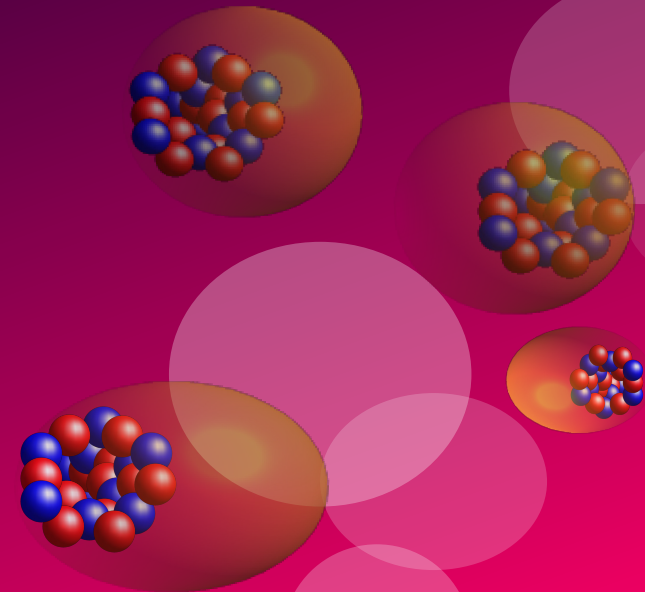
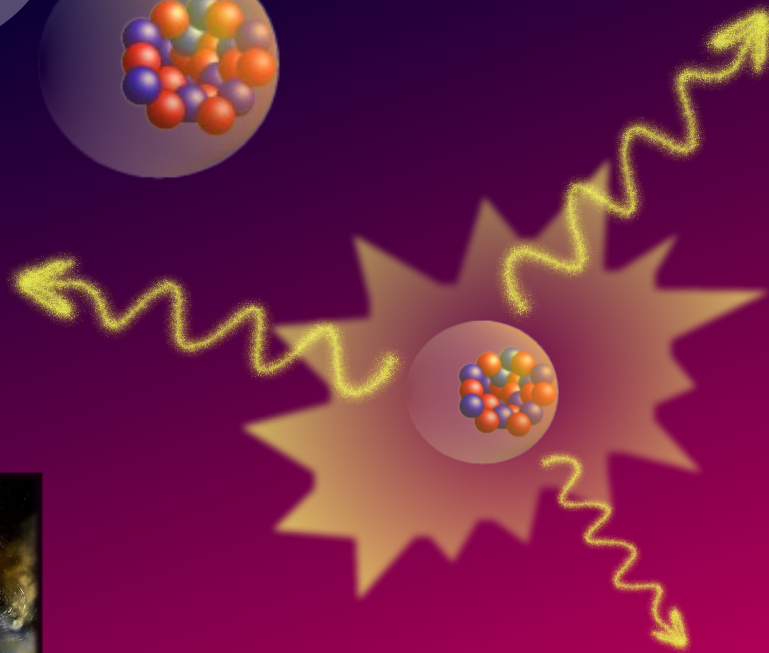
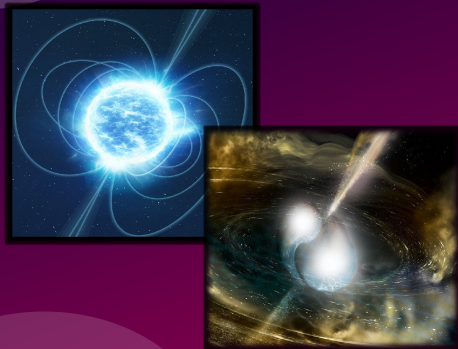
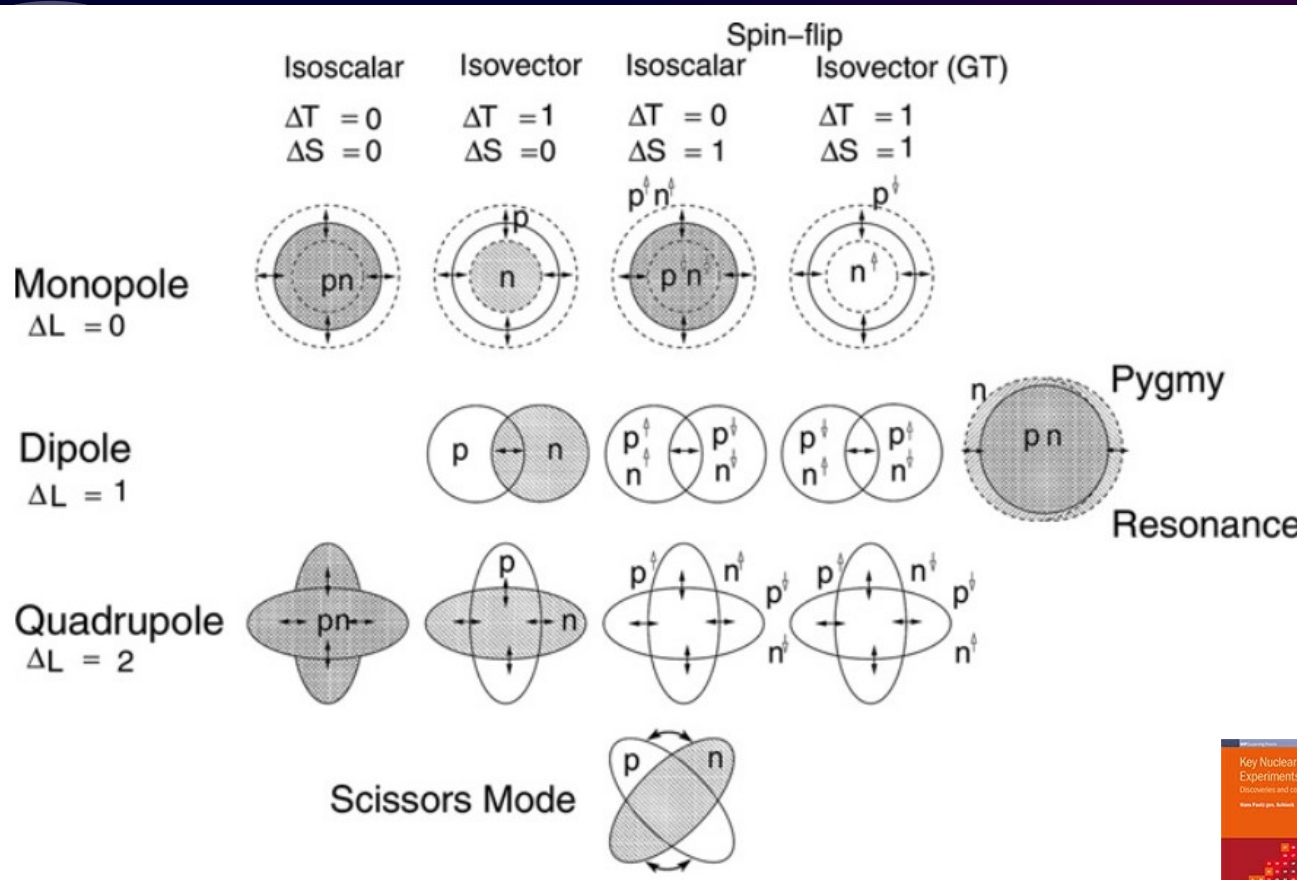


# Study of Pygmy Dipole Resonances and their potential role in nuclear astrophysics

Iolanda Matea Macovei (IJCLab)



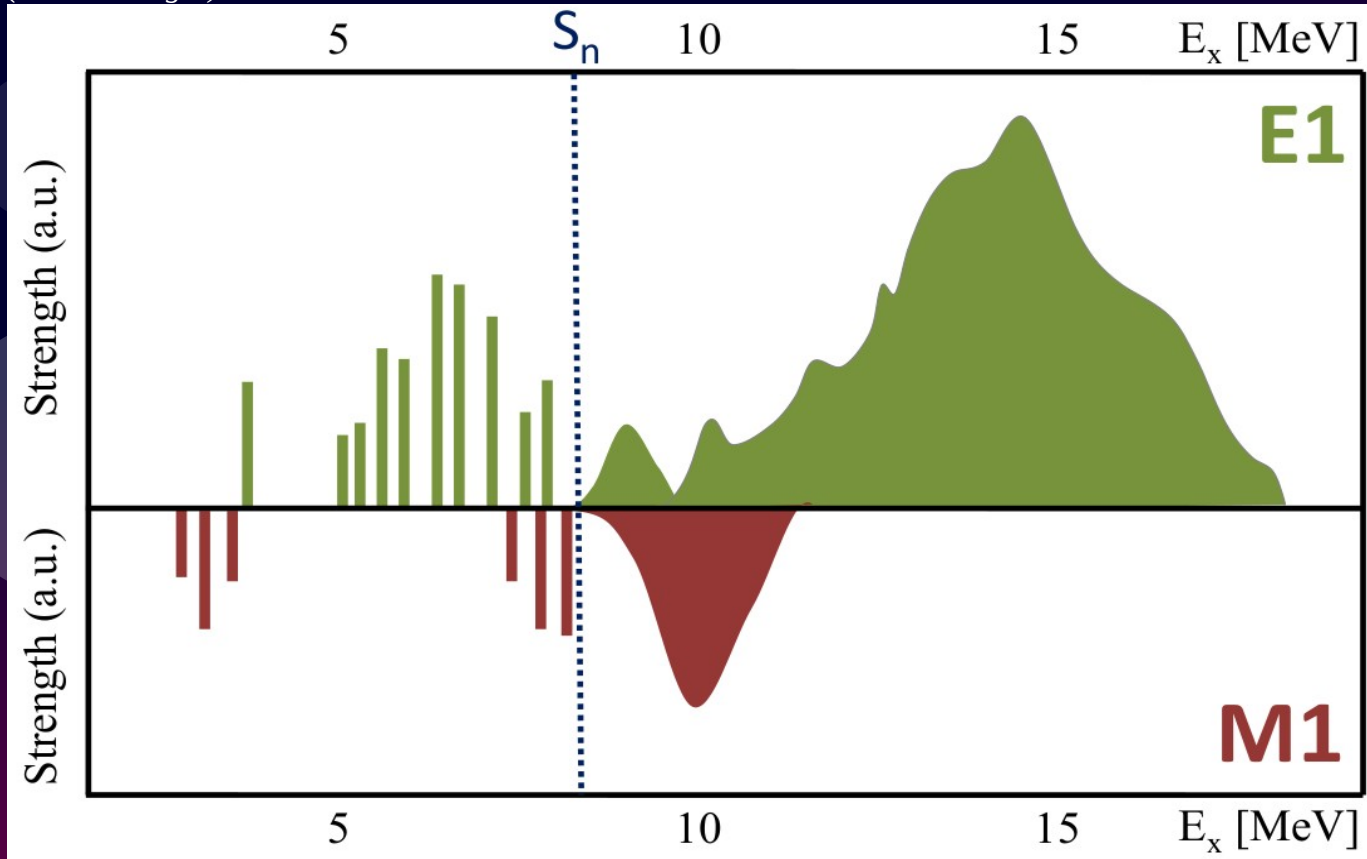
## $\Delta T, \Delta S, \Delta L$



$\Delta L = 1 \rightarrow$

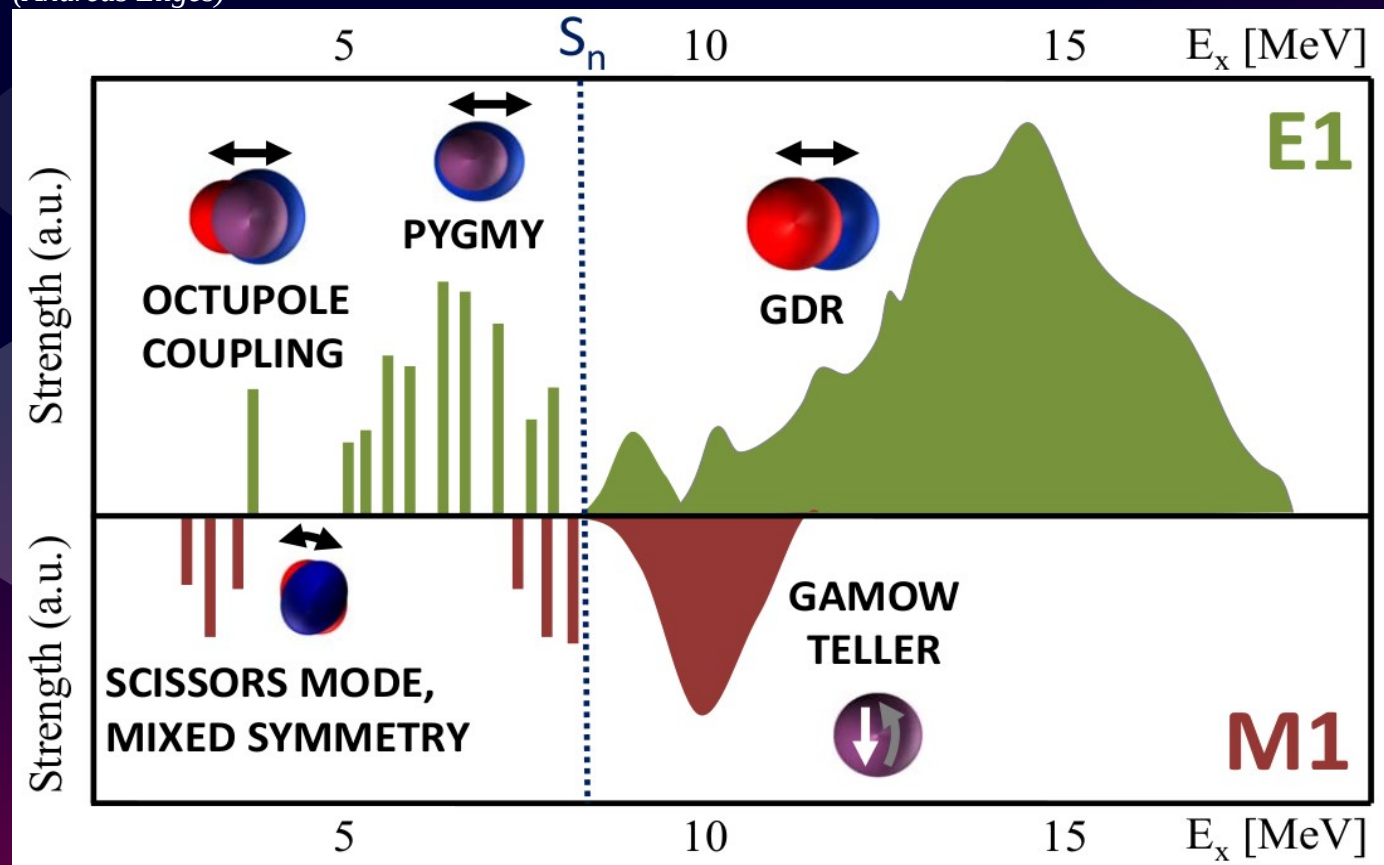


(Andreas Zilges)

