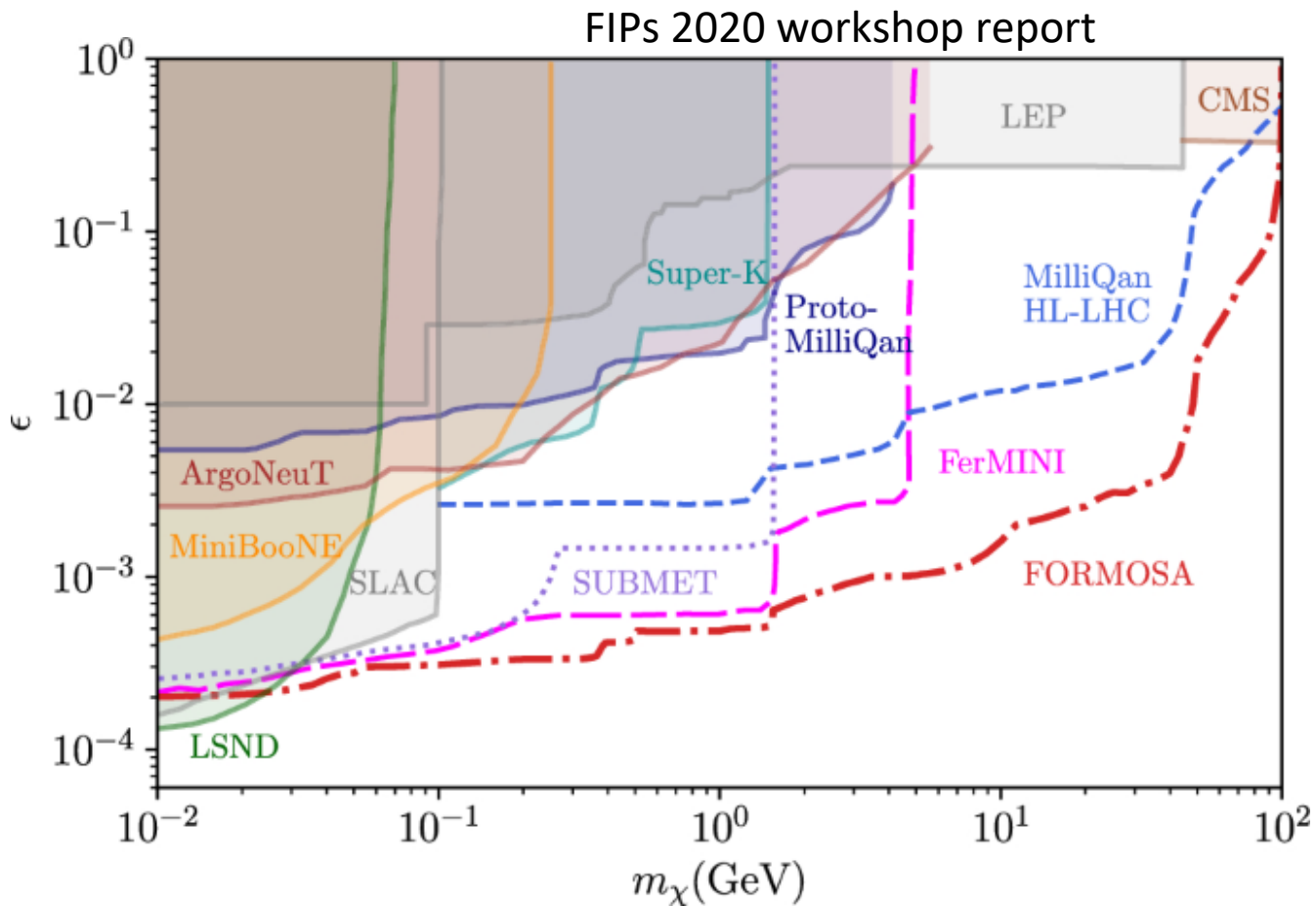
*+ massive dark  
photon  
+ EW effects*

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# Accelerator searches are in full swing

