



# The fundamental role of fission during the r-process nucleosynthesis

Part II:

The rapid neutron-capture process

- Astrophysical aspects
- Nuclear physics aspects & fission
  - Nuclear physics uncertainties
  - Models of relevance for r-process calculation
  - Role of fission
  - Calculation of fission observables & their impact

# Another uncertainty: nuclear physics input $(n,\gamma) - (\gamma,n) - \beta$ competition & Fission

Still many open questions

- β-decay rates
- $(n,\gamma)$  and  $(\gamma,n)$  rates
- Fission (nif, sf,  $\beta$ df) rates
- Fission Fragments Distributions  $\int$  some 5000 nuclei with  $Z \le 110$  involved on the n-rich side



## Nuclear physics input (n, $\gamma$ ) – ( $\gamma$ ,n) – $\beta$ competition & Fission

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- (n, $\gamma$ ) and ( $\gamma$ ,n) rates  $\rightarrow$  Masses
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Still many open questions

#### Nuclear uncertainties

Some Progress in considering "theoretical uncertainties" in NA

Two types of uncertainties affecting nuclear inputs (e.g Masses)



Model or parameter variations must be constrained by experimental data e.g. mass models with  $\sigma_{\rm rms} < 0.8$  MeV or  $(n,\gamma)$  models with  $f_{\rm rms} \le 2$ 

But what about their impact on astrophysical observables?

How to propagate such NP uncertainties into astrophysics simulations ??



#### Uncorrelated MC approach

Model-correlated approach



- Rates within an arbitray factor of 2, 10, 100
- Neglect correlations between uncertainties
- Overestimates impact if not exp-constrained
- Coherent model-correlated uncertainties
- Only parameter or model uncertainties
- Overestimates impact if not exp-constrained

#### In all cases, propagation must be applied to a large representative sample of trajectories

#### Impact of nuclear model uncertainties on the composition of NSM ejecta

15 different "acceptable" sets of nuclear inputs (masses, β-decay, n-capture, fission) Kullmann et al. (2022) Prompt dynamical ejecta: SFHo 135-135

Single trajectory

Multiple trajectories



Global & Local discrepancies

Local (correlated) discrepancies

Astrophysical models evolve and may still not be robust (what is "important" today may not be next year)

> → Remain extremely critical about propagation of nuclear uncertainties into astrophysical models



#### Nuclear inputs to nuclear reaction & decay calculations



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#### Nuclear inputs to nuclear reaction & decay calculations



"Microscopic" approach is a necessary but not a sufficient condition ! "(Semi-)Microscopic" models must be competitive in reproducing exp. data !

#### Role of fission in the *dynamical* ejecta of NS-NS Merger

using multiple (~500) trajectories and correlated model uncertainties



	total	dynamical	NS-torus	BH-torus
$M_{\rm ej}  [{ m M_o}]$	0.024	0.012	0.0	0.012
<ye></ye>	0.27	0.26	0.0	0.28
<i>v</i> <sub>ej</sub> [c]	0.14	0.24	0.0	0.05

#### Role of fission in the dynamical ejecta of NS-BH Merger

Fission expected to be even more important in the dynamical ejecta of NS-BH merger

NS-BH (no weak interaction)

NS-NS (with weak interactions)



Fission processes and fission fragment distribution of relevance for estimating the

1. termination point of the r-process (recycling, prod of SH, Universality ?)

After a few hundred ms in the dynamical ejecta

![](_page_12_Figure_4.jpeg)

Fission processes and fission fragment distribution of relevance for estimating the

- 1. termination point of the r-process (recycling, prod of SH)
- 2. production of light species ( $A \sim 110-160$ ) by fission recycling

![](_page_13_Figure_4.jpeg)

Fission processes and fission fragment distribution of relevance for estimating the

- 1. termination point of the r-process (recycling, prod of SH)
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- 3. heating of the ejecta during the expansion

![](_page_14_Figure_5.jpeg)

Fission processes and fission fragment distribution of relevance for estimating the

- 1. termination point of the r-process (recycling, prod of SH)
- 2. production of light species ( $A \sim 110-160$ ) by fission recycling
- 3. heating of the ejecta during the expansion
- 4. production of Pb-peak elements
- 5. production of radiocosmochronometers (U, Th)

![](_page_15_Figure_7.jpeg)

## Nuclear fission

Three fission modes play an important role during the r-process nucleosynthesis:

- spontaneous fission: strongly depends on the fission barrier height
- neutron-induced fission: strongly depends on  $S_n B_f$  (for  $E_n = kT \sim 100 \text{ keV}$ )
- $\beta$ -delayed fission, *i.e.* fission following a  $\beta$ -decay: strongly depends on  $Q_{\beta} B_f$

![](_page_16_Figure_5.jpeg)

#### **Fission properties and the r-process in NSM**

![](_page_17_Figure_1.jpeg)

#### $90 \le Z \le 101$ Elemental abundances expected in the dynamical ejecta HFB-14 fission barriers

![](_page_18_Figure_1.jpeg)

Significant production of actinides up to  $Z \sim 100$ 

P.S. overall yields mass-averaged over hundreds of ejected mass elements

#### $100 \le Z \le 122$ elemental abundances expected in the dynamical ejecta HFB-14 fission barriers

![](_page_19_Figure_1.jpeg)

Significant production of actinides up to  $Z \sim 100$ Non-zero production of super-heavies up to  $Z \sim 120$ ... if fission allows it ...

