

# Femtoscopy analysis of ultrasoft pion trap at energies available at the CERN Large Hadron Collider

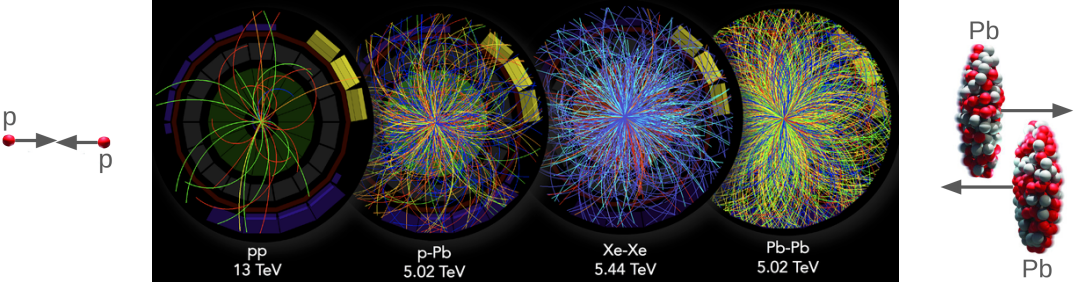
Wioleta Rzęsa (Warsaw University of Technology)

WPCF 2024 (Toulouse, France)

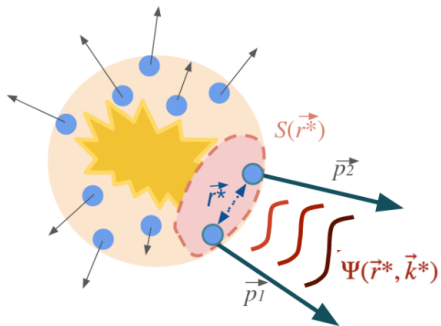
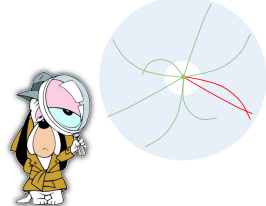
WUT



# Collisions of relativistic particles and heavy ions



# Is there a "pattern" between the tracks? Femtoscscopy



$$C(k^*) = \int S(r^*) |\Psi(k^*, r^*)|^2 d^3r^*$$



Source function

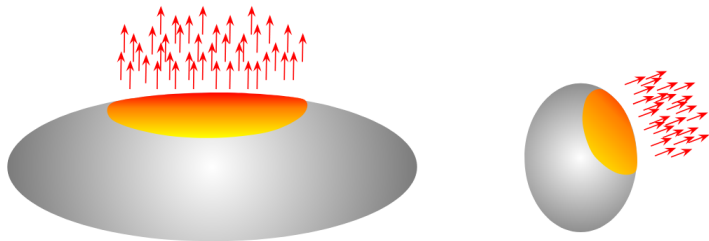


Wave function

# Homogeneity

In femtoscopy we study the space-time spread of particles from the emitter identified as its **homogeneity region**.

**Homogeneity region** → area of the source from which particles are emitted with similar velocities and directions → **it is not the overall source!**



# How to describe the homogeneity region?

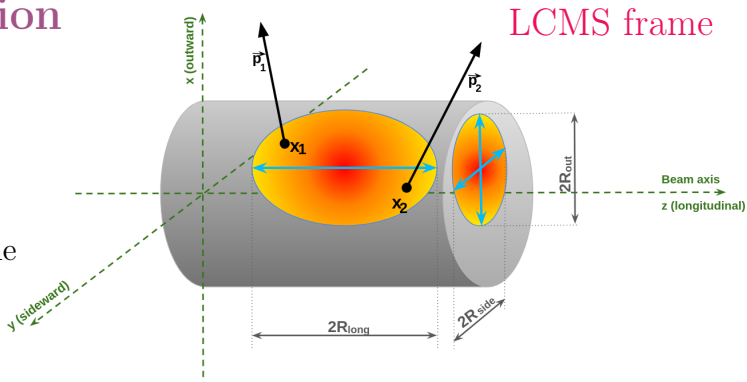
## – the source function

→ Direct information about the size in each direction.

