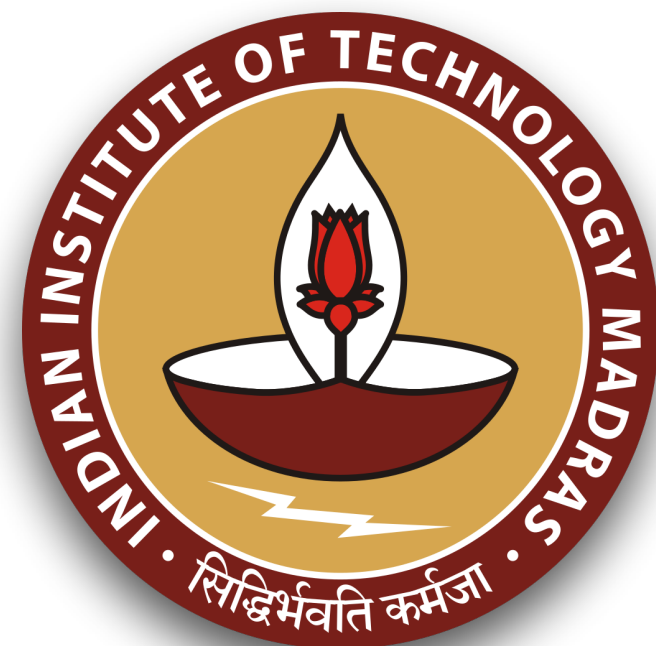


# Strange particle femtoscopy in PbPb collisions at 5.02 TeV with the CMS detector

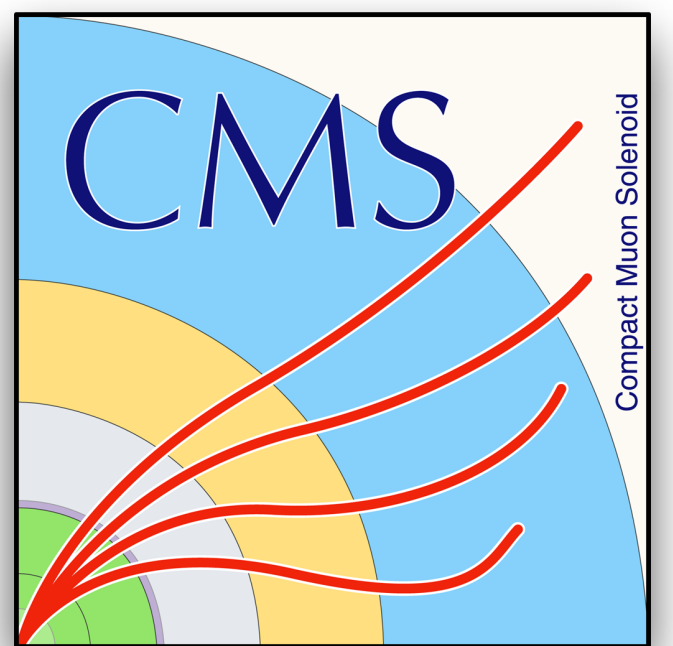
**Raghunath Pradhan**

*for the CMS collaboration*

**University of Illinois, Chicago**



**WPCE 2024**



# Introduction: femtoscopy

- **Femtoscopy:** Powerful tool to probe space-time dimensions of the particle emitting source region on the femtometer scale
- Use final state particle correlations

## Femtoscopic correlation

Sensitive to

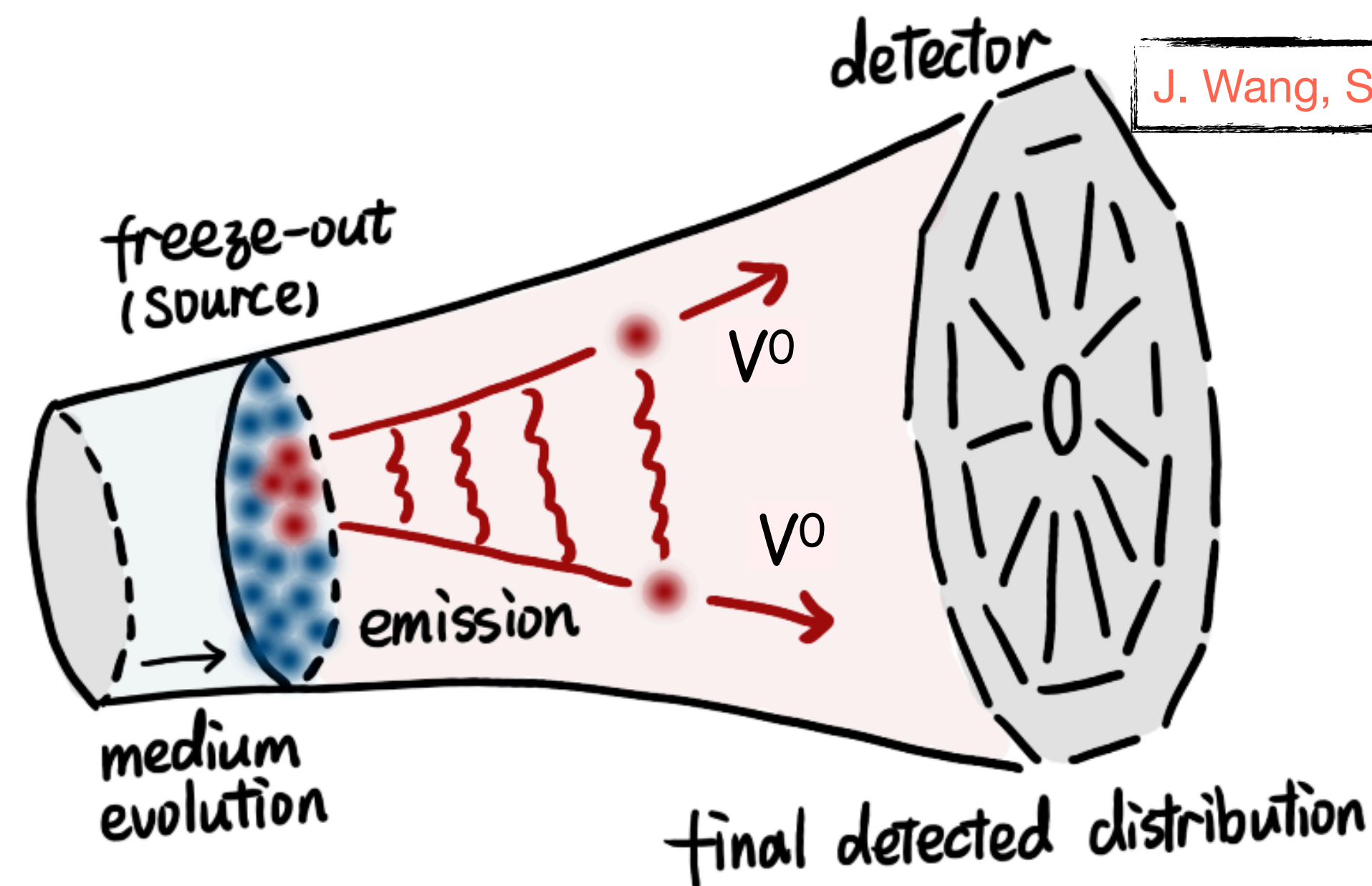
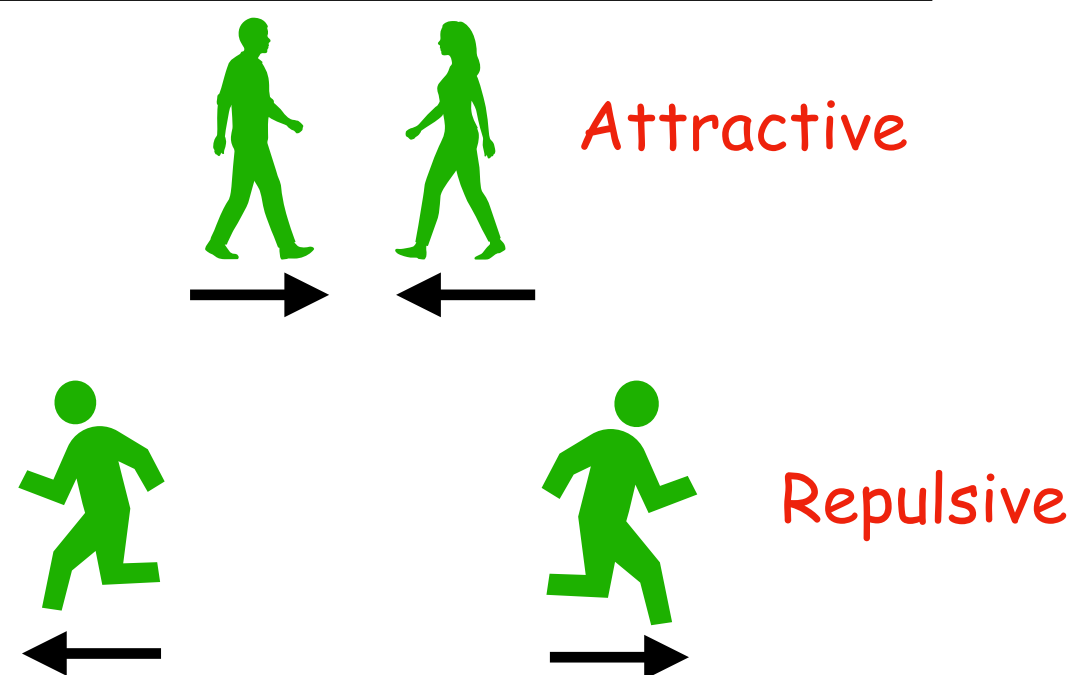
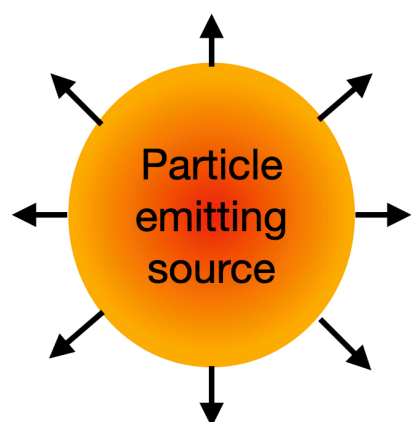
Quantum statistical effect (QS) :  
Bose - Einstein or Fermi - Dirac

Possible final state interaction (FSI) :  
Strong, Coulomb,..

Help us to  
measure

Size and shape of the  
particles emitting source  
at kinetic freeze-out

Which type of interaction?  
Attractive or repulsive?

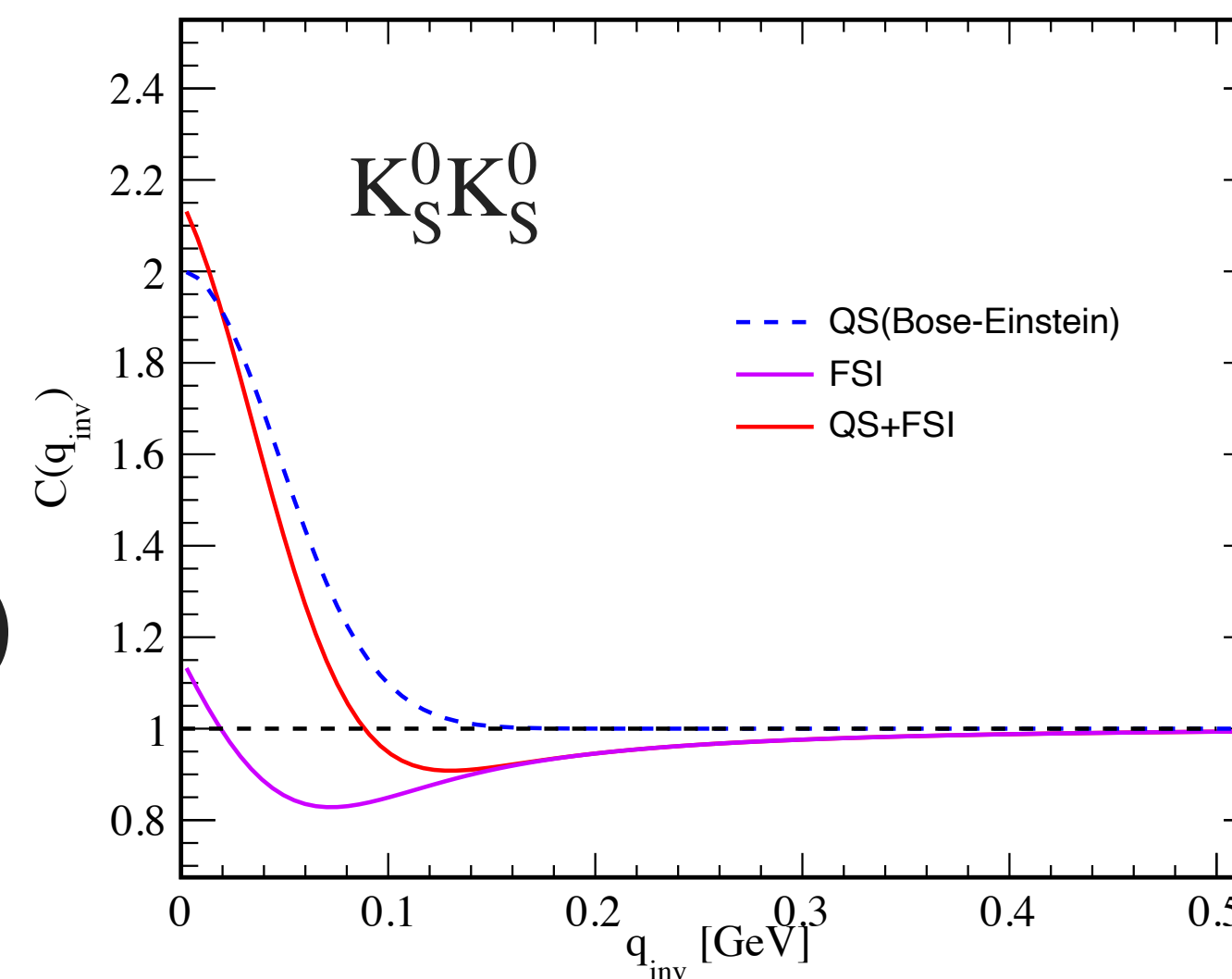


J. Wang, SQM 2022

# Motivation: $V^0$ femtoscopy

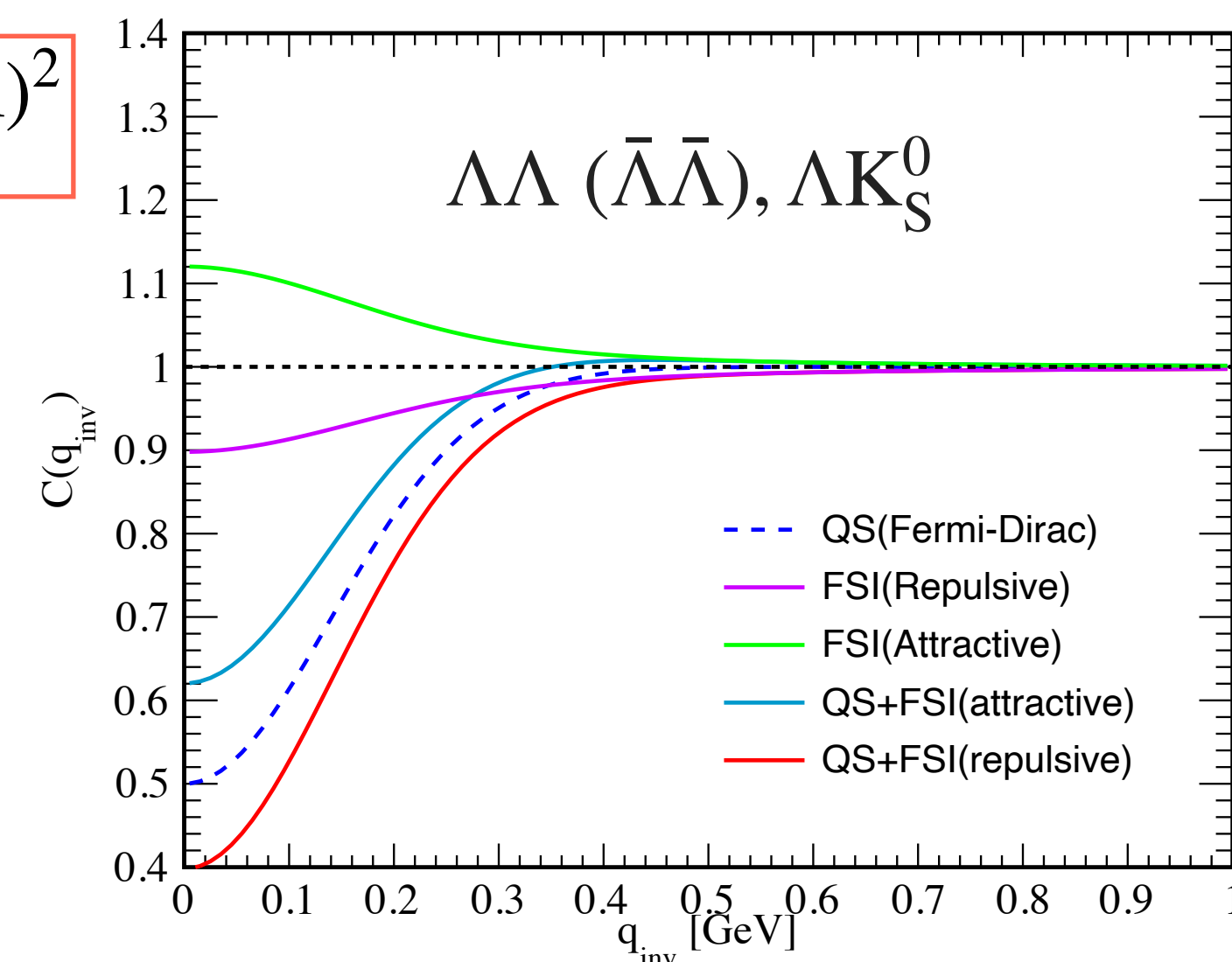
- **Why study  $V^0$  particles ( $\Lambda(\bar{\Lambda})$  &  $K_S^0$ ) femtoscopic correlation?**
  - No coulomb interaction
  - Quantum statistical (QS) effect and strong final state interaction (FSI)
  - Less resonance contribution (less feed down contribution)
  - Size of the particles emitting source
  - Interaction between baryons and mesons
    - ▶ Strong interaction scattering parameters
      - ➔ Scattering length and effective range
  - $\Lambda\Lambda(\bar{\Lambda}\bar{\Lambda})$  correlation is relevant for searching bound H-dibaryon