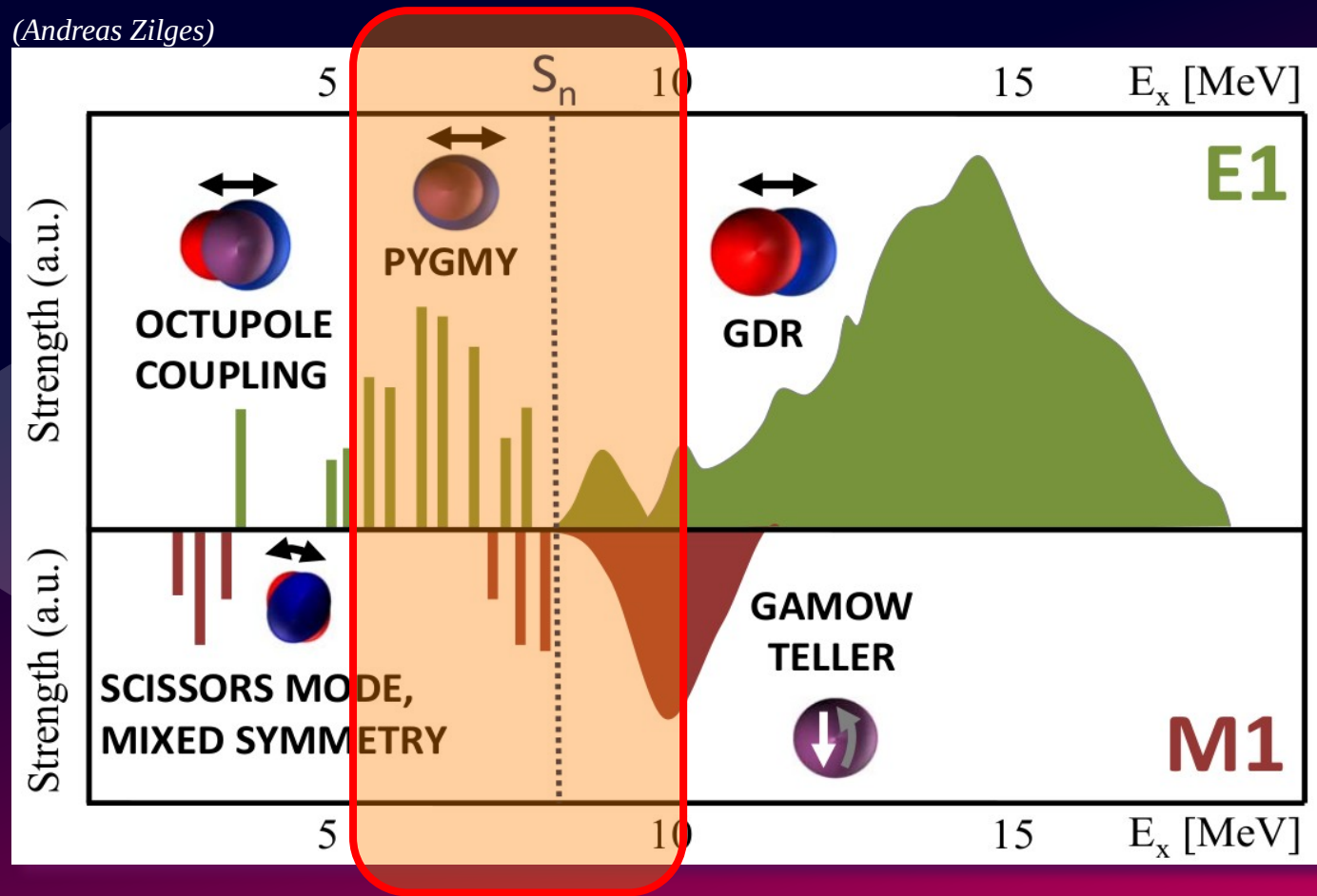


This strength distribution is related to the average probability of  $\gamma$ -ray absorption or emission as a function of  $\gamma$ -ray energy :  $\gamma$ SE

(Andreas Zilges)



(Andreas Zilges)



Only neutron-rich nuclei will be considered throughout the lecture

## PDR

(n, $\gamma$ ) capture

$E^*$  and evolution  
with N/Z ratio

Collective vs.  
single particle  
structure

Excitation x-sections  
with different probes

B(E1) distribution  
(EWSR)

Temperature  
dependence

Dipole  
Polarizability

Isospin character

$\gamma$ -decay branching  
above  $S_n$

*A.Bracco et al, PPNP106 (2019)*  
*E.G. Lanza et al, PPNP129 (2023)*  
*D. Savran et al, PPNP70 (2013)*  
*N. Paar et al, Rep.Prog.Phys70 (2007)*

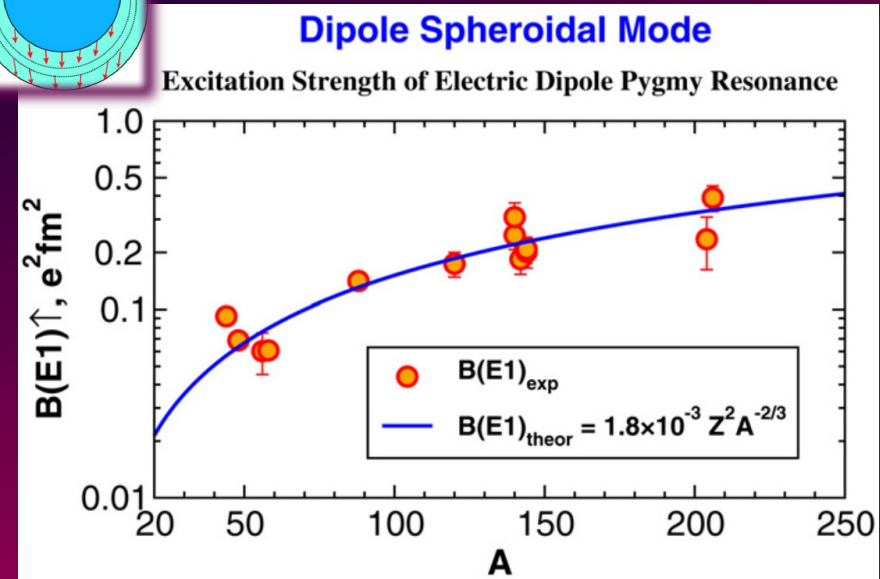
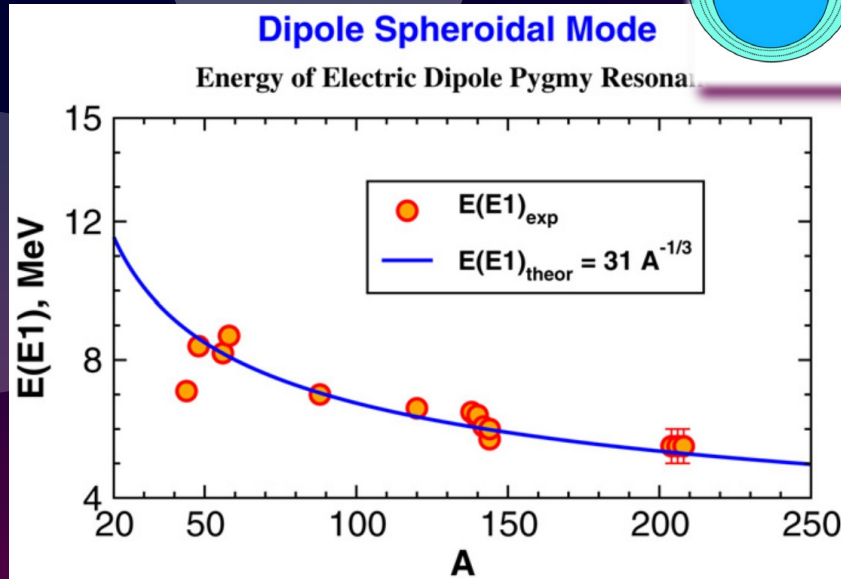
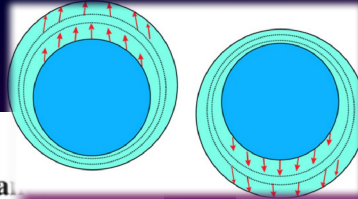
- linear displacement of core/neutron excess density distributions
- density distributions : Wood-Saxon
- restoring force : np interaction

$$B(E1)_{PDR} = \frac{(N - N_c) Z \cdot E_{GDR}}{(Z + N_c) N E_{PDR}} \cdot B(E1)_{GDR}$$

$$E_{PDR}^* = \left[ \frac{Z}{3(Z + N_c)} \right]^{1/2} \cdot f(a_n, a_p, R_n, R_p) \cdot E_{GDR}^*$$

*P. van Isacker et al, PRC45 (1992)*

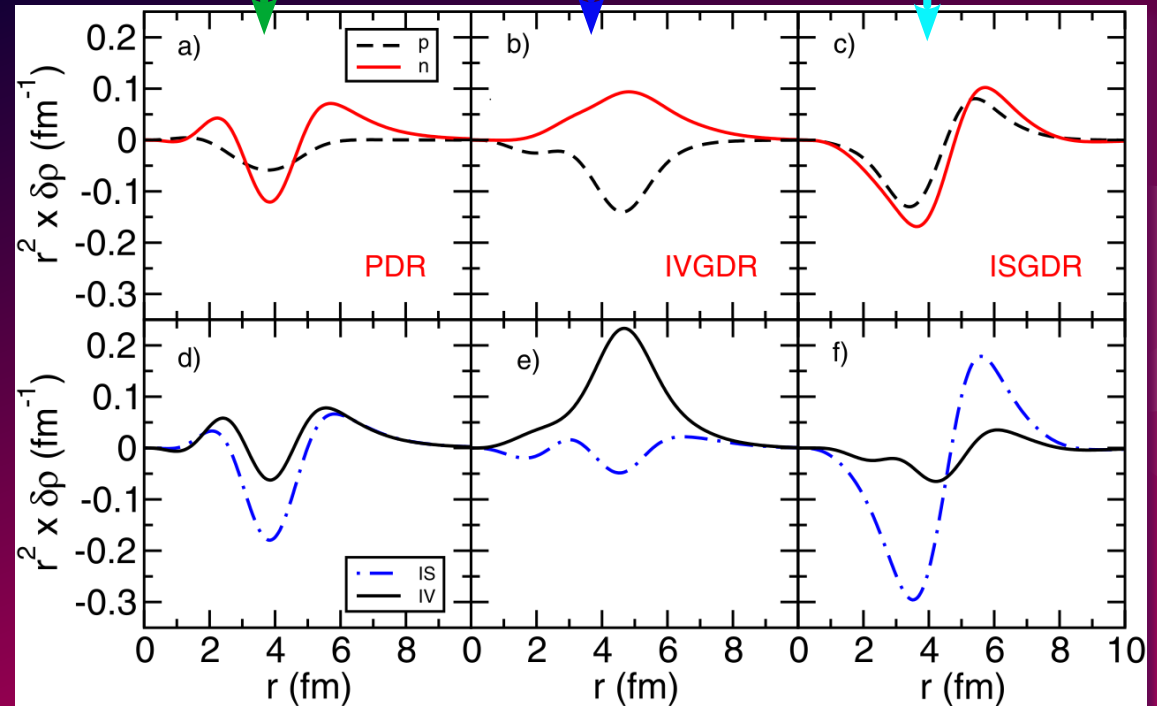
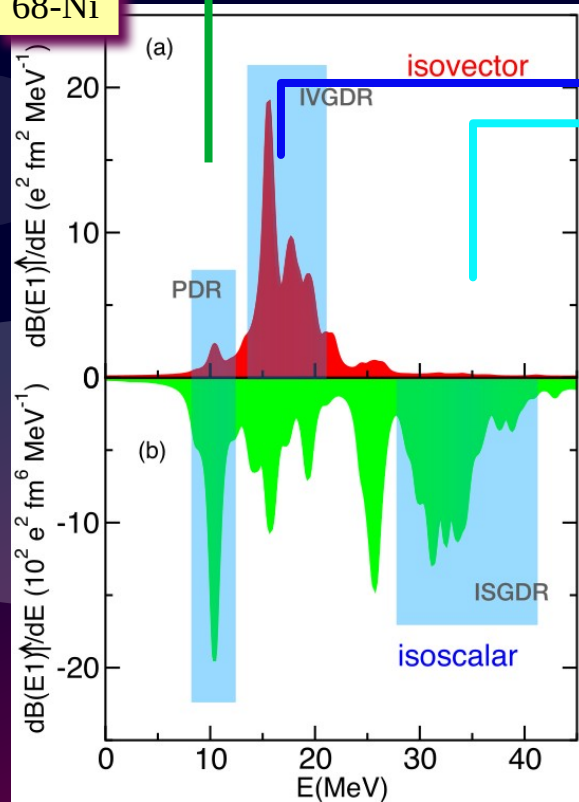
- elasto-dynamic excitation of the electric PDR
- elastic layer oscillation against elastic continuous core



