

DETERMINATION OF THE n - ^{17}B SCATTERING LENGTH

Emeline OLIVEIRA - LPC Caen

WPCF - November 2024



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Case of interest: $n+^{17}\text{B}$

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➤ Both are unstable !

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EPJ Web Conf. 311, 00006 (2024)

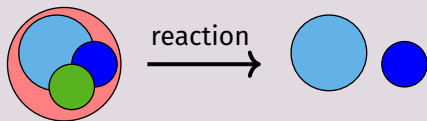
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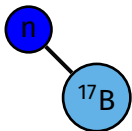
→ Knockout reaction

✓ Sudden removal



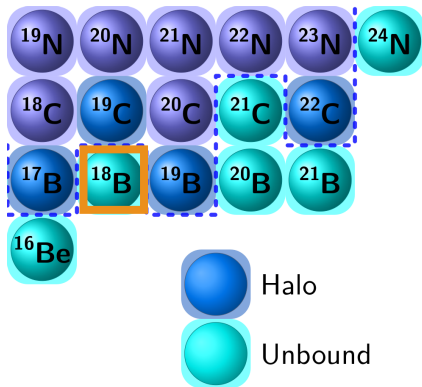
$$\sigma(k) \approx k \left| \int_0^\infty \psi_i(r) \varphi_k^*(r) r^2 dr \right|^2$$

THE STRUCTURE OF ^{18}B & ^{19}B

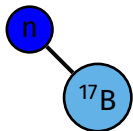


► Why $n+^{17}\text{B}$?

$$(a_s, r_e) \rightarrow \delta_0 \rightarrow V(n-^{17}\text{B})$$



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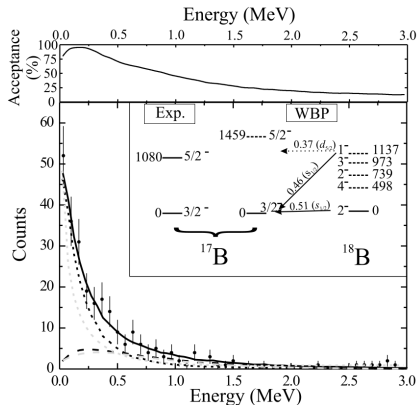
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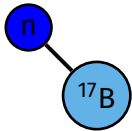
► Why $V(n-^{17}\text{B})$?

Only measure: $a_s < -50 \text{ fm} !$

📄 Spyrou, Phy. Let. B 683 (2010)



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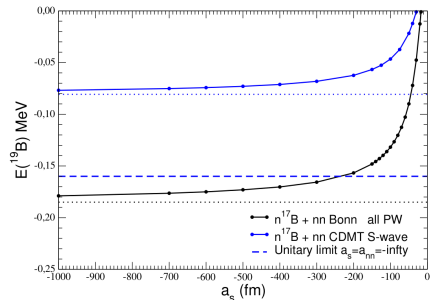
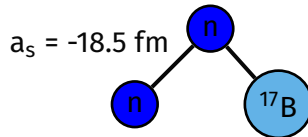
► Why $V(n-^{17}\text{B})$?

Only measure: $a_s < -50$ fm !

► $^{17}\text{B}+n+n$: Borromean binding from two large a_s ...

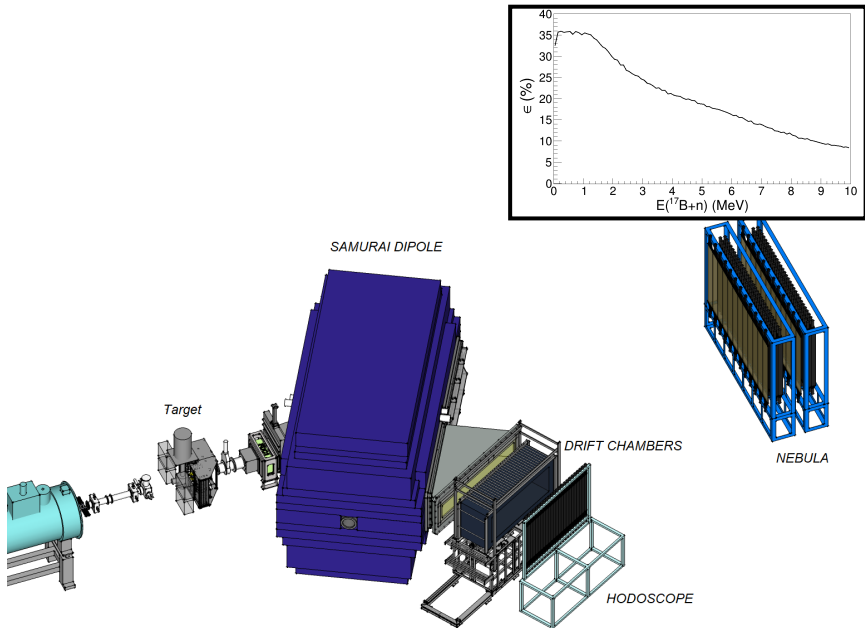
Efimov states ?

