

ICISE

The 30th Vietnam School of Physics: Particle and Dark Matter

Department of Theoretical Physics, University of Science
Vietnam National University, HCM city

Student Seminar

Theoretical studies the $e^+e^- \rightarrow \nu\bar{\nu}H$ process in the ILC

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1. Introduction

The field of Physics has witnessed remarkable advancements in recent times, with the emergence of numerous novel theories and models.

The Standard Model, a mathematical framework that describes the fundamental interactions of nature, has garnered significant attention and interest.

The Standard Model has been extensively validated through numerous experiments, particularly at high-energy particle accelerators such as ILC, LHC, RHIC, ATLAS,...

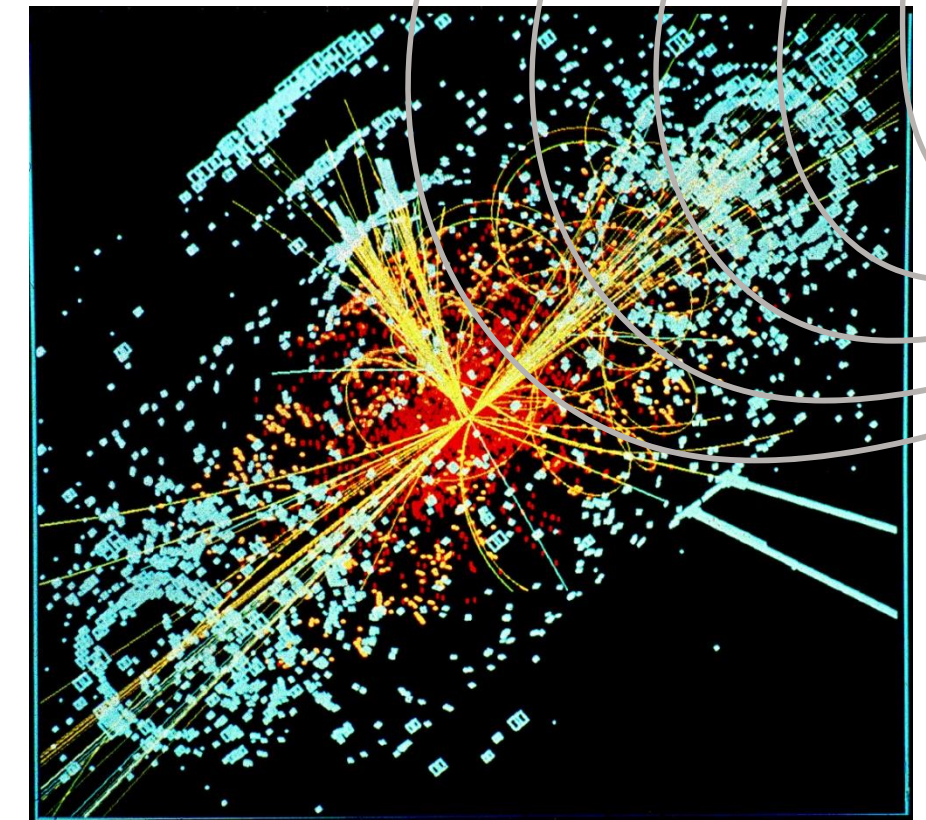


Figure 1. Simulating the Higgs boson decay event after the collision of protons occurring at the CERN.

2. Objective

Starting from the Standard Model, this seminar will focus on studying the process $e^+e^- \rightarrow \nu\bar{\nu}H$.

The project will be carried out:

Calculate the scattering amplitude $|\mathcal{M}|^2$.

Construct the phase space for the scattering process.

Perform numerical calculations of scattering cross sections, obtain results and discuss physical distributions.



Contributes to providing more useful information to verify the Standard Model

3. Theoreticle Framework

3.1. Overview of the Standard Model

The Standard Model is a Physics model built based on Quantum Field Theory through the group symmetry property $SU_c(3) \times SU_L(2) \times U_Y(1)$, this model is to describe the physical particles and interactions between them.

Table 1: Elementary particles in the Standard Model

	I	II	III
Lepton	e	μ	τ
	ν_e	ν_μ	ν_τ
Quark	u	c	t
	d	s	b

Gauge boson	Higgs boson
W^\pm	H
Z	
γ	
g	

The Standard Model of Particle Physics

	Fermions			Bosons	
Quarks	 up	 charm	 top	 photon	Force Carriers
	 down	 strange	 bottom	 Z boson	
	Leptons	 electron neutrino	 muon neutrino	 tau neutrino	
 electron		 muon	 tau	 gluon	
 Higgs boson					



3.2. Lagrangian for the Standard Model

The Lagrangian of the Standard Model is given by

$$\mathcal{L}_M = \mathcal{L}_{Kin} + \mathcal{L}_{YM} + \mathcal{L}_{Higgs} + \mathcal{L}_{Yukawa}$$

The Lagrangian of the Standard Model consist of components

$$\mathcal{L}_{Kin} = i\bar{Q}_L^i \gamma^\mu D_\mu Q_L^i + i\bar{u}_R^i \gamma^\mu D_\mu u_\mu^i + i\bar{d}_R^i \gamma^\mu D_\mu d_R^i + i\bar{L}^i \gamma^\mu D_\mu L^i + i\bar{e}_R^i \gamma^\mu D_\mu e_R^i$$

$$\mathcal{L}_{YM} = -\frac{1}{4}W^{\mu\nu}W_{\mu\nu} - \frac{1}{4}B^{\mu\nu}B_{\mu\nu} - \frac{1}{4}G^{\mu\nu}G_{\mu\nu}$$

$$\mathcal{L}_{Higgs} = (D_\mu \phi)^\dagger (D_\mu \phi) - V(\phi)$$

$$\mathcal{L}_{Yukawa} = -h_e(\bar{L}^i \phi e_R^i + h.c) - h_d(\bar{Q}^i \phi d_R^i + h.c) - h_u(\bar{Q}^i \phi_c u_R^i + h.c)$$

Standard Model of Elementary Particles

	three generations of matter (fermions)			interactions / force carriers (bosons)	
	I	II	III		
mass	=2.2 MeV/c ²	=1.28 GeV/c ²	=173.1 GeV/c ²	0	=124.97 GeV/c ²
charge	2/3	2/3	2/3	0	0
spin	1/2	1/2	1/2	1	0
QUARKS	u up	c charm	t top	g gluon	H higgs
	d down	s strange	b bottom	γ photon	
	e electron	μ muon	τ tau	Z Z boson	
LEPTONS	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