#### Elemental abundances expected in the dynamical ejecta

![](_page_20_Figure_1.jpeg)

100% of matter experiences fission

~10% of matter experiences fission

#### Significant production of lanthanides and actinides

P.S. overall yields mass-averaged over hundreds of ejected mass elements

#### **Calculation of the fission probability**

cf talk of D. Regnier

![](_page_21_Figure_2.jpeg)

$$T(E, J, \pi) = \frac{T_A T_B}{T_A + T_B}$$

$$T(E, J, \pi) = \sum_d P(E - \varepsilon_d^{J, \pi}) + \int_{E_d}^E P(E - \varepsilon)\rho(\varepsilon, J, \pi)d\varepsilon$$

$$P_j(E) = \frac{1}{1 + \exp(2K_j)} \quad \text{WKB}$$

$$K_j = \pm \int_{a_j}^{b_j} [2\mu(E - V_j(\beta))/\hbar^2]^{1/2} d\beta \quad \stackrel{+\text{ for } E > B_j}{-\text{ for } E < B_j}$$

 $\mu \sim 0.054 \text{ A}^{5/3}$  (inertial mass) often assumed deformation-independent

### **Fission in the Nucleosynthesis Context**

Complicate NP input associated with HFB-type calculation of

- Potential Energy Surface → fission path, inertial mass
- NLD at the saddle points (& isomeric wells)

 $\rightarrow$  WKB Fission probability

- - Spontaneous fission: Least-action path integral (2D HFB PES)
    - + coupling with competitive n-,  $\gamma$ -,  $\beta$ -channels
  - Neutron-induced fission: HF reaction code
  - $\beta$ -delayed fission probability: QRPA GT  $\beta$ strength + HF reaction code
  - Real effort needed to improve *predictions* of fission properties (Still far from being achieved, even for U and Th !)

![](_page_22_Figure_11.jpeg)

## **Calculation of the fission path and barriers** The fundamental role of octupole deformations

Two-dimensional Potential Energy Surface based on mean-field approach

![](_page_23_Figure_2.jpeg)

#### The fundamental role of triaxiality

![](_page_24_Figure_1.jpeg)

#### Effects of triaxiality on *both*

• Triaxial inner barrier

![](_page_25_Figure_2.jpeg)

• Triaxial- and octupole-deformed outer barrier

(also for odd-*A* and odd-odd nuclei)

![](_page_25_Figure_5.jpeg)

Only a few models applied to the  $\sim 2000$  nuclei of r-process interest

![](_page_26_Figure_1.jpeg)

 $\sigma(45 B_f) > 0.6 \text{ MeV}$ 

#### New BSkG3 predictions

Grams et al. (2023)

#### New Skyrme-HFB mass model: BSkG3

- Triaxiality, time-reversal symmetry breaking & octupole GS deformation
- Microscopic pairing from "realistic" calc.
- Stiff EoS (NS  $M_{\text{max}} \sim 2.26 M_{\text{o}}$ )
- Accurate masses:  $\sigma(2457M)=0.63$  MeV
- Accurate fission barriers  $\sigma(45B_f)=0.33$  MeV including triaxial & octupole deformations simultaneously

	BSkG3	
σ( <i>M</i> ) [MeV]	0.63	
σ( <i>E</i> <sub>l</sub> ) [MeV]	0.33	
σ( <i>E</i> <sub>II</sub> ) [MeV]	0.51	
$\sigma(E_{iso})$ [MeV]	0.36	

![](_page_27_Figure_9.jpeg)

Including odd-A & odd-odd

 $\rightarrow$  being extended now to the calculation of 2000 nuclei...