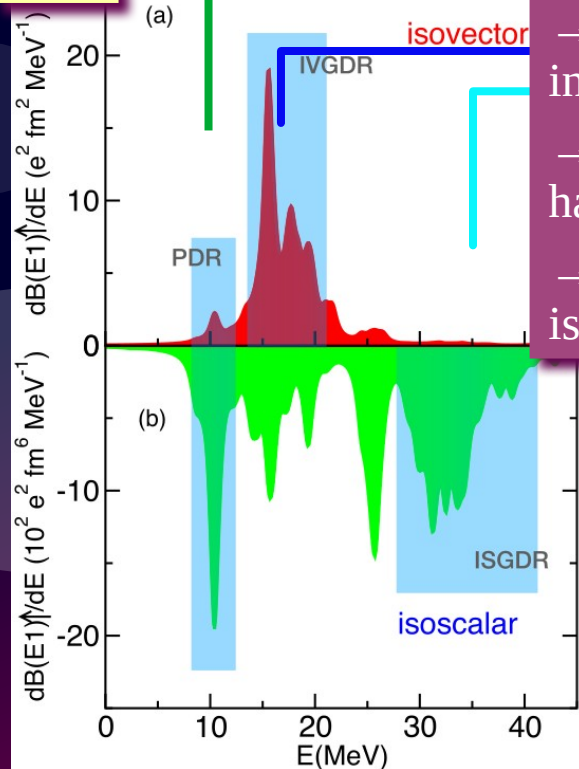
*S.I. Bastrokov et al, PLB664 (2008)*



68-Ni

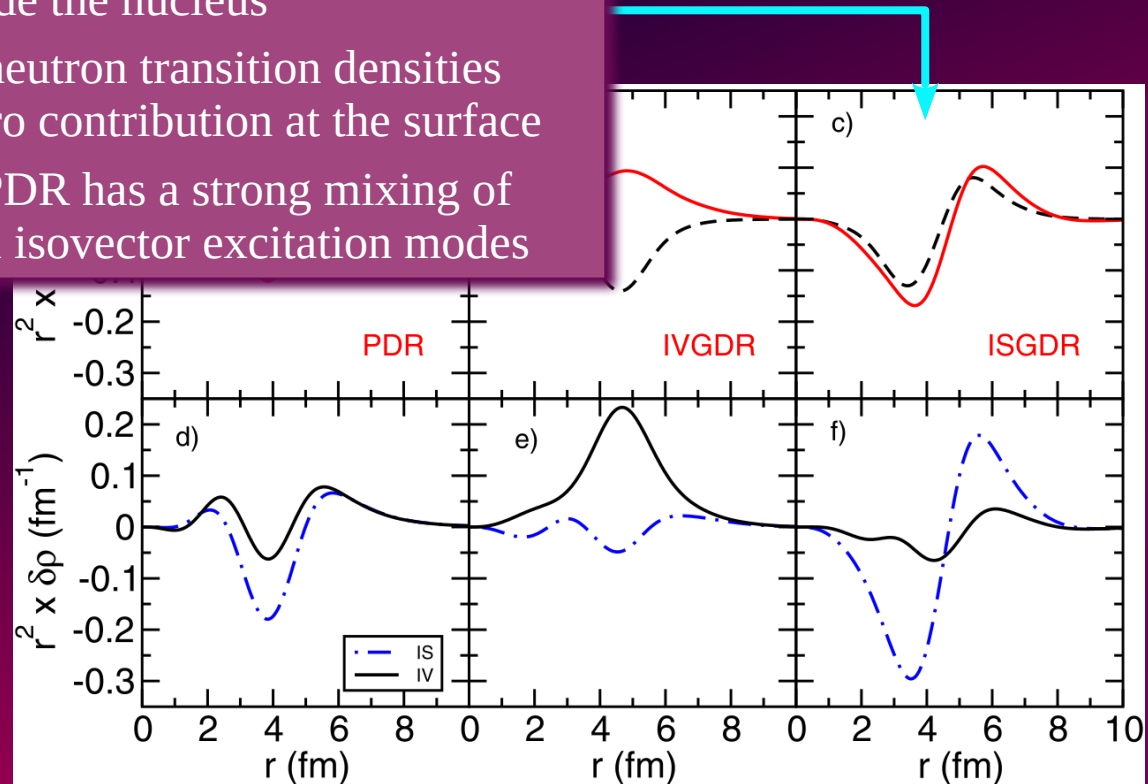


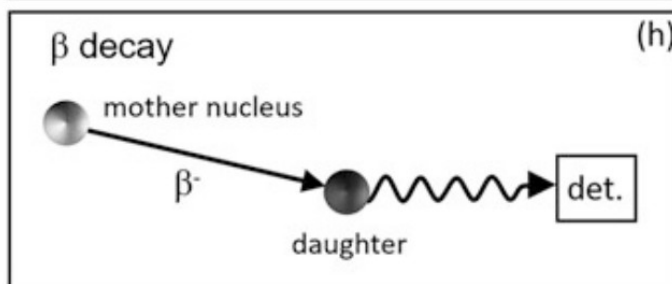
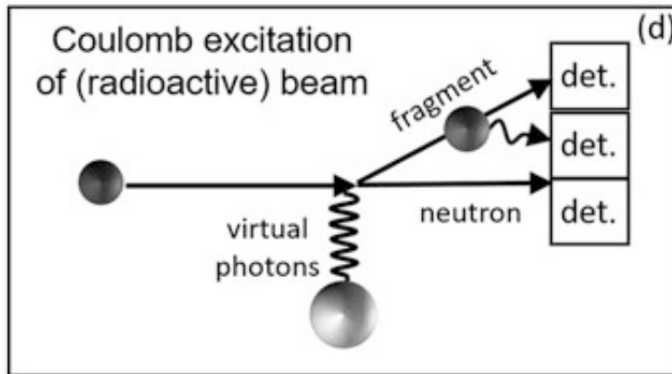
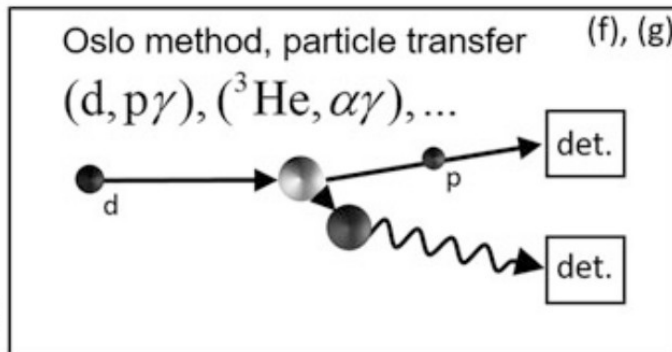
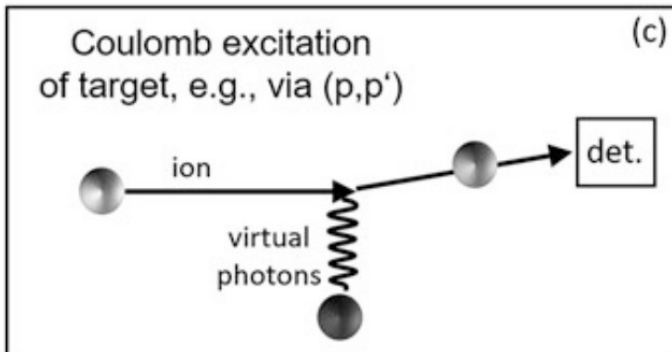
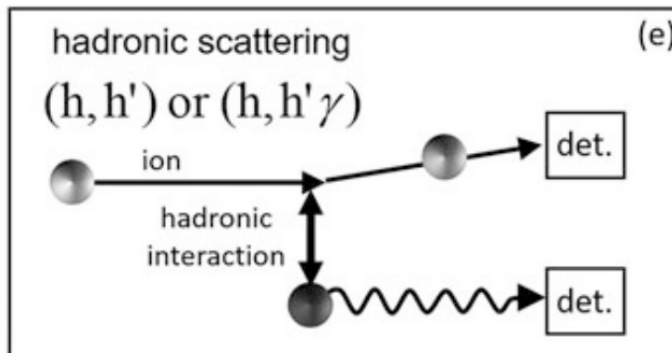
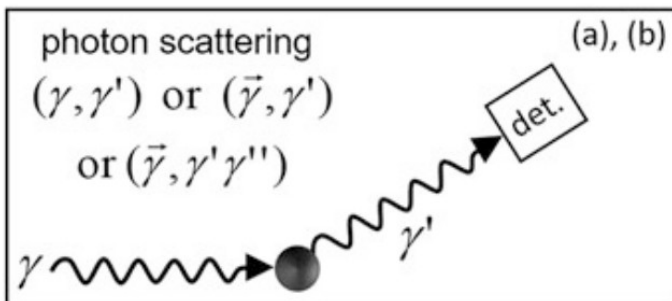
**<sup>68</sup>Ni**



## Definition of PDR:

- proton/neutron transition densities are in phase inside the nucleus
- only the neutron transition densities have non-zero contribution at the surface
- as such, PDR has a strong mixing of isoscalar and isovector excitation modes

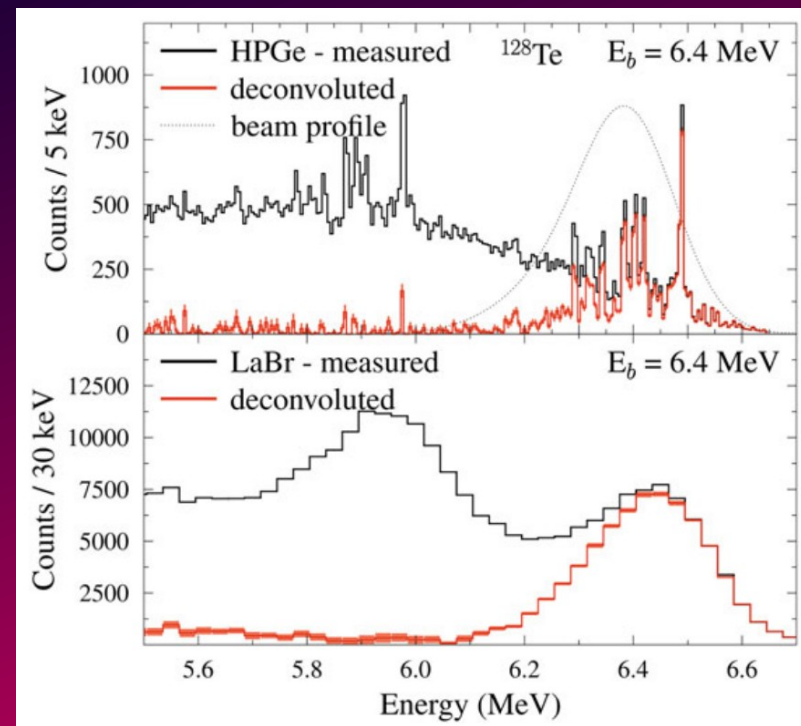
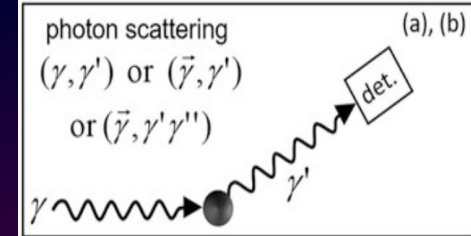




Selected remarks and examples follow ...

