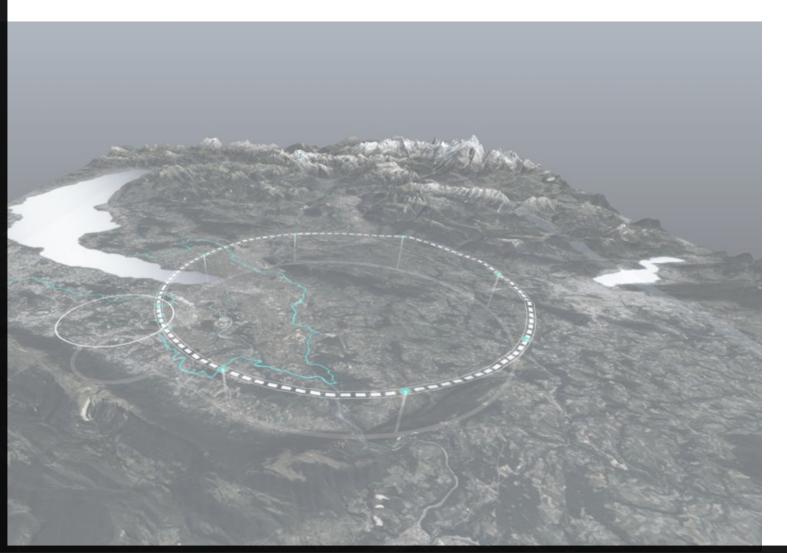
Future Circular Collider —Physics Case—

LPNHE, May 6, 2024









Christophe Grojean

DESY (Hamburg)
Humboldt University (Berlin)
CERN

(christophe.grojean@desy.de)

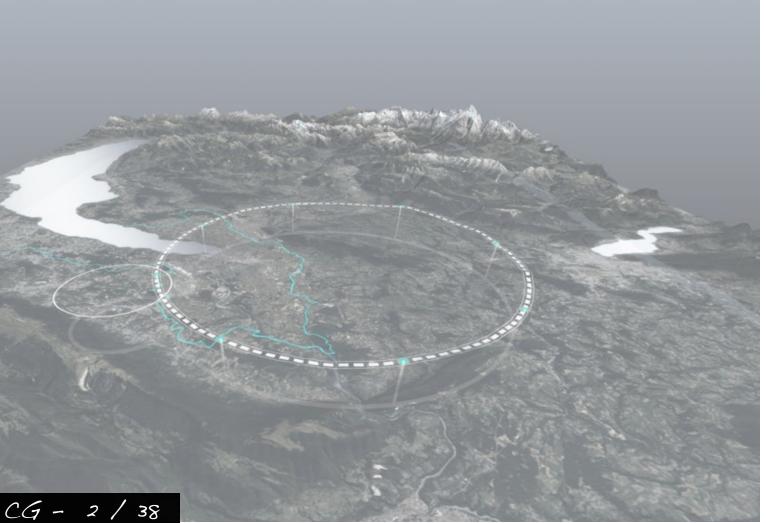
— on behalf of the FCC team—

Future Circular Collider

- A versatile particle collider housed in a 91km underground ring
- Implemented in several stages:
 - an e+e- "Higgs/EW/Flavour/top/QCD" factory running at 90-365 GeV



followed by a high-energy pp collider reaching 100 TeV



Outline

- 1. Why do we need a new collider?
- 2. FCC feasibility study
- 3. FCC-ee: much more than a Higgs factory
- 4. FCC-ee/hh as a Higgs/electroweak factory
- 5. FCC-ee as a flavour factory
- 6. FCC-hh: the broadest exploration potential at high-energy
- 7. FCC-ee→FCC-hh: complementarity and synergy
- 8. Conclusion

1. Why do we need a new collider?

CG - 3 / 38

The LHC Legacy (so far).

- ▶ Standard Model (SM) confirmed to high accuracy up to energies of several TeV (thanks to a firm control of exp. & th. syst. uncertainties, the LHC became a precision machine)
- ▶ Higgs boson discovered at the mass predicted* by LEP precision EW measurements

*within the Standard Model

TeV-scale Naturalness might not explain DM/baryogenesis

Traditional New Physics models are under siege

Absence of new physics

New approaches: relaxion, Nnaturalness, clockwork...

Cosmology might settle the vacuum of the SM -

We need a broad, versatile and ambitious programme that

- 1. sharpens our knowledge of already discovered physics
- 2. pushes the frontiers of the unknown at high and low scales
 - together FCC-ee & FCC-hh combine these 2 aspects —

more PRECISION and more ENERGY, for more SENSITIVITY to New Physics

G-4/38 May 6, 2024

Precision as a discovery tool.

Many historical examples

- Uranus anomalous trajectory Neptune
- ▶ Mercury perihelion → General Relativity
- ▶ Z/W interactions to quarks and leptons → Higgs boson

▶ ...

Sometimes, these discoveries were expected based on theoretical arguments

(e.g. Rayleigh-Jeans UV catastrophe for QM, unitarity breakdown for the Higgs)

but precision gave valuable additional clues.

In any case, experimentalists shouldn't lean too heavily on theorist priors/prejudices

(remember discovery of CP violation).

At times when we don't have a precise theoretical guidance, we need powerful experimental tools to make progress.

The FCC project offers unprecedented opportunities on many different fronts. No LHC/SSC-like **no-lose theorem** but a **promise** of making significant steps forward in our understanding of the fundamental laws of Nature.

CG - 5 / 38 May 6, 2024

The Higgs requires more precision.

"The Higgs isn't everything; it's the only thing!"*
The scalar discovery in 2012 has been an important milestone for HEP.
Many of us are still excited about it. Others should be too.

Higgs = **new forces** of different nature than the interactions known so far

- No underlying local symmetry.
- No quantised charges.
- Deeply connected to the space-time vacuum structure.
- The discovery of the Higgs opens new deep questions —
- What is the origin of the Higgs boson?
- Is it elementary and isolated, or does it emerge from a deeper underlying dynamics?
- Which role did the Higgs play during the big bang, and how did it influence the evolution of the Universe?
- Does the Higgs boson play a role in explaining other fundamental open questions in particle physics which the SM cannot address (flavour, DM, baryogenesis, inflation...)

CG - 6 / 38

The Higgs requires more precision.

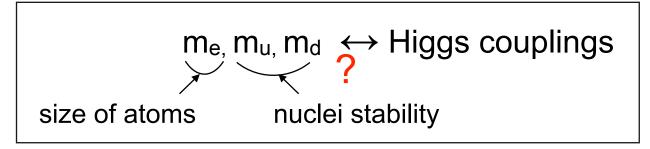
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Higgs = **new forces** of different nature than the interactions known so far

- No underlying local symmetry.
- No quantised charges.
- Deeply connected to the space-time vacuum structure.

The knowledge of the values of the **Higgs couplings** is essential to understand the deep structure of matter/Universe:

EWSB @
$$t\sim10^{-10}s \leftrightarrow \frac{\text{Higgs self-coupling(s)}}{\text{Higgs(es) potential}}$$



```
matter/anti-matter ↔ CPV in Higgs sector
```

CG - 6/38

The Higgs requires more precision.

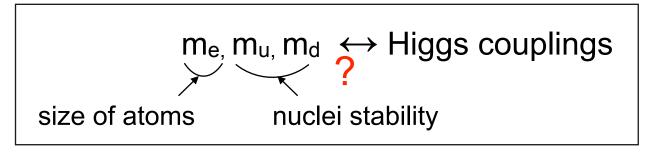
(HL)-LHC will make remarkable progress.

But it won't be enough.

A new collider is needed!

The knowledge of the values of the **Higgs couplings** is essential to understand the deep structure of matter/Universe:

```
EWSB @ t\sim10^{-10}s \leftrightarrow \frac{\text{Higgs self-coupling(s)}}{\text{Higgs(es) potential}}
```



matter/anti-matter ↔ CPV in Higgs sector

CG - 6/38

2. FCC feasibility study

CG - 7/38

The launch of the feasibility study.



"An **electron-positron** Higgs factory is the highest-priority next collider. For the longer term, the European particle physics community has the ambition to operate a **proton-proton** collider at the highest achievable energy."

— CERN council approved the Strategy and CERN management implemented it — FCC Feasibility Study (FS) started in 2021 and will be completed in 2025.
Mid-term review in 2023.

CG - 8 / 38 May 6, 2024

Objectives of FCC feasibility study.

- Demonstration of the geological, technical, environmental and administrative feasibility of the tunnel and surface areas and optimisation of placement and layout of the ring and related infrastructure.
- Pursuit, together with the Host States, of the preparatory administrative processes required for a potential project approval to identify and remove any showstopper.
- Optimisation of the design of the colliders and their injector chains, supported by R&D to develop the needed key technologies.
- Elaboration of a sustainable operational model for the colliders and experiments in terms of human and financial resource needs, as well as environmental aspects and energy efficiency.
- Development of a consolidated cost estimate, as well as the funding and organisational models needed to enable the project's technical design completion, implementation and operation.
- Identification of substantial resources from outside CERN's budget for the implementation of the first stage of a
 possible future project (tunnel and FCC-ee).
- Consolidation of the physics case and detector concepts for both colliders.

CG - 9 / 38 May 6, 2024

Optimized placement and layout.

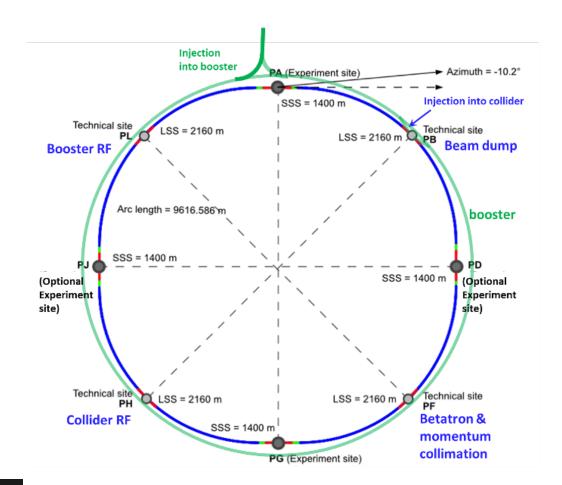
M. Benedikt @ CERN 13.02.24

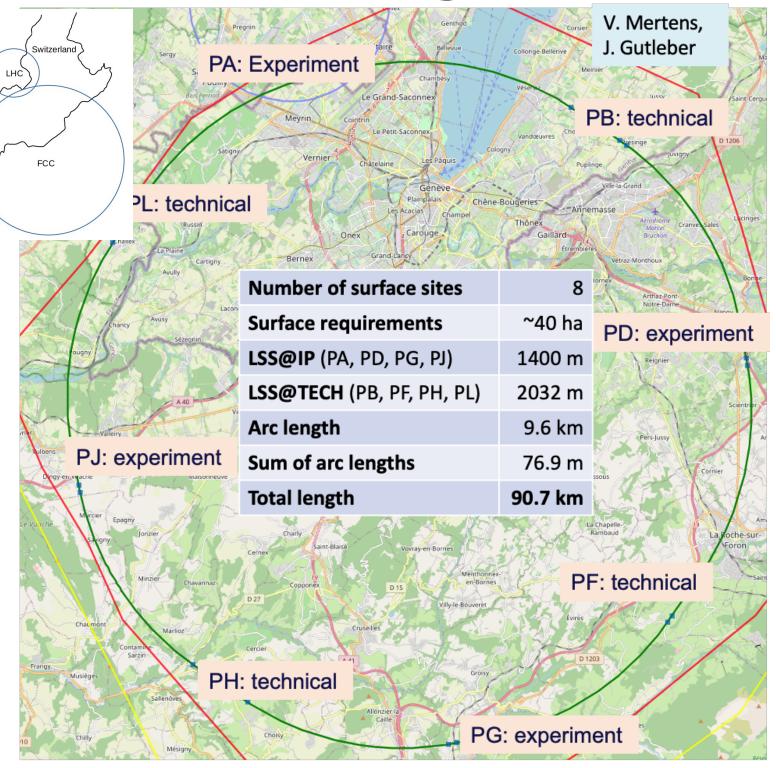
Layout chosen out of ~ 100 initial variants, based on **geology** and **surface constraints** (land availability, access to roads, etc.), **environment**, (protected zones), **infrastructure** (water, electricity, transport), **machine performance** etc.

"Avoid-reduce-compensate" principle of EU and French regulations

Overall lowest-risk baseline: 90.7 km ring, 8 surface points.

Whole project now adapted to this placement





CG - 10 / 38

Optimized placement and layout.

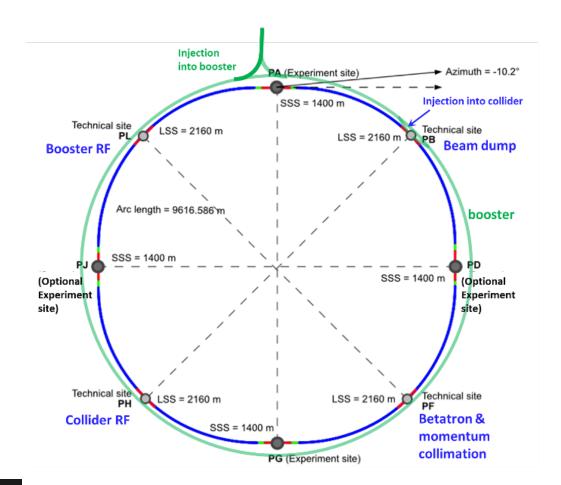
M. Benedikt @ CERN 13.02.24

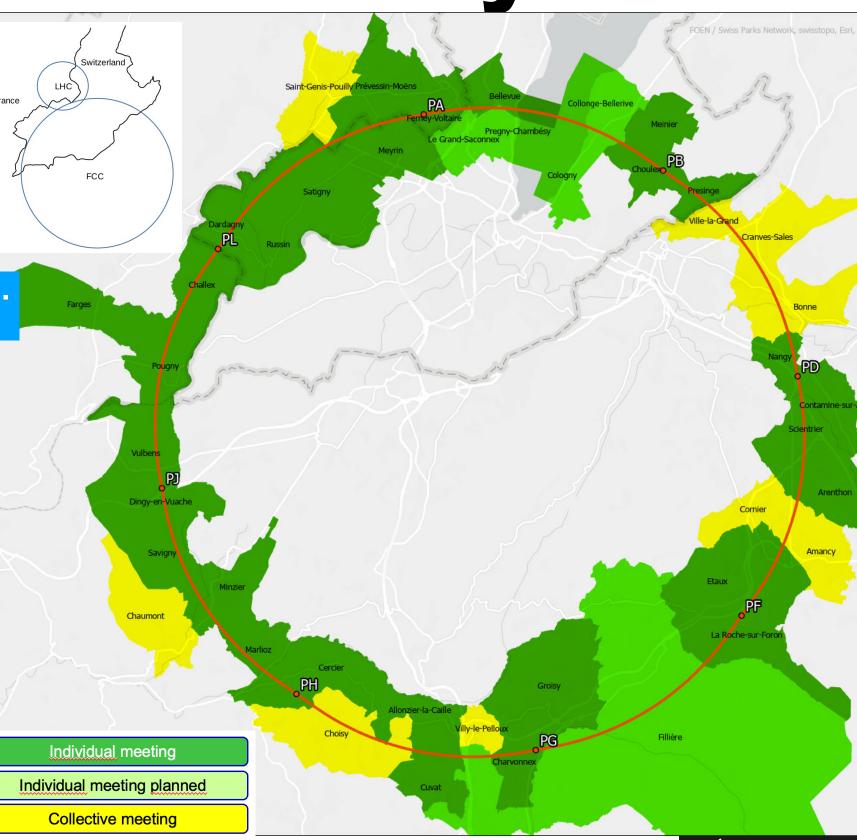
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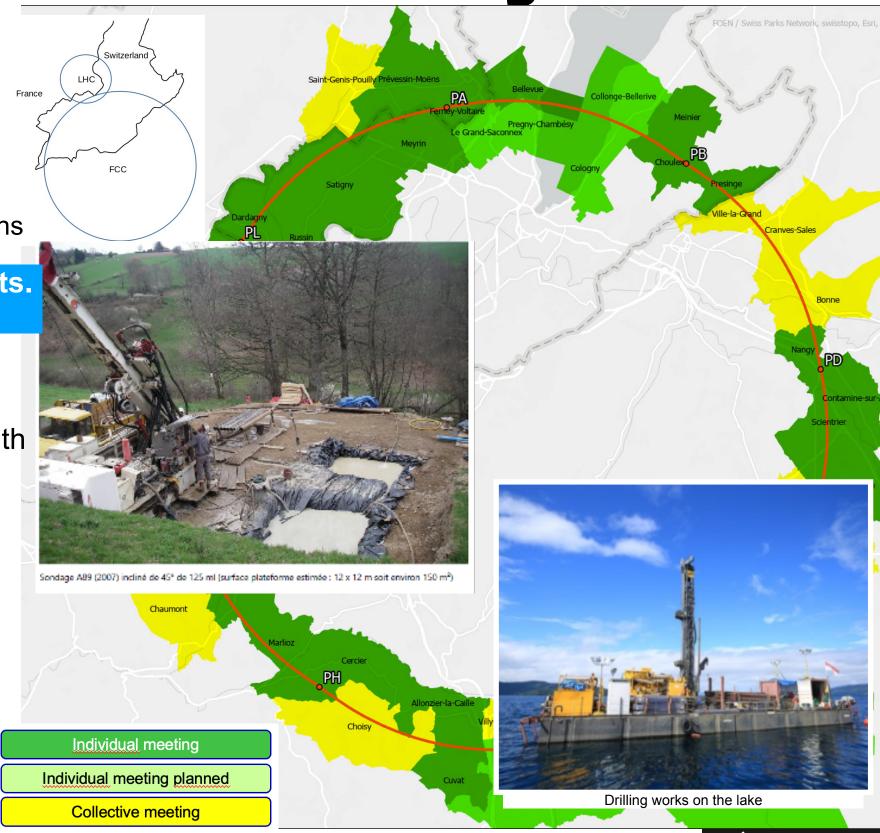
Overall lowest-risk baseline: 90.7 km ring, 8 surface points.

Whole project now adapted to this placement

- Site investigations in areas with uncertain geological conditions:
 - ► Optimisation of localisation of drilling locations ongoing with site visits since end 2022.
 - ► Alignment with FR and CH on the process for obtaining autorisation procedures. Ongoing for start of drillings in Q2/2024

Contracts Status:

- ► Contract for engineering services and role of Engineer during works, active since July 2022
- ► Site investigations tendering ongoing towards contract placement in December 2023 and mobilization from January 2024



Environmental considerations.

M. Benedikt @ CERN 13.02.24

- Excavated material from FCC subsurface infrastructures: 6.5 Mm³ in situ, 8.4 Mm³ excavated
- Priority : reuse, minimize disposal
- 2021-2022: International competition "Mining the Future", launched with the support of the EU Horizon 2020 grant, to find innovative and realistic ideas for the reuse of molasse (96% of excavated materials)
- 2023: "OpenSky Laboratory" project: Objective Develop and test an innovative process to transform
 sterile "molasse" into fertile soil for agricultural use and
 afforestation. launched in Jan. 2024: 5500m² near LHC
 P5 in Cessy (FR). Trial with 5 000t of excavated local
 molasse → convert it to arable soil (agricultural/forestry)
- Heat:
 - heating for local houses
 - cheese factories in Jura and Haute-Savoie expressed special interest

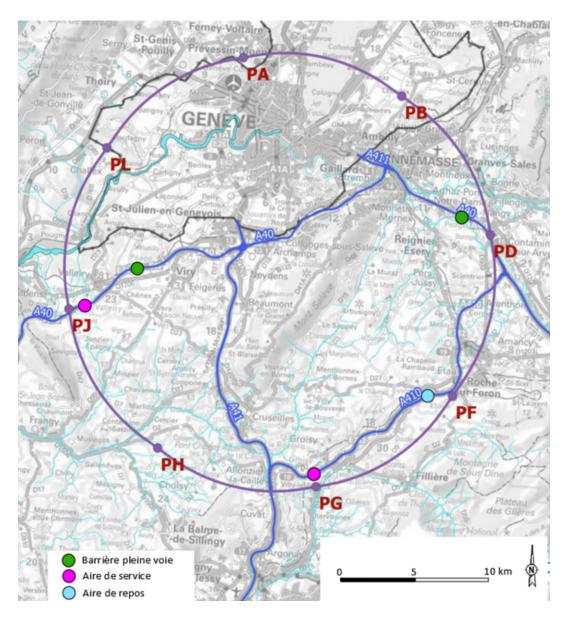


Accelerated soil transformation with funghi

Connections with local infrastructre.

M. Benedikt @ CERN 13.02.24

- Road accesses developed for all 8 surface sites
 - ► Four possible highway connections defined
 - ► Less than 4 km new departmental roads required

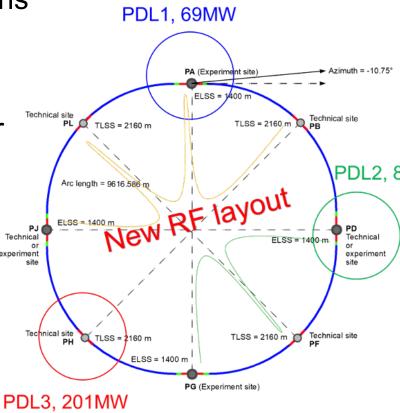


Connections to electrical grid

► Electrical connection concept studied by RTE (French electrical grid operator) → requested loads have no significant impact on grid

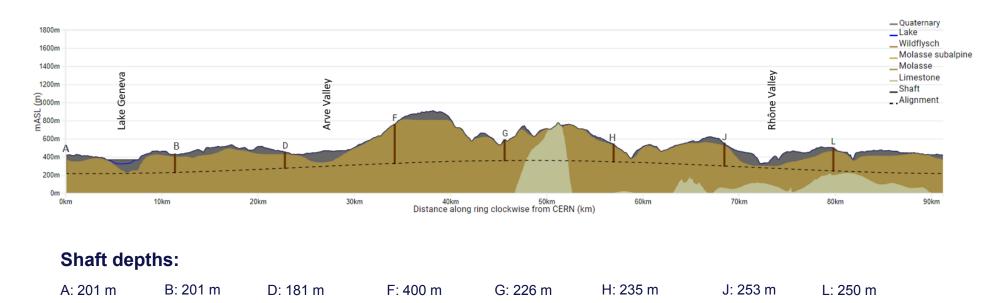
► Powering concept and power rating of the three sub-stations compatible with FCC-hh

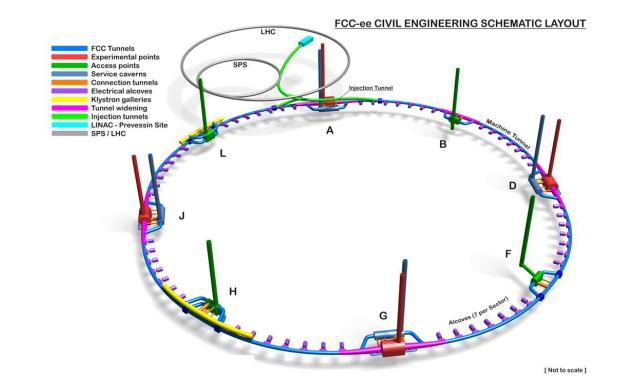
 R&D efforts aiming at further reduction of the energy consumption of FCCee and FCC-hh



Civil engineering

T. Watson @ Annecy FCC Physics '24

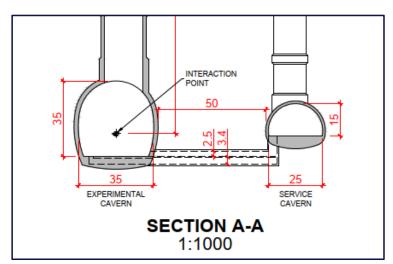


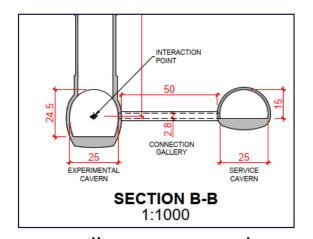




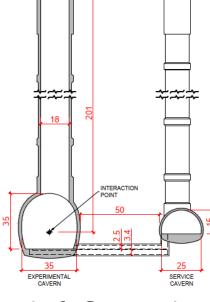
Tunnel Boring Machine (TBM)

- Tunnel Boring Machines are designed to work almost continuously 24/7 other than periodic maintenance. Rate of 18m/day in the Molasse. 21-27 months to complete one sector \rightarrow 8 years with two TBMs.
- 13 shafts
- 2/2 large/small caverns



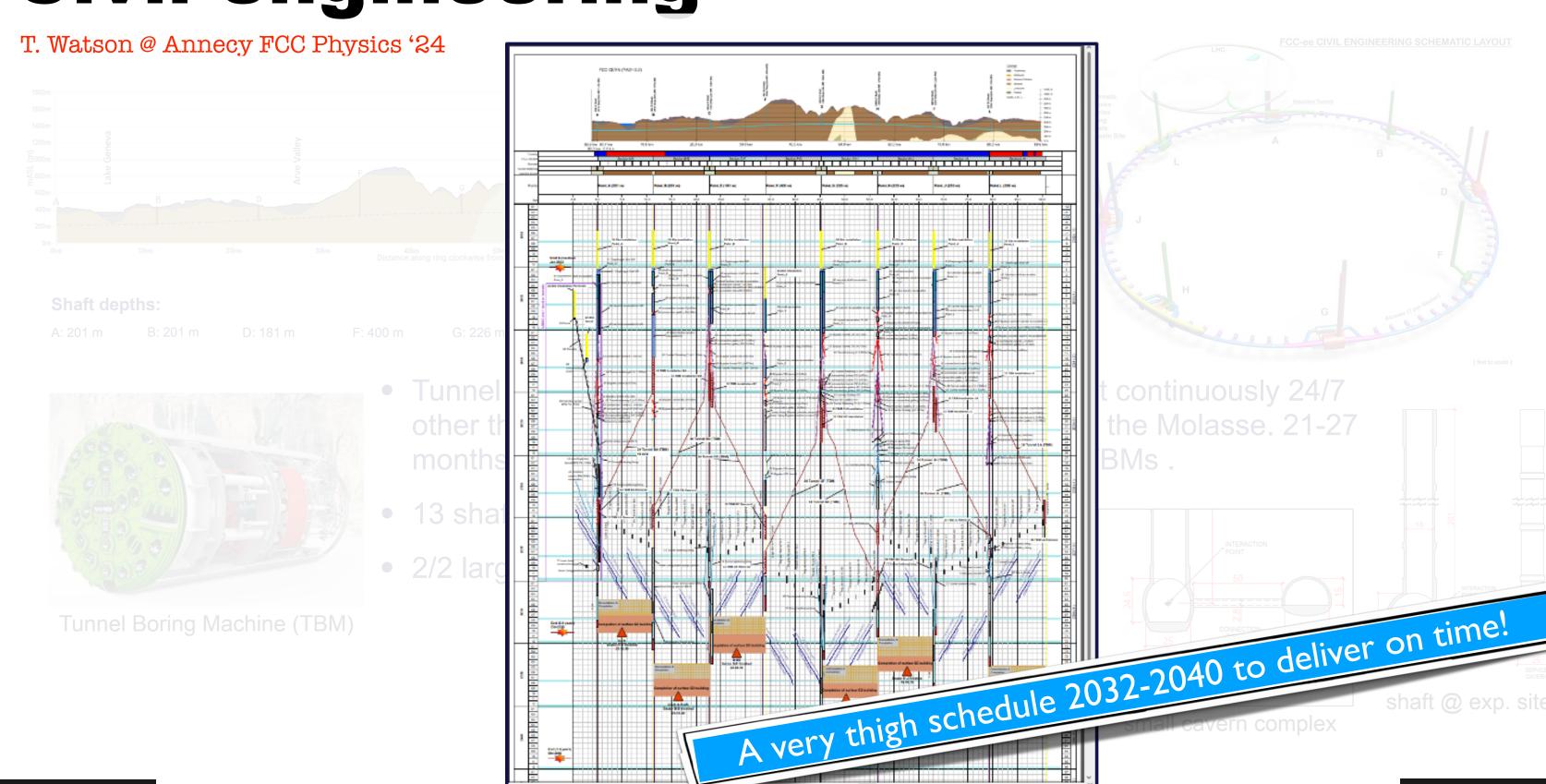


small cavern complex



shaft @ exp. site

Civil engineering



FCC feasibility mid-term report.