Phys. Rev. Lett. 38, 195



Assumed Gaussian source

$$S(r) \sim e^{-(r/R)^2}$$



# CMS detector and $V^0$ decay

## CMS DETECTOR

Total weight : 14,000 tonnes  
 Overall diameter : 15.0 m  
 Overall length : 28.7 m  
 Magnetic field : 3.8 T

STEEL RETURN YOKE  
 12,500 tonnes

SILICON TRACKERS  
 Pixel (100x150  $\mu\text{m}$ )  $\sim 16\text{m}^2 \sim 66\text{M}$  channels  
 Microstrips (80x180  $\mu\text{m}$ )  $\sim 200\text{m}^2 \sim 9.6\text{M}$  channels

SUPERCONDUCTING SOLENOID  
 Niobium titanium coil carrying  $\sim 18,000\text{A}$

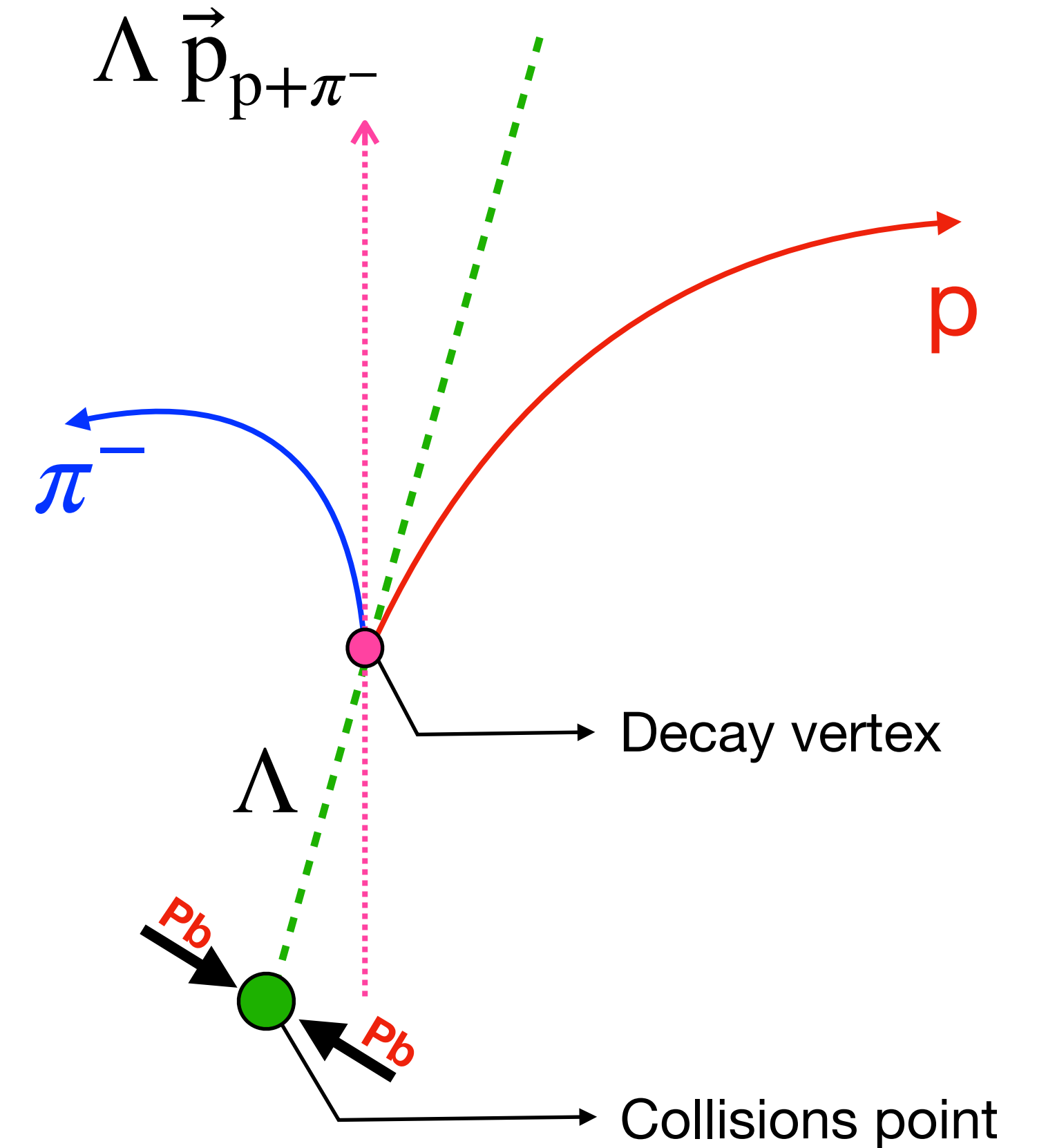
MUON CHAMBERS  
 Barrel: 250 Drift Tube, 480 Resistive Plate Chambers  
 Endcaps: 468 Cathode Strip, 432 Resistive Plate Chambers

PRESHOWER  
 Silicon strips  $\sim 16\text{m}^2 \sim 137,000$  channels

FORWARD CALORIMETER  
 Steel + Quartz fibres  $\sim 2,000$  Channels

CRYSTAL ELECTROMAGNETIC CALORIMETER (ECAL)  
 $\sim 76,000$  scintillating  $\text{PbWO}_4$  crystals

HADRON CALORIMETER (HCAL)  
 Brass + Plastic scintillator  $\sim 7,000$  channels

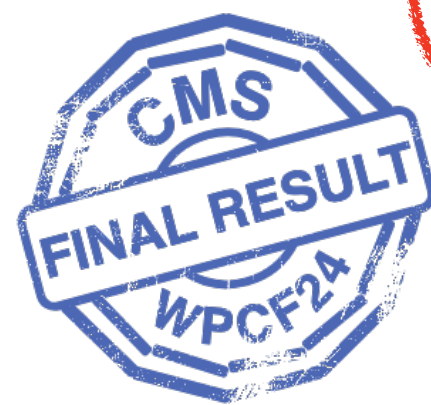


- $\Lambda \rightarrow p + \pi^-$  [(63.9  $\pm$  0.5)%]
- $K_S^0 \rightarrow \pi^+ + \pi^-$  [(69.20  $\pm$  0.05)%]

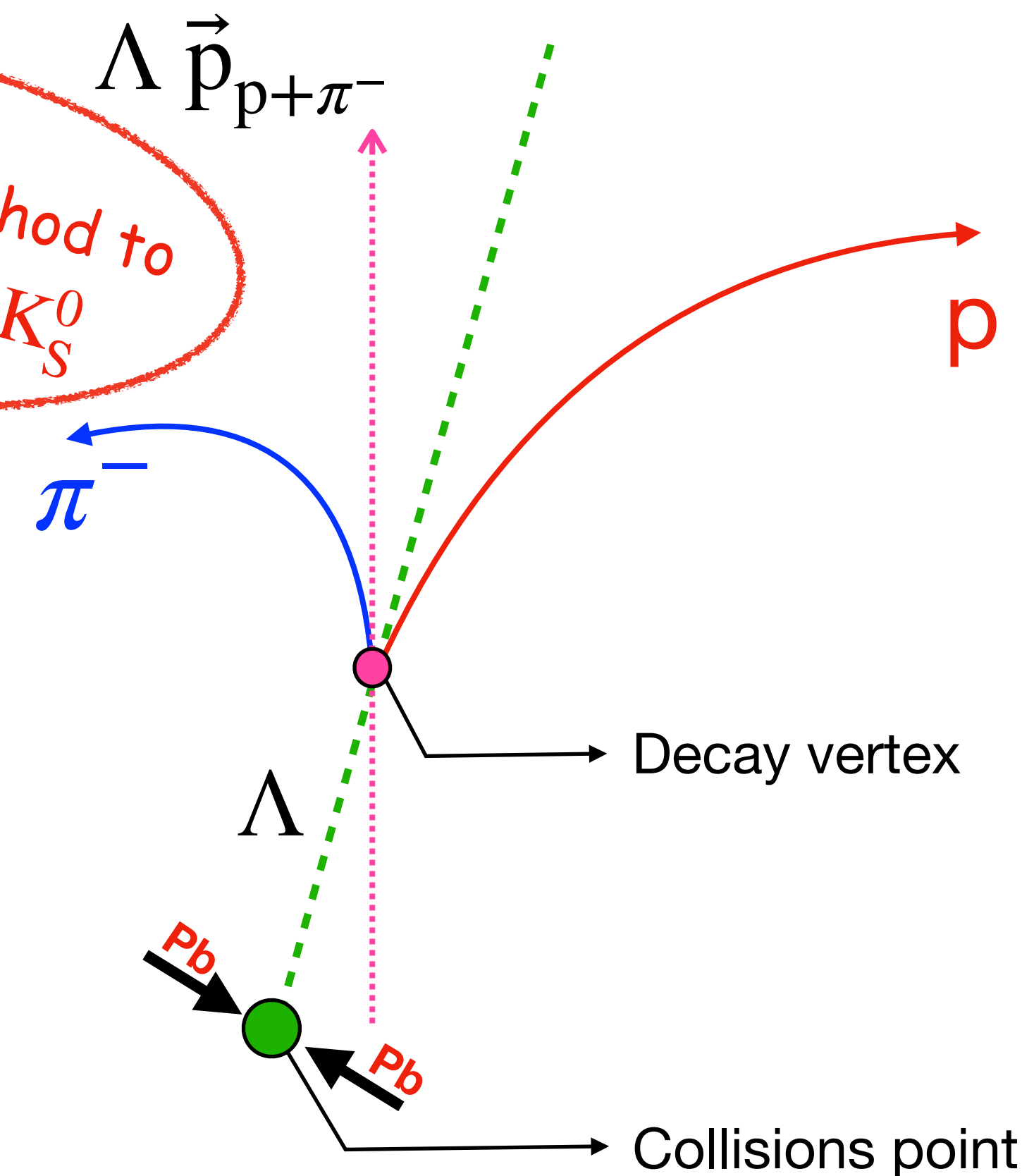
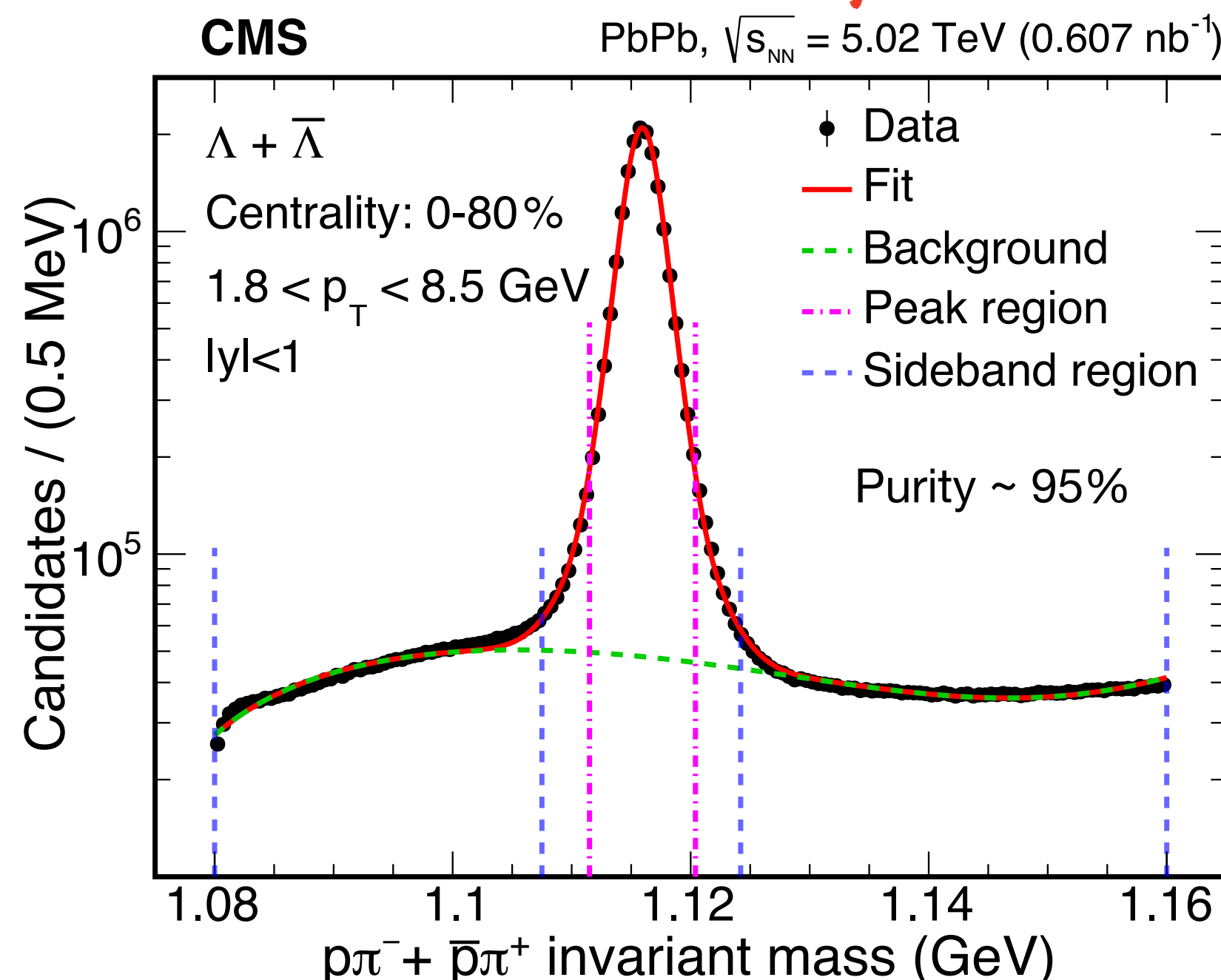
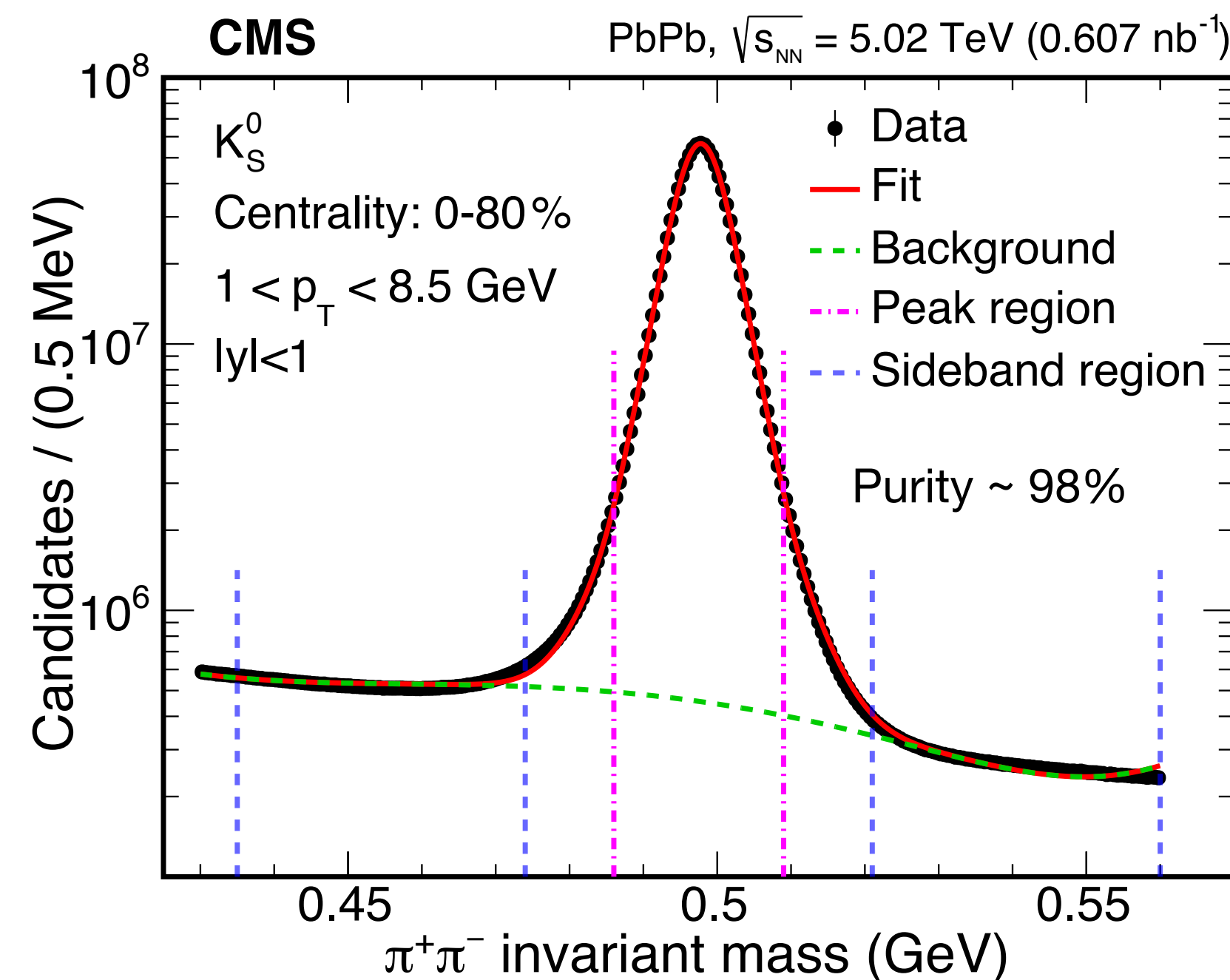
# $V^0$ particles reconstruction

