



Universitat  
de les Illes Balears

IAC3

Institut d'Aplicacions  
Computacionals  
de Codi Comunitari

RINGDOWN

INSPIRAL

MERGER

# Waveform modelling

Eleanor Hamilton



---

LISA SCHOOL FOR EARLY CAREER SCIENTISTS – 09.10.25

# Modelling approaches

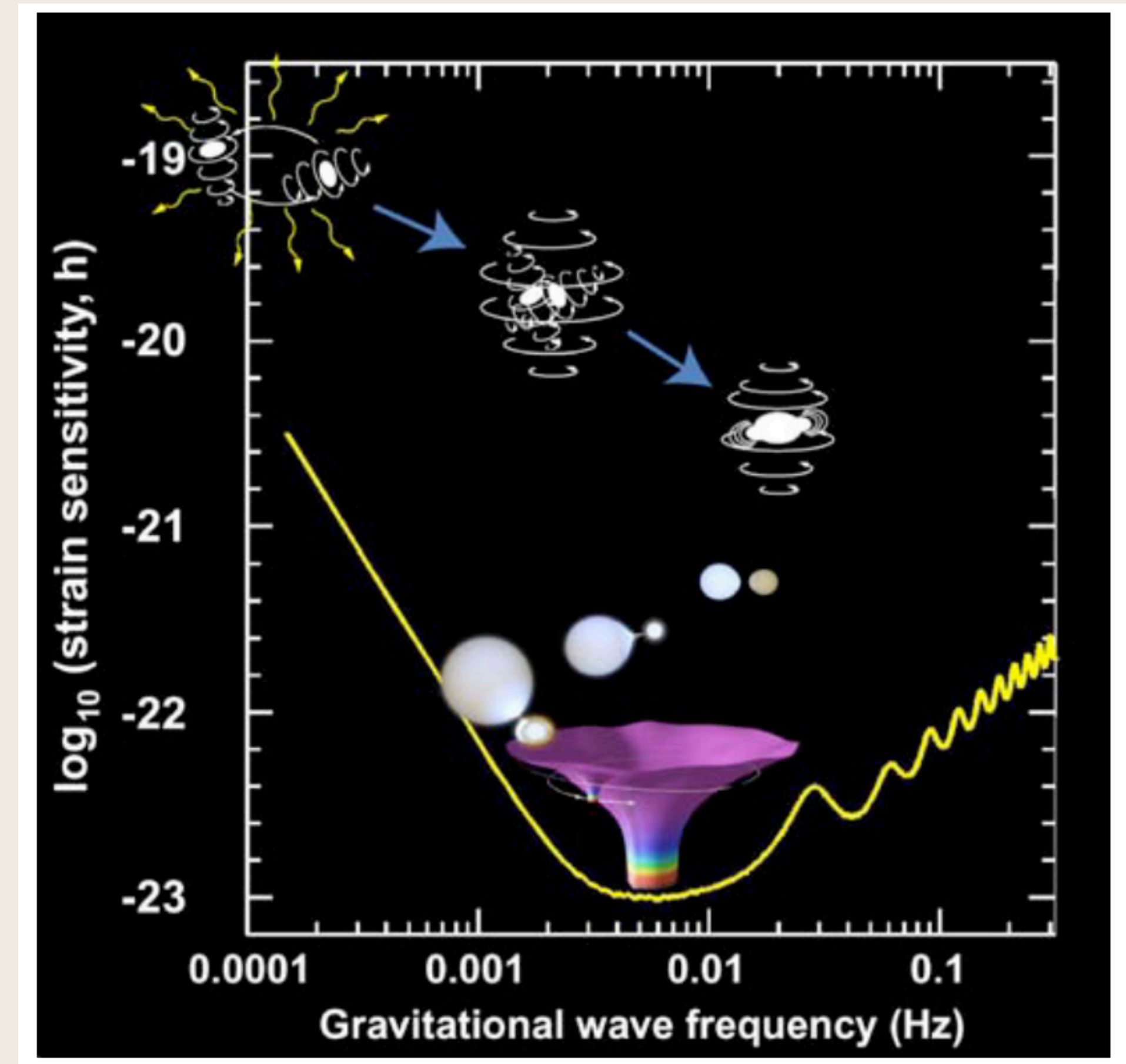
For different sources

---



# LISA sources

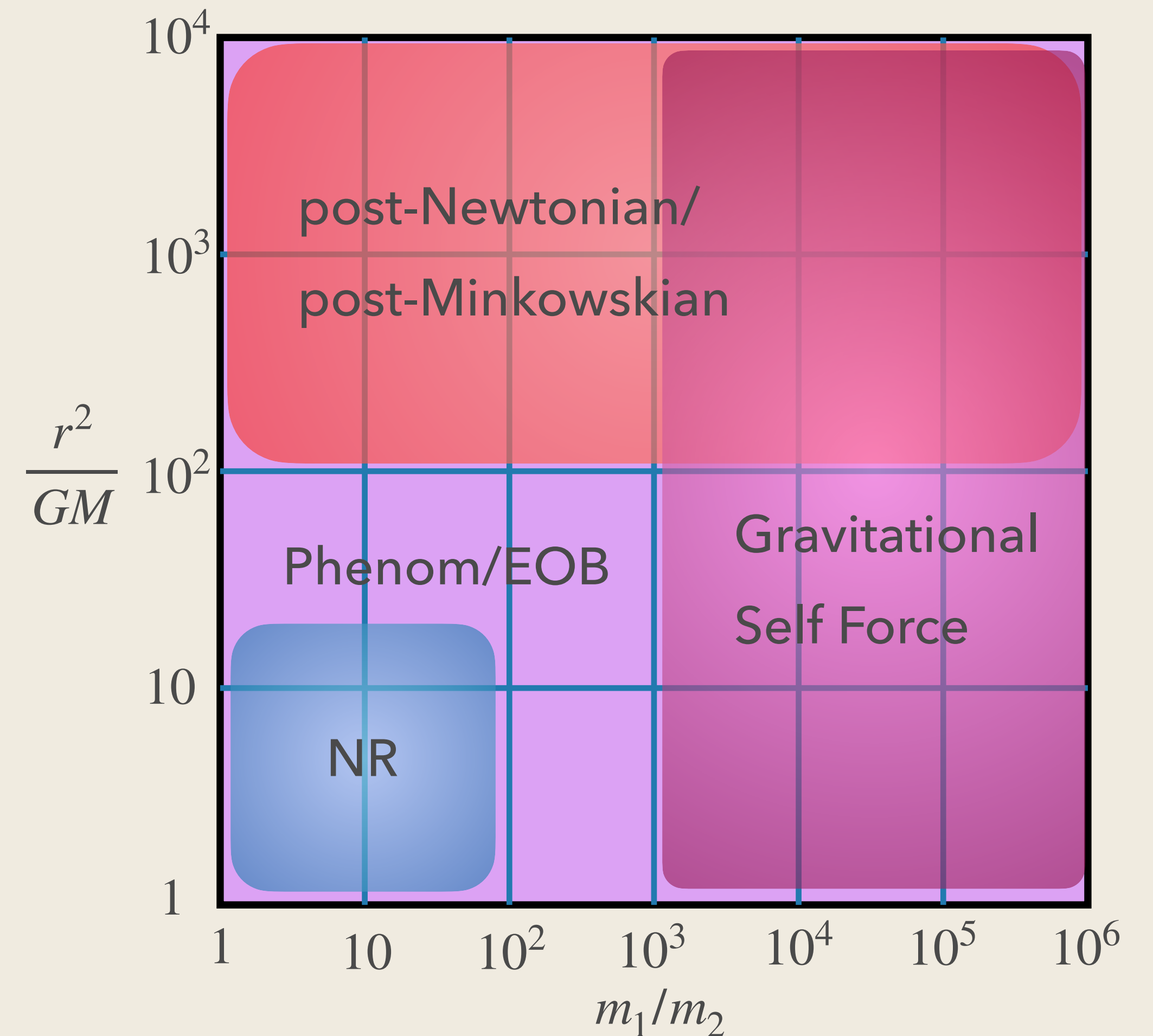
- Main sources for LISA:
  - Massive Black Hole Binaries (MBHBs)
  - Galactic binaries
  - Extreme Mass Ratio Inspirals



LISA-LIST team. LISA: Probing the universe with gravitational waves.  
Technical report, LISA Mission Science Office, 2007.

# Modelling approaches

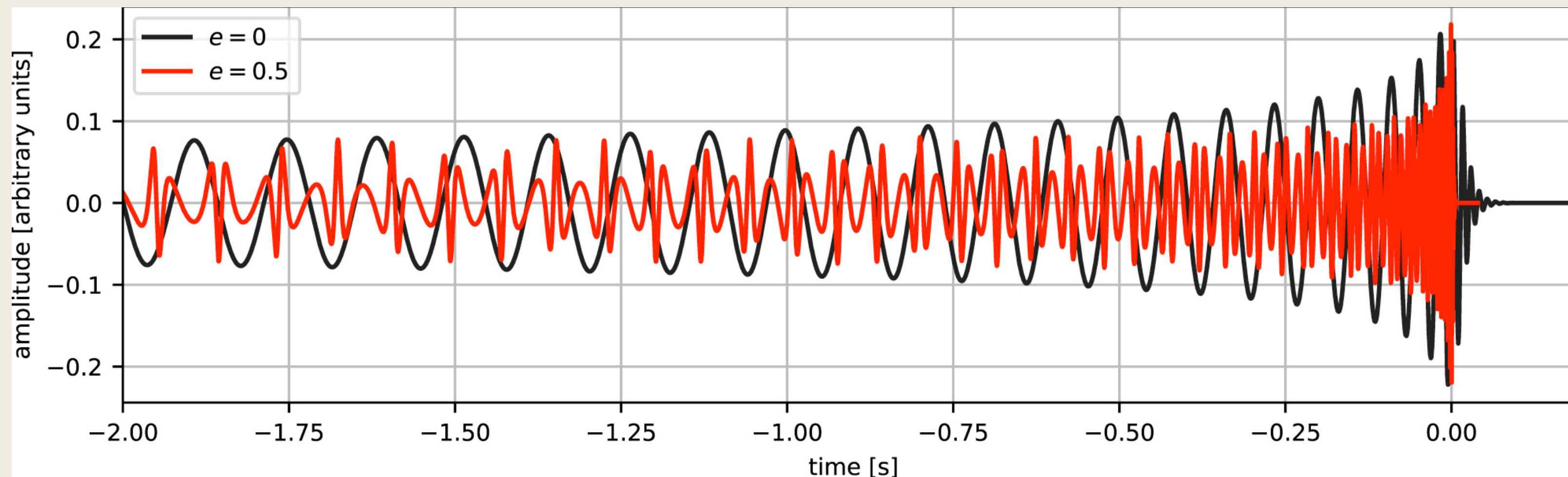
- 3 main approaches to resolving the spacetime around radiating binaries:
  - Numerical relativity
  - post-Newtonian/post-Minkowskian theory
  - Gravitational self-force
- Effective formulations to cover larger regions of parameter space:
  - Effective-one-body
  - Phenom



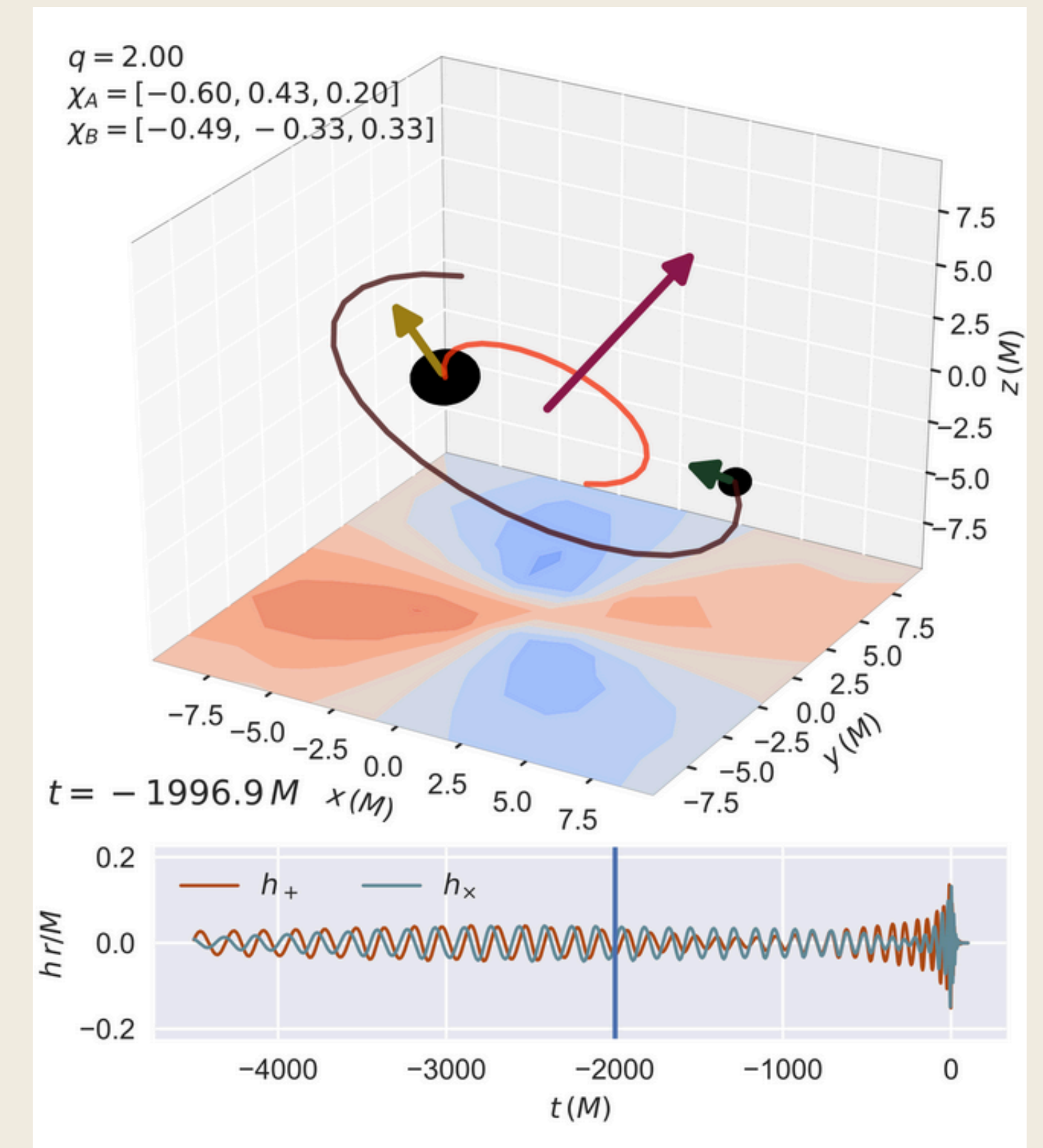


# Physical Phenomena

- Non-spinning/aligned-spin systems
- Eccentric orbits
- Arbitrarily oriented spins
- Fully generic binaries
- Asymmetric masses



Abbott, B. P., et al. "Search for eccentric binary black hole mergers with Advanced LIGO and Advanced Virgo during their first and second observing runs." *The Astrophysical Journal* 883.2 (2019): 149.



Varma, Vijay, Leo C. Stein, and Davide Gerosa. "The binary black hole explorer: on-the-fly visualizations of precessing binary black holes." *Classical and Quantum Gravity* 36.9 (2019): 095007.