# Real photon scattering

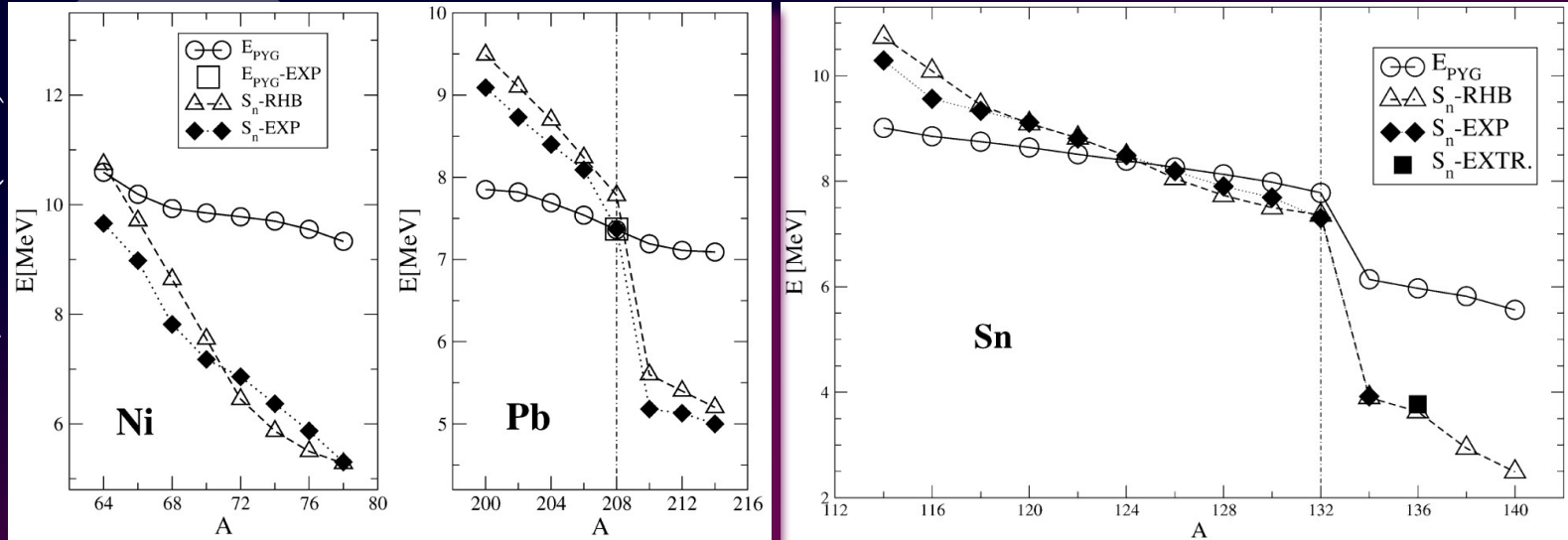
- L=1 excitations dominant
- photon : EM interaction, under control
  - NRF (nuclear resonance fluorescence) with bremsstrahlung  $\gamma$
  - tagged photon NRF : lower bandwidth
  - laser Compton back-scattered  $\gamma$  beam (LCB)
    - small bandwidth + polarized
- in-volume interaction
- iso-vector probe
- measure the strength up to  $S_n$
- limited to stable nuclei
- detection usually HPGe with  $W(\theta, \varphi)$  measure =>  
clear measurement L and  $\pi$  (polarized) of  $\gamma$ -radiation



*J. Isaak et al, PRC103 (2021)*

One remark about  $(\gamma, \gamma')$  experiments :

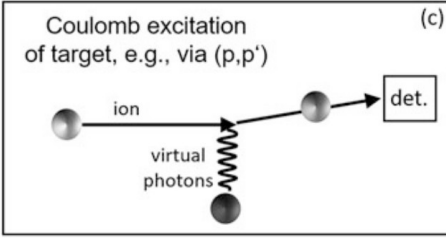
N. Paar et al., PLB606 (2005)



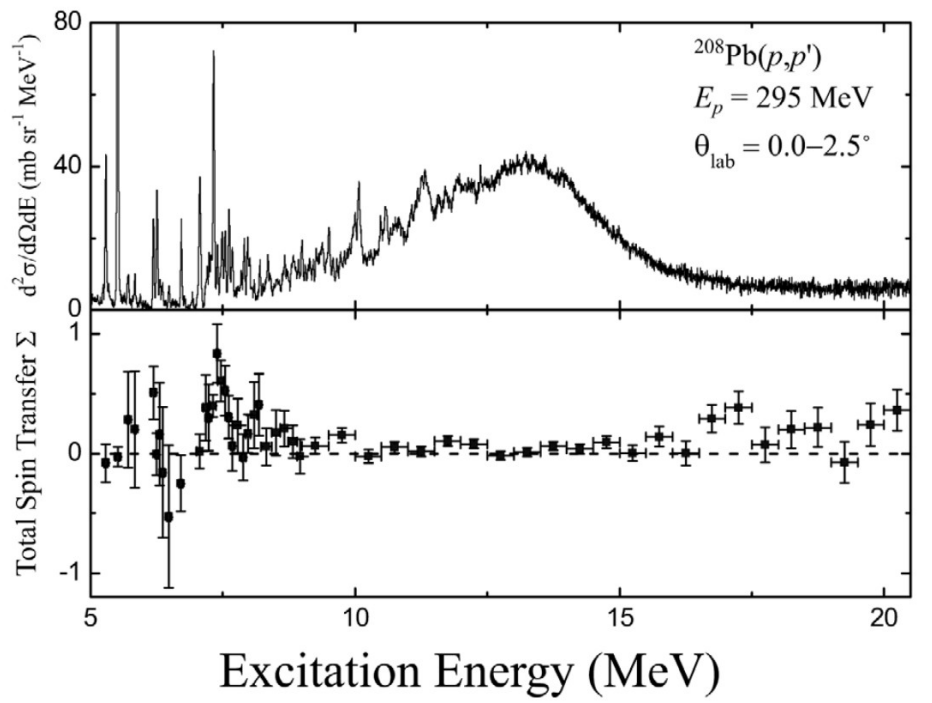
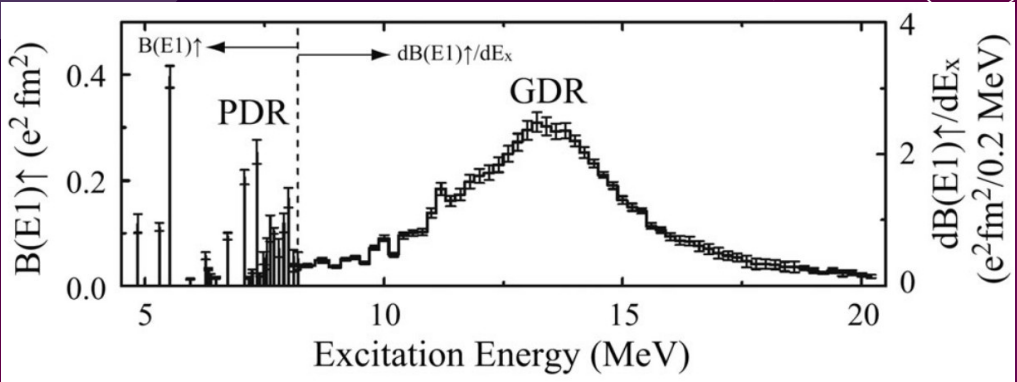
PDR is well above  $S_n$  in exotic n-rich nuclei  $\Rightarrow$   $\gamma$ -decay is too slow ...

# Virtual photon scattering

- mixed  $L=1,2,3\dots$  excitations
- (polarized) proton at  $\sim 300\text{-}400\text{MeV}$  scattered at extreme forward angles
  - detection of the scattered proton => measure the strength below and above  $S_n$
  - $L=1$  excitations favored, but not only ... MDA needed.
  - $\pi$  measurements (polarized proton beam)
- in-volume interaction
- iso-vector probe
- limited to stable nuclei

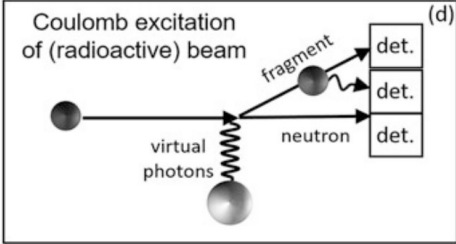


A. Tamii et al, EPJA50 (2014)

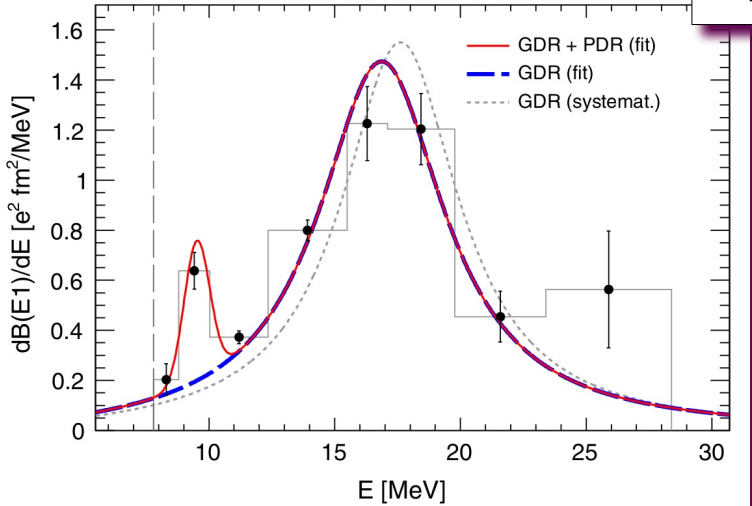


# Virtual photon scattering

- radioactive nuclei
- inverse kinematics => less resolved  $\gamma$  spectra
- lower statistics
- relativistic beam energies => L=1 dominates
  - can measure also above  $S_n$  (if neutrons are detected)

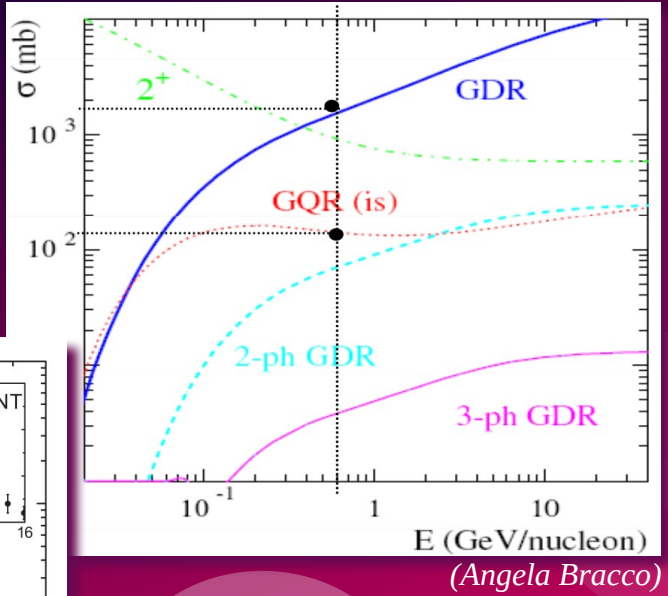
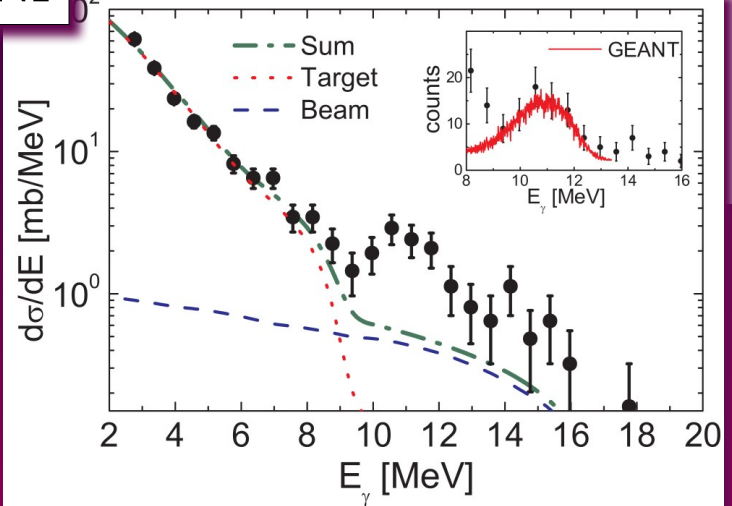


*D.M. Rossi et al, PRL111 (2013)*



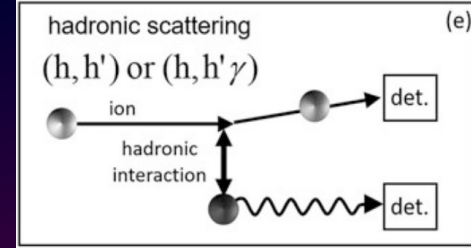
$^{68}Ni$

*O. Wieland et al, PRL102 (2009)*



# Different isospin response

- e.g.,  $\alpha$  or  $^{12}\text{C}$   $\rightarrow$  iso-scalar response
- non selective population of states; to pin down the E1 character  $\rightarrow$   $(h, h'\gamma)$
- function of the energy of the probe:
  - 30MeV neutrons : iso-vector
  - $\sim 1\text{GeV}$  neutrons : iso-scalar
- surface interaction (ignoring the charge)
- most relevant studies were done below  $S_n$
- direct or inverse kinematics  $\Rightarrow$  stable and radioactive nuclei