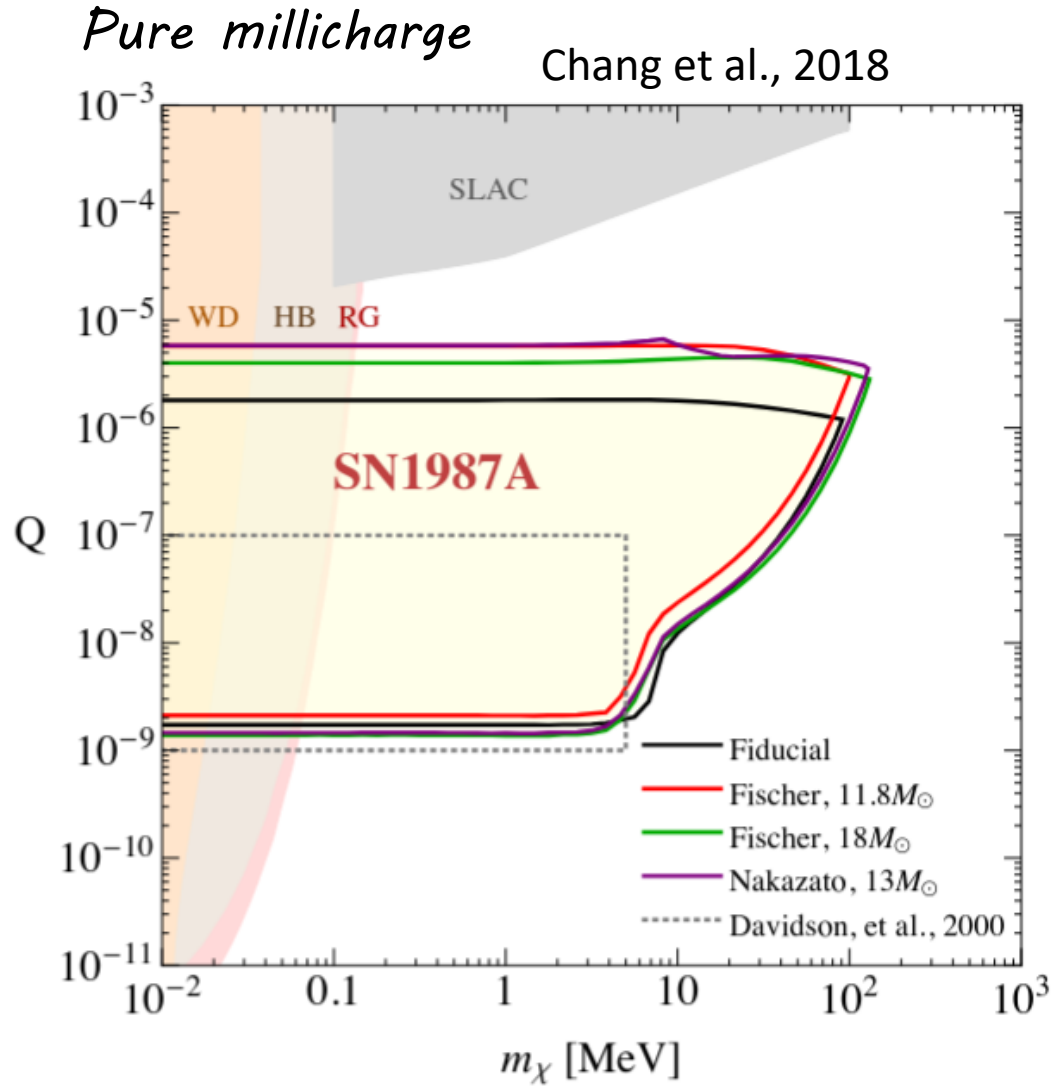


- Accelerator experiments sensitive to millicharges as small as  $10^{-4}$ 
  - much weaker bounds for heavier MCPs
- For couplings in this range, MCPs can be efficiently produced in the early universe and stellar cores

- Are these parameter regions consistent with astrophysics/cosmology ?

# Astrohpysical and cosmological constraints on MCPs

# MeV-range : SN1987 cooling



- For  $Q > 10^{-5}$  MCPs are efficiently trapped in the SN core with a mean free path smaller than neutrinos
- Negligible contribution to cooling and heat transport
- Stellar cooling constraints only apply to MCPs with mass below 100 keV

