Cosmic Inflation at the Crossroads

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Annecy, 6/11/2024

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Outline

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J. Martin, CR and V. Vennin: arXiv:2404.10647, arXiv:2404.15089, arXiv:1303.3787v4

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- A predictive, testable and tested early universe paradigm
 - Accelerated expansion of the universe at $E_{inf} > MeV$ (BBN)
 - Addresses some unexplainable features of the Friedmann-Lemaître model
- For the simplest incarnation of inflation...
 - Historically introduced to dilute monopoles formed at GUT
 - Flatness of the spatial sections ($\Omega_{\rm K} = 0.0009 \pm 0.0018$)
 - Statistical isotropy of the observable universe (horizon problem)
 - Origin of CMB and LSS (quantum fluctuations)
 - Gaussianities of the cosmological perturbations $(f_{\rm NL} < -0.9 \pm 5)$
 - Adiabaticity of the cosmological perturbations (isocurv. < 1%)
 - Almost scale invariance $(n_{\rm s} = 0.9649 \pm 0.004)$

Triggering and stopping acceleration

The simplest way: single-field inflation

$$S = \int dx^4 \sqrt{-g} \left[\frac{1}{2\kappa^2} R + \mathcal{L}(\phi) \right] \quad \text{with} \quad \mathcal{L}(\phi) = -\frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi)$$

Inflation occurs in the plateau and is followed by a reheating era



• The reheating stage: everything after ϕ_{end} till radiation domination

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Cosmological perturbations of quantum origin



Model testing: reheating effects must be included!

Redshift at which inflation ends

Depends on the redshift of reheating

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$$1 + z_{\text{end}} = \frac{a_0}{a_{\text{end}}} = \frac{a_{\text{reh}}}{a_{\text{end}}} (1 + z_{\text{reh}}) = \frac{a_{\text{reh}}}{a_{\text{end}}} \left(\frac{\rho_{\text{reh}}}{\tilde{\rho}_{\gamma}}\right)^{1/4} = \frac{1}{R_{\text{rad}}} \left(\frac{\rho_{\text{end}}}{\tilde{\rho}_{\gamma}}\right)^{1/4}$$

• The reheating parameter
$$R_{\rm rad} \equiv \frac{a_{\rm end}}{a_{\rm reh}} \left(\frac{\rho_{\rm end}}{\rho_{\rm reh}}\right)^{1/4}$$

• Encodes any observable deviations from a radiation-like or instantaneous reheating $R_{rad} = 1$

• $R_{
m rad}$ can be expressed in terms of $(
ho_{
m reh}, \overline{w}_{
m reh})$ or $(\Delta N_{
m reh}, \overline{w}_{
m reh})$

$$\ln R_{\rm rad} = \frac{\Delta N_{\rm reh}}{4} (3\overline{w}_{\rm reh} - 1) = \frac{1 - 3\overline{w}_{\rm reh}}{12(1 + \overline{w}_{\rm reh})} \ln \left(\frac{\rho_{\rm reh}}{\rho_{\rm end}}\right)$$

• Scary Astrophysics (early universe)

$$10^{10} < z_{\rm end} < 10^{28}, \qquad -46 < \ln R_{\rm rad} < 15$$

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New version (published in PDU 10/2024) (arXiv:1303.3787v4)

- Deals with accurate slow-roll predictions for 287 models
- Comes with a public runtime library ASPIC



• Computes the Hubble-flow functions from the model parameters $heta_{ ext{inf}}$

 $(\boldsymbol{\theta}_{\mathrm{inf}}, R_{\mathrm{rad}}) \longrightarrow \mathsf{ASPIC} \longrightarrow \boldsymbol{\epsilon_i}(\boldsymbol{\theta}_{\mathrm{inf}}, R_{\mathrm{rad}}) \equiv \frac{\mathrm{d}\ln|\boldsymbol{\epsilon_{i-1}}|}{\mathrm{d}\ln a}, \quad \boldsymbol{\epsilon_0} \propto \frac{1}{H}$