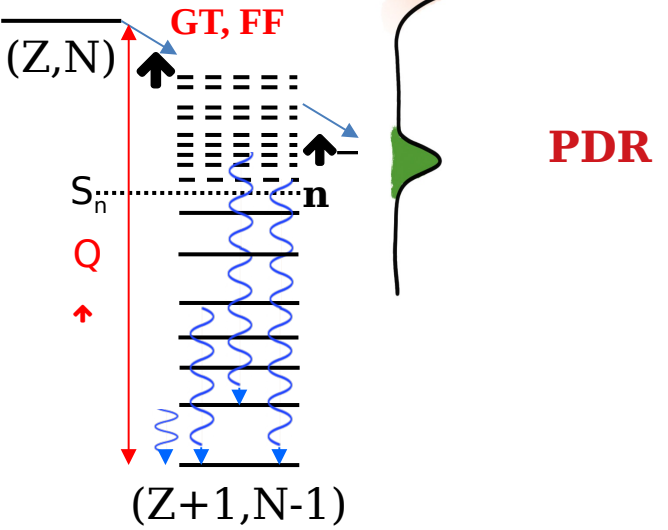


**Mainly for the PDR structure studies**  
(see Perine's poster)

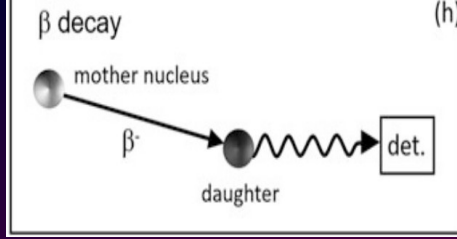


# β-decay

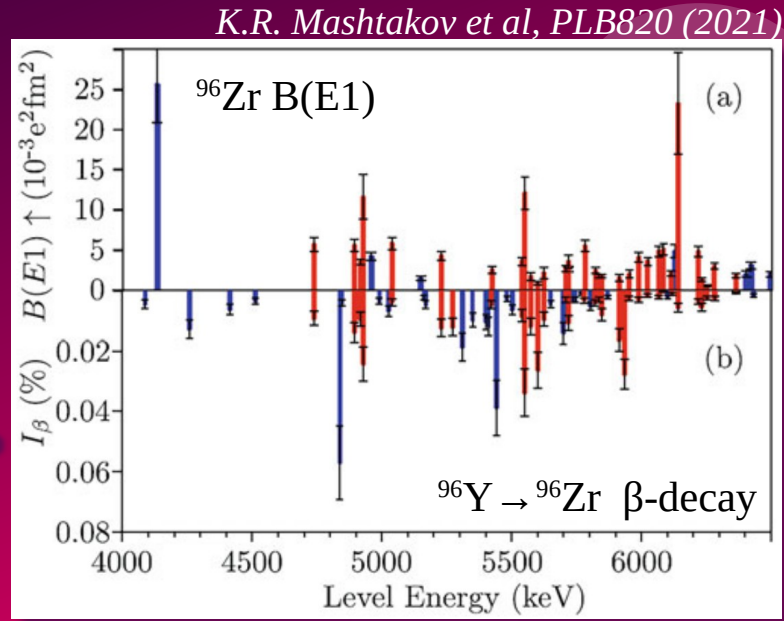
## E1 Strength Distribution



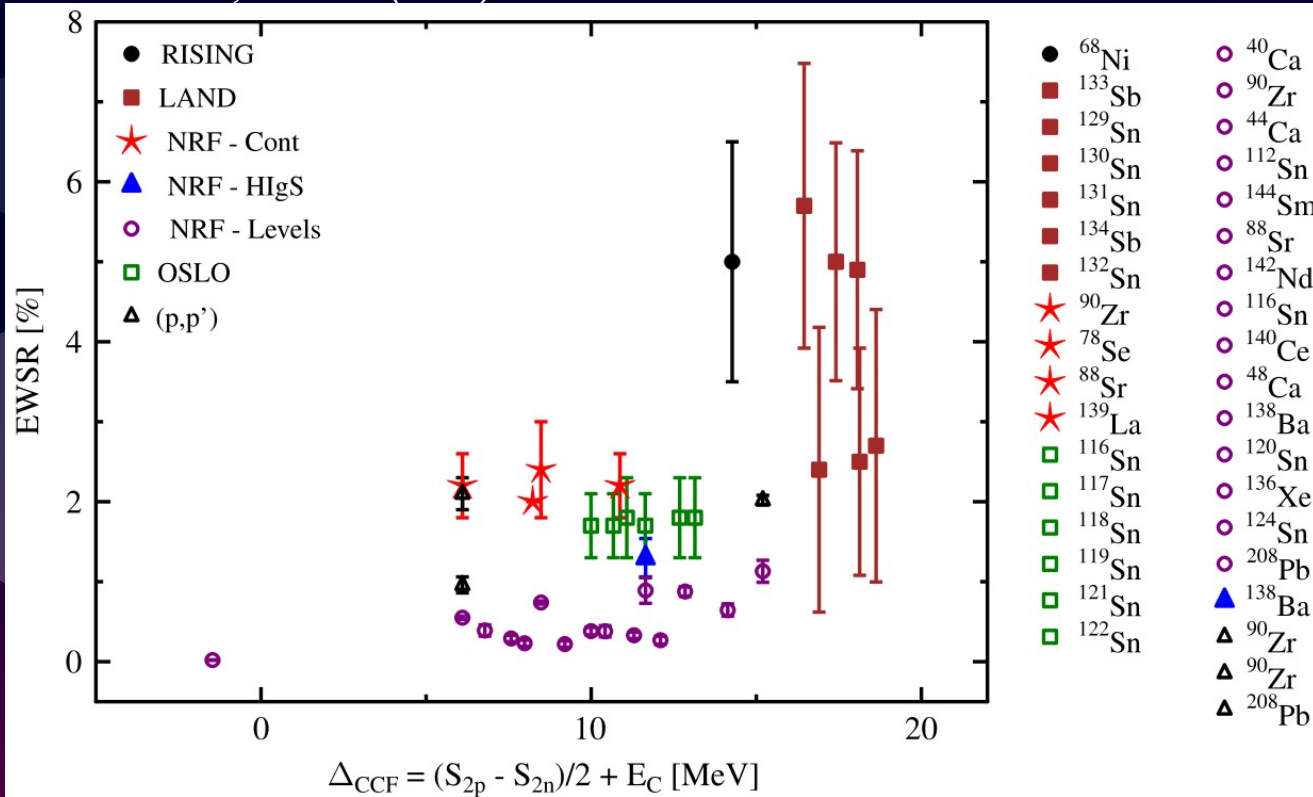
- weak interaction selection rules (GT, FF)
  - selective population of states around  $S_n$
- very clean  $\gamma$  and  $(\gamma, \gamma)$ -coincidence spectra
  - but (very) low statistics
- access to very exotic nuclei, but a very partial picture
  - beta decay of  $^{82}\text{Ga}$  (L. Al Ayoubi, PhD Univ. Paris Saclay, 2023)
  - beta decay of  $^{80}\text{Ga}$  (R. Li, PhD Univ. Paris Saclay, 2022)



Exemple of NRF vs  $\beta$ -decay measurements : controlled interpretation available only on few stable nuclei



D. Savran et al, PPNP70 (2013)

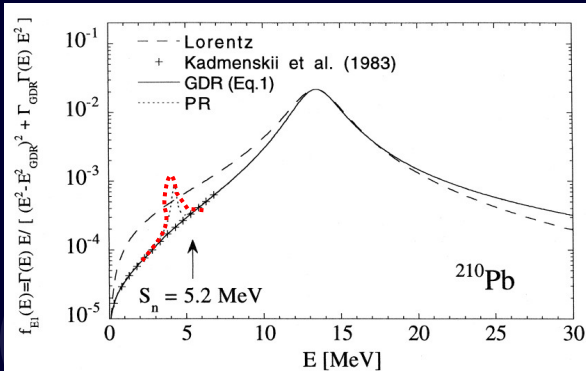


$E^* \sim 5-10$  MeV

EWSR (% of TRK SR)  $\sim 1-5\%$

(TRK SR =  $14.8 \cdot (NZ)/A$  ( $e^2\text{fm}^2\text{MeV}$ ))

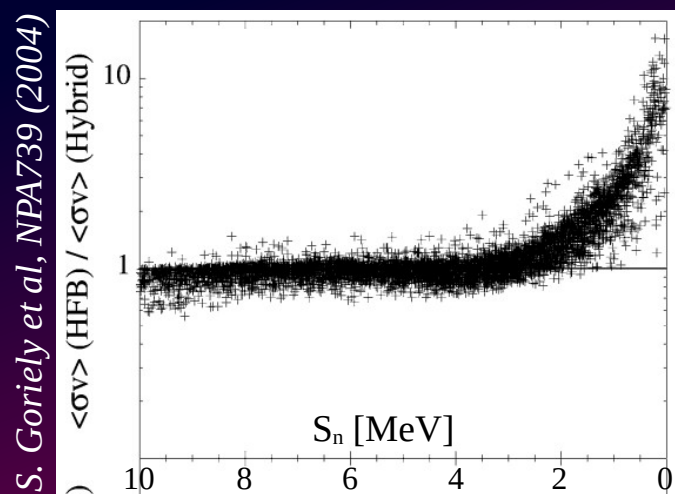
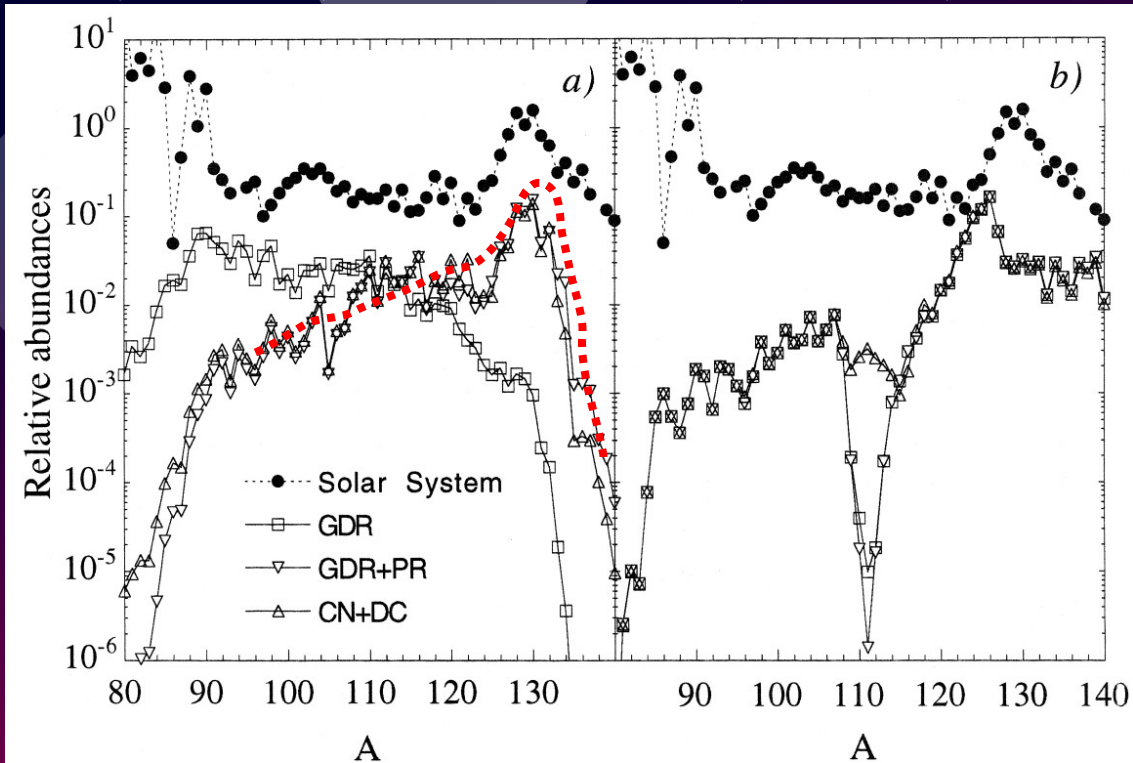
# (n, $\gamma$ ) capture



S. Goriely, PLB436 (1998)

$T=10^9$  K;  $N_n = 10^{20}$  cm $^{-3}$ ;  $t_{irr} = 2.4$  sec

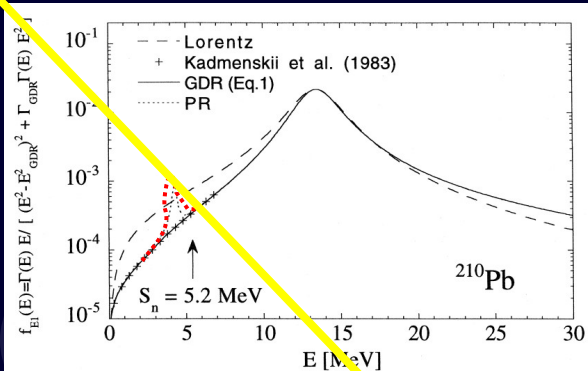
$T=1.5 \times 10^9$  K;  $N_n = 10^{28}$  cm $^{-3}$ ;  $t_{irr} = 0.3$  sec