→ Gaussian-based definition

$$S(r) \sim \exp\left(-\frac{r_{out}^2}{2R_{ABout}^2} - \frac{r_{side}^2}{2R_{ABside}^2} - \frac{r_{long}^2}{2R_{ABlong}^2}\right),$$

where  $R_{AB} = \sqrt{R_A^2 + R_B^2}$  is the homogeneity length for particles A and B of the pair.



# How to reflect 3D shape with only one R?

PRF frame\*

A center-of-mass of  
the pair is at rest

→ Not trivial average over 3 directions.

→ Gaussian based definition:

$$S(r) \sim \exp\left(-\frac{r_{\text{out}}^2}{2R_{\text{A}B\text{out}}^2} - \frac{r_{\text{side}}^2}{2R_{\text{A}B\text{side}}^2} - \frac{r_{\text{long}}^2}{2R_{\text{A}B\text{long}}^2}\right)$$

↓

$$S(r^*) \sim r^{*2} \exp\left(\frac{r^{*2}}{2R_{\text{AB}}^2}\right)$$

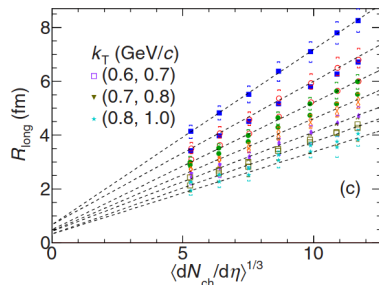
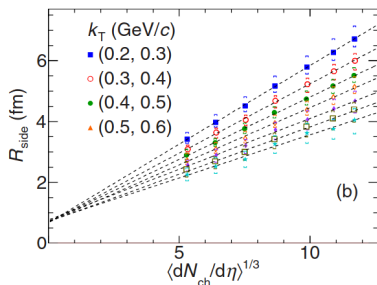
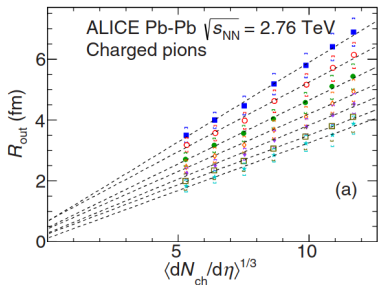
$$k_{\text{A}}^* = -k_{\text{B}}^*$$

# Observed scalings of the source sizes

# A linear scaling

→ A linear scaling of the femtosopic radii is universally observed versus cube root of the final state mean particle multiplicity  $\langle dN_{ch}/d\eta \rangle^{1/3}$ .

→ Observed for different momentum and particle types.



Phys. Rev. C 93, 024905

The more particles produced in the collision  
the bigger femtosopic volume.



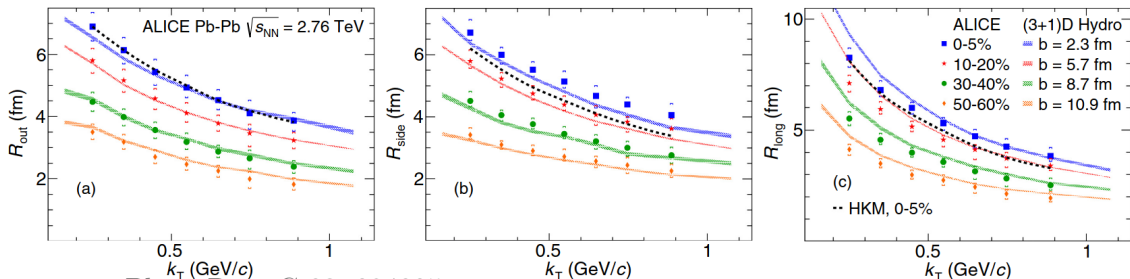
# A power-law scaling

→ The radii versus the pair transverse momentum exhibit a power-law-like scaling in pair transverse mass  $m_T$ .