*!! Highly dependent on dark sector structures !!*

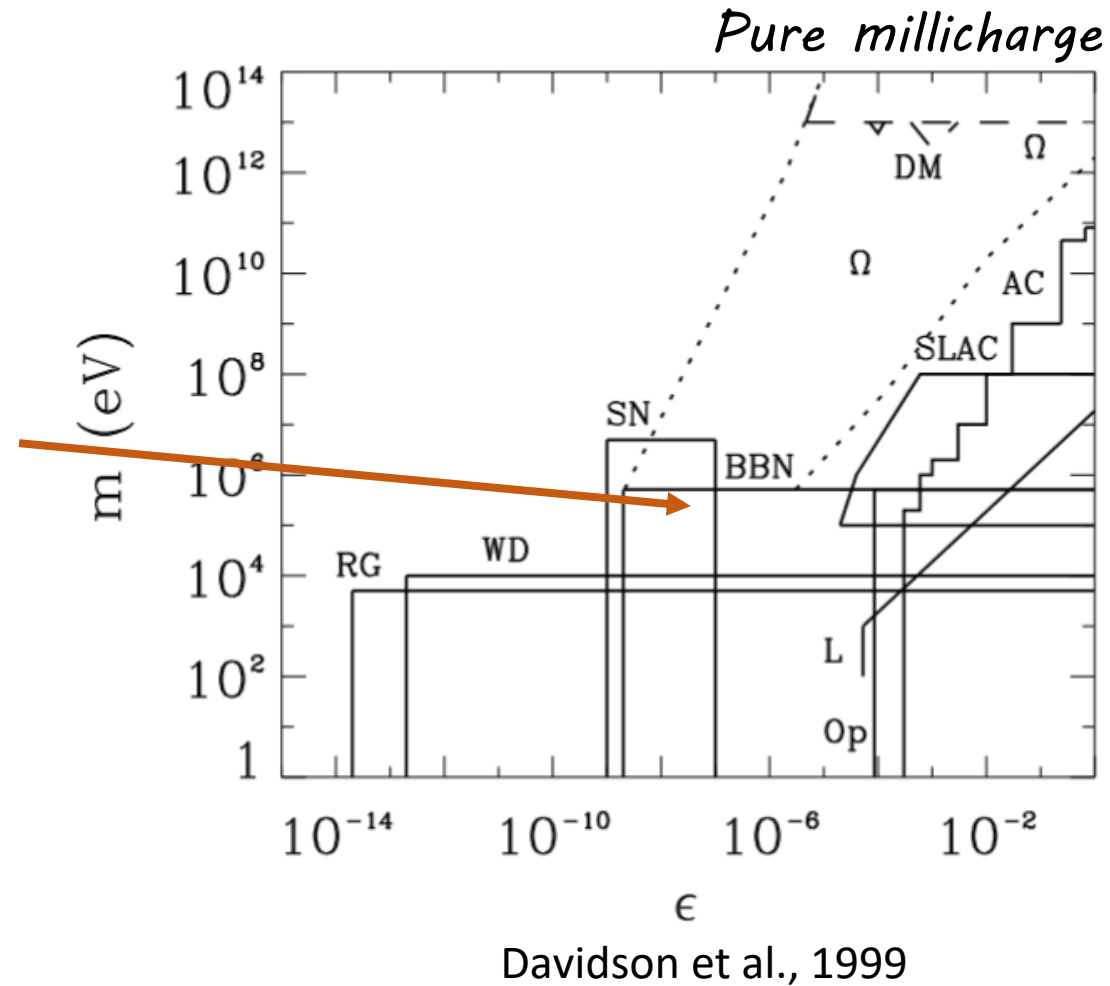


# Relativistic degrees of freedom

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  - Significant contribution to the total energy density (and hence the expansion rate)
  - Robustly excluded by BBN measurements



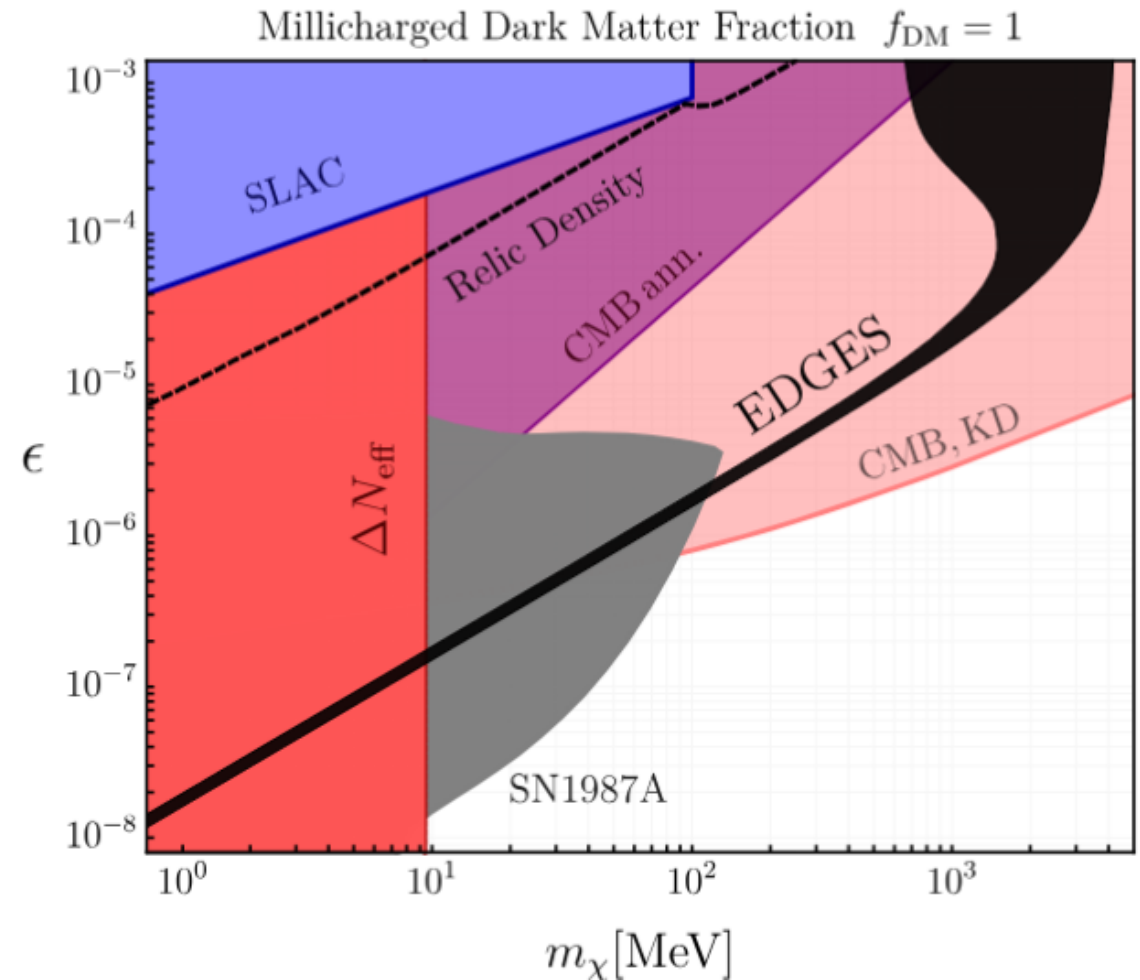
# MCP relic density : freeze-out

- Since they are stable and thermalised, can they be DM ?
  - For heavier MCP masses, the particles become Boltzmann-suppressed and freeze out
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*Pure millicharge*