- 703 pages: 7 chapters (cost and financial feasibility is a separate document) + refs.
 - Placement scenario (75 pages)
 - Civil engineering (50 pages)
 - Implementation with the host states (45 pages)
 - Technical infrastructure (110 pages)
 - FCC-ee collider design and performance (170 pages)
 - FCC-hh accelerator (60 pages)
 - (Cost and financial feasibility)
 - Physics and experiments (110 pages)
 - References (70 pages)
- Executive summary: 44 pages
- Reviewed by
 - Scientific Advisory Committee and Cost Review Panel on Oct. 16-18
 - Scientific Policy Committee and Financial Committee on Nov. 21-22
 - CERN Council Feb. 2

Future Circular Collider Midterm Report

February 2024

296 authors 16 editors

Edited by:

B. Auchmann, W. Bartmann, M. Benedikt, J.P. Burnet, P. Craievich, M. Giovannozzi, C. Grojean, J. Gutleber, K. Hanke, P. Janot, M. Mangano, J. Osborne, J. Poole, T. Raubenheimer, T. Watson, F. Zimmermann



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This document has been produced by the organisations participating in the FCC feasibility study. The studies and technical concepts presented here do not represent an agreement or commitment of any of CERN's Member States or of the European Union for the construction and operation of an extension to CERN's existing research infrastructures.

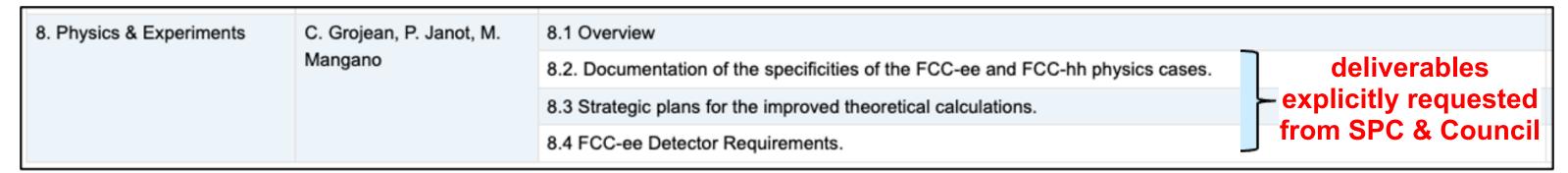
The midterm report of the FCC Feasibility Study reflects work in progress and should therefore not be propagated to people who do not have direct access to this document.

confidential documents (work in progress) available to CERN personnel

CG - 14 / 38 May 6, 2024

Physics, Experiments, Detectors.

FCC Feasibility Study PED deliverables for mid-term review



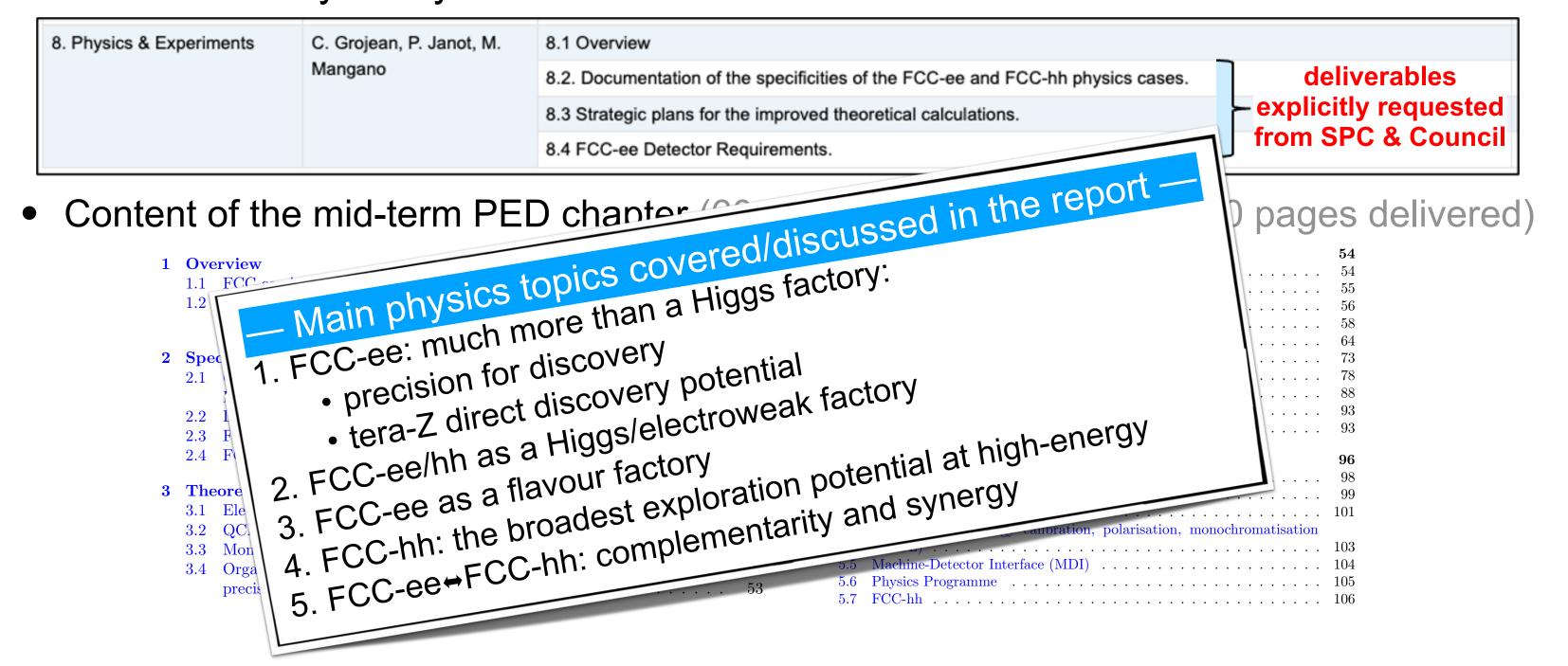
Content of the mid-term PED chapter (60 pages were expected → 110 pages delivered)

	0	n	4 D	etector requirements	$\bf 54$
L	Overview	3	4.	1 Introduction	54
	1.1 FCC-ee: A great Higgs factory, and so much more	4	4.	2 Machine-detector interface	
	1.2 FCC-hh: The energy-frontier collider with the broadest exploration		4.	3 The current detector concepts	
	potential	13	4.	4 Measurement of the tracks of charged particles	
			4.	5 Requirements on the vertex detector	
2	Specificities of the FCC physics case	15	4.	6 Requirements on charged hadron particle identification	
	2.1 Characterisation of the Higgs boson: role of EW measurements and of		4.	7 Requirements on electromagnetic calorimetry	78
	FCC-hh	16	4.	8 Requirements on the hadronic calorimeter	88
	2.2 Discovery landscape	_	4.	9 Requirements on the muon detector	93
	2.3 Flavour advancement	3/1	4.	10 Precise timing measurements	93
	2.4 FCC-hh specificities compared to lepton colliders	36			
	2.4 FCC-III specificities compared to repton coniders	30		utlook and further steps	96
_			5.	1 Software and Computing	98
3	Theoretical calculations	42	5.	2 Physics Performance	99
	3.1 Electroweak corrections	44	5.		
	3.2 QCD precision calculations	46	5.	4 Centre-of-mass energy calibration, polarisation, monochromatisation	
	3.3 Monte Carlo event generators	50		(EPOL)	103
	3.4 Organization and support of future activities to improve theoretical		5.	5 Machine-Detector Interface (MDI)	104
	precision	53	5.	6 Physics Programme	105
	precision	00	5.	7 FCC-hh	106

CG - 15 / 38 May 6, 2024

Physics, Experiments, Detectors.

• FCC Feasibility Study PED deliverables for mid-term review

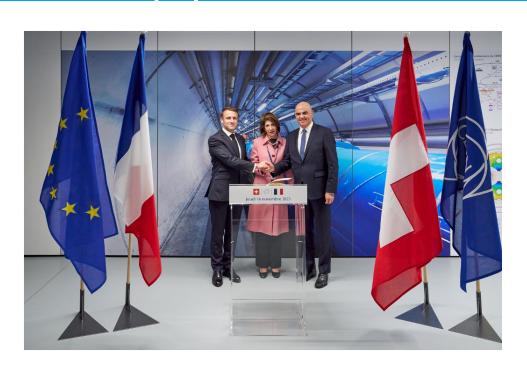


CG - 15 / 38

Feedback.

Andy **Parker** (SAC chair), Norbert **Holtkamp** (CRP chair), Hugh **Montgomery** (SPC chair), Laurent **Salzarulo** (FC chair), Eliezer **Rabinovici** (Council president) "many thanks for the work done, congratulations for the results, impressive quality of the study…"

"Financial Committee underlines the need to make the project attractive from the physics viewpoint and takes the view that it would be unfortunate to sacrifice the attractiveness of the physics for the sake of reducing costs."



"Si j'ai voulu venir là aujourd'hui c'est pour témoigner ma confiance aux équipes et notre volonté, notre ambition de conserver la première place dans ce domaine." ["My visit here bears witness to my trust in CERN personnel and France's will and ambition to keep the leadership in this domain."]

E. Macron, CERN 16.11.2023

CG - 16 / 38

US Statement of Intent

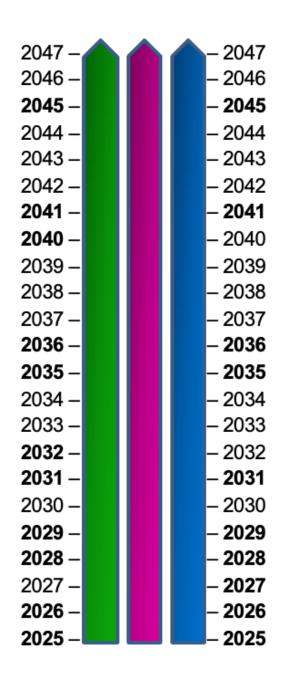


"Should the CERN Member States determine the FCC-ee is likely to be CERN's next world-leading research facility following the highluminosity Large Hadron Collider, the United States intends to collaborate on its construction and physics exploitation, subject to appropriate domestic approvals."

White House, April 26, 2024

CG - 17 / 38 May 6, 2024

The way forward.



We are only 20 years away from the first collisions!

May 6, 2024

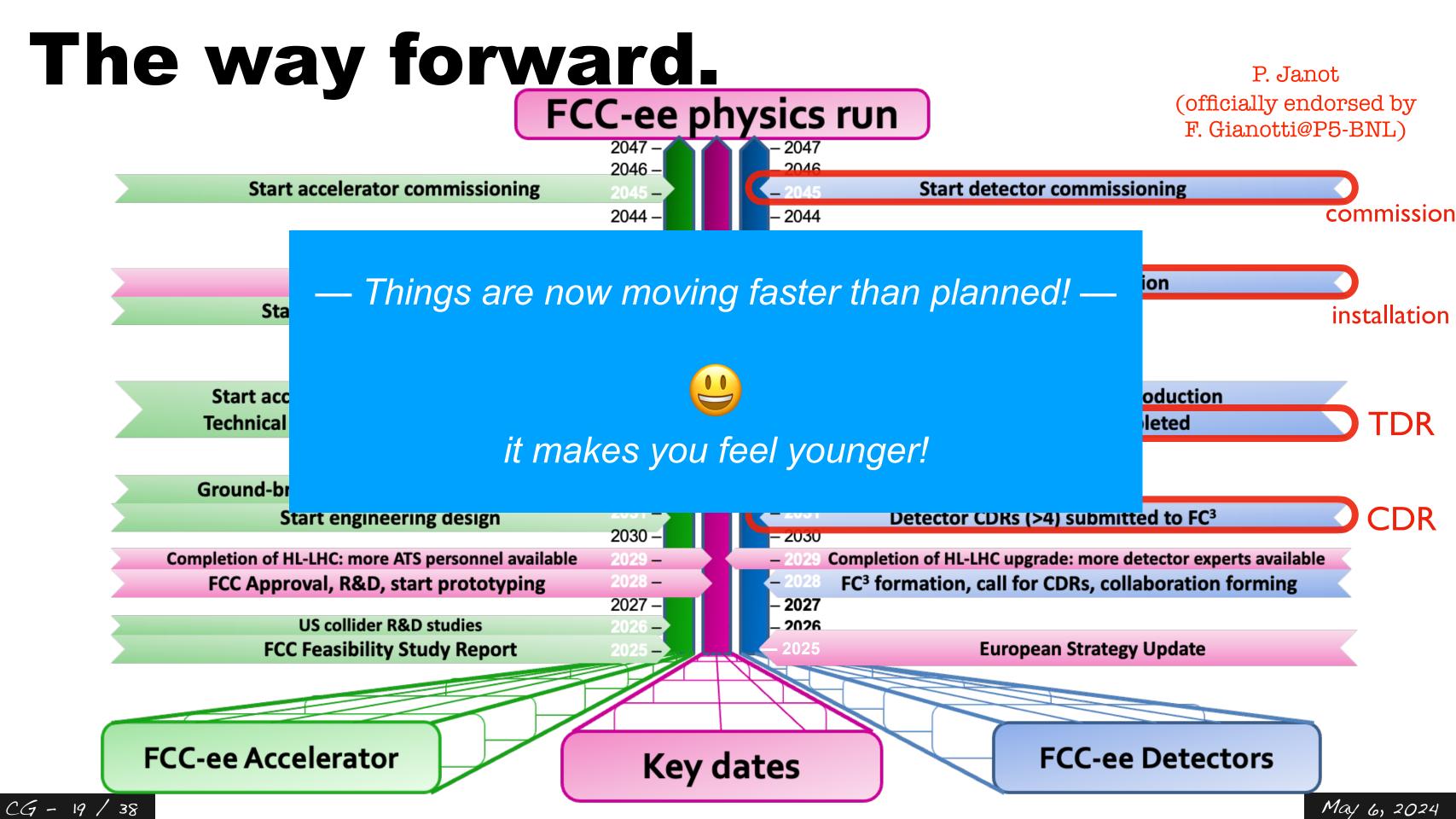


The way forward. P. Janot FCC-ee physics run (officially endorsed by F. Gianotti@P5-BNL) 2046 2046 **Start detector commissioning** Start accelerator commissioning commission - 2044 2044 -2043 -- 2043 2042 -2042 Start detector installation **End of HL-LHC operation** - 2040 Start accelerator installation installation 2039 --20392038 -- 2038 - 2037 2037 -Start accelerator component production Start detector component production **Technical design & prototyping completed** TDR Four detector TDRs completed 2034 -**– 2034** 2033 -- 2033 Ground-breaking and start civil engineering - 2032 CDR Detector CDRs (>4) submitted to FC³ Start engineering design 2030 -- 2030 Completion of HL-LHC: more ATS personnel available Completion of HL-LHC upgrade: more detector experts available FCC Approval, R&D, start prototyping FC³ formation, call for CDRs, collaboration forming 2027 -**European Strategy Update US collider R&D studies** Detector EoI submission by the community Eol **FCC Feasibility Study Report FCC-ee Accelerator**

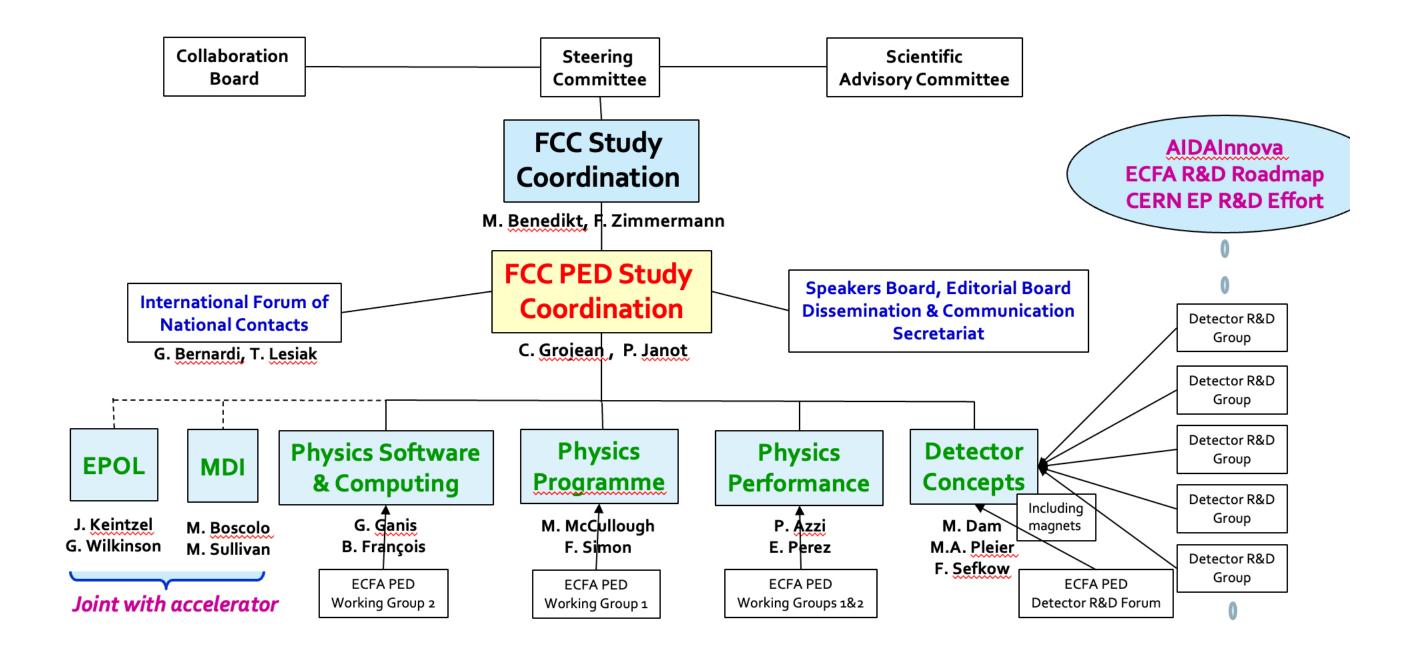
Key dates

CG - 19 / 38

FCC-ee Detectors



How to contribute?

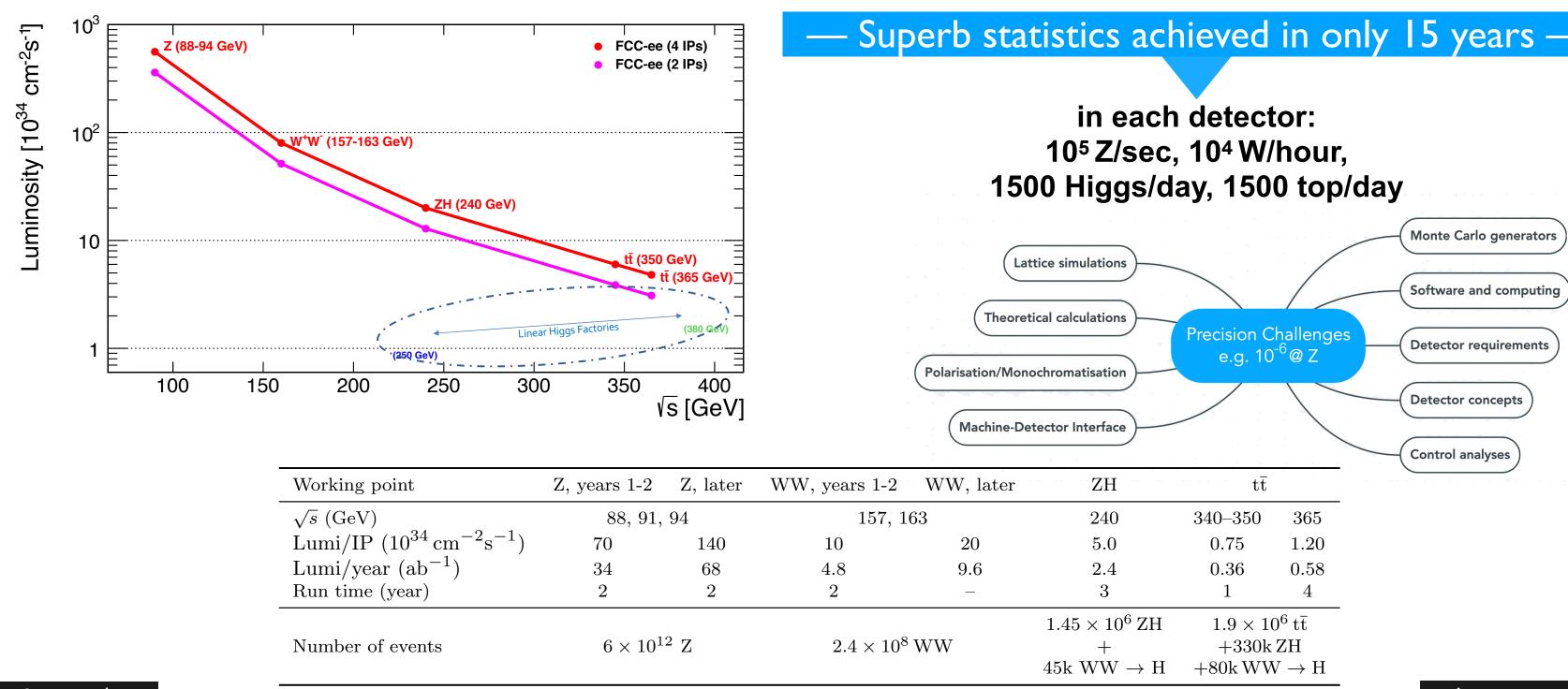


CG - 20 / 38 May 6, 2024

3. FCC-ee: much more than a Higgs factory

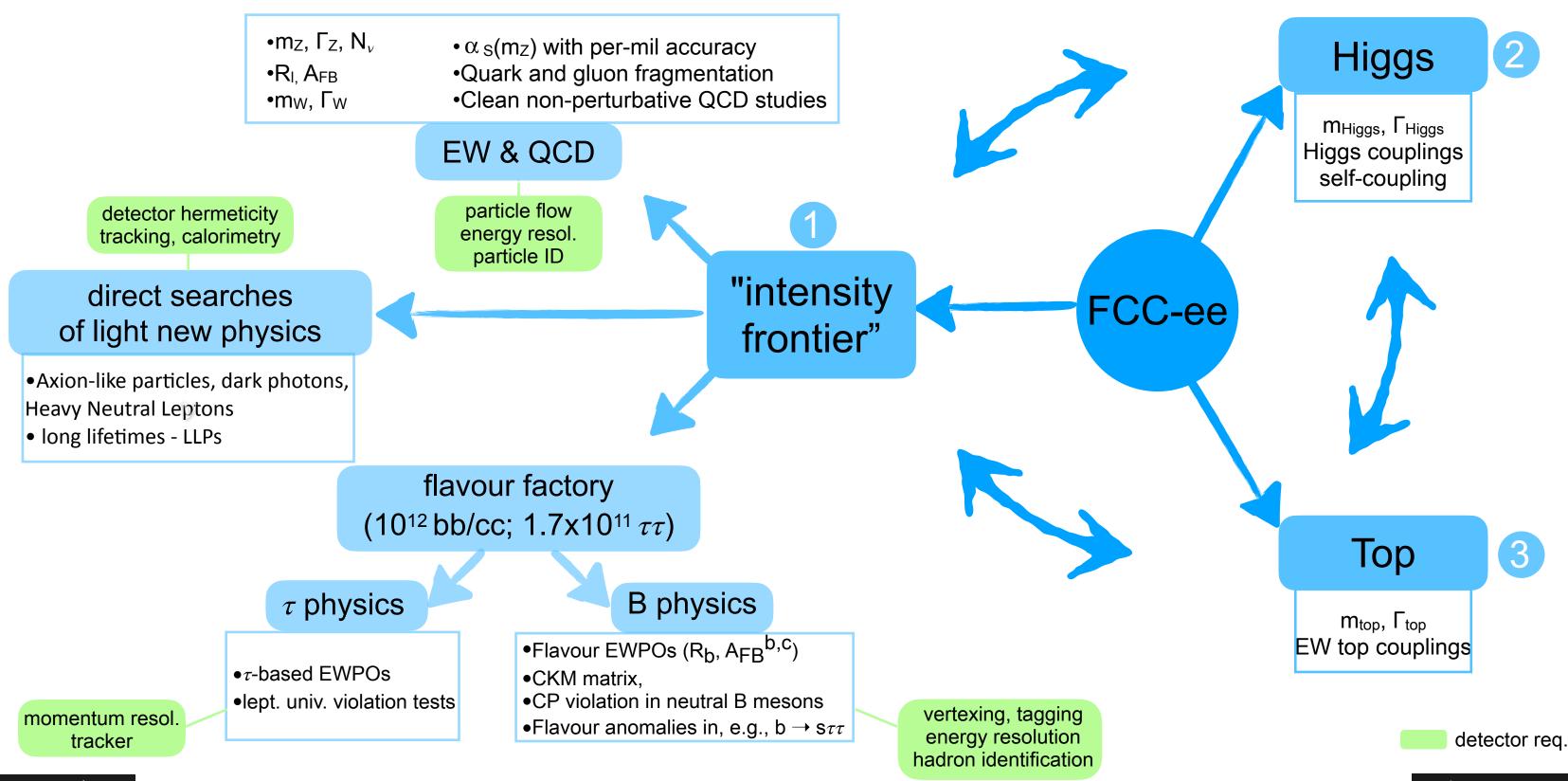
FCC-ee Run Plan.

LEP1 data accumulated in **every 2 mn**. Exciting & diverse programme with different priorities every few years. (order of the different stages still subject to discussion/optimisation)



CG - 22 / 38 May 6, 2024

FCC-ee Physics Programme.



CG - 23 / 38

May 6, 2024

FCC-ee Physics Programme.

- • m_Z , Γ_Z , N_{ν}
 - $\alpha_{s}(m_z)$ with per-mil accuracy
- •R_{I,} A_{FB} •m_W, Γ_W
- Quark and gluon fragmentationClean non-perturbative QCD studies

EW & QCD



Higgs

m_{Higgs}, F_{Higgs} Higgs couplings self-coupling

detector hermetici tracking, calorimet

direct search of light new ph

Axion-like particles, darkHeavy Neutral Leptonslong lifetimes - LLPs

baseline i GG de actector periorinarie	baseline	FCC-ee	detector	performance
--	----------	--------	----------	-------------

track momentum	$\frac{\sigma_p}{p} = 0.02 \cdot 10^{-3} \cdot p_T(\text{GeV}) \oplus 1 \cdot 10^{-3}$
	15 um

track impact parameter
$$\sigma_{d_0} = rac{15\,\mu\mathrm{m}}{\sin^{3/2} heta} \oplus 5\,\mu\mathrm{m}$$

electromagnetic energy
$$\dfrac{\sigma_{E_{\gamma}}}{E_{\gamma}}=\dfrac{15\%}{E_{\gamma}}\oplus 1\%$$

electromagnetic energy
$$xy$$
 position $\sigma_{\gamma,xy} = \frac{6\,\mathrm{mm}}{E(\mathrm{GeV})} \oplus 2\,\mathrm{mm}$



physics B physic

●r-based EWPOs

pt. univ. violation tests

•Flavour EWPOs (R_b, A_{FB}^{b,c})

CKM matrix,

CP violation in neutral B mesons

•Flavour anomalies in, e.g., b ightarrow sau au

m_{top}, Γ_{top} EW top coupl

detector req

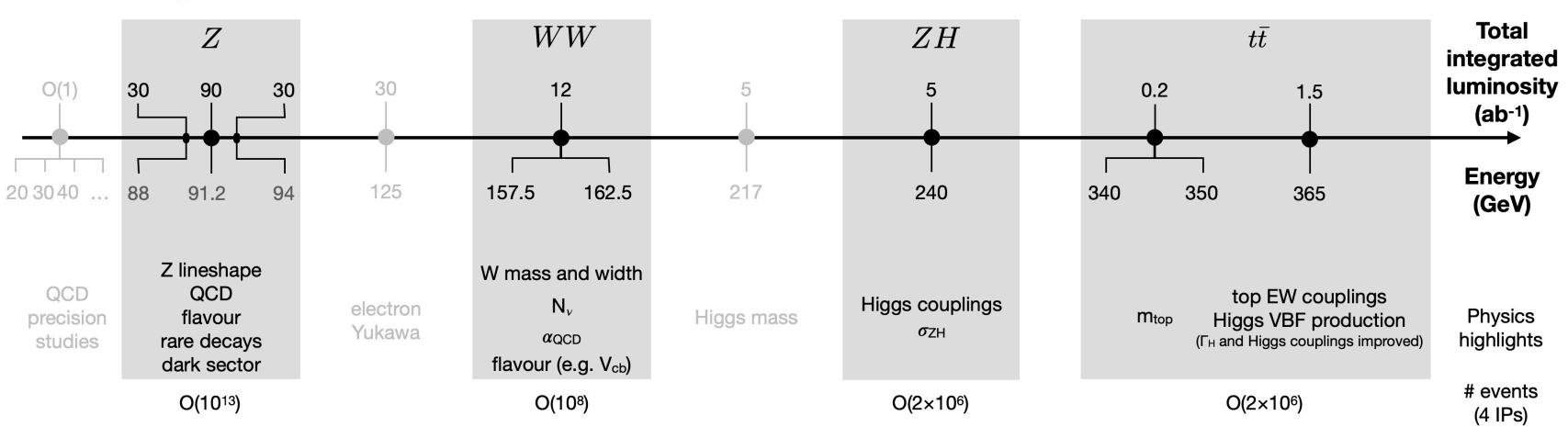
momentum re tracker

CG - 23 / 38

May 6, 2024

Collider Programme (and beyond).