- 2018 PbPb collisions @ 5.02 TeV
  - ▶ Minimum Bias
  - ▶ ~ 4 B events

PLB 857 (2024) 138936



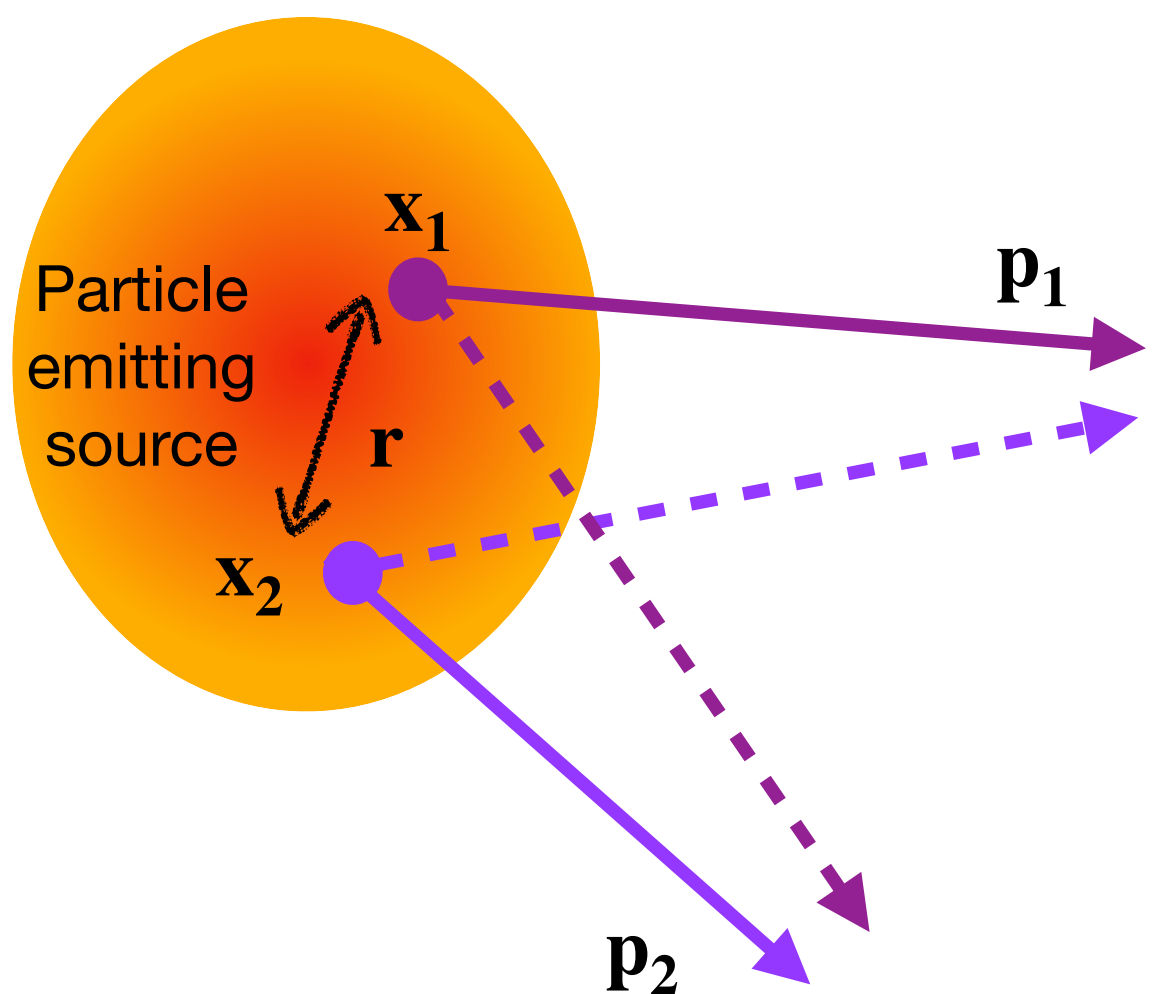
Applied BDT method to select  $\Lambda(\bar{\Lambda})$  &  $K_S^0$



- Signal : triple Gaussian
- Combinatorial background : 4<sup>th</sup> order polynomial

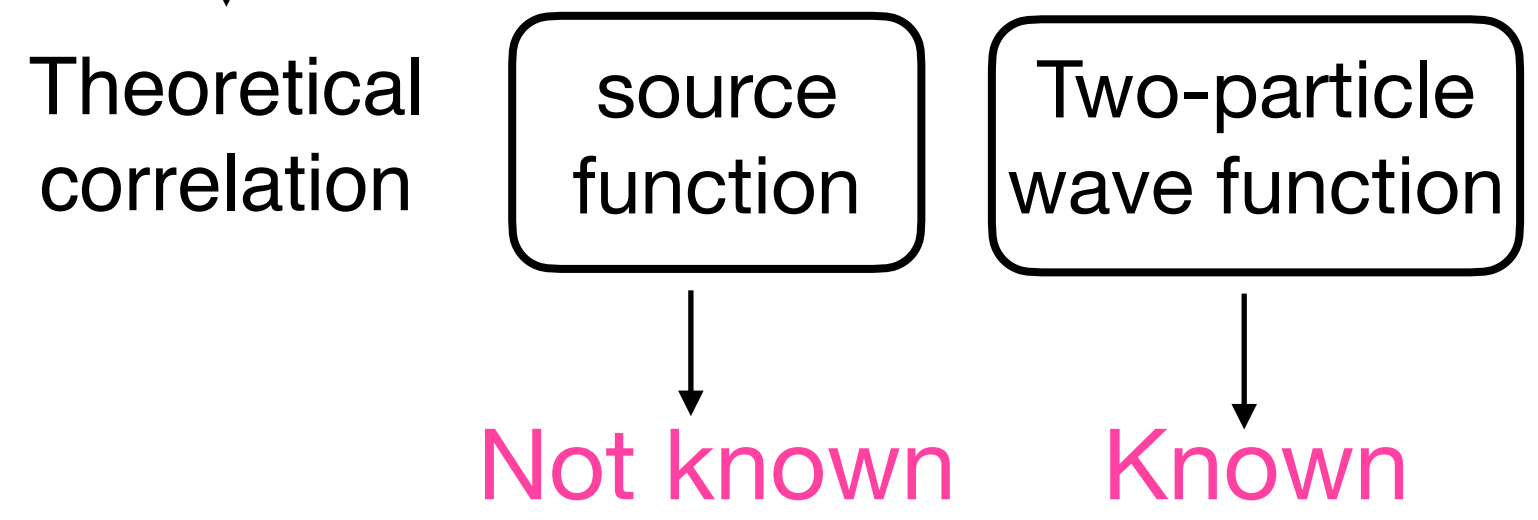
- $\Lambda \rightarrow p + \pi^-$  [(63.9 ± 0.5)%]
- $K_S^0 \rightarrow \pi^+ + \pi^-$  [(69.20 ± 0.05)%]

# Correlation function



● **In theory :**

$$C_K(q) = \int S(r) |\Psi_{1,2}(q, r)|^2 d^3r = 1 \pm C_{QS}(q) + C_{FSI}(q)$$



+ for identical bosons  
- for identical fermions

Generally we assume Gaussian source function

$$S(r) \sim e^{-(r/R)^2}$$

$$q = p_1 - p_2$$

$$K = \frac{p_1 + p_2}{2}$$

$$r = x_1 - x_2$$

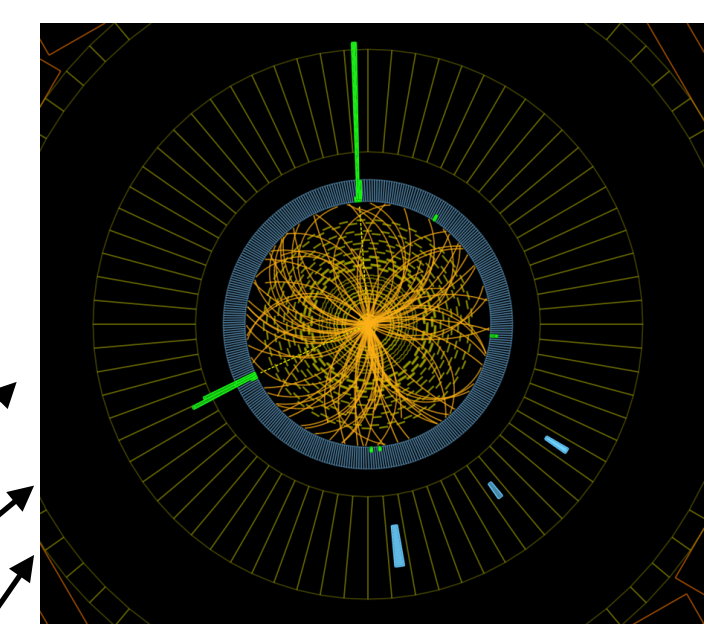
● **In the experiment :**

$$C(q_{inv}) = N \left[ \frac{A(q_{inv})}{B(q_{inv})} \right], \quad q_{inv} = |q^\mu|, \quad q^\mu = k^\mu - \frac{k \cdot P}{P^2} P^\mu,$$

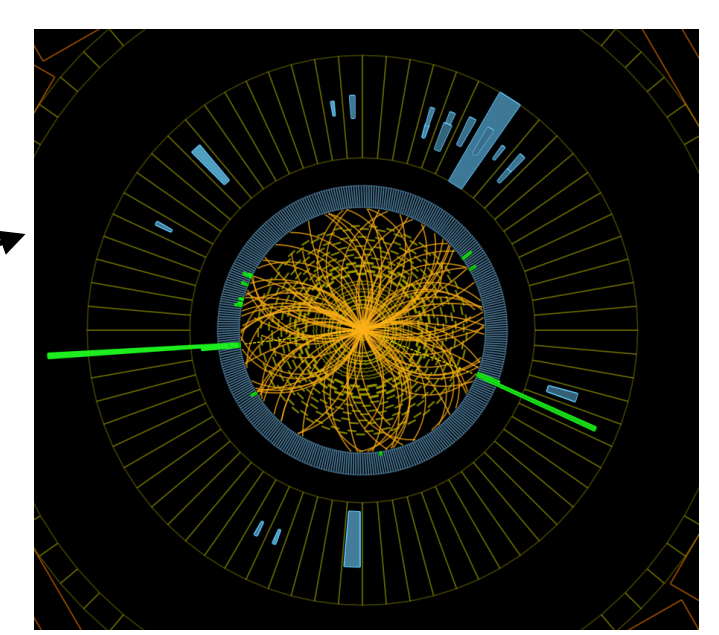
Ann.Rev.Nucl.Part.Sci.55:357-402,2005

$$k = p_1 - p_2, \quad P = p_1 + p_2$$

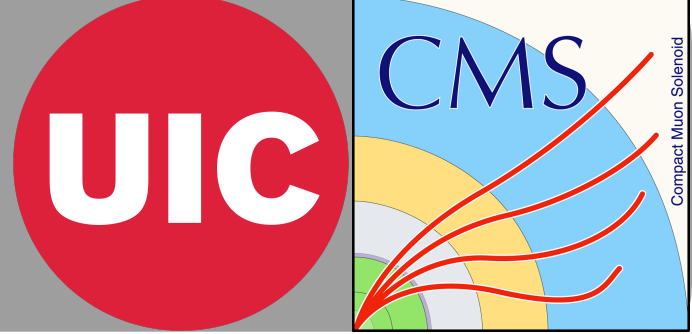
A(q<sub>inv</sub>): Signal distribution of pair from same event  
 B(q<sub>inv</sub>): Reference distribution of pair from mixed events  
 N: Normalization constant



[cds.cern.ch/record/2736135](https://cds.cern.ch/record/2736135)



# Results: correlation and fitting



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CMS

$1 < p_T < 8.5$  GeV  
 $0 < k_T < 2$  GeV

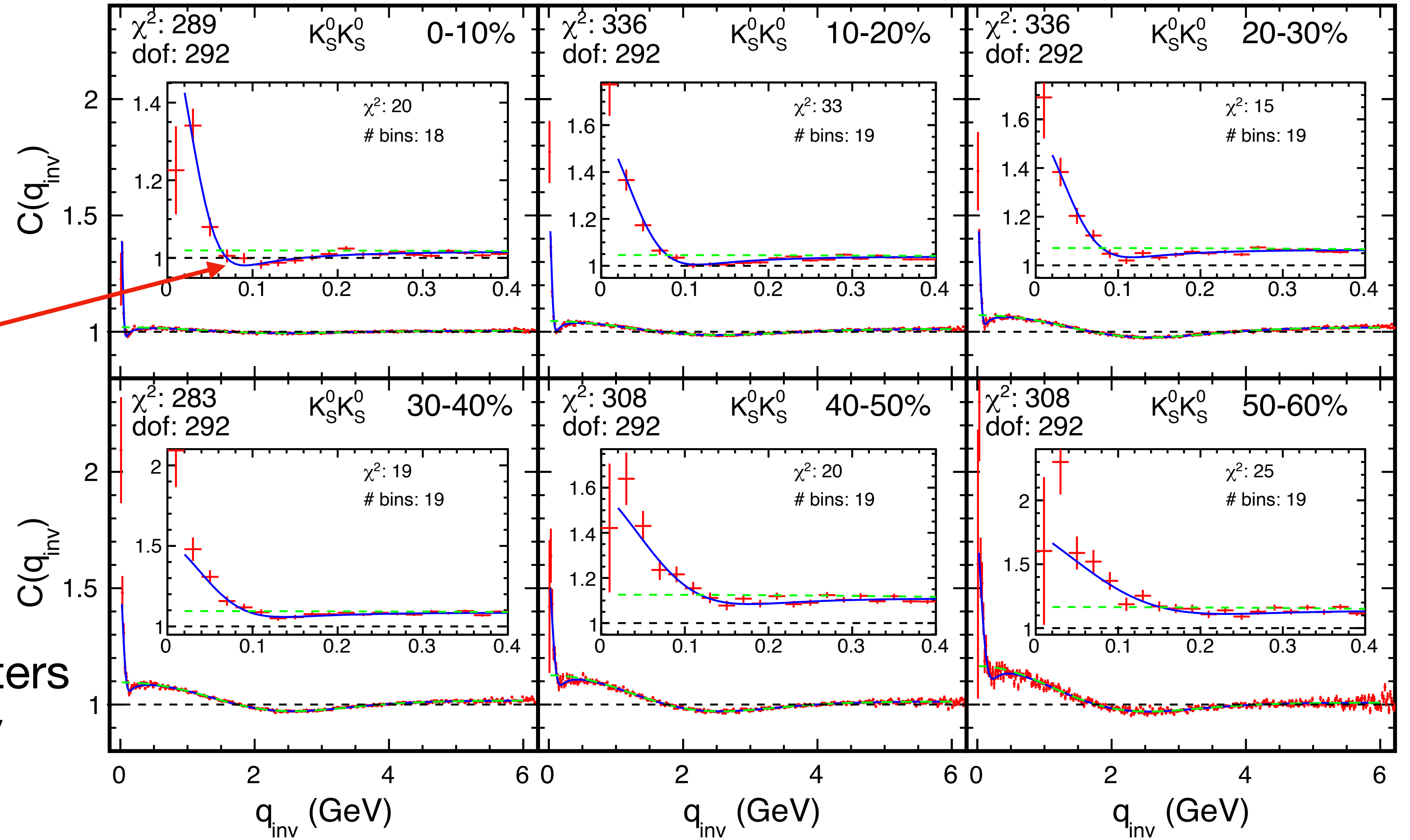
+ Data  
 — Full fit  
 - - - Nonfemto

PbPb,  $\sqrt{s_{NN}} = 5.02$  TeV ( $0.607$  nb $^{-1}$ )



$K_S^0 K_S^0$

QS (Bose-Einstein)  
 + strong FSI  
 (repulsive)