Efimov physics ! $\frac{|a_s|}{r_e} \gg 1$

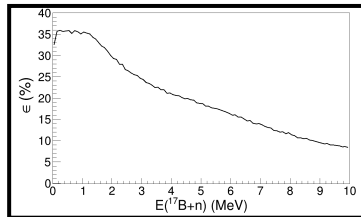
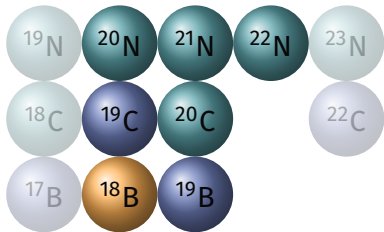


Hiyama PRC 100, 011603(R) (2019)

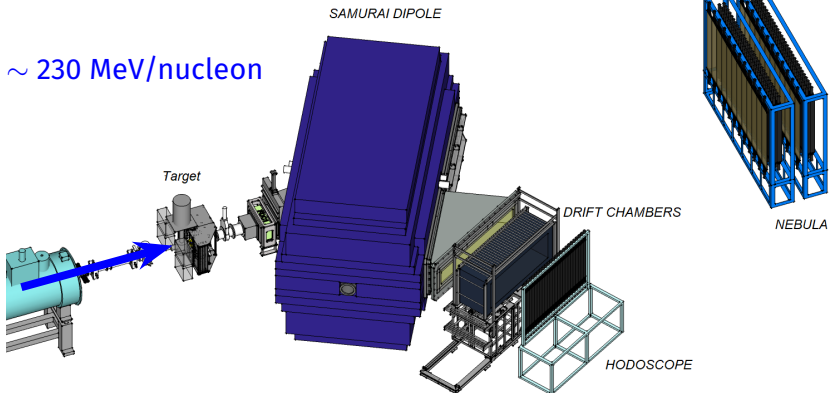
EXPERIMENTAL SETUP: SAMURAI DayOne (RIKEN)



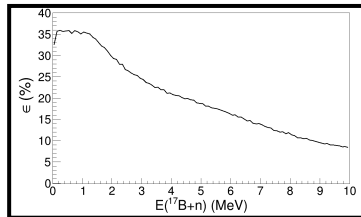
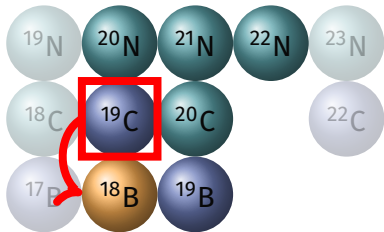
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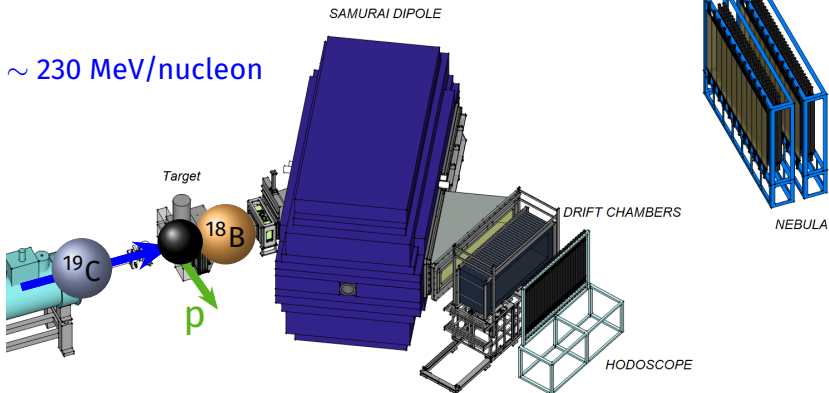
~ 230 MeV/nucleon



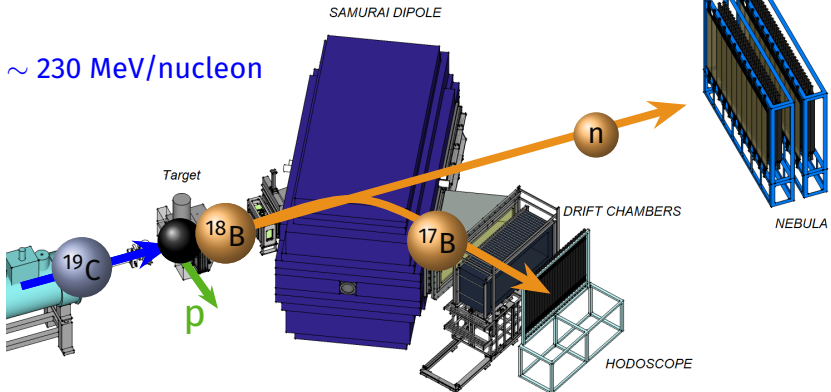
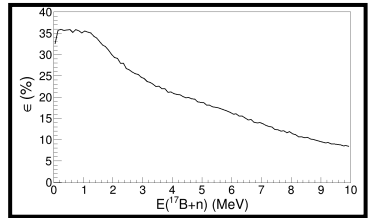
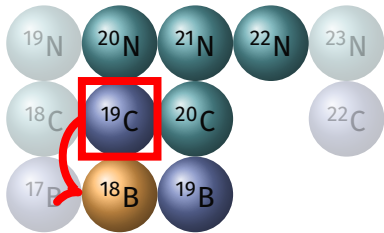
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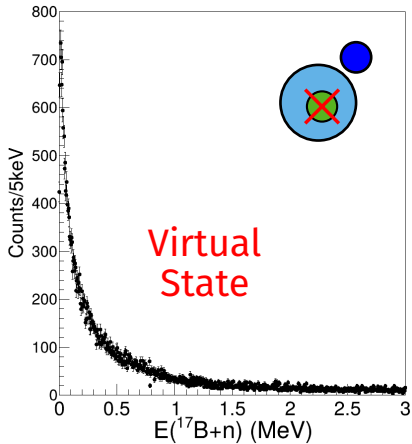
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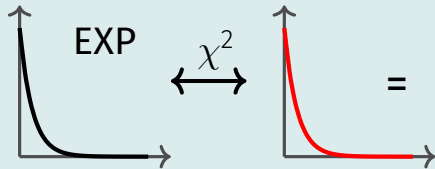
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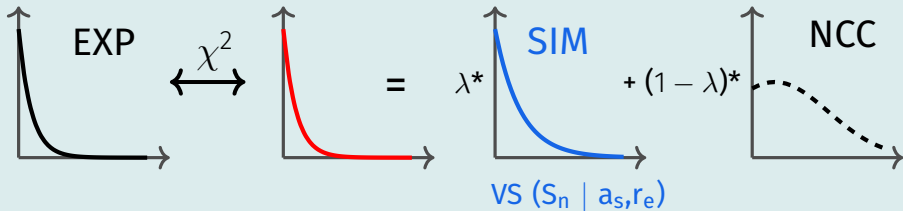
PARAMETERS CHARACTERISATION METHOD