- CDR baseline runs (2IPs)
- Additional opportunities



- **Opportunities** beyond the baseline plan (√s below Z, 125GeV, 217GeV; larger integrated lumi...)
- Opportunities to exploit FCC facility differently (to be studied more carefully):
 - using the electrons from the injectors for beam-dump experiments,
 - extracting electron beams from the booster,
 - reusing the synchrotron radiation photons.

CG - 24 / 38 May 6, 2024

FCC-ee: Explore & Discover.

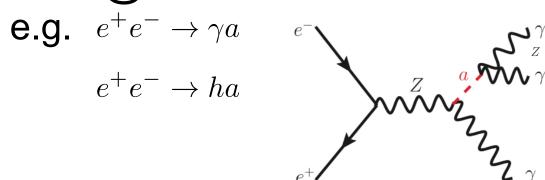
PED @ CERN-SPC 2022

- **EXPLORE INDIRECTLY** the 10-100 TeV energy scale with precision measurements
 - \bullet From the correlated properties of the Z , b, c, τ , W, Higgs, and top particles
 - ► Up to 20-50-fold improved precision on ALL electroweak observables (EWPO)
 - \rightarrow m_Z , m_W , m_{top} , Γ_Z , sin² θ_W^{eff} , R_b , $\alpha_{QED}(m_Z)$, $\alpha_s(m_Z m_W m_t)$, top EW couplings ...
 - ► Up to 10 × more precise and model-independent Higgs couplings (width, mass) measurements
 - → Access the Higgs potential and infer the vacuum structure of the Universe
 - → Reveals the dynamics of the EW phase transition and infer the fate of the EW vacuum
- **DISCOVER** that the Standard Model does not fit
 - New Physics! → Pattern of deviations may point to the source.
- **DISCOVER** a violation of flavour conservation / universality
 - $Z \rightarrow \tau \mu$ in $5x10^{12}$ Z decays; $\tau \rightarrow \mu \nu / e \nu$ in 2×10^{11} τ decays; $B^0 \rightarrow K^{*0} \tau^+ \tau^-$ or $B_S \rightarrow \tau^+ \tau^-$ in 10^{12} bb evts
- DISCOVER dark matter, e.g., as invisible decays of Higgs or Z
- **DISCOVER DIRECTLY** elusive (aka feebly-coupled) particles
 - in the 5-100 GeV mass range, such as right-handed neutrinos, dark photons, light Higgs-like scalars, dilaton, ALPs, relaxions...

See Bonus Slides
for examples and plots

FCC-ee: Explore & Discover.

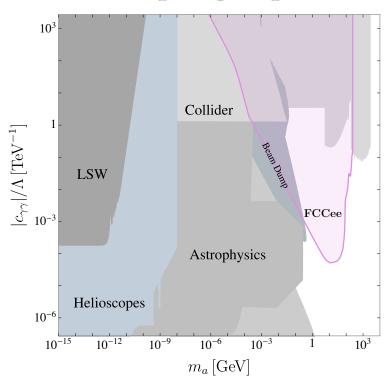
ALPs@ colliders



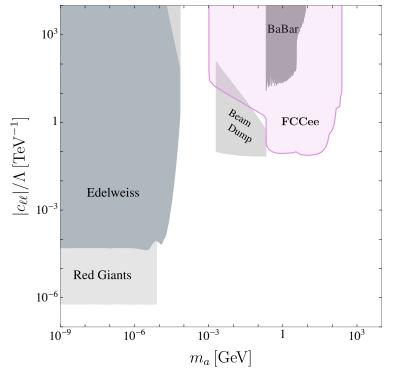
Knapen, Thamm arXiv:2108.08949

Astro/Cosmo → long-lived ALPs colliders → short-lived ALPs MeV+

ALP coupling to photons

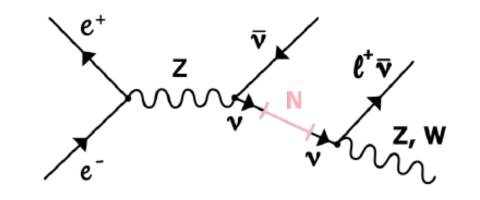


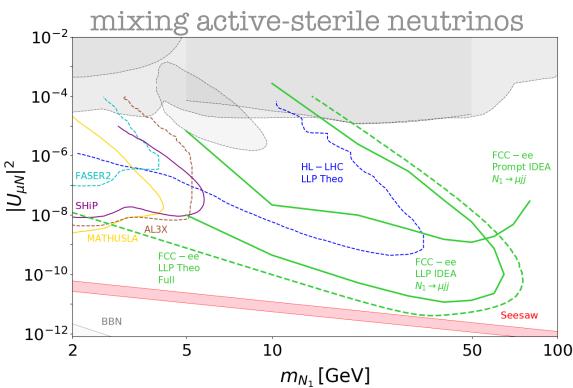
ALP coupling to electrons



• Search for u_{RH} .

Direct observation in Z decays from LH-RH mixing

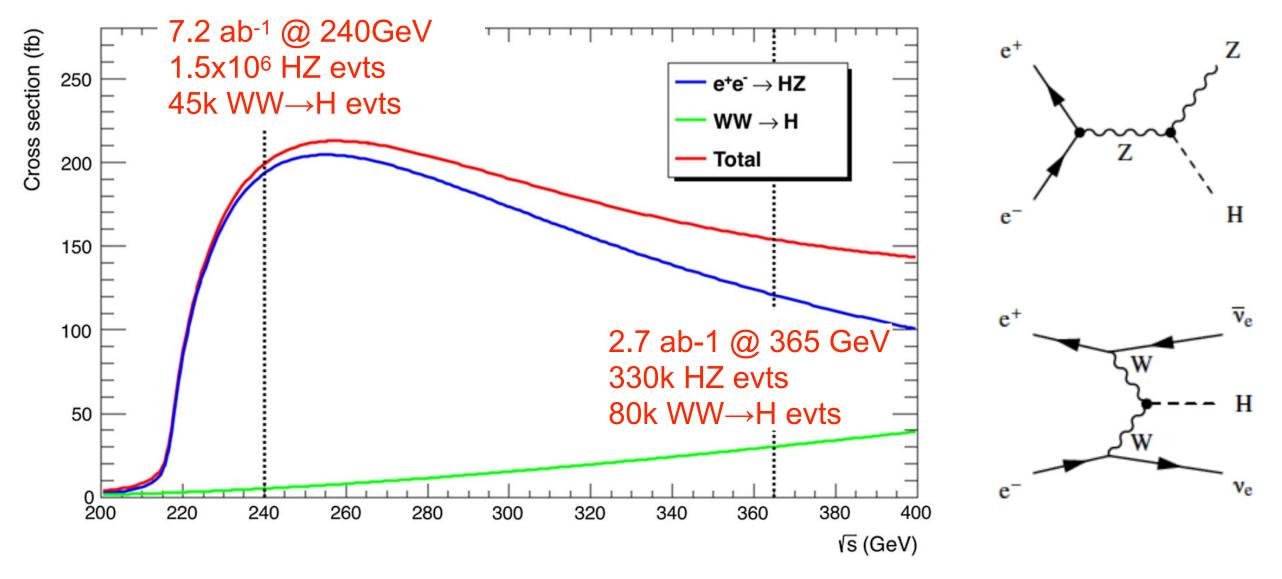




4. FCC-ee/hh as a Higgs/electroweak factory

Higgs @ FCC-ee.

Central goal of FCC-ee: model-independent measurement of Higgs width and couplings with (<)% precision. Achieved through operation at two energy points.



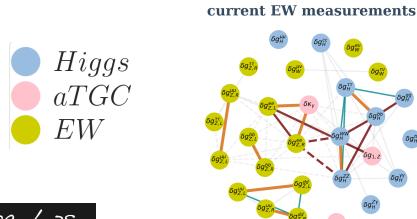
Sensitivity to both processes very helpful in improving precision on couplings.

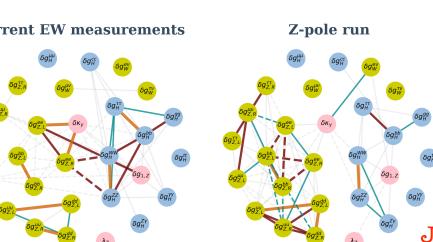
Complementarity with 365GeV on top of 240GeV improvement factor: $\infty/3/2/1.5/1.2$ on $\kappa_{\lambda}/\kappa_{W}/\kappa_{b}/\kappa_{g}$, $\kappa_{c}/\kappa_{\gamma}$ (plot in bonus)

Higgs @ FCC-ee.

- Absolute normalisation of couplings (by recoil method)
- Measurement of width (from ZH>ZZZ* and WW>H)
- $\delta\Gamma_H\sim 1\%, \delta m_H\sim 3\,{
 m MeV}$ (resp. 25%, 30 MeV @ HL-LHC)
- Model-independent coupling determination and improvement factor up to 10 compared to LHC
- (Indirect) sensitivity to new physics up to 70 TeV (for maximally strongly coupled models) $(\delta \kappa_X = v^2/f^2 ~\&~ m_{\rm NP} = g_{\rm NP} f)$
- Unique access to electron Yukawa

— Higgs programme needs Z-pole —





Higgs coupling sensitivity

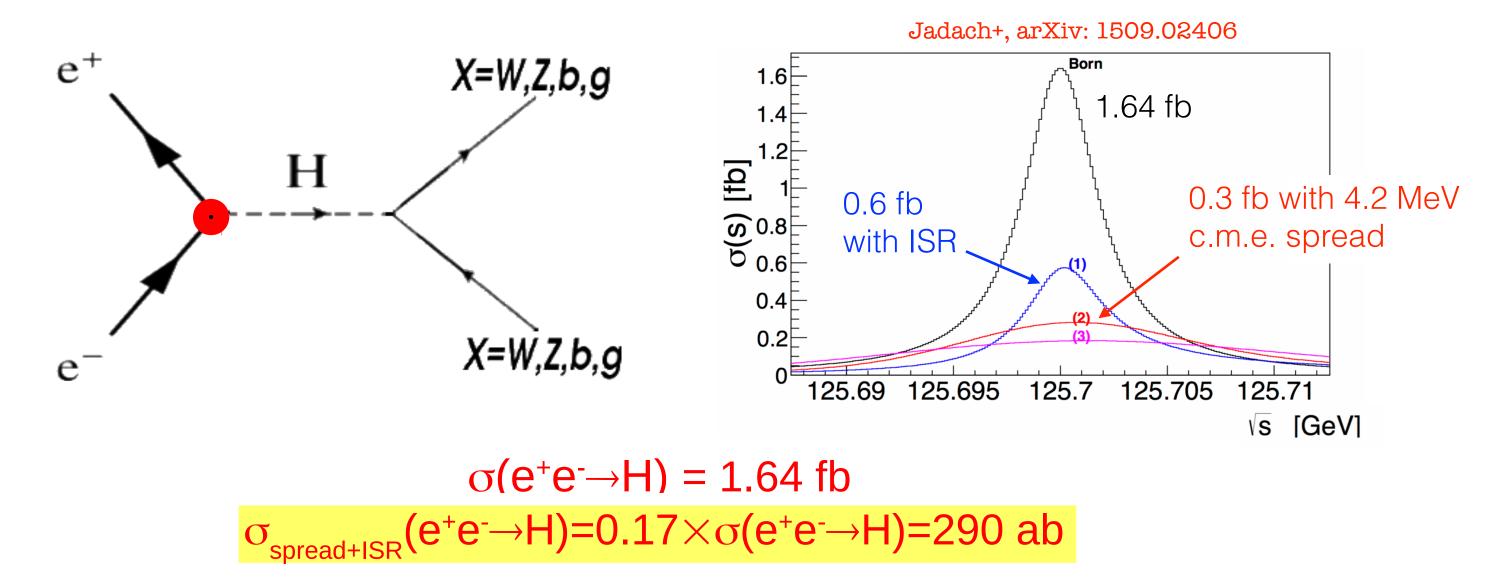
Coupling	HL-LHC	FCC-ee (240–365 GeV) 2 IPs / 4 IPs
$egin{aligned} \kappa_W & [\%] \ \kappa_Z & [\%] \ \kappa_g & [\%] \ \kappa_{\gamma} & [\%] \ \kappa_{Z\gamma} & [\%] \ \kappa_c & [\%] \ \kappa_t & [\%] \ \kappa_b & [\%] \end{aligned}$	1.5* 1.3* 2* 1.6* 10* - 3.2* 2.5*	$0.43 \ / \ 0.33$ $0.17 \ / \ 0.14$ $0.90 \ / \ 0.77$ $1.3 \ / \ 1.2$ $10 \ / \ 10$ $1.3 \ / \ 1.1$ $3.1 \ / \ 3.1$ $0.64 \ / \ 0.56$
$\kappa_{\mu} \ [\%]$ $\kappa_{\tau} \ [\%]$ $\text{BR}_{\text{inv}} \ (<\%, 95\% \text{ CL})$ $\text{BR}_{\text{unt}} \ (<\%, 95\% \text{ CL})$	4.4* 1.6* 1.9* 4*	3.9 / 3.7 0.66 / 0.55 0.20 / 0.15 1.0 / 0.88

Table from mid-term report

$$\kappa_X = \frac{g_{hXX}}{g_{hXX}^{\text{SM}}}$$

The stuff we are made of: Ye.

The high luminosity, the precise control of the beam √s, the clean reconstruction of final states make it possible to observe:



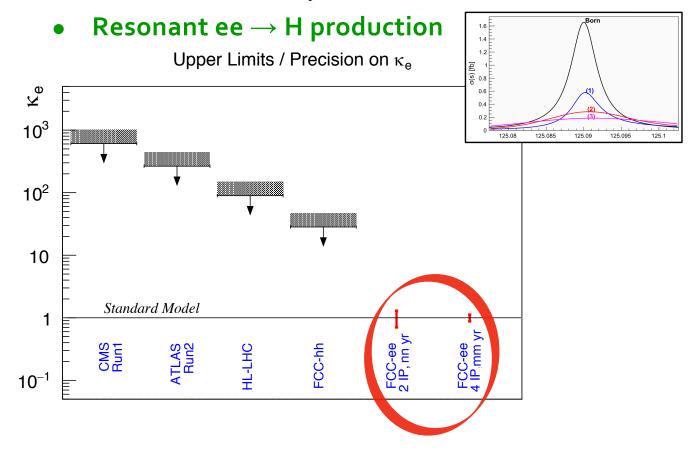
FCC-ee is in the unique position to establish that the Higgs is responsible for the mass of the stable elementary particles ordinary matter is made of.

CG - 30 / 38 May 6, 2024

The stuff we are made of: Ye.

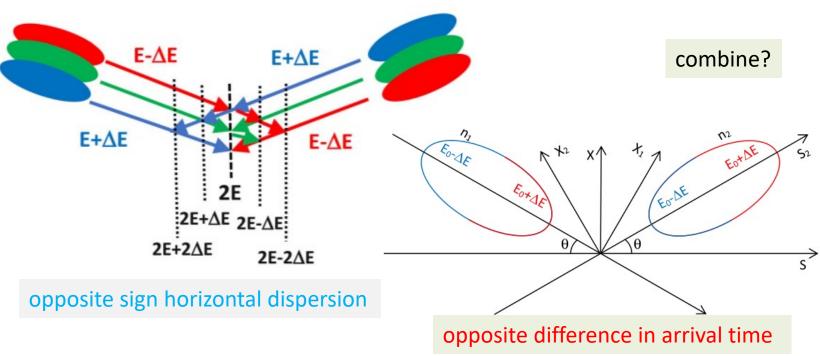
The high luminosity, the precise control of the beam √s, the clean reconstruction of final states make it possible to observe:

- ♦ 20 ab⁻¹/year at \sqrt{s} = 125 GeV (not in baseline FCC-ee)
- Monochromatization $\sigma_{\sqrt{s}} \sim 1-2 \times \Gamma_{H} \sim 6$ to 10 MeV



Monochromatisation

Monochromatization: **UNDER STUDY** taking advantage of the separate e+ and e- rings, one can distribute in opposite way high and low energies in the beam (in x, z time)

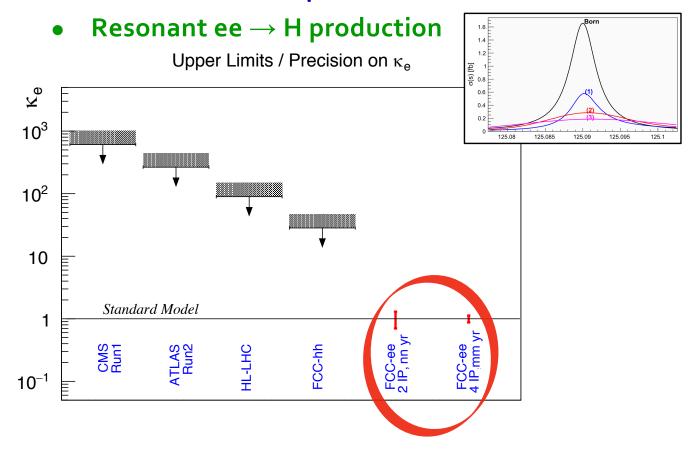


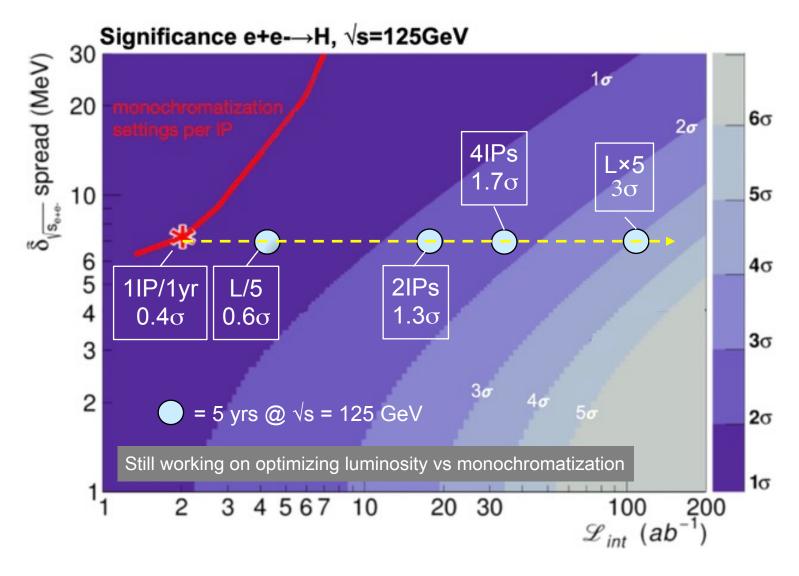
CG - 30 / 38 May 6, 2024

The stuff we are made of: Ye.

The high luminosity, the precise control of the beam √s, the clean reconstruction of final states make it possible to observe:

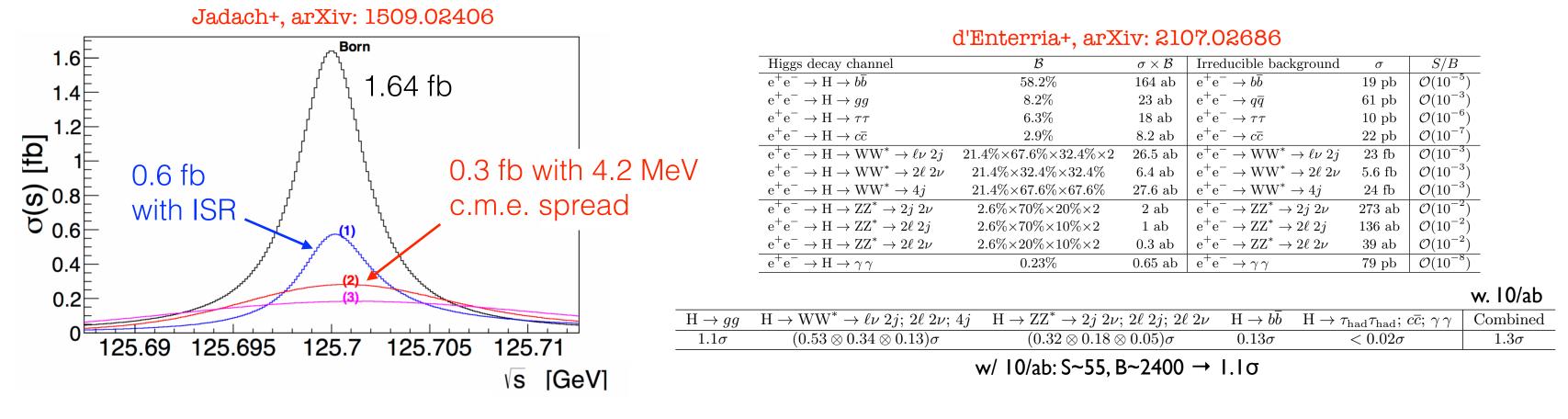
- 20 ab⁻¹ / year at $\sqrt{s} = 125$ GeV (not in baseline FCC-ee)
- Monochromatization $\sigma_{\sqrt{s}} \sim 1-2 \times \Gamma_{H} \sim 6$ to 10 MeV





CG-30/38 May 6, 2024

The stuff we are made of: Y_e.



FCC-ee is in the unique position to establish that the Higgs is responsible for the mass of the stable elementary particles ordinary matter is made of.

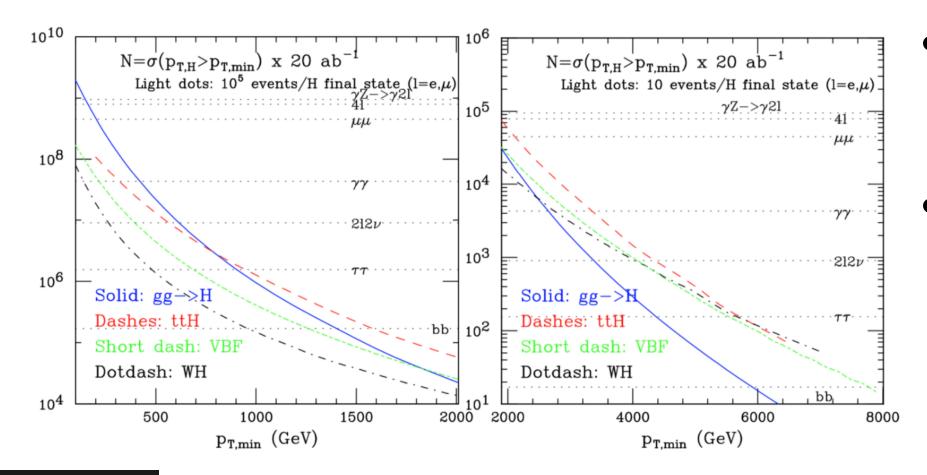
CG - 31 / 38

May 6, 2024

Higgs @ FCC-hh.

	$ggH (N^3LO)$	$\mid \text{VBF (N}^2\text{LO)} \mid$	WH (N^2LO)	$ZH (N^2LO)$	$ t\bar{t}H (N^2LO) $	HH (NLO)
N100	24×10^9	2.1×10^9	4.6×10^{8}	3.3×10^{8}	9.6×10^{8}	3.6×10^{7}
-N100/N14	180	170	100	110	530	390

$$(N100 = \sigma_{100 \text{ TeV}} \times 30 \text{ ab}^{-1} \& N14 = \sigma_{14 \text{ TeV}} \times 3 \text{ ab}^{-1})$$



- Large rate (> 10¹⁰H, > 10⁷ HH)
 - unique sensitivity to rare decays
 - few % sensitivity to self-coupling
- Explore extreme phase space:
 - e.g. 10⁶ H w/ pT>1 TeV
 - clean samples with high S/B
 - small systematics

CG - 32 / 38 May 6, 2024

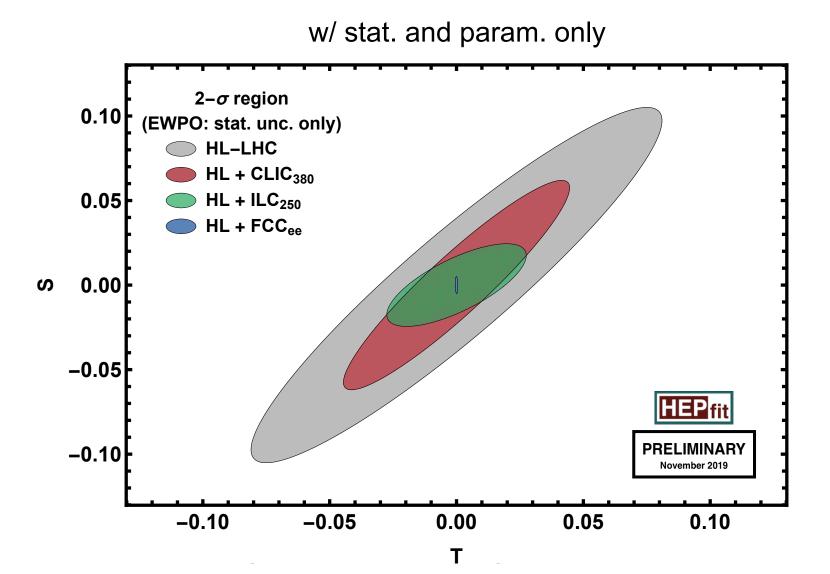
Tera-Z EW precision measurements.

- ▶ The target is to reduce syst. uncertainties to the level of stat. uncertainties. (exploit the large samples and innovative control analyses)
- ▶ Exquisite √s precision (100keV@Z, 300keV@WW) reduces beam uncertainties





~50 times better precision than LEP/LSD on EW precision observables



Indirect sensitivity
to 70TeV-scale sector
connected to EW/Higgs

(For the impact of the theory uncertainties on the EW fit, see bonus slides)

Tera-Z EW precision measurements.