Primordial power spectra

Uniquely determined from the Hubble-flow functions

$$\begin{aligned} \mathcal{P}_{\zeta} &= \frac{0}{8\pi^2 M_{\rm Pl}^2 \epsilon_{1*}} \left\{ 1 - 2(1+C)\epsilon_{1*} - C\epsilon_{2*} + \left(\frac{\pi^2}{2} - 3 + 2C + 2C^2\right) \epsilon_{1*}^2 \\ &+ \left(\frac{7\pi^2}{12} - 6 - C + C^2\right) \epsilon_{1*} \epsilon_{2*} + \left(\frac{\pi^2}{8} - 1 + \frac{C^2}{2}\right) \epsilon_{2*}^2 + \left(\frac{\pi^2}{24} - \frac{C^2}{2}\right) \epsilon_{2*} \epsilon_{3*} \\ &+ \left[- 2\epsilon_{1*} - \epsilon_{2*} + (2+4C)\epsilon_{1*}^2 + (-1+2C)\epsilon_{1*} \epsilon_{2*} + C\epsilon_{2*}^2 - C\epsilon_{2*} \epsilon_{3*} \right] \ln \left(\frac{k}{k_*}\right) \\ &+ \left[2\epsilon_{1*}^2 + \epsilon_{1*} \epsilon_{2*} + \frac{1}{2}\epsilon_{2*}^2 - \frac{1}{2}\epsilon_{2*} \epsilon_{3*} \right] \ln^2 \left(\frac{k}{k_*}\right) + \cdots \right\}, \end{aligned}$$

$$\mathcal{P}_h &= \frac{2H_*^2}{\pi^2 M_{\rm Pl}^2} \left\{ 1 - 2(1+C)\epsilon_{1*} + \left[-3 + \frac{\pi^2}{2} + 2C + 2C^2 \right] \epsilon_{1*}^2 + \left[-2 + \frac{\pi^2}{12} - 2C - C^2 \right] \epsilon_{1*} \epsilon_{2*} \\ &- \left[-2\epsilon_{1*} + (2+4C)\epsilon_{1*}^2 + (-2-2C)\epsilon_{1*} \epsilon_{2*} \right] \ln \left(\frac{k}{k_*}\right) + \left(2\epsilon_{1*}^2 - \epsilon_{1*} \epsilon_{1*} \right) \ln^2 \left(\frac{k}{k_*}\right) + \cdots \right\}. \end{aligned}$$

• Currently known at third order: arXiv:2205.12608 (involves ϵ_{4*})

Reheating consistent model predictions

Quick check for which reheating history a model is compatible with the data



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To speed-up data analysis for 287 models

$$\mathcal{L}_{\text{eff}}(\boldsymbol{D}|P_*,\varepsilon_1,\varepsilon_2,\varepsilon_3) \propto \int P(\boldsymbol{D}|\boldsymbol{\theta}_{\text{s}},P_*,\varepsilon_1,\varepsilon_2,\varepsilon_3,\varepsilon_4) \,\pi(\varepsilon_4) \,\pi(\boldsymbol{\theta}_{\text{s}}) \,\mathrm{d}\varepsilon_4 \mathrm{d}\boldsymbol{\theta}_{\text{s}}.$$

The full likelihood has 55 parameters and is built upon

- Planck 2020 post-legacy TT, TE, EE data (PR4/NPIPE maps)
- Large scale EE polarization (lowE)
- BICEP/Keck *B*-mode 2018 (arXiv:2110.00483)
- Small scale TE and EE from SPT-3G (arXiv:2103.13618)
- Baryon Acoustic Oscillations (SDSS collaboration)
- MCMC exploration of the 55 dimensions
 - COSMOMC up to 25 million samples $(R 1 < 10^{-3})$
 - GetDist marginalization in 4D: $P_*, \varepsilon_1, \varepsilon_2, \varepsilon_3$
- Basic machine learning: 1 hidden layer with 300 nodes

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Marginalized posteriors

Effective likelihood vs exact: 2D



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Effective likelihood vs exact: 1D



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Computing bayesian evidences

Probability of a model ${\cal M}$ to explain the data $oldsymbol{D}$

$$P(\mathcal{M}|\mathbf{D}) = \frac{\mathcal{E}(\mathbf{D}|\mathcal{M}) P(\mathcal{M})}{P(\mathbf{D})}$$

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$$P(\mathcal{M}|\mathbf{D}) = \frac{\mathcal{E}(\mathbf{D}|\mathcal{M}) P(\mathcal{M})}{P(\mathbf{D})}$$

Bayesian evidence

$$\mathcal{E}(\boldsymbol{D}|\mathcal{M}) \propto \int \mathcal{L}_{\mathrm{eff}}(\boldsymbol{D}|P_*, \varepsilon_1, \varepsilon_2, \varepsilon_3) \, \pi(\boldsymbol{\theta}_{\mathrm{inf}}, R_{\mathrm{rad}}) \, \mathrm{d}\boldsymbol{\theta}_{\mathrm{inf}} \, \mathrm{d}R_{\mathrm{rad}}$$