Minimization

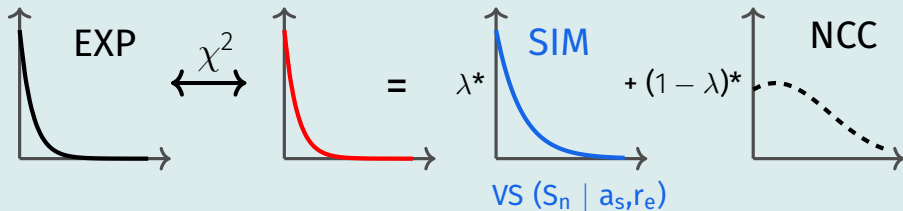


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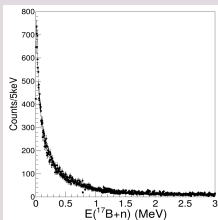


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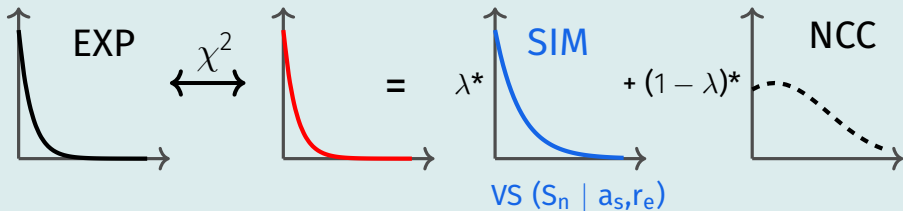
Non-correlated contribution



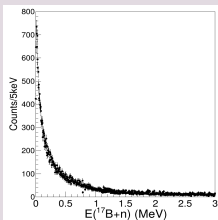
- ^{18}B ($^{17}\text{B}+n$)
- ^{17}B & n

PARAMETERS CHARACTERISATION METHOD

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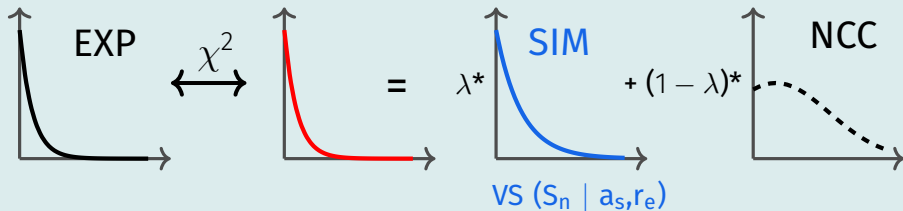


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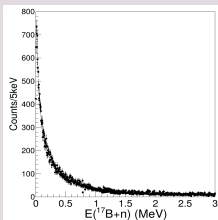


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Non-correlated contribution



$\left\{ \begin{array}{l} \blacktriangleright \text{}^{18}\text{B} (\text{}^{17}\text{B} + n) \\ \blacktriangleright \text{}^{17}\text{B} \ \& \ n \end{array} \right.$

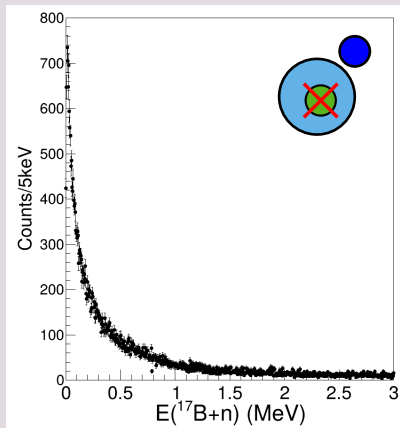
Event
 Mixing

i	f	n
1	●	●
2	●	●
...
N	●	●

✓ Shape

✓ Measured
events

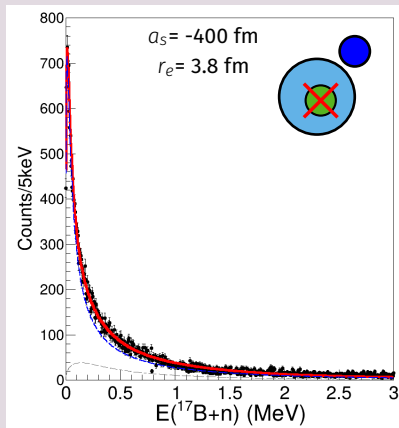
Fit Example



$$\text{EXP} \xleftrightarrow{\chi^2} \text{Fit} = \lambda^* \text{SIM} + (1 - \lambda)^* \text{NCC}$$

The diagram illustrates the fit equation: $\text{EXP} \xleftrightarrow{\chi^2} \text{Fit} = \lambda^* \text{SIM} + (1 - \lambda)^* \text{NCC}$. The 'Fit' is represented by a red curve, 'SIM' by a blue curve, and 'NCC' by a dashed black curve.