$$\begin{aligned}
 \mathcal{L}_{SM} = & -\frac{1}{2}\partial_\mu g_\nu^\alpha \partial_\mu g_\nu^\alpha - g_s f^{abc} \partial_\mu g_\nu^a g_\nu^b g_\nu^c - \frac{1}{2}g_s^2 f^{abc} f^{ade} g_\nu^b g_\nu^c g_\nu^d g_\nu^e - \partial_\mu W_\nu^+ \partial_\mu W_\nu^- - \\
 & M^2 W_\nu^+ W_\nu^- - \frac{1}{2}\partial_\mu Z_\nu^\alpha \partial_\mu Z_\nu^\alpha - \frac{1}{2}M^2 Z_\nu^\alpha Z_\nu^\alpha - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - ig_{cw} (\partial_\mu Z_\nu^\alpha (W_\nu^+ W_\nu^- - \\
 & W_\nu^- W_\nu^+) - Z_\nu^\alpha (W_\nu^+ \partial_\mu W_\nu^- - W_\nu^- \partial_\mu W_\nu^+) + Z_\nu^\alpha (W_\nu^+ \partial_\mu W_\nu^- - W_\nu^- \partial_\mu W_\nu^+)) - \\
 & ig_{sw} (\partial_\mu A_\nu (W_\nu^+ W_\nu^- - W_\nu^- W_\nu^+) - A_\nu (W_\nu^+ \partial_\mu W_\nu^- - W_\nu^- \partial_\mu W_\nu^+) + A_\nu (W_\nu^+ \partial_\mu W_\nu^- - \\
 & W_\nu^- \partial_\mu W_\nu^+)) - \frac{1}{2}g^2 W_\nu^+ W_\nu^- W_\nu^+ W_\nu^- + \frac{1}{2}g^2 W_\nu^+ W_\nu^- W_\nu^+ W_\nu^- + g^2 c_w^2 (Z_\nu^\alpha W_\nu^+ Z_\nu^\alpha W_\nu^- - \\
 & Z_\nu^\alpha Z_\nu^\alpha W_\nu^+ W_\nu^-) + g^2 s_w^2 (A_\nu W_\nu^+ A_\nu W_\nu^- - A_\nu A_\nu W_\nu^+ W_\nu^-) + g^2 s_w c_w (A_\nu Z_\nu^\alpha (W_\nu^+ W_\nu^- - \\
 & W_\nu^- W_\nu^+) - 2A_\nu Z_\nu^\alpha W_\nu^+ W_\nu^-) - \frac{1}{2}\partial_\mu H \partial_\mu H - 2M^2 \alpha_h H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \\
 & \beta_h \left(\frac{2M^2}{\Lambda^2} H + \frac{2M}{\Lambda} H + \frac{1}{\Lambda} (H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) \right) + \frac{2M^4}{\Lambda^4} \alpha_h - \\
 & g_{\alpha h} M (H^3 + H \phi^0 \phi^0 + 2H \phi^+ \phi^-) - \\
 & \frac{1}{2}g^2 \alpha_h (H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2) - \\
 & g M W_\nu^+ W_\nu^- H - \frac{1}{2}g \frac{M^2}{\Lambda^2} Z_\nu^\alpha Z_\nu^\alpha H - \\
 & \frac{1}{2}ig (W_\nu^+ (\partial^\mu \partial_\mu \phi^- - \phi^- \partial_\mu \partial^\mu) - W_\nu^- (\partial^\mu \partial_\mu \phi^+ - \phi^+ \partial_\mu \partial^\mu)) + \\
 & \frac{1}{2}g (W_\nu^+ (H \partial_\mu \phi^- - \phi^- \partial_\mu H) + W_\nu^- (H \partial_\mu \phi^+ - \phi^+ \partial_\mu H)) + \frac{1}{2}g \frac{1}{c_w} (Z_\nu^\alpha (H \partial_\mu \phi^0 - \phi^0 \partial_\mu H) + \\
 & M (\frac{1}{c_w} Z_\nu^\alpha \partial_\mu \phi^0 + W_\nu^- \partial_\mu \phi^- + W_\nu^+ \partial_\mu \phi^+) - ig \frac{M}{c_w} Z_\nu^\alpha (W_\nu^- \phi^- - W_\nu^+ \phi^+) + ig s_w M A_\nu (W_\nu^+ \phi^- - \\
 & W_\nu^- \phi^+) - ig \frac{1-2c_w^2}{2c_w} Z_\nu^\alpha (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) + ig s_w A_\nu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \\
 & \frac{1}{2}g^2 W_\nu^+ W_\nu^- (H^2 + (\phi^0)^2 + 2\phi^+ \phi^-) - \frac{1}{2}g^2 \frac{1}{c_w} Z_\nu^\alpha Z_\nu^\alpha (H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-) - \\
 & \frac{1}{2}g^2 \frac{1}{c_w} Z_\nu^\alpha Z_\nu^\alpha (W_\nu^+ \phi^- + W_\nu^- \phi^+) - \frac{1}{2}ig^2 \frac{1}{c_w} Z_\nu^\alpha Z_\nu^\alpha H (W_\nu^+ \phi^- - W_\nu^- \phi^+) + \frac{1}{2}g^2 s_w A_\nu (W_\nu^+ \phi^- + \\
 & W_\nu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\nu H (W_\nu^+ \phi^- - W_\nu^- \phi^+) - g^2 \frac{2c_w^2 - 1}{2c_w} Z_\nu^\alpha A_\nu \phi^+ \phi^- - \\
 & g^2 s_w^2 A_\nu A_\nu \phi^+ \phi^- + \frac{1}{2}ig_s \lambda_1 (\bar{q}^i \gamma^\mu q^j) g_\mu^i - e^i (\gamma^\mu + m_i^2) e^i - \bar{\nu}^i (\gamma^\mu + m_i^2) \nu^i - \bar{u}^i (\gamma^\mu + \\
 & m_i^2) u^i - \bar{d}^i (\gamma^\mu + m_i^2) d^i + ig s_w A_\nu (-e^i \gamma^\mu e^i + \frac{2}{3}(\bar{u}^i \gamma^\mu u^i) - \frac{1}{3}(\bar{d}^i \gamma^\mu d^i)) + \\
 & \frac{ig}{c_w} Z_\nu^\alpha ((\bar{\nu}^i \gamma^\mu (1 + \gamma^5) \nu^i) + (e^i \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^i) + (\bar{d}^i \gamma^\mu (\frac{2}{3}s_w^2 - 1 - \gamma^5) d^i) + \\
 & (\bar{u}^i \gamma^\mu (1 - \frac{2}{3}s_w^2 + \gamma^5) u^i)) + \frac{ig}{2c_w} W_\nu^+ ((\bar{\nu}^i \gamma^\mu (1 + \gamma^5) U^{ij} \nu^j) + (\bar{u}^i \gamma^\mu (1 + \gamma^5) C_{ij} d^j)) + \\
 & \frac{ig}{2c_w} W_\nu^- ((\bar{e}^i U^{ij} \nu^j (1 + \gamma^5) \nu^i) + (\bar{d}^i C_{ij} u^j (1 + \gamma^5) u^i)) + \\
 & \frac{ig}{2M^2} \phi^+ (-m_e^2 (\bar{\nu}^i U^{ij} \nu^j (1 - \gamma^5) e^i) + m_\mu^2 (\bar{\nu}^i U^{ij} \nu^j (1 + \gamma^5) e^i) + \\
 & \frac{ig}{2M^2} \phi^- (m_e^2 (\bar{e}^i U^{ij} \nu^j (1 + \gamma^5) \nu^i) - m_\mu^2 (\bar{e}^i U^{ij} \nu^j (1 - \gamma^5) \nu^i) - \frac{2}{3} \frac{M^2}{\Lambda^2} H (\bar{\nu}^i \nu^i) - \\
 & \frac{2}{3} \frac{M^2}{\Lambda^2} H (\bar{e}^i e^i) + \frac{ig}{2} \frac{M^2}{\Lambda^2} \phi^0 (\bar{\nu}^i \gamma^5 \nu^i) - \frac{ig}{2} \frac{M^2}{\Lambda^2} \phi^0 (\bar{e}^i \gamma^5 e^i) - \frac{1}{2} e_h M_h^2 (1 - \gamma_5) \bar{\nu}_e - \\
 & \frac{1}{2} \bar{\nu}_h M_h^2 (1 - \gamma_5) \nu_h + \frac{ig}{2M^2} \phi^+ (-m_d^2 (\bar{u}^i C_{ij} (1 - \gamma^5) d^j) + m_u^2 (\bar{u}^i C_{ij} (1 + \gamma^5) d^j)) + \\
 & \frac{ig}{2M^2} \phi^- (m_d^2 (\bar{d}^i C_{ij} (1 + \gamma^5) u^j) - m_u^2 (\bar{d}^i C_{ij} (1 - \gamma^5) u^j) - \frac{2}{3} \frac{M^2}{\Lambda^2} H (\bar{u}^i u^i) - \\
 & \frac{2}{3} \frac{M^2}{\Lambda^2} H (\bar{d}^i d^i)) + \frac{ig}{2} \frac{M^2}{\Lambda^2} \phi^0 (\bar{u}^i \gamma^5 u^i) - \frac{ig}{2} \frac{M^2}{\Lambda^2} \phi^0 (\bar{d}^i \gamma^5 d^i) + G^a \partial^\mu G^a + g_s f^{abc} \partial_\mu G^a G^b G^c + \\
 & \bar{X}^+ (\partial^\mu - M^2) X^+ + \bar{X}^- (\partial^\mu - M^2) X^- + \bar{X}^0 (\partial^\mu - \frac{M^2}{2}) X^0 + \bar{Y} \partial^\mu Y + ig_{cw} W_\nu^+ (\partial_\mu \bar{X}^0 X^- - \\
 & \partial_\mu \bar{X}^+ X^0) + ig_{sw} W_\nu^+ (\partial_\mu \bar{Y} X^- - \partial_\mu \bar{X}^+ Y) + ig_{cw} W_\nu^- (\partial_\mu \bar{X}^- X^0 - \\
 & \partial_\mu \bar{X}^0 X^+) + ig_{sw} W_\nu^- (\partial_\mu \bar{X}^- Y - \partial_\mu \bar{Y} X^+) + ig_{cw} Z_\nu^\alpha (\partial_\mu \bar{X}^+ X^- - \\
 & \partial_\mu \bar{X}^- X^+) + ig_{sw} A_\nu (\partial_\mu \bar{X}^+ X^- - \\
 & \partial_\mu \bar{X}^- X^+) - \frac{1}{2}gM (\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{2} \bar{X}^0 X^0 H) + \frac{1-2c_w^2}{2c_w} igM (\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-) + \\
 & \frac{1}{2c_w} igM (\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-) + igM s_w (\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-) + \\
 & \frac{1}{2}igM (\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0) .
 \end{aligned}$$