▶ The target is to reduce syst. uncertainties to the level of stat. uncertainties. (exploit the large samples and innovative control analyses)

▶ Exquisite √s precision (100keV@Z, 300keV@WW) reduces beam uncertainties





~50 times better precision than LEP/LSD on EW precision observables

Need TH results to fully exploit Tera-Z

Quantity	Current precision	FCC-ee stat. (syst.) precision	Required theory input	Available calc. in 2019	Needed theory improvement †
$m_{ m Z}$ $\Gamma_{ m Z}$ $\sin^2 heta_{ m eff}^\ell$	$2.1 \mathrm{MeV}$ $2.3 \mathrm{MeV}$ 1.6×10^{-4}	$0.004~(0.1)\mathrm{MeV}$ $0.004~(0.025)\mathrm{MeV}$ $2(2.4)\times10^{-6}$	non-resonant $e^+e^- \rightarrow f\bar{f}$, initial-state radiation (ISR)	NLO, ISR logarithms up to 6th order	NNLO for $e^+e^- \to f\bar{f}$
m_W	$12\mathrm{MeV}$	$0.25~(0.3){ m MeV}$	lineshape of $e^+e^- \rightarrow WW$ near threshold	NLO (ee \rightarrow 4f or EFT framework)	NNLO for $ee \rightarrow WW$, $W \rightarrow ff$ in EFT setup
HZZ coupling		0.2%	cross-sect. for $e^+e^- \to ZH$	NLO + NNLO QCD	NNLO electroweak
$m_{ m top}$	$100\mathrm{MeV}$	17 MeV	threshold scan $e^+e^- \rightarrow t\bar{t}$	N ³ LO QCD, NNLO EW, resummations up to NNLL	Matching fixed orders with resummations, merging with MC, α_s (input)

[†]The listed needed theory calculations constitute a minimum baseline; additional partial higher-order contributions may also be required.

Indirect sensitivity
to 70TeV-scale sector
connected to EW/Higgs

5. FCC-ee as a flavour factory

Flavour potential.

At present (Z/h/NewPhysics) FCNCs mostly constrained by low energy observables. The large statistics of FCC will open on-shell opportunities.

Particle production (10 ⁹)	$B^0 \ / \ \overline{B}^0$	B^+ / B^-	$B_s^0 \ / \ \overline{B}_s^0$	$\Lambda_b \ / \ \overline{\Lambda}_b$	$c\overline{c}$	τ^-/τ^+
Belle II	27.5	27.5	n/a	n/a	65	45
FCC-ee	300	300	80	80	600	150

FCC-ee 10 x Belle II

Mont	Decay mode/Experiment	Belle II $(50/ab)$	LHCb Run I	LHCb Upgr. (50/fb)	FCC-ee
√¶o	$\overline{\mathrm{EW}/H}$ penguins				
•	$B^0 \to K^*(892)e^+e^-$	~ 2000	~ 150	~ 5000	~ 200000
Ω 0	$\mathcal{B}(B^0 \to K^*(892)\tau^+\tau^-)$	~ 10	_	_	~ 1000
See	$B_s o \mu^+\mu^-$	n/a	~ 15	~ 500	~ 800
01	$B^0 o \mu^+ \mu^-$	~ 5	_	~ 50	~ 100
	$\mathcal{B}(B_s o au^+ au^-)$				
	Leptonic decays				
out of reach	$B^+ o \mu^+ \nu_{mu}$	5%	_	_	3%
	$B^+ \to \tau^+ \nu_{tau}$	7%	_	_	2%
at LHCb/Belle	$B_c^+ o au^+ u_{tau}$	n/a	_	-	5%
	CP / hadronic decays				
	$B^0 o J/\Psi K_S \; (\sigma_{\sin(2\phi_d)})$	$\sim 2. * 10^6 (0.008)$	$41500 \ (0.04)$	$\sim 0.8 \cdot 10^6 \ (0.01)$	$\sim 35 \cdot 10^6 \ (0.006)$
	$B_s o D_s^{\pm} K^{\mp}$	n/a	6000	~ 200000	$\sim 30 \cdot 10^6$
	$B_s(B^0) \to J/\Psi \phi \; (\sigma_{\phi_s} \; \mathrm{rad})$	n/a	$96000 \ (0.049)$	$\sim 2.10^6 \ (0.008)$	$16 \cdot 10^6 \ (0.003)$

boosted b's/ τ 's at FCC-ee

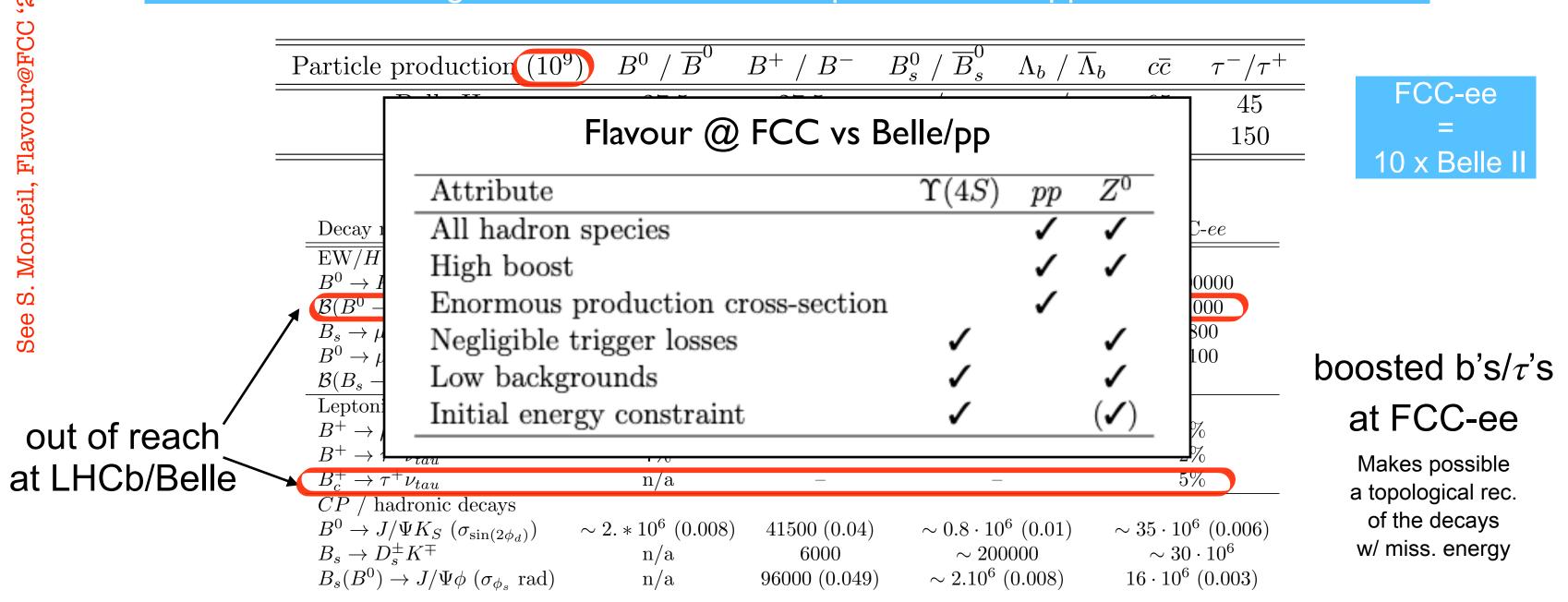
Makes possible a topological rec. of the decays w/ miss. energy

CG - 35 / 38 May 6, 2024

Flavour potential.

At present (Z/h/NewPhysics) FCNCs mostly constrained by low energy observables.

The large statistics of FCC will open on-shell opportunities.



Flavour defines shared (vertexing, tracking, calorimetry) and specific (hadronic PID) detector requirements.

CG - 35 / 38 May 6, 2024

FCC-ee flavour opportunities.

- CKM element V_{cb} (critical for normalising the Unitarity Triangle) from WW decays
- **Tau physics** (>10¹¹ pairs of tau's produced in Z decays)
 - test of lepton flavour universality: G_F from tau decays @ 10 ppm @ FCC-ee (0.5 ppm from muon decays)
 - lepton flavour violation:
 - ► $\tau \rightarrow \mu \gamma$: 4x10⁻⁸ @Belle2021 \rightarrow 10⁻⁹ @ FCC-ee
 - ► $\tau \to 3\mu$: 2x10⁻⁸ @Belle $\to 3$ x10⁻¹⁰ @Belle II $\to 10^{-11}$ @ FCC-ee
 - tau lifetime uncertainty:
 - ▶ 2000 ppm → 10 ppm
 - tau mass uncertainty:
 - ► 70 ppm → 14 ppm
- Semi-leptonic mixing asymmetries as_{sl} and ad_{sl}

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May 6, 2024

6.

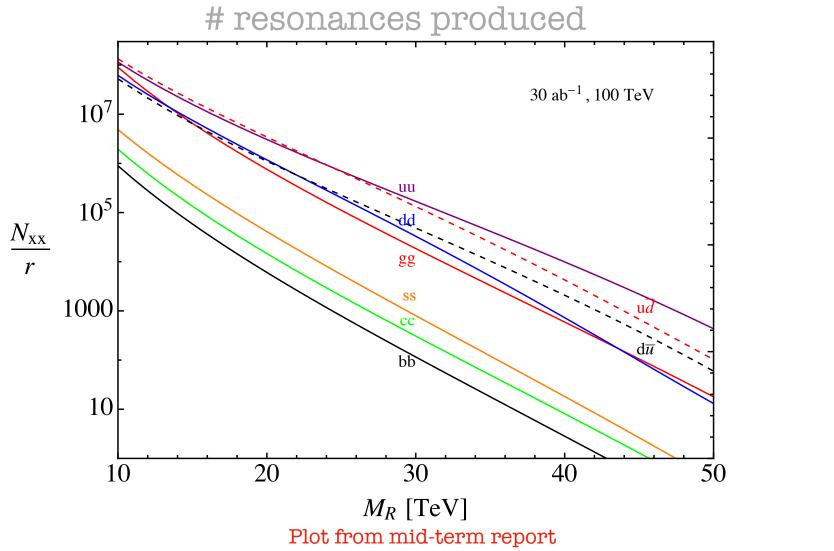
FCC-hh: the broadest exploration potential at high-energy

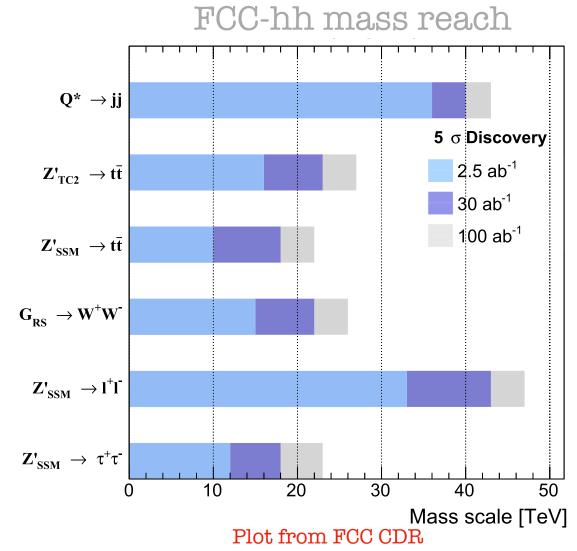
CG - 37 / 38

Resonance production.

Protons are made of 5 quarks, gluons, photons, W/Z

FCC-hh effectively collides 196 different initial states = perfect exploratory machine

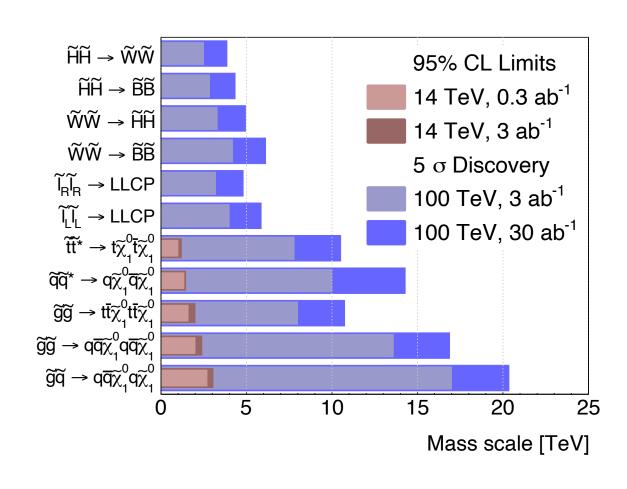




FCC-hh allows the direct exploration of new physics at energy scales up to 40 TeV, including any physics that may be indirectly indicated by precision Higgs and EW measurements at FCC-ee.

CG - 38 / 38

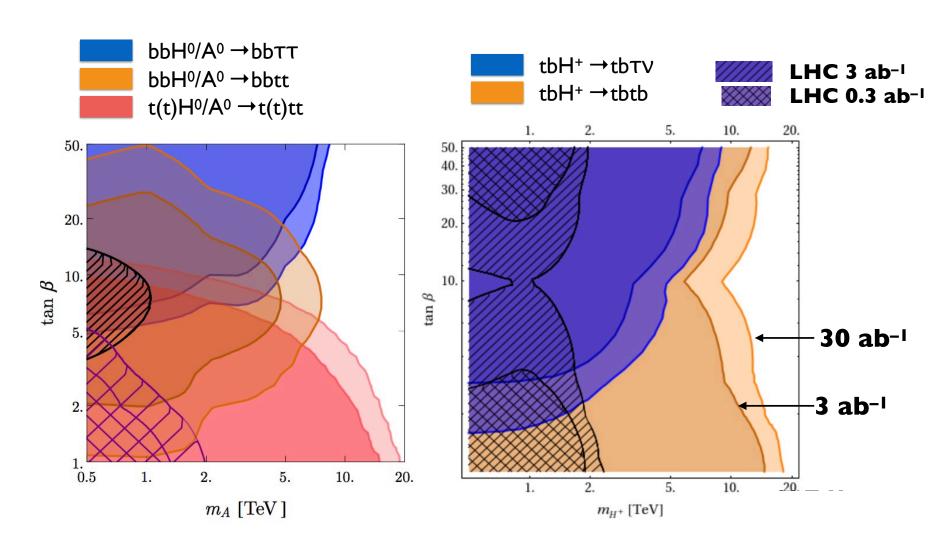
Pushing limits of SUSY.



Plot from arXiv:1606.00947

15-20TeV squarks/gluinos require kinematic threshold 30-40TeV:

FCC-hh is more than a √ŝ~10TeV factory



Plot from arXiv:1605.08744 and arXiv:1504.07617

Factor 10 increase on the HL-LHC limits.

CG - 39 / 38

7. Examples of complementarity & synergy FCC-ee⇔FCC-hh

Synergy ee+hh.

FCC-hh without ee could bound BR_{inv} but it could say nothing about BR_{untagged} (FCC-ee needed for absolute normalisation of Higgs couplings)

0.0 1.5 3.0 4.5 6.0 7.5

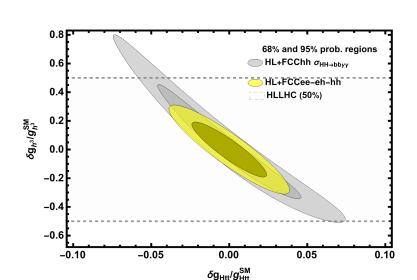
FCC-hh is determining top Yukawa through ratio tth/ttZ

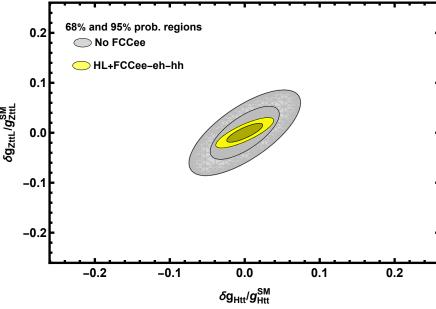
So the extraction of top Yukawa heavily relies on the knowledge of ttZ from FCC-ee

Mangano+'15							
	$\sigma(t \bar{t} H) [ext{pb}]$	$\sigma(t \bar{t} Z) [ext{pb}]$	$rac{\sigma(tar{t}H)}{\sigma(tar{t}Z)}$				
13 TeV	$0.475^{+5.79\%}_{-9.04\%} + 3.33\%$	$0.785^{+9.81\%}_{-11.2\%}{}^{+3.27\%}_{-3.12\%}$	$0.606^{+2.45\%}_{-3.66\%}{}^{+0.525\%}_{-0.319\%}$				
100 TeV	$33.9^{+7.06\%}_{-8.29\%}^{+2.17\%}_{-2.18\%}$	$57.9^{+8.93\%+2.24\%}_{-9.46\%-2.43\%}$	$0.585^{+1.29\%}_{-2.02\%}{}^{+0.314\%}_{-0.147\%}$				

(uncertainty drops in ratio)

Subsequently, the 1% sensitivity on tth is essential to determine h³ at O(5%) at FCC-hh





Plots from mid-term report

FCC-hh tunnel is great for FCC-ee.

- 80-100 km is needed to accelerate pp up to 100 TeV
- 80-100 km is also exactly what is needed
 - to get enough luminosity (5 times more than in 27 km) to maybe get sensitivity to the Higgs self coupling, the electron Yukawa coupling, or sterile neutrinos,
 - to make TeraZ a useful flavour factory,
 - for transverse polarisation to be available all the way to the WW threshold (allowing a precise W mass measurement)
 - for the top threshold to be reached and exceeded.

CG - 42 / 38

Conclusions & Outlook

A circular "Higgs factory" like FCC-ee has a rich potential:

- * Search directly and indirectly for New Physics
- ★ Establish new organising principles of Nature (LEP→ gauge symmetries, FCC→??)
- * Probe the **HEP-Cosmo connections** thanks to the high statistics of the **Z-pole run** (omitting this exploration would be ignoring the outcome of LHC).

And FCC-ee is an essential part of an integrated programme to probe the energy frontier.

We have profound questions and we need to create opportunities to answer them.

FCC will for sure contribute.

We can learn a lot from nice pictures/observations but experiments remain the driver of physics.

Colliders are the most powerful microscopes we have to study Nature at the smallest scales and also from the early moments of the Universe. They'll keep providing a quantitative understanding to progress forward.

CG - 43 / 38 May 6, 2024

Acknowledgement.

This project is supported from the European Union's Horizon 2020 research and innovation programme under grant agreement No 951754.



CG - 44 / 38

BONUS

Some work ahead of us.

- Development of a common software and the estimate of the computing needs
- Evaluation of the physics performance and requirements for detectors
- Conceptualisation of detectors capable of delivering these requirements
- Mitigation of the interaction region constraints on detectors and vice versa
- Design of methods and tools for centre-of-mass energy calibration, beam polarisation, and monochromatization
- Understanding and optimisation of the physics programm
- Exploration of the physics opportunities
- Development of the theoretical tools and observables needed to meet the measurement targets

May 6, 2024

Higgs and EW measurements

CG -47 / 38

Experimental Inputs.

A circular ee Higgs factory starts as a Z/EW factory (**TeraZ**)

A linear ee Higgs factory operating above Z-pole can also preform EW measurements via **Z-radiative** return

A linear ee Higgs factory could also operate on the Z-pole though at lower lumi (**GigaZ**)

	Higgs	aTGC	EWPO	Top EW
FCC-ee	Yes (μ, σ _{ZH}) (Complete with HL-LHC)	Yes (aTGC dom.) Warning	Yes	Yes (365 GeV, Ztt)
ILC	Yes (μ, σ _{ZH}) (Complete with HL-LHC)	Yes (HE limit) Warning	LEP/SLD (Z-pole) + HL-LHC + W (ILC)	Yes (500 GeV, Ztt)
CEPC	Yes (μ, σ _{ZH}) (Complete with HL-LHC)	Yes (aTGC dom) Warning	Yes	No
CLIC	Yes (μ, σ _{ZH})	Yes (Full EFT parameterization)	LEP/SLD (Z-pole) + HL-LHC + W (CLIC)	Yes
HE-LHC	Extrapolated from HL-LHC	N/A → LEP2	LEP/SLD + HL-LHC (M _W , sin ² θ _w)	-
FCC-hh	Yes (µ, BR _i /BR _j) Used in combination with FCCee/eh	From FCC-ee	From FCC-ee	-
LHeC	Yes (μ)	N/A → LEP2	LEP/SLD + HL-LHC (M _W , sin ² θ _w)	-
FCC-eh	Yes (µ) Used in combination with FCCee/hh	From FCC-ee	From FCC-ee + Zuu, Zdd	-