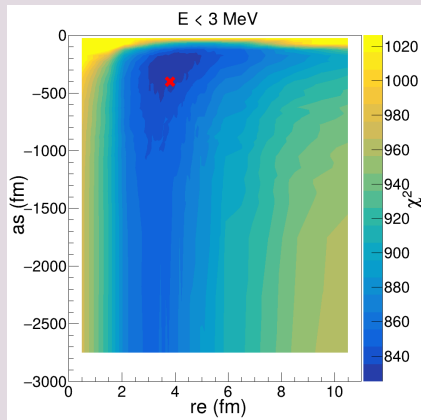
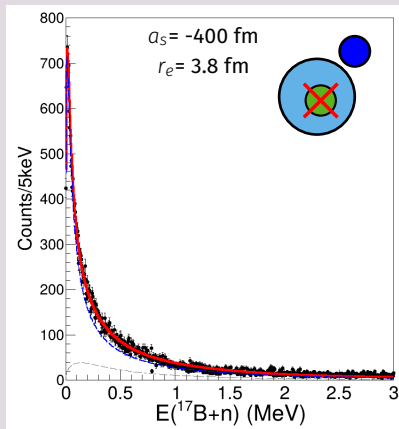
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$$\begin{array}{c} \uparrow \\ \text{EXP} \\ \downarrow \end{array} \longleftrightarrow \chi^2 \begin{array}{c} \uparrow \\ \text{red curve} \\ \downarrow \end{array} = \lambda^* \begin{array}{c} \uparrow \\ \text{SIM} \\ \downarrow \end{array} + (1 - \lambda)^* \begin{array}{c} \uparrow \\ \text{NCC} \\ \downarrow \end{array}$$

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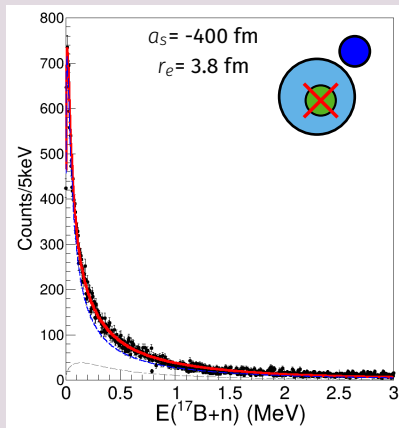
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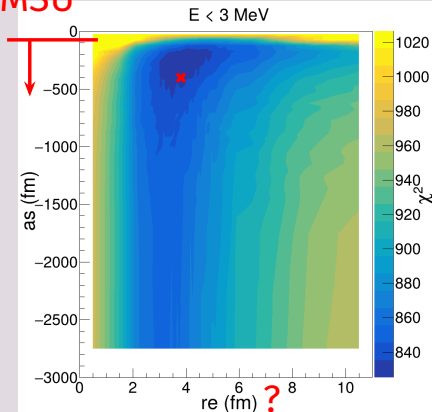
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The diagram illustrates the parameter characterization method. It shows the experimental data (EXP) being compared to a fit (Fit) using the χ^2 statistic. The fit is a combination of two simulated models: SIM (Simulation) and NCC (Neutron Capture Cross-section). The fit is expressed as $\lambda^* \text{SIM} + (1 - \lambda)^* \text{NCC}$, where λ^* is the fraction of the fit that is SIM.

Fit Example



MSU



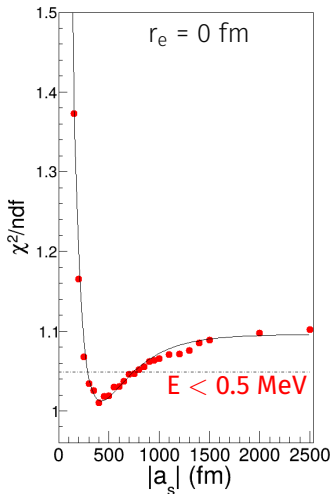
$$\text{EXP} \xleftrightarrow{\chi^2} \text{Fit} = \lambda^* \text{SIM} + (1 - \lambda)^* \text{NCC}$$

The diagram illustrates the decomposition of the experimental fit into two simulated components: SIM (Simulation) and NCC (Neutron Capture Cross-section). The fit is a red curve, SIM is a blue curve, and NCC is a dashed black curve. The parameter λ^* represents the fraction of the fit that is simulated.