4. Methodology

In this project, we will use the Mathematica with FeynArts and FormCalc package to calculation program and numerical solution using Monte Carlo to perform calculations for the scattering process $e^+e^- \rightarrow \nu\bar{\nu}H$.

Mathematica

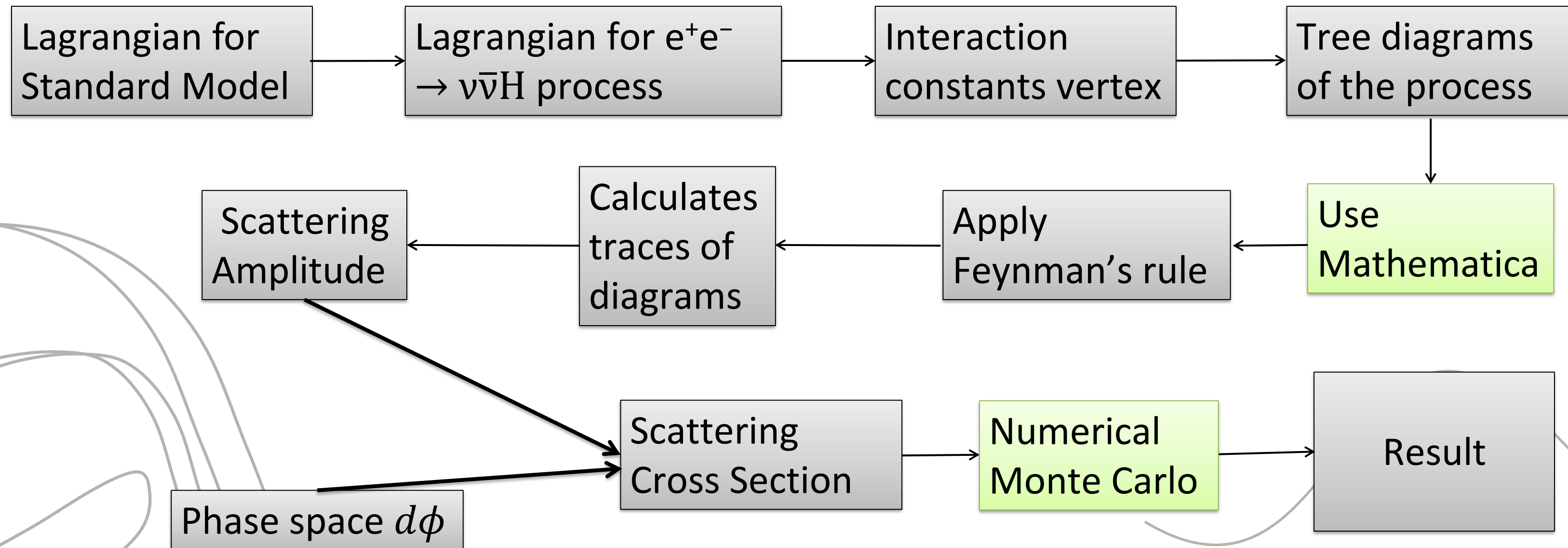
Tool to calculate complex expressions, used to calculate diagrams of the scattering process

Monte Carlo

Random number seeding method plays an important role in numerical solution for multi-dimensional integral

5. Implementation

Going from the Lagrangian for the Standard Model, use Mathematica to calculate the process diagrams to find the scattering amplitude $|\mathcal{M}|^2$ of the $e^+e^- \rightarrow \nu\bar{\nu}H$ process. Then build the phase space to give the scattering cross section. Finally, proceed to numerically solve the process using Monte Carlo.