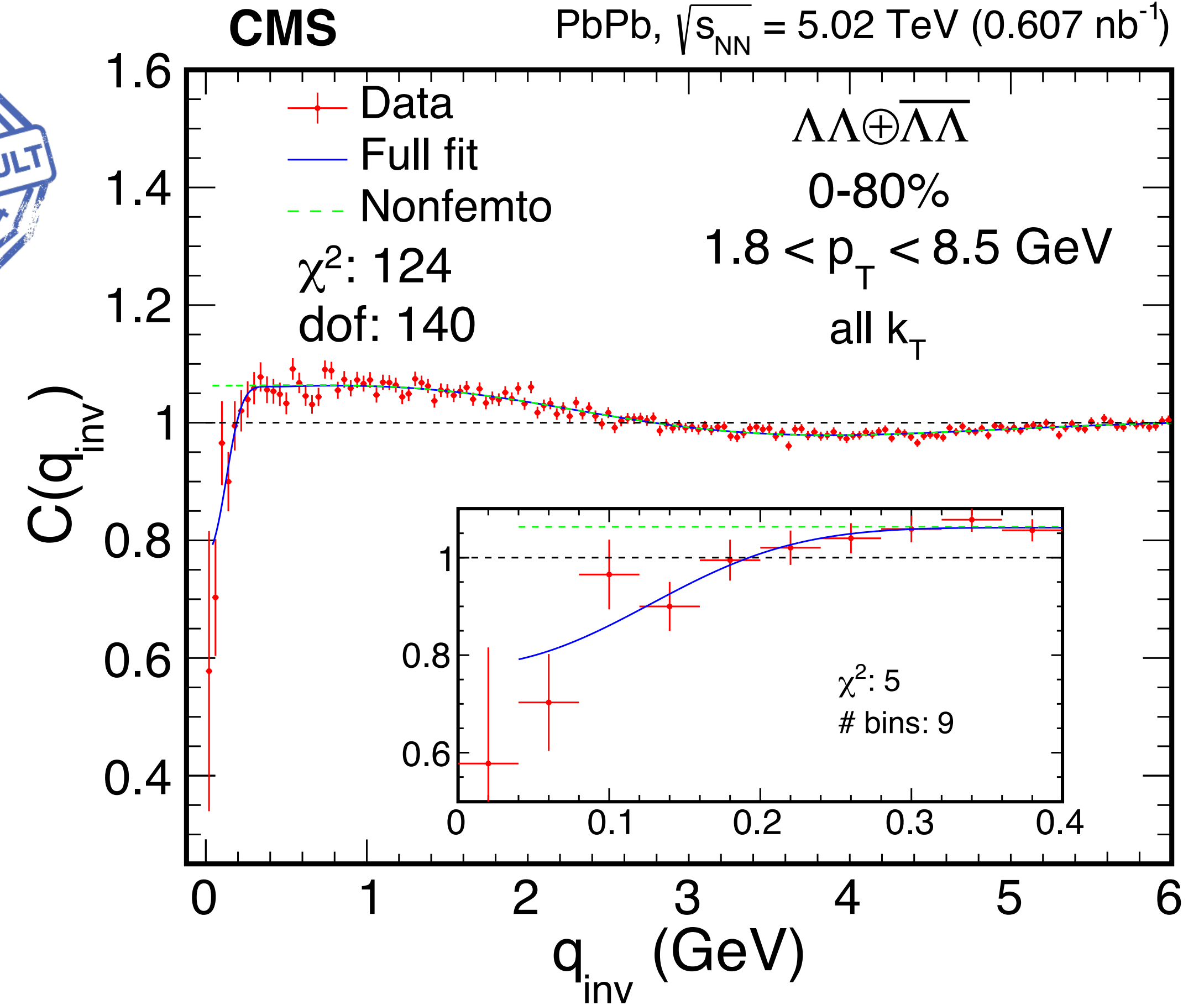
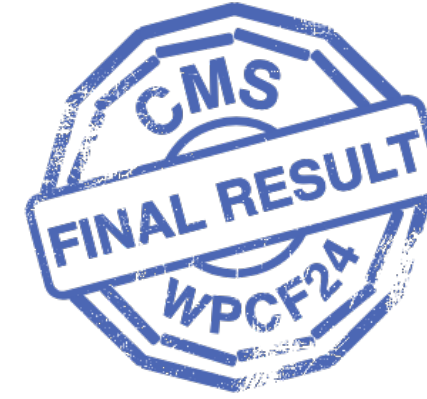
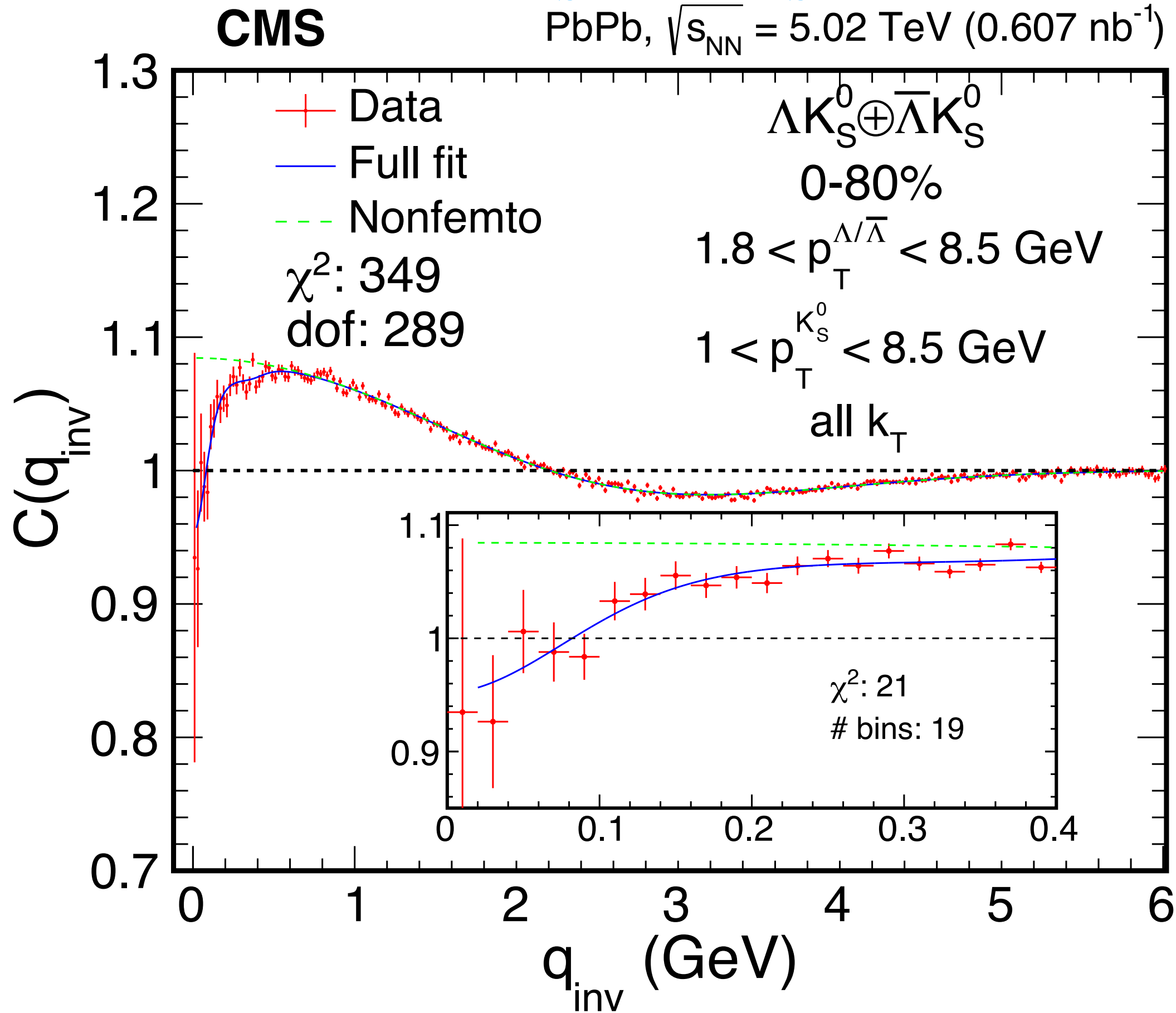
• Strong interaction parameters were fixed from low energy scattering

# Results: correlation and fitting

$$\Lambda K_S^0 \oplus \bar{\Lambda} K_S^0$$

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$$\Lambda\Lambda \oplus \bar{\Lambda}\bar{\Lambda}$$



QS + strong FSI [ non-identical ]

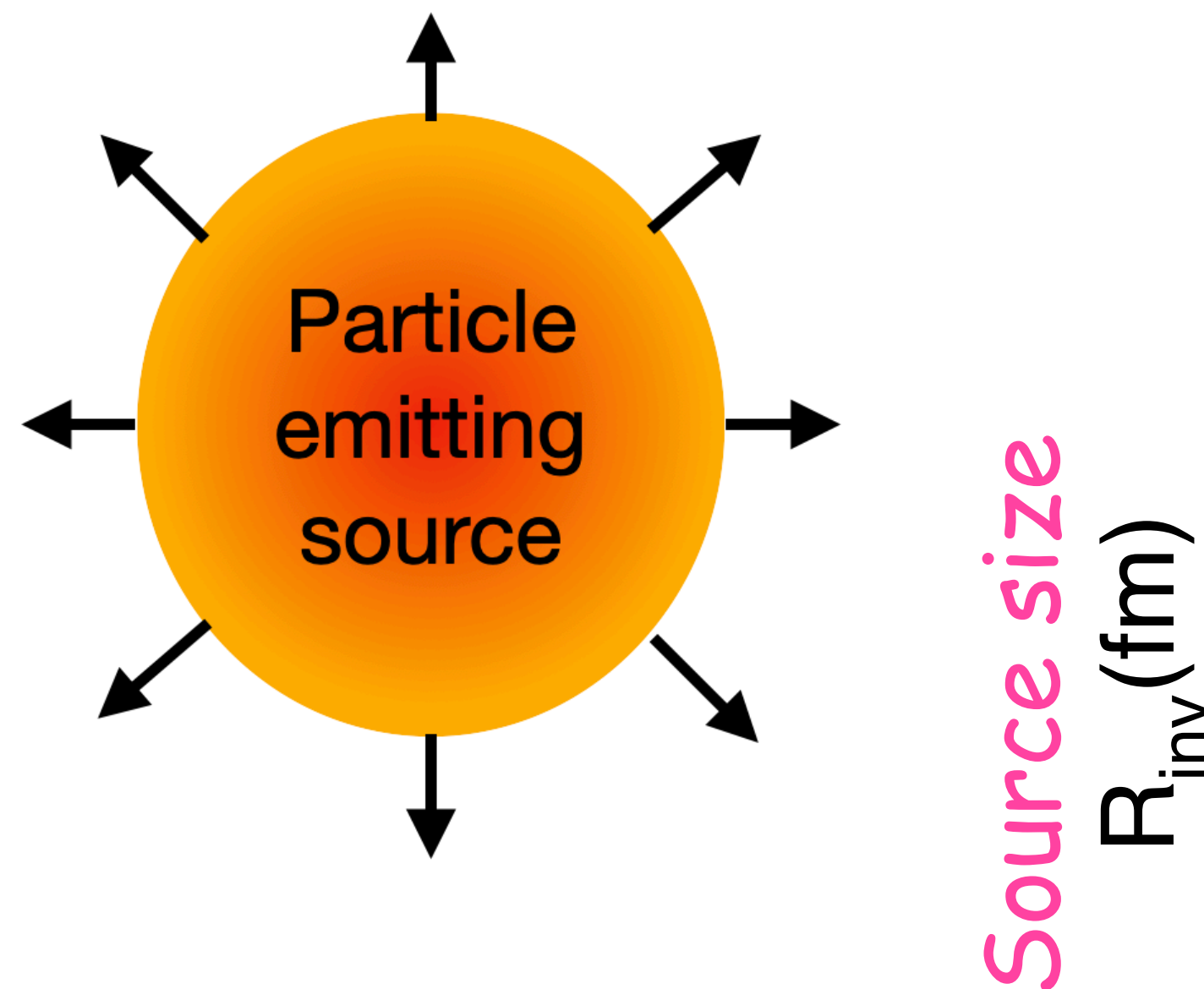
QS (Fermi-Dirac) + strong FSI

- Different pairs have different shape depending on their correlation features.



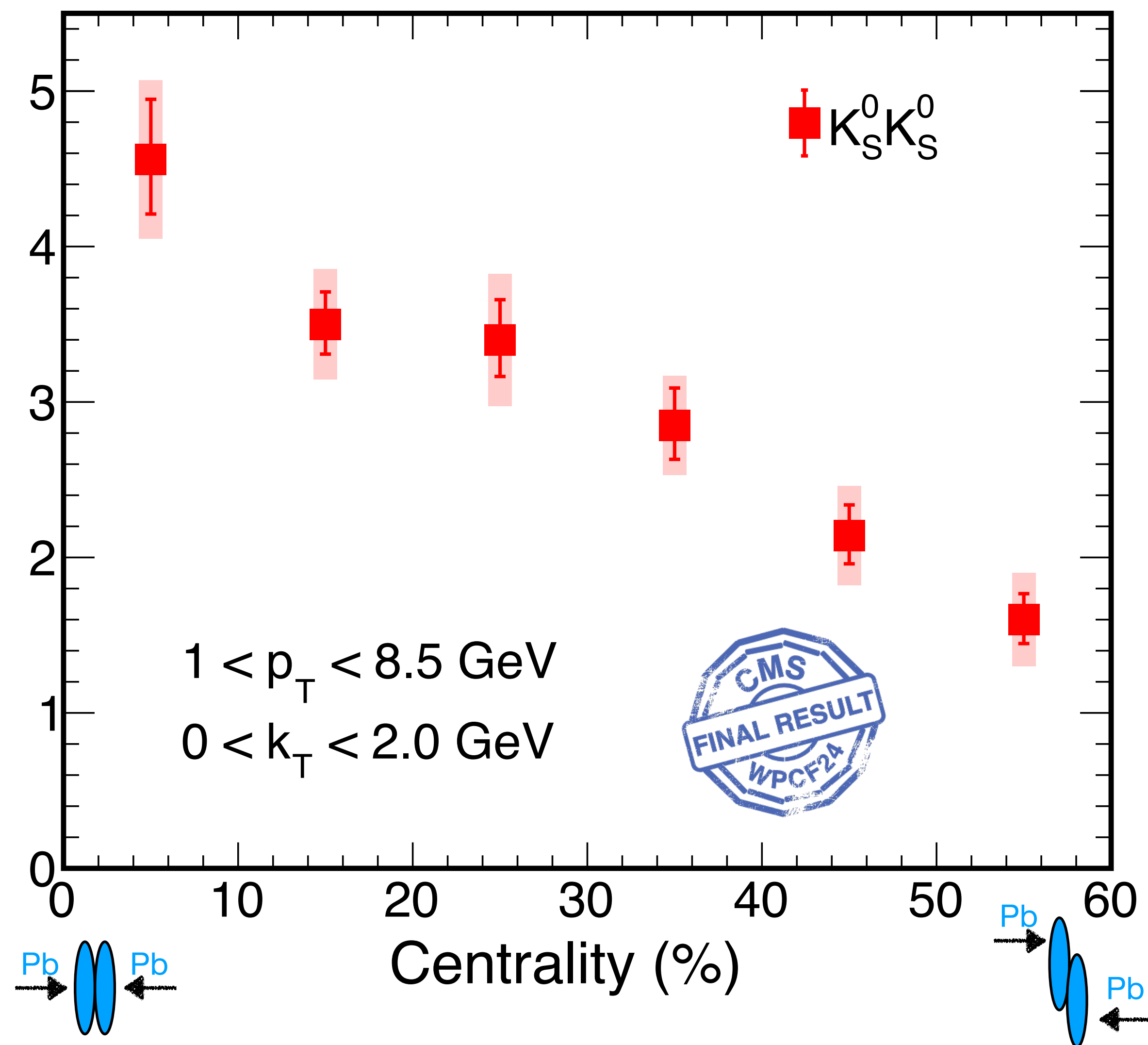
# Results: source size from $K_S^0 K_S^0$

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CMS

PbPb,  $\sqrt{s_{NN}} = 5.02$  TeV ( $0.607$  nb $^{-1}$ )



- Source size ( $R_{inv}$ ) decreases from central to peripheral collisions
  - expected from a simple geometric picture