PARAMETERS SENSITIVITY

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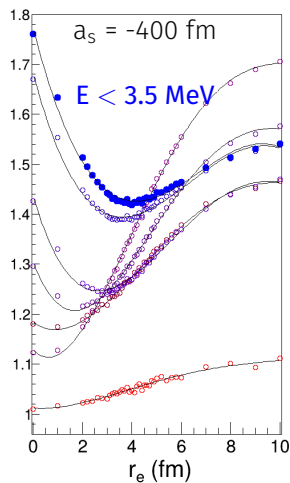
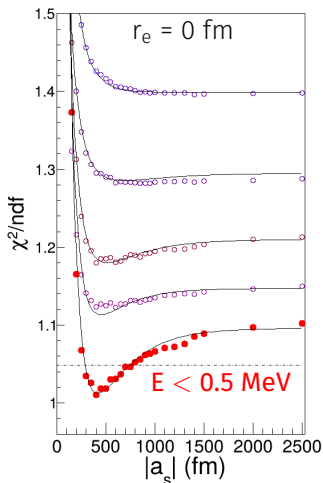
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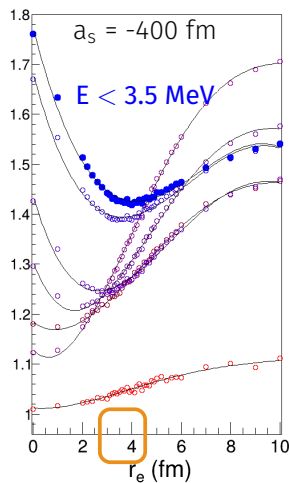
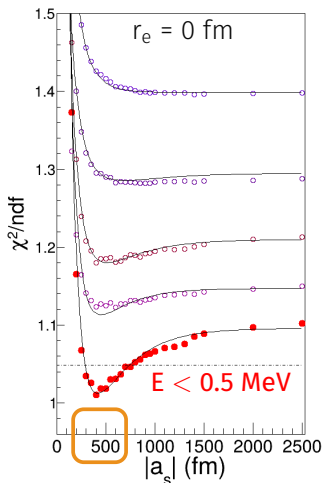
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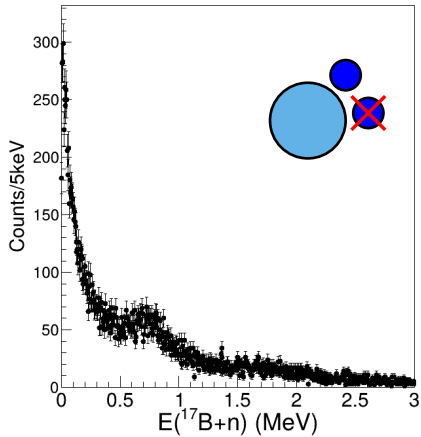
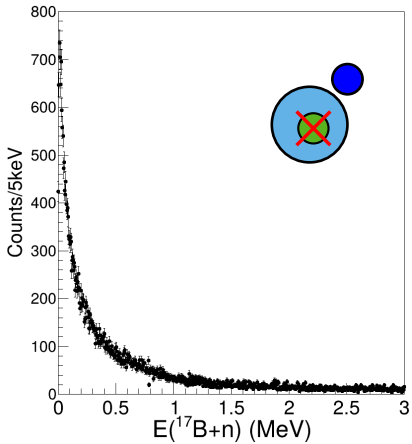
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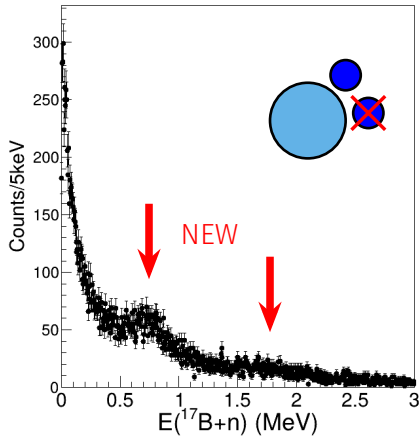
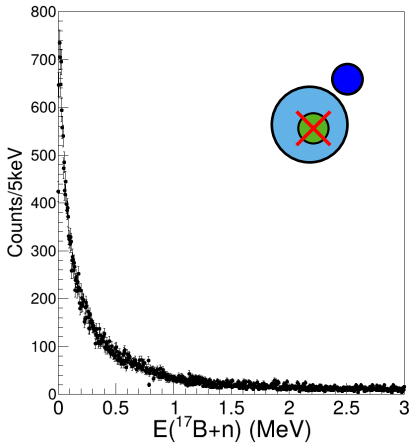


$$\frac{|a_s|}{r_e} \approx 100 !$$

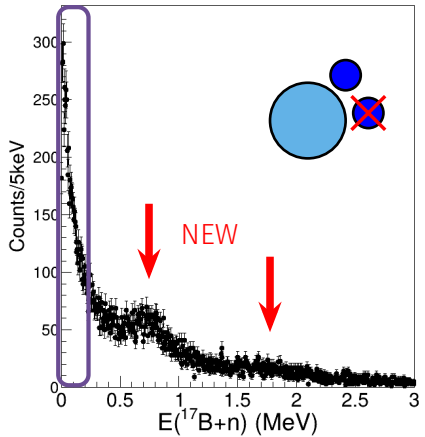
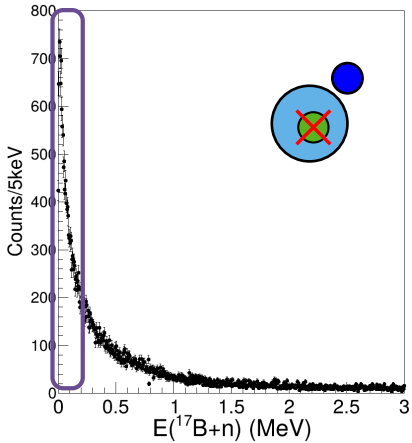
EXPLORATION OF ^{19}B CHANNEL



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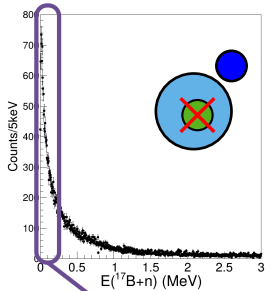


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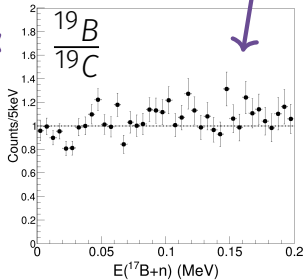
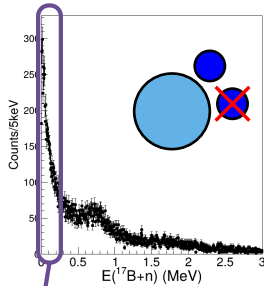


EXPLORATION OF ^{19}B CHANNEL

$$S_n(^{19}\text{C}) = 580 \pm 90 \text{ keV}$$

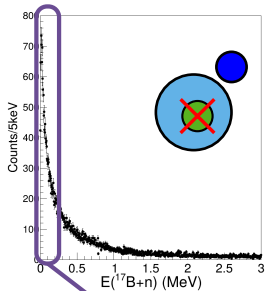


$$S_n(^{19}\text{B}) = 90 \pm 560 \text{ keV}$$

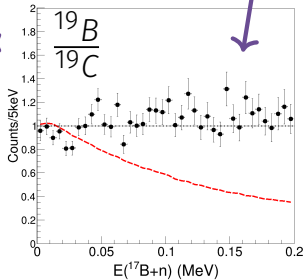
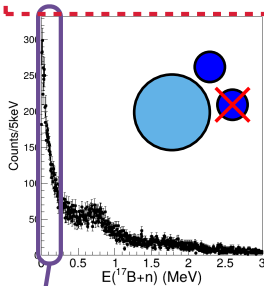


