

# **Measurements of (anti)(hyper)nuclei with ALICE**

### **Ramona Lea for the ALICE Collaboration**

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### Light nuclei in heavy-ion collisions





- The study of light (anti)(hyper)nuclei is very important:
	- Production mechanism is not well understood
		- How/when do they form?
			- "early" at chemical freeze-out (thermal production)
			- or "late" at kinetic freeze-out (coalescence)?
		- Do they suffer for the dissociation by rescattering?
	- Low binding energy (few MeV) "Snowballs in hell": nuclei formation is very sensitive to chemical freeze-out conditions and to the dynamics of the emitting source
	- Baseline for exotic bound state searches
	- Light nuclei measurements in high energy physics can be used to estimate the background of secondary anti-nuclei in dark matter search

### Antinuclei production



- Antinuclei can be a sign of Dark Matter annihilation:
	- Background: production in the collisions between cosmic rays (CR) and the interstellar medium (ISM) (pp and p-A collisions)
		- Nuclei production must be known very well!





ALT-PUB-532052

### Hypernuclei production

- Hypernuclei can be used to study nucleon-hyperon (N-Y) interaction
	- Production of exotic bound states
	- Determination of the equation of state
		- Application to neutron stars





**[D. Lonardoni et al., PRL 114, 092301 \(2015\)](https://journals.aps.org/prl/abstract/10.1103/PhysRevLett.114.092301)** 

[D. Logoteta et al., EPJA 55 \(2019\) 11, 207](https://doi.org/10.1140/epja/i2019-12909-9)

### Production models



- Statistical hadronization models (SHMs)
	- $\circ$  describe the yields of light-flavoured hadrons  $\frac{1}{3}$ by requiring thermal and hadron-chemical equilibrium
		- **■** Parameters:  $(T, V, \mu_B)$
	- light (anti)(hyper)**nuclei** are treated as **point-like objects**





### **07/11/2024 Ramona Lea - WPCF 2024**

### Production models: Statistical hadronization models

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	- describe the yields of light- flavoured hadrons by requiring thermal and hadron-chemical equilibrium
		- **■** Parameters:  $(T, V, \mu_B)$
	- light (anti)(hyper)**nuclei** are treated as **point-like objects**
	- Canonical ensemble (CSM): local conservation of quantum numbers (S, Q and B)
	- Central Xe–Xe collisions: π, K,  $\phi$ , p, d, <sup>3</sup>He
		- $\circ$   $T_{\text{chem}} = (154.2 \pm 1.1)$  MeV (Similar to the one obtained in Pb–Pb collision)
		- $V = (3626 \pm 298)$  fm<sup>3</sup>
		- $\degree$  x<sup>2</sup>/NDF = 0.83





### Production models: Coalescence

### **● Coalescence**

- Nuclei are formed by nucleons emitted at freeze-out hypersurface
- Coalescence calculations incorporate the **size of nuclei**
	- convolution between nucleon phase-space distribution and Wigner function of the nucleus



**[J. I. Kapusta, PRC 21, 1301 \(1980\)](https://journals.aps.org/prc/abstract/10.1103/PhysRevC.21.1301)** 

**[Mahlein et al., EPJC 83 \(2023\) 9, 804](https://link.springer.com/article/10.1140/epjc/s10052-023-11972-3)** 

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- Coalescence calculations incorporate the **size of nuclei**
	- convolution between nucleon phase-space distribution and Wigner function of the nucleus
- Coalescence parameter  $B_A$ , related to formation probability via coalescence:



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[Mahlein et al., EPJC 83 \(2023\) 9, 804](https://link.springer.com/article/10.1140/epjc/s10052-023-11972-3)

$$
E_A \frac{\mathrm{d}^3 N_A}{\mathrm{d} p_A^3} = B_A \left( E_p \frac{\mathrm{d}^3 N_p}{\mathrm{d} p_p^3} \right)^2
$$



ALICE particle identification capabilities are unique. Almost all known techniques are exploited: specific energy loss (d*E*/d*x*), time of flight, transition radiation, Cherenkov radiation, calorimetry and decay topology (V0, cascade).