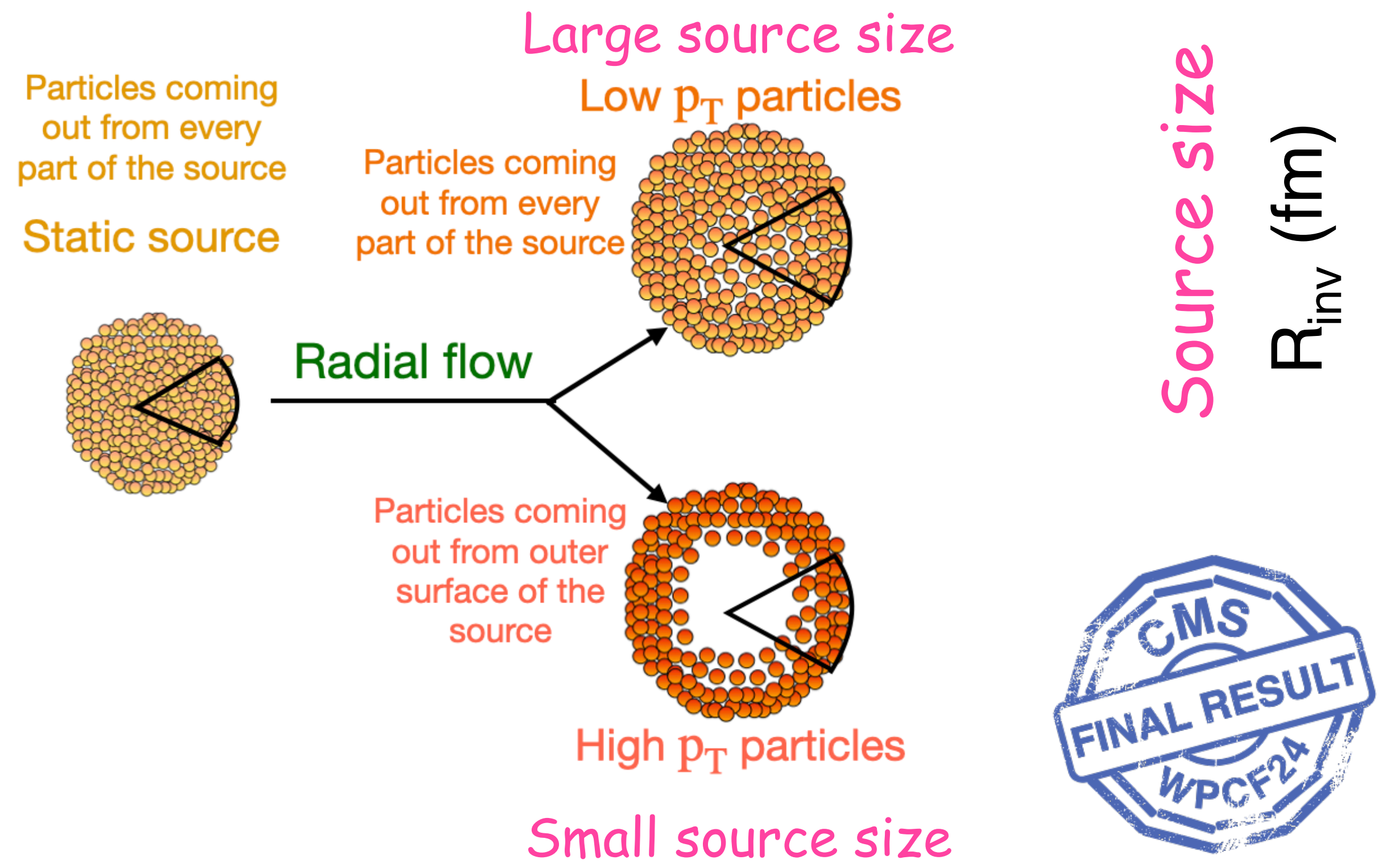
# Results: comparison with ALICE



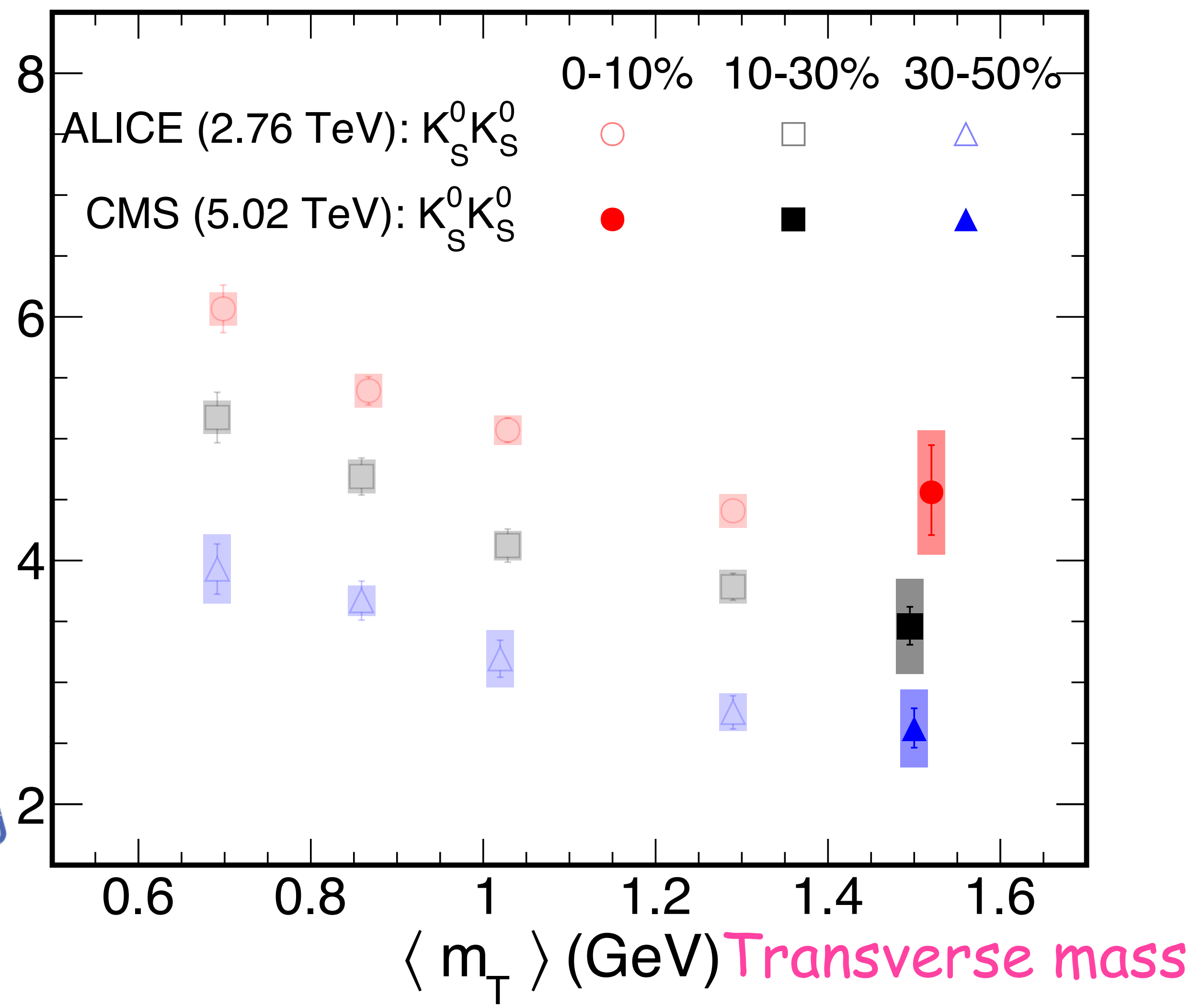
PLB 857 (2024) 138936

## $K_S^0 K_S^0$ source size comparison with ALICE

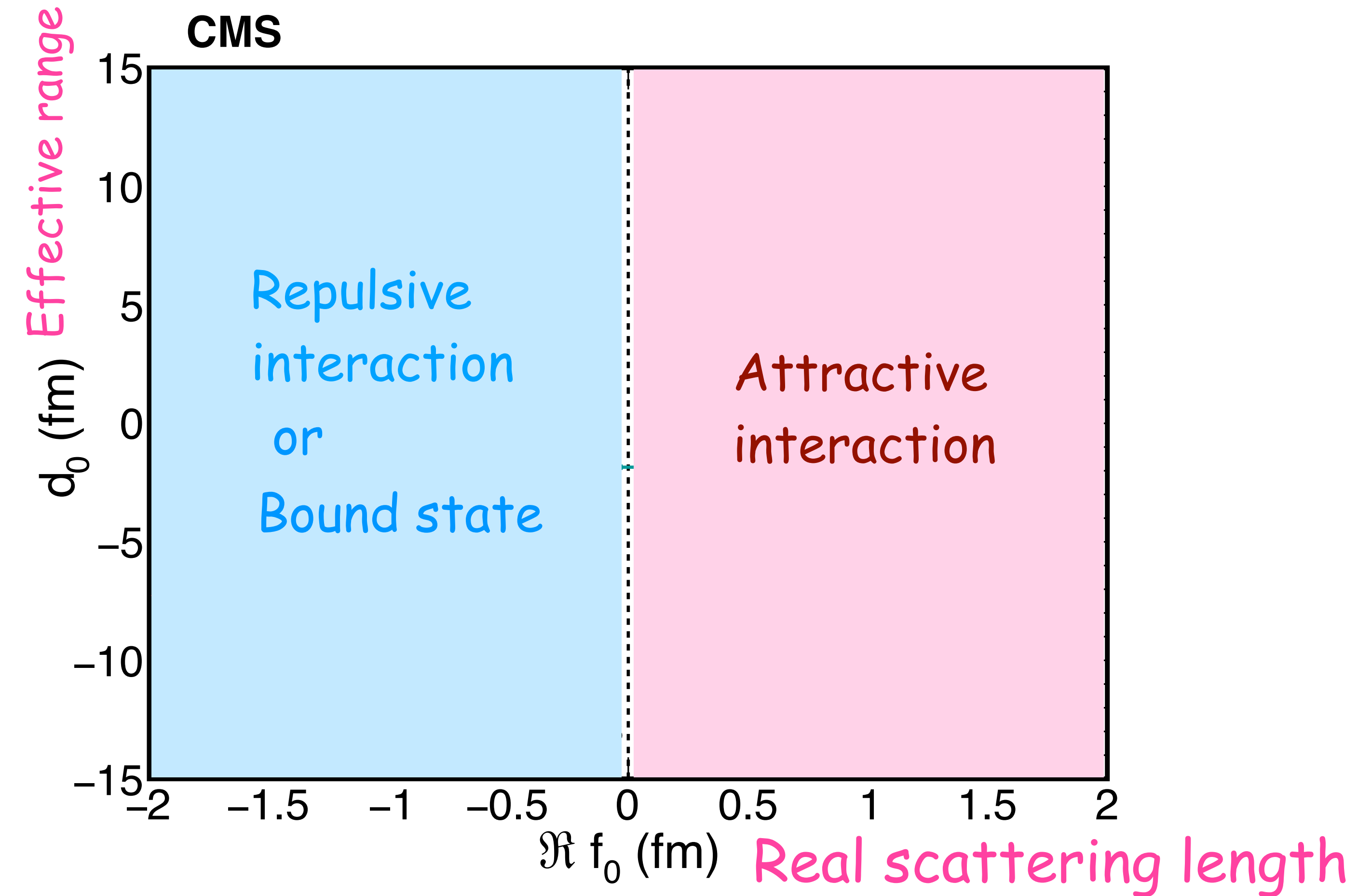
- Source size is decreasing with increasing  $\langle m_T \rangle$
- Following the trend measured by ALICE



CMS

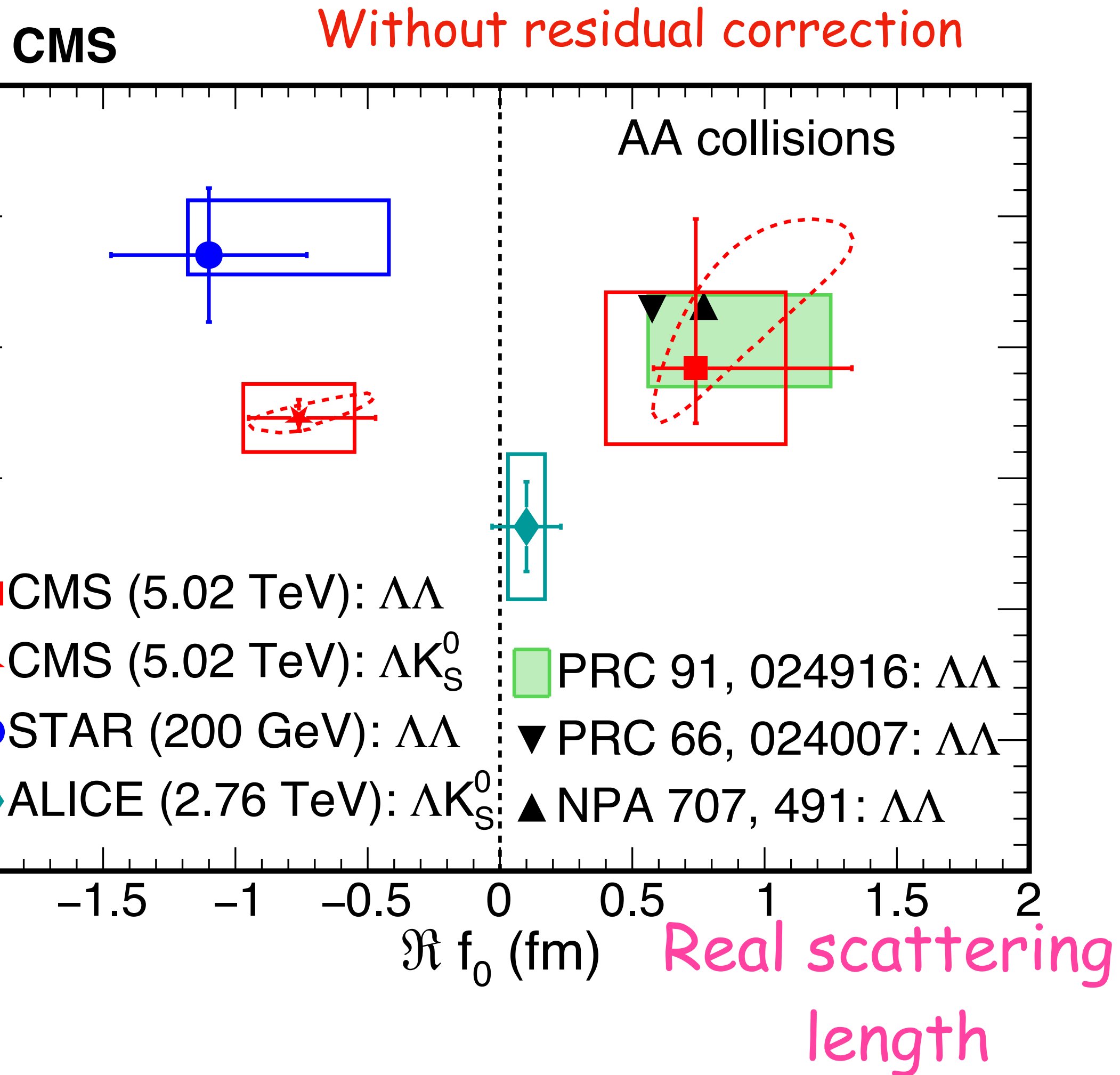
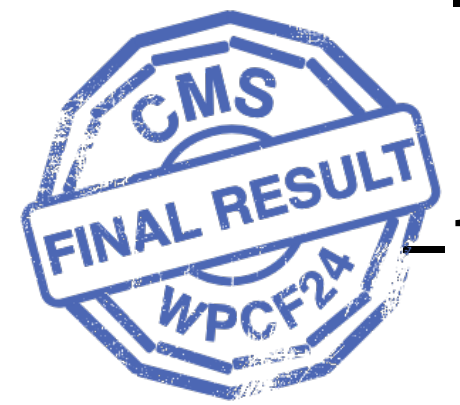
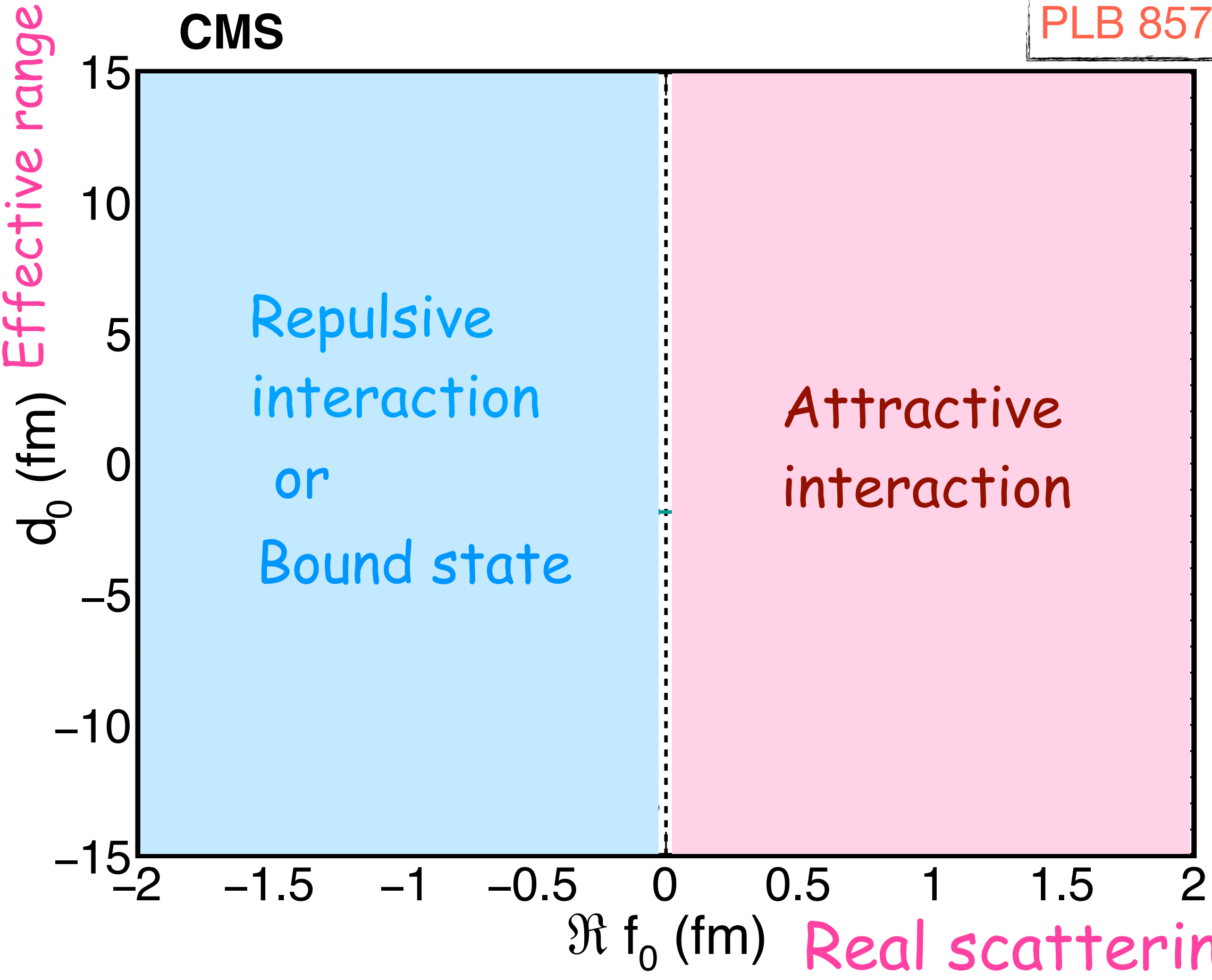


# Results: scattering parameters



# Results: scattering parameters

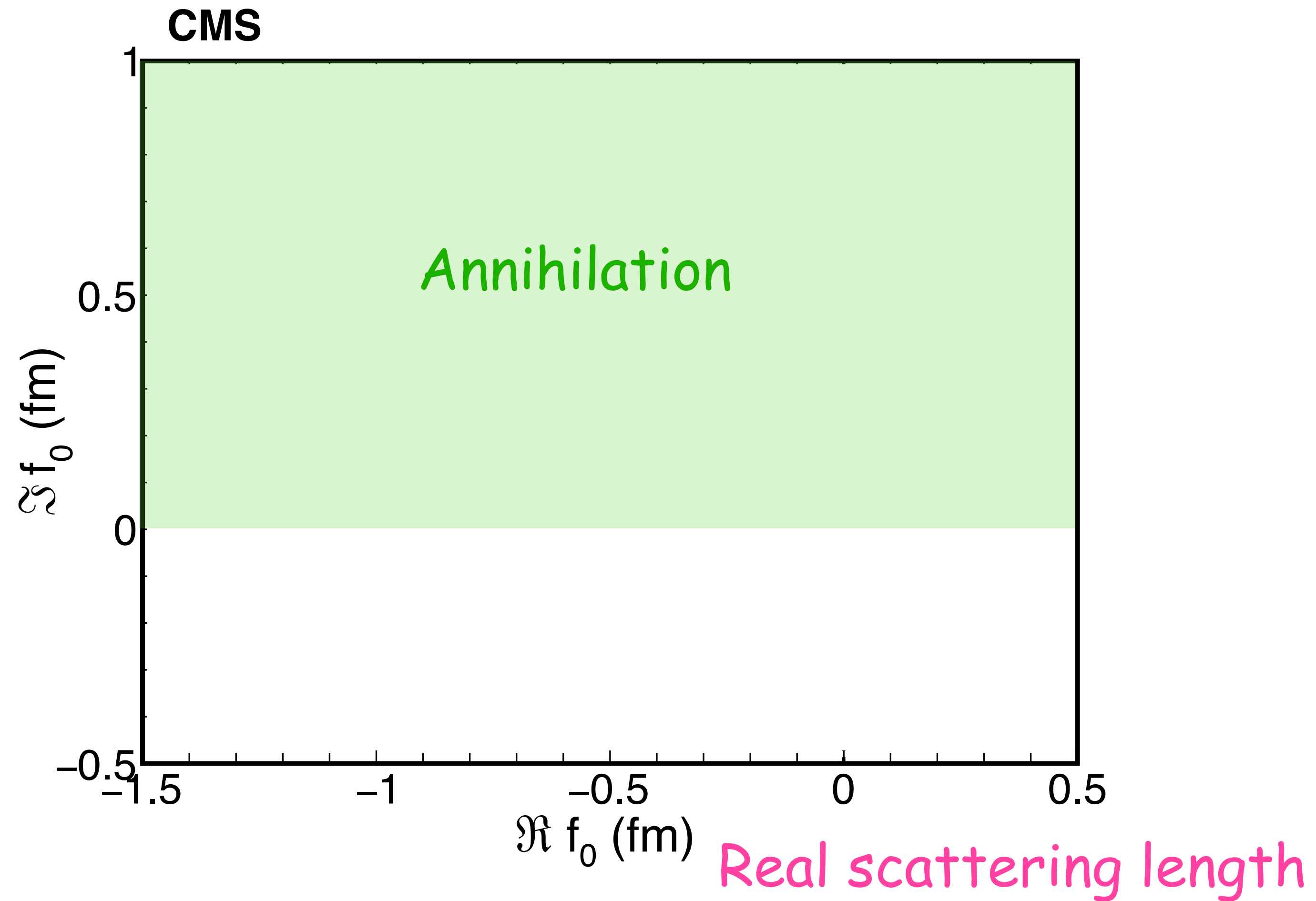
PLB 857 (2024) 138936



★  $\Re f_0 < 0$  for  $\Lambda K_S^0 \oplus \bar{\Lambda} K_S^0$  → Repulsive interaction