→ Observed for all femtoscopic radii and all particle types.

$m_T = \sqrt{\langle k_T \rangle^2 + m_0^2}$ , where  $k_T$  is pair mean transverse momentum

$k_T = |\mathbf{p}_{T1} + \mathbf{p}_{T2}|/2$ , and  $m_0$  is particle mass)

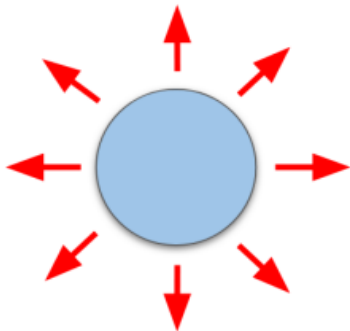


Phys. Rev. C 93, 024905

But where this  $m_T$  scaling come from?

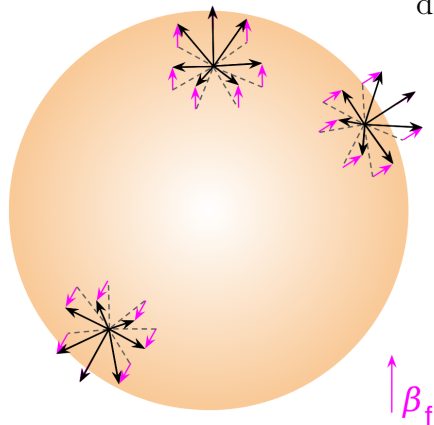
## Collective flow

The particles gain velocity due to hydrodynamic expansion originating from the non zero pressure gradient of the fireball.



# Collective motion

All particles emitted close to each other, in radial direction pointing outwards from the center gain a **common velocity**,  $\beta_f$  from the collectivity.



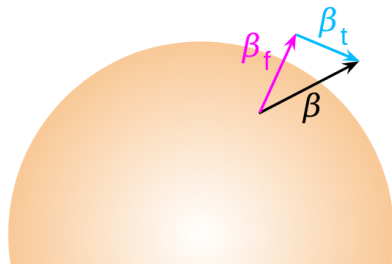
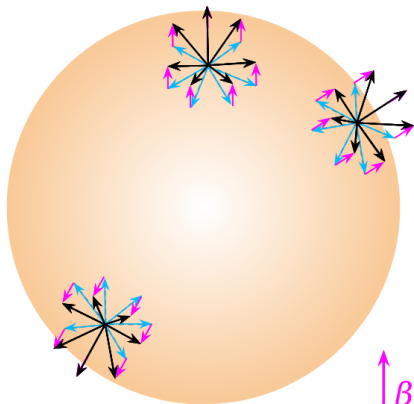
The process drives the space-time evolution of the system. It should apply to all particles.

# Collective motion

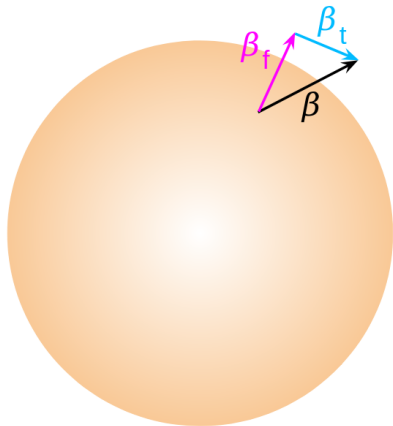
All particles emitted close to each other, in radial direction pointing outwards from the center gain a **common velocity**,  $\beta_f$  from the collectivity.

The process drives the space-time evolution of the system. It should apply to all particles.

In addition each particle gain a random **thermal velocity**,  $\beta_t$ .