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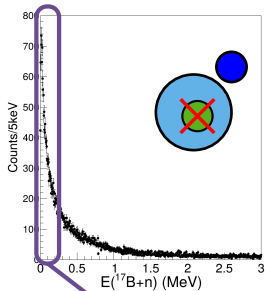


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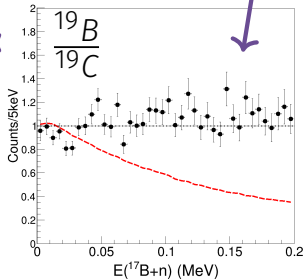
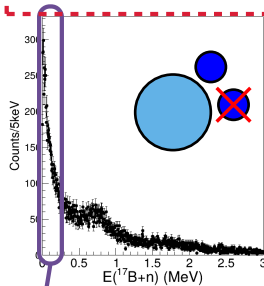


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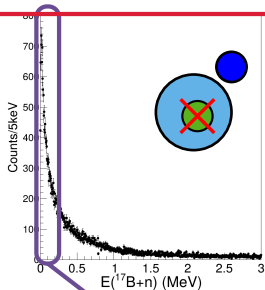


$\triangleright S_n(^{19}\text{B}) \sim 0.5 \text{ MeV} ?$

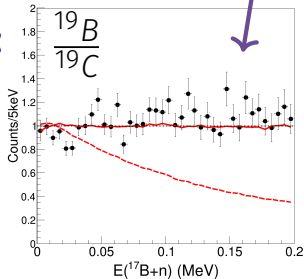
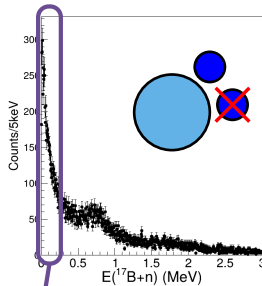
 Cook PRL 124, 212503 (2020)

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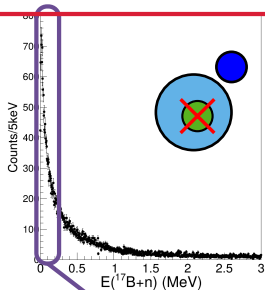


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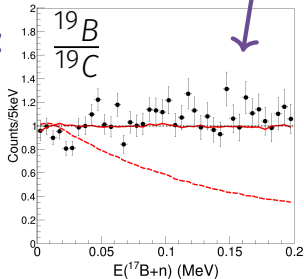
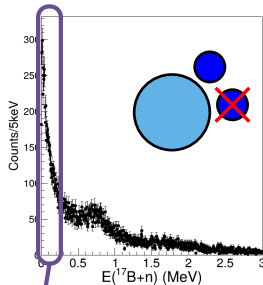
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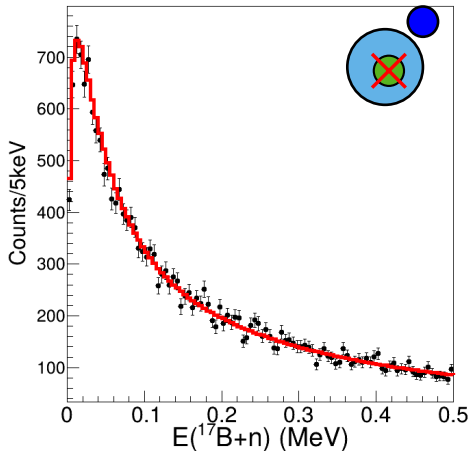
Re-estimation of $S_n(^{19}\text{B}) ?$

Low energy scattering

- Formalism
- Exotic nuclei (& n) unstable
- **n- ^{17}B** case: ^{18}B & ^{19}B
 - **scattering length**

Exp. program @ RIKEN

- Luminosity \otimes Acceptance:
 - **effective range**
- Multiple channels: (a_s, r_e)
 - $\frac{|a_s|}{r_e} \sim 100!$
- Measure of ^{19}B mass ?
- First resonances in ^{18}B



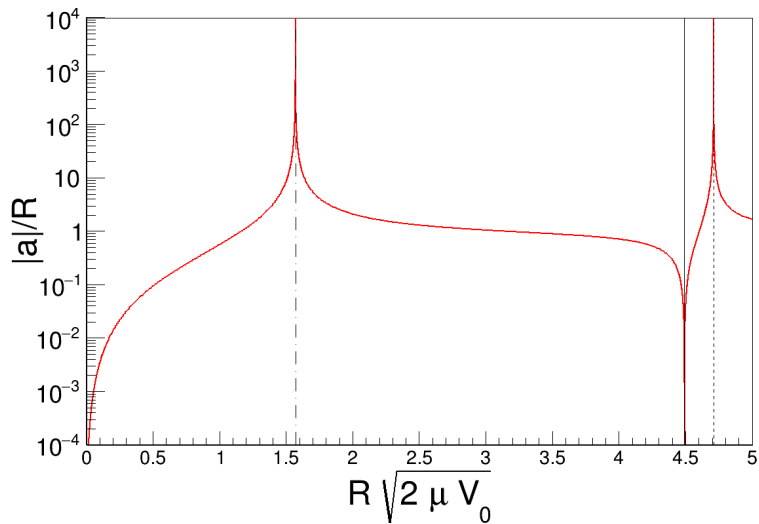
DETERMINATION OF THE n - ^{17}B SCATTERING LENGTH