ALICE Collaboration Int. J. Mod. Phys. A 29 (2014) 1430044



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### **TPC: d***E***/d***x* **in gas**

Separation of (anti)nuclei thanks to their large mass (and charge)



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TPC: d*E*/d*x* in gas

- Separation of (anti)nuclei thanks to their large mass (and charge)
- **TOF: measurements of velocity** *β* = *v*/*c*
	- *p* = *γβm* → mass



ALI-PERF-141622

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## Production spectra of nuclei

- ALICE measured production spectra of nuclei in pp, p–Pb, Xe–Xe and Pb–Pb collisions at mid-rapidity
- Measurements in classes of multiplicity or centrality
	- related to system size

[arXiv:2405.19826](https://arxiv.org/abs/2405.19826)





- Smooth evolution of  $d/p$  and 3He/p ratios with the system size
	- A=2 : multiplicity dependence is well reproduced by both **CSM** and **coalescence**
	- A=3 : ratio fairly described by the **coalescence** approach at low and high charged-particle multiplicity densities. Tension at intermediate multiplicities (10-40 charged particles)

## Constraining nuclei wave function



- A recent study shows that ALICE measurements in HM collisions of:
	- proton production yields
	- proton source radius
- Proton production yields<br>
roton source radius<br>
allow for the prediction of the deuteron spectrum via  $\sum_{n=1}^{\infty}$ event-by-event coalescence with no free parameters!





### Hypertriton production

- Lightest known hypernucleus consisting of  $(p, n, \Lambda)$
- Mass =  $2.991 \text{ GeV}/c^2$
- $B_{\Lambda} = 0.13 \pm 0.05$  MeV ( $B_{d} = 2.2$  MeV,  $B_{t} = 8.5$  MeV,  $B_{3He} = 7.7$  MeV)
- $\frac{3}{4}$ H has a large size:
	- $\circ$  d<sub>d−</sub> = 10.79 fm, r (d) = 1.96 fm



<https://hypernuclei.kph.uni-mainz.de/>

[F. Hildenbrand and H.-W. Hammer, Phys. Rev. C 100, 034002](https://journals.aps.org/prc/abstract/10.1103/PhysRevC.100.034002)



## Hypertriton production in pp

**● SHM** and **coalescence** predictions for <sup>3</sup><sub>Λ</sub>H are very different at low multiplicity



- [Phys. Rev. Lett. 128 \(2022\) 252003](https://doi.org/10.1103/PhysRevLett.128.252003)
- K.-J. Sun, et al. [Phys. Lett. B 792, 132 \(2019\)](http://dx.doi.org/10.1016/j.physletb.2019.03.033)
- **N.** V. Vovchenko, et al. [Phys. Lett. B 785, 171 \(2018\).](http://dx.doi.org/10.1016/j.physletb.2018.08.041)



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- $\frac{3}{\Lambda}$ H measured in Run 3 by ALICE with good precision





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- $\frac{3}{\Lambda}$ H measured in Run 3 by ALICE with good precision
- $\frac{3}{\Lambda}$ H/Λ is compared with the prediction of CSM and coalescence model
	- Two-body coalescence model provides the best description of data



➤ **Hypertriton in pp clearly favours coalescence**



● <sup>3</sup><sub> $\Lambda$ </sub>H has also been recently measured in Pb–Pb collisions at  $\sqrt{s_{_{\rm NN}}}$  = 5.02 TeV



**[arXiv:2405.19839](https://arxiv.org/abs/2405.19839)** 







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- $\bullet$   $p_{\tau}$ -differential measurement is also in agreement with **blast-wave** with common parameters with other nuclei
	- Large statistical uncertainties → Ongoing  $p_{_{\rm T}}$ -differential analyses with Run 3 data are fundamental to disentangle the two models

### **Q** [arXiv:2405.19839](https://arxiv.org/abs/2405.19839)

### 4He in Pb–Pb collisions

- ALICE has measured the production spectra for (anti)<sup>4</sup>He in Pb-Pb
- <sup>4</sup>He is more bound and compact than lighter nuclei:
	- $\circ$   $E_{\rm B}$  (<sup>4</sup>He) ~ 28 MeV, *r*(<sup>4</sup>He) ~ 1.7 fm
- $\bullet$   $p_{\tau}$  spectra are well reproduced by a blast-wave function, using common parameters with the other nuclei



[Phys. Lett. B 858 \(2024\) 138943](https://doi.org/10.1016/j.physletb.2024.138943)





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- $\bullet$   $p_{\tau}$  spectra are well reproduced by a blast-wave function, using common parameters with the other nuclei
- The parameter  $B_4$  is compared with **SHM+blast wave** and **coalescence** predictions
	- **SHM describes nuclei with A = 4 better**  $\Box$  [Phys. Lett. B 858 \(2024\) 138943](https://doi.org/10.1016/j.physletb.2024.138943)