# Collective motion



$$\langle x_{out} \rangle = \frac{\langle r \beta_f \rangle}{\langle \sqrt{\beta_t^2 + \beta_f^2} \rangle}$$

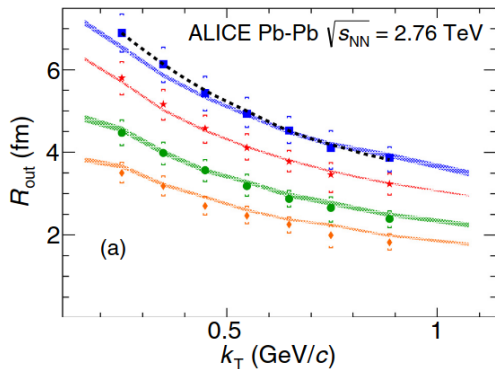
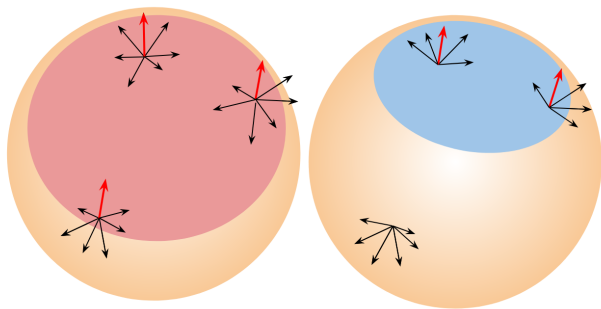
If  $\beta_f$  is dominant  $\rightarrow$  all particles are emitted from the **edge** of the source

If  $\beta_t$  is dominant  $\rightarrow$  all particles are emitted from the **whole** source

If  $\beta_f$  and  $\beta_t$  are similar  $\rightarrow$  for different  $p_T$  we should observe different source sizes and asymmetry (shift).

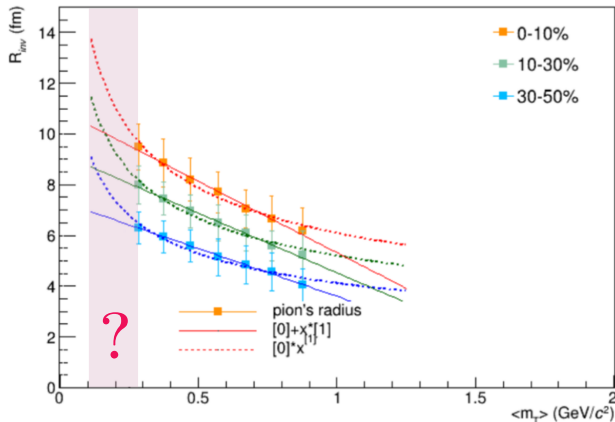
# Consequence of flow

As observed  $p_T$  grows, the region where pairs with similar velocity and direction can be emitted gets smaller and shifted toward the edge of the source.



# Power-law?

But what if power-law is not so universal description?







## Simulation study of ultra-soft pions

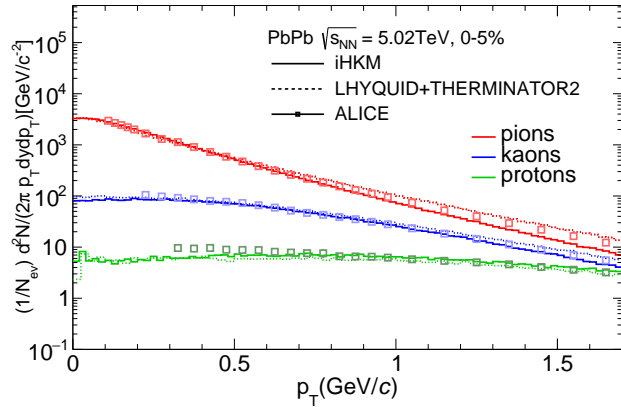
# Modeling the hydrodynamic of heavy-ion collisions