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  - For heavier MCP masses, the particles become Boltzmann-suppressed and freeze out
  - For  $Q \sim 10^{-4} - 10^{-3}$  MCPs could be all of DM
- **Problem: such scenario leads to strong CMB constraints**
  - Dominant limit comes from requiring DM decoupling at the time of baryon-photon recombination



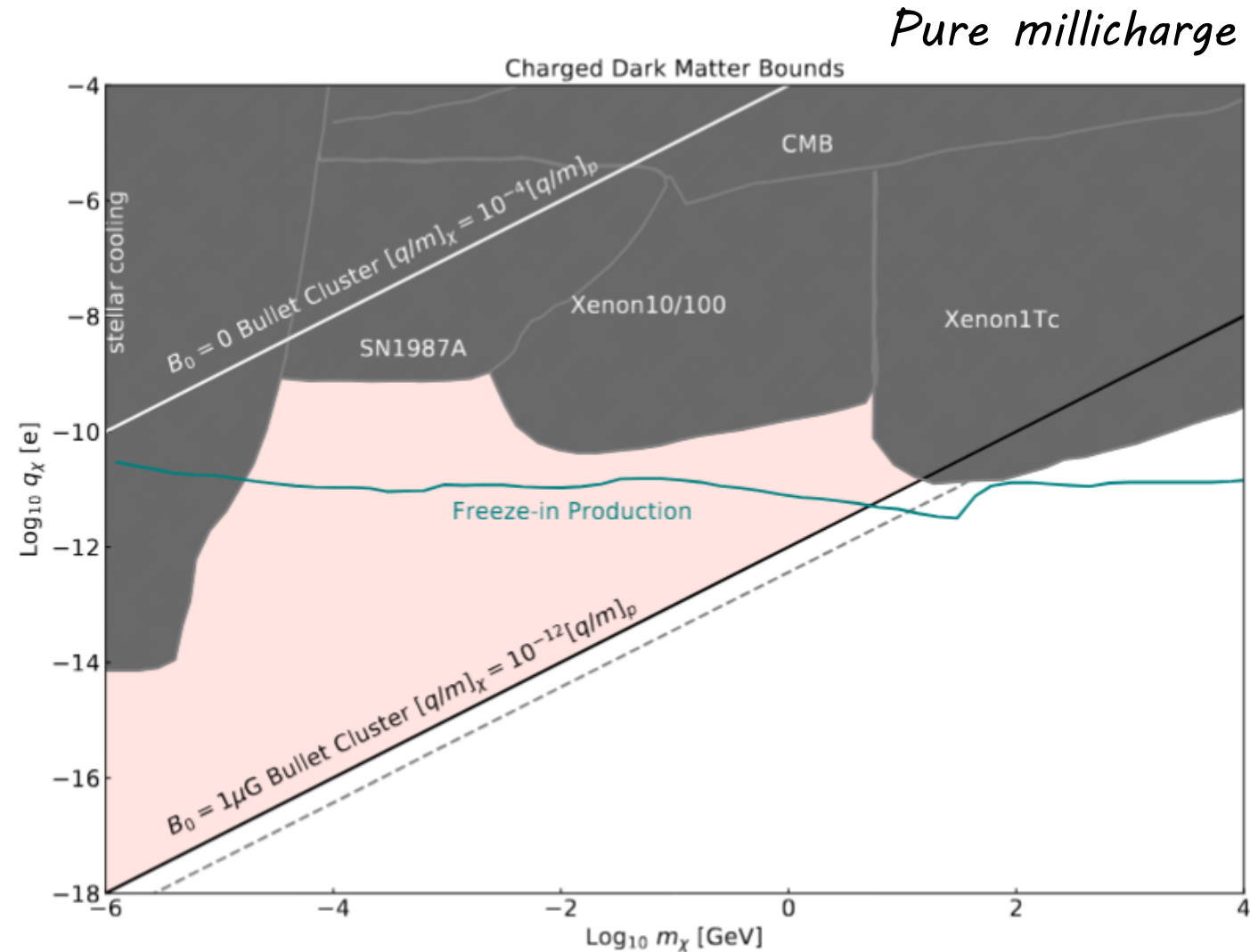
# Can MCPs actually be dark matter ?