CG - 48 / 38

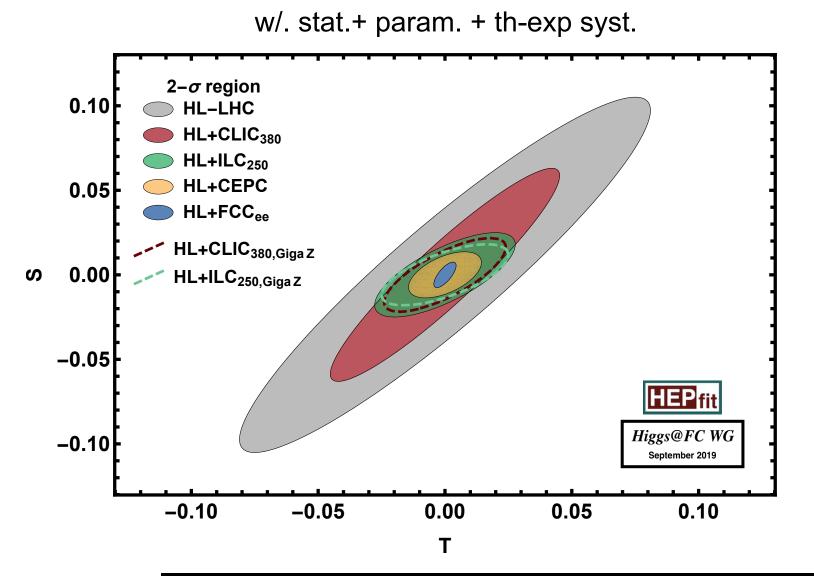
m_Z (keV) 91186700 ± 2200 4 100 From Z line shape sean Beam energy calibration Γ_Z (keV) 2495200 ± 2300 4 25 From Z line shape sean search geam energy calibration $\sin^2 \theta_W^{eff}(\times 10^6)$ 231480 ± 160 2 2.4 From A_{EB}^{eff} at Z peak Beam energy calibration $1/\alpha_{\rm QED}(m_Z^2)(\times 10^3)$ 128952 ± 14 3 small From A_{EB}^{eff} at Z peak $\alpha_{\rm CPD}$ $\alpha_{\rm CP$	Observable	value	preser ±	nt error	FCC-ee Stat.	FCC-ee Syst.	Comment and leading erro
	$m_{\mathrm{Z}} (\mathrm{keV})$	91186700		2200		-	From Z line shape scar
Beam energy calibratio $1/\alpha_{\rm QED}({\rm m_Z^2})(\times 10^3)$ 128952 \pm 14 3 small From Λμμ of Peach of	$\Gamma_{\rm Z}~({\rm keV})$	2495200	±	2300	4	25	_
$R_{\ell}^{2} (\times 10^{3}) \qquad 20767 \qquad \pm \qquad 25 \qquad \textbf{0.06} \qquad 0.2-1 \qquad \text{Ratio of hadrons to lepton } \\ \alpha_{s}(m_{Z}^{2}) (\times 10^{4}) \qquad 1196 \qquad \pm \qquad 30 \qquad \textbf{0.1} \qquad 0.4-1.6 \qquad \textbf{From } R_{l}^{2} \\ \sigma_{had}^{0} (\times 10^{3}) (\text{nb}) \qquad 41541 \qquad \pm \qquad 37 \qquad \textbf{0.1} \qquad 4 \qquad \textbf{Peak hadronic cross-section } \\ Luminosity measuremen \\ N_{\nu} (\times 10^{3}) \qquad 2996 \qquad \pm \qquad 7 \qquad \textbf{0.005} \qquad 1 \qquad Z \text{peak cross-section } \\ Luminosity measuremen \\ R_{b} (\times 10^{6}) \qquad 216290 \qquad \pm \qquad 660 \qquad \textbf{0.3} \qquad < 60 \qquad \textbf{Ratio of bib to hadron } \\ Stat. \text{extrapol, from SLI} \\ A_{FB}^{b}, 0 (\times 10^{4}) \qquad 992 \qquad \pm \qquad 16 \qquad \textbf{0.02} \qquad 1-3 \qquad \textbf{b-quark asymmetry at } Z \text{pol} \\ From \text{jet charg} \\ A_{FB}^{\text{pol,}\tau} (\times 10^{4}) \qquad 1498 \qquad \pm \qquad 49 \qquad \textbf{0.15} \qquad < 2 \qquad \tau \text{polarization asymmetry } \\ \tau \text{decay physic} \\ \tau \text{lifetime (fs)} \qquad 290.3 \qquad \pm \qquad 0.5 \qquad \textbf{0.001} \qquad 0.04 \qquad \textbf{Radial alignmen} \\ \tau \text{mass (MeV)} \qquad 1776.86 \qquad \pm \qquad 0.12 \qquad \textbf{0.004} \qquad 0.04 \qquad \textbf{Momentum scal} \\ \tau \text{leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{B.R. (\%)} \qquad 17.38 \qquad \pm \qquad 0.04 \qquad \textbf{0.0001} \qquad 0.003 \qquad \text{e/p/hadron separation} \\ m_{W} (\text{MeV}) \qquad 2085 \qquad \pm \qquad 15 \qquad \textbf{0.25} \qquad 0.3 \qquad \text{From WW threshold scan} \\ Beam \text{energy calibration} \\ \alpha_{s}(m_{W}^{2})(\times 10^{4}) \qquad 1010 \qquad \pm \qquad 270 \qquad \textbf{3} \qquad \text{small} \qquad \textbf{Ratio of invis. to leptoni} \\ n_{V} (\times 10^{3}) \qquad 2920 \qquad \pm \qquad 50 \qquad \textbf{0.8} \qquad \text{small} \qquad \textbf{Ratio of invis. to leptoni} \\ n_{Top} (\text{MeV}) \qquad 172740 \qquad \pm \qquad 500 \qquad 17 \qquad \text{small} \qquad \textbf{From } t\bar{t} \text{threshold scan} \\ QCD \text{errors dominat} \\ \lambda_{Lop} \lambda_{Lop}^{\text{SM}} \qquad 1.2 \qquad \pm \qquad 0.3 \qquad \textbf{0.10} \qquad \text{small} \qquad \textbf{From } t\bar{t} \text{threshold scan} \\ QCD \text{errors dominat} \\ \lambda_{Lop} \lambda_{Lop}^{\text{SM}} \qquad 1.2 \qquad \pm \qquad 0.3 \qquad \textbf{0.10} \qquad \text{small} \qquad \textbf{From } t\bar{t} \text{threshold scan} \\ QCD \text{errors dominat} \\ \lambda_{Lop} \lambda_{Lop}^{\text{SM}} \qquad 1.2 \qquad \pm \qquad 0.3 \qquad \textbf{0.10} \qquad \text{small} \qquad \textbf{From } t\bar{t} \text{threshold scan} \\ QCD \text{errors dominat} \\ \lambda_{Lop} \lambda_{Lop}^{\text{SM}} \qquad 1.2 \qquad \pm \qquad 0.3 \qquad \textbf{0.10} \qquad \text{small} \qquad \textbf{From } t\bar{t} \text{threshold scan} \\ QCD \text{errors dominat} \\ \lambda_{Lop} \lambda_{Lop}^{\text{SM}} \qquad 1.2 \qquad \pm \qquad 0.3 \qquad \textbf{0.10} \qquad \text{small} \qquad \textbf{From } t\bar{t} threshold $	$\sin^2 \theta_{\rm W}^{\rm eff}(\times 10^6)$	231480	±	160	2	2.4	
$\frac{\alpha_s(m_Z^2) \left(\times 10^4\right)}{\alpha_s(m_Z^2) \left(\times 10^4\right)} \qquad 1196 \pm 30 \qquad \textbf{0.1} \qquad 0.4\text{-}1.6 \qquad \textbf{From R}; \\ \frac{\sigma_{\text{had}}^0 \left(\times 10^3\right) \left(\text{nb}\right)}{\sigma_{\text{had}}^0 \left(\times 10^3\right) \left(\text{nb}\right)} \qquad 41541 \pm 37 \qquad \textbf{0.1} \qquad 4 \qquad \textbf{Peak hadronic cross-section Luminosity measuremen} \\ N_{\nu}(\times 10^3) \qquad 2996 \pm 7 \qquad \textbf{0.005} \qquad 1 \qquad \textbf{Z peak cross-section Luminosity measuremen} \\ R_b \left(\times 10^6\right) \qquad 216290 \pm 660 \qquad \textbf{0.3} \qquad < 60 \qquad \textbf{Ratio of bb to hadron Stat. extrapol. from SLI} \\ A_{\text{PB}}^b, 0 \left(\times 10^4\right) \qquad 992 \pm 16 \qquad \textbf{0.02} \qquad 1\text{-}3 \qquad \textbf{b-quark asymmetry at Z pol From jet charg} \\ A_{\text{PB}}^{\text{pol,}\tau} \left(\times 10^4\right) \qquad 1498 \pm 49 \qquad \textbf{0.15} \qquad < 2 \qquad \tau \text{ polarization asymmetr} \\ \tau \text{ decay physic} \\ \tau \text{ lifetime (fs)} \qquad 290.3 \pm 0.5 \textbf{0.001} \qquad 0.04 \qquad \textbf{Radial alignmen} \\ \tau \text{ mass (MeV)} \qquad 1776.86 \pm 0.12 \qquad \textbf{0.004} \qquad 0.04 \qquad \textbf{Momentum scal} \\ \tau \text{ leptonic } \left(\mu\nu_{\mu}\nu_{\tau}\right) \text{ B.R. (\%)} \qquad 17.38 \pm 0.04 \textbf{0.0001} \qquad 0.003 \qquad e/p/\text{hadron separation} \\ m_W \left(\text{MeV}\right) \qquad 80350 \pm 15 \qquad \textbf{0.25} \qquad 0.3 \qquad \textbf{From WW threshold scal Beam energy calibration} \\ \Gamma_W \left(\text{MeV}\right) \qquad 2085 \pm 42 \qquad \textbf{1.2} \qquad 0.3 \qquad \textbf{From WW threshold scal Beam energy calibration} \\ n_{\text{V}}\left(\text{MeV}\right) \qquad 1010 \pm 270 \qquad \textbf{3} \text{small} \qquad \textbf{Ratio of invis. to leptonic in radiative Z return} \\ m_{\text{top}} \left(\text{MeV}\right) \qquad 172740 \pm 500 \qquad \textbf{17} \qquad \text{small} \qquad \textbf{From tit threshold scal QCD errors dominat} \\ \Gamma_{\text{top}} \left(\text{MeV}\right) \qquad 1410 \pm 190 \qquad \textbf{45} \qquad \text{small} \qquad \textbf{From tit threshold scal QCD errors dominat} \\ n_{\text{top}} \left(\text{MeV}\right) \qquad 1410 \pm 190 \qquad \textbf{45} \qquad \text{small} \qquad \textbf{From tit threshold scal QCD errors dominat} \\ n_{\text{top}} \left(\text{MeV}\right) \qquad 1410 \pm 190 \qquad \textbf{45} \qquad \text{small} \qquad \textbf{From tit threshold scal QCD errors dominat} \\ n_{\text{top}} \left(\text{MeV}\right) \qquad 1410 \pm 190 \qquad \textbf{45} \qquad \text{small} \qquad \textbf{From tit threshold scal QCD errors dominat} \\ n_{\text{top}} \left(\text{MeV}\right) \qquad 1410 \pm 190 \qquad \textbf{45} \qquad \text{small} \qquad \textbf{From tit threshold scal QCD errors dominat} \\ n_{\text{top}} \left(\text{MeV}\right) \qquad 1410 \pm 190 \qquad \textbf{45} \qquad \text{small} \qquad \textbf{From tit threshold scal QCD errors dominat} \\ n_{\text{top}} \left(\text{MeV}\right$	$\frac{1/\alpha_{\rm QED}(m_{\rm Z}^2)(\times 10^3)}{1/\alpha_{\rm QED}(m_{\rm Z}^2)(\times 10^3)}$	128952	±	14	3	small	
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$R_{\ell}^{\rm Z} \ (\times 10^3)$	20767	±	25	0.06	0.2-1	-
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\frac{\alpha_{\rm s}(m_{\rm Z}^2) \ (\times 10^4)}{\alpha_{\rm s}(m_{\rm Z}^2) \ (\times 10^4)}$	1196	土	30	0.1	0.4-1.6	From R
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\sigma_{\rm had}^0 \ (\times 10^3) \ ({\rm nb})$	41541	±	37	0.1	4	
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$N_{\nu}(\times 10^3)$	2996	土	7	0.005	1	=
From jet charg $A_{\mathrm{FB}}^{\mathrm{pol},\tau} \ (\times 10^4)$ 1498 \pm 49 0.15 $<$ 2 τ polarization asymmetr τ decay physic τ lifetime (fs) 290.3 \pm 0.5 0.001 0.04 Radial alignment τ mass (MeV) 1776.86 \pm 0.12 0.004 0.04 Momentum scal τ leptonic ($\mu\nu_{\mu}\nu_{\tau}$) B.R. (%) 17.38 \pm 0.04 0.0001 0.003 e/μ /hadron separation $\mathrm{m_W} \ (\mathrm{MeV})$ 80350 \pm 15 0.25 0.3 From WW threshold scan Beam energy calibration $\sigma_{\mathrm{s}} (\mathrm{m_W}^2) (\times 10^4)$ 1010 \pm 270 3 small From WW threshold scan Beam energy calibration $\sigma_{\mathrm{s}} (\mathrm{m_W}^2) (\times 10^4)$ 1010 \pm 270 3 small Ratio of invis. to leptoni in radiative Z return $\mathrm{m_{top}} \ (\mathrm{MeV})$ 172740 \pm 500 17 small From tt threshold scan QCD errors dominat $\Gamma_{\mathrm{top}} \ (\mathrm{MeV})$ 1410 \pm 190 45 small From tt threshold scan QCD errors dominat $\Lambda_{\mathrm{top}}/\Lambda_{\mathrm{top}}^{\mathrm{SM}}$ 1.2 \pm 0.3 0.10 small From tt threshold scan QCD errors dominat	$R_{\rm b} \ (\times 10^6)$	216290	土	660	0.3	< 60	
$\frac{\tau \text{ decay physics}}{\tau \text{ lifetime (fs)}} \frac{\tau \text{ lifetime (fs)}}{290.3} \pm 0.5 \textbf{0.001} 0.04 \text{Radial alignment}}{\tau \text{ mass (MeV)}} \frac{\tau \text{ lifetime (fs)}}{1776.86} \pm 0.12 \textbf{0.004} 0.04 \text{Momentum scal}}{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ B.R. (\%)}} \frac{17.38}{17.38} \pm 0.04 \textbf{0.0001} 0.003 \text{e/p/hadron separation}}{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ B.R. (\%)}} \frac{17.38}{17.38} \pm 0.04 \textbf{0.0001} 0.003 \text{e/p/hadron separation}}{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ B.R. (\%)}} \frac{17.38}{15} \pm 15 \textbf{0.25} 0.3 \text{From WW threshold scal Beam energy calibration}}{\tau \text{ Beam energy calibration}} \frac{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ B.R. (\%)}}{\tau \text{ log of invision}} \frac{10.00 \pm 270 3 \text{small}}{\tau \text{ small}} \frac{\tau \text{ From WW threshold scal Peam energy calibration}}{\tau \text{ log of invision}} \frac{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \frac{\tau \text{ log of invision}}{\tau \text{ log of invision}} \tau \text{ log of inv$	$A_{\rm FB}^{\rm b}, 0 \ (\times 10^4)$	992	±	16	0.02	1-3	
$\frac{\tau \; \text{mass (MeV)}}{\tau \; \text{mass (MeV)}} 1776.86 \; \pm \; 0.12 \textbf{0.004} 0.04 \text{Momentum scal}}{\tau \; \text{leptonic } (\mu\nu_{\mu}\nu_{\tau}) \; \text{B.R. (\%)}} 17.38 \; \pm \; 0.04 \textbf{0.0001} 0.003 \text{e/μ/hadron separation}}{\tau \; \text{leptonic } (\mu\nu_{\mu}\nu_{\tau}) \; \text{B.R. (\%)}} 17.38 \; \pm \; 0.04 \textbf{0.0001} 0.003 \text{e/μ/hadron separation}}{\tau \; \text{mw} \; \text{MeV}} 10.25 0.3 \text{From WW threshold scal}} \frac{\tau \; \text{Beam energy calibration}}{\tau \; \text{Beam energy calibration}} \frac{\tau \; \text{MeV}}{\tau \; \text{Beam energy calibration}} \frac{\tau \; \text{Single of invis. to leptonic in radiative Z return}}{\tau \; \text{MeV}} 10.10 \; \pm \; 2.085 \; \pm \; 4.2 1.2 0.3 \text{Single of invis. to leptonic in radiative Z return}} \frac{\tau \; \text{MeV}}{\tau \; \text{MeV}} 10.10 \; \pm \; 5.00 1.7 \text{Single of invis. to leptonic in radiative Z return}} \frac{\tau \; \text{MeV}}{\tau \; \text{MeV}} 10.10 \; \pm \; 10.0 10.0 \text{Single of invis. to leptonic in radiative Z return}} \frac{\tau \; \text{MeV}}{\tau \; \text{MeV}} 10.00 \; \frac{\tau \; \text{MeV}}{\tau \; MeV$	$\overline{A_{FB}^{\text{pol},\tau} (\times 10^4)}$	1498	土	49	0.15	<2	
$\frac{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ B.R. } (\%)}{\text{mw}} 17.38 \pm 0.04 \textbf{0.0001} 0.003 \text{e/μ/hadron separation} \\ \frac{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ B.R. } (\%)}{\text{mw}} 18.350 \pm 15 \textbf{0.25} 0.3 \text{From WW threshold scar} \\ \frac{\tau \text{ Beam energy calibration}}{\text{Beam energy calibration}} \\ \frac{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ B.R. } (\%)}{\text{MeV}} 2085 \pm 42 \textbf{1.2} 0.3 \text{From WW threshold scar} \\ \frac{\tau \text{ Beam energy calibration}}{\text{Beam energy calibration}} \\ \frac{\tau \text{ leptonic } (\mu\nu_{\mu}\nu_{\tau}) \text{ B.R. } (\%)}{\text{Beam energy calibration}} \\ \frac{\tau \text{ leptonic } (\mu\nu_{\tau}) lept$	τ lifetime (fs)	290.3	±	0.5	0.001	0.04	Radial alignmen
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\tau \text{ mass (MeV)}$	1776.86	土	0.12	0.004	0.04	Momentum scal
$\Gamma_{\rm W} ({\rm MeV}) \qquad \qquad 2085 \pm 42 \qquad \qquad \textbf{1.2} \qquad \qquad \textbf{0.3} \qquad \begin{array}{c} \text{From WW threshold scal} \\ \text{Beam energy calibration} \\ \alpha_{\rm S} (m_{\rm W}^2) (\times 10^4) \qquad \qquad 1010 \pm 270 \qquad \textbf{3} \qquad \text{small} \qquad \begin{array}{c} \text{From R}_{\ell}^{\rm V} \\ \text{N}_{\nu} (\times 10^3) \qquad \qquad 2920 \pm 50 \qquad \textbf{0.8} \qquad \text{small} \qquad \begin{array}{c} \text{Ratio of invis. to leptoni} \\ \text{in radiative Z return} \\ \text{M}_{\rm top} ({\rm MeV}) \qquad \qquad 172740 \pm 500 \qquad \textbf{17} \qquad \text{small} \qquad \begin{array}{c} \text{From t\bar{t} threshold scal} \\ \text{QCD errors dominat} \\ \text{QCD errors dominat} \\ \\ \lambda_{\rm top} / \lambda_{\rm top}^{\rm SM} \qquad \qquad 1.2 \pm 0.3 \qquad \textbf{0.10} \qquad \text{small} \qquad \begin{array}{c} \text{From t\bar{t} threshold scal} \\ \text{QCD errors dominat} \\ \end{array}$	$\overline{\tau}$ leptonic $(\mu\nu_{\mu}\nu_{\tau})$ B.R. $(\%)$	17.38	土	0.04	0.0001	0.003	e/µ/hadron separation
Beam energy calibration $\alpha_{\rm s}({\rm m_W^2})(\times 10^4)$ 1010 \pm 270 ${\bf 3}$ small From R _{ℓ} N _{ν} ($\times 10^3$) 2920 \pm 50 ${\bf 0.8}$ small Ratio of invis. to leptoni in radiative Z return m _{top} (MeV) 172740 \pm 500 ${\bf 17}$ small From tt threshold scalar QCD errors dominate $\Gamma_{\rm top}$ (MeV) 1410 \pm 190 ${\bf 45}$ small From tt threshold scalar QCD errors dominate $\lambda_{\rm top}/\lambda_{\rm top}^{\rm SM}$ 1.2 \pm 0.3 ${\bf 0.10}$ small From tt threshold scalar QCD errors dominate $\lambda_{\rm top}/\lambda_{\rm top}^{\rm SM}$ 200 \pm 200 \pm 300	$m_{ m W}~({ m MeV})$	80350	±	15	0.25	0.3	
$N_{\nu}(\times 10^3)$ 2920 \pm 50 0.8 small Ratio of invis. to leptoni in radiative Z return $m_{\rm top}$ (MeV) 172740 \pm 500 17 small From $t\bar{t}$ threshold scan QCD errors dominat $\Gamma_{\rm top}$ (MeV) 1410 \pm 190 45 small From $t\bar{t}$ threshold scan QCD errors dominat $\lambda_{\rm top}/\lambda_{\rm top}^{\rm SM}$ 1.2 \pm 0.3 0.10 small From $t\bar{t}$ threshold scan QCD errors dominat	$\Gamma_{ m W} \ ({ m MeV})$	2085	±	42	1.2	0.3	
$\frac{\text{in radiative Z return}}{\text{m}_{\text{top}} \; (\text{MeV})} = \frac{172740}{12740} \; \pm \; 500 = \frac{17}{12740} \; \text{small} \qquad \frac{\text{From } t\bar{t} \; \text{threshold scan}}{\text{QCD errors dominat}} = \frac{\Gamma_{\text{top}} \; (\text{MeV})}{\Gamma_{\text{top}} \; (\text{MeV})} = \frac{1410}{12740} \; \pm \; 190 = \frac{120}{12740} = \frac$	$\frac{\alpha_s(m_W^2)(\times 10^4)}{\alpha_s(m_W^2)(\times 10^4)}$	1010	±	270	3	small	From R_ℓ^V
$\Gamma_{ m top} ({ m MeV})$ 1410 \pm 190 45 small From $t\bar{t}$ threshold scalar QCD errors dominate $\lambda_{ m top}/\lambda_{ m top}^{ m SM}$ 1.2 \pm 0.3 0.10 small From $t\bar{t}$ threshold scalar QCD errors dominate $\lambda_{ m top}/\lambda_{ m top}^{ m SM}$	$N_{\nu}(\times 10^3)$	2920	土	50	0.8	small	•
$\lambda_{ m top}/\lambda_{ m top}^{ m SM}$ 1.2 \pm 0.3 0.10 small From ${ m t\bar{t}}$ threshold scar QCD errors dominat	$m_{\mathrm{top}} \; (\mathrm{MeV})$	172740	±	500	17	small	
QCD errors dominat	$\Gamma_{\rm top} ({ m MeV})$	1410	±	190	45	small	
ttZ couplings \pm 30% 0.5 – 1.5 % small From $\sqrt{s} = 365\mathrm{GeV}$ rus	$\lambda_{ m top}/\lambda_{ m top}^{ m SM}$	1.2	±	0.3	0.10	small	
	ttZ couplings		±	30%	0.5 - 1.5 %	small	From $\sqrt{s} = 365 \mathrm{GeV}$ rus

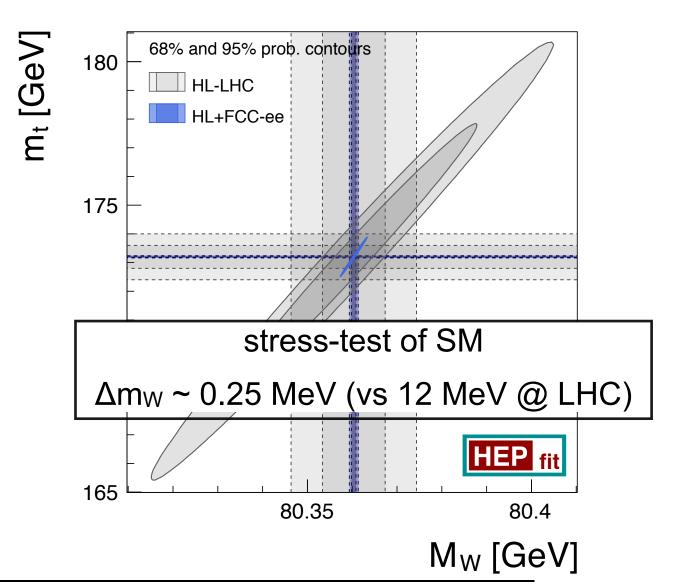
EW Precision Measurements at FCC-ee

Table from mid-term report

Improvements of EW measurements

Exquisite measurements of m_Z (100 keV), Γ_Z (25 keV), m_W (<500 keV), $\alpha_{QED}(m_Z)$ (3.10-5) (all unique to FCC-ee)





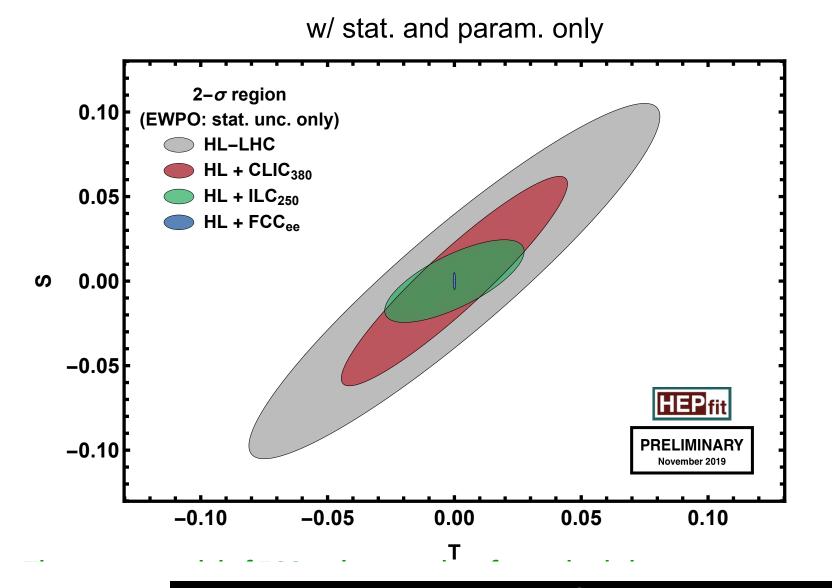
The importance of improved EW measurements is threefold:

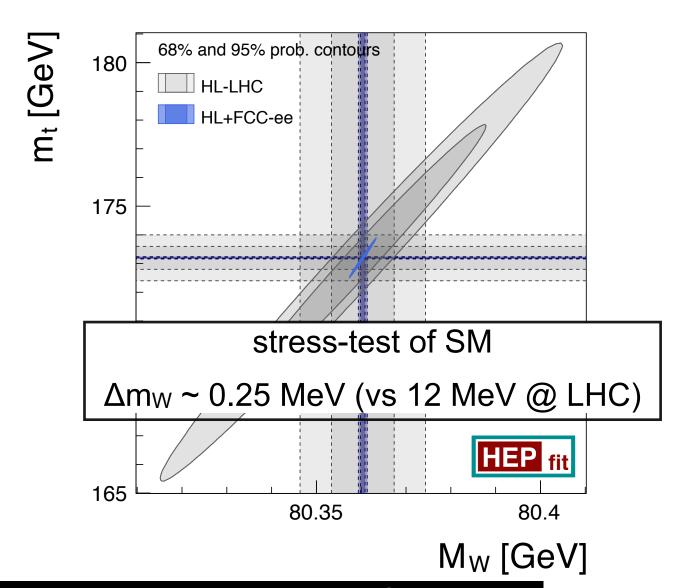
- 1) improve mass reach in indirect search for NP (S~10-2 → M~70 TeV)
 - 2) reduced parametric uncertainties for other measurements
 - 3) reduced degeneracies in a global fit for Higgs couplings

CG-50 / 38

Improvements of EW measurements

Exquisite measurements of m_Z (100 keV), Γ_Z (25 keV), m_W (<500 keV), $\alpha_{QED}(m_Z)$ (3.10-5) (all unique to FCC-ee)





The importance of improved EW measurements is threefold:

- 1) improve mass reach in indirect search for NP (S~10-2 → M~70 TeV)
 - 2) reduced parametric uncertainties for other measurements
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CG-50 / 38

Systematics vs. Statistics.

PED @ CERN-SPC '2022

We often hear that more Z pole statistics is useless, because they are systematics-limited

FCC-ee

- This is a passive attitude, which leads to pessimistic expectations and wrong conclusions/planning
 - Experience shows that a careful experimental systematic analysis boils down to a statistical problem
 - If well prepared, theory will go as far as deemed useful: this preparation starts today (and needs SUPPORT)
 - We are working in the spirit of matching systematic errors to expected statistics for all precision measurements
- Take the Z lineshape

$\alpha_{\text{QED}}(\mathbf{m}_{\text{Z}})$: Stat. 3×10^{-5}



- Enters as a limiting parametric uncertainties in the new physics interpretation many past and future measurements.
- Is statistics limited and will directly benefit from more luminosity
- No useful impact on $\alpha_{QED}(m_z)$ with five times less luminosity
 - ♦ Most of the work is (will be) on systematics
 - But huge statistics will turn into better precision
 - → A real chance for discovery

 $\sin^2\theta_W^{eff}$ and Γ_Z (also m_W vs m_Z): Stat. 2×10⁻⁶ and 4 keV Error dominated by point-to-point energy uncertainties.

Based on in-situ comparisons between √s (e.g. with muon pairs), with measurements made every few minutes (100's times per day)

Boils down to

- statistics (the more data the better, scales down as $1/\sqrt{L}$)
- detector systematics (uncorrelated between experiments, scales down a $1/\sqrt{N_{\text{experiments}}}$)

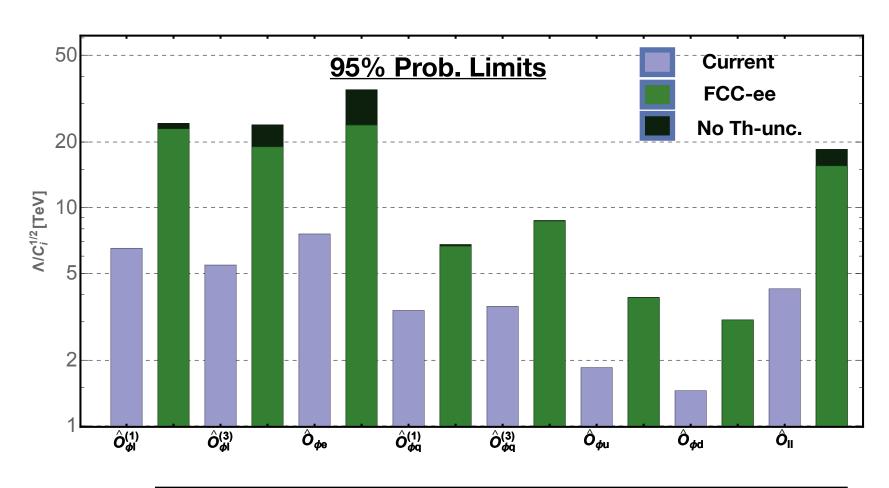
Z (and W) mass: Stat. 4 keV (250 keV)

Error dominated by \sqrt{s} determination with resonant depolarization. As more understanding is gained, progress are made at a constant pace, and this error approaches regularly the statistical limit

CG - 51 / 38

Impact of TH uncertainties.

J. de Blas, FCC CDR overview '19



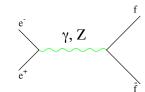
	Cur	rent	\mathbf{FCCee}				
	Exp.	\mathbf{SM}	Exp.	SM (par.)	SM (th.)		
$\overline{\delta M_W \; [{ m MeV}]}$	±15	±8	<u>±1</u>	$\pm 0.6/\pm 1$	<u>±1</u>		
$\delta \Gamma_Z \; [{ m MeV}]$	± 2.3	± 0.73	± 0.1	± 0.1	± 0.2		
$\delta \mathcal{A}_\ell \left[imes 10^{-5} ight]$	± 210	± 93	± 2.1	$\pm 8/\pm 14$	± 11.8		
$\delta R_b^0 \left[imes 10^{-5} ight]$	± 66	± 3	± 6	± 0.3	± 5		

CG - 52 / 38

2021 week Wilkinson

Some EW measurements @ Tera

measure $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$ and $A_{FB}^{\mu\mu}$ at (a) judicious \sqrt{s}

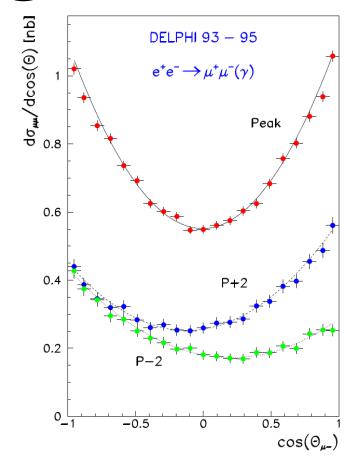


The γ exchange term is proportional to $\alpha^2_{OFD}(\sqrt{s})$ The Z exchange term is proportional to G_{F}^2 , hence independent of α_{OFD} The γZ interference is proportional to $\alpha_{OED}(\sqrt{s}) \times G_F$

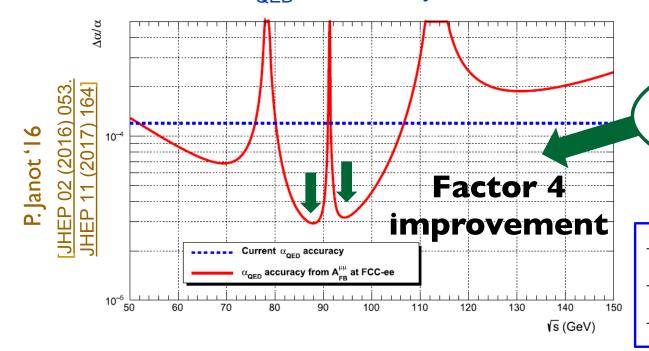
Excellent experimental control of off-peak di-muon asymmetry motivates campaign to collect 50-80 ab-1 off peak to gain highest sensitivity to Z-γ interference

strongly depends on
$$\sqrt{s}$$
 \leftarrow direct measurement of $\alpha_{\rm QED}(s)$ at \sqrt{s} != $m_{\rm Z}$ \leftarrow $A_{\rm FB}^{\mu\mu}(s) \simeq \frac{3}{4} \mathcal{A}_{\rm e} \mathcal{A}_{\mu} \times \left[1 + \frac{8\pi\sqrt{2}\alpha_{\rm QED}(s)}{m_{\rm Z}^2 G_{\rm F} \left(1 - 4\sin^2\theta_{\rm W}^{\rm eff}\right)^2} \frac{s - m_{\rm Z}^2}{2s}\right]$ measure $\sin^2\!\theta_{\rm W}$ to high precision \leftarrow

Allows for clean determination of $\alpha_{OED}(m_Z^2)$, which is a *critical* input for m_w closure tests (see later).



relative α_{QED} uncertainty with 80 ab⁻¹



This dependence, & location of half-integer spin tunes, guides the choice of off-peak energies: 87.8 & 93.9 GeV.

- \rightarrow Measure $\alpha_{QED}(m_Z^{-2})$ to $3x10^{-5}$ rel. precision (currently $1.1x10^{-4}$)
- \rightarrow Stat. dominated; syst. uncertainties < 10⁻⁵ (dominated by \sqrt{s} calib)
- → Theoretical uncertainties ~ 10⁻⁴, higher order calcs needed

CG - 53 / 38

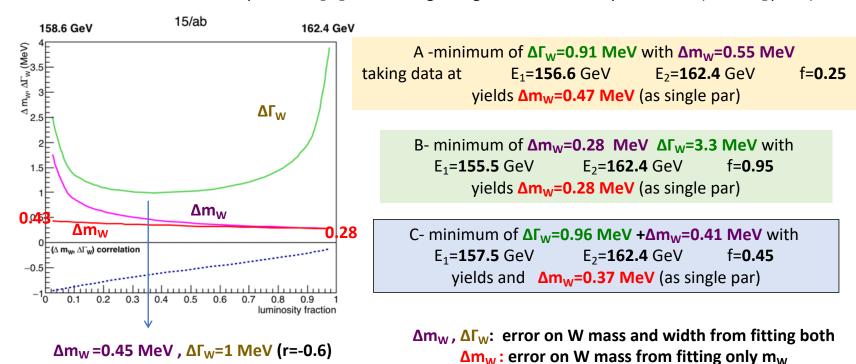
M_{W.}

 $\Delta m_{W} = 0.35 \text{ MeV}$

FCC workshop - 27 Jan 2023

- Two independent W mass and width measurements @FCCee:
- **1. The** m_W and Γ_W determinations from the WW threshold cross section lineshape, with 12/ab at $E_{CM} \simeq 157.5-162.5$ GeV $\Delta m_W = 0.4$ MeV $\Delta \Gamma_W = 1$ MeV
- 2. Other measurements of m_W and Γ_W from the decay products kinematics at $E_{CM} \simeq 162.5-240-365$ GeV Δm_W , $\Delta \Gamma_W = 2-5$ MeV?