6. Calculate

The scattering cross sector of process $e^+e^- \rightarrow \nu\bar{\nu}H$ has form

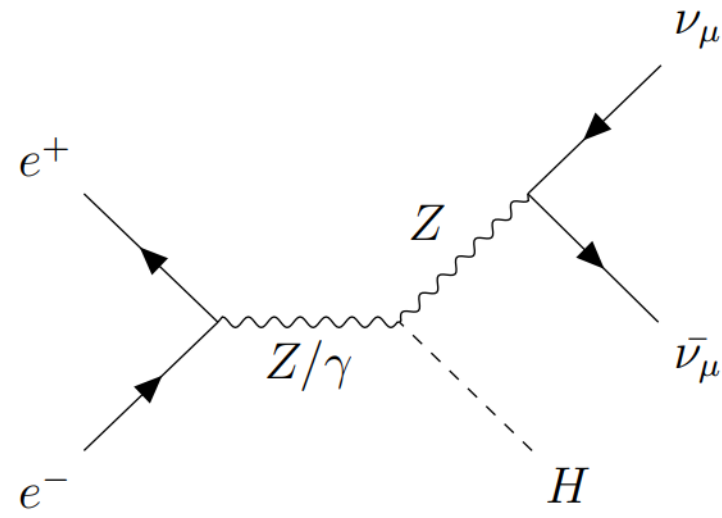
$$\sigma = \int d\phi_3 |\mathcal{M}|^2$$

with

$$d\Phi_3 = J dx_1 dx_2 dx_3 dx_4 dx_5$$

Perform calculations for the process $e^+e^- \rightarrow \nu\bar{\nu}H$ the input parameters:
 $m_W = 80.3767$; $m_Z = 91.1876$; $m_e = m_\nu = 0$; $\Gamma_Z = 2.4952$; $\sqrt{s} = 500$. With the
integration area $\Omega = [0,1]^5$ for the independent variables x_1, x_2, x_3, x_4, x_5 .

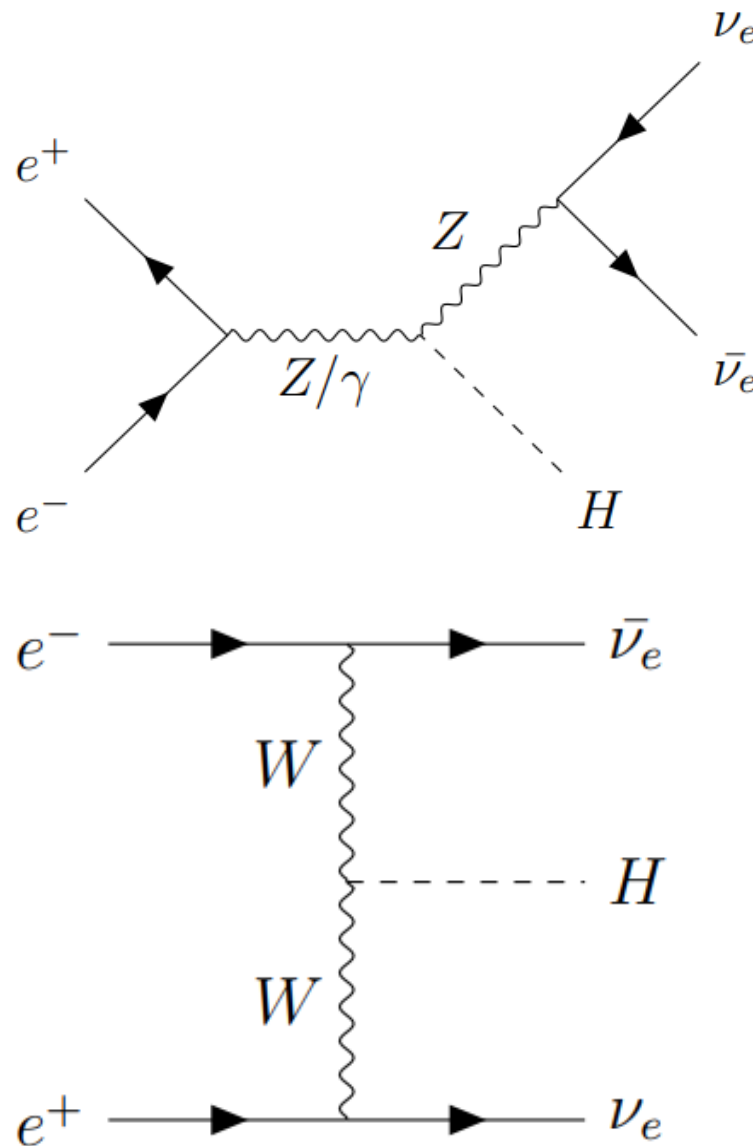
6.1. Calculate the scattering amplitude $|\mathcal{M}|^2$



We have the relation for the diagram $e^+e^- \rightarrow \nu_\mu \bar{\nu}_\mu H$

$$|\mathcal{M}_1|^2 = \frac{16\alpha^3 M_W^2 \pi^3 \left(4p_1 p_3 \times p_2 p_4 \times \sin^4 \theta_W + p_1 p_4 \times p_2 p_3 [1 - 2\sin^2 \theta_W]^2 \right)}{\cos^8 \theta_W \left((M_Z^2 - 2p_3 p_4)^2 + \Gamma_Z^2 M_Z^2 \right) (M_Z^2 - s)^2 \sin^6 \theta_W}$$

We have the relation for the diagram $e^+e^- \rightarrow \nu_e \bar{\nu}_e H$



$$|\mathcal{M}_2|^2 = 16\alpha^3 M_W^2 \pi^3 \left(p_1 p_4 \times p_2 p_3 (M_W^4 + 4p_1 p_3 \times p_2 p_4 + 2M_W^2 (p_1 p_3 + p_2 p_4) - 2\cos^4 \theta_W \right. \\ \times (M_Z^2 - 2p_3 p_4)(M_Z^2 - s))^2 - 4(M_W^2 + 2p_1 p_3) p_1 p_4 \times p_2 p_3 (M_W^2 + 2p_2 p_4) (M_W^4 + 4p_1 p_3 \\ \times p_2 p_4 + 2M_W^2 (p_1 p_3 + p_2 p_4) - 2\cos^4 \theta_W (M_Z^2 - 2p_3 p_4)(M_Z^2 - s)) \sin^2 \theta_W + 4(M_W^2 + 2p_1 p_3)^2 \\ \left. \times (M_W^2 + 2p_2 p_4)^2 (p_1 p_4 \times p_2 p_3 + p_1 p_3 \times p_2 p_4) \sin^4 \theta_W \right) / \left(\cos^8 \theta_W (M_W^2 + 2p_1 p_3)^2 \right. \\ \left. \times (M_W^2 + 2p_2 p_4)^2 \left[(M_Z^2 - 2p_3 p_4)^2 + \Gamma_Z^2 M_Z^2 \right] (M_Z^2 - s)^2 \sin^6 \theta_W \right)$$