→ integrated HydroKinetic Model: **iHKM**

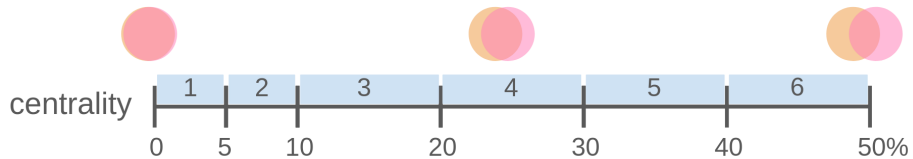
→ hydrodynamic model LHYQUID coupled to the statistical hadronisation code THERMINATOR2: **LHYQUID+THERMINATOR2**

The two models differ in the initial conditions and the way of modeling the system evolution.

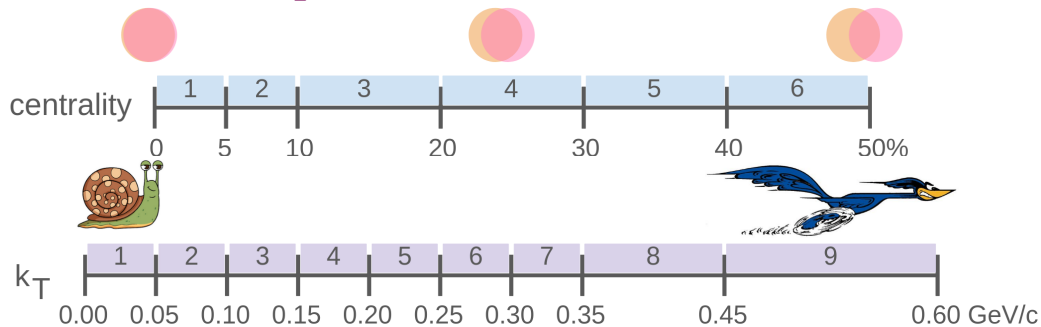
Both well describe measured data.



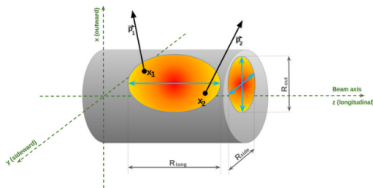
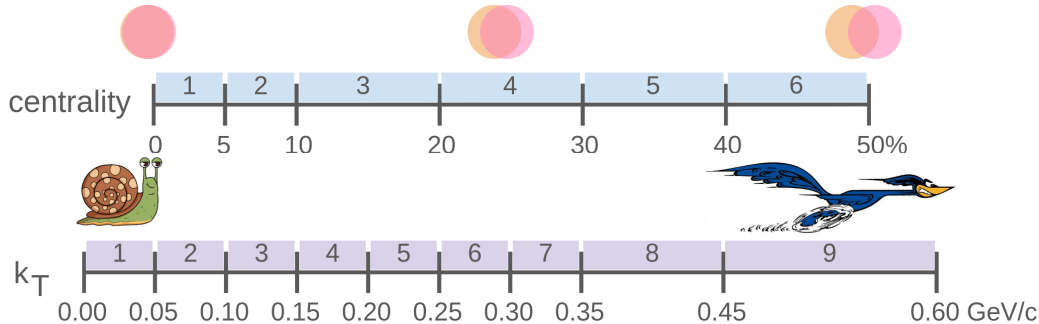
$\pi^+\pi^+ \oplus \pi^-\pi^-$  pairs