# Effective-One-Body formalism

- Maps two-body dynamics onto that of an effective test mass/spin in a deformed Schwarzschild/Kerr background

1. Produce **Hamiltonian** describing **conservative** dynamics

$$H_{\text{EOB}} = M \sqrt{1 + 2\nu \left( \frac{H_{\text{eff}}}{\mu} - 1 \right)}$$

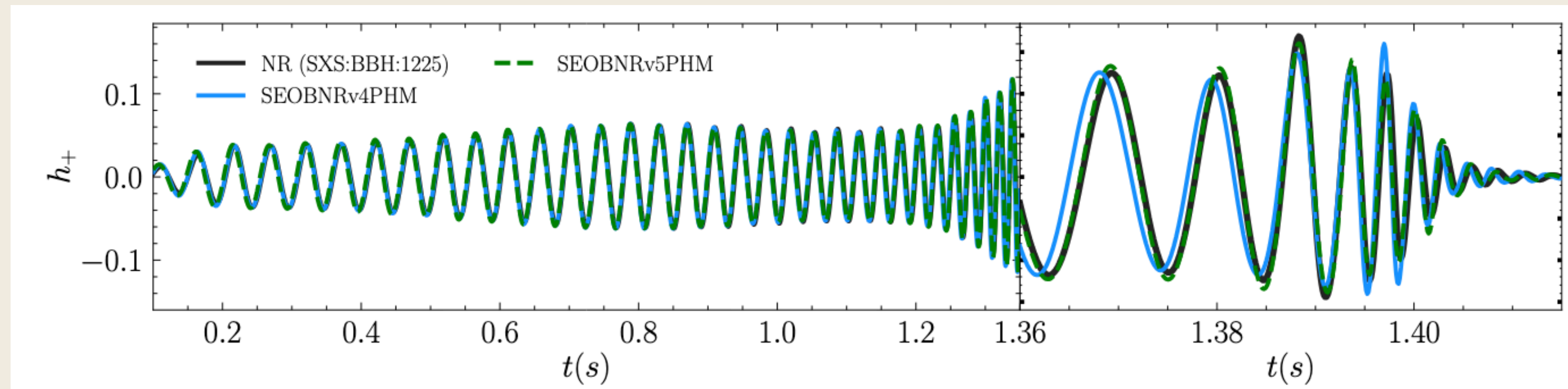
for an effective Hamiltonian  $H_{\text{eff}}$  of a test mass  $\mu$  moving in a Kerr spacetime of mass  $M$  deformed by amount  $\nu$

2. Introduce **radiation-reaction**



# Effective-One-Body: SEOB

- Combines information from analytical methods with numerical relativity data
- Includes:
  - Tidal and generic spin interactions
  - Generic orbits
- Suitable for IMRIs/MBHBs

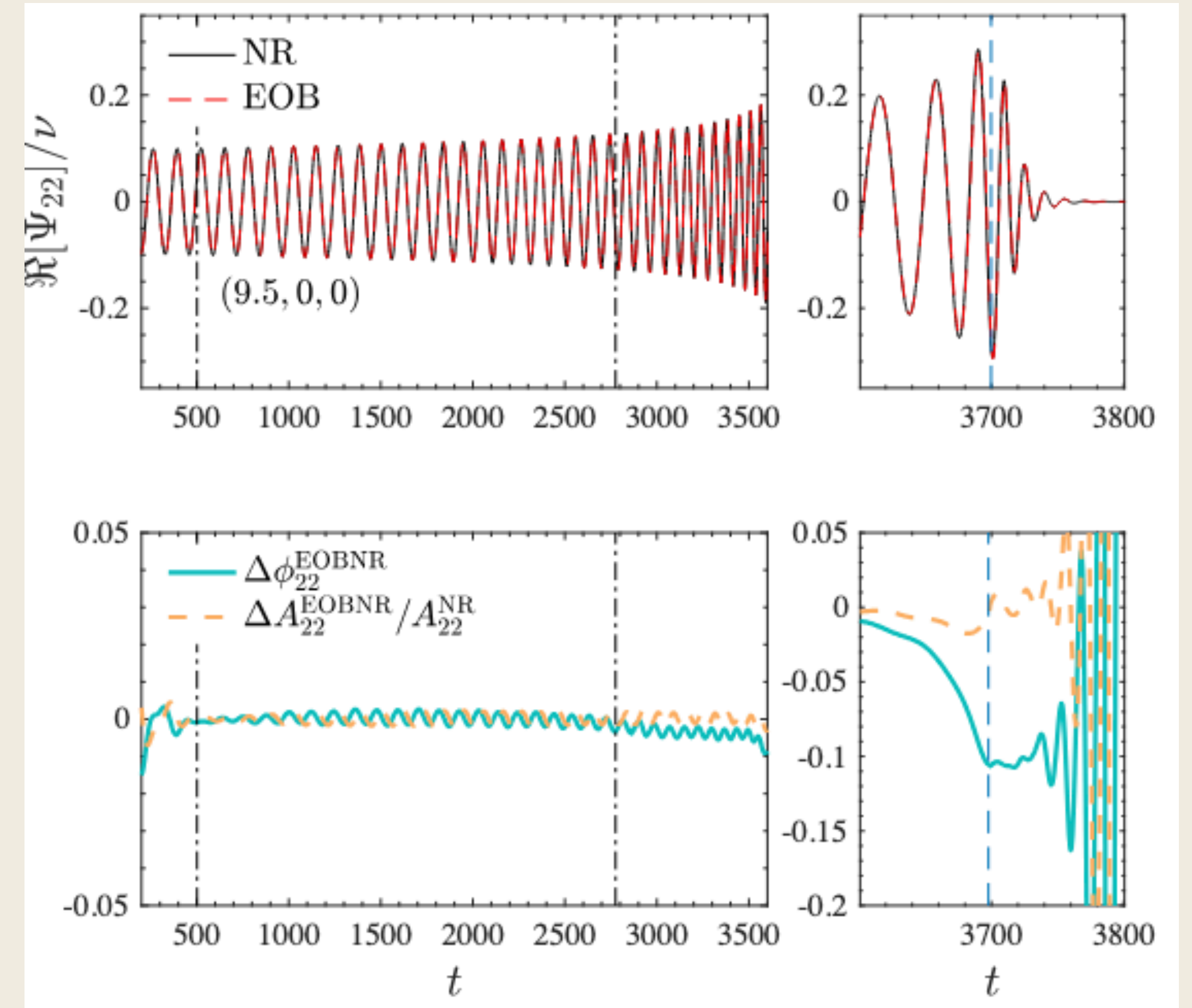


Ramos-Buades, Antoni, et al. "Next generation of accurate and efficient multipolar precessing-spin effective-one-body waveforms for binary black holes." *Physical Review D* 108.12 (2023): 124037.



# Effective-One-Body: TEOB

- Includes systematic resumption of EOB potentials via Padé approximants
- Includes:
  - Tidal and generic spin interactions
  - Generic orbits (arbitrary eccentricity, hyperbolic, scattering)
  - Gravitational self-force
- Minimal NR calibration

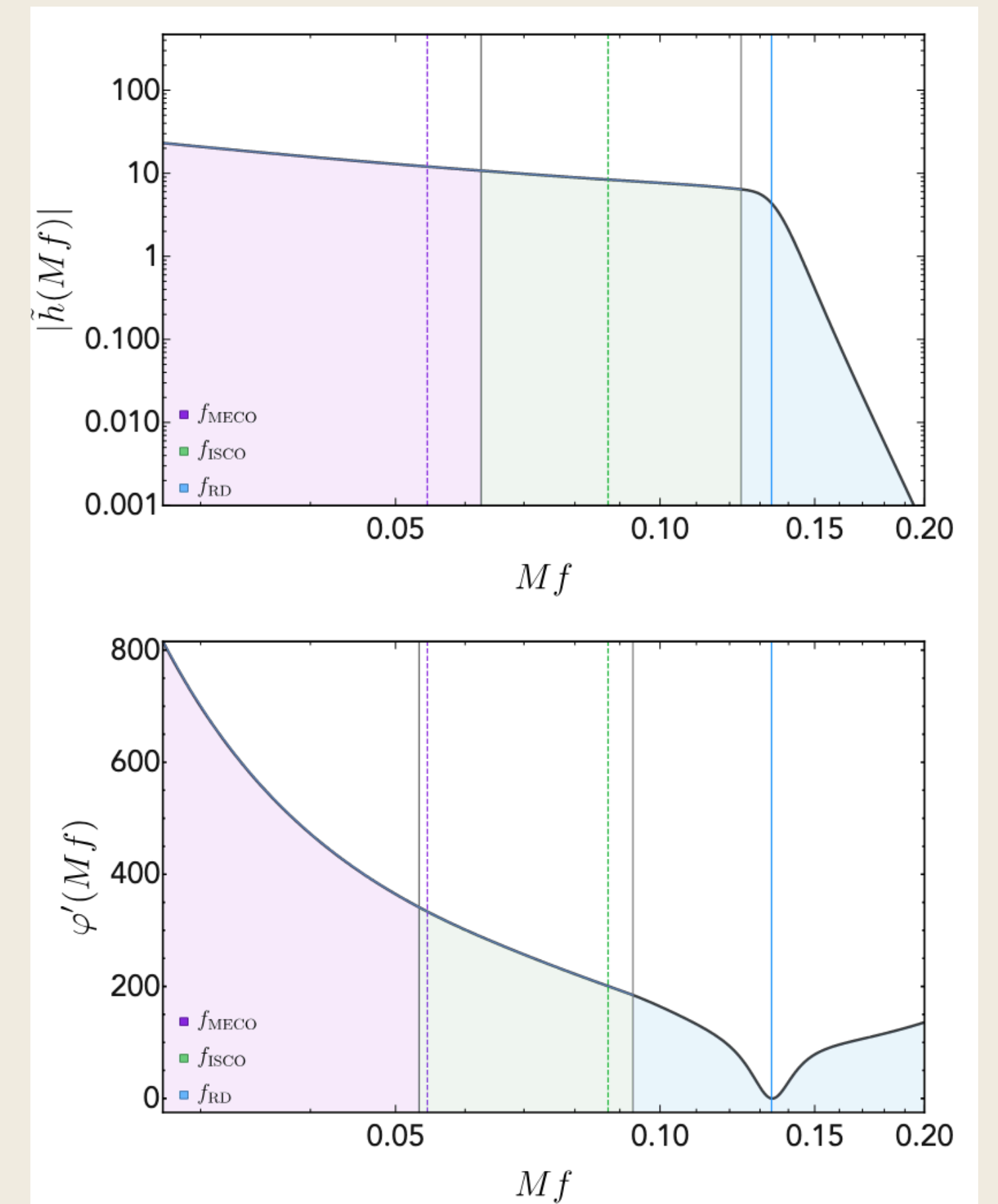


Nagar, Alessandro, et al. "Analytic systematics in next generation of effective-one-body gravitational waveform models for future observations." *Physical Review D* 108.12 (2023): 124018.



# Phenomenological models

- Complete IMR model; evaluates down to arbitrary frequencies
- Largely relies on closed-form expressions to ensure efficiency
- Construction of model:
  - Appropriate piece-wise ansatz
  - Fit co-efficients to calibration data set
  - Co-efficients interpolated across parameter space
- Suitable for MBHBs/IMRIs

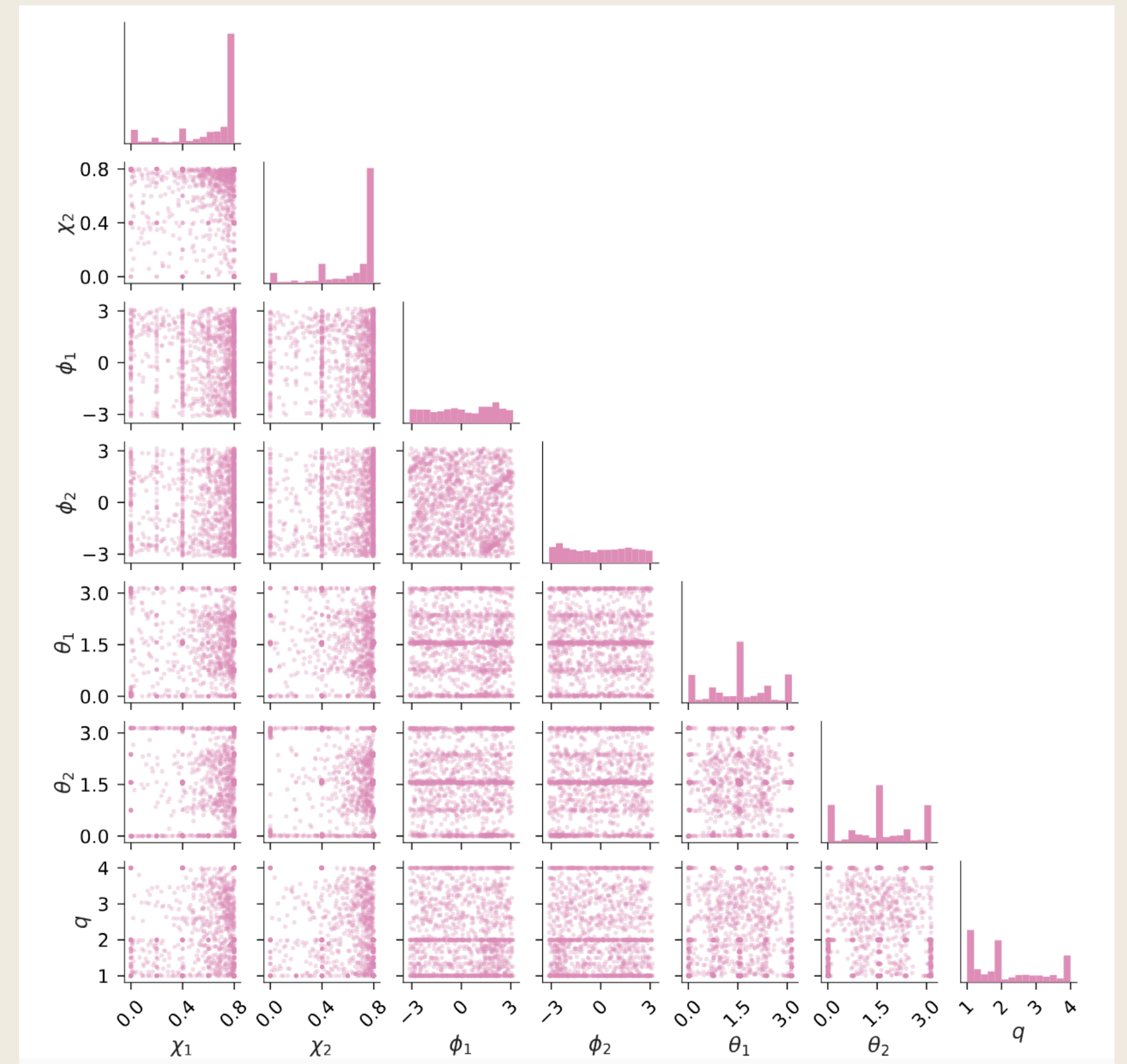


Pratten, G, et al. Physical Review D 102.6 (2020): 064001.



# Surrogate models

- A **data driven** approach
- Interpolates large data sets with no assumptions about the phenomenology of the underlying waveforms
  - **Accurate**
  - **Finite length**
- Range of available models:
  - NRSur7dq4 (MBHBs)
  - BHPTNRSurrogate (MBHB/IMRI)



Varma, V., et al. "Surrogate models for precessing binary black hole simulations with unequal masses." *Phys Rev Research* 1.3 (2019): 033015.



