# SEARCH FOR Z' RADIATING FROM THE DARK MATTER AT THE LHC

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#### Outline:

- 1. Introduction
- 2. The Model
- 3. Methods for Simulation and Analysis
- 4. Results
- 5. Summary

#### Work with R. Primulando, Q.M.B. Soesanto, B. Dirgantara, B.E. Gunara

The details can be found in arXiv : 2404.12781

#### PROBLEMS IN SM

Incomplete to be the description of the Universe!

Several major problems to answer:

- 1. Dark Matter (26% of mass-energy present in the Universe, SM only accounts for ~5%) and Dark Energy  $(69\%)^1$
- 2. Neutrino masses, observed from experiments on neutrino oscillation
- 3. Matter-antimatter symmetry
- 4. Gravity is not included

LHC is suited to scrutinize the coupling of Dark Matter with quarks.

DM itself is invisible.

- The detector detects the recoils of DM with visible objects : Monojet, monophoton, etc
- ATLAS search for Monojet with axial-vector mediator  $Z<sub>A</sub>$  (Phys. Rev. D 103, 112006)
- Robust constraint but substantial background.
- Dark sector can have dark gauge group, DM radiates dark gauge bosons. Can resemble monojet.
- If dark gauge bosons have interaction with SM gauge bosons, it can decay to SM particles.

#### PROPOSED MODEL

Fermionic DM  $\chi$  and new corresponding gauge group  $U(1)_D$ .

$$
\mathcal{L}_{\text{Dark}} = g_D X_\mu \bar{\chi} \gamma^\mu \chi
$$

Kinetic mixing term of  $U(1)_D$  with SM  $U(1)_Y$ ,

$$
\mathcal{L}_{\text{kin,mixing}} = -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} + \frac{\epsilon}{2\cos\theta_W}B_{\mu\nu}X^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu}
$$

This mixing will yields (after diagonalization) two massive neutral gauge bosons,  $Z$  and  $Z'$ .

We introduce new variables (used also in Phys. Rev. D 106, 015027) :

$$
\kappa = \frac{\epsilon}{\cos \theta_W},
$$

$$
\eta = \frac{\kappa}{\sqrt{1 - \kappa^2}},
$$

We have bounds on  $\epsilon$ , upper limit is in the order of  ${\rm O}(10^{-3})$  (Phys. Rev. Lett. 124, 131802).

Due to the mixing, in case of  $m_{Z}$ ,  $< 2 m_{\gamma}$ ,  $Z'$  will exclusively decay into SM particles. The decay width of  $Z' \rightarrow f^+f^-$  is given by:

$$
\Gamma_{Z' \to f^+ f^-} = \frac{m_{Z'}}{12\pi} N_c \sqrt{1 - \frac{4m_f^2}{m_{Z'}^2} \left[ \left( c_V^2 + c_A^2 \right) \left( 1 - \frac{m_f^2}{m_{Z'}^2} \right) + 3(c_V^2 - c_A^2) \frac{m_f^2}{m_{Z'}^2} \right]}
$$

We are interested in the case where the final state are dilepton pairs + MET.

The production of DM is through a simplified model, similar to the one used in LHC monojet analysis (Phys. Rev. D 103, 112006).

The Lagrangian is given by:

$$
\mathcal{L}_{q-\chi} = g_Q(\bar{u}\gamma^\mu \gamma^5 u + \bar{d}\gamma^\mu \gamma^5 d) Z_A^{\ \mu} + g_\chi \bar{\chi} \gamma^\mu \gamma^5 \chi Z_A^{\ \mu}
$$

The mediator  $Z_A$  will connect the quark and DM sectors.

Our focus is directed towards a process in which DM radiates  $Z'$ , succeeded by the subsequent decay of  $Z'$ into lepton pairs. This process is termed as "Darkstrahlung" (JHEP 1509 (2015) 079).



Fig 1. The Darkstrahlung Process

### SEARCH FOR DILEPTON + MET FINAL STATE AT THE LHC

- ATLAS Search for Dilepton Resonance ([https://cds.cern.ch/record/2870113\)](https://cds.cern.ch/record/2870113)
- ATLAS Slepton Search (*Eur. Phys. J. C* 80, 123, 2020)
- CMS Slepton Search (*J. High Energ. Phys.* 04, 2021, 123)

Important selection criteria on these experiments:



#### TOOLS USED IN THIS WORK



## LIKELIHOOD AND CHI-SQUARE VALUES

In order to calculate the bounds for our model, we follow the strategy of GAMBIT (*Eur. Phys. J. C* 77, 795, 2017).

The likelihood of observing a signal in particular bin  $i$ ,  $\overline{s_i}$  is

$$
\mathcal{L}_i(n_i|s_i, b_i) = \int_0^\infty \frac{(\xi(s_i + b_i))^{n_i} e^{-\xi(s_i + b_i)}}{n_i!} P_i(\xi) d\xi,
$$

With  $n$ ,  $b$  are number of observed and predicted background, respectively. The probability is defined as :

$$
P_i(\xi) = \frac{1}{\sqrt{2\pi}\sigma_i} \frac{1}{\xi} \exp\left[-\frac{1}{2}\left(\frac{\ln\xi}{\sigma_i}\right)\right]
$$

and finally the chi-square values are given by :

$$
\chi^2 = -2 \sum_i (\ln \mathcal{L}(n_i | s_i, b_i) - \ln \mathcal{L}(n_i | s_i = 0, b_i))
$$

#### RESULTS

Recasted bounds from ATLAS dilepton resonance, ATLAS Slepton, and CMS Slepton is shown below. We chose a benchmark of  $m_{Z_A} = 2$  TeV,  $m_\chi = 120$  GeV,  $g_Q = 0.25$  and  $g_\chi = 1$ , corresponding to ATLAS monojet analysis.



 $m_{Z'}$  (GeV)

Logically, ATLAS Dilepton Resonance search should gives stronger constraints! Not the case, because of low  $E^T_{miss}$ 

We proposed that ATLAS Dilepton Resonance could be improved by implementing a stronger cut on missing energy.

This is the simulation result with additional cuts : $\bigl| |m_{ll} - m_{Z\prime} | < 5$  GeV and  $E^{T}_{miss} > 250$  GeV.



This optimization also gives stronger bounds on  $g_D$  for various values of  $m_Z$ , and  $m_Y$ .





(b) Optimized bounds

### SUMMARY

- We discussed LHC constraints for our model, where the dark sector consists of a dark matter charged under new hidden  $U(1)$  symmetry.
- Recasting of LHC searches (ATLAS Dilepton Resonance + ATLAS Slepton + CMS Slepton) has been performed.
- Higher cut on  $E^T_{miss}$  and choosing narrow region for  $m_{ll}$  can improve the constraints.



#### **BACKUP**









CMS Slepton **CMS** Slepton **ATLAS Monojet Search** 