So we have amplitude

$$|\mathcal{M}|^2 = 2|\mathcal{M}_1|^2 + |\mathcal{M}_2|^2$$

6.2. Phase space for the scattering process

Phase space for $e^+e^- \rightarrow \nu\bar{\nu}H$ process

$$d\Phi_3 = \frac{1}{8^2(2\pi)^5} dm_{34}^2 d\cos\theta_5 d\phi_5 d\cos\theta_{34} d\phi_{34} \frac{\sqrt{\lambda(s, m_{34}^2, m_5^2)}}{s} \frac{\sqrt{\lambda(m_{34}^2, m_3^2, m_4^2)}}{m_{34}^2}$$

Transform expressions by selecting distributions

$$p(m_{34}^2) = \begin{cases} \frac{1}{q_{max} - q_{min}}, & q_{min} \leq m_{34}^2 \leq q_{max} \\ 0, & \neq \end{cases}$$

$$p(\cos\theta_5) = \begin{cases} \frac{1}{2}, & \cos\theta_5 \in [-1, 1] \\ 0, & \neq \end{cases}$$

$$p(\cos\theta_{34}) = \begin{cases} \frac{1}{2}, & \cos\theta_{34} \in [-1, 1] \\ 0, & \neq \end{cases}$$

with $(m_3 + m_4)^2 \leq m_{34}^2 \leq (\sqrt{s} - m_5)^2$

$$p(\phi_5) = \begin{cases} \frac{1}{2\pi}, & \phi_5 \in [0, 2\pi] \\ 0, & \neq \end{cases}$$

$$p(\phi_{34}) = \begin{cases} \frac{1}{2\pi}, & \phi_{34} \in [0, 2\pi] \\ 0, & \neq \end{cases}$$

Let J be the Jacobian quantity, now we have the Jacobian corresponding to the 2-particles phase space of the process $p_1 + p_2 \rightarrow p_{34} + p_5$

$$J_1 = \frac{1}{4\pi} \frac{\sqrt{\lambda(s, m_{34}^2, m_5^2)}}{2s}$$

We have the Jacobian corresponding to the phase space of two particles in the process $p_{34} = p_3 + p_4$

$$J_2 = \frac{1}{4\pi} \frac{\sqrt{\lambda(m_{34}^2, m_3^2, m_4^2)}}{2m_{34}^2}$$

We get the 3-particle phase space of the scattering process rewritten in terms of the 5 independent variables with a new integration region $\Omega = [0,1]^5$

$$d\Phi_3 = J dx_1 dx_2 dx_3 dx_4 dx_5 \quad \text{with} \quad J = \frac{q_{max} - q_{min}}{2\pi} J_1 J_2$$

7. Result

7.1. Scattering Cross Section

m_H [GeV]	σ_{tree} [fb]	σ [fb]	δ [%]
150	61.12	61.08	-0.06
200	37.33	37.44	0.29
250	21.17	21.13	-0.18
300	10.76	10.83	0.65
350	4.603	4.605	0.05

<https://arxiv.org/abs/hep-ph/0211268>

Note: Result at 2002

Table 2: Scattering cross section of the $e^+e^- \rightarrow \nu\bar{\nu}H$ process under the dependence of the Higgs mass at $\sqrt{s} = 500$ GeV

7.2. Distribution

Compare the distribution for the $e^+e^- \rightarrow \nu\bar{\nu}H$ process at $m_H = 350$ [GeV], $\sqrt{s} = 500$ GeV with previous research.