---

LISA SCHOOL FOR EARLY CAREER SCIENTISTS – 09.10.25

# Modelling requirements

From data analysis

---



# Requirements

- Ideal:
  - Infinitely accurate
  - Highly computationally efficient
- Errors in models:
  - Hamper detection and identification of sources
  - Affect tests of GR, rate and population estimates, formation models...
  - Introduce residuals when subtracting sources



# Accuracy measures

- Can look at dephasing, errors in the fits, etc.
- When using waveforms, the more interesting quantity when assessing accuracy is the match

$$M(h_1, h_2) = \max \left[ 4\text{Re} \int \frac{h_1(f)h_2^*(f)}{S_n(f)} df \right]$$

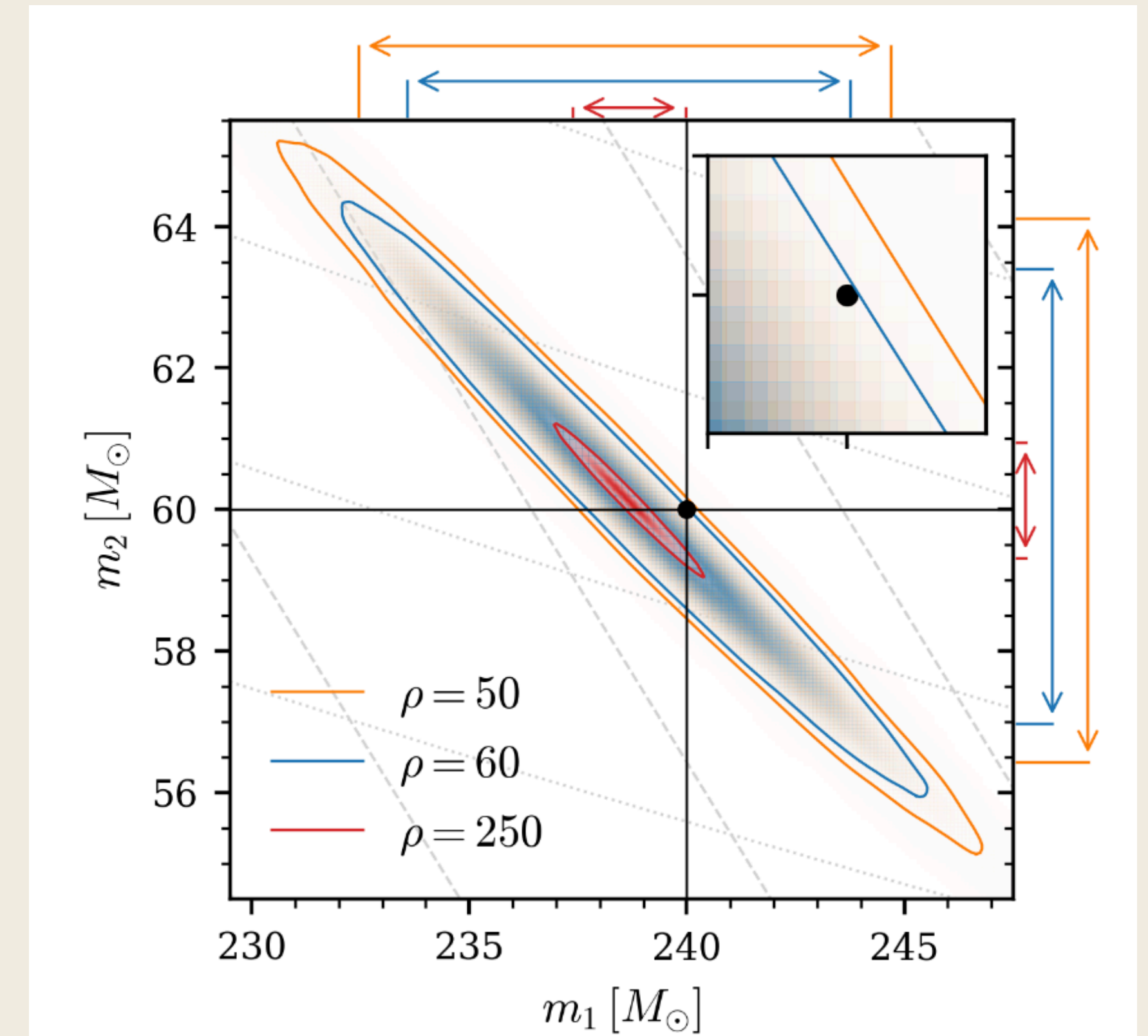
$$\mathcal{M} = 1 - M(h_1, h_2)$$

- Ideally, all models would include a measure of their error across the parameter space
- Need to understand how traditional accuracy measure relate to bias in parameter estimation



# Accuracy measures and biases

- Ideally, all models would include a measure of their error across the parameter space
- Need to understand how traditional accuracy measure relate to bias in parameter estimation  
e.g. distinguishability criterion



Thompson, Jonathan E., et al. "Use and interpretation of signal-model indistinguishability measures for gravitational-wave astronomy." *Physical Review D* 112.6 (2025): 064011.



---

LISA SCHOOL FOR EARLY CAREER SCIENTISTS – 09.10.25

# Case study:

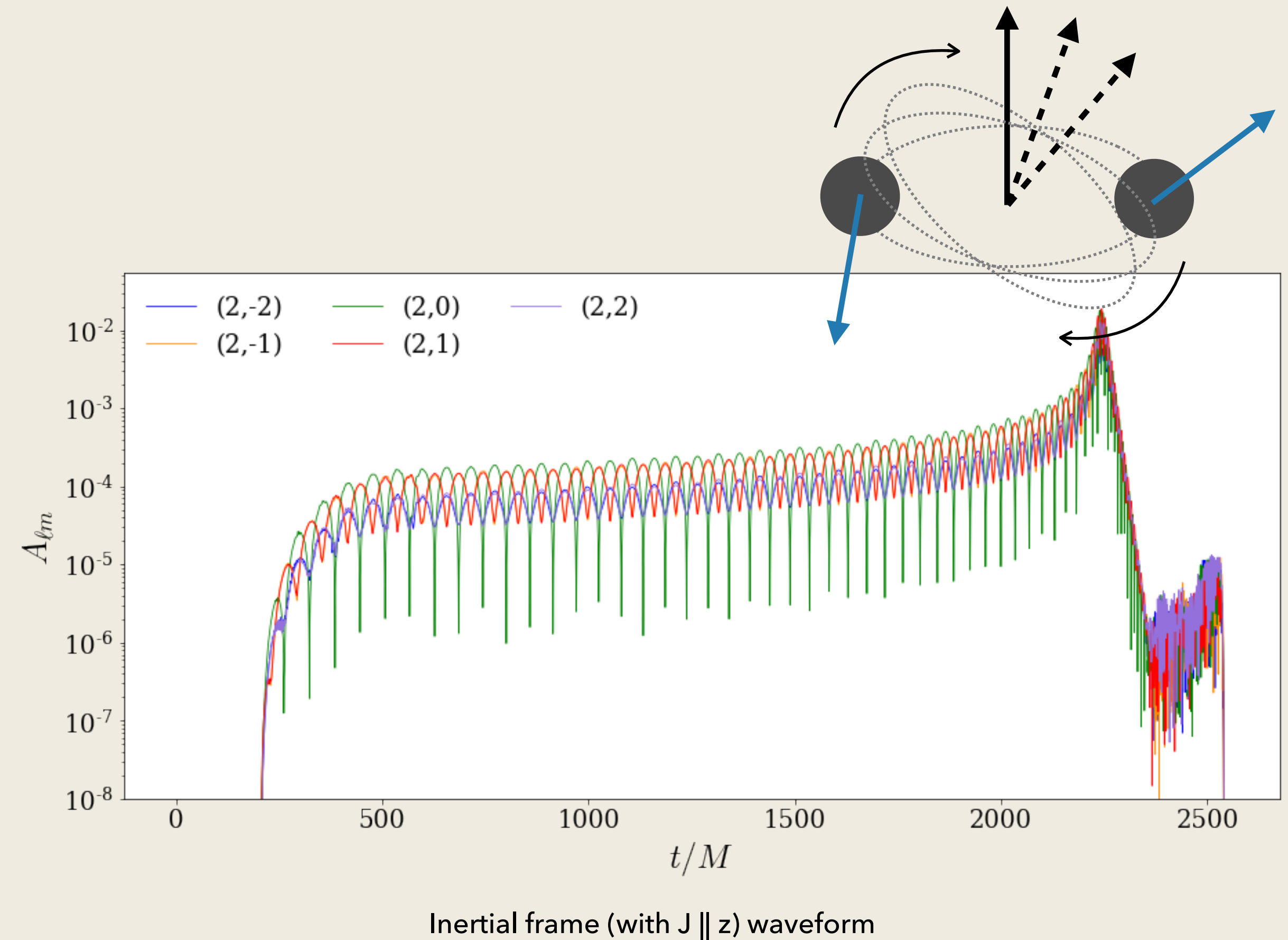
Creating a phenomenological  
precessing model

---



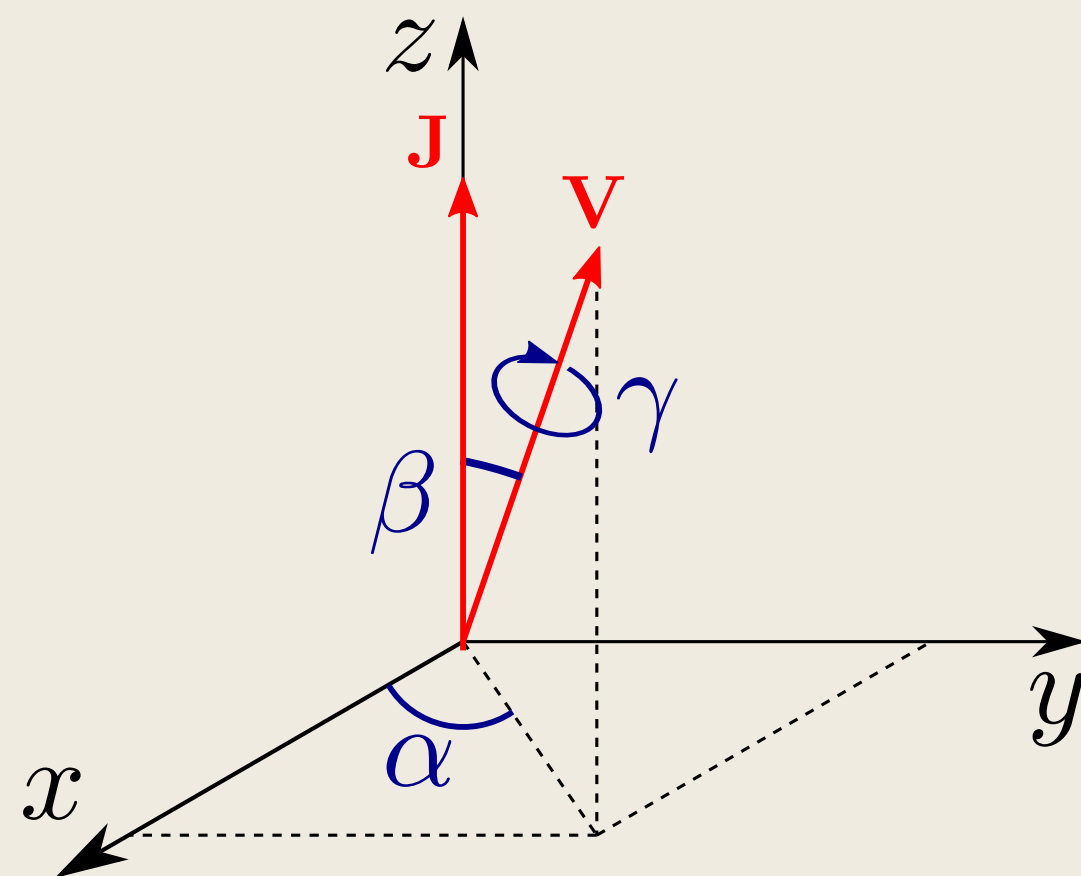
# Precessing systems

- A **precessing** binary is one where the spin on one or both of the black holes is misaligned with the orbital angular momentum
- This introduces **oscillations** in the amplitude and phase of the waveform, making it harder to model

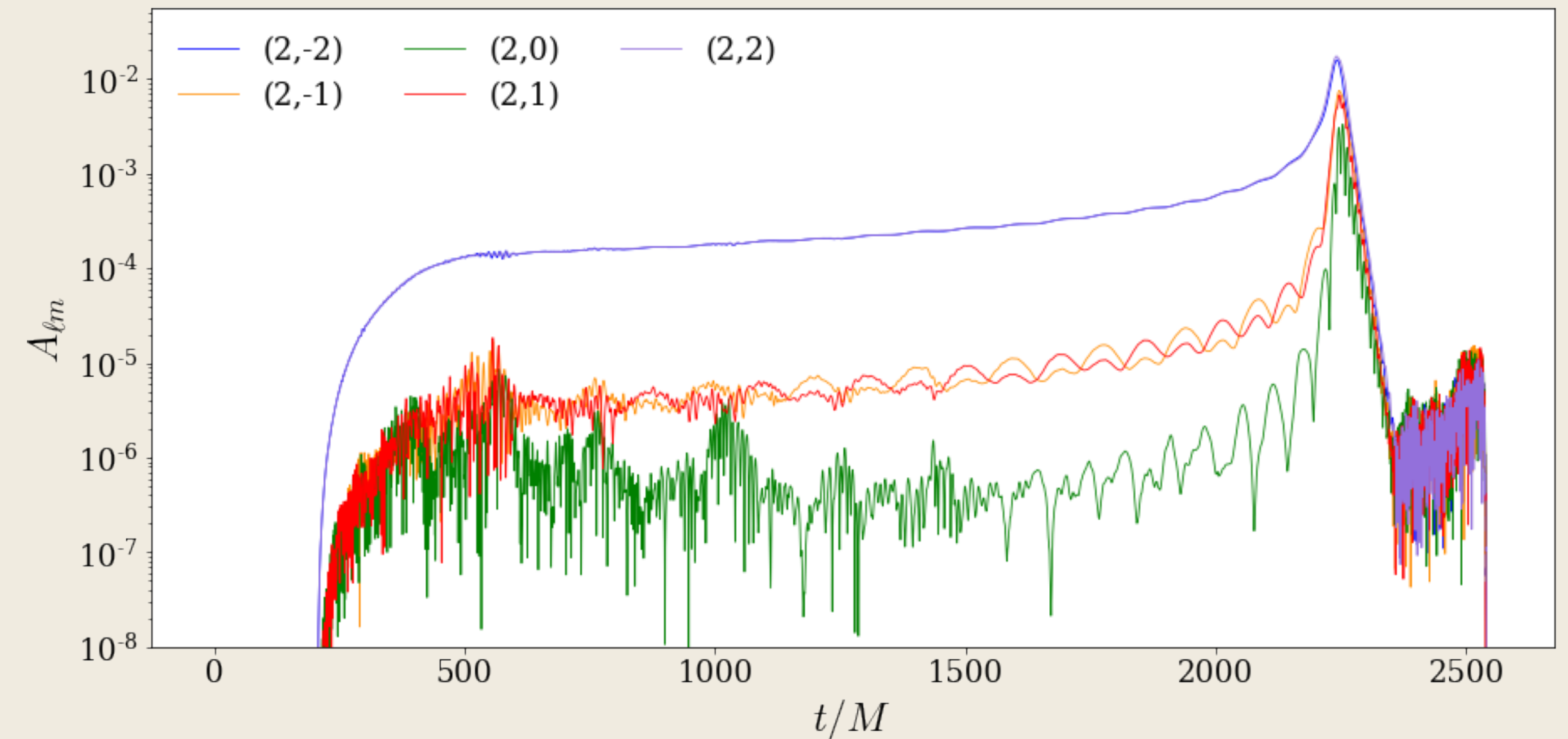


# Modelling precessing systems

- To simplify the modelling of a precessing system, we use the **co-precessing frame**, where the waveform resembles a non-precessing one
- We now need to model the **co-precessing waveform** and the **precession dynamics**



Representation of the three precession angles.  $V$  indicates the optimal emission direction (or the orbital angular momentum)



Co-precessing frame waveform



# Co-processing waveform

- **Co-processing model**
  - Based on PhenomXAS
    - Calibrated to hybrid EOB + NR hybrids
    - Employs PN TaylorF2 in inspiral with pseudo-PN terms calibrated to hybrid data
    - Phenomenological merger-ringdown
    - Connected via collocation points