- Various papers have pointed out that millicharged dark matter experiences plasma instabilities
  - Mostly the product of having a long-range interaction available ...
- Extremely strong constraints on the charge of dark matter from the Bullet Cluster

Cruz & McQinn, arXiv:2202.12464

Lasenby, arXiv:2007.00667

Medvedev & Loeb, arXiv:2406.15750



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# Suppressing the MCP abundance (I)

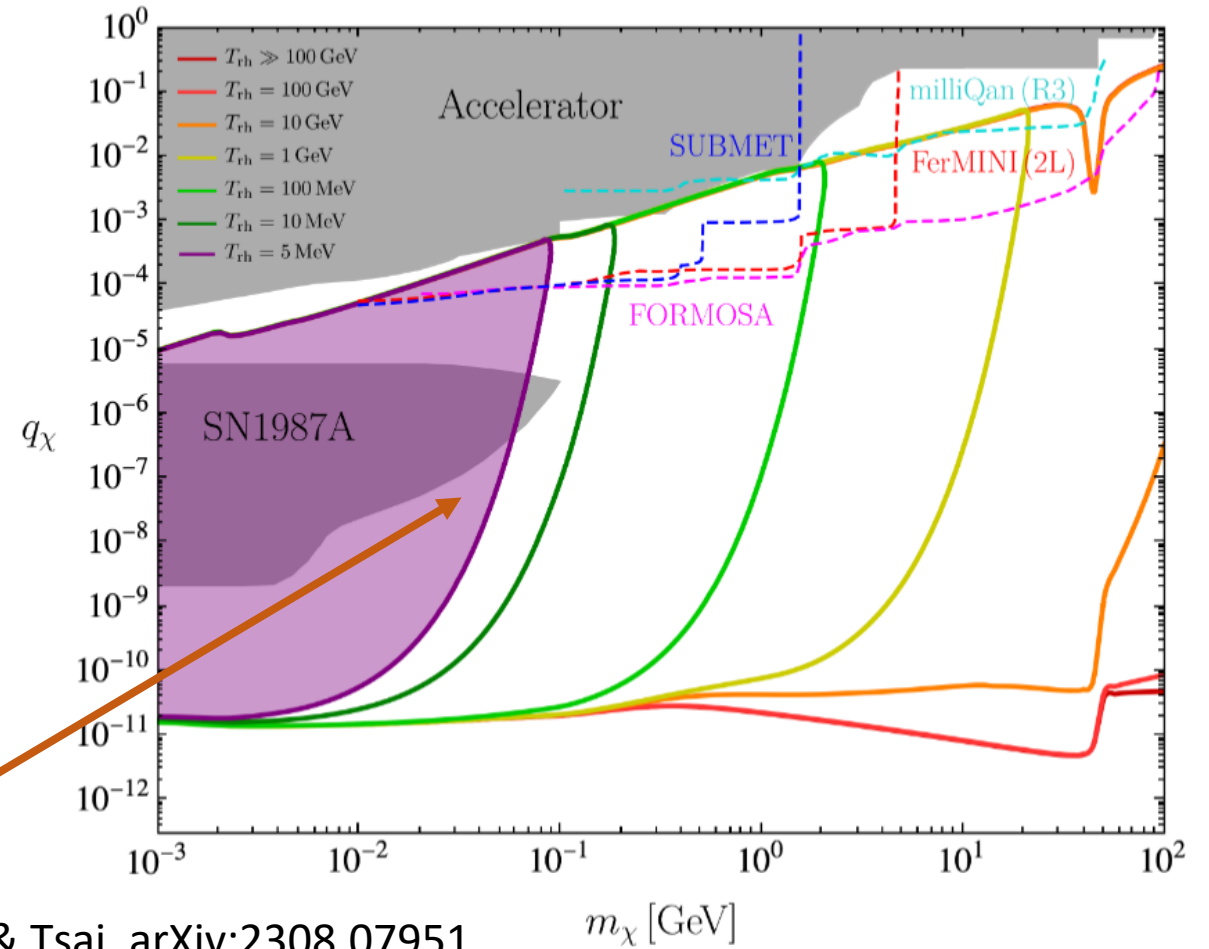
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# Suppressing the MCP abundance (I)

- Clearly, we need ensure a low abundance of MCPs to open a collider-reachable parameter space
- Possible solution: Modification of standard cosmology
- Example: Low reheating temperature  
→ Exponential suppression of MCP abundance if their mass is greater than  $T_{RH}$