# Results: scattering parameters

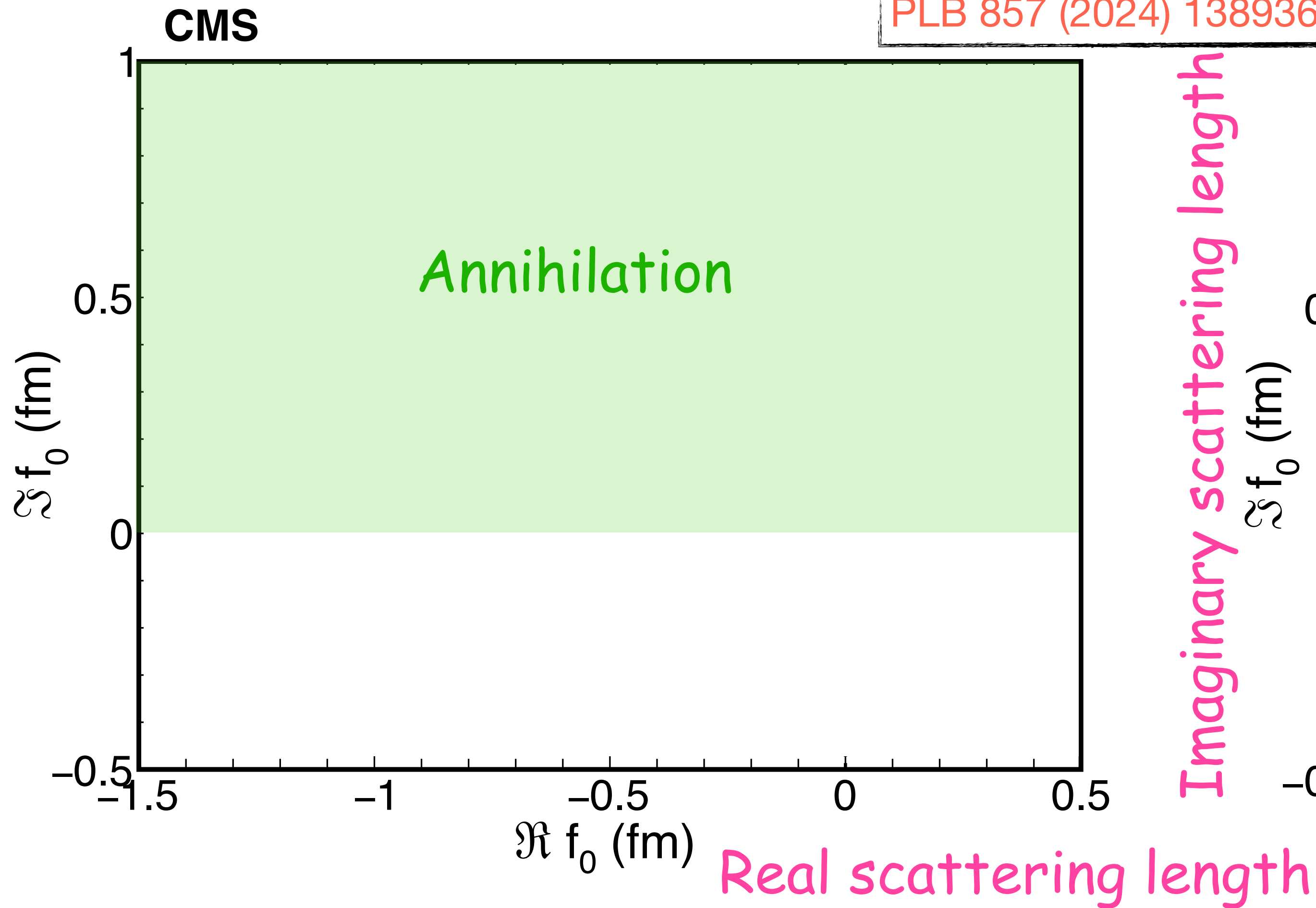
Imaginary scattering length



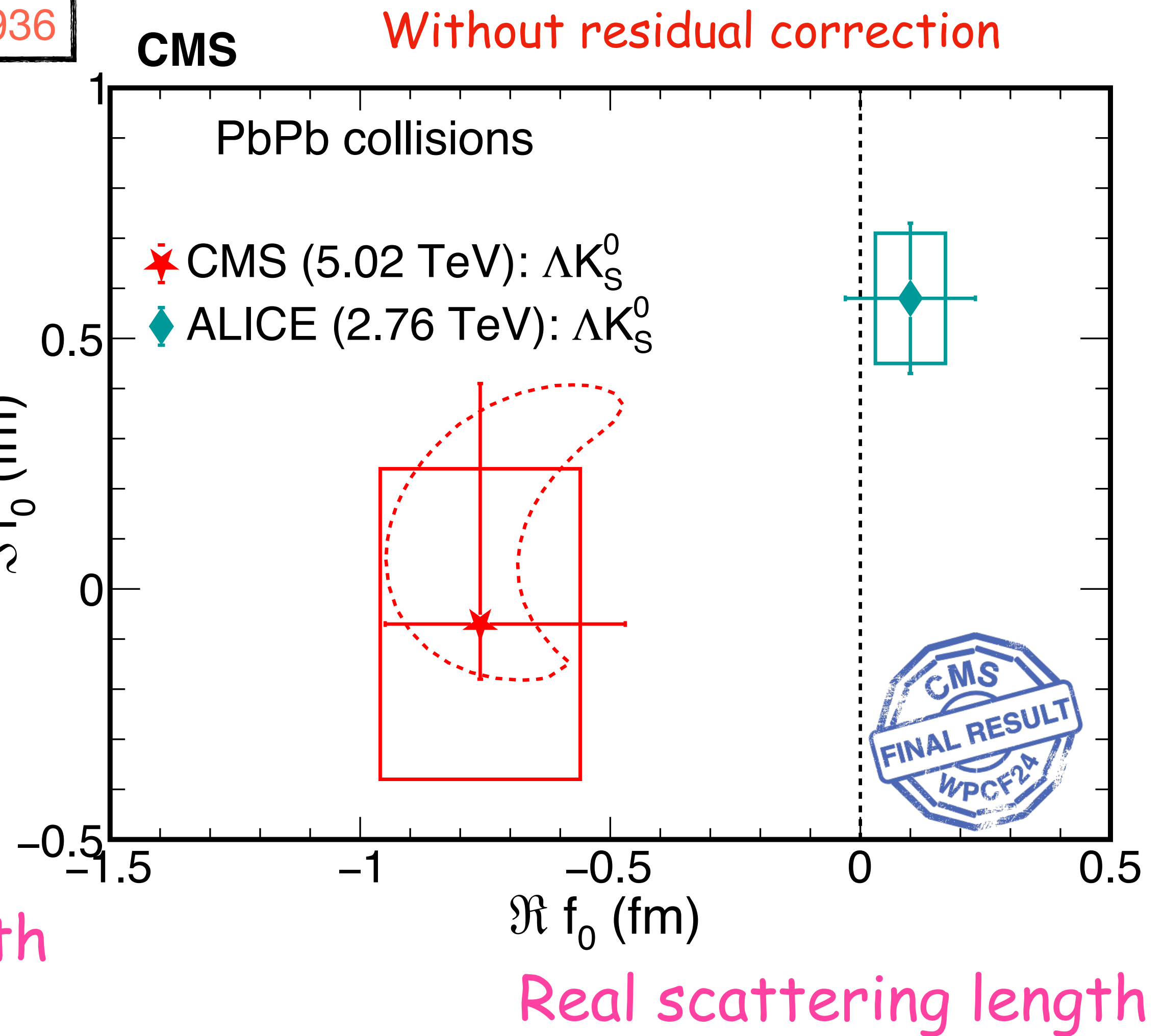
# Results: scattering parameters

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Imaginary scattering length



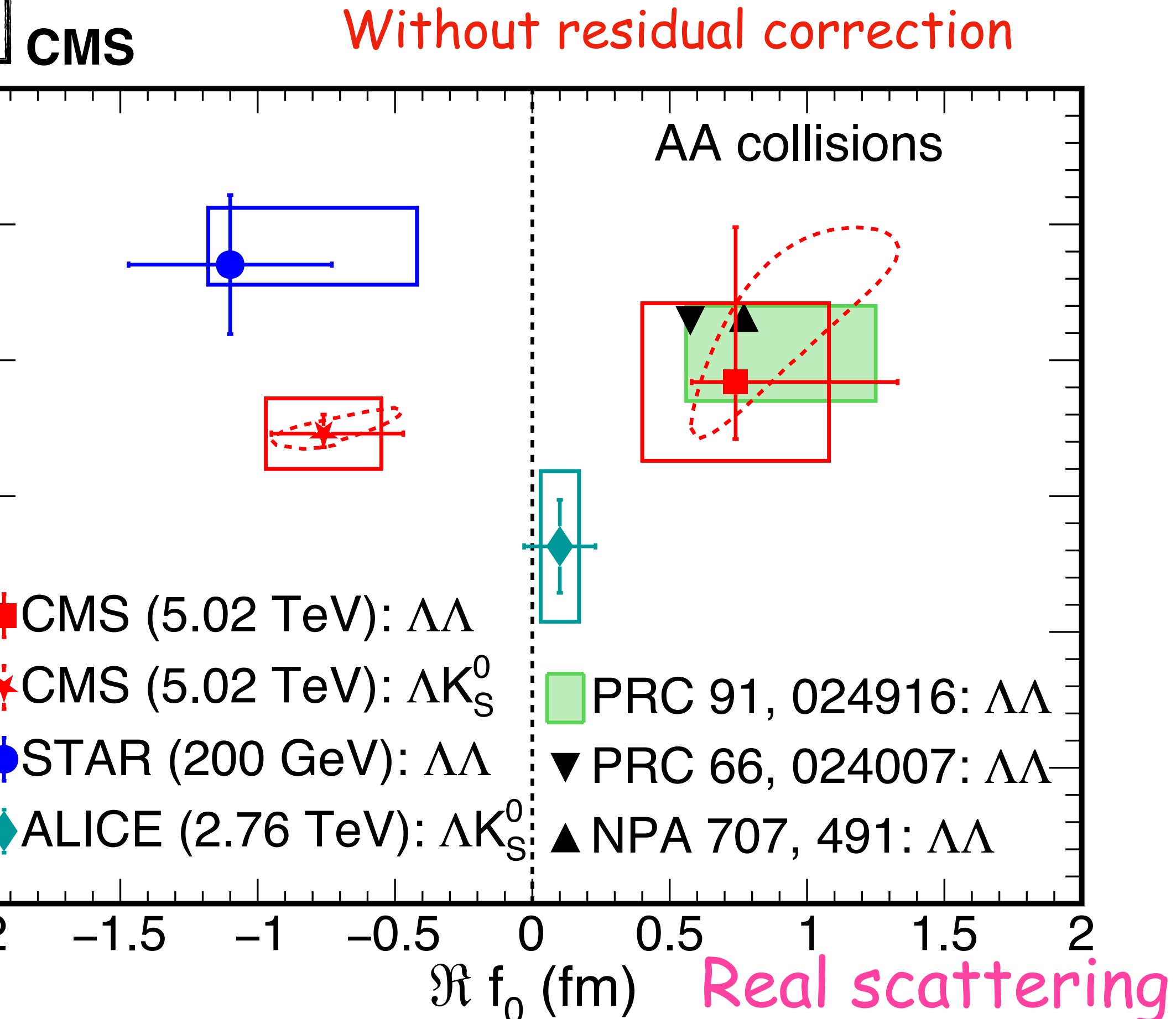
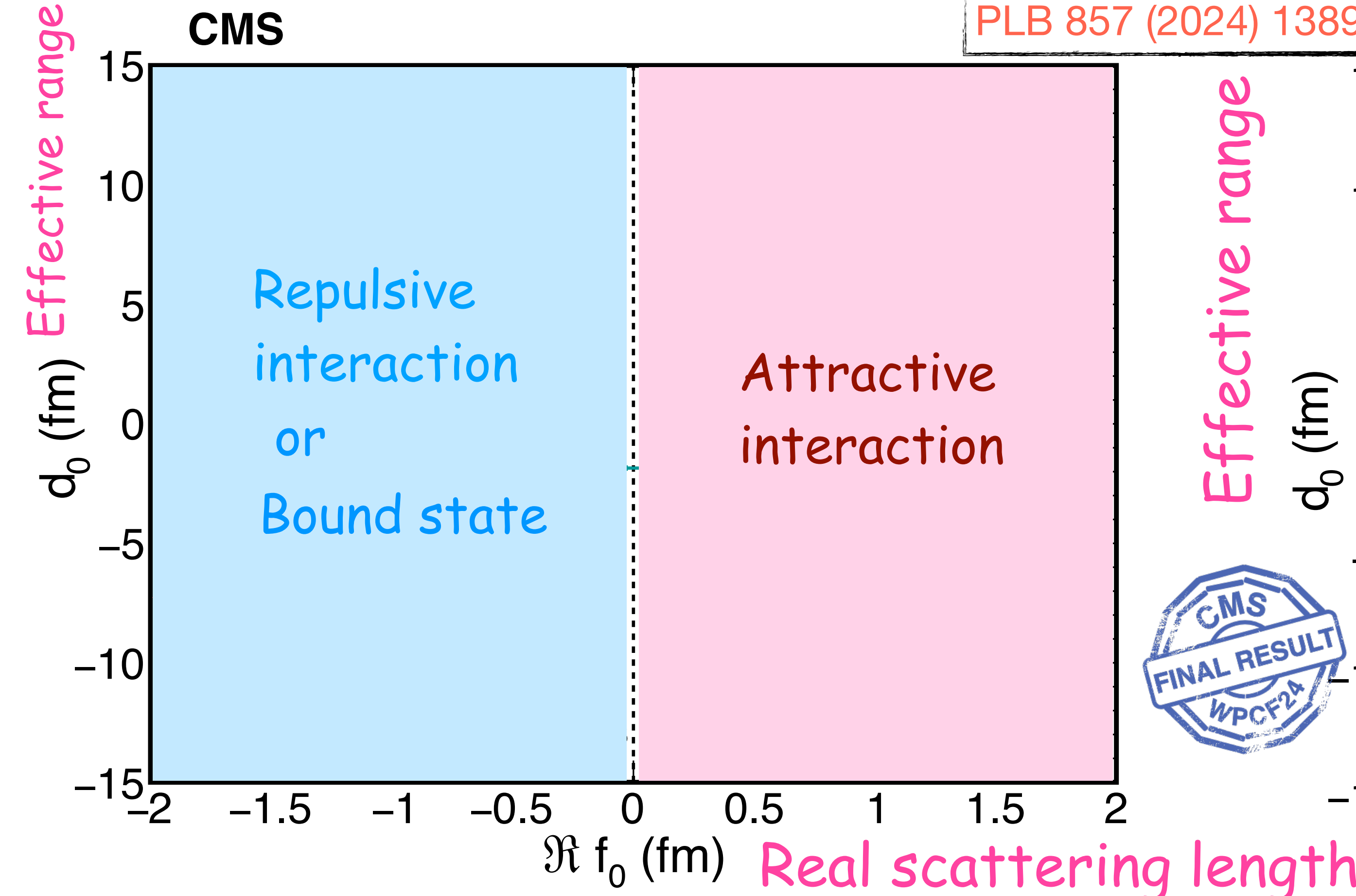
Imaginary scattering length



★  $\Im f_0$  for  $\Delta K_S^0 \oplus \bar{\Delta K_S^0}$  consistent with zero within error bar