Emeline OLIVEIRA - LPC Caen

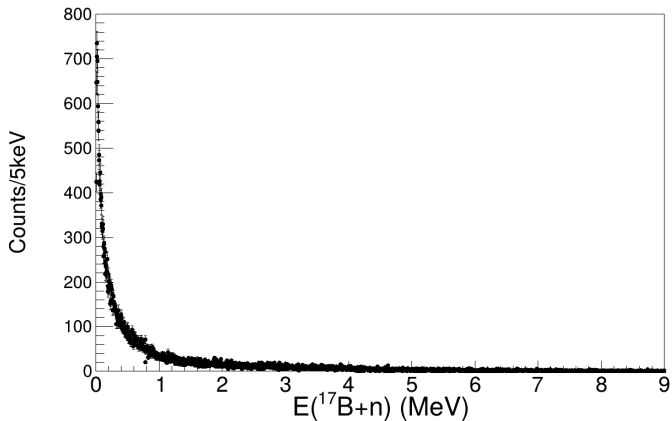
WPCF - November 2024



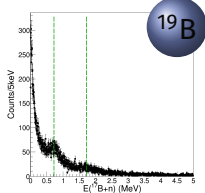
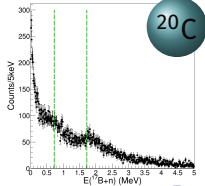
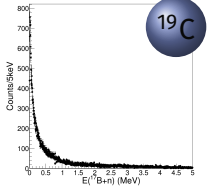
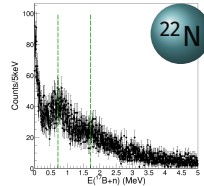
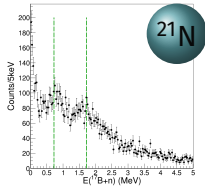
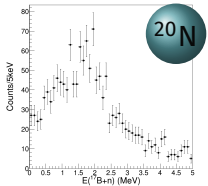
Particle off square well



Relative Energy: ^{19}C channel

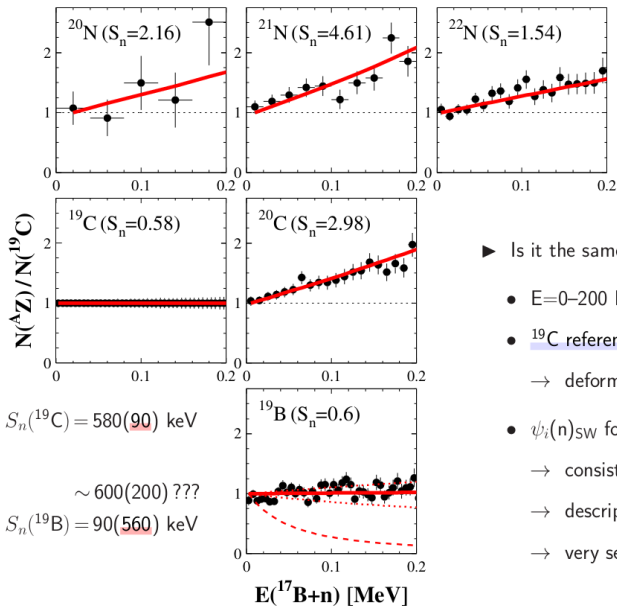


Relative energy for every channel



^{17}B
n

Effect of the initial state



► Is it the same virtual state?

- $E=0-200$ keV (no other resonances)
- ^{19}C reference ('pure' state)
 - deformation increases with S_n
- $\psi_i(n)_{\text{SW}}$ for $(a_s, r_c) = (-600, 4)$ fm
 - consistent with same state!
 - description of container effects
 - very sensitive to $S_n(^{19}\text{B})$!!!

Non Correlated Contribution

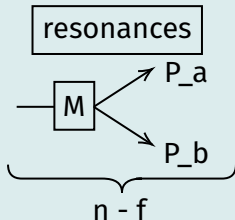
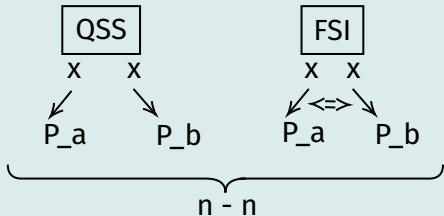
Non correlated events

Independent particles: $\begin{cases} a \\ b \end{cases} \Rightarrow \begin{cases} P_a \\ P_b \end{cases} \Rightarrow \frac{d\sigma}{dP_a} \frac{d\sigma}{dP_b}$

If (a,b) emitted together: $\frac{d^2\sigma}{dP_a dP_b} = \frac{d\sigma}{dP_a} \frac{d\sigma}{dP_b} \times C(P_a, P_b)$

$C(P_a, P_a) ? \rightarrow \frac{d\sigma}{dP_a} \frac{d\sigma}{dP_b}$

Mechanisms modifying (P_a, P_b)



Event Mixing

Independent particles \rightarrow virtual pairs $\rightarrow E_{rel}$

i	a	b
1	●	●
2	●	●
...
N	●	●

$$\begin{aligned} \frac{d\sigma_{\otimes}}{dP_a} &= \int \frac{d^2\sigma}{dP_a dP_b} dP_b = \frac{d\sigma}{dP_a} \int C(P_a, P_b) \frac{d\sigma}{dP_b} dP_b \\ &= \frac{d\sigma}{dP_a} \langle C \rangle (P_a) \end{aligned}$$