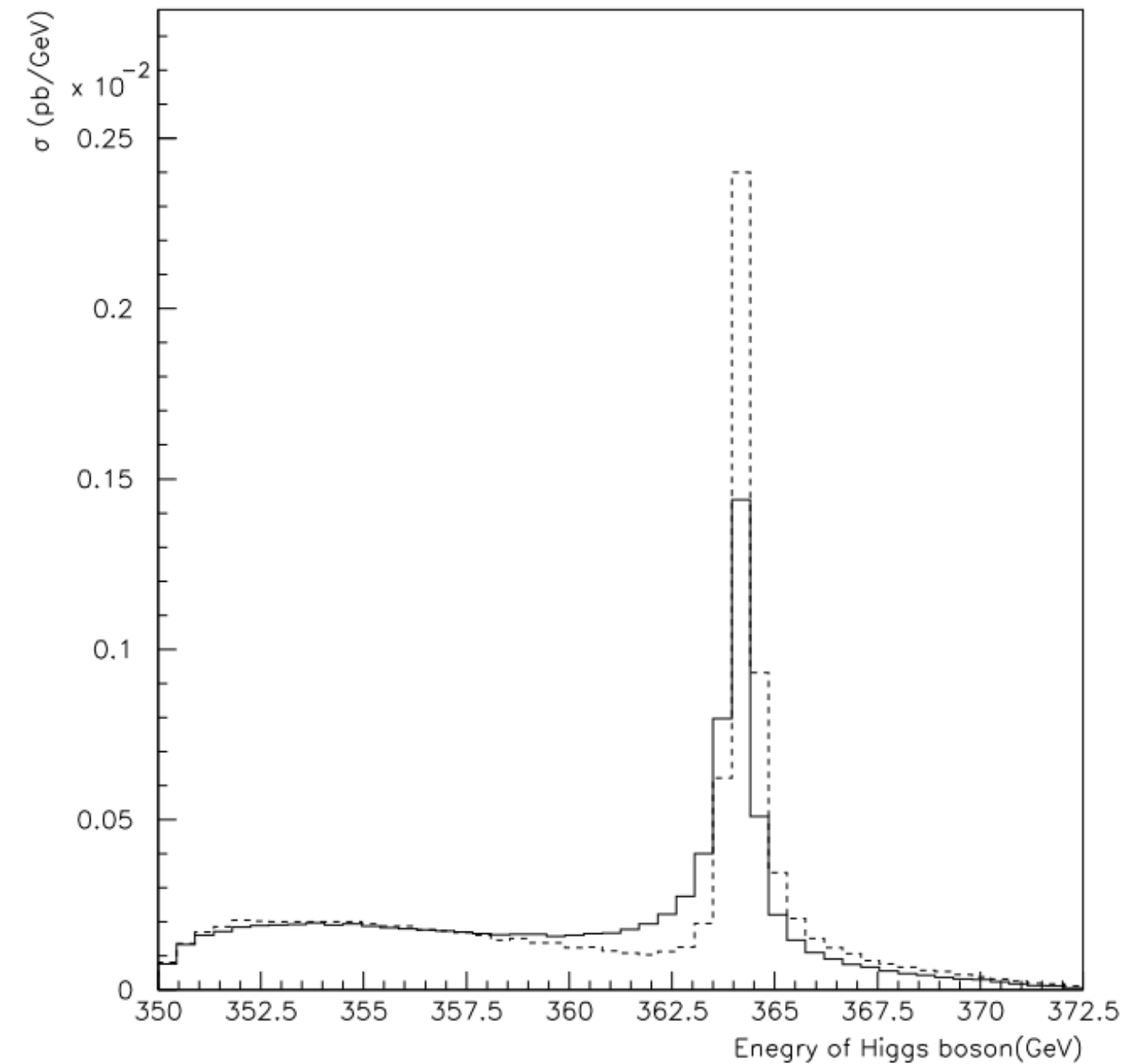
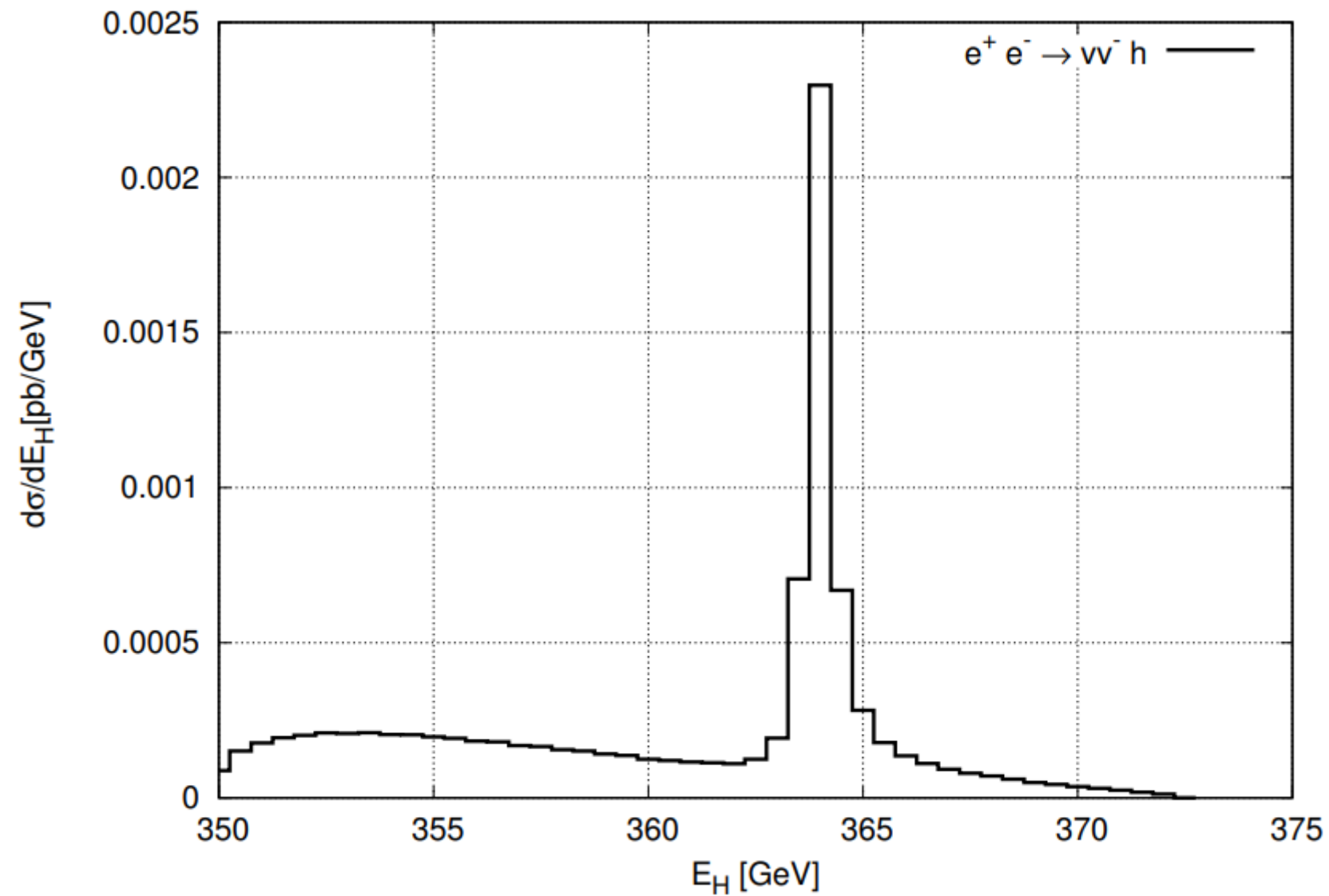


Figure 2. Higgs energy distribution

<https://arxiv.org/abs/hep-ph/0211268>

Result at 2002

7.2. Distribution

Result of Higgs boson energy and $m_{\nu\nu^-}$ for the $e^+e^- \rightarrow \nu\bar{\nu}H$ process at $m_H = 350$ [GeV] with $\sqrt{s} = 500$ GeV