- **SHM** predicts hypernuclei with  $A = 4$  in Pb–Pb collisions
	- they are rare:
		- $\blacksquare$  penalty factor for increasing A:  $\sim$  300
		- suppression due to strangeness content



A. Adronic, private communication, based on ◯◯ [A. Andronic et al., PLB 697 \(2011\) 203-207](https://www.sciencedirect.com/science/article/pii/S0370269311001006?via%3Dihub)



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- Some factors may enhance the yield  $(x 4)$ :
	- $\circ$  larger binding energy wrt A = 3
	- existence of excited states  $dN$  $\frac{dN}{dy} \propto 2J+1$
	- spin degeneracy



[M. Schäfer et al., PRC 106, L031001 \(2022\)](https://journals.aps.org/prc/abstract/10.1103/PhysRevC.106.L031001)

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	- $^{4}$ <sub>Λ</sub>H  $\rightarrow$  <sup>4</sup>He +  $\pi^{-}$
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	- $^{4}$ <sub> $\Lambda$ </sub>He → <sup>4</sup>He + p +  $\pi^{-}$
- Yields in agreement with **SHM** prediction that includes feed-down from excited states
	- ➤ **SHM describes hypernuclei with A = 4 well**



1 [arXiv:2410.17769](https://arxiv.org/abs/2410.17769)





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	- It is sensitive to a different production in-plane and out-of-plane
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	- Pb–Pb collisions at  $\sqrt{s_{NN}}$  = 5.36 TeV<br>
	 more differential both in  $p_{\text{T}}$  and centrality, more precise than in Run 2





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- Data are compared with the predictions of blast wave and coalescence model
	- ➤ **coalescence is favoured**
- Flow of hypertriton has been measured for the first time:
	- **○ compatible with 3He, but large uncertainties currently**





- In cosmic rays interstellar medium collisions (anti)nuclei are mainly produced at forward rapidity:
	- important to study nuclear production vs rapidity



**A. K. Blum Phys. Rev. C 109, L031904** 

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	- $\circ$  important to study nuclear production vs rapidity
- Measurement of p and d production in rapidity intervals (|y|<0.7)





[arXiv:2407.10527](https://arxiv.org/abs/2407.10527)



- In cosmic rays interstellar medium collisions (anti)nuclei are mainly produced at forward rapidity:
	- $\circ$  important to study nuclear production vs rapidity
- Measurement of p and d production in rapidity intervals (|*y*|<0.7)
- $\bullet$  *B*<sub>2</sub> is measured as a function of  $p_{\text{T}}$  and *y*:
	- data are compared with predictions from coalescence (PYTHIA and EPOS+**Coalescence**)
	- $\blacktriangleright$  the shape is correctly reproduced, the magnitude not





[arXiv:2407.10527](https://arxiv.org/abs/2407.10527)



- Production of (anti)(hyper)nuclei measured at mid-rapidity in pp, p–Pb, Xe–Xe and Pb–Pb
	- light nuclei ( $E_B$  ~ MeV) are reproduced by both **SHM** and **coalescence**
	- loosely bound objects such 3 ΛH (*B*<sup>Λ</sup> ~ 100 keV ) are described better by **coalescence** as it includes nuclei size in the estimation
- With Run 3, some measurements that were possible only in Pb–Pb collisions will be accessible also in small systems
	- Measurements will help to disentangle the different production models providing a clearer understanding of the dynamics underlying nuclei formation dynamics

**Backup slides**

### Event-by-event coalescence

Possible to implement event-by-event coalescence, with probability:

$$
\mathcal{P}(r_0,q)=\int d^3r_\mathrm{d}\int d^3r H_\mathrm{pn}(\vec{r},\vec{r}_\mathrm{d};r_0)\mathcal{D}(\vec{q},\vec{r})
$$

- $\circ$   $r_{\text{o}}$  is the size of the emitting source
- *q* is the relative p-n momentum
- Two-particle emitting source: average two-particle distance
- Wigner transform of the deuteron wavefunction
- production measurements to constrain the nuclear wave function



### Rapidity dependence of coalescence

- Model predictions based on ALICE measurements are used as input to calculate antideuteron flux from cosmic rays
- **But typically ALICE measurements** cover midrapidity (|y|<0.5) while astrophysical models extrapolate to the forward region



