



Simulation studies of a micro bulk XY-Micromegas detector for neutron-induced charged particle tracking

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Neutron detectors based on MICRO-Mesh Gaseous Structure (MICROMEGAS)



XY – micromegas detector

Drift HV



→ 60 x 60 strips (6x6 cm²) → Mesh hole ~ 60 μ m → Pitch: 100 μ m



Fig. 2. (Colour online) Photo of the first 6×6 cm² segmented mesh microbulk detector produced, mounted on the thick PCB.



Fig. 1. (Colour online) Schematic view of the segmented mesh microbulk detector. The holes of the micromesh are arranged in matrices with a fixed number of holes/column in the overlapping region of mesh and anode strips.



Development of a novel segmented mesh MicroMegas detector for neutron beam profiling, https://doi.org/10.1016/j.nima.2018.06.019, Nuclear Inst. and Methods in Physics Research, A 903 (2018) 46–55

XY – micromegas detector at CERN







3000 Beam profile 20 3 (all n energies) 2500 10 -(mm) 2000 1500 X -10 1000 -20 500 -30 30 -30 -20-10 0 10 20 X (mm)

Development of a novel segmented mesh MicroMegas detector for neutron beam profiling, https://doi.org/10.1016/j.nima.2018.06.019, Nuclear Inst. and Methods in Physics Research, A 903 (2018) 46–55

Experimental results in Orphée reactor (LLB)



 \rightarrow Study the performance in the detection of the neutrons and the reconstruction of neutron beam profiles

→ Neutron kinetic energy of 3 meV

→ Simulate the experiment
 → Reproduce the amplitude
 distribution histograms





Fig. 11. Reconstructed total amplitude distribution histogram, by adding the amplitudes of all the strip signals in each event, for the anode (up) and the mesh (down), from all the events (black) and only from the selected ones with the criteria applied (grey).

Table 2

Characteristics of the masks used. In most cases, PMMA plates were used in order to reduce the neutron fluence (by a factor of 16).

| Mask shape | Dimensions (mm) |
|------------------|-----------------|
| Circular hole | ø 5 |
| Circular hole | ø 2 |
| Square hole | 5 × 5 |
| Rectangular hole | 1×5 |

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Basic set-up: approximation – FLUKA simulations



Drift : 1 cm Size: 6 x 6 cm²

Frontiers in Physics 9, 788253 (2022). Annals of Nuclear Energy 82, 10-18 (2015). https://fluka.cern, https://flair.cern/



⁶LiF target: Ion Penetration Distance – SRIM vs SDTrimSP



Fig. 7. ${}^{6}Li(n,\alpha)^{3}H$ reaction product energy loss in a ${}^{6}LiF$ film as described by the Bragg distribution.

doi:10.1016/s0168-9002(02)02078-8

A. Mutzke et al.(2019), SDTrimSP Version 6.00 (IPP 2019-02). Garching: Max-Planck-Institut für Plasmaphysik. doi:10.17617/2.3026474.

alpha / triton yield - FLUKA simulations



- Differential yield → solid angle in steradian
- Angles with respect of the beam direction (polar angles in degrees) and results normalized as double differential, expressed in [particles GeV⁻¹sr⁻¹/primary]

→ For a given neutron interaction point the direction the two capture fragments are emitted is isotropically distributed.
 → The detector response, depending on its geometry, will be different according to the particle ejection direction.
 → What is measured, when the detector is exposed to a point source, is an average over all the possibilities that results into a distribution with its own FWHM that defines the resolution due to the gas.

alpha / triton fluence – FLUKA simulations



→ Track-length density in particles/cm² per primary

FLUKA vs Geant4





Energy deposition spectra

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dE/dx: FLUKA vs SRIM



TRITON and 4-HELIUM Stopping Power in Ar-CO₂ (10%), 1 Atm

- Triton: FLUKA and SRIM are in perfect agreement
- Alpha: Underestimation in FLUKA ("different" stopping ٠ power recipe)

ŰĽ

MGDRAW: Energy deposition hits – FLUKA simulations



Comparison with experimental data



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Experimental data



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New XY-mmegas design

10 x 10 cm²

Micromegas chamber



Ar-iC₄H₁₀ (5%), 1 Atm

drift, ring Al 1 mm thick, ^{10}B deposit on mylar 10 μm



- \rightarrow 100 x 100 strips (10x10 cm²)
- \rightarrow Mesh hole ~ 60 μm
- → Pitch: 100 µm

- Larger detector
- VMM3 electronics to replace AGET
 - n_TOF facility (CERN) (thermal-GeV)
 - GELINA (IRMM) (1meV-20MeV)
 - NFS (GANIL) (1-40MeV)

Experimental set-up and first tests



Penetration depth – SRIM calculations



Energy deposition spectra – FLUKA simulations (i)



Energy deposition spectra – FLUKA simulations (ii)



• signal representation etc...

Garfield++ simulation of micromegas



Garfield++: methods to be used



→ **DriftLineRKF** calculates the path of an electron or ion by numerical integration of the drift velocity vector. In the absence of a magnetic field, the drift lines will follow the electric field lines (using the previously computed tables of transport parameters to calculate drift lines and multiplication). The method is well adapted to fields that are smooth, such as analytic potentials.

→ AvalancheMicroscopic simulates electron trajectories using a "microscopic" Monte Carlo simulation based on the electron-atom/molecule scattering crosssections where the