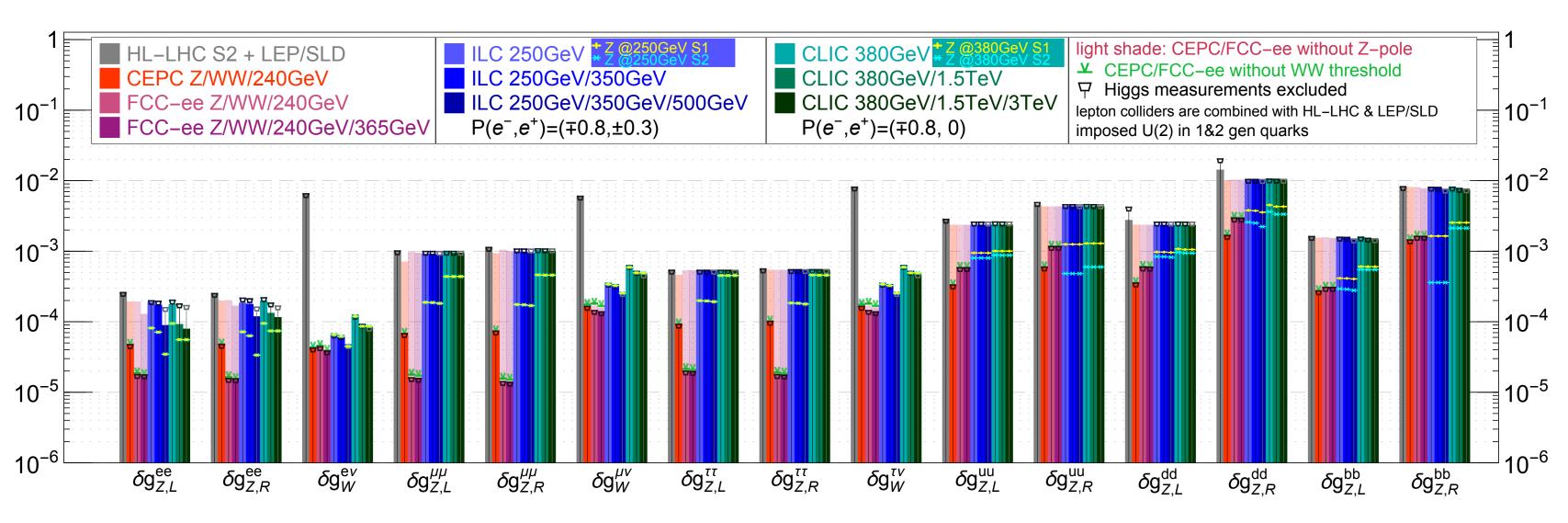
Scans of possible E₁ E₂ data taking energies and luminosity fractions f (at the E₂ point)



P.Azzurri - W mass and width

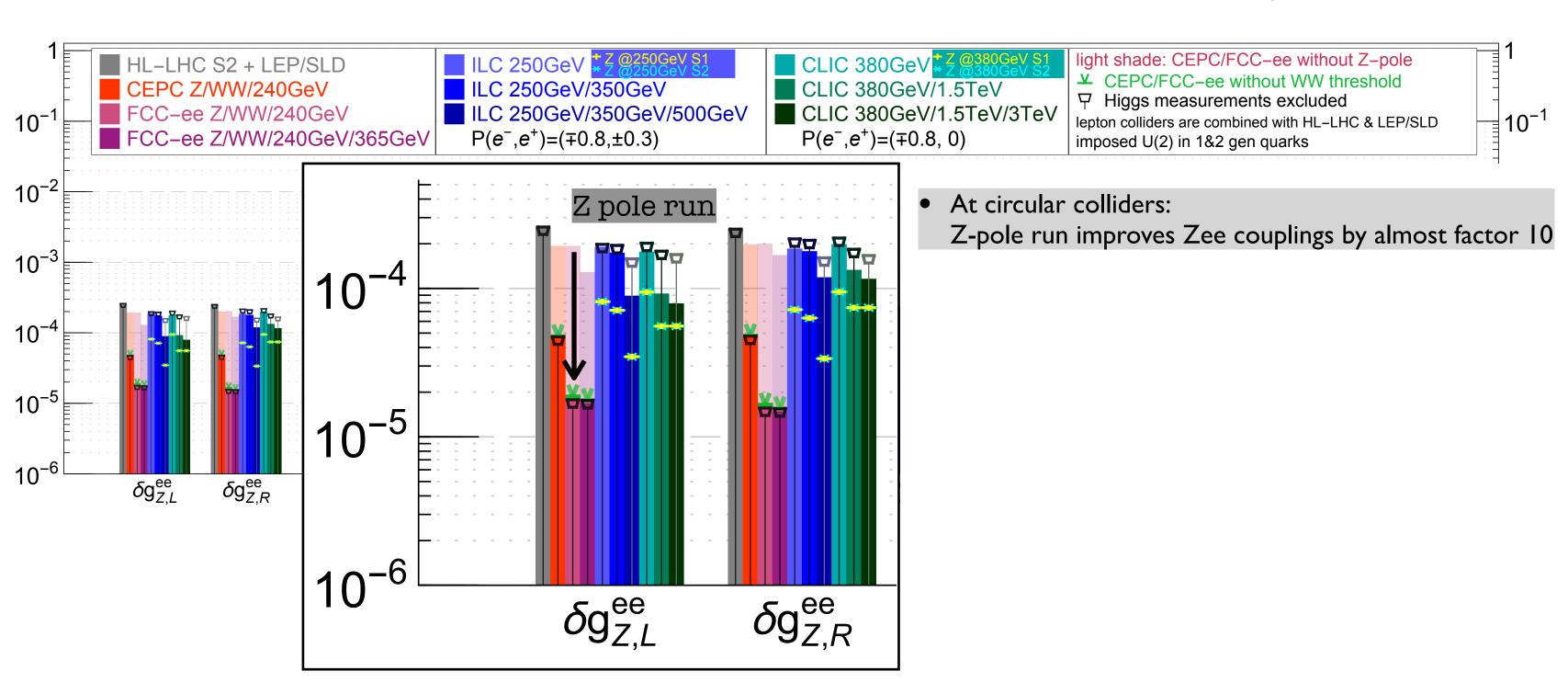
CG - 54 / 38 May 6, 2024

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



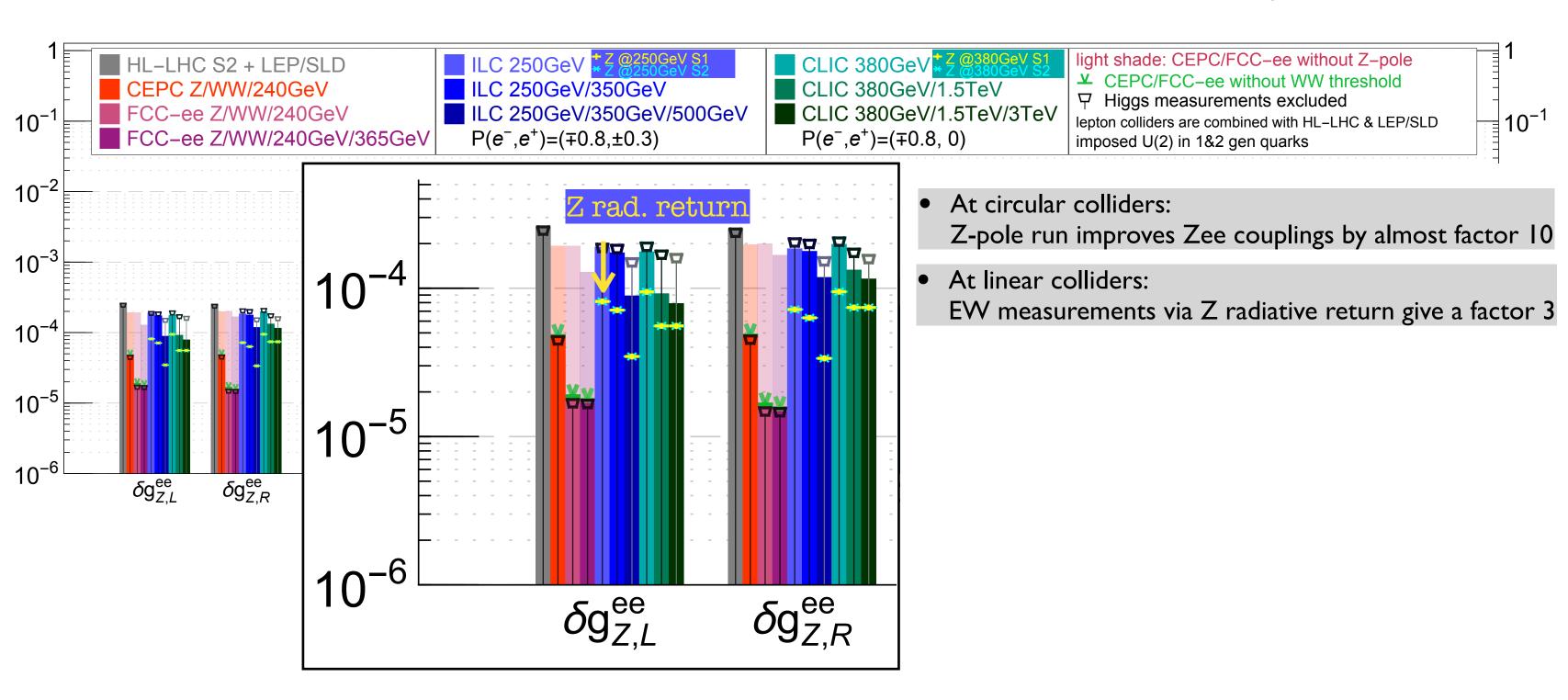
CG - 55 / 38 May 6, 2024

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



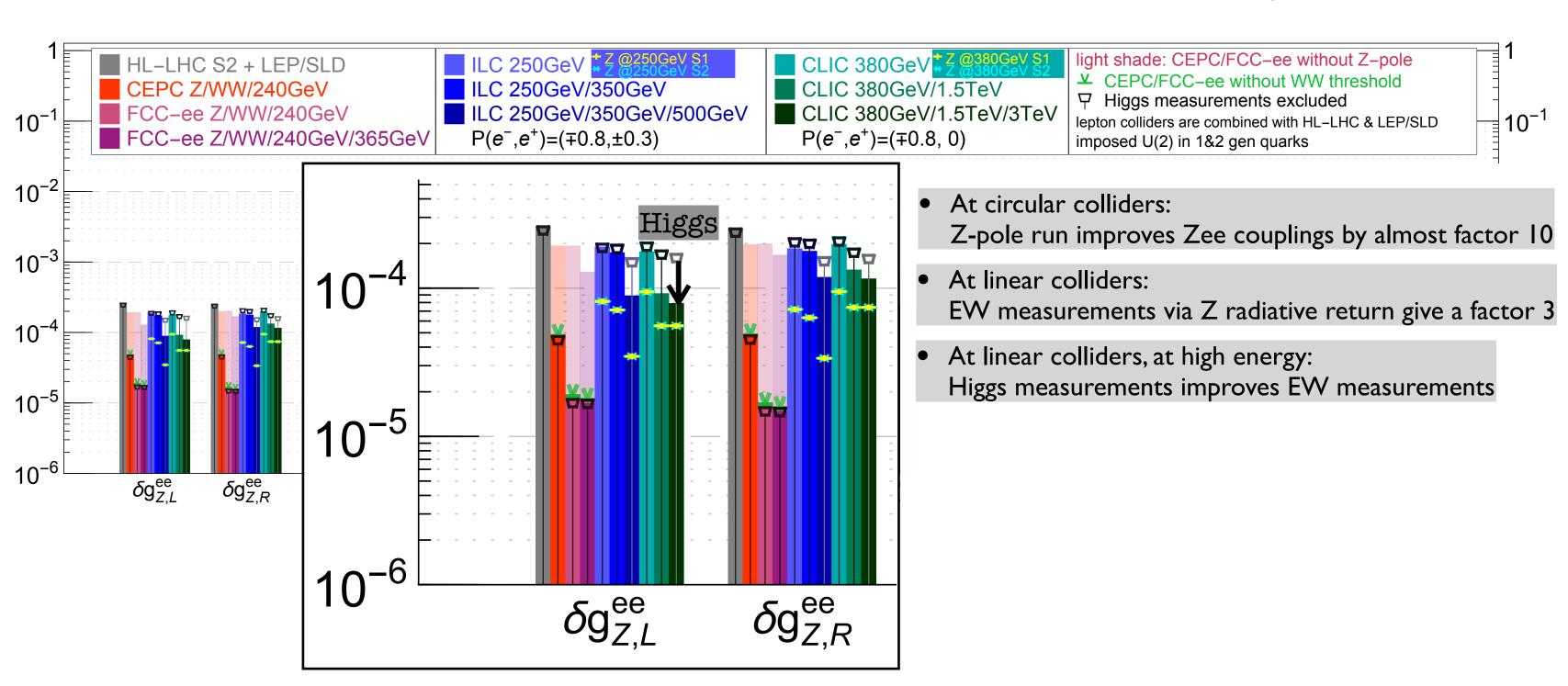
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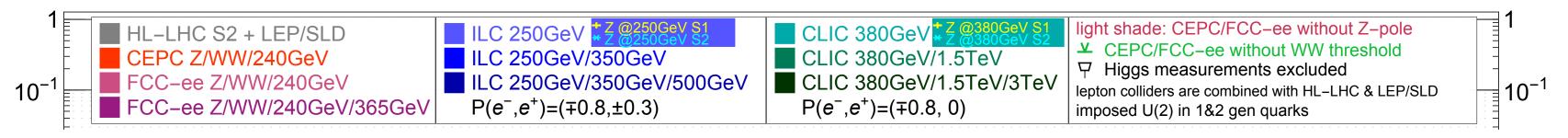
CG - 55 / 38

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



CG - 55 / 38 May 6, 2024

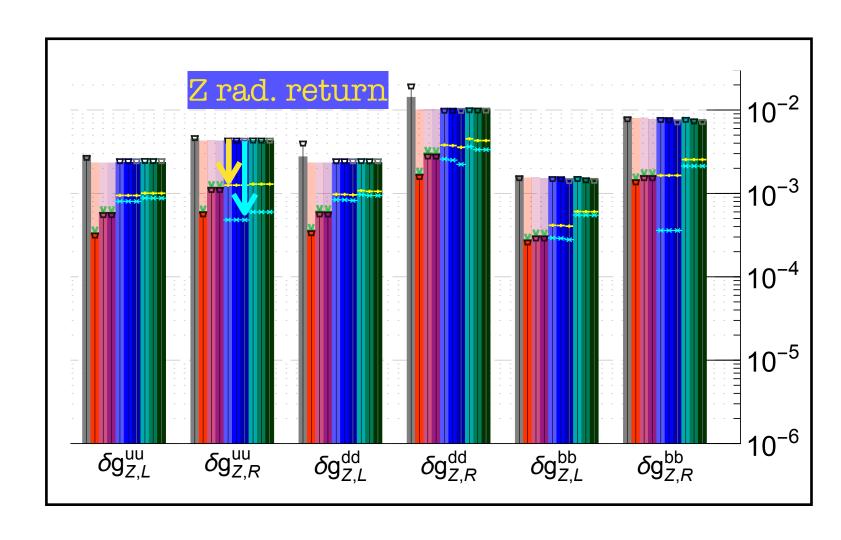
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- At linear colliders, at high energy:
 EW measurements via Z-radiative return has a large impact on Zqq couplings
- Improvements depend a lot on hypothesis on systematic uncertainties

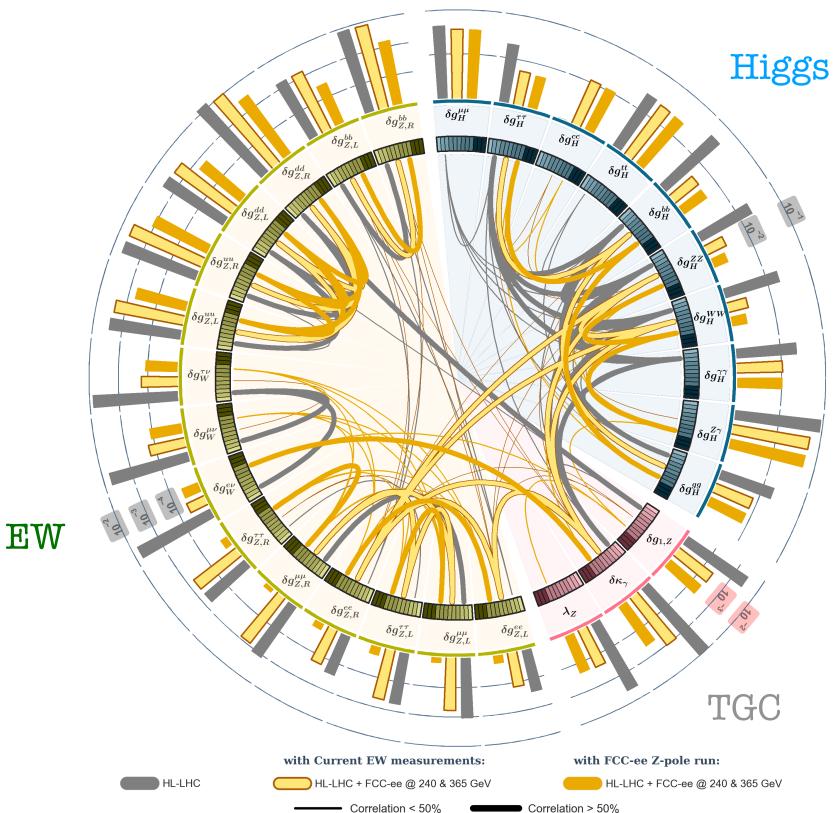
Yellow: LEP/SLD systematics / 2

Blue: small EXP and TH systematics



CG-55/38 May 6, 2024

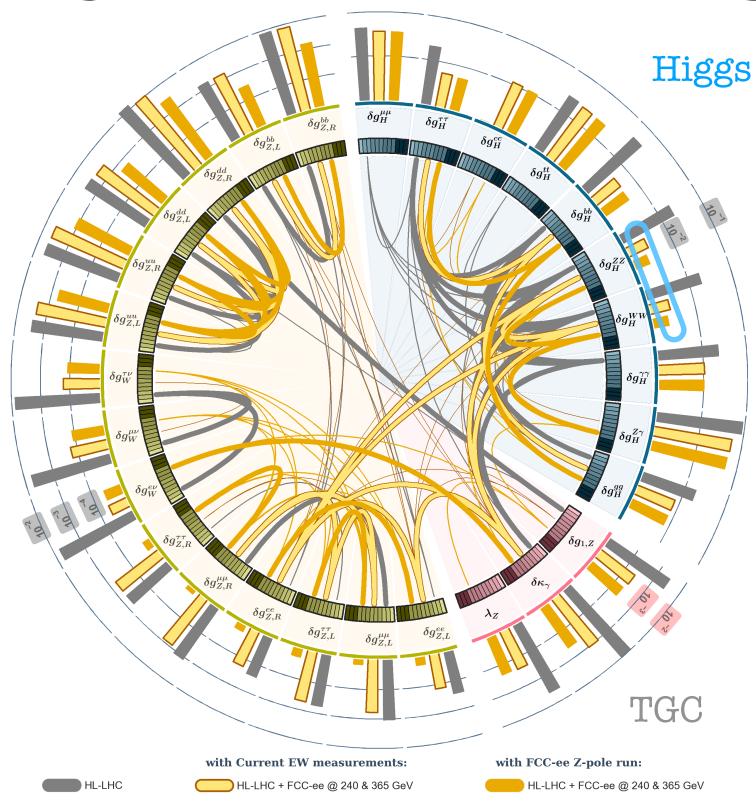
Why Z-pole for Higgs?



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CG - 56 / 38 May 6, 2024

Why Z-pole for Higgs?



Correlation < 50%

Correlation > 50%

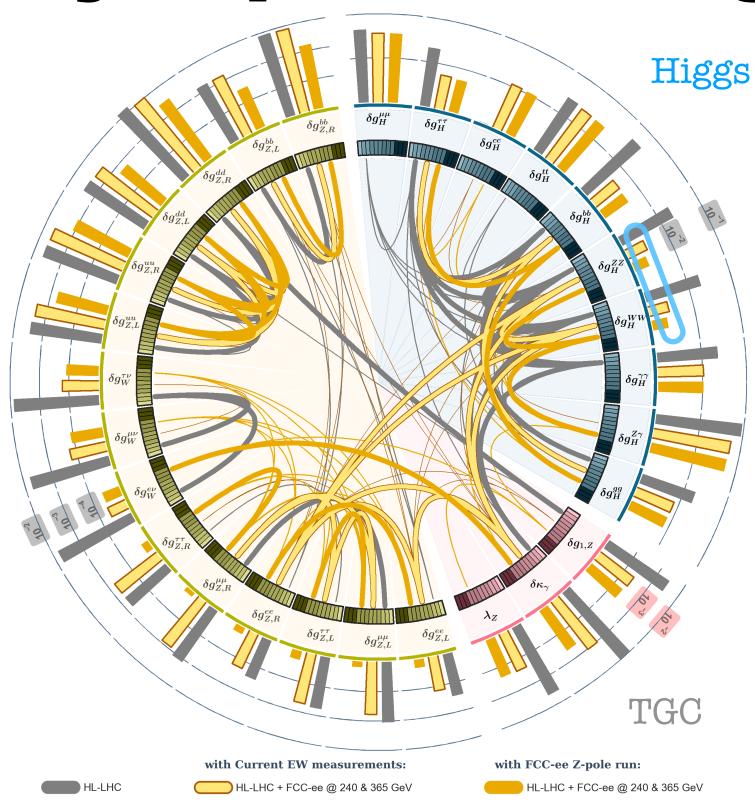
EW

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With Z-pole measurements, Higgs coupling determination improves by up to 50%

CG-56/38

Why Z-pole for Higgs?

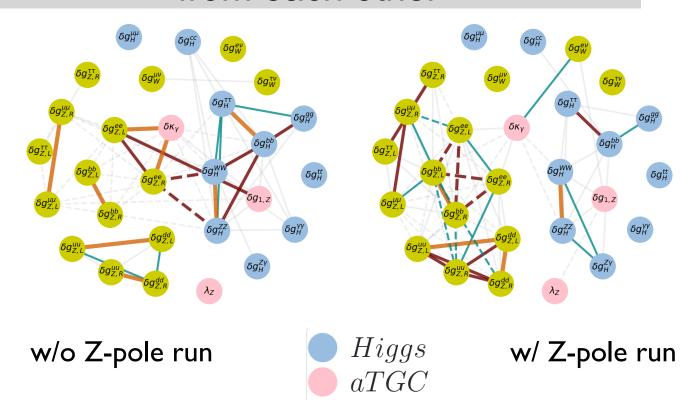


EW

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With Z-pole measurements, Higgs coupling determination improves by up to 50%

Z-pole run at circular colliders decorrelates EW and Higgs sectors from each other

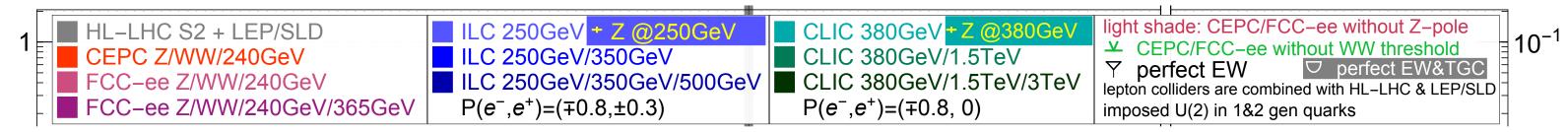


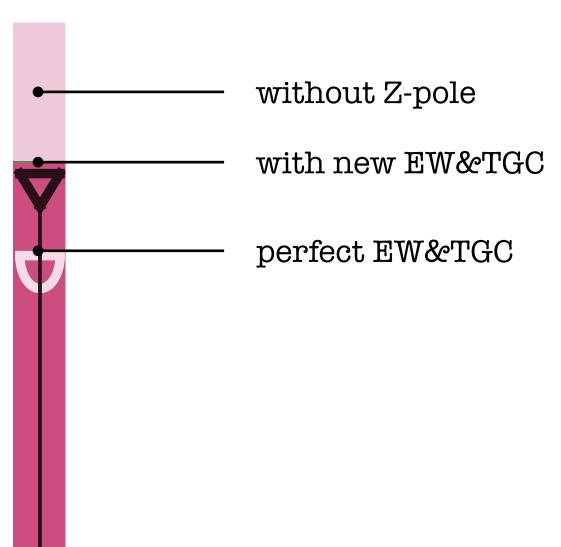
EW

CG - 56 / 38

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Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements



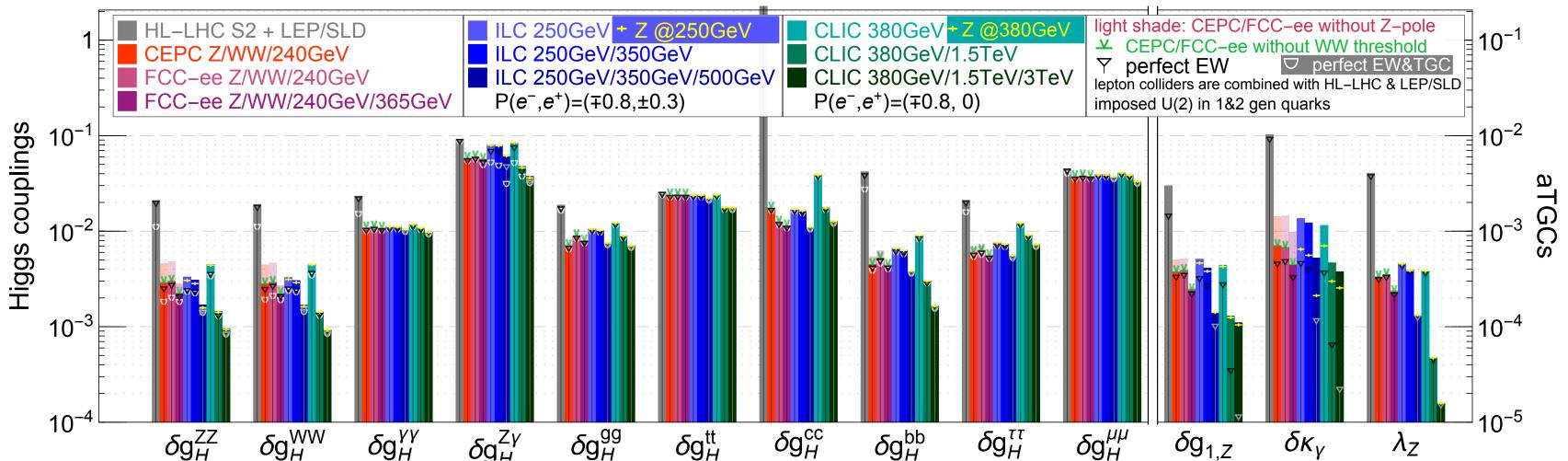


FCC-ee Z/WW/240GeV

May 6, 2024

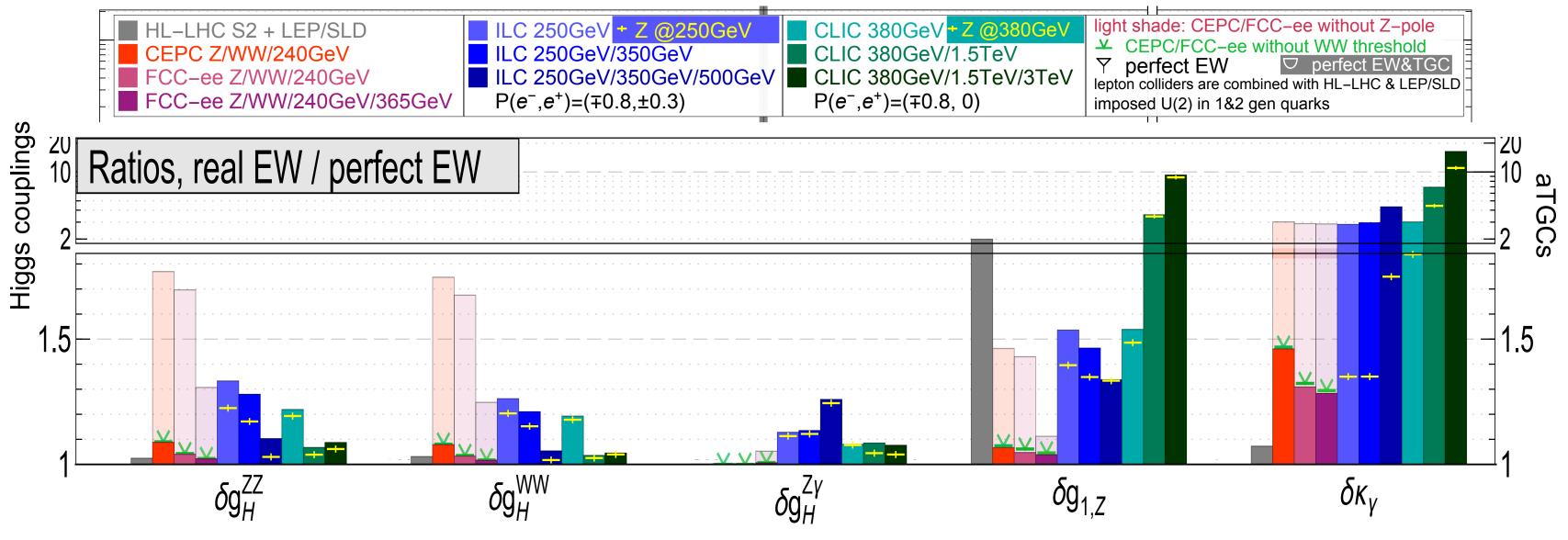
J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements



J. De Blas et al. 1907.04311

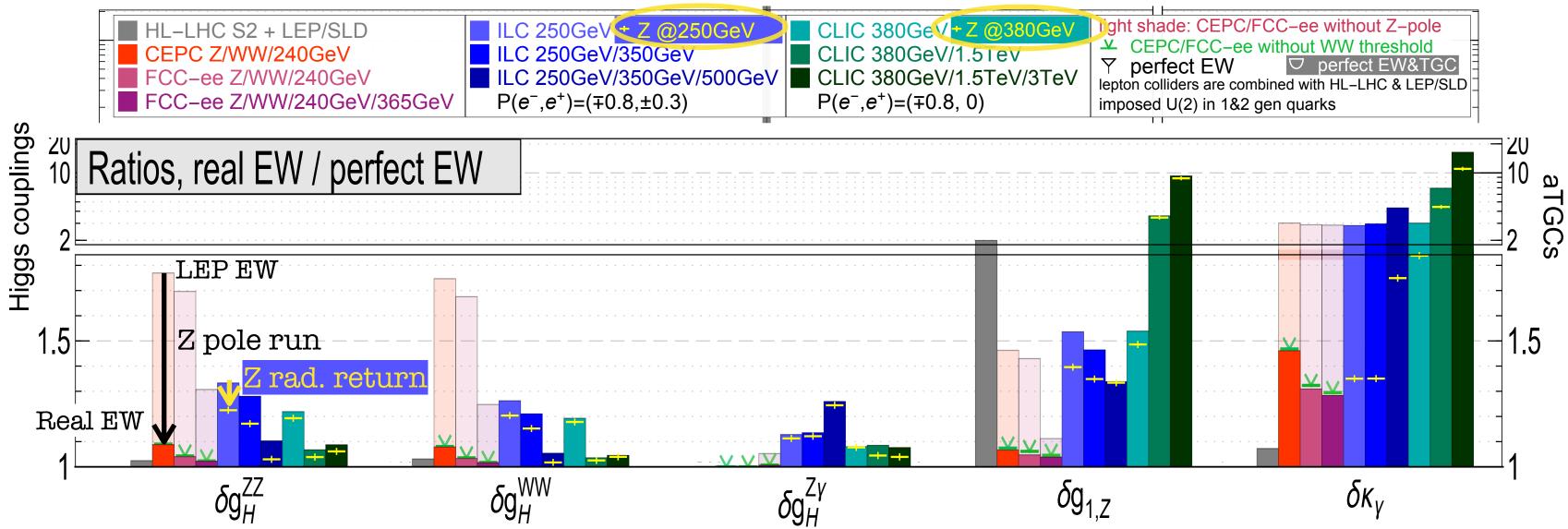
Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements



CG-57/38

J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements

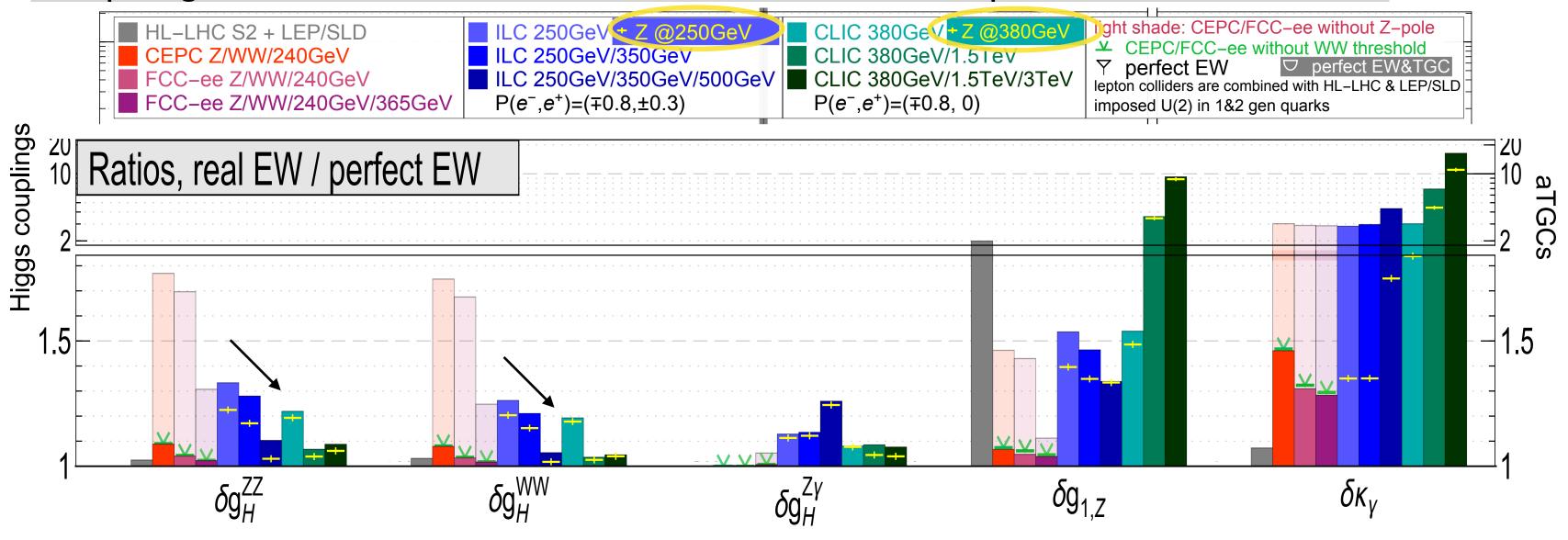


- FCC-ee and CEPC benefit a lot (>50% on HVV) from Z-pole run
- FCC-ee and CEPC EW measurements are almost perfect for what concerns Higgs physics (<10%).
- LEP EW measurements are a limiting factor (\sim 30%) to Higgs precision at ILC, especially for the first runs But EW measurements at high energy (via Z-radiative return) help mitigating this issue

CG-57/38

J. De Blas et al. 1907.04311

Comparing 3 EW scenarios: LEP/SLD, actual EW measurements, perfect EW measurements

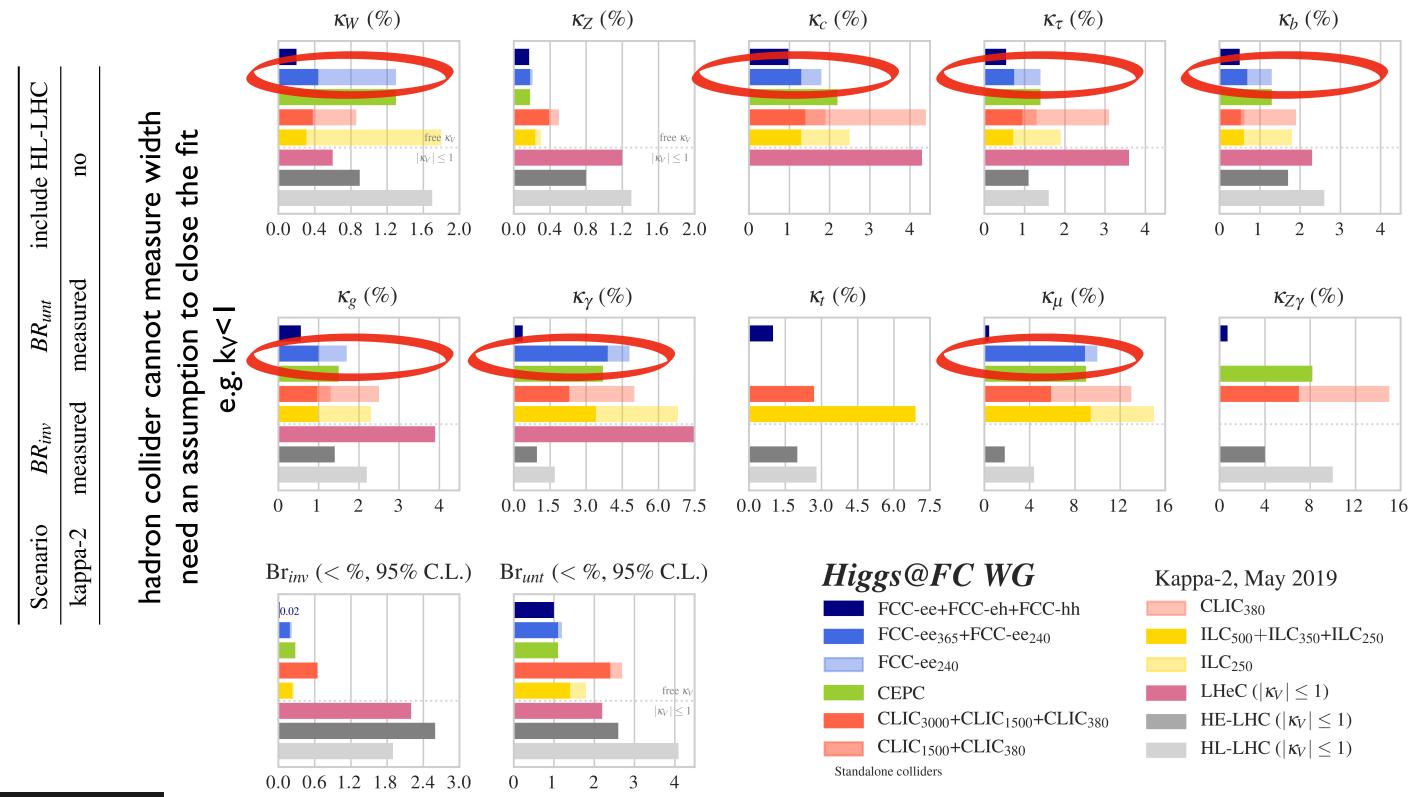


• Higher energy runs reduce the EW contamination in Higgs coupling extraction

CG-57/38

Complementarity 240+365 GeV.

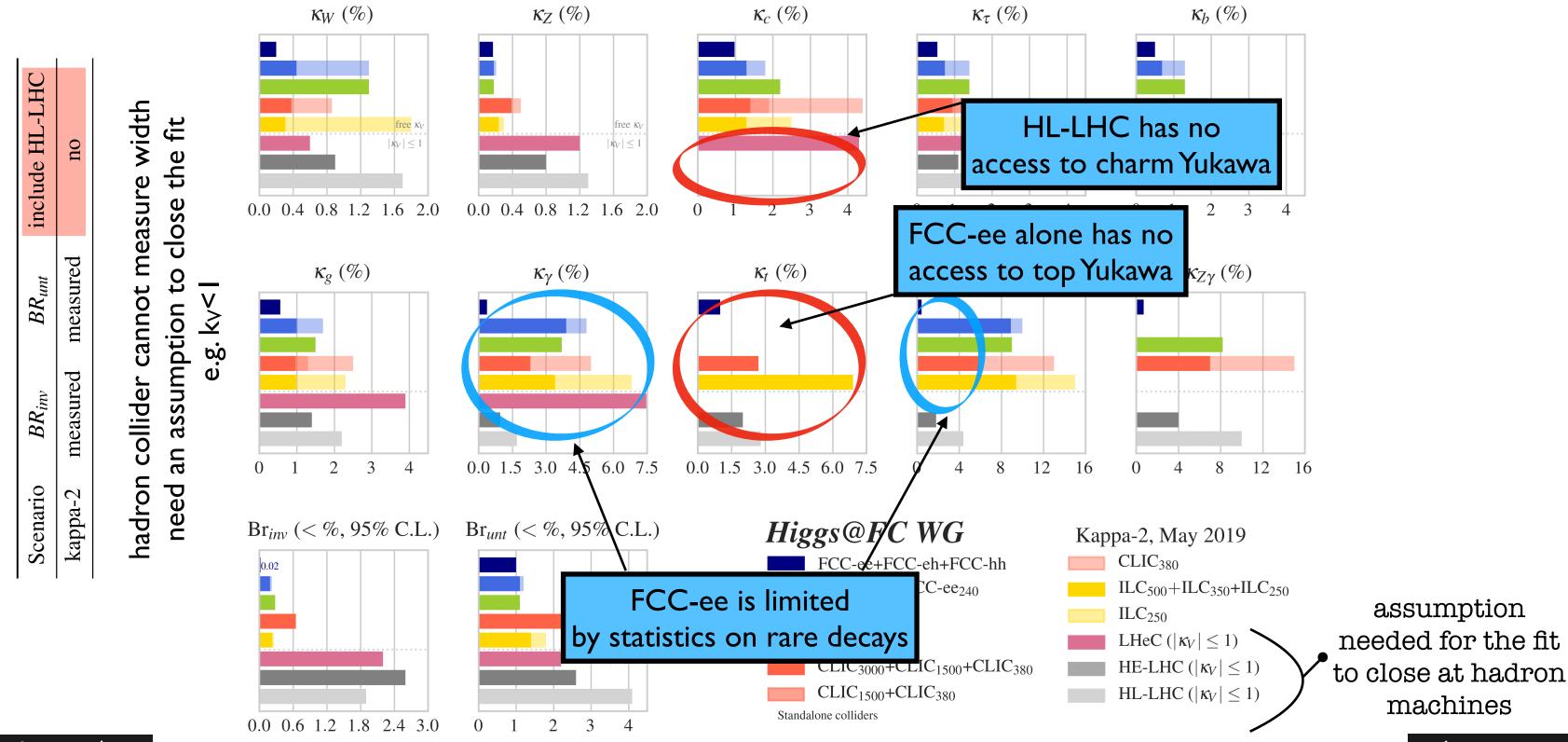
ECFA Higgs study group '19



CG - 58 / 38

Complementarity FCC-ee+HL-LHC.

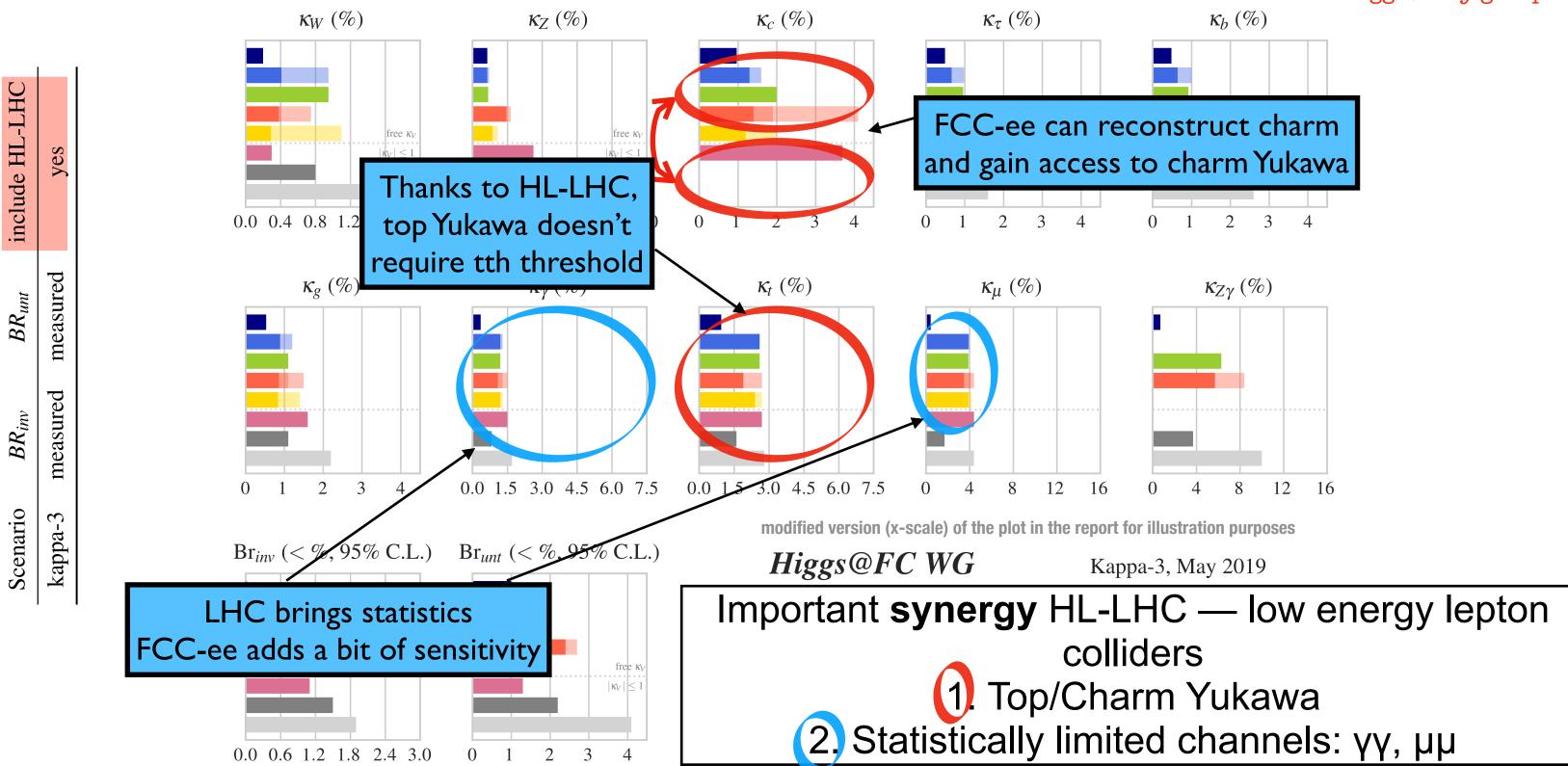
ECFA Higgs study group '19



CG - 59 / 38

Complementarity FCC-ee+HL-LHC.

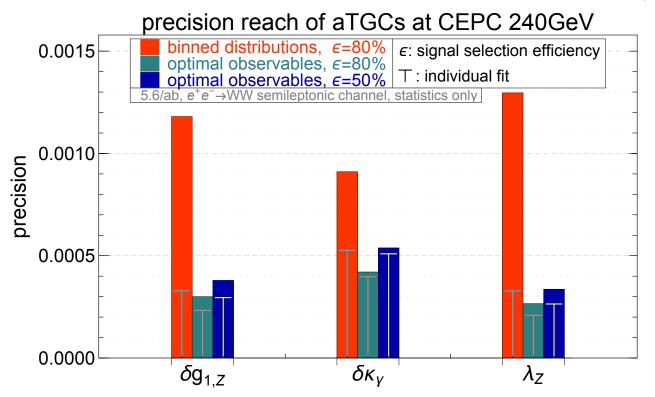
ECFA Higgs study group '19



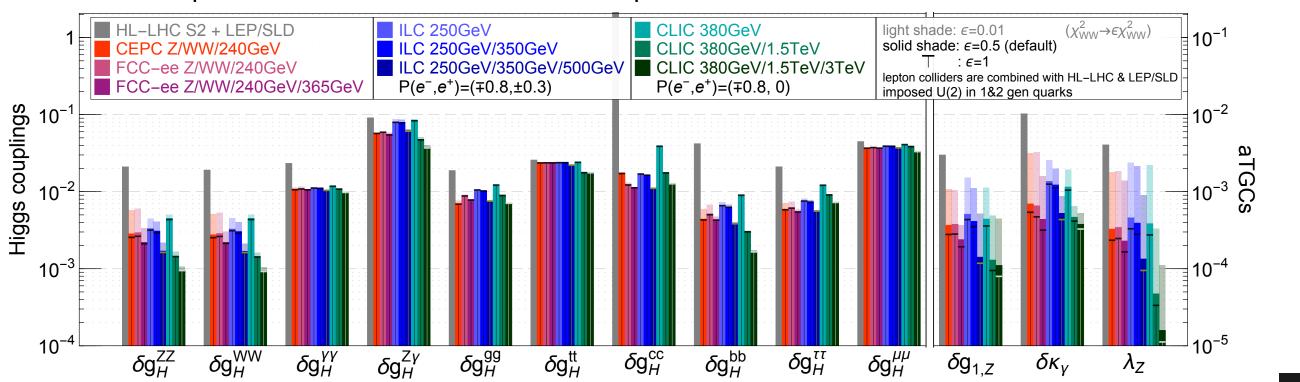
CG-60 / 38

Impact of Diboson Systematics.

J. De Blas, G. Durieux, C. Grojean, J. Gu, A. Paul 1907.04311



precision reach with different assumptions on $e^+e^-\rightarrow WW$ measurements

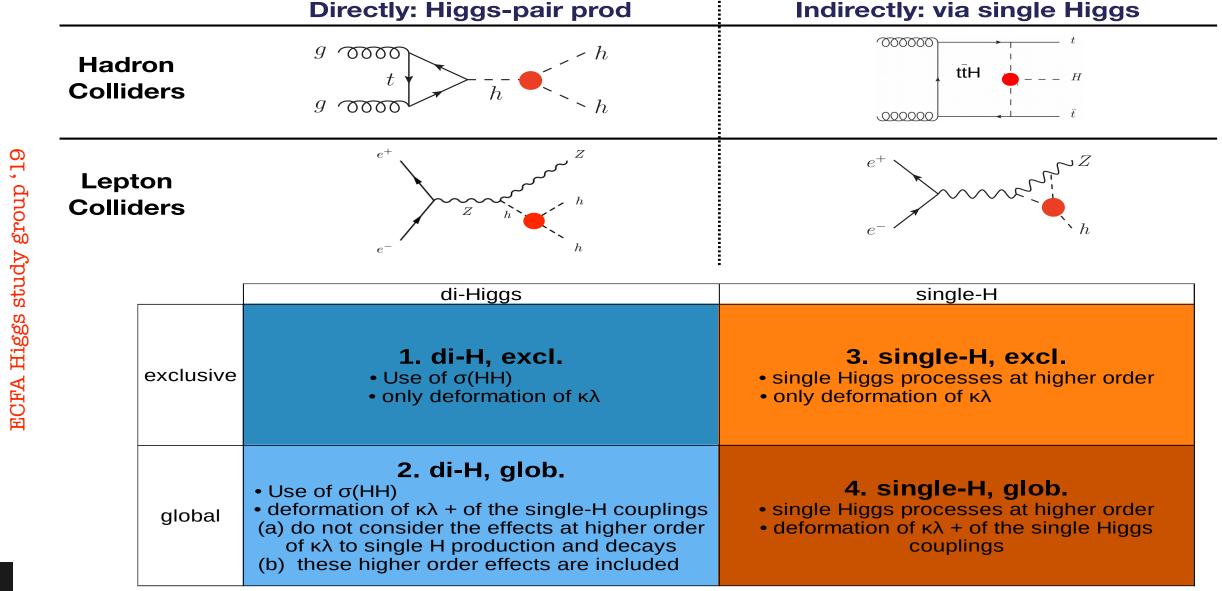


May 6, 2024

Higgs self-coupling.

Higgs self-couplings is very interesting for a multitude of reasons (vacuum stability, hierarchy, baryogenesis, GW, EFT probe...).

How much can it deviate from SM given the tight constraints on other Higgs couplings? Do you need to reach HH production threshold to constrain h³ coupling?



CG - 62 / 38

May 6, 2024

Large self-coupling scenarios.

Generically:

$$\left| \frac{\delta_{h^3}}{\delta_{\text{single }h}} \right| \sim O(1)$$

 $\left| \frac{\delta_{h^3}}{\delta_{\text{single }h}} \right| \sim O(1)$ (composite Higgs/susy)

Particular exceptions: Higgs DM-portal models or custodial EW quadruplet

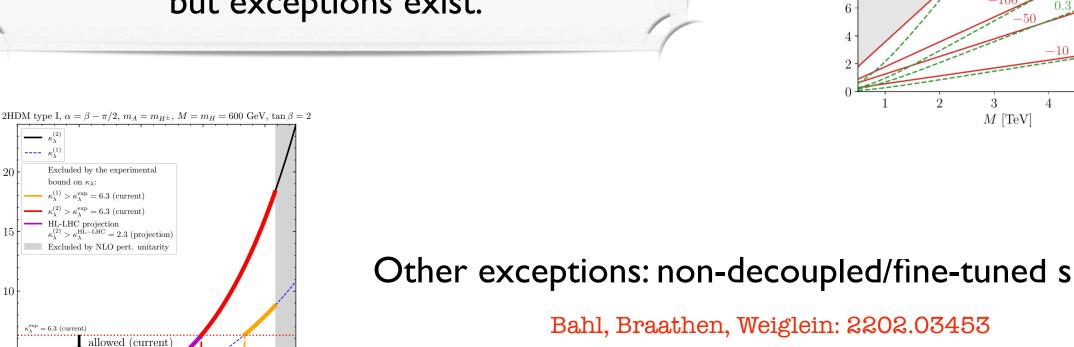
DiVita et al,: 1704.01953

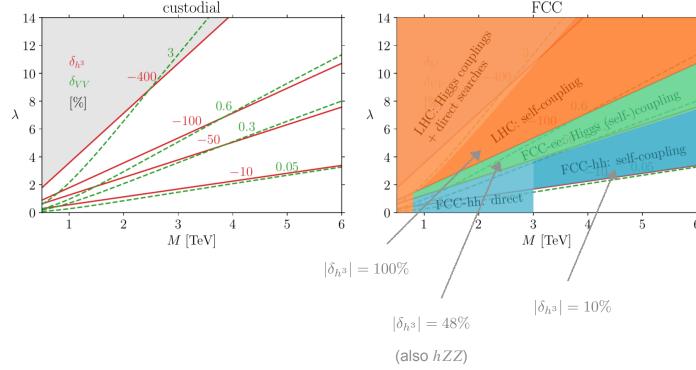
 m_A [GeV]

Falkowski, Rattazzi: 1902.05936

Durieux, McCullough, Salvioni: 2209.00666

h³ generically is not a tool to discover BSM but exceptions exist.