Event Mixing

Independent particles \rightarrow virtual pairs $\rightarrow E_{\text{rel}}$

i	a	b
1	●	●
2	●	●
...
N	●	●

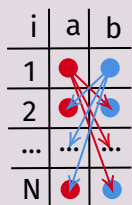
$$\frac{d\sigma_{\otimes}}{dP_a} = \int \frac{d^2\sigma}{dP_a dP_b} dP_b = \frac{d\sigma}{dP_a} \int C(P_a, P_b) \frac{d\sigma}{dP_b} dP_b$$

$$= \frac{d\sigma}{dP_a} \langle C \rangle (P_a)$$

$$\frac{d\sigma_{\otimes}}{dP_b} = \frac{d\sigma}{dP_b} \langle C \rangle (P_b)$$

Event Mixing

Independent particles \rightarrow virtual pairs $\rightarrow E_{rel}$



$$\frac{d\sigma_{\otimes}}{dP_a} = \int \frac{d^2\sigma}{dP_a dP_b} dP_b = \frac{d\sigma}{dP_a} \int C(P_a, P_b) \frac{d\sigma}{dP_b} dP_b$$

$$= \frac{d\sigma}{dP_a} \langle C \rangle (P_a)$$

$$\frac{d\sigma_{\otimes}}{dP_b} = \frac{d\sigma}{dP_b} \langle C \rangle (P_b)$$

if $\langle C \rangle (P) \gg 1 \Rightarrow C(P_a, P_b) \geq \frac{d^2\sigma/dP_a dP_b}{(d\sigma_{\otimes}/dP_a)(d\sigma_{\otimes}/dP_b)}$

Event Mixing

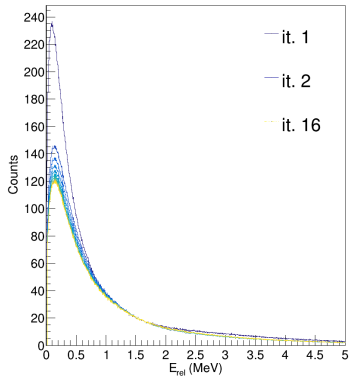


Figure: All distribution of event mixing procedure

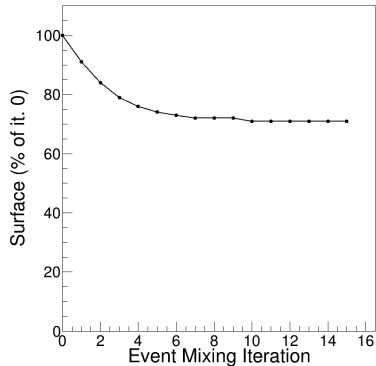
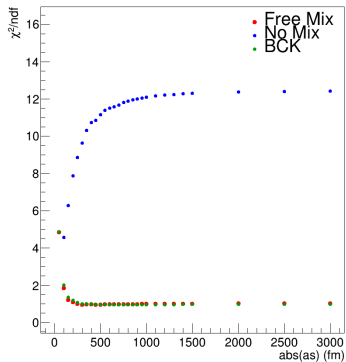


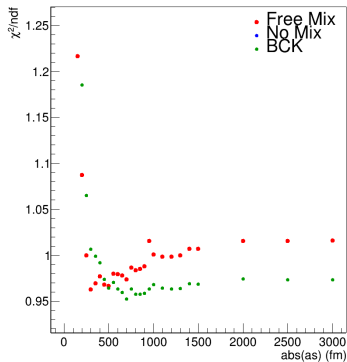
Figure: Convergence of the event mixing procedure

Non Correlated Contribution

Section in $r_e = 0$ (fm), $E < 0.5$ MeV



Section in $r_e = 0$ (fm), $E < 0.5$ MeV



Non-Correlated Distribution

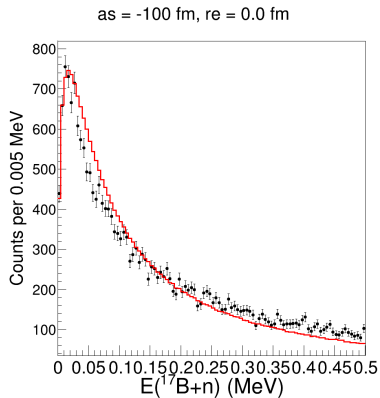


Figure: Best Fit without Mix, $\chi^2 = 4.57$

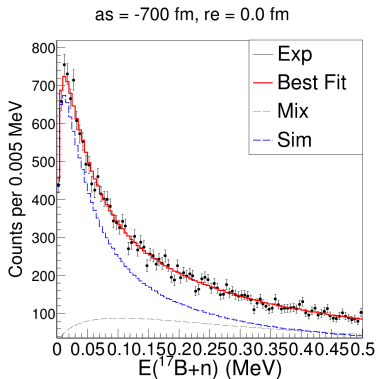


Figure: Best Fit with simulated Mix, $\chi^2 = 0.95$

Induced long-range interaction

→ Effective interaction is mediated between two particles by the third particle moving back and forth between the two. It is thus possible for the three particles to feel their influence at **distances much larger than the range of interactions**, typically up to **distances on the order of the scattering length**.

Discrete scale invariance

→ **infinite series of bound states**, the Efimov trimers, whose properties such as size and energy are related to each others' by a scale transformation with a **universal scaling factor**.

Borromean binding

When the interaction is not strong enough to support a two-body bound state, it may nonetheless support one, up to infinitely many, Efimov trimers. This possibility of **binding N particles, while the $N-1$ subsystems are unbound** is called 'Borromean' binding.

Alignment

$$\text{Min}(\langle E_{rel} \rangle) \quad \langle \beta_n - \beta_f \rangle = 0 \quad \langle P_z^f(n) \rangle = 0 \quad \text{Max}(E_{rel})$$