*Also include CMB limits from earlier on*

Overproduction Bounds for “Pure” mCP



# Suppressing the MCP abundance (II)

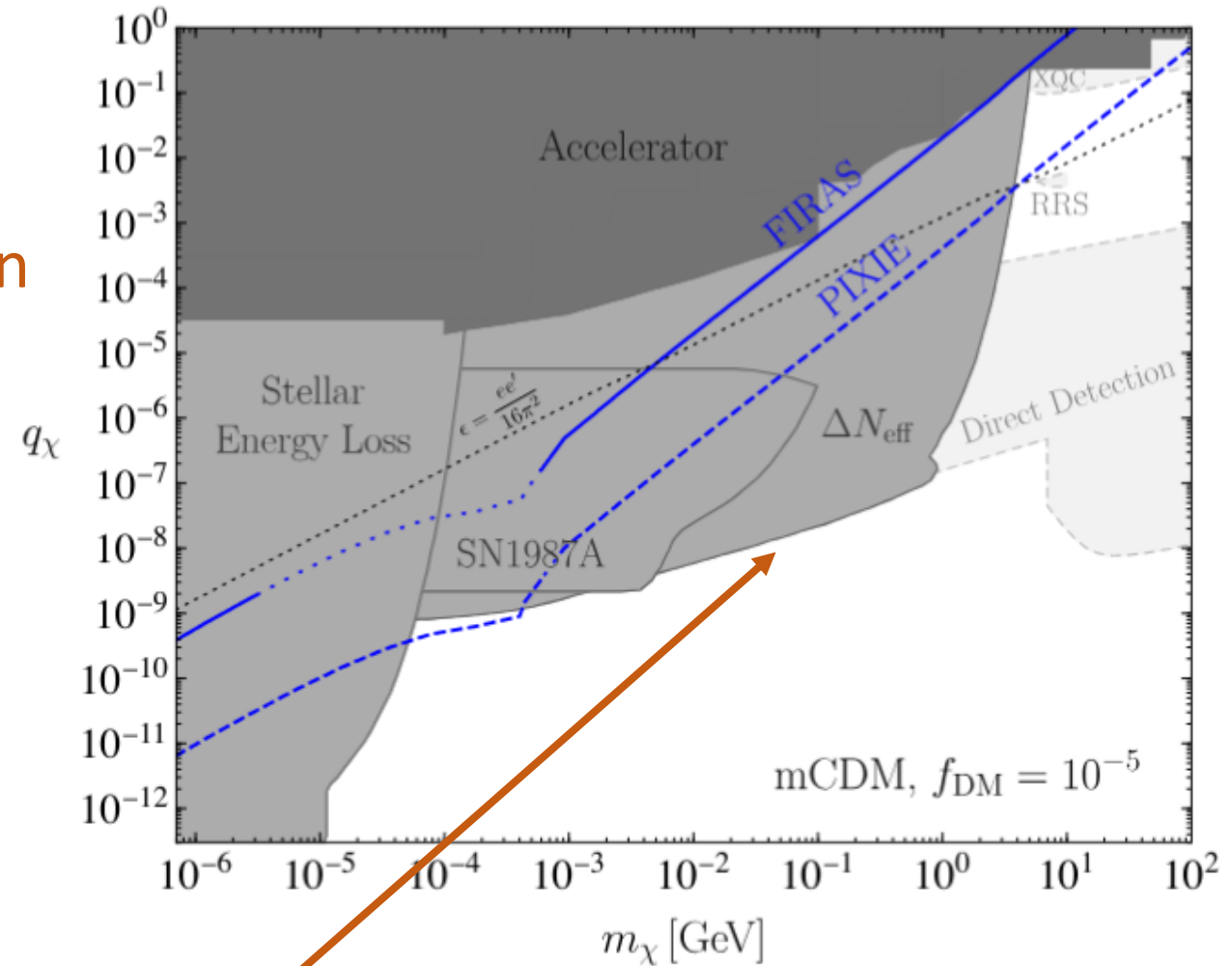
- Alternative avenue: use additional dark sector states (dark photon, dark Higgs boson) as additional annihilation channels to deplete MCP abundance



# Suppressing the MCP abundance (II)

*Kinetic mixing millicharge*

- Alternative avenue: **use additional dark sector states (dark photon, dark Higgs boson) as additional annihilation channels to deplete MCP abundance**
- Challenge: Lighter states need to decay into SM to avoid reappearance of  $\Delta N_{\text{eff}}$  constraint
  - Prediction of additional light (and long-lived?) particles
- **Relevant parameter space opens for accelerator searches !**