# Results: scattering parameters

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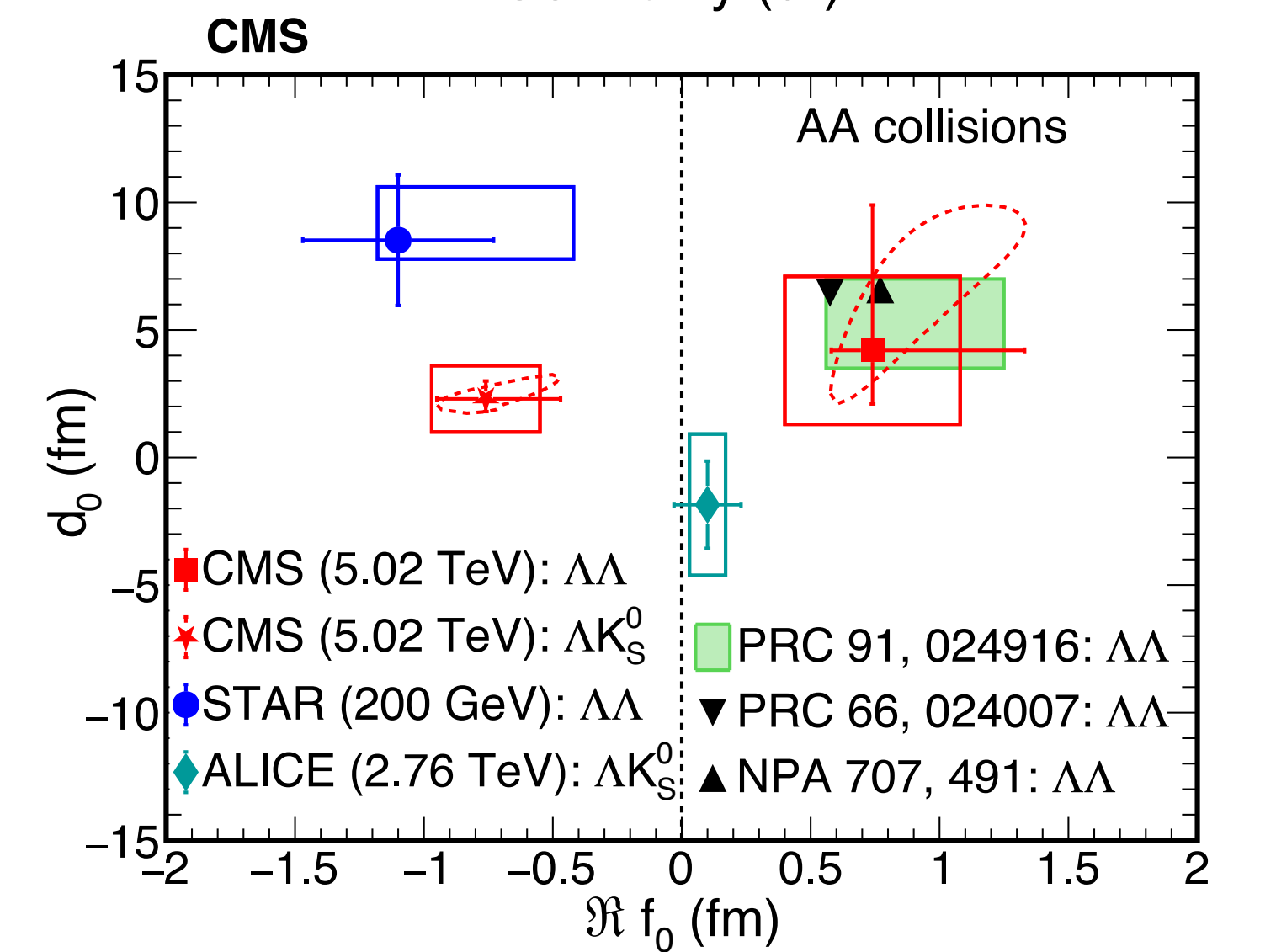
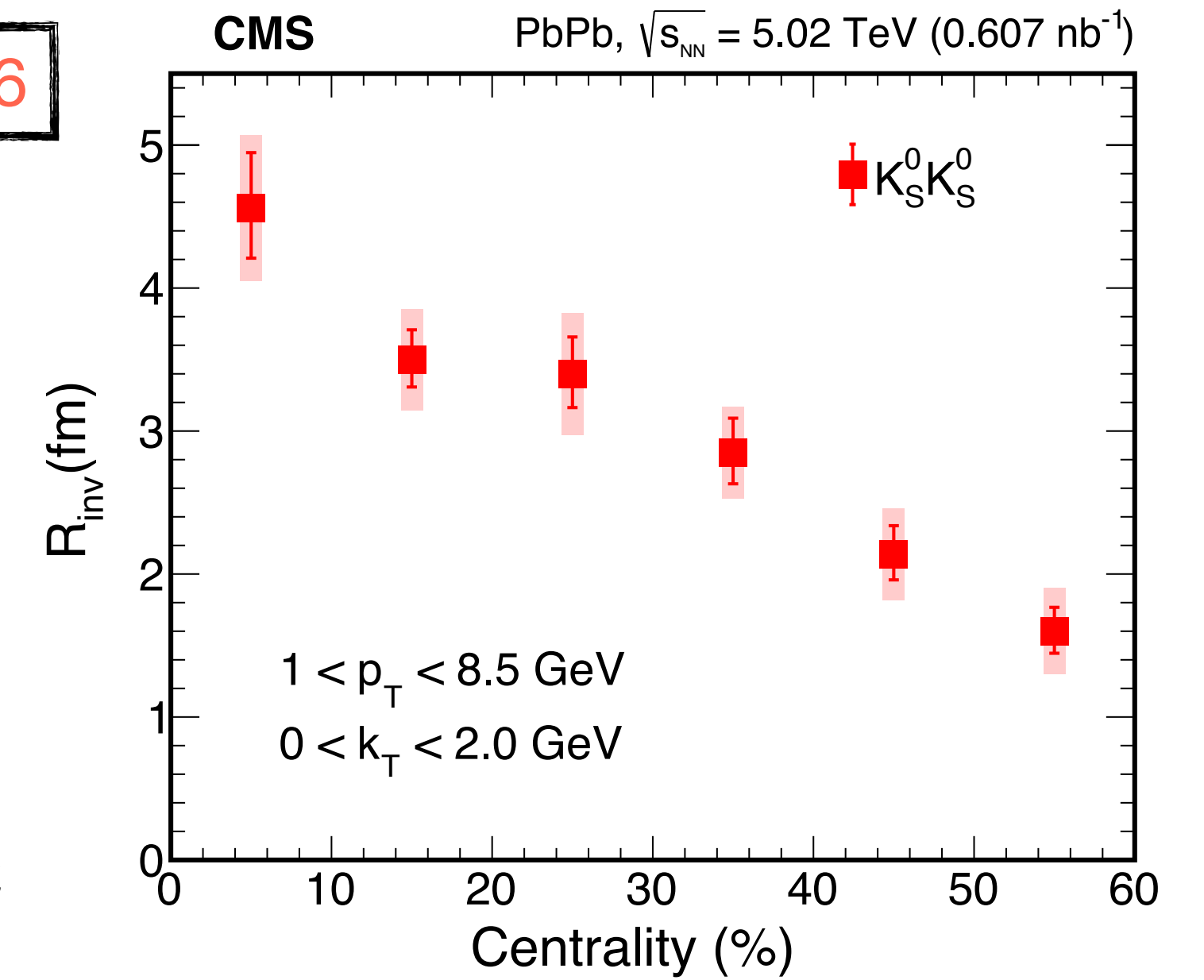


■  $\Re f_0 > 0$  for  $\Lambda\Lambda \oplus \bar{\Lambda}\bar{\Lambda}$  **Attractive interaction**

Indicating no bound state between two  $\Lambda(\bar{\Lambda})$

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- Source size is extracted from  $K_S^0 K_S^0$  correlation and it increases from peripheral to central collisions as expected.
- First measurement of  $\Lambda\Lambda \oplus \bar{\Lambda}\bar{\Lambda}$  correlation in PbPb collisions at LHC
  - $\Lambda\Lambda \oplus \bar{\Lambda}\bar{\Lambda}$  interaction : Attractive
    - Indicating non-existence of bound H-dibaryon of two  $\Lambda(\bar{\Lambda})$
- $\Lambda K_S^0 \oplus \bar{\Lambda} K_S^0$  interaction : Repulsive





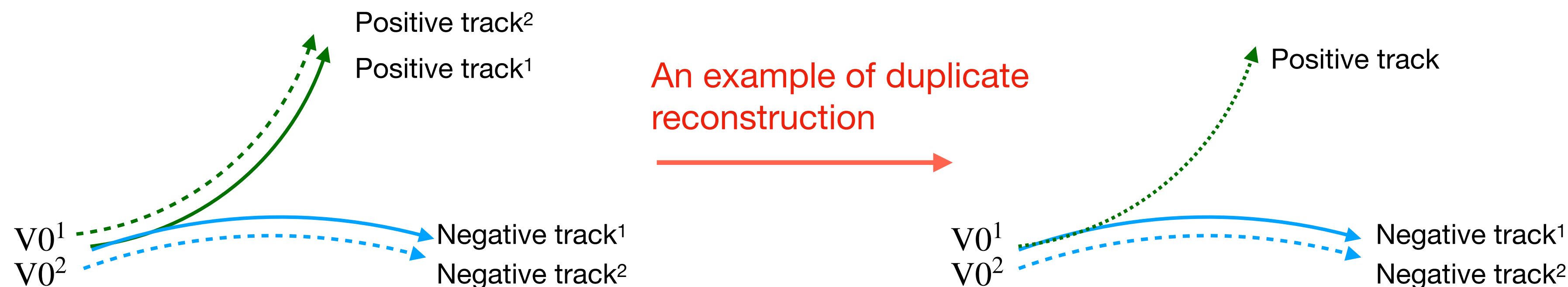
Thank you for your kind attention !



# Backup

# Duplicate V0 removal

- Removed duplicate V0 (sharing common daughters):
- If  $|\Delta\chi^2/ndf| = 0$  between V0 daughters with same charge, remove one V0 randomly



# Correction to the correlation

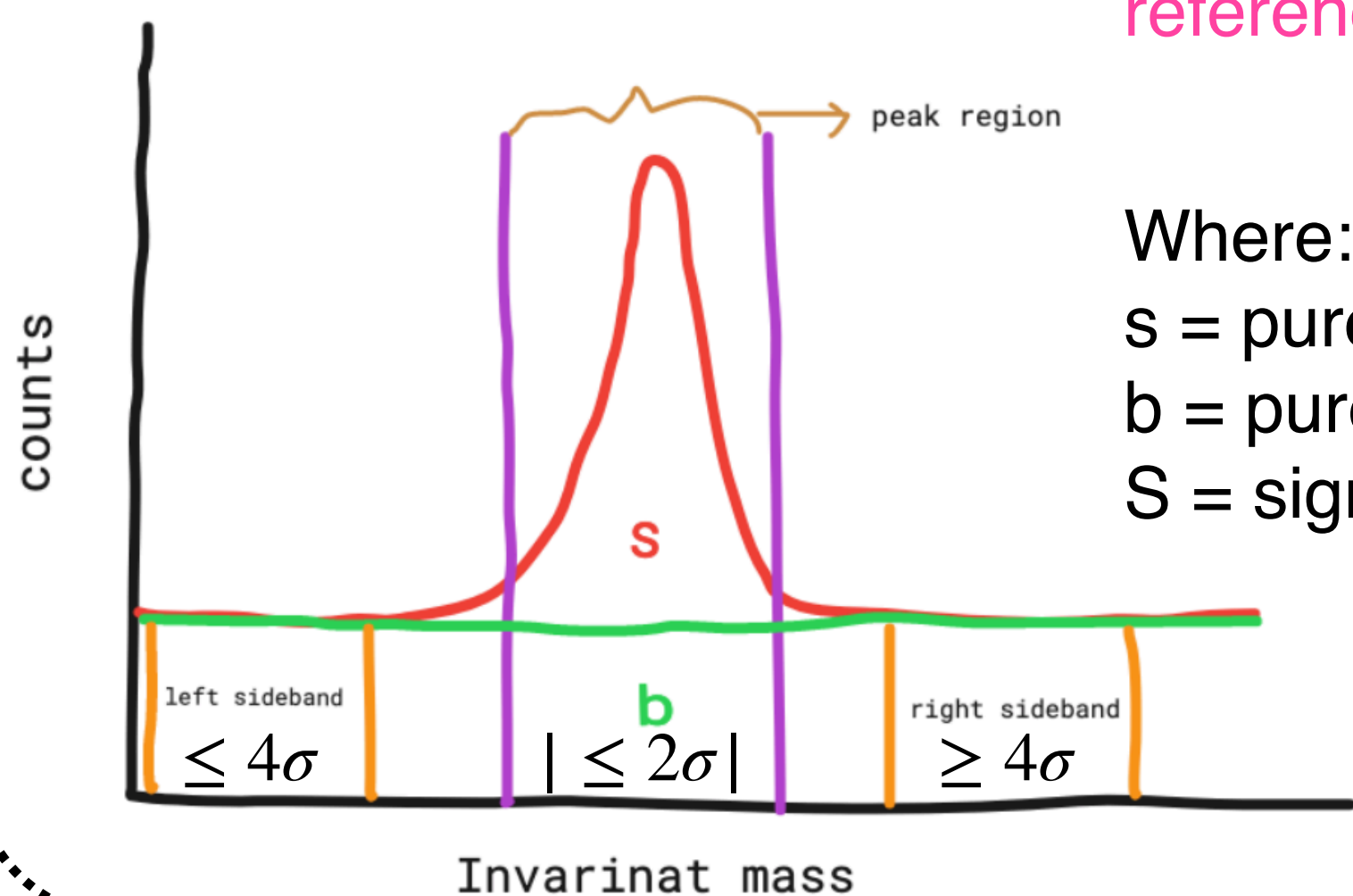
## Pair purity

$$A^{\text{measured}}(q_{\text{inv}}) = f^{\text{ss}} A^{\text{ss}}(q_{\text{inv}}) + f^{\text{bb}} A^{\text{bb}}(q_{\text{inv}}) + f^{\text{Sb}} A^{\text{Sb}}(q_{\text{inv}})$$

$$A^{\text{ss}}(q_{\text{inv}}) = [A^{\text{measured}}(q_{\text{inv}}) - f^{\text{bb}} A^{\text{bb}}(q_{\text{inv}}) - f^{\text{Sb}} A^{\text{Sb}}(q_{\text{inv}})] / f^{\text{ss}}$$

$$f^{\text{ss}} = \frac{\binom{s}{2}}{\binom{s+b}{2}}, \quad f^{\text{bb}} = \frac{\binom{b}{2}}{\binom{s+b}{2}}, \quad f^{\text{Sb}} = 1 - f^{\text{ss}} - f^{\text{bb}}$$

Applied on signal and reference samples



Where:  
 s = pure signal,  
 b = pure background  
 S = signal + background (s+b)

## Non-femtoscopic background

$$\Omega(q_{\text{inv}}) = \mathcal{N}(1 + \alpha_1 e^{-(q_{\text{inv}} R_1)^2})(1 - \alpha_2 e^{-(q_{\text{inv}} R_2)^2})$$

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Total correlation function will be:

$$C_{\text{Fit}}(q_{\text{inv}}) = \Omega(q_{\text{inv}}) \times C'_{\text{Fit}}(q_{\text{inv}}) \rightarrow \text{Theoretical fitted function}$$

Sov. J. Nucl. Phys. 35 (1982) 770.

- All the non-femto parameters,  $\mathcal{N}$ ,  $\alpha_1$ ,  $\alpha_2$ ,  $R_1$ , and  $R_2$ , were treated as free parameters during fitting

# Fitting function: Lednicky model

$$K_S^0 K_S^0$$



$$C'_{\text{Fit}}(q_{\text{inv}}) = N \left[ 1 + \lambda \left( \exp(-q_{\text{inv}}^2 R_{\text{inv}}^2) + \frac{1}{2} \left| \frac{f(q_{\text{inv}}/2)}{R_{\text{inv}}} \right|^2 + \frac{2\Re f(q_{\text{inv}}/2)}{\sqrt{\pi} R_{\text{inv}}} F_1(q_{\text{inv}} R_{\text{inv}}) - \frac{\Im f(q_{\text{inv}}/2)}{R_{\text{inv}}} F_2(Q_{\text{inv}} R_{\text{inv}}) \right) \right]$$

$R_{\text{inv}}$ ,  $\lambda$  and  $N$  are the free parameters

$$F_1(q_{\text{inv}} R_{\text{inv}}) = \int_0^{q_{\text{inv}} R_{\text{inv}}} dx \frac{e^{x^2 - q_{\text{inv}}^2 R_{\text{inv}}^2}}{q_{\text{inv}} R_{\text{inv}}}$$

$$F_2(q_{\text{inv}} R_{\text{inv}}) = \frac{1 - e^{-q_{\text{inv}}^2 R_{\text{inv}}^2}}{q_{\text{inv}} R_{\text{inv}}}$$

$$f(q_{\text{inv}}/2) = \frac{f_{f_0} + f_{a_0}}{2}$$

$$f_{f_0, a_0}(q_{\text{inv}}/2) = \gamma_{f_0, a_0} / [m_{f_0, a_0}^2 - s - i\gamma_{f_0, a_0} q_{\text{inv}}/2 - i\gamma'_{f_0, a_0} k'_{f_0, a_0}]$$

$$QS$$

$$FSI$$

$$\Lambda K_S^0 \oplus \bar{\Lambda} K_S^0$$



$$C'_{\text{Fit}}(q_{\text{inv}}) = N \left[ 1 + \lambda \left( \frac{1}{2} \left| \frac{f(q_{\text{inv}}/2)}{R_{\text{inv}}} \right|^2 \left( 1 - \frac{d_0}{2\sqrt{\pi} R_{\text{inv}}} \right) + \frac{2\Re f(q_{\text{inv}}/2)}{\sqrt{\pi} R_{\text{inv}}} F_1(q_{\text{inv}} R_{\text{inv}}) - \frac{\Im f(q_{\text{inv}}/2)}{R_{\text{inv}}} F_2(Q_{\text{inv}} R_{\text{inv}}) \right) \right]$$

$R_{\text{inv}}$ ,  $d_0$ ,  $\Re f_0$ ,  $\Im f_0$ ,  $\lambda$  and  $N$  are the free parameters

$$f(q_{\text{inv}}/2) = \left( \frac{1}{f_0} + \frac{1}{8} d_0 q_{\text{inv}}^2 - i \frac{q_{\text{inv}}}{2} \right)^{-1}$$

$$\Lambda \Lambda \oplus \bar{\Lambda} \bar{\Lambda}$$



$$C'_{\text{Fit}}(q_{\text{inv}}) = N \left[ 1 + \lambda \left( -\frac{1}{2} \exp(-q_{\text{inv}}^2 R_{\text{inv}}^2) + \frac{1}{4} \left| \frac{f(q_{\text{inv}}/2)}{R_{\text{inv}}} \right|^2 \left( 1 - \frac{d_0}{2\sqrt{\pi} R_{\text{inv}}} \right) + \frac{\Re f(q_{\text{inv}}/2)}{\sqrt{\pi} R_{\text{inv}}} F_1(q_{\text{inv}} R_{\text{inv}}) - \frac{\Im f(q_{\text{inv}}/2)}{2R_{\text{inv}}} F_2(Q_{\text{inv}} R_{\text{inv}}) \right) \right]$$

$R_{\text{inv}}$ ,  $d_0$ ,  $\Re f_0$ ,  $\lambda$  and  $N$  are the free parameters

$$f(q_{\text{inv}}/2) = \left( \frac{1}{f_0} + \frac{1}{8} d_0 q_{\text{inv}}^2 - i \frac{q_{\text{inv}}}{2} \right)^{-1}$$