S. Goriely et al, NPA739 (2004)

- under certain conditions ( $T$ ,  $N_n$ , time), PDR (few % of TRK) is able to produce the  $A \sim 130$  peak in relative abundances where Lorentz-shaped GDR can't
- microscopic models (HFB+QRPA) for E1 distribution in nuclei predict an even more pronounced modification in the (n, $\gamma$ ) x-sections for very exotic nuclei
- closer to stability, low energy E1 strength can also influence locally the abundance of i-process nuclei (M. Markova et al, PRC109 2024)

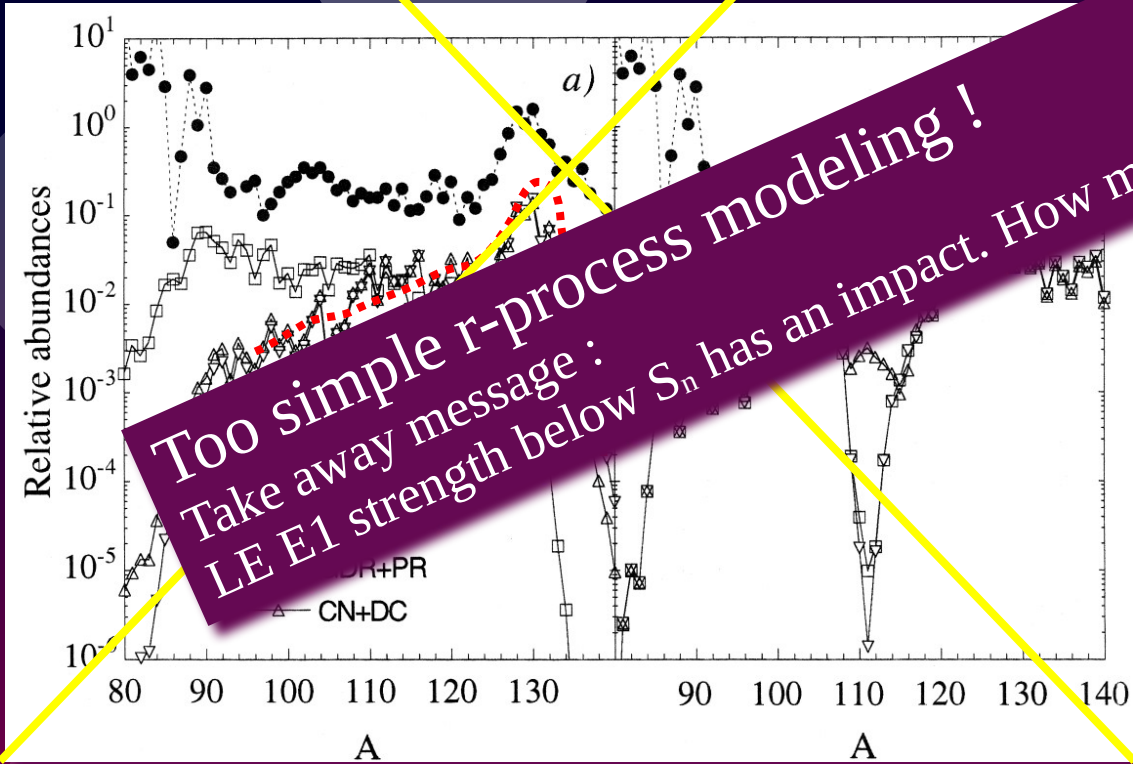
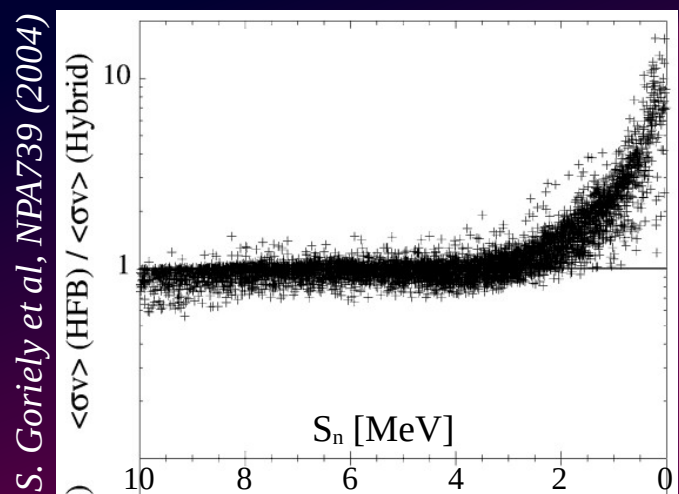
# (n,γ) capture



S. Goriely, PLB436 (1998)

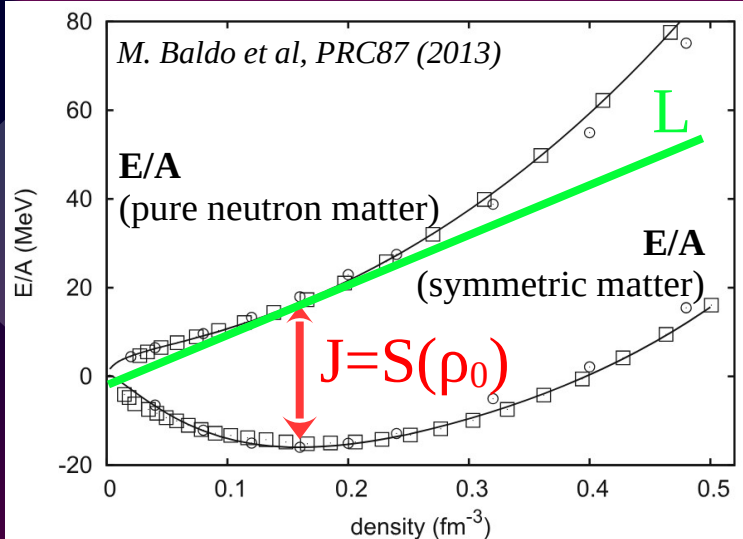
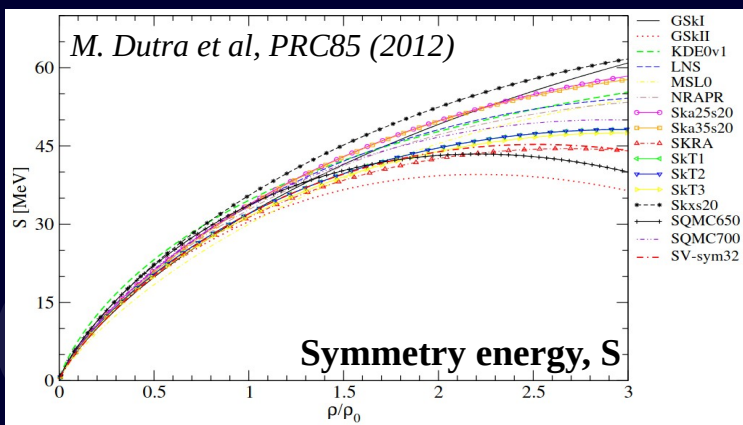
$T=10^9 \text{ K}$ ;  $N_n = 10^{20} \text{ cm}^{-3}$ ;  $\text{tirr} = 2.4 \text{ sec}$

$T=1.5 \times 10^9 \text{ K}$ ;  $N_n = 10^{28} \text{ cm}^{-3}$ ;  $\text{tirr} = 0.3 \text{ sec}$



Under certain conditions ( $T$ ,  $N_n$ , time), PDR (few % of TRK) is able to produce the  $A \sim 130$  peak in relative abundances where Lorentz-shaped GDR can't

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symmetry energy

$$\frac{E}{A}(\rho, \delta) = \frac{E}{A}(\rho, 0) + \underbrace{S(\rho)}_{\text{symmetry energy}} \cdot \delta^2 + \mathcal{O}(\delta^4) \quad \text{with } \delta = \frac{\rho_n - \rho_p}{\rho}$$

energy per nucleon of symmetric NM

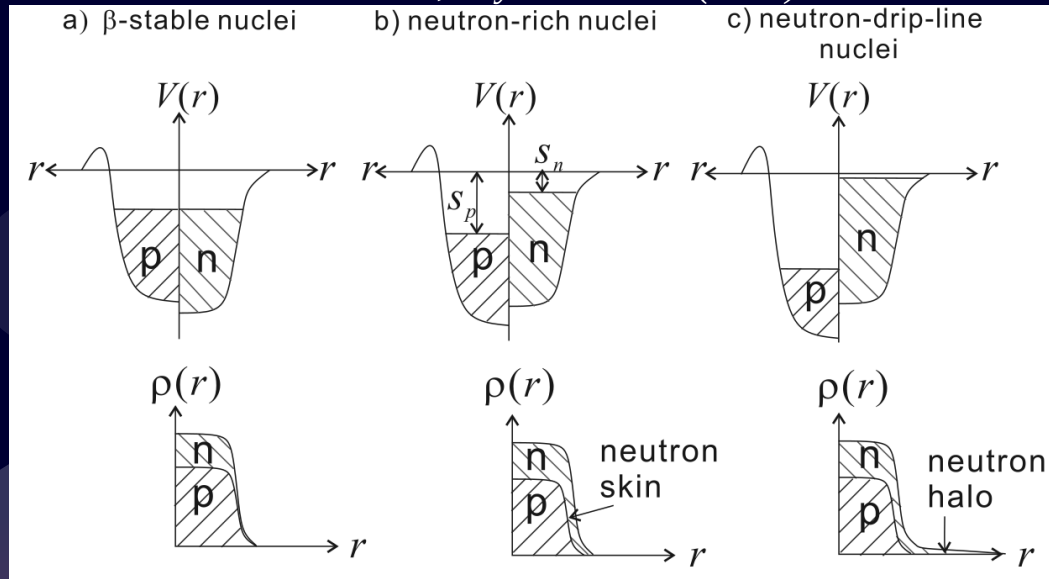
$$\frac{E}{A}(\rho, 0) = \frac{E}{A}(\rho_0, 0) + \frac{1}{2} k \cdot \varepsilon^2 + \mathcal{O}(\varepsilon^3) \quad \text{with } \varepsilon = \frac{\rho_0 - \rho}{3\rho_0}$$

$$S(\rho) = S(\rho_0) - L \cdot \varepsilon + \frac{1}{2} k_{\text{sym}} \cdot \varepsilon^2 + \mathcal{O}(\varepsilon^3)$$

**Astrophysical systems (like NS) require the knowledge of EoS of asymmetric matter, related to the isovector parameters of the symmetry energy : J and L**

**The natural choice for finding physical quantities that could be sensitive to the symmetry energy is to look at the iso-vector elementary nuclear excitations**