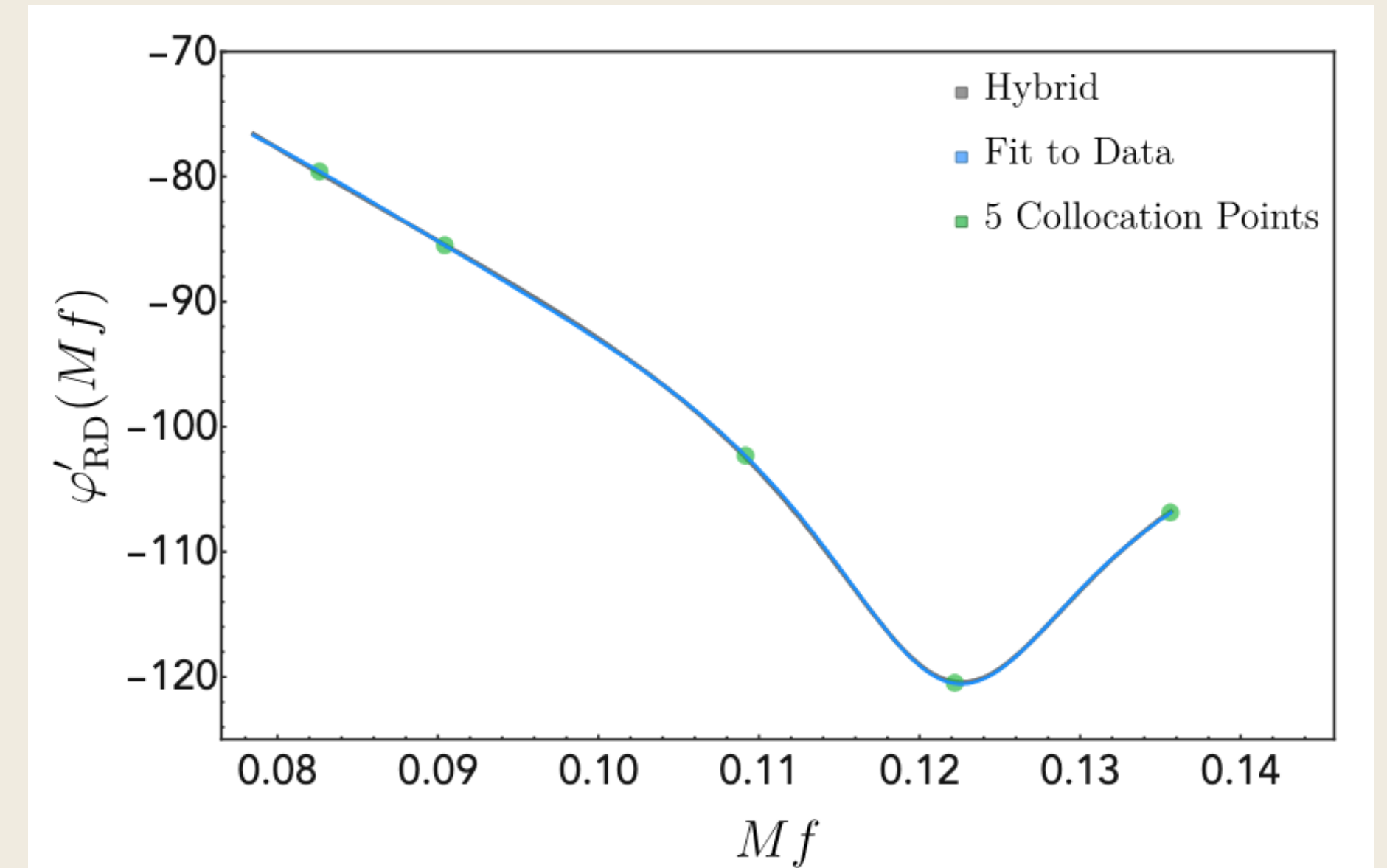
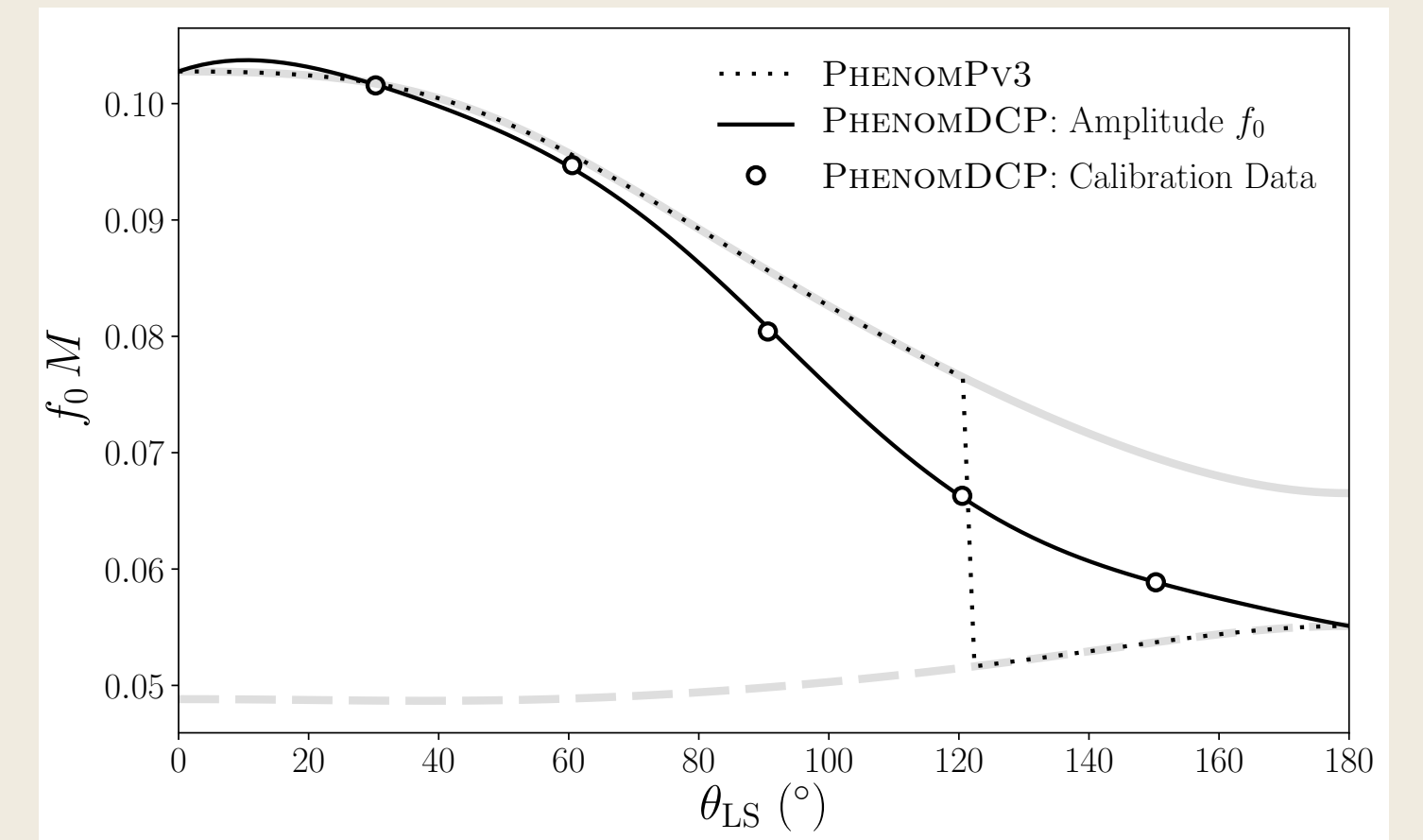


FIG. 10. Phase derivative for SXS-BBH-0153 hybridized against SEOBNRv4 in the merger-ringdown region along with the fit to the data (blue) reconstructed from a system of collocation points (green).

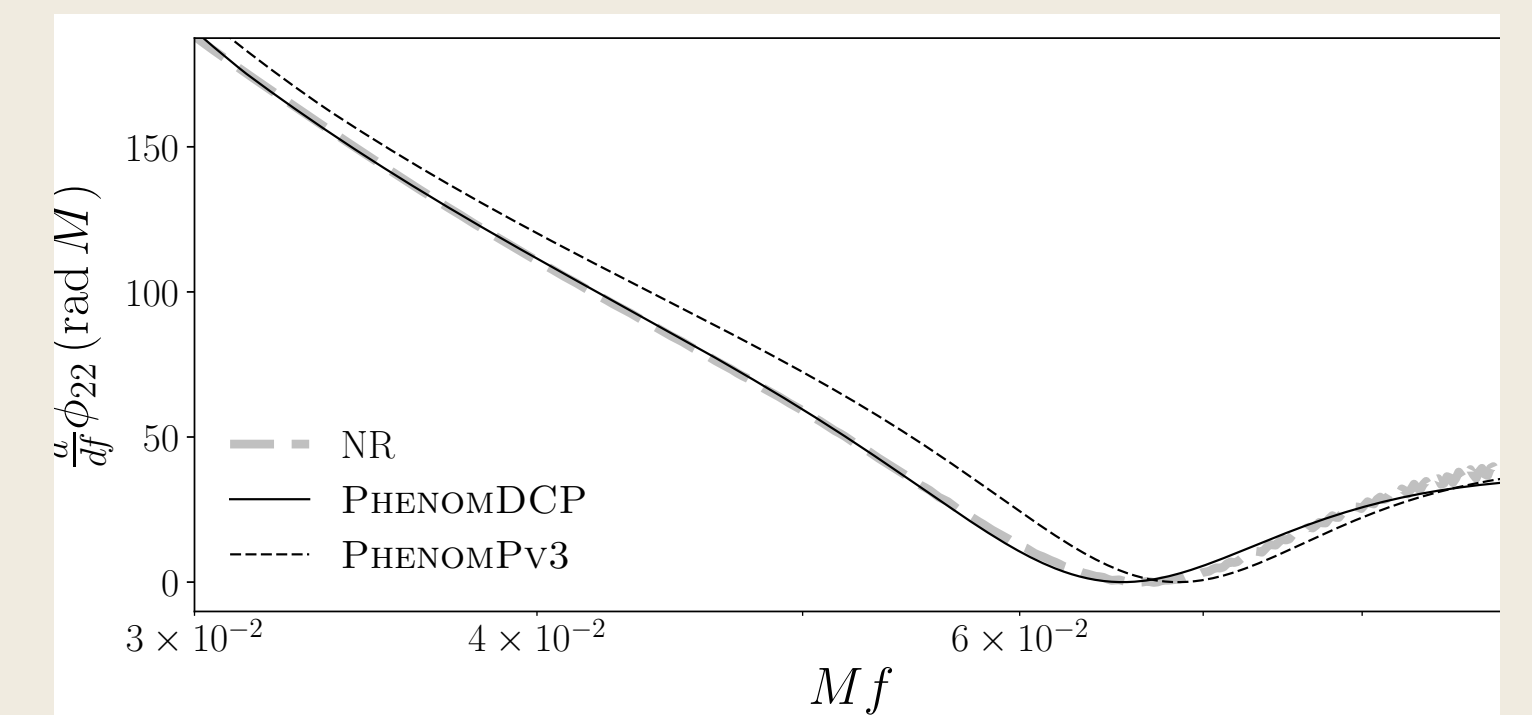
Pratten, Geraint, et al. "Setting the cornerstone for a family of models for gravitational waves from compact binaries: The dominant harmonic for nonprecessing quasicircular black holes." *Physical Review D* 102.6 (2020): 064001.

# Co-precessing waveform

- **Co-precessing model**
  - Tuned parameters in merger-ringdown ansatz to precessing simulations
  - Key feature: effective ringdown frequency
  - Includes calibrated multipole asymmetry



Evolution of the ringdown frequency across the parameter space for cases with mass ratio 8, dimensionless spin 0.8

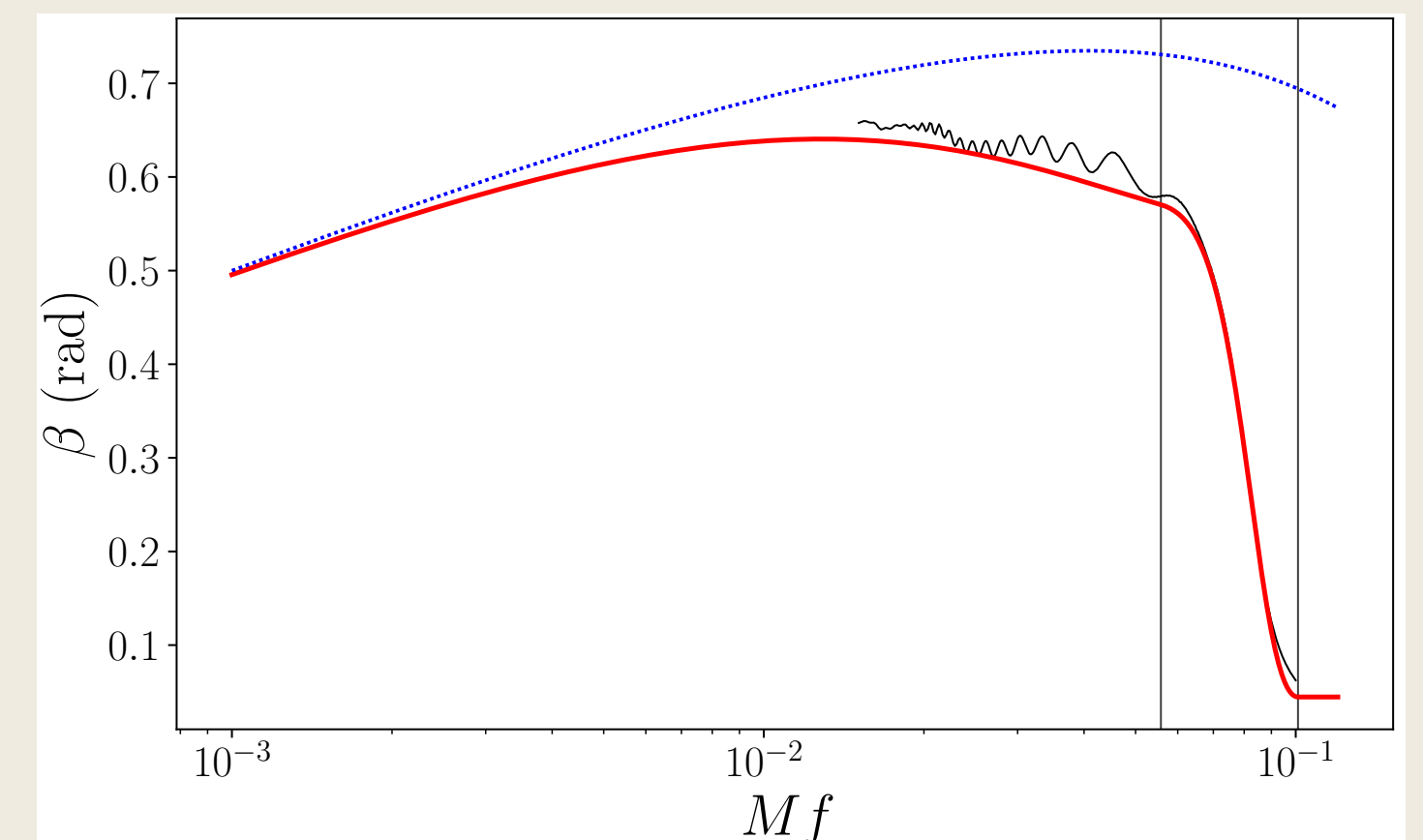
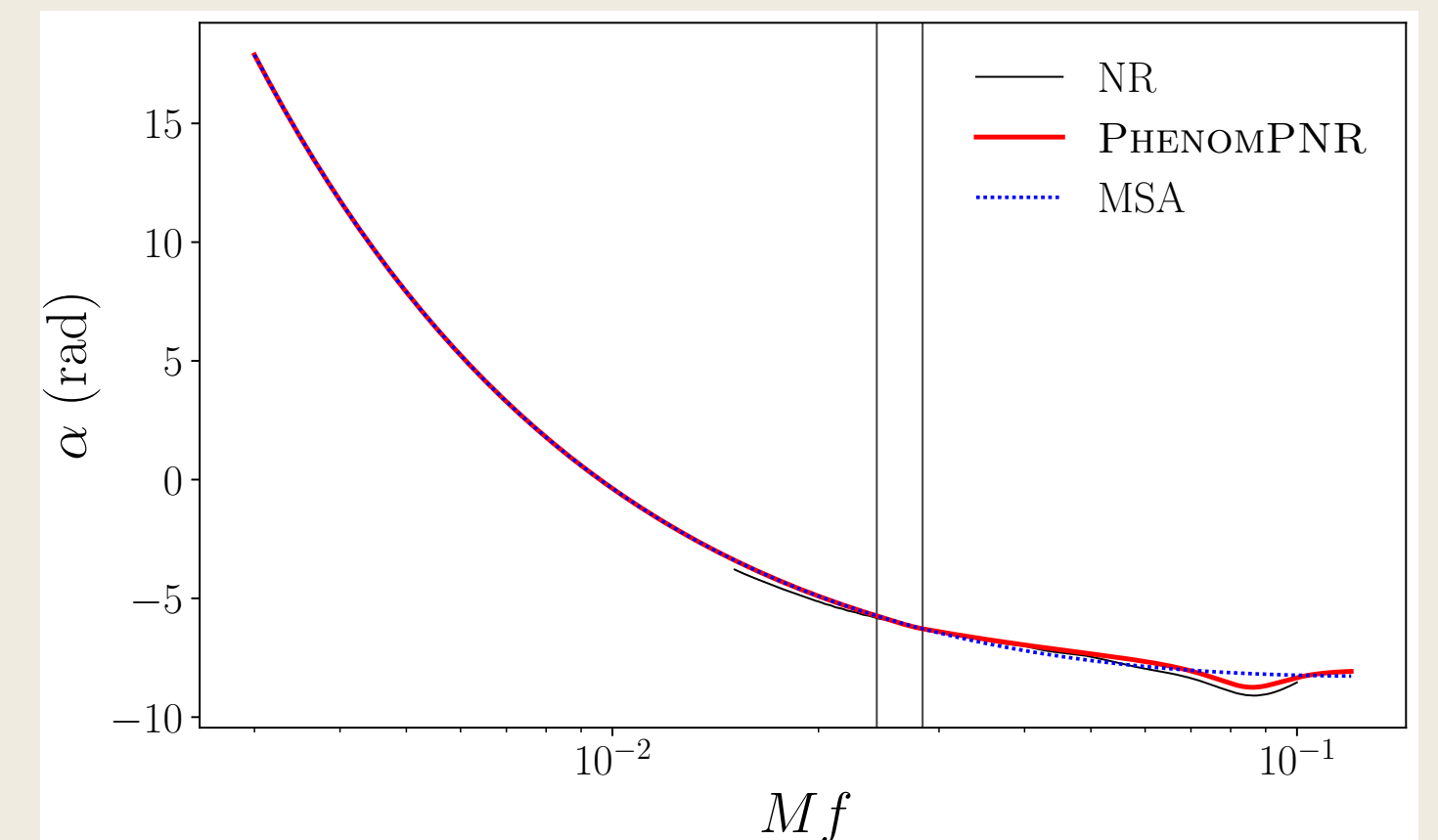