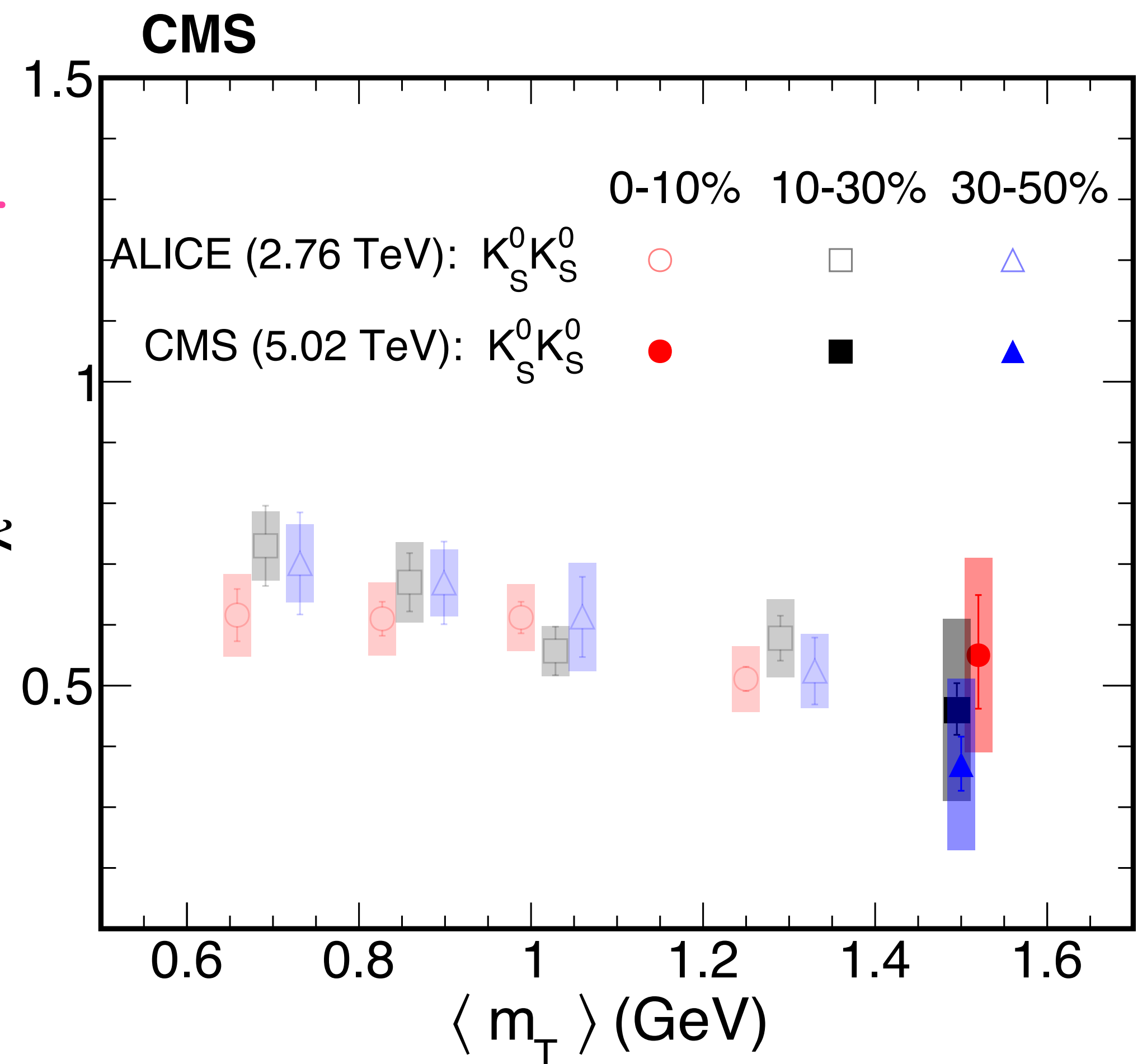
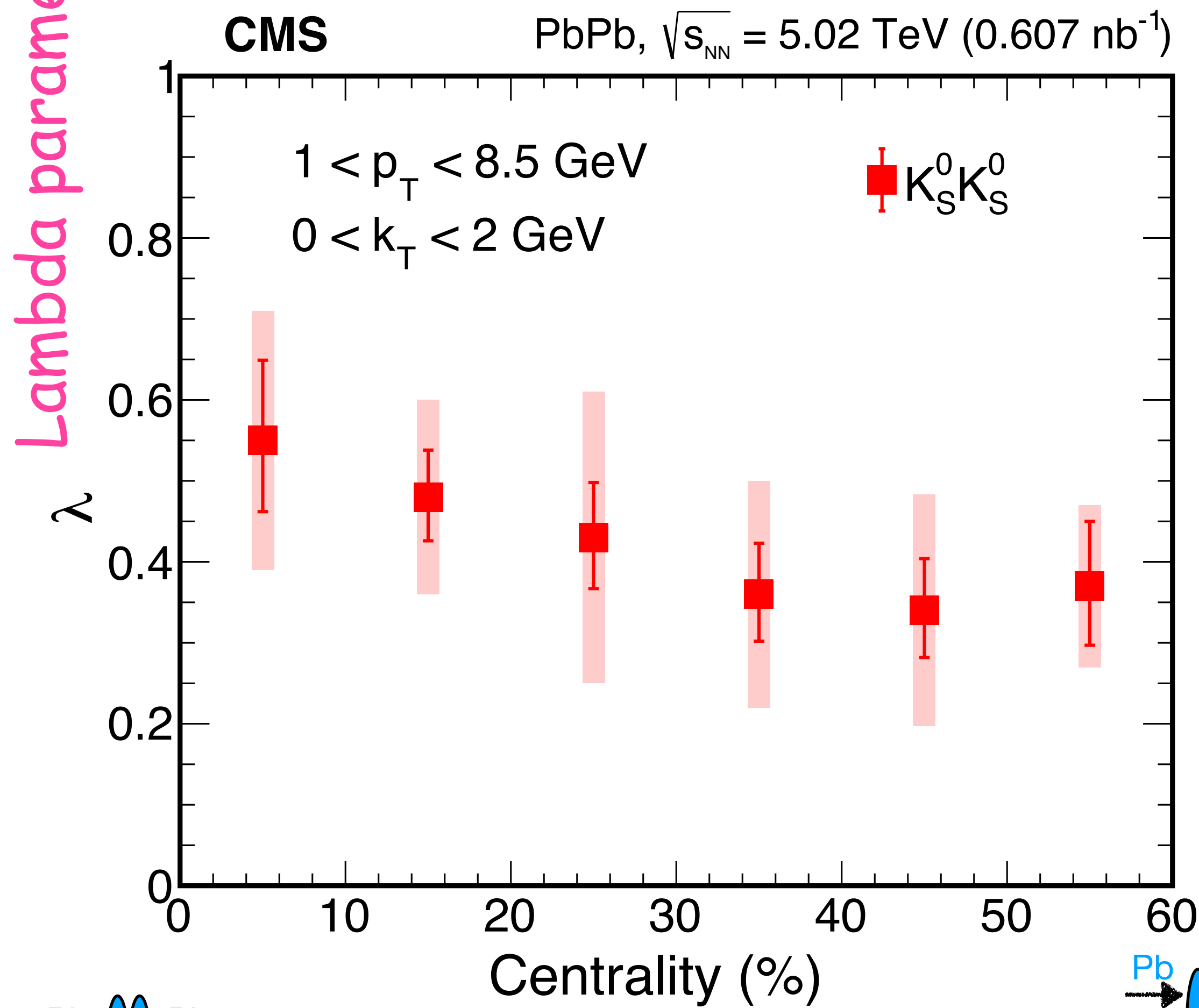
Sov. J. Nucl. Phys. 35 (1982) 770.

# Lambda parameter

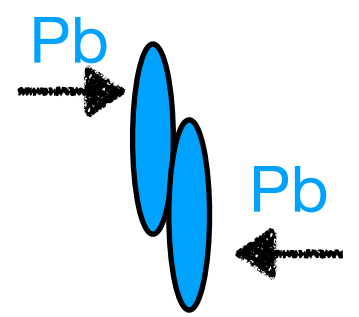
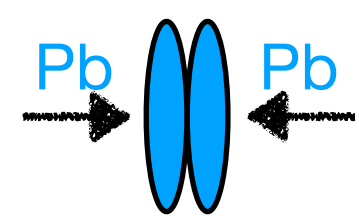
PLB 857 (2024) 138936

Lambda parameter

Lambda parameter



Transverse mass

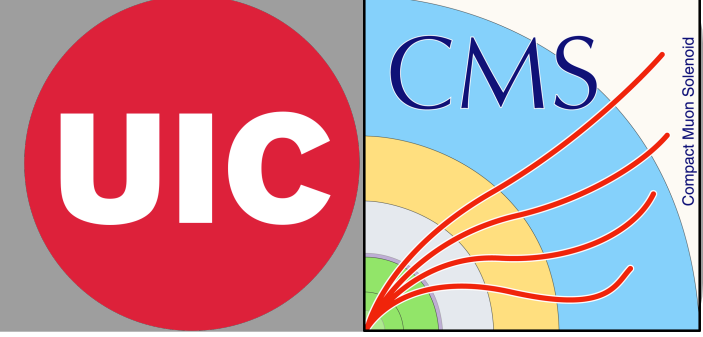


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Table 3: Extracted values of the  $R_{\text{inv}}$ ,  $\Re f_0$ ,  $\Im f_0$ ,  $d_0$ ,  $\lambda$ , and  $\langle m_T \rangle$  parameters from the  $K_S^0 K_S^0$ ,  $\Lambda K_S^0$ , and  $\Lambda\Lambda$  combinations in the 0–80% centrality range. The first and second uncertainties are statistical and systematic, respectively.

Parameter	$K_S^0 K_S^0$	$\Lambda K_S^0$	$\Lambda\Lambda$
$R_{\text{inv}}$ (fm)	$3.40 \pm 0.11 \pm 0.37$	$2.1_{-0.5}^{+1.4} \pm 0.8$	$1.3_{-0.2}^{+0.4} \pm 0.3$
$\Re f_0$ (fm)	—	$-0.76_{-0.19}^{+0.29} \pm 0.20$	$0.74_{-0.16}^{+0.59} \pm 0.33$
$\Im f_0$ (fm)	—	$-0.07_{-0.11}^{+0.48} \pm 0.32$	—
$d_0$ (fm)	—	$2.3_{-0.5}^{+0.7} \pm 1.3$	$4.2_{-2.1}^{+5.7} \pm 2.9$
$\lambda$	$0.43 \pm 0.03 \pm 0.13$	$0.34_{-0.12}^{+0.41} \pm 0.17$	$1.5_{-1.1}^{+1.2} \pm 1.4$
$\langle m_T \rangle$ (GeV)	1.50	2.09	2.60

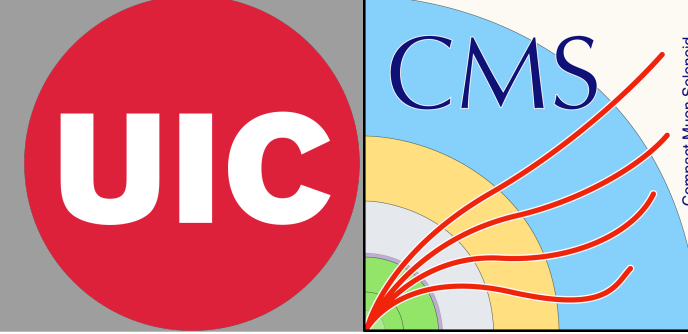
# Non-prompt fraction



- HYDJET: 85%  $\Lambda(\bar{\Lambda})$  produce directly and 15%  $\Lambda(\bar{\Lambda})$  from secondary decay
- HIJING: 39%  $\Lambda(\bar{\Lambda})$  produce directly and 61%  $\Lambda(\bar{\Lambda})$  from secondary decay

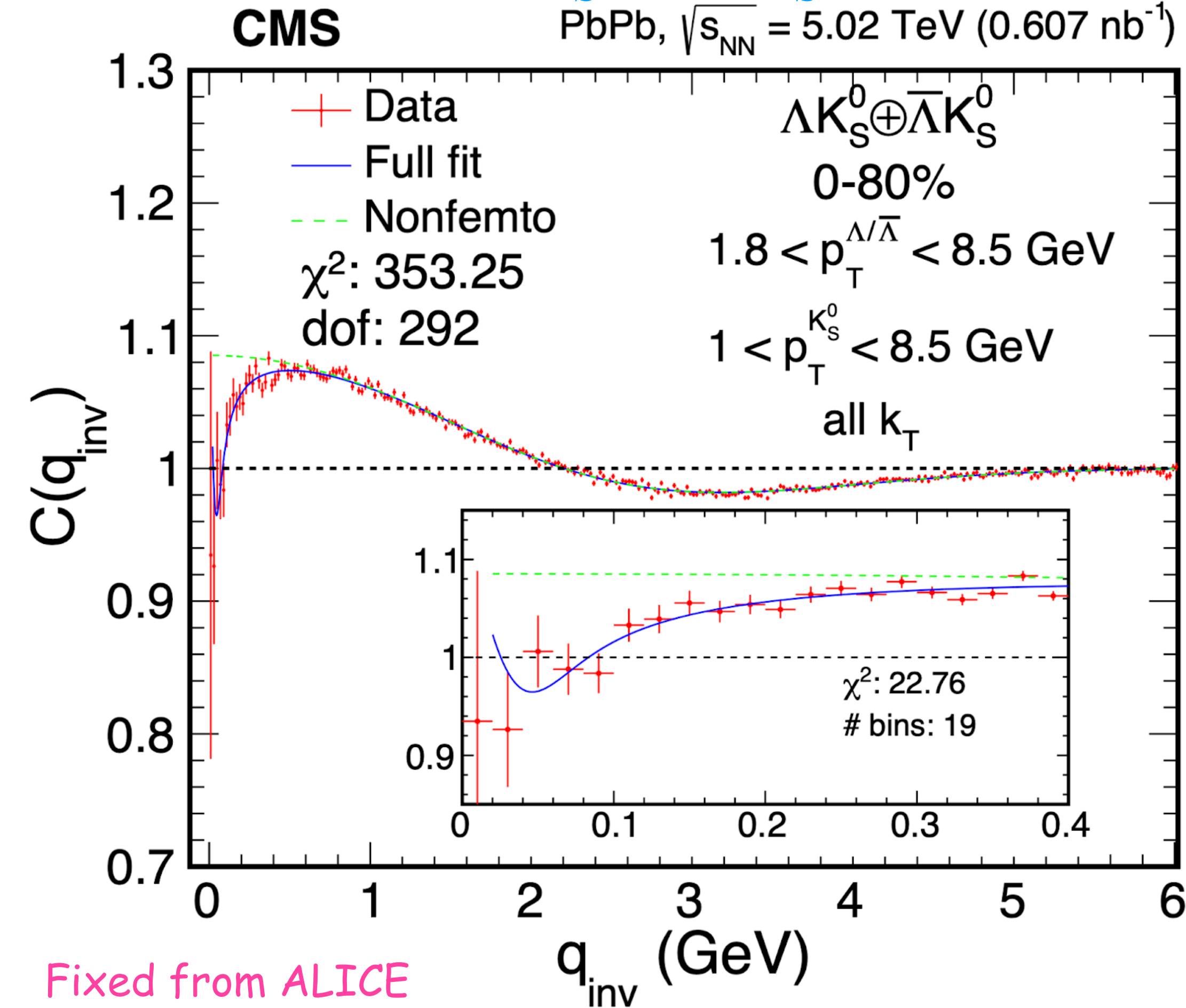


# Strong parameters fixed



$\Lambda K_S^0 \oplus \bar{\Lambda} K_S^0$

PbPb,  $\sqrt{s_{NN}} = 5.02$  TeV ( $0.607 \text{ nb}^{-1}$ )

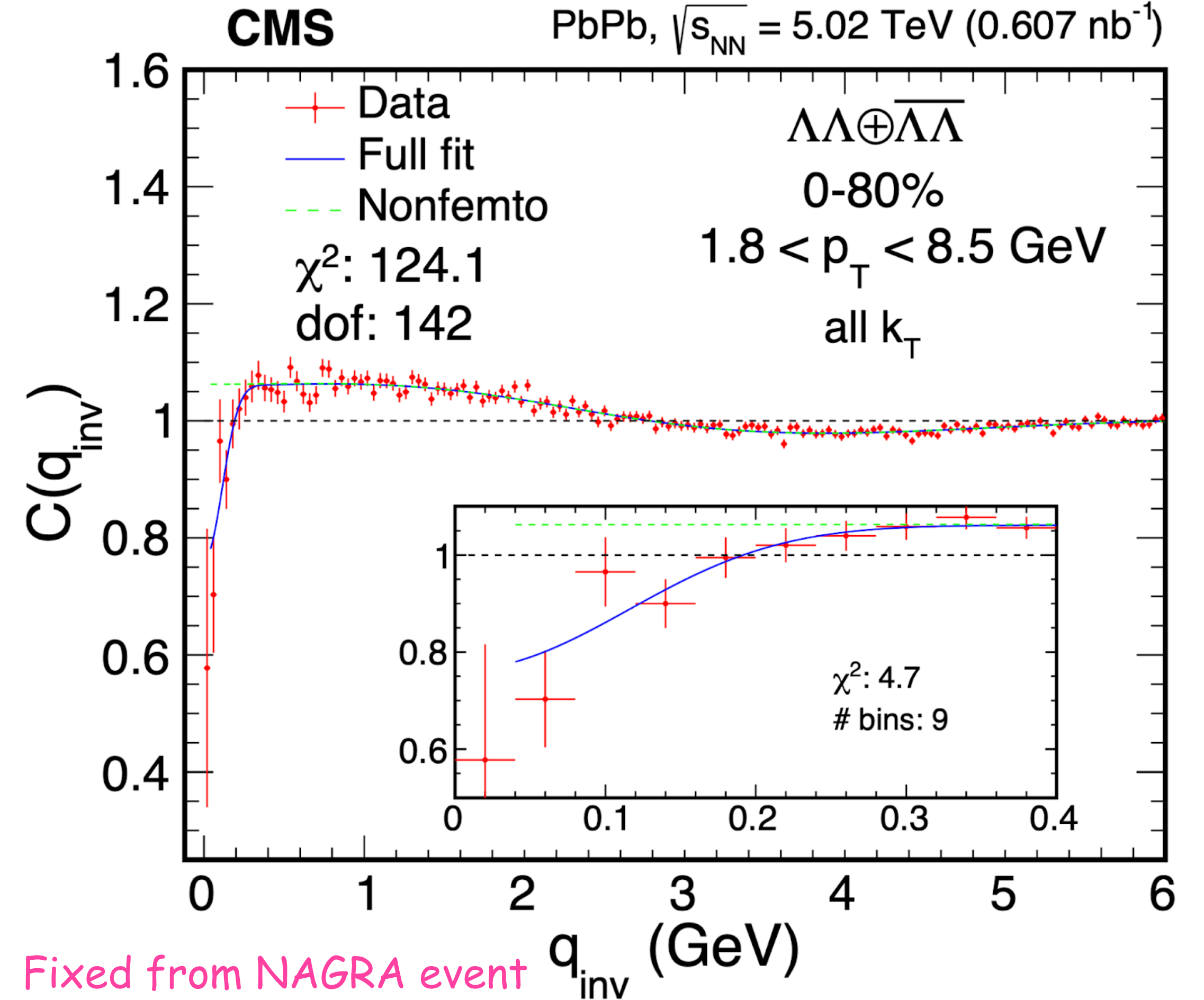


$$R_{inv} = 5.27 + 2.209 - 1.146$$

$$\lambda = 2.13 + 1.893 - 0.766$$

$\Lambda\Lambda \oplus \bar{\Lambda}\bar{\Lambda}$

PbPb,  $\sqrt{s_{NN}} = 5.02$  TeV ( $0.607 \text{ nb}^{-1}$ )



$$R_{inv} = 1.34 + 0.208 - 0.173$$

$$\lambda = 1.03 + 0.289 - 0.261$$

# Armenteros-Podolanski plot

$$\alpha = (p_{1L} - p_{2L}) / (p_{1L} + p_{2L})$$

$$p_{iL} = (\vec{p}_{V^0} \cdot \vec{p}_i) / |\vec{p}_{V^0}|$$

$$Q_T = |\vec{p}_1 \times \vec{p}_2| / |\vec{p}_{V^0}|$$