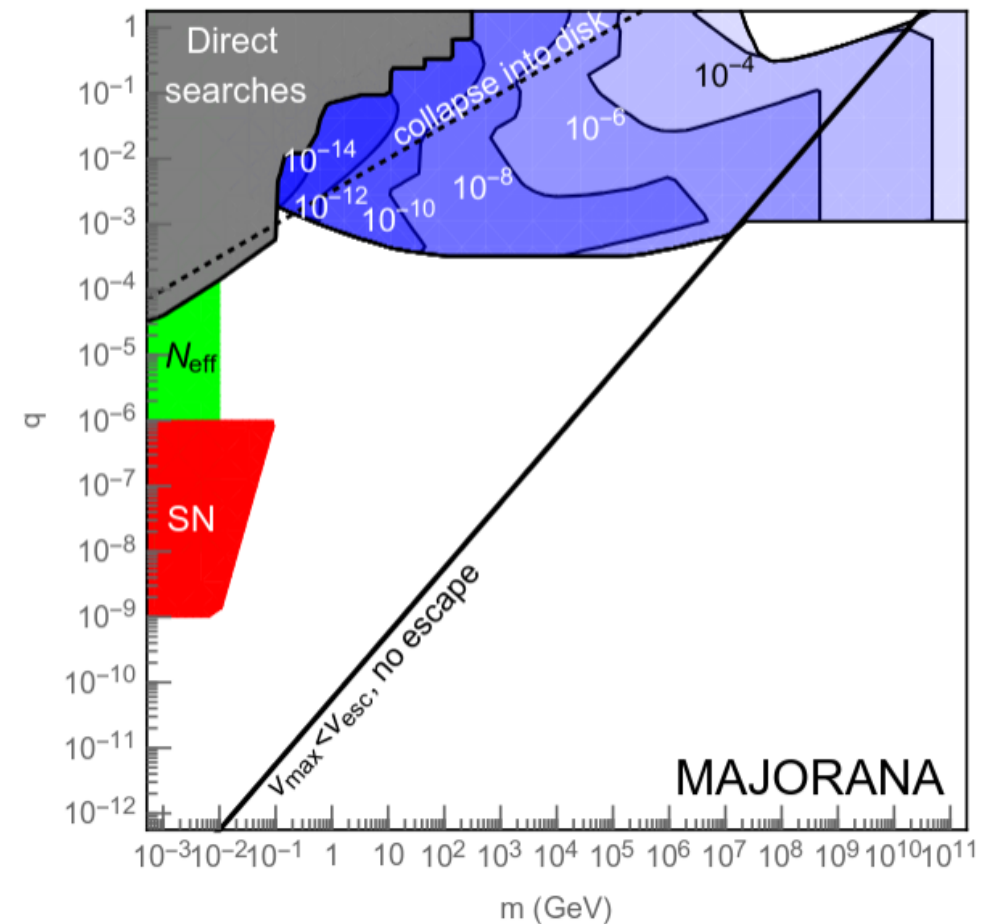
*Model with massless dark photon*

Berlin et al., arXiv:2211.05139

# Direct detection constraints?

- The evolution of a cosmological population of MCPs depends on many factors
  - Diffusion in magnetic fields
  - Expulsion through supernova shocks
  - Collapse into disk
- Rock overburden may attenuate the flux of MCPs in experiments
- In non-minimal models, a large self-interaction component may arise

→ Calculation of direct detection constraints highly non-trivial



Building millicharge particles  
dark sectors

# Kinetically mixed millicharge

- We use a new *unbroken*  $U(1)_D$  symmetry along with a “kinetic mixing” term
  - Could be loop generated, or seen as a convenient way of expressing that the  $U(1)_D$  current has a small hypercharge component

$$\mathcal{L}_{A'} = -\frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} - \boxed{\frac{1}{2} \frac{\varepsilon}{\cos \theta_w} B_{\mu\nu} F'^{\mu\nu}}$$

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*Kinetic mixing term*

- We then re-define the gauge bosons to get a properly normalised fields
  - We actually have a  $SO(2)$  redefinition freedom afterwards since both the photon and the dark photon are massless

$$\begin{pmatrix} A_\mu \\ Z_\mu \\ X_\mu \end{pmatrix} \rightarrow \begin{pmatrix} 1 & 0 & -\frac{c\theta \epsilon_B}{\sqrt{1-\epsilon_B^2}} \\ 0 & 1 & \frac{s\theta \epsilon_B}{\sqrt{1-\epsilon_B^2}} \\ 0 & 0 & \frac{1}{\sqrt{1-\epsilon_B^2}} \end{pmatrix} \begin{pmatrix} A_\mu \\ Z_\mu \\ X_\mu \end{pmatrix}$$

# Two (equivalent) choices

- We are thus left with two equivalent descriptions of the interactions

The standard « Dark photon » convention (which is similar to the massive case)

$$\mathcal{L} \subset e j_{\mu}^{em} A_{\mu} + (g_X j_{\mu}^X - c_W \epsilon_B j_{\mu}^{em}) X^{\mu}$$

The usual « millicharge » convention

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- In both cases the **dark photon does NOT decouple from the physics of the problem since it couples directly to the millicharge particles.**

→ Thus the millicharge relics created using this mechanism have in any case a extra « dark QED » interaction that may dramatically affect their evolution

# A massive dark photon: higgsed millicharge

- If we give a mass to the dark photon, we cannot rotate to the « millicharge convention » and we have a zero-charge light DM model
- Idea: we can use a new dark Higgs singlet, with dark charge  $q_D$   
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Kinetic mixing term

Dark Higgs potential

*(suppose for now the dark charges chosen to avoid dark Yukawa coupling)*

# Precision constraints

- A new Higgs scalar singlet with non-zero hypercharge  $q_Y$  is an open-door to EW precision physics problems ...
  - We expect to have to take  $v_S \ll v_{EW}$  et  $q_Y \ll 1$
  - More complex Higgs sectors (or dark-charged SM-like Higgs doublets) may have more relaxed constraints

*As a first estimate,  
we can take the shift  
in the Z boson mass*

$$\rho_0 - 1 = \frac{m_W^2}{m_Z^2 c_W^2} - 1 \sim -6q_Y^2 \frac{v_S^2}{v_{EW}^2} \longrightarrow \frac{m_X q_Y}{m_Z q_D} \leq O(0.005)$$

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- We thus need either a light dark photon or a small ratio  $\delta = \frac{q_Y}{q_D}$ 
  - One still needs to tune the « millicharge ratio » :  $\delta$
- Millicharge particles are a direct prediction of modified EW-breaking sector