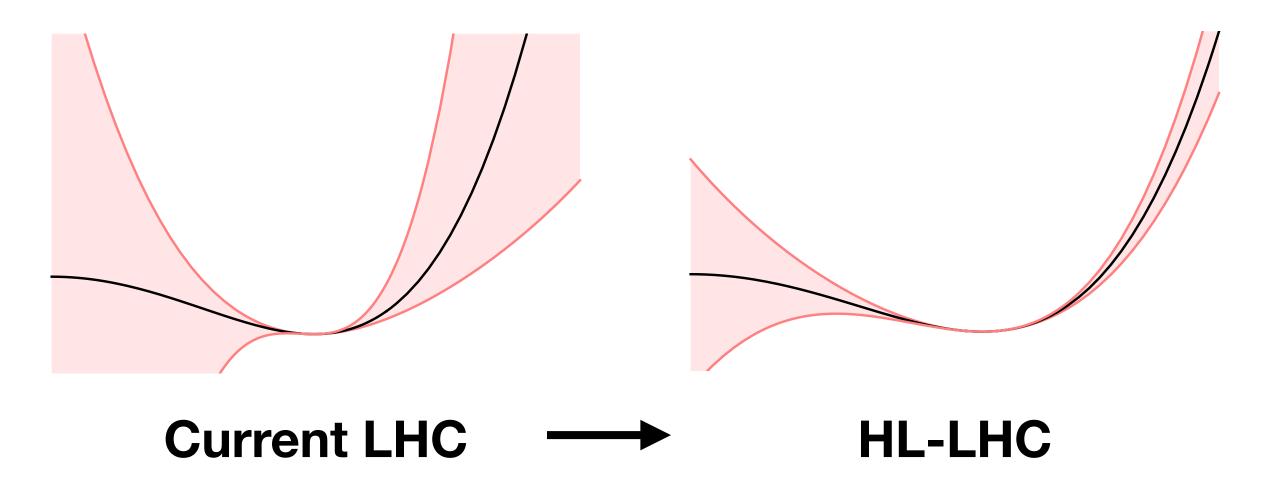




Other exceptions: non-decoupled/fine-tuned spectra

Large self-coupling scenarios.

It is true that we haven't "measured" the Higgs potential but there are only peculiar physics scenarios that produce large deviations in the shape of the potential without leaving imprints elsewhere.



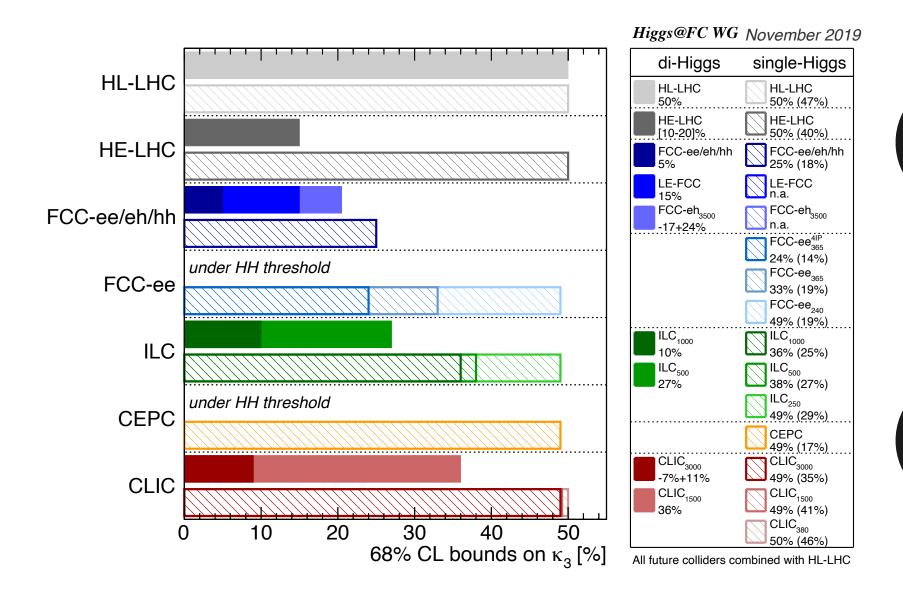
R. Petrossian-Byrne/N. Craig @ LCWS'23

Important to understand which dynamics is really probed when embarking into challenging measurements. Actually, double Higgs production is also interesting to probe new physics in its tail rather than near threshold (where the sensitivity to Higgs self-coupling comes from).

CG - 63 / 38

ECFA Higgs study group '19

Higgs self-coupling.



Don't need to reach HH threshold

to have access to h³.

Z-pole run is very important
if the HH threshold cannot be reached

The determination of h^3 at FCC-hh relies on HH channel, for which FCC-ee is of little direct help. But the extraction of h^3 requires precise knowledge of y_t .

1% $y_t \leftrightarrow 5\%$ h^3 Precision measurement of y_t needs ee

50% sensitivity: establish that h³≠0 at 95%CL

20% sensitivity: 5σ discovery of the SM h³ coupling

5% sensitivity: getting sensitive to quantum corrections to Higgs potential

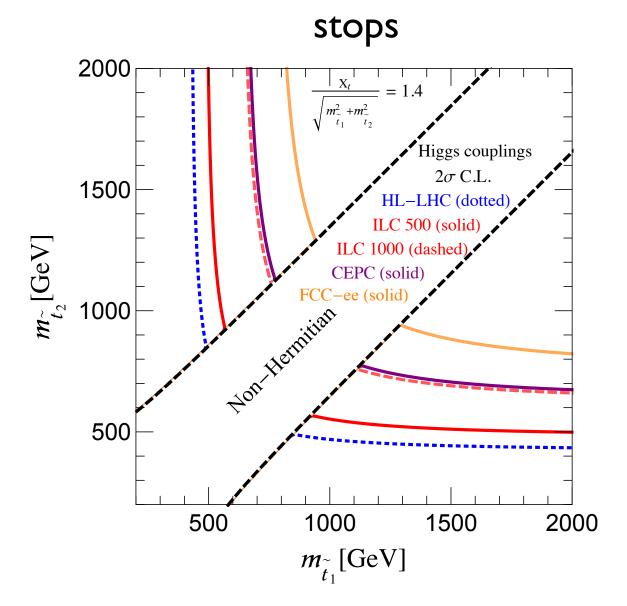
CG-64/38

Discovery potential beyond LHC

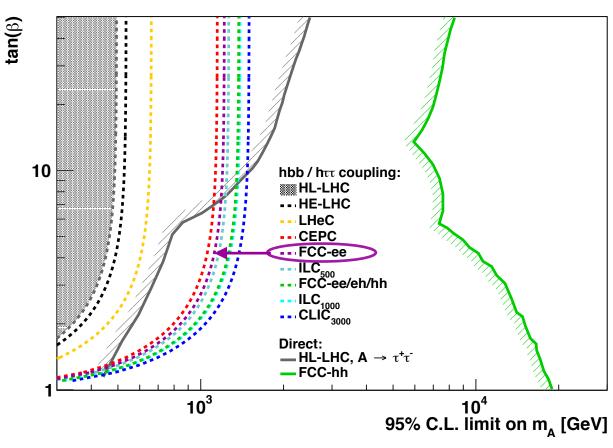
CG - 65 / 38

Discovery Potential Beyond LHC.

Precisely measured EW and Higgs observables are sensitive to heavy New Physics Examples of improved sensitivity wrt direct reach @ HL-LHC: SUSY





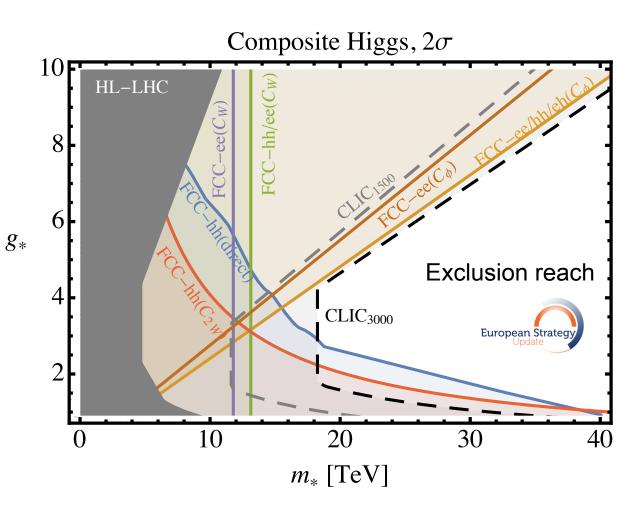


Fan, Reece, Wang '14

ESU Physics BB '19

Discovery Potential Beyond LHC.

Precisely measured EW and Higgs observables are sensitive to heavy New Physics Examples of improved sensitivity wrt direct reach @ HL-LHC: Composite Higgs



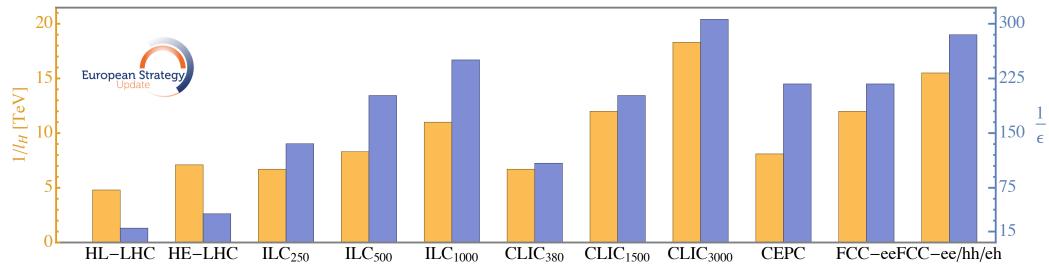


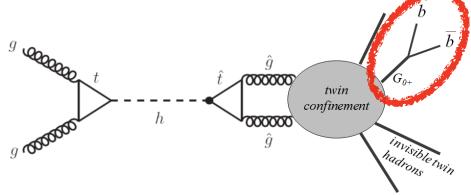
Fig. 8.5: Exclusion reach of different colliders on the inverse Higgs length $1/\ell_H = m_*$ (orange bars, left axis) and the tuning parameter $1/\varepsilon$ (blue bars, right axis), obtained by choosing the weakest bound valid for any value of the coupling constant g_* .

ESU Physics BB '19

Direct Searches for Elusive New Physics

- LLP searches with displaced vertices
 - e.g. in twin Higgs models glueballs that mix with the Higgs and decay back to b-quarks

Craig et al, arXiv:1501.05310



• Rare decays

Gori et al arXiv:2005.05170

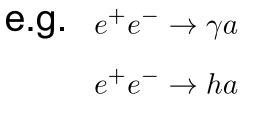
e.g. ALP mixing w/ SM mesons:

$$K_L \to \pi^0 a \to \pi^0 \gamma \gamma \text{ (KOTO)}$$

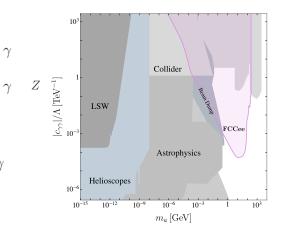
$$K^+ \to \pi^+ a \to \pi^+ \gamma \gamma \text{ (NA62)}$$

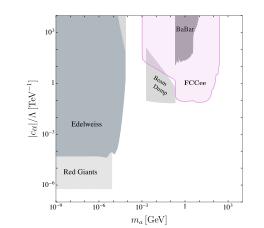
$$\mathcal{L} = \frac{\alpha_s}{8\pi F_a} a G_{\mu\nu} \tilde{G}^{\mu\nu}$$

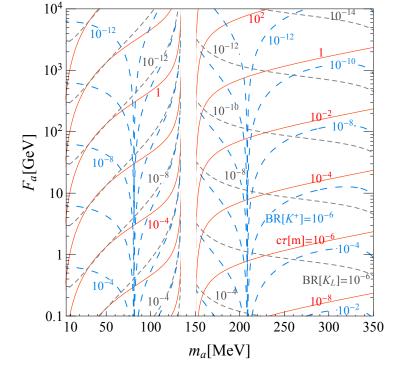




Knapen, Thamm arXiv:2108.08949





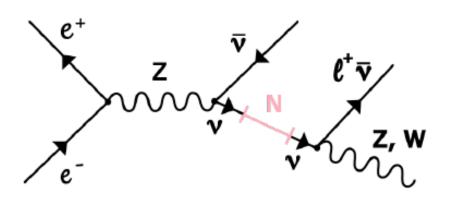


Astro/Cosmo → long-lived ALPs colliders → short-lived ALPs MeV+

May 6, 2024

Search for VRH.

Direct observation in Z decays from LH-RH mixing



Important to understand

- 1. how neutrinos acquired mass
- 2. if lepton number is conserved
 - 3. if leptogenesis is realised

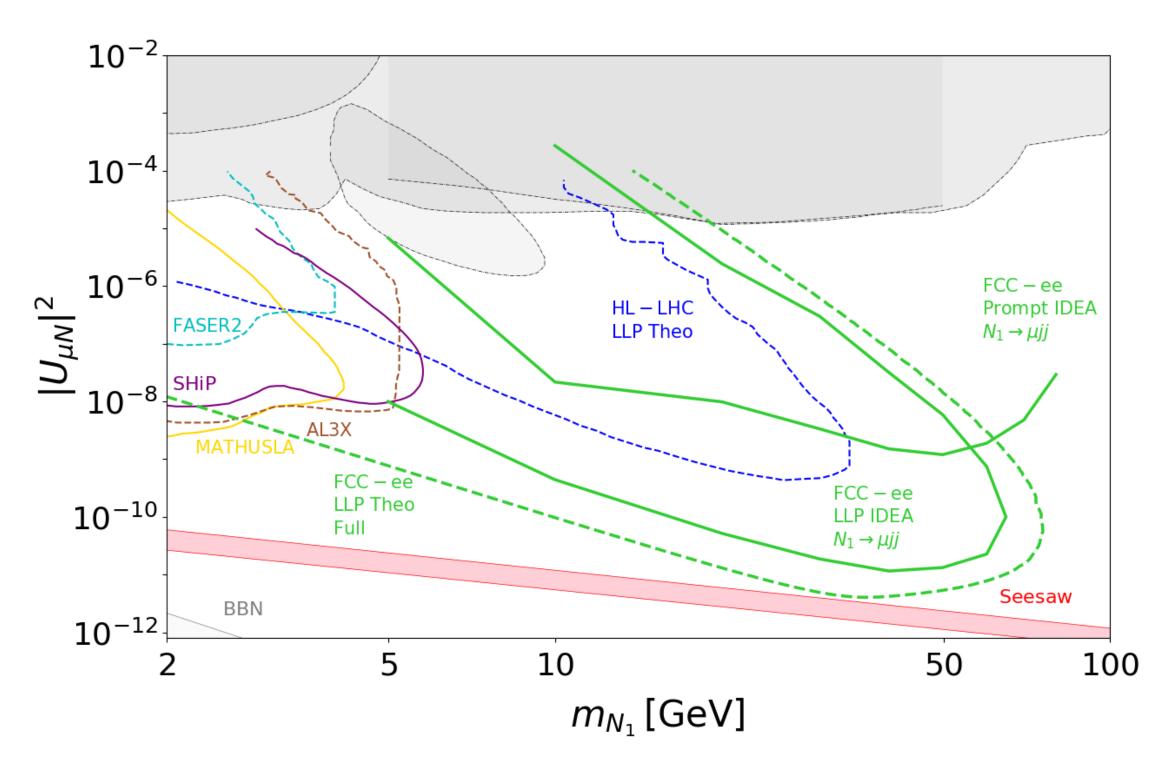
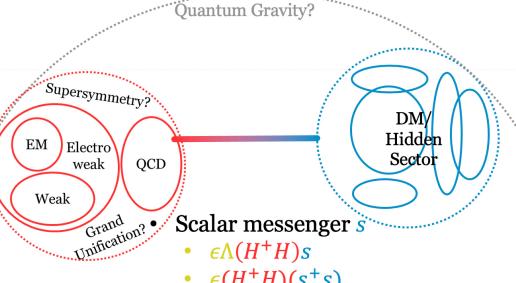


Fig. from mid-term report

CG - 69 / 38

Exotics/Long Lived Particles.

Z. Liu @ CEPC 2020



The Higgs could be a good portal to Dark Sector

— rich exotic signatures —

Sc	alar messenger s
•	$\epsilon \Lambda (H^+H)s$
•	$\epsilon(H^+H)(s^+s)$
• Ve	ctor messenger A'_{μ}
•	$\epsilon F^{\mu\nu}F'_{\mu\nu}$
•	$\epsilon J_{SM}^{\mu} A_{\mu}^{\prime}$
Ne	eutrino messenger N
•	$\epsilon(LH)N$
· Ax	ion messenger <i>a</i>
•	$\frac{a}{a} \left(\frac{\alpha_3}{a} G \tilde{G} + \frac{\alpha_2}{a} W \tilde{W} \right)$

Decay Topologies	Decay mode \mathcal{F}_i	Decay Topologies	Decay mode \mathcal{F}_i
h o 2	$h \to E_{\mathrm{T}}$	h o 2 o 4	$h o (b \bar b) (b \bar b)$
h o 2 o 3	$h o \gamma + E_{ m T}$		$h \rightarrow (b\bar{b})(\tau^+\tau^-)$
	$h ightarrow (bar{b}) + ot\!\!\!/ _{ m T}$		$h o (b\bar{b})(\mu^+\mu^-)$
	$h ightarrow (jj) + ot\!\!\!/_{ m T}$		$h \rightarrow (\tau^+ \tau^-)(\tau^+ \tau^-)$
	$h ightarrow (au^+ au^-) + E_{ m T}$	-	$h \rightarrow (\tau^+ \tau^-)(\mu^+ \mu^-)$
	$h o (\gamma \gamma) + ot\!\!\!E_{ m T}$		$h \to (jj)(jj)$
	$h \rightarrow (\ell^+\ell^-) + \cancel{E}_{\mathrm{T}}$		$h \rightarrow (jj)(\gamma\gamma)$
h o 2 o 3 o 4	$h ightarrow (bar{b}) + ot\!\!\!\!\!E_{ m T}$		$h \rightarrow (jj)(\mu^+\mu^-)$
	$h o (jj) + \cancel{E}_{\mathrm{T}}$		$h \rightarrow (\ell^+\ell^-)(\ell^+\ell^-)$
	$h \rightarrow (\tau^+ \tau^-) + \cancel{E}_{\mathrm{T}}$		$h o (\ell^+\ell^-)(\mu^+\mu^-)$
	$h \to (\gamma \gamma) + \cancel{E}_{\mathrm{T}}$	_	$h \to (\mu^+ \mu^-)(\mu^+ \mu^-)$
	$h \rightarrow (\ell^+\ell^-) + \cancel{E}_{\mathrm{T}}$		$h o (\gamma \gamma) (\gamma \gamma)$
h > 2 > (1 + 2)	$h \rightarrow (\mu^+\mu^-) + \cancel{E}_T$		$h o \gamma \gamma + E_{ m T}$
$h \rightarrow 2 \rightarrow (1+3)$	$egin{aligned} h ightarrow bar{b} + ot \!$	$h \rightarrow 2 \rightarrow 4 \rightarrow 6$	$h ightarrow (\ell^+\ell^-)(\ell^+\ell^-) + E_{ m T}$
	$h o ff + \not\!\!\!E_{ m T}$ $h o au^+ au^- + \not\!\!\!E_{ m T}$		$h \rightarrow (\ell^+\ell^-) + \cancel{E}_{\mathrm{T}} + X$
	$h \rightarrow \gamma \gamma + \cancel{E}_{\mathrm{T}}$	h o 2 o 6	$h ightarrow \ell^+\ell^-\ell^+\ell^- + E_{ m T}$
	$h o \ell^+\ell^- + \cancel{E}_{\mathrm{T}}$	\leftarrow	$h \rightarrow \ell^+\ell^- + E_T + X$

LHC's strength
Hard at LHC due to
missing energy
Hard at LHC due to
hadronic
background

Lepton colliders' strength

CG-70 / 38 May 6, 2024

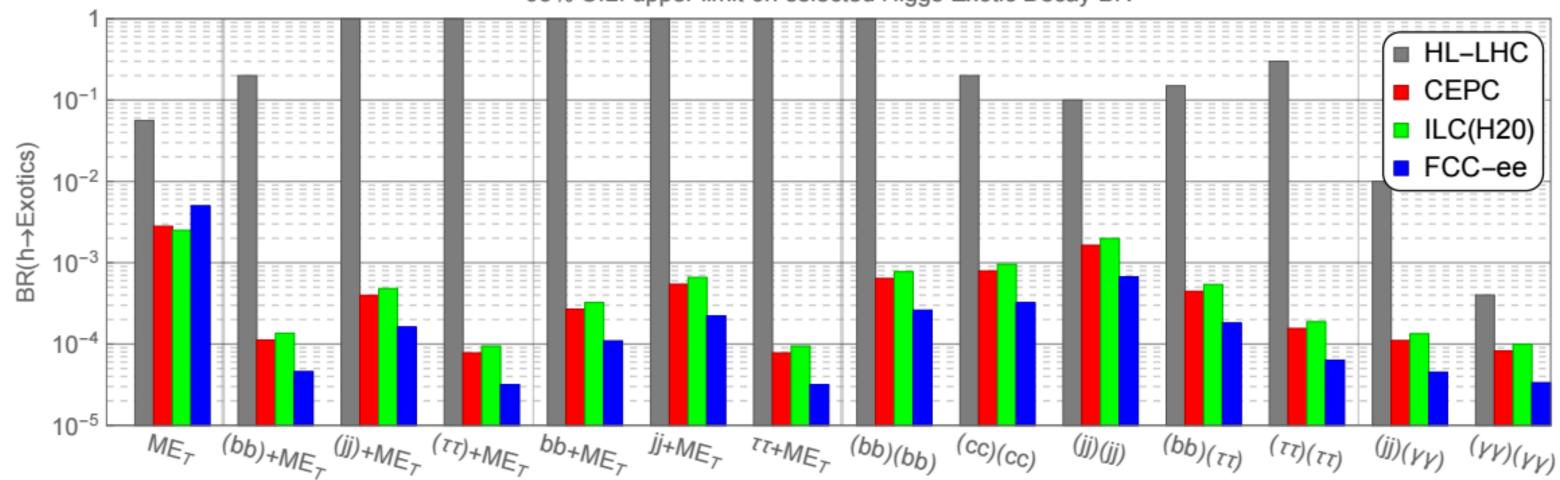
Exotics/Long Lived Particles.

Z. Liu @ CEPC 2020

The Higgs could be a good portal to Dark Sector

— rich exotic signatures —





How to improve?

> Dedicated detectors, see e.g. talk by R. Gonzalez Suarez @ FCC week 2021

CG -70 / 38

Cost of Operation

CG - 71 / 38

Energy and carbon footprint.

FCC-ee total instantaneous power demand at each centre-of-mass energies

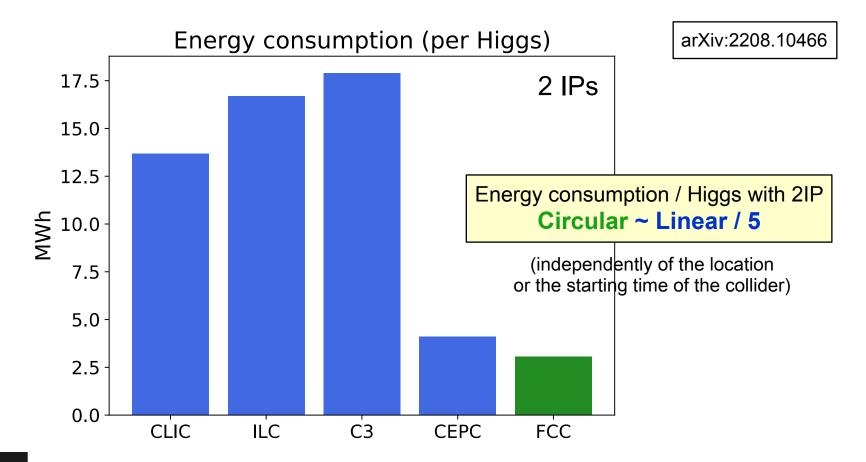
				JP. Burnet, Fo	CC VVEEK ZZ	
		Z	W	Н	TT	
Beam energy (GeV)		45.6	80	120	182.5	
Magnet current		25%	44%	66%	100%	
Power ratio		6%	19%	43%	100%	
PRF EL (MW)	Storage	146	146	146	146	Ongoing R&D
PRFb EL (MW)	Booster	2	2	2	2	
Pcryo (MW)	all	1,3	12,6	15,8	47,5	
Pcv (MW)	all	33	34	36	40.2	Ongoing R&D
PEL magnets (MW)	Stroage	6	17	39	89	
PEL magnets (MW)	Booster	1	3	5	11	
Experiments (MW)	Pt A & G	8	8	8	8	Potential energy sayings
Data centers (MW)	Pt A & G	4	4	4	4	Potential energy savings
General services (MW)		36	36	36	36	
Power during beam operation (MW)		237	262	291	384	

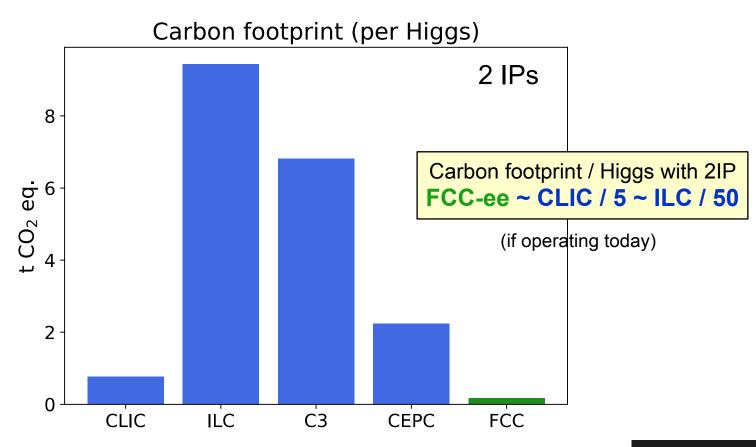
- ◆ At 240 GeV, the instantaneous power of FCC-ee amounts to 291 MW
 - As a comparison, $P(ILC_{250})=140$ MW, $P(CLIC_{380})=110$ MW: less power hungry than FCC-ee?
 - ► Not clear: both produce (2 to 4 times) less Higgs than FCC-ee₂₄₀, with (3 to 6 times) longer running time

CG - 72 / 38

Energy and carbon footprint.

- Our first responsibility (as particle physicists) is to do the maximum of science
 - With the minimal energy consumption and the minimal environmental impact for our planet
 - Should become one of our top-level decision criteria for design, choice and optimization of a collider
- All Higgs factories have a "similar" physics outcome (ESU'20 and Snowmass'21)
 - Natural question: what is their energy consumption or carbon footprint for the same physics outcome?
 - Circular colliders have a much larger instantaneous luminosity and operate several detectors
 - FCC-ee is at CERN, where electricity is already almost carbon-free (and will be even more so in 2048)





March 8, 2024

Cost of Operation.

The total electrical energy consumption over the fourteen years of the FCC-ee research programme is estimated to be around 27 TWh [58], corresponding to an average electricity consumption of 1.9 TWh/year over the entire operation programme, to be compared with the 1.2 TWh/year consumed by CERN today and the expected 1.4 TWh/year for HL-LHC⁹. At the CERN electricity prices from 2014/15, the electricity cost for FCC-ee collider operation would be about 85 MEuro per year. In the HZ running mode, about one million Higgs bosons are expected to be produced in three years, which sets the price of each FCC-ee Higgs boson at 255 Euros. A similar exercise can be done for the first stage of CLIC, expected to consume 0.8 TWh/year over 8 years at 380 GeV to produce about 150,000 Higgs bosons, which sets the price of a CLIC Higgs boson at about 2000 Euros. Finally, with the official ILC operation cost in Japan of 330 MEuro per year [10], its 11.5 to 18.5 years of operation (Section 5), and the 500,000 Higgs bosons produced in total, the price of an ILC Higgs boson is between 7,000 and 12,000 Euros, i.e., between 30 and 50 times more expensive than at FCC-ee. These operation costs are summarized in Table 8.

Table 8: Operation costs of low-energy Higgs factories, expressed in Euros per Higgs boson.

Collider	ILC_{250}	$CLIC_{380}$	$FCC-ee_{240}$
Cost (Euros/Higgs)	7,000 to 12,000	2,000	255