# $\pi^+\pi^+ \oplus \pi^-\pi^-$ pairs



# $\pi^+\pi^+ \oplus \pi^-\pi^-$ pairs



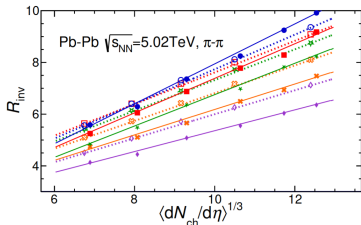
Analysis

$R_{out}, R_{side}, R_{long},$

$R_{inv}$

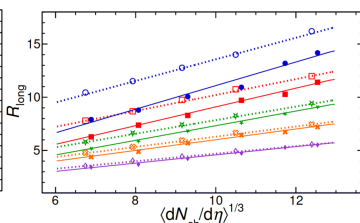
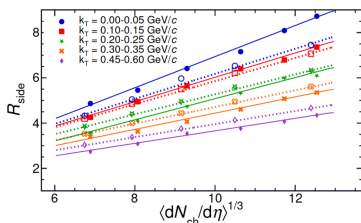
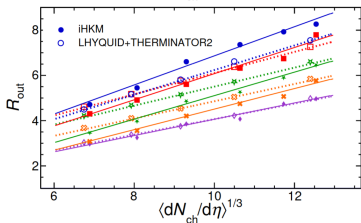
# $\langle dN/d\eta \rangle^{1/3}$ dependence

1D



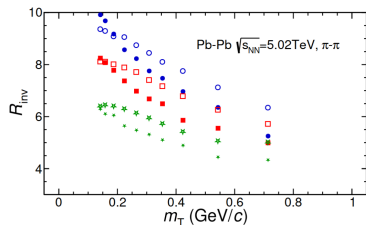
Linear scaling preserved in all  $k_T$  bins.

3D



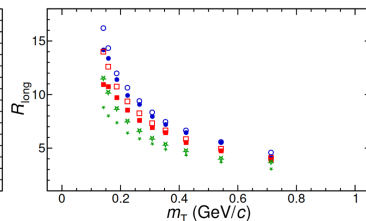
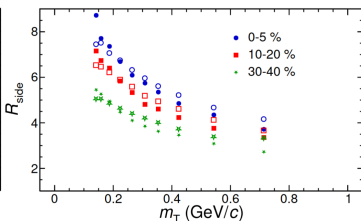
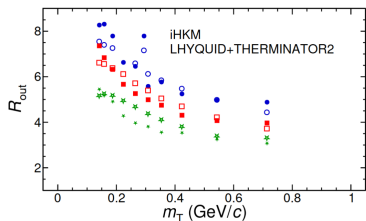
# $m_T$ dependence

1D

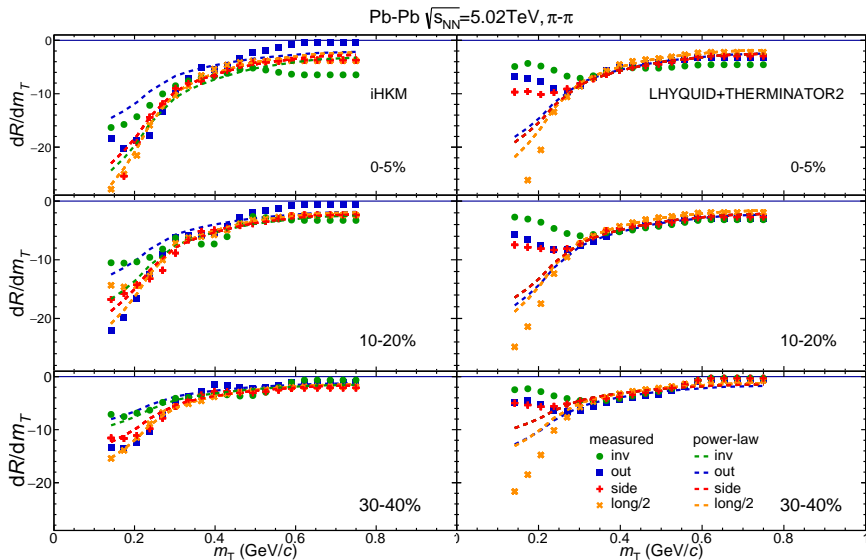


Radii in low  $m_T$  visibly change the behavior in most of the cases.