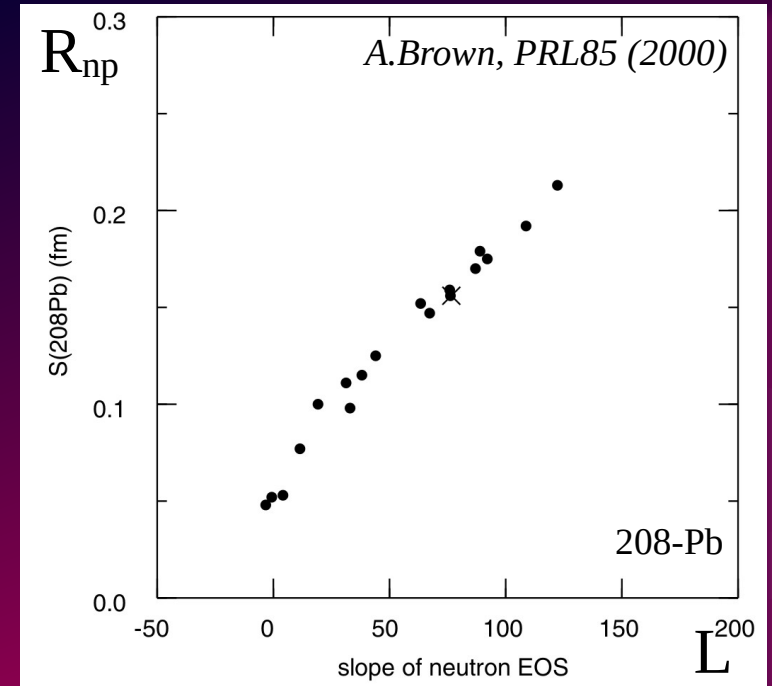
**GDR and iso-vector part of the PDR**



$\sim 8 \text{ MeV}$   $S_n$  (neutron separation energy)  $\sim 0 \text{ MeV}$

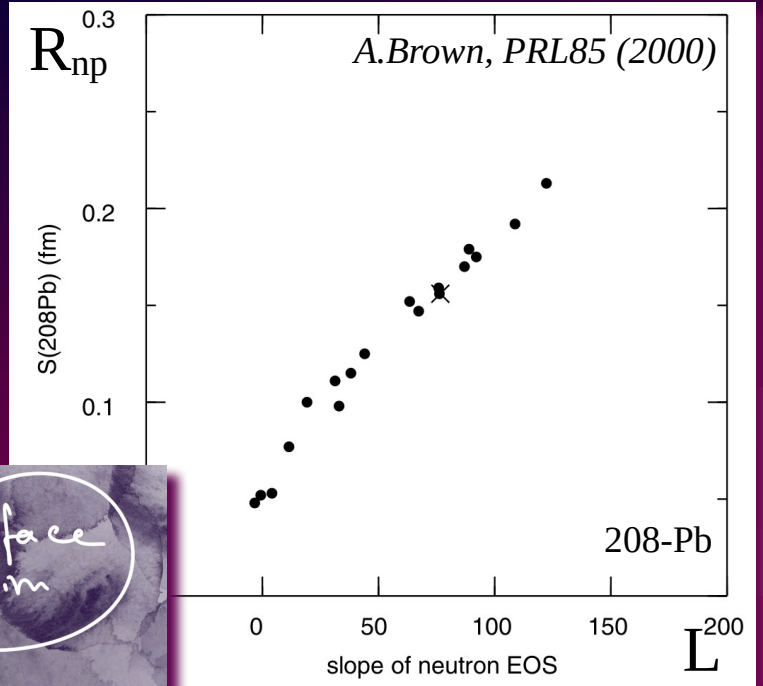
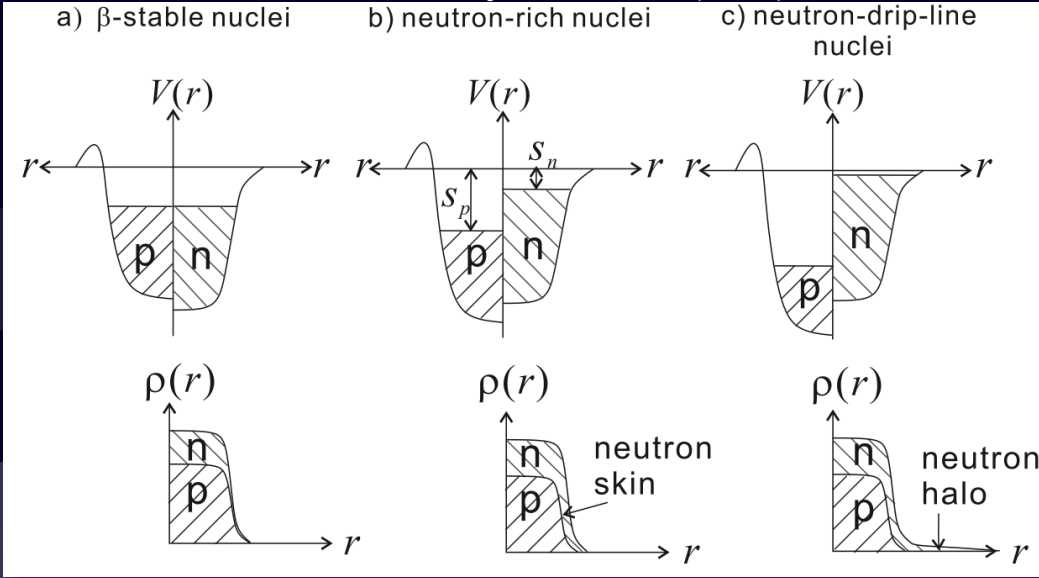
# Connection to EoS and NS

Neutron skin ( $R_{np}$ ) and  $L$

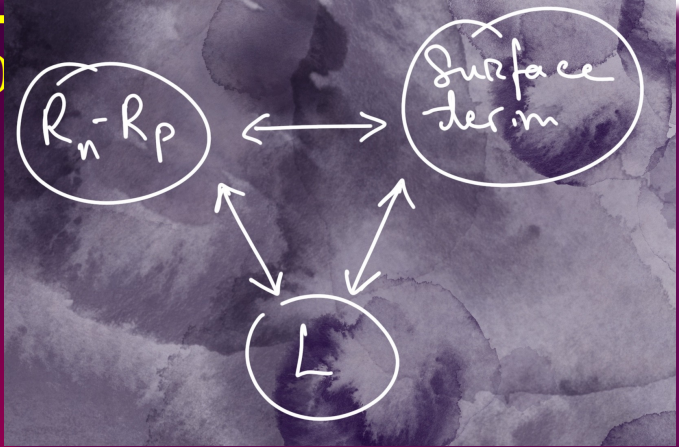


**PDR is related with neutron skin !**

Neutron skin ( $R_{np}$ ) and  $L$



~8 MeV  $S_n$  (neutron separation energy)



related with neutron skin !



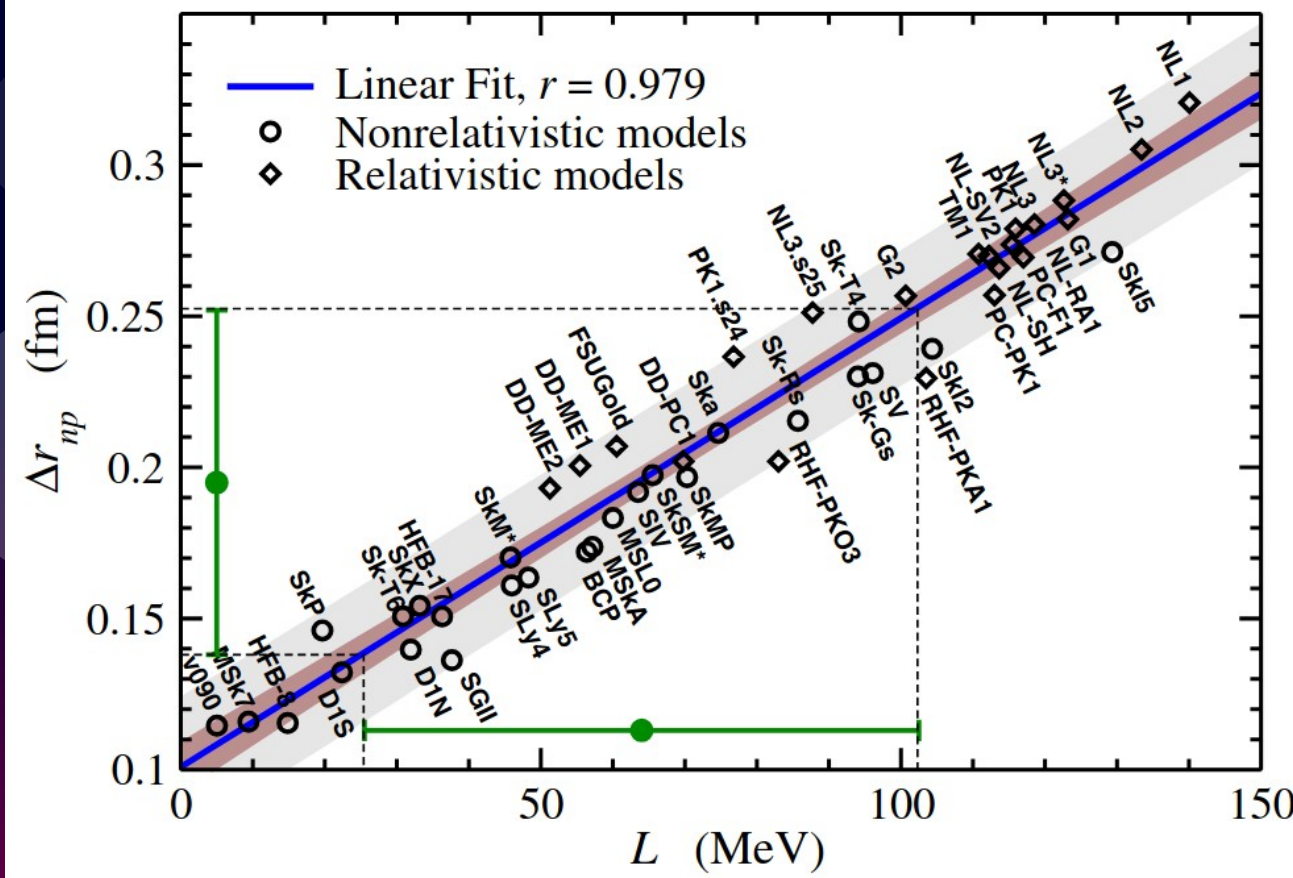
### PREX experiment

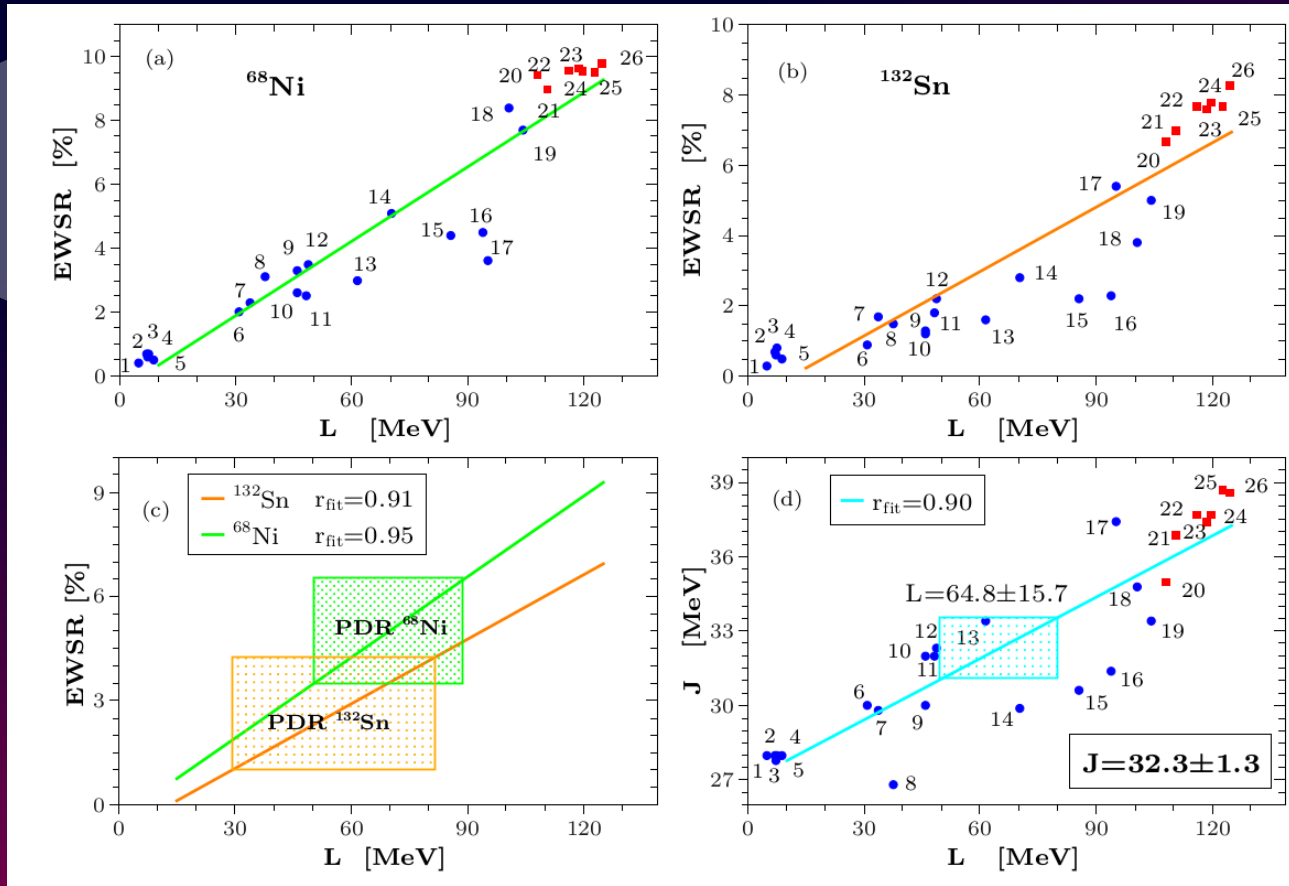
- electroweak parity-violating e scattering
- asymmetry measurement
- $\Delta r_{np} (^{208}\text{Pb}) = 0.283(71)$  fm – PREX II

BUT :

tension with polarizability measurements  
and with the CREX measurement  
(not the end of the story ...)