# The resulting currents

- In the mass basis we finally obtain the gauge interactions of the form

$$\mathcal{L}_{\text{int}} = -(A^\mu, Z^\mu, Z'^\mu) J_\mu^{\text{phys}}$$

- If the dark photon mass is much smaller than the Z mass we get :

*Millicharge generation*

$$J_\mu^{\text{phys}} = \left( \begin{array}{c} e(j_\mu^{\text{em}} - \delta j_\mu^x) \\ s_\theta \epsilon_B g_x j_\mu^x + g_z j_\mu^Z \\ g_x j_\mu^x - \left( \frac{e^2 \delta}{g_x} - c_\theta e \epsilon_B \right) j_\mu^{\text{em}} \end{array} \right) \cdot$$

*Dark sector interactions* ←

→ *Kinetic mixing*

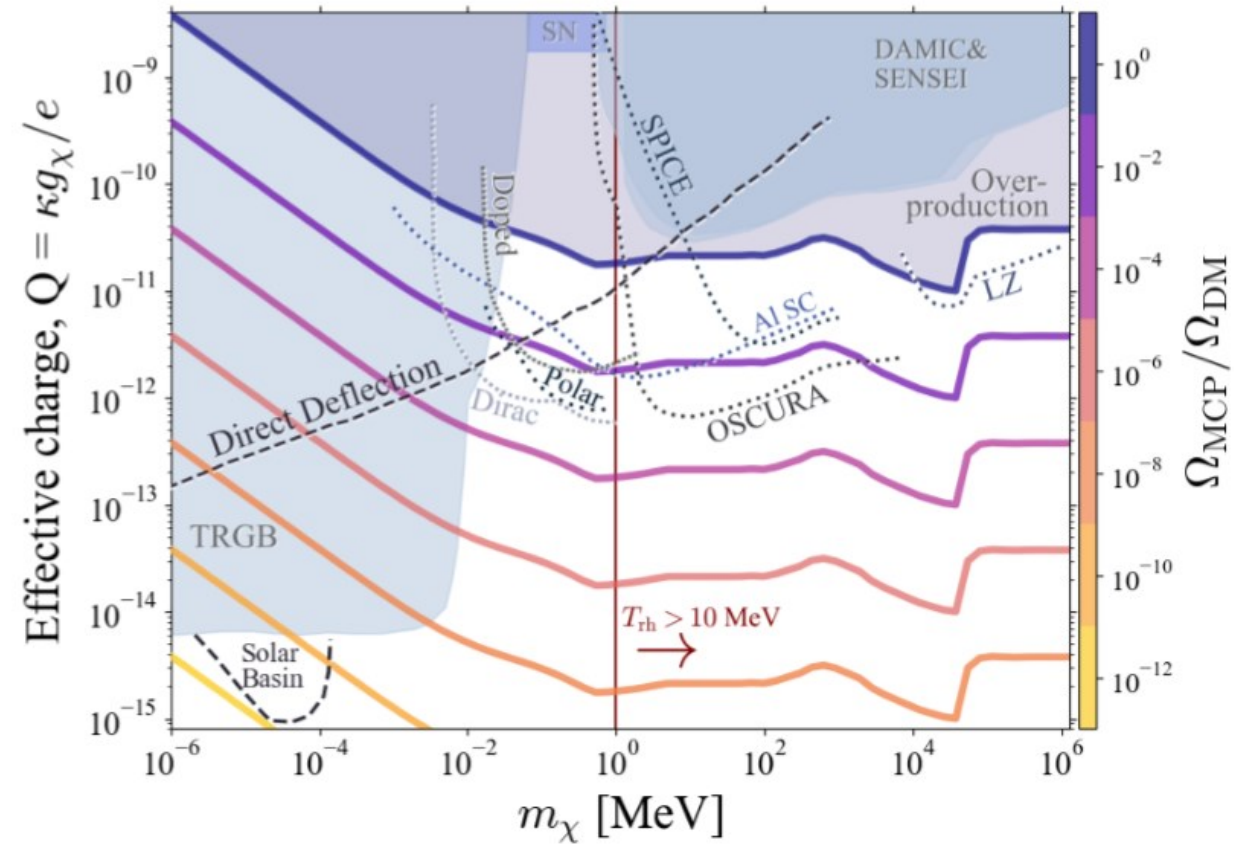
# Summary

- For accelerator searches to actually be relevant, we must ensure that the MCP relic abundance is depleted
  - Preferentially via something that can further decay into SM particles and not dark radiation
- Complete constructions based on dark Higgs mechanisms could work
  - Copious signatures at accelerators since the dark sector is richer
  - They however do bring quite a lot of additional constraints to be fully explored and still feature the tuning of a small charge.
  - Relic abundance currently being estimated to check if the density of MCP can be sufficiently depleted

Backup

# Direct detection constraints?

- Direct detection experiments place strong constraints even on the interactions of dark matter subcomponents
- Necessary to deplete abundance by many orders of magnitude?



# No CMB limits for smaller relic fraction

- CMB constraints at recombination falls off very quickly if relic density smaller the