3D



# $dR/dm_T$ dependence



Constant slope in low  $m_T$  in  $R_{inv}$  and transverse directions of 3D

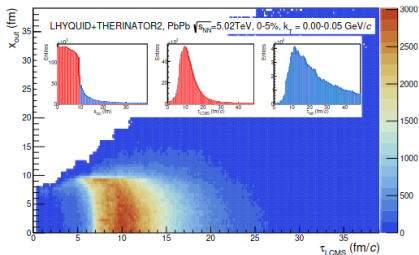
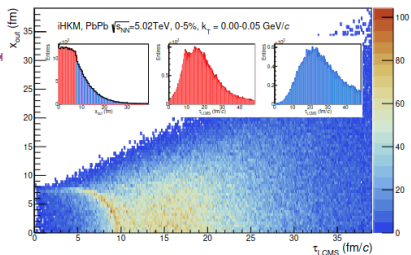


# Emission position vs time

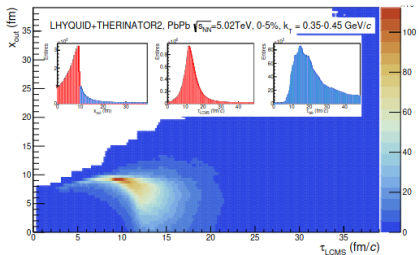
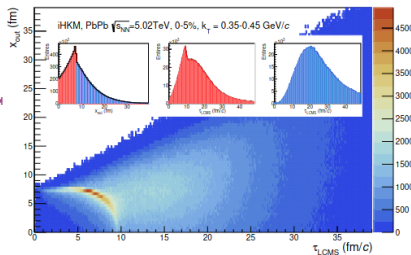
model 1

model 2

ultra-soft pions



normal pions



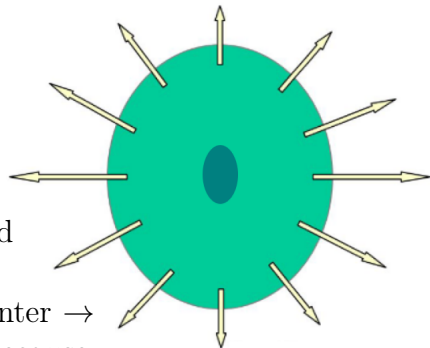
Soft pions are kept inside and escape the system later than other pions.

# Trap?

The gradient of density is not equal in the radial directions: it is more strong at the periphery, and less strong in the center, where soft hadrons mostly come from.

The hadrons stay together for a longer time

- Very low transverse collective velocity;
- Smaller density gradient in the center as compared to non-central and periphery parts.
- The initially highest density in the geometrical center → this high-mass-density region is difficult to expand because of the relatively small transverse pressure gradient in the center.



# Condensation of pions?

Bose-Einstein condensation of pions in heavy-ion collisions at the  
CERN Large Hadron Collider (LHC) energies

Viktor Begun<sup>1,2</sup> and Wojciech Florkowski<sup>3</sup>

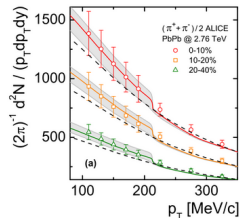
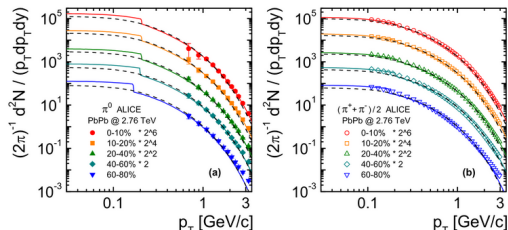
<sup>1</sup>*Institute of Physics, Jan Kochanowski University, PL-25406 Kielce, Poland*