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Bayesian evidence

$$\mathcal{E}(\boldsymbol{D}|\mathcal{M}) \propto \int \mathcal{L}_{\mathrm{eff}}(\boldsymbol{D}|P_*, \varepsilon_1, \varepsilon_2, \varepsilon_3) \, \pi(\boldsymbol{ heta}_{\mathrm{inf}}, R_{\mathrm{rad}}) \, \mathrm{d} \boldsymbol{ heta}_{\mathrm{inf}} \, \mathrm{d} R_{\mathrm{rad}}$$

- Computed with BAYASPIC running over 287 models
 - ◆ BAYASPIC \equiv ASPIC + PolyChord + $\mathcal{L}_{ ext{eff}}$
 - A few cpu-hours per model \mathcal{M} (but up to 2 days for some)
- $1 \,\mathrm{TB}$ of data output (nested chains, posteriors, plots,...)

Bayes factors for all models



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Information gain on the reheating

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Kullback-Leibler divergence between the prior and posterior

$$D_{\rm KL}^{\rm rad} = \int P(\ln R_{\rm rad} | \boldsymbol{D}) \ln \left[\frac{P(\ln R_{\rm rad} | \boldsymbol{D})}{\pi(\ln R_{\rm rad})} \right] d\ln R_{\rm rad},$$



Information gain on the reheating

Kullback-Leibler divergence between the prior and posterior

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Within the space of the single-field models $\mathcal{I}_{ ext{mod}} \equiv \{\mathcal{M}_i\}$

Posterior probability of the running $\alpha_s \simeq -\epsilon_{2*} \left(2\epsilon_{1*} + \epsilon_{3*} \right)$

$$P(\alpha_{s}|\boldsymbol{D}, \mathcal{I}_{mod}) = \sum_{i} P(\alpha_{s}|\boldsymbol{D}, \mathcal{M}_{i}) P(\mathcal{M}_{i}|\boldsymbol{D})$$

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| Qualitative Inflation Quantitative inflation Bayesian inference Predictions in model space & Running of the spectral index & Model space vs slow-roll space @Reheating energy density and equation of state Conclusion | • | Non-trivial pred The sign of α_s |
|--|---|--|
| | | |

- Non-trivial prediction coming from both theoretical prior + data
- The sign of $\alpha_{\rm S}$ remains undetermined otherwise

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Non-trivial prediction coming from both theoretical prior + data

- The sign of $\alpha_{
 m s}$ remains undetermined otherwise
 - Posterior by assuming "just slow-roll" (no model) + data



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Non-trivial prediction coming from both theoretical prior + data

- The sign of $\alpha_{
 m s}$ remains undetermined otherwise
 - Model space prior (no data input)



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Non-trivial prediction coming from both theoretical prior + data

- The sign of $lpha_{
 m s}$ remains undetermined otherwise
- The field evolution in single-field models creates a correlation between the sign of $n_{\rm s} 1$ and the sign of the running $\alpha_{\rm s}!$



← Models with $n_{\rm S} \lesssim 1$ have all an accelerated Hubble radius in the observable window $\implies \alpha_{\rm S} < 0$

Reheating energy density and equation of state

Posterior probability for $(\ln \rho_{\rm reh}, \overline{w}_{\rm reh})$

$$P(\ln \rho_{\rm reh}, \overline{w}_{\rm reh} | \boldsymbol{D}, \mathcal{I}_{\rm mod}) = \sum_{i} P(\ln \rho_{\rm reh}, \overline{w}_{\rm reh} | \boldsymbol{D}, \mathcal{M}_{i}) P(\mathcal{M}_{i} | \boldsymbol{D})$$

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$$P(\ln \rho_{\rm reh}, \overline{w}_{\rm reh} | \boldsymbol{D}, \mathcal{I}_{\rm mod}) = \sum_{i} P(\ln \rho_{\rm reh}, \overline{w}_{\rm reh} | \boldsymbol{D}, \mathcal{M}_{i}) P(\mathcal{M}_{i} | \boldsymbol{D})$$



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Bayesian data analysis in model space \mathcal{I}_{mod}

- Enforces model consistency + new insights on the reheating era
- Predicts: $\langle \alpha_{\rm s} \rangle = -7.3 \times 10^{-4}$
- Data constraining power is winning against theoretical proposals



Looking forward to the Euclid, LSS & CMB-S4 data!