Phase derivative



# Dominant multipole precession dynamics

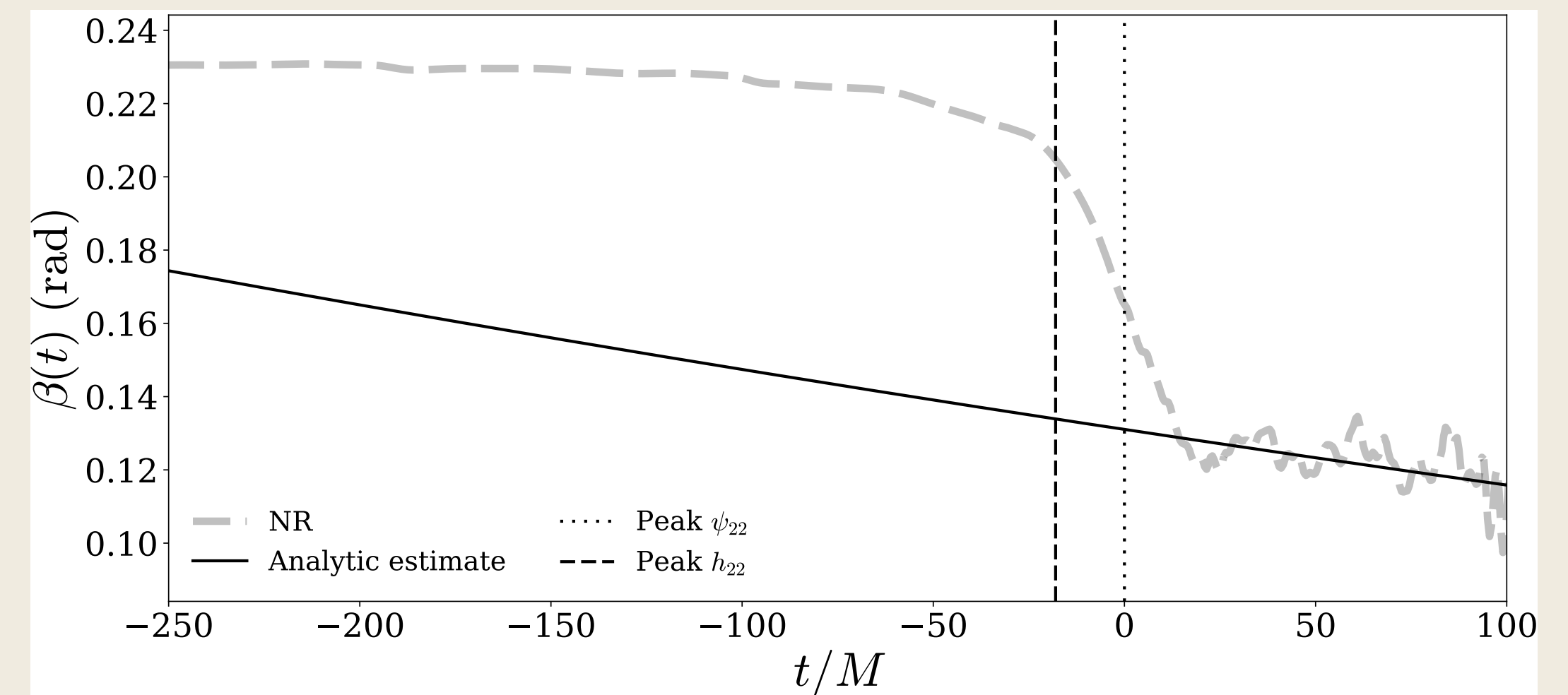
- **Precession angle model**
  - Use SpinTaylor PN expressions for angles in inspiral
  - Phenomenological ansatz describes merger-ringdown behaviour
  - Captures key physical features
    - Ringdown value of  $\beta$
    - Mode hierarchy
  - $\gamma$  calculated using minimal rotation condition





# Collapse of $\beta$ through merger

- $\beta$  drops rapidly during merger
- Inspiral can be described by PN theory and ringdown can be described by perturbation theory
- Drop itself cannot be described by either theory



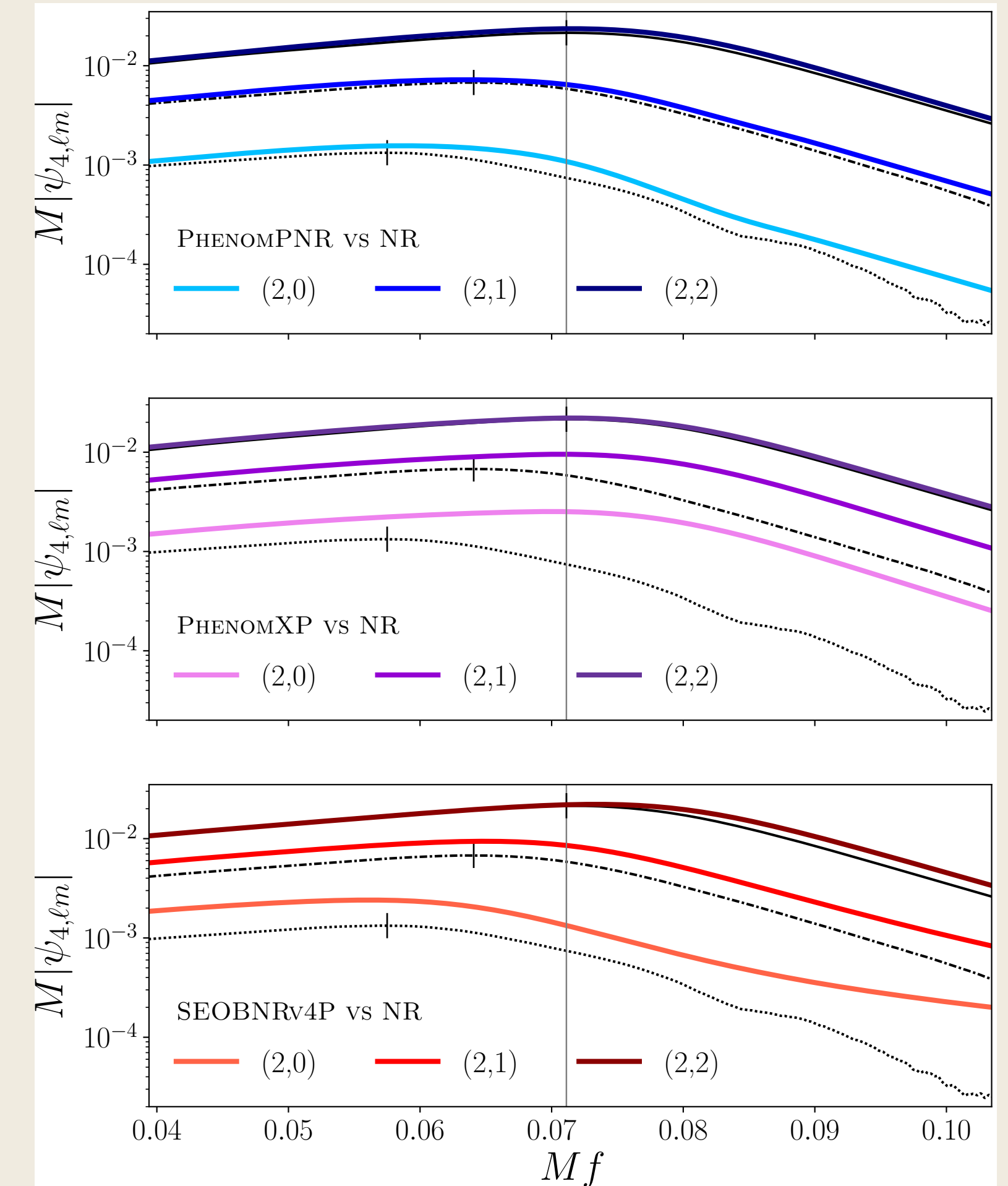
$$q = 8, \chi = 0.4, \theta_{LS} = 30^\circ$$



# Mode hierarchy

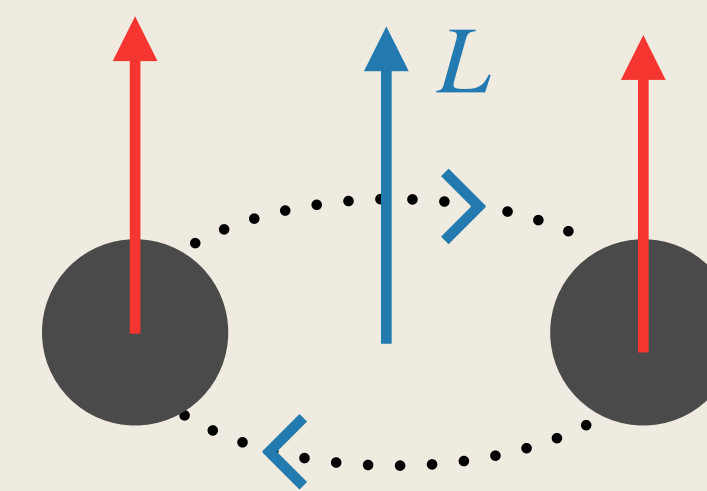
- In precessing waveforms we see a hierarchy in the turnover frequency of precessing waveforms
- This is related to the rapid drop off in beta since

$$\beta = 2 \arctan \frac{|h_{21}^J|}{|2h_{22}^J|}$$

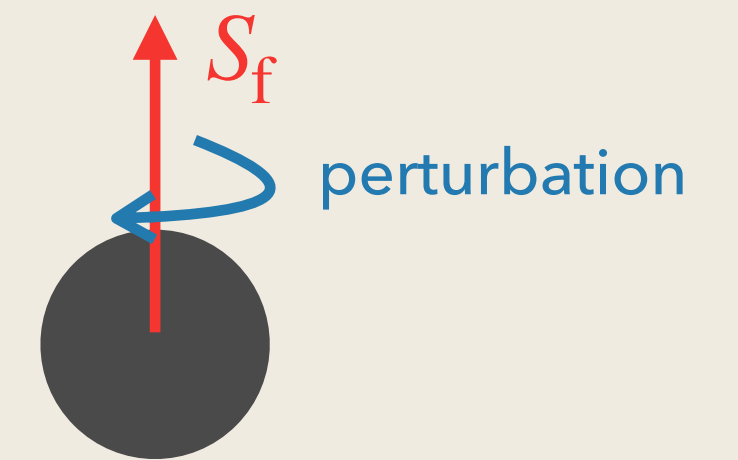


# Ringdown frequency

- The remnant black hole formed after a binary-black-hole merger emits gravitational waves at characteristic frequencies, known as quasi-normal modes (QNMs)
- Frequencies determined by perturbation theory but amplitudes unknown
- For each QNM we have both a prograde and a retrograde frequency
- Basic picture: if the perturbation is in the same direction as the final spin, prograde frequencies are excited and vice versa, where the perturbation is in the same direction as orbital direction