X. Roca-Maza et al, PRL106 (2011)



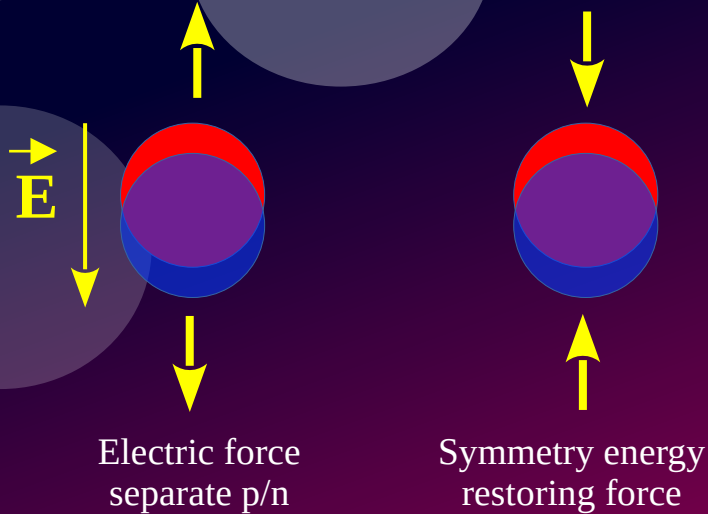


A. Carbone et al, PRC81 (2010)

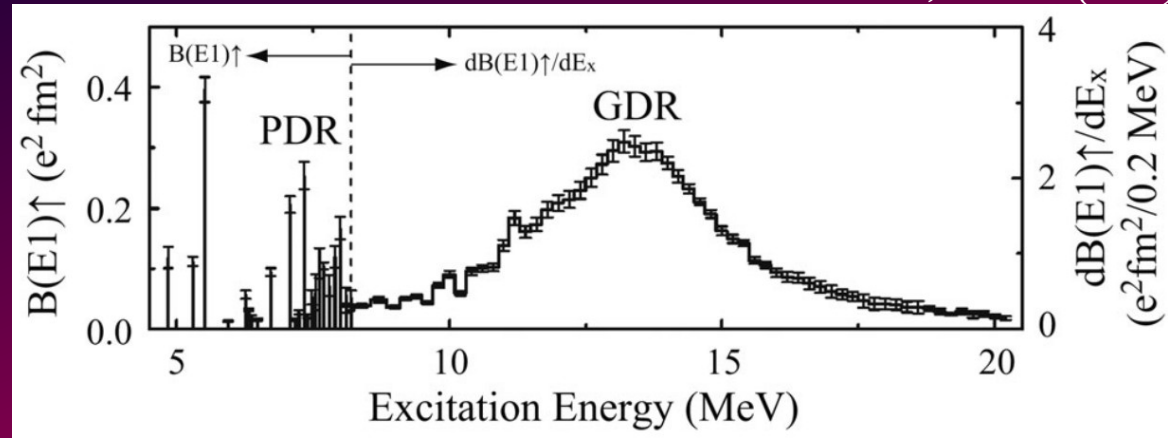
# Connection to EoS and NS

## Dipole Polarizability and (L,J)

$$\alpha_D = \frac{\hbar c}{2\hbar^2} \int \frac{G_{\text{ass}}^{E1}}{\omega^2} d\omega = \frac{8\pi e^2}{9} \int \frac{dB(E1)}{\omega}$$



(figure inspired by A. Tamii)

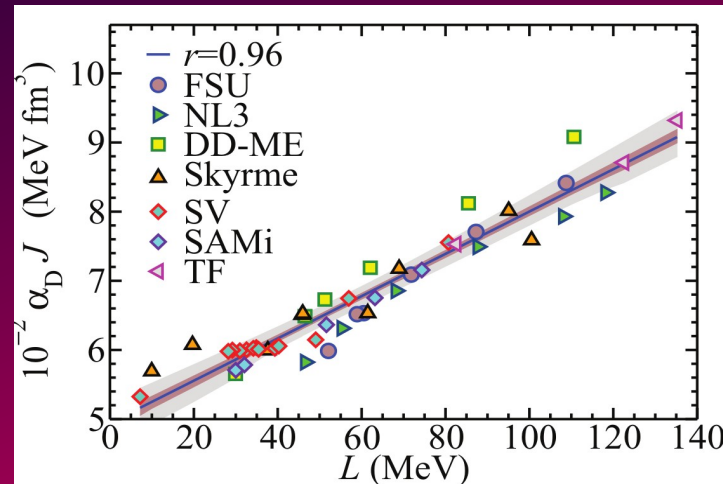
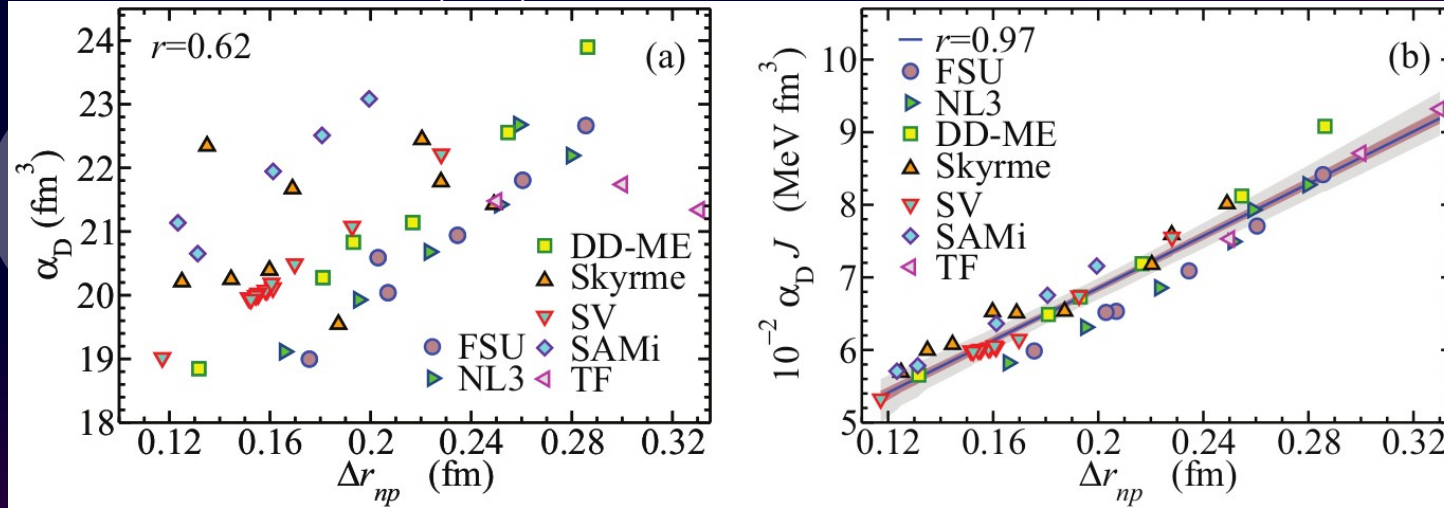


$$\alpha_D \text{ (fm}^3\text{)} = \mathbf{2.7} \quad \mathbf{16.2} \quad \mathbf{1.2}$$

# Connection to EoS and NS

## Dipole Polarizability and (L,J)

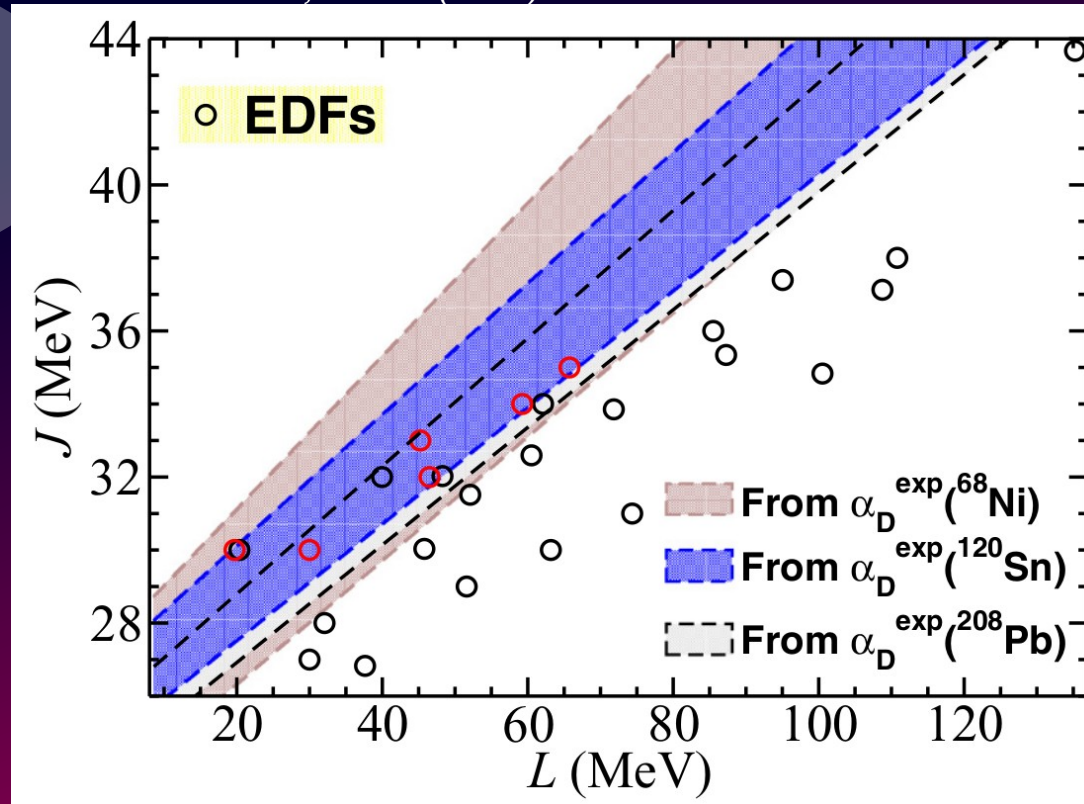
X. Roca-Maza et al, PRC88 (2013)



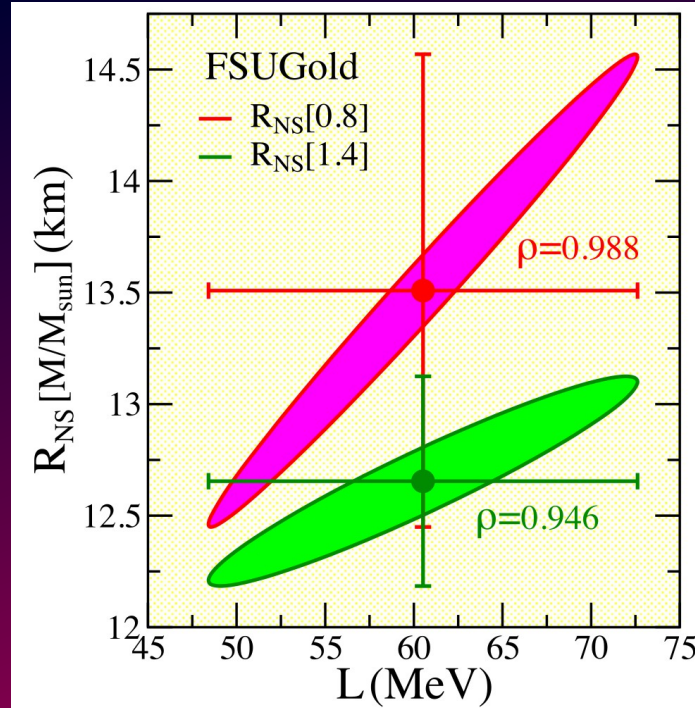
$^{208}\text{Pb}$

# Connection to EoS and NS Dipole Polarizability and (L,J)

X. Roca-Maza et al, PRC92 (2015)

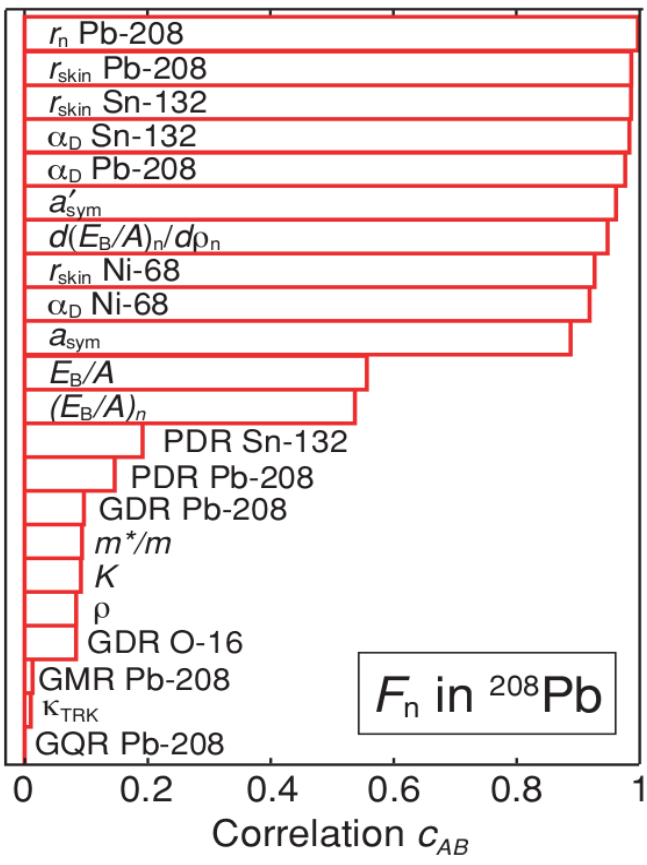


M. Thiel et al, *J. Phys. G: Nucl. Part. Phys.* 46 (2019)



**Information content of a new observable: The case of the nuclear neutron skin**

P.-G. Reinhard<sup>1</sup> and W. Nazarewicz<sup>2,3,4,5</sup>



# Connection to EoS and NS

Maybe not that clear ...

$^{208}\text{Pb}$  study case :

EWSR of PDR  $\rightarrow$  not strongly correlated with  $\Delta r_{np}$  !  
 Polarizability  $\rightarrow$  strong correlation with  $\Delta r_{np}$   
 (even better,  $\alpha_D^*J$ )

But there are indications that these correlations are model dependent (cf X. Roca-Maza studies)

Not the end of the story ...