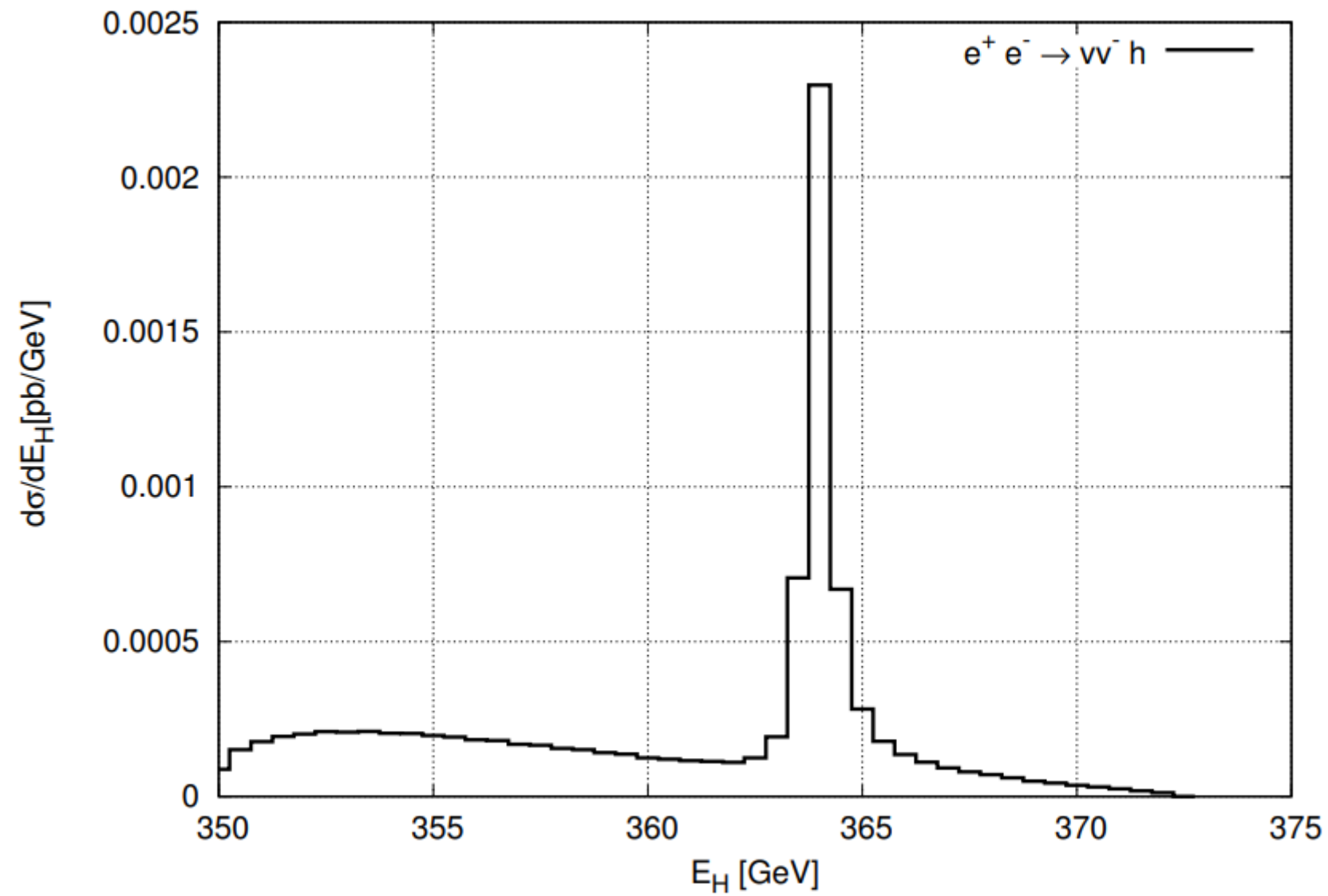


Figure 2. Higgs energy distribution

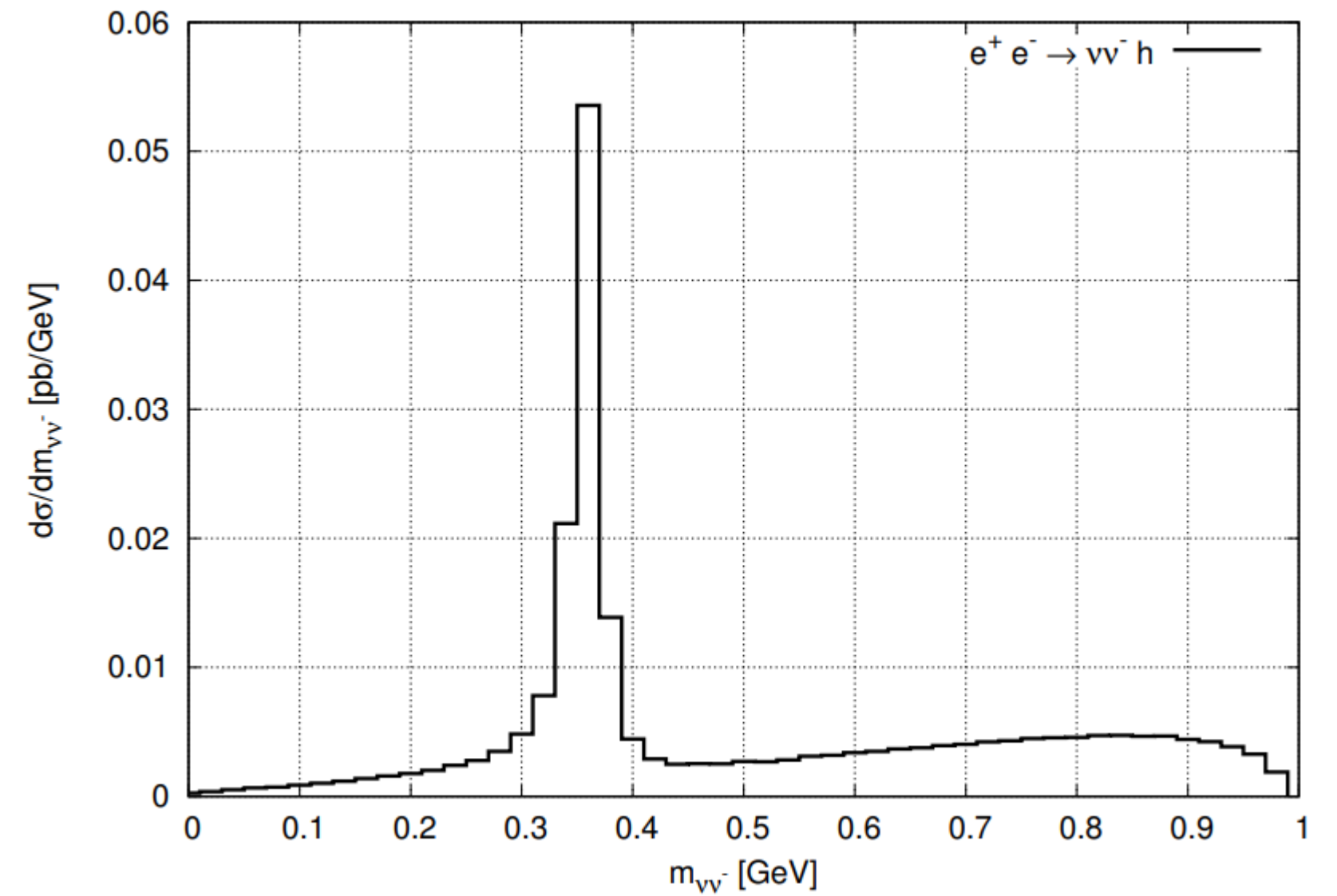


Figure 3. Energy distribution of $m_{\nu\nu^-}$

7.2. Distribution

Perform calculations with Higgs boson mass $m_H = 125$ [GeV] for the $e^+e^- \rightarrow \nu\bar{\nu}H$ process with $\sqrt{s} = 500$ GeV.