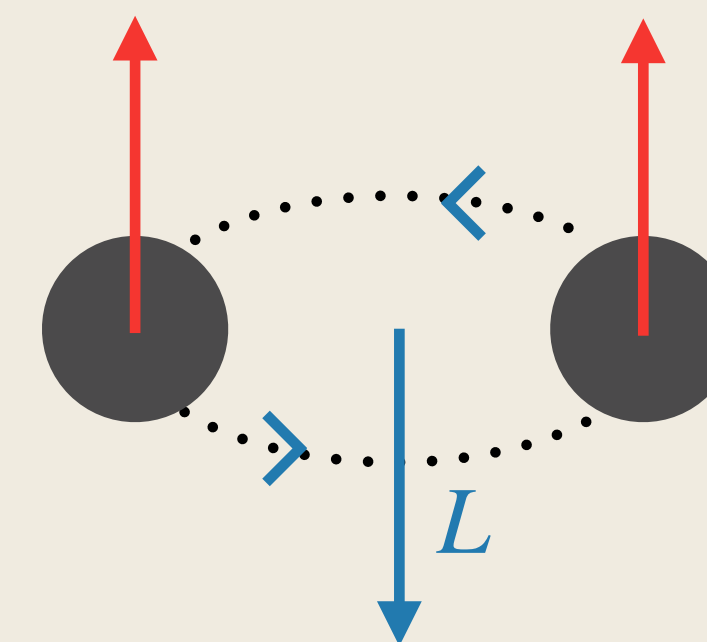


Initial binary

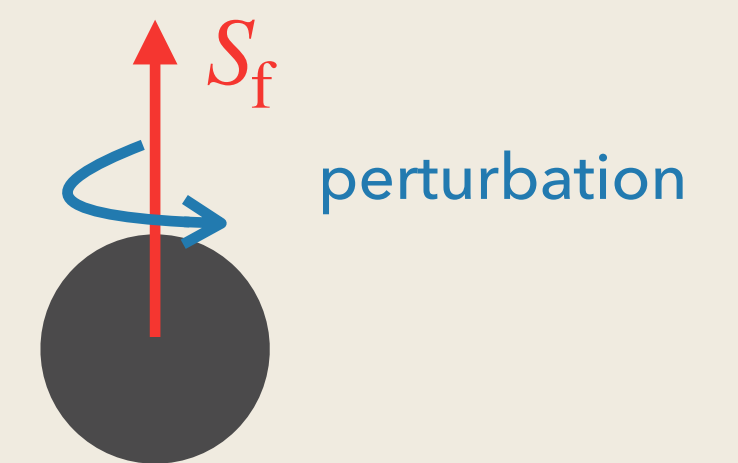


Remnant

**Prograde** frequencies excited



Initial binary



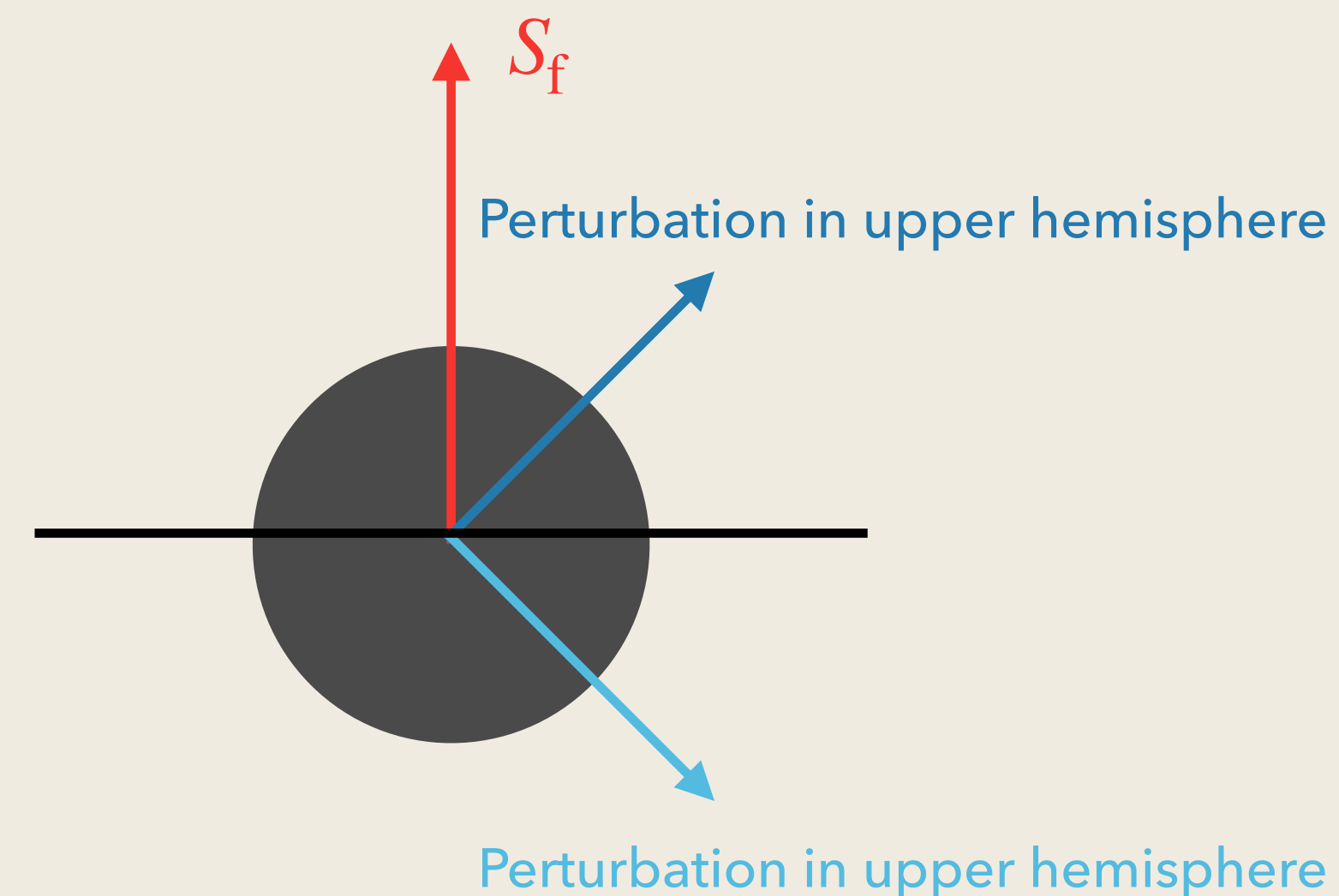
Remnant

**Retrograde** frequencies excited



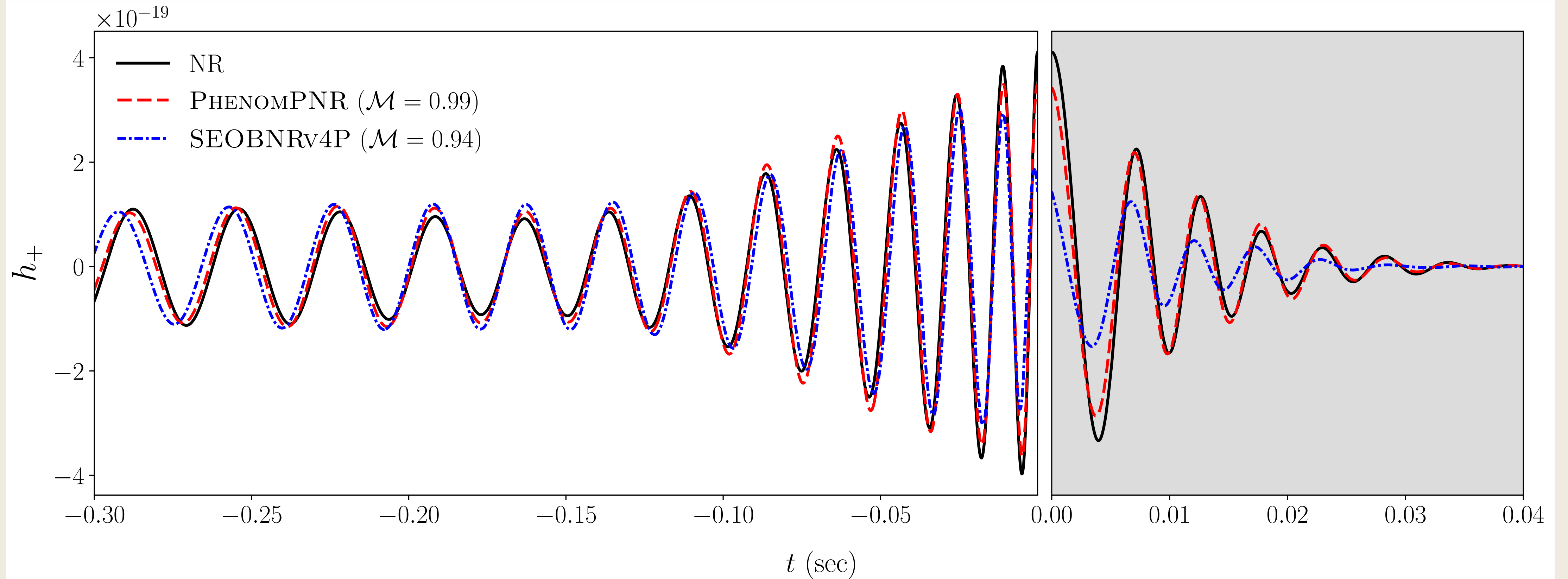
# Precessing picture

- In precessing systems:
  - Axis of the perturbation is not aligned with the final spin
  - Difficult to quantify the direction of the perturbation
- **Both** prograde **and** retrograde modes are excited
- In co-precessing frame get a "mean frequency"
 
$$\omega'_{\ell m'} = \omega_{\ell m'} - m'(1 - |\cos \beta|)(\omega_{22} - \omega_{21})$$



- "positive final spin"
- Prograde frequencies dominate
- "negative final spin"
- Retrograde frequencies dominate

# Time domain waveforms



Time domain comparison of  
 $q = 8, \chi = 0.8, \theta_{LS} = 60^\circ$



---

LISA SCHOOL FOR EARLY CAREER SCIENTISTS – 09.10.25

# Current IMR models

MBHB sources

---



# Quasi-circular models

	Modeling approach	Complete inspiral	Calibration (aligned sector)	Calibration (precessing sector)	Extrapolation region
<b>NRSur7dq4</b>	Surrogate	×	$q \leq 4;  \chi_{1,2}  \leq 0.8$	$q \leq 4;  \chi_{1,2}  \leq 0.8$	$q \leq 6;  \chi_{1,2}  \leq 1.0$
<b>PhenomTPHM</b>	Semi-analytic	✓	$q \leq 8;  \chi_{1,2}  \leq 0.8$	×	$q \leq 20;  \chi_{1,2}  \leq 1.00$
<b>PhenomXPHM</b>	Semi-analytic	✓	$q \leq 18;  \chi_{1,2}  \leq 0.95$	×	$q \leq 20;  \chi_{1,2}  \leq 1.00$
<b>PhenomXPNR</b>	Semi-analytic	✓	$q \leq 18;  \chi_{1,2}  \leq 0.95$	$q \leq 8;  \chi_{1,2}  \leq 0.8$	$q \leq 20;  \chi_{1,2}  \leq 1.00$
<b>SEOBNRv5PHM</b>	Semi-analytic	✓	$q \leq 20;  \chi_{1,2}  \leq 0.95$	×	$q \leq 20;  \chi_{1,2}  \leq 1.00$
<b>TEOBResumS-Giotto</b>	Semi-analytic	✓	×	×	—



# Modelling challenges: higher modes

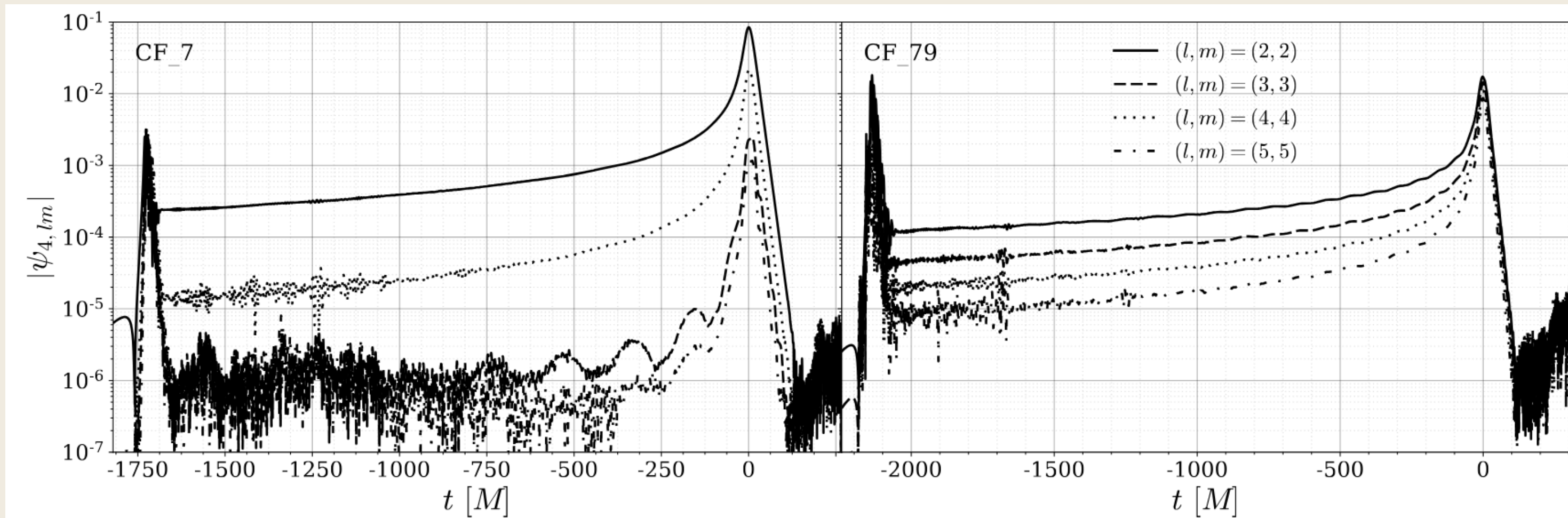


FIG. 4. Comparison of copreprocessing frame time domain amplitudes for the  $l = m$  modes for  $l \in \{2, 3, 4, 5\}$ . The left panel shows CF\_7 with initial parameters  $(q, \chi_2, \theta) = (1, 0.4, 60)$  and the right panel shows CF\_79 with parameters  $(q, \chi_2, \theta) = (8, 0.8, 120)$ .

Hamilton, E., et al. Physical Review D 109.4 (2024): 044032.

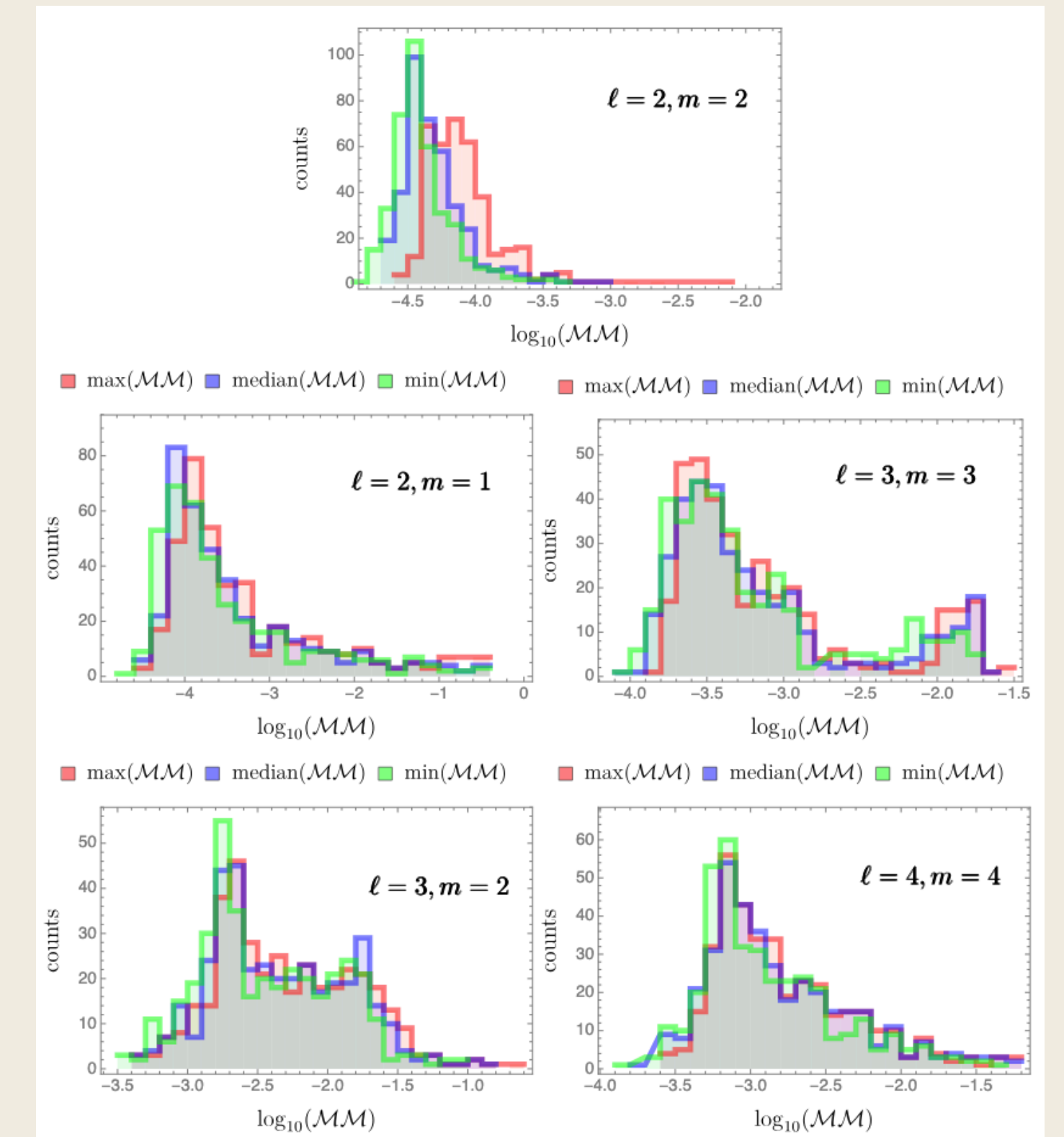


FIG. 13: Mode-by-mode mismatches between IMRPHENOMXHM and a validation set of hybrids built using the latest release of the SXS collaboration catalog. Each plot shows the maximum (red), median (blue) and minimum (green) mismatch over a range of total masses

García-Quirós, C., et al. PhysRevD 102.6(2020): 064002.