#### Projection of the static path along the quadrupole deformation parameter $\beta_2$

![](_page_28_Figure_1.jpeg)

![](_page_29_Figure_0.jpeg)

#### **Fission properties mainly depend on the primary fission barriers**

![](_page_30_Figure_1.jpeg)

#### Fission barriers calculated for all nuclei with $90 \le Z \le 120$

Giuliani et al. (2018)

#### Impact of the barrier height on the r-process on the composition of the dynamical NS-BH ejecta

![](_page_31_Figure_1.jpeg)

#### FISSION CROSS SECTIONS (Microscopic approaches)

#### Nuclear level densities at the saddle points

HFB model provides at each deformation (including saddle points) all nuclear properties needed to estimate the NLD

single-particle scheme

![](_page_32_Figure_3.jpeg)

Possibility to estimate NLD at the saddle point within the HFB+Combinatorial model

#### FISSION CROSS SECTIONS (Microscopic approaches) Nuclear Level Density at Saddle Points

- Fission Barriers and saddle point deformations (Q,O,H) determined within HFB method
- Nuclear properties (spl, pairing) at the inner and outer saddle points with constrained HFB model
- NLD in the framework of the microscopic combinatorial model based on HFB single-particle level and pairing predictions at the HFB saddle points (plus collective rotational and vibrational enhancement)

All ingredients described on the basis of the same Skyrme effective interaction (BSk14) at GS and Saddle Points

NLD in a table format at inner and outer saddle points

(~2000 nuclei : 2/3 saddle points & 1/2 shape isomers)

For inner barrier, usually predicted to be triaxial:  $\rho_{triax} = \sqrt{\frac{\pi}{2}} \sigma_{\perp} \times \rho_{Comb}$  Bjornholm & Lynn (1980) For outer barrier, usually predicted to be left-right asymmetric:  $\rho_{asym} = 2 \times \rho_{Comb}$ 

#### FISSION CROSS SECTIONS (HFB-14 fission path & NLD)

![](_page_34_Figure_1.jpeg)

Fission barriers adjusted for each target

Fission barriers adjusted for each type of target - odd-odd

- odd-even
- even-odd
- even-even

 $\Rightarrow$  Quite "decent" (astrophysically speaking) after few adjustments.

 $\Rightarrow$  Default calculations not sufficient for applications.

## Fit to <sup>23x</sup>U+n cross sections in the 0.001 – 30 MeV range 1 UNIQUE set of "microscopic" nuclear ingredients

![](_page_35_Figure_1.jpeg)

#### **Fission properties and the r-process**

![](_page_36_Figure_1.jpeg)

Accumulation of material at the N=184 closed shell  $\rightarrow$  Special emphasis on the Fission Fragment Distribution for the  $A\sim278$  isobars

## **Fission Fragment Distributions**

#### **GEF (ABLA): semi-empirical macroscopic-microscopic scission point model**

- microscopic properties are essentially determined by the shell effects of the fragments
- only the macroscopic properties of the fissioning system are taken into account.
- inclusion of dynamical effects to include the impact of inertia along the fission trajectory
   Schmidt and Jurado (2010-2012)

#### SPY: Scission Point model based on Gogny potential energy surfaces and NLD

- ONLY based on fission fragments & first-chance fission
- Evolution (quasi-static) from saddle to scission point is neglected
- Isolated fragments at rest
- Well-defined fragments properties  $(Z, N, \beta)$
- Fragmentation probability  $\propto$  number of available states

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Lemaitre et al. (2015, 2018)
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#### **SPY** vs **GEF** predictions of the FFD

![](_page_38_Figure_1.jpeg)

 $A: 70 \longrightarrow 190$ 

#### Sensitivity of dynamical composition to the fission fragment distribution

along the A=278 isobar (from the N=184 closed shell)

![](_page_39_Figure_2.jpeg)

#### Gogny-HFB calculation of <sup>278</sup>Cf PES

![](_page_40_Figure_1.jpeg)

Missing dynamical calculations (Goutte et al. 2005, Bernard et al. 2011, Scamps et al. 2018) (e.g. the time-dependent GCM; cf talks Regnier & Pillet )

#### Mean number of emitted neutrons

![](_page_41_Figure_1.jpeg)

- during n-irradiation:
  - Neutrons emitted are rapidly recaptured by existing nuclei
    - $\rightarrow$  no major impact on the final abundance distribution
- after n-irradiation:
  - Neutrons add up to  $\beta$ -delayed neutrons
  - Neutrons are captured by the most abundant species ( $N \sim 126$ )
  - Possible shift of the abundances by a few units

Impact of the neutron emission on the r-process on the composition of the dynamical NS-BH ejecta

![](_page_42_Figure_1.jpeg)

## Conclusions

The role of fission remains important: 10 to 100% of ejected matter experience fission in compact object mergers but quantitatively conclusions remain sensitive to the complex and still uncertain

- r-process site & weak interaction processes
  - define the nbr of free neutrons available
  - possible actinide/SH production
  - fission recycling & heating
- prediction of fission probabilities (sf,  $\beta$ df, nif) & FFD !
  - potential energy surface for ~2000 nuclei
  - nuclear level densities at saddles/wells
  - fission fragment distributions
  - neutrons emitted

![](_page_43_Figure_11.jpeg)

![](_page_43_Figure_12.jpeg)