[Kfir Blum](https://journals.aps.org/prc/abstract/10.1103/PhysRevC.109.L031904) [Phys. Rev. C 109, L031904](https://journals.aps.org/prc/abstract/10.1103/PhysRevC.109.L031904) 

# <sup>3</sup><sub>Λ</sub>H *p*<sub>T</sub> spectra in Pb–Pb





[arXiv:2405.19839](https://arxiv.org/abs/2405.19839)M





ALI-PUB-577118

[arXiv:2407.10527](https://arxiv.org/abs/2407.10527)

### Deuteron production in PYTHIA

- PYTHIA 8.3:
	- d production via ordinary reactions
	- Energy dependent cross sections parameterized based on data

Reactions:

- $p + n \rightarrow \gamma + d$
- $\bullet$  p + n  $\rightarrow \pi^0$  + d
- $p + n \rightarrow \pi^0 + \pi^0 + d$
- $\bullet$  p + n  $\rightarrow$   $\pi$ <sup>+</sup> +  $\pi$ <sup>-</sup> + d
- $\bullet$  p + p  $\rightarrow \pi^+$ + d
- $p + p \rightarrow \pi^+ + \pi^0 + d$
- $\bullet$  n + n  $\rightarrow$   $\pi$ <sup>-</sup> + d
- $\bullet$  n + n  $\rightarrow$   $\pi$ <sup>-</sup> +  $\pi$ <sup>0</sup> + d





### Light nuclei flow measurement





Angular distribution of reconstructed charged particles can be expanded into a Fourier series w.r.t. symmetry plane  $\psi_{\boldsymbol{n}}$ :

$$
E\frac{d^3N}{dp^3} = \frac{1}{2\pi} \frac{d^2N}{p_T dp_T dy} \left( 1 + \sum_{n=1}^{\infty} 2v_n \cos\left(n\left(\phi - \psi_n\right)\right) \right)
$$

$$
v_n = \left\langle \cos\left(n\left(\phi - \psi_n\right)\right) \right\rangle
$$

- Elliptic flow  $(v_2)$  is sensitive to the system evolution:
	- It probes initial conditions and constraints particle production mechanisms
- The measurement of light nuclei  $v_n$  will help in the understanding of particle production mechanisms
	- Do light nuclei follow the mass ordering observed for lighter particles?
	- Do light nuclei follow a quark/baryon number scaling (coalescence) or follow mass scaling (hydro)?

### Flow analysis method

- $\bullet$  v<sub>2</sub> is measured using the scalar product method
	- Hits measured by V0A (2.8 < η < 5.1) and V0C (-3.7 < η < -1.7) are used as reference particles
	- $\circ$  Deuteron candidates are the particles of interest (|η|<0.8)
- The contribution to the measured elliptic flow  $(v_2^{\text{Tot}})$  due to misidentified deuterons  $(v_2^{\text{Bkg}})$ was removed by studying the azimuthal correlations versus ΔM (ΔM =  $m_{\tau$ <sub>OF</sub> -  $m_{\sf d}}$ )

$$
v_n(\Delta M)=v_n^{\text{sig}}\frac{N^{\text{sig}}}{N^{\text{tot}}}(\Delta M)+v_n^{\text{bkg}}(\Delta M)\frac{N^{\text{bkg}}}{N^{\text{tot}}}(\Delta M)\frac{\sum\limits_{\substack{\hat{\alpha}\\ \text{g. 0.3}}}^{\widehat{\alpha}}\widehat{\sigma}_{0.3}}{\sum\limits_{\substack{\hat{\alpha}\\ \text{g. 0.2}}}^{\widehat{\alpha}}\widehat{\sigma}_{0.3}}\widehat{\sigma}_{0.3}^{\text{ALICE Pb-Pb}\sqrt{s_{NN}}=5.02\text{ TeV}}
$$

The yields  $N^{Sig}$  and  $N^{Bkg}$  are extracted from fits of the invariant mass distribution obtained with the TOF detector

**[Phys. Rev. C 102 \(2020\) 055203](https://link.aps.org/doi/10.1103/PhysRevC.102.055203)** 

$$
24 \qquad \qquad \blacksquare
$$

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 $v_n\{\text{SP}\} = \frac{\left\langle u_{n,i}(p_T,\eta)\cdot \frac{Q_n^*}{M}\right\rangle}{\sqrt{\left\langle \frac{Q_{n,A}^*}{M_A}\cdot \frac{Q_{n,B}^*}{M_B}\right\rangle}}$ 

 $2.2 \le p_{\tau} < 2.4$  GeV/c

 $0.1$