# QC models: model accuracy

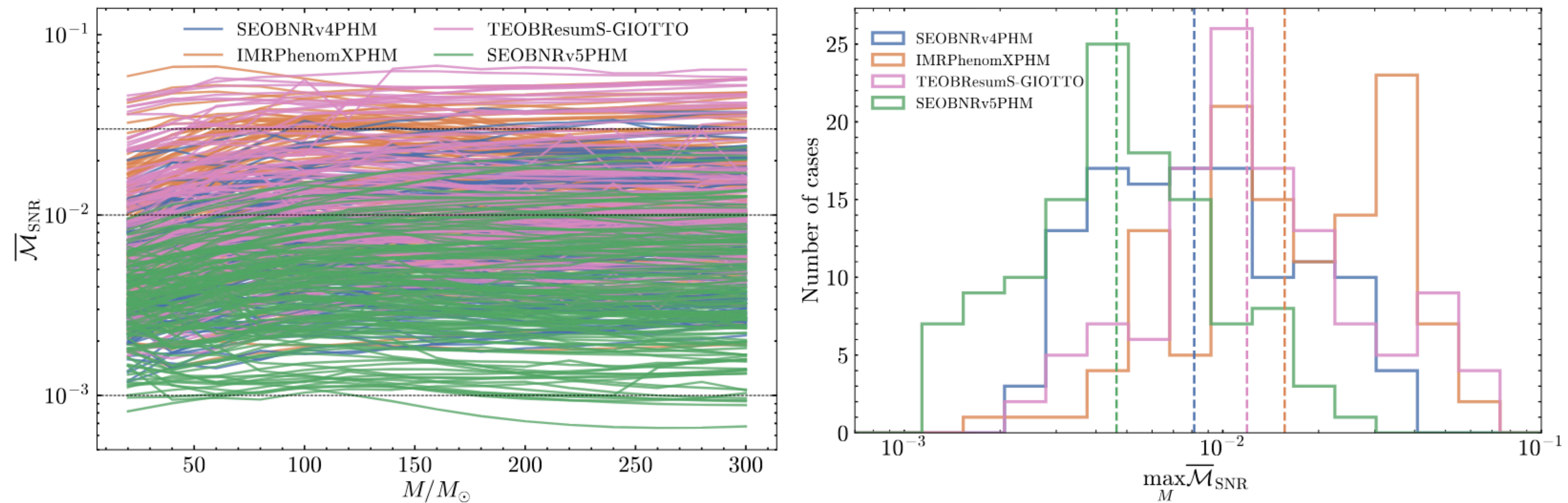


Figure 4. *Left panel:* Sky-and-polarization averaged, SNR-weighted unfaithfulness in the total mass range between  $[20 - 300]M_{\odot}$  for an inclination  $\iota = \pi/3$ , between SEOBNRv4PHM (blue), IMRPhenomXPHM (orange), TEOBResumS-GIOTTO (pink) and SEOBNRv5PHM (green) against NR for the 118 highly precessing-spin BBH simulations from Ref. [98]. The dashed horizontal vertical lines correspond to the  $10^{-3}$ , 0.01 and 0.03 unfaithfulness values. *Right panel:* Distribution of the maximum unfaithfulness over the total mass range for each NR simulation considered in the left plot. The vertical dashed lines indicate the median values of the distribution.

Ramos-Buades, A., et al. "Next generation of accurate and efficient multipolar precessing-spin effective-one-body waveforms for binary black holes." *Physical Review D*, 108.12 (2023): 124037.



# QC models: parameter recovery

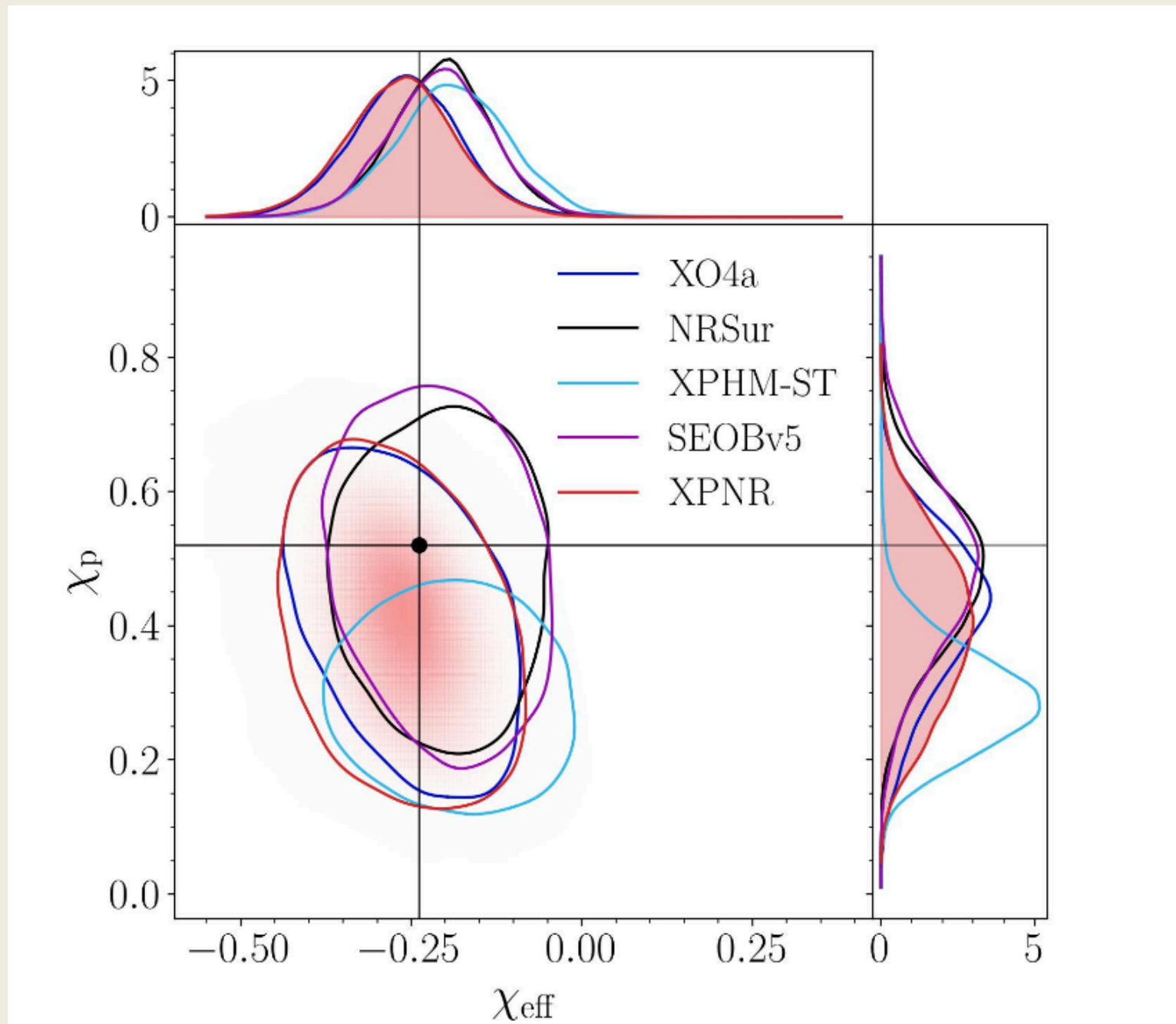


FIG. 11. Comparison between the two-dimensional marginalized posterior distributions for the effective parallel spin  $\chi_{\text{eff}}$  and effective perpendicular spin  $\chi_p$  when analysing the CF\_54 numerical relativity simulation [39]. The contours encase 90% probability. The black crosshairs show the true values.

Hamilton, E., Colleoni, M. et al. arXiv:2507.02604 (2025).

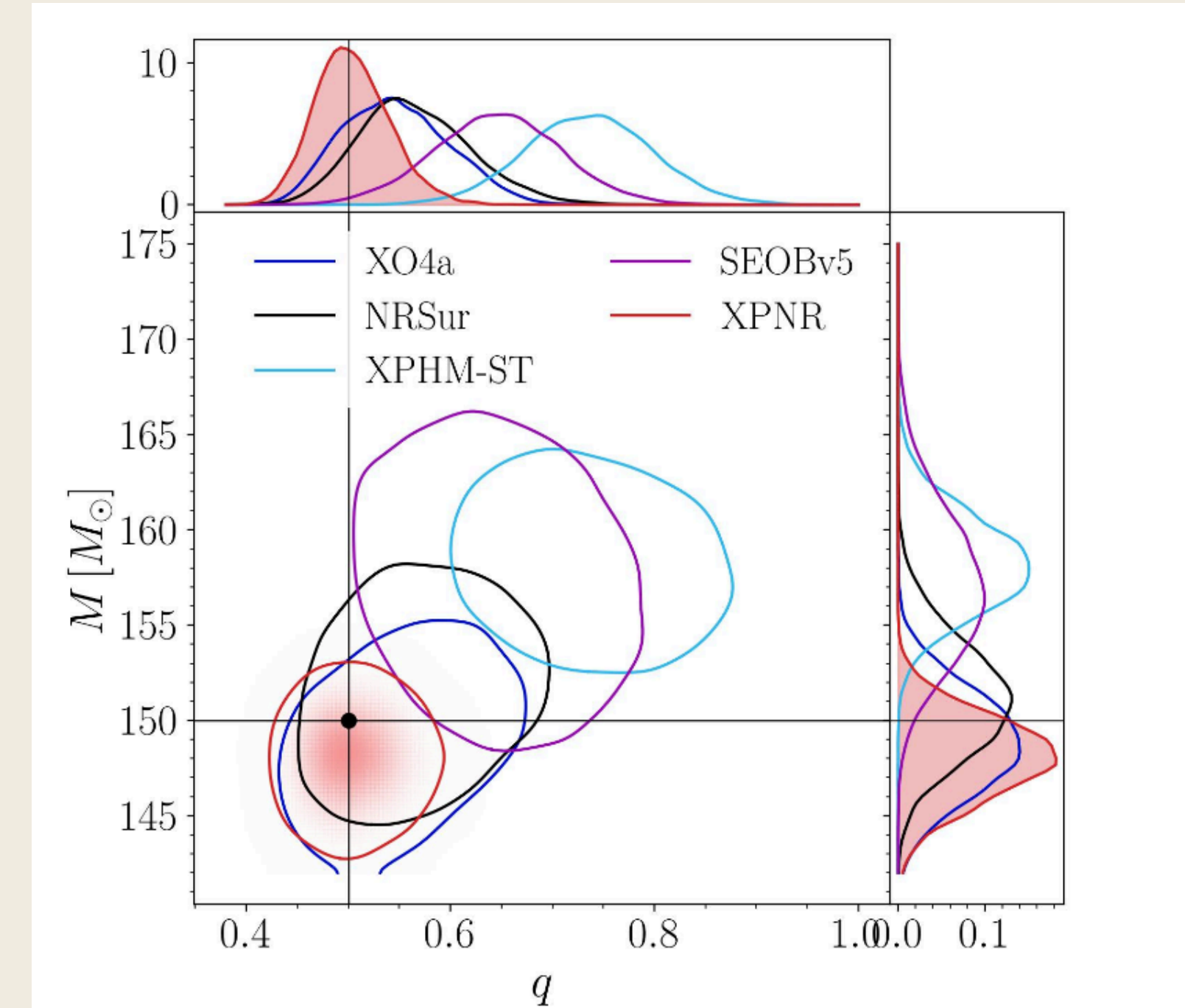


FIG. 12. Comparison between the two-dimensional marginalized posterior distributions for the binary total mass  $M$  and mass ratio  $q$  when analysing an NRSur7DQ4 injection. The contours encase 90% probability. The black cross hairs show the true values.

Hamilton, E., Colleoni, M. et al. arXiv:2507.02604 (2025).

# QC models: timing

- Frequency domain models typically most efficient
- Phenom family more efficient
- Adding in additional physics slows down model evaluation

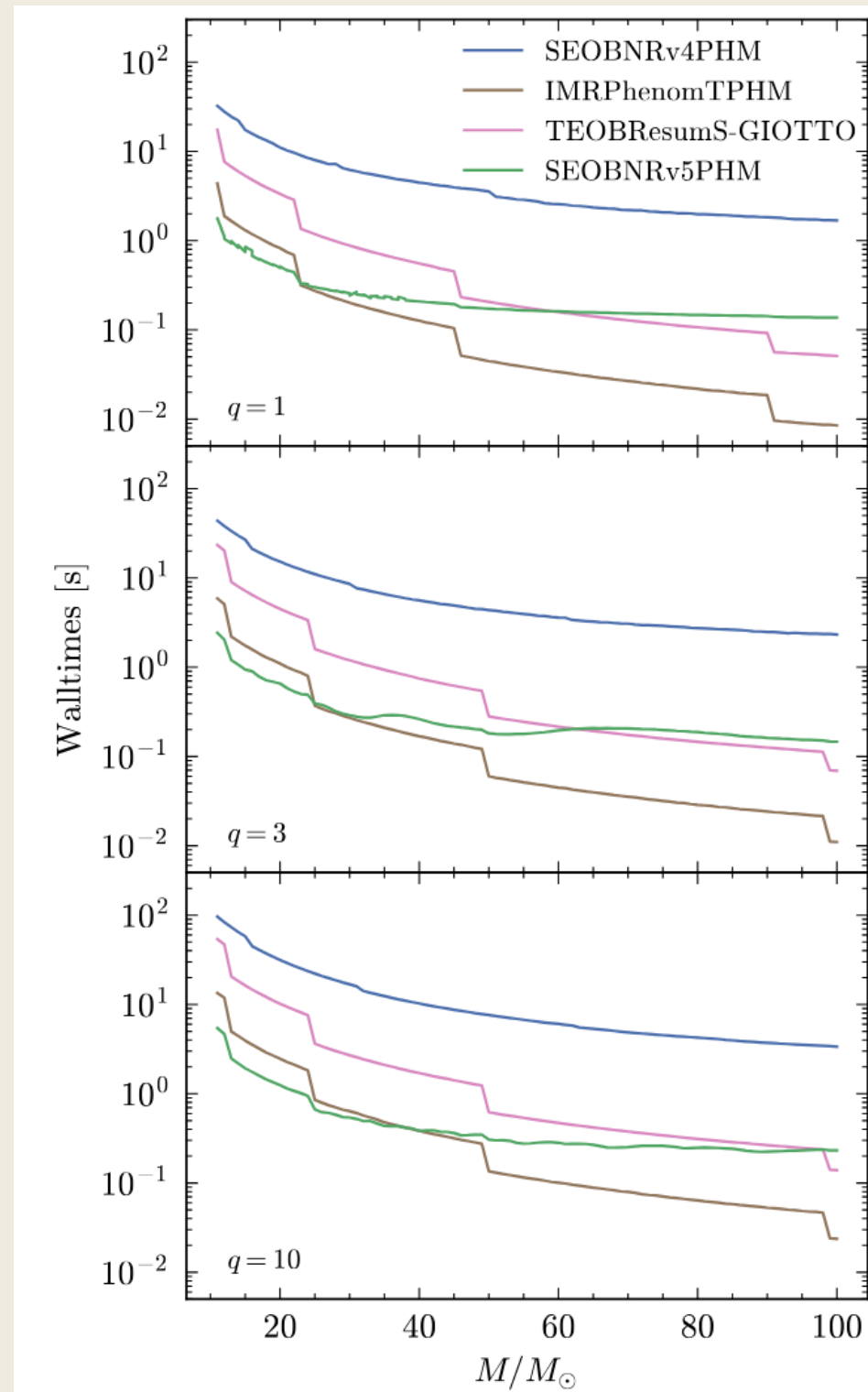


Figure 11. Walltimes of the SEOBv4PHM, IMRPhenomTPHM, TEOBResumS-GIOTTO and SEOBv5PHM models for a configuration with dimensionless spins  $\chi_1 = [0.5, 0, 0.8]$ ,  $\chi_2 = [0, 0.5, 0.3]$ , total mass range  $M \in [10, 100]M_\odot$ , starting frequency  $f_{\text{start}} = 10\text{Hz}$  and three different mass ratios 1 (top panel), 3 (mid panel) and 10 (bottom panel).

Ramos-Buades, A., et al. (2023).  
Physical Review D, 108(12), 124037.