<sup>2</sup>*Bogolyubov Institute for Theoretical Physics, 03680 Kiev, Ukraine*

<sup>3</sup>*The H. Niewodniczański Institute of Nuclear Physics,  
Polish Academy of Sciences, PL-31342 Kraków, Poland*

(Dated: July 6, 2018)

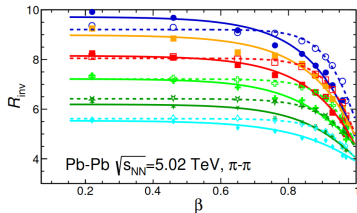
“Our study shows that the pion condensate rate in the present data is 2% in the most central collisions and 7 – 8% in peripheral collisions (with a sudden increase to 19% in the most peripheral collisions).”



# Velocity dependence

$$\beta = k_T/m_T, \text{ fitted with: } R(\beta) = \frac{a}{e^{(\beta-b)/c} + 1}$$

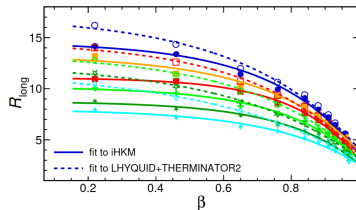
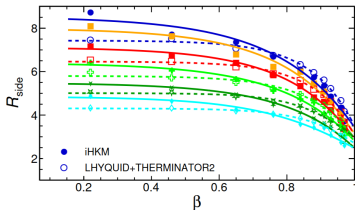
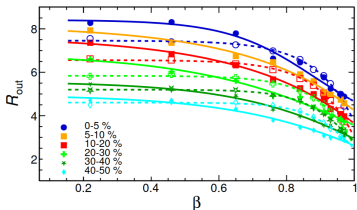
1D



$\beta$  decreases  $\rightarrow$  R became similar

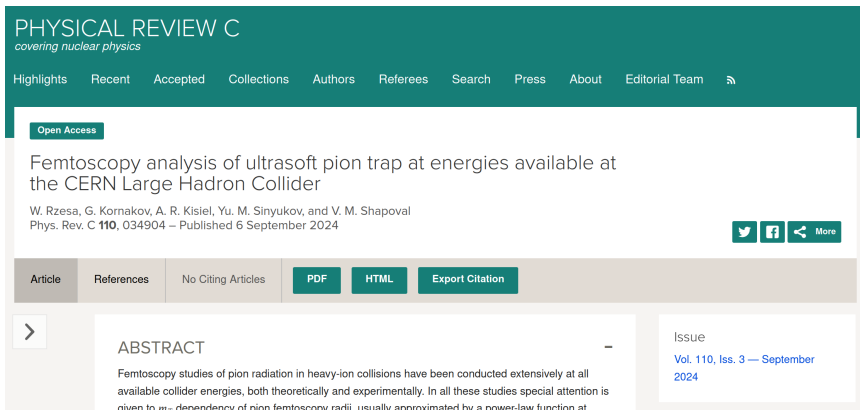
Idea: such value could be used to determine an extrapolated value of the radii for ultra-soft pions, instead of the historically used power-law  $m_T$  scaling.

3D



# Conclusions from the ultra-soft pion's simulations

The power-law description of pion's source size  $m_T$ -scaling is not the best solution for ultra-soft region → especially in  $1D$  source study the description with the power-law is not satisfactory.



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Femtoscopia analysis of ultrasoft pion trap at energies available at the CERN Large Hadron Collider

W. Rzesza, G. Kornakov, A. R. Kisiel, Yu. M. Sinyukov, and V. M. Shapoval  
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ABSTRACT

Femtoscopia studies of pion radiation in heavy-ion collisions have been conducted extensively at all available collider energies, both theoretically and experimentally. In all these studies special attention is given to  $m_T$  dependency of pion femtoscopia radii, usually approximated by a power-law function at

Thank you for your attention!