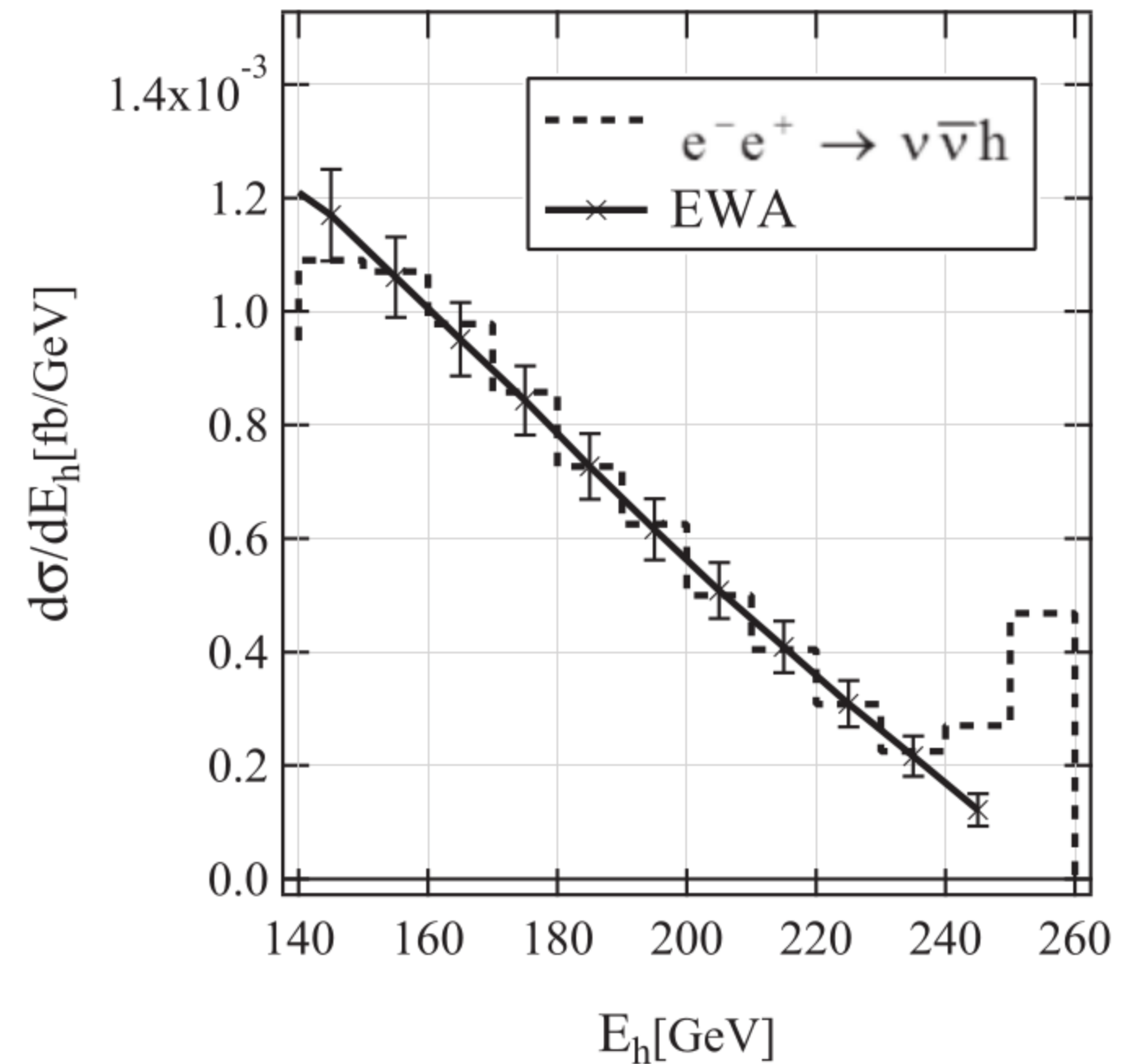
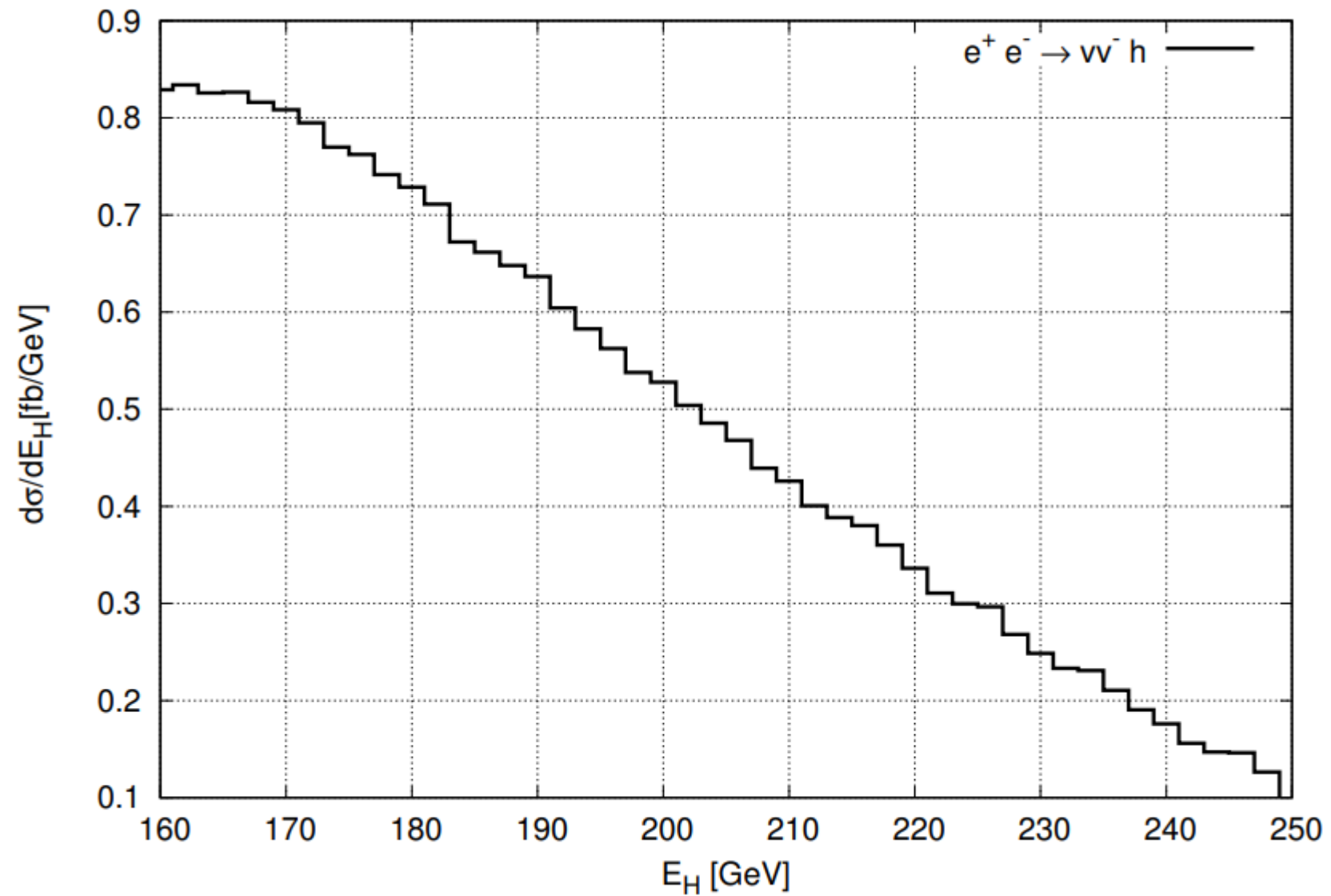


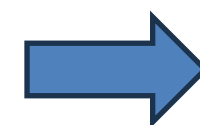
Figure 4. Higgs energy distribution at $m_H = 125$ [GeV]

<https://arxiv.org/abs/1807.01911>

8. Summary

In summary, we have performed calculations to verify the Standard Model, studying the Physical properties through $e^+e^- \rightarrow \nu\bar{\nu}H$ process. From here, we get the comparison results and Physical distributions necessary to discuss the topic achieved.

1. Calculate the scattering amplitude $|\mathcal{M}|^2$ for $e^+e^- \rightarrow \nu\bar{\nu}H$ process
2. Apply phase space to build an expression to calculate the scattering cross section for the process
3. Calculate the scattering cross sections using the Monte Carlo numerical solution method
4. Comparison with existing results and present the distribution



Through these steps, we can verify the theoretical results from the Standard Model and compare them with results obtained from other studies



Thank you for watching