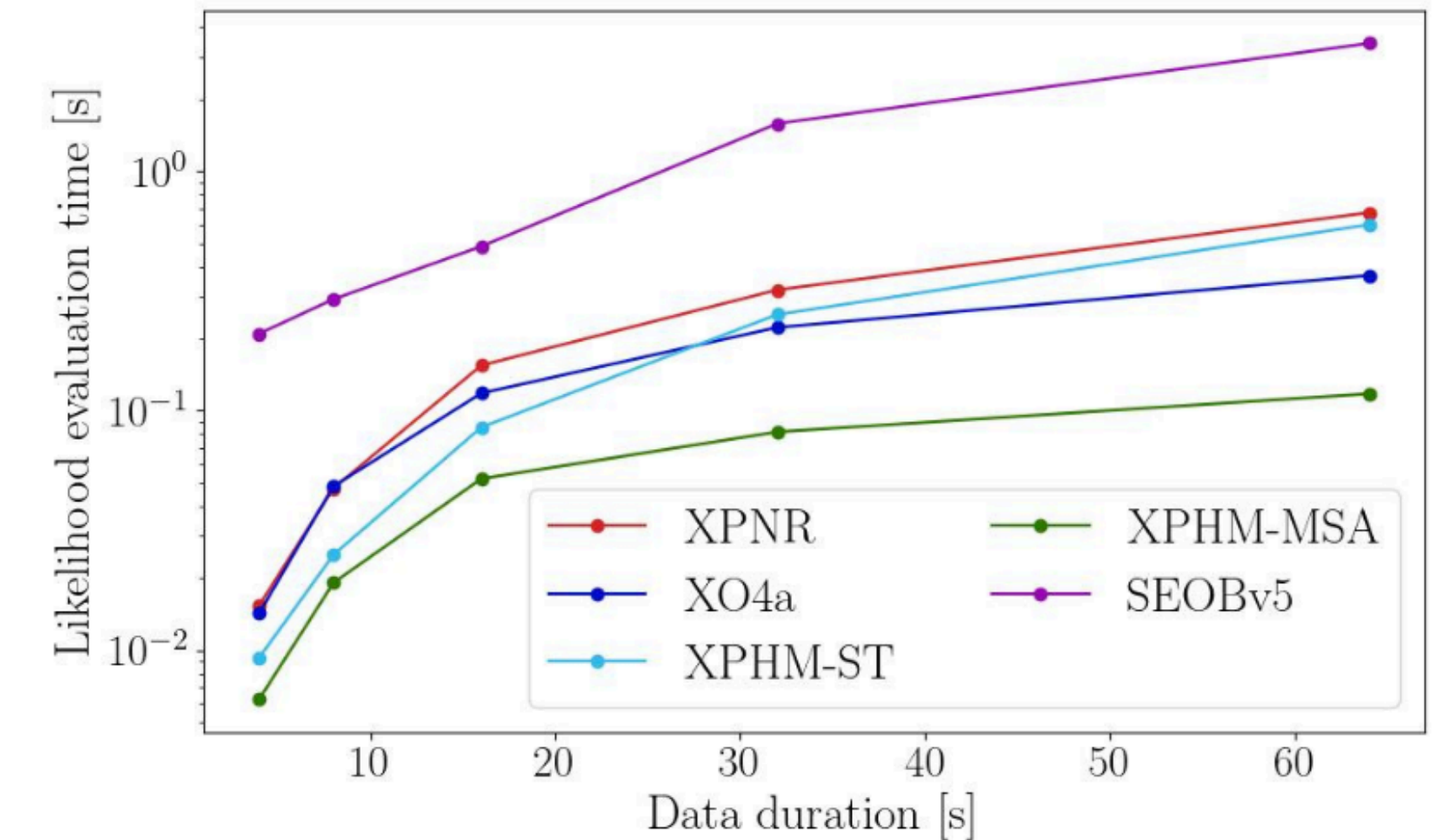


FIG. 16. Comparison between likelihood evaluation times for varying data duration. We compare the evaluation times for four different FD models: PHENOMXPNR, PHENOMXO4A, PHENOMXPHM-SPINTAYLOR, and the TD model: SEOBv5PHM. Each data point represents the mean of 10,000 likelihood evaluations calculated via BILBY. Binary parameters were drawn from priors distributions that cut the chirp mass parameter space to ensure that the gravitational waves fit within the specified data duration: larger data durations correspond to smaller chirp mass binaries.

Hamilton, E., Colleoni, M. et al. arXiv:2507.02604 (2025).



# Eccentric models

	Modeling approach	Complete inspiral	Calibration to eccentric systems	Precessing extension
<b>NRSurE_q4NoSpin_22</b>	Surrogate	×	$q \leq 1; e \leq 0.25$	×
<b>NRSur2dq1Ecc</b>	Surrogate	×	$q = 1; e \leq 0.2$	×
<b>PhenomTEHM</b>	Semi-analytic	✓	×	×
<b>PhenomXE</b>	Semi-analytic	✓	×	×
<b>SEOBNRv5EHM</b>	Semi-analytic	✓	×	×
<b>TEOBResumS-Dali</b>	Semi-analytic	✓	×	×

# Eccentric models: model accuracy

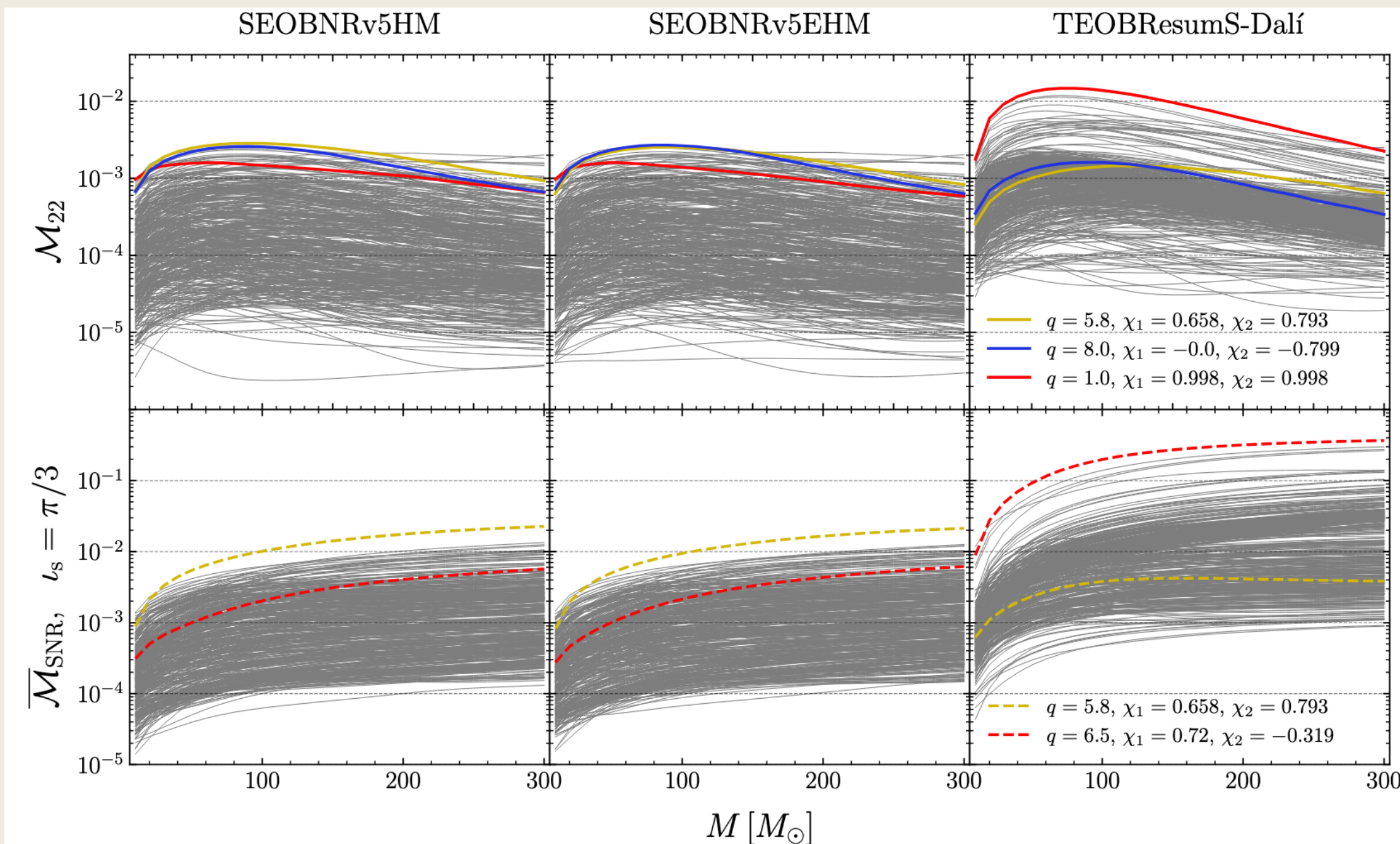


Figure 1. Waveform mismatches between different aligned-spin approximants and the 441 SXS QC NR simulations used in this work, calculated over a range of total masses between 10 and 300  $M_\odot$ . The first column corresponds to the QC SEOBNRv5HM model, the second column to the eccentric SEOBNRv5EHM model, and the third column to the eccentric TEOBResumS-Dalí model. The top panels show the  $(2,2)$ -mode mismatch  $\mathcal{M}_{22}$  for each NR waveform, and the bottom panels show the sky-and-polarization averaged, SNR-weighted mismatch  $\mathcal{M}_{\text{SNR}}$  (which takes into consideration higher-order modes) for inclination  $\iota_s = \pi/3$ . The colored lines highlight cases with the worst maximum mismatch for each model (solid lines for  $\mathcal{M}_{22}$  and dashed lines for  $\mathcal{M}_{\text{SNR}}$ ) with the legends indicating the parameters of the worst-case mismatch for the corresponding model.

Gamboa, Aldo, et al. "Accurate waveforms for eccentric, aligned-spin binary black holes: The multipolar effective-one-body model SEOBNRv5EHM." *Physical Review D* 112.4 (2025): 044038.

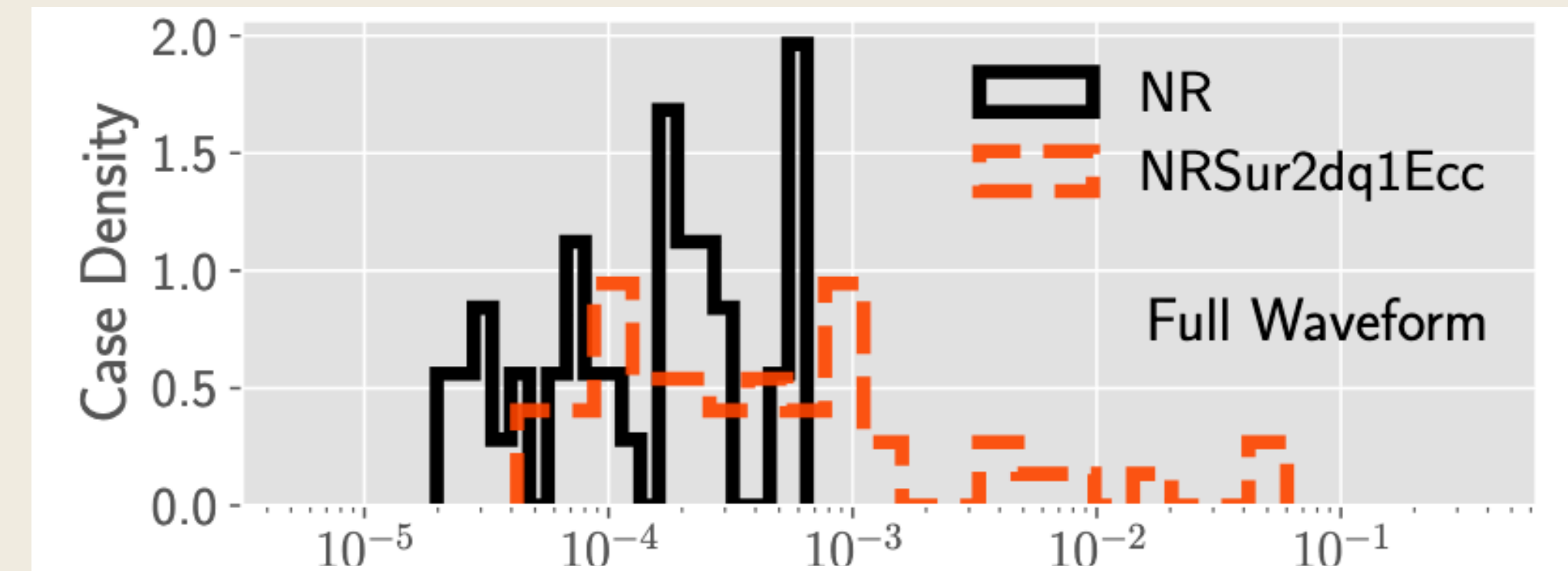
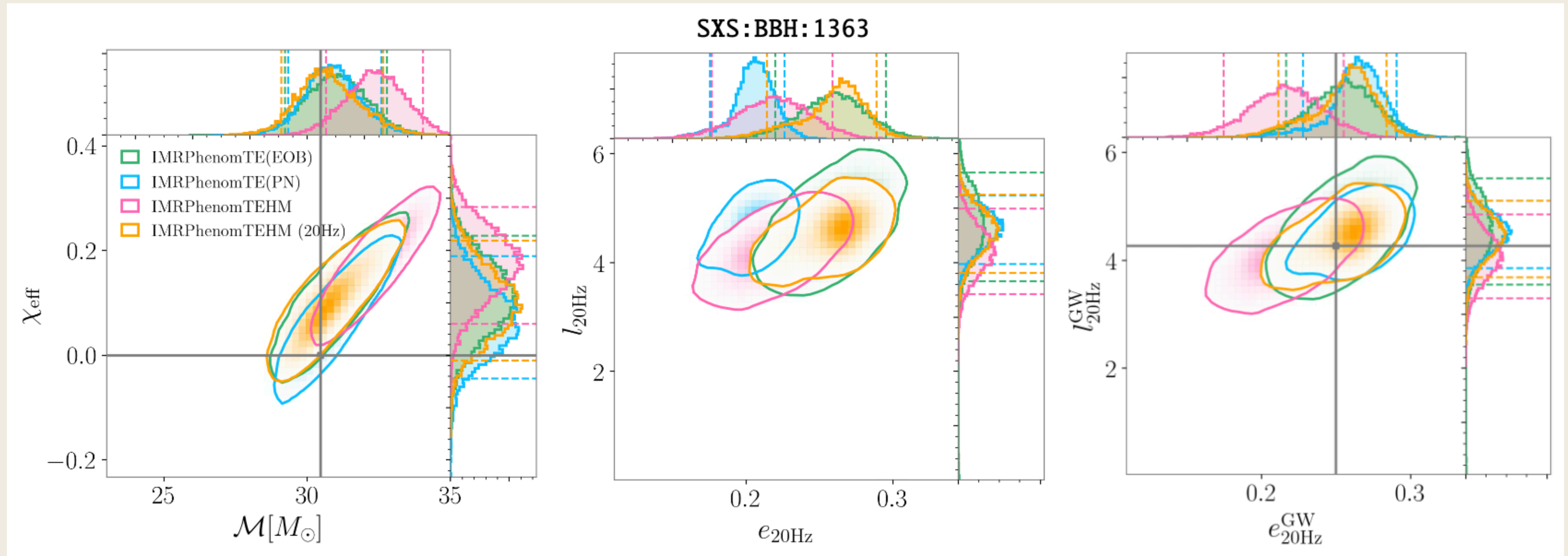


Figure 6. Time-domain leave-one-out errors  $\mathcal{E}$ , defined in Eq. (17), for the full waveform as well as the individual modes considered in the model. For comparison, we also show the NR error between the two highest resolutions. The largest errors are found near the parameter domain's boundary where the trial surrogate, built as part of the cross-validation study, is extrapolating.

Islam, Tousif, et al. "Eccentric binary black hole surrogate models for the gravitational waveform and remnant properties: Comparable mass, nonspinning case." *Physical Review D* 103.6 (2021): 064022.



# Eccentric models: parameter recovery



Planas, Maria de Lluc, et al. "Time-domain phenomenological multipolar waveforms for aligned-spin binary black holes in elliptical orbits." arXiv preprint arXiv:2503.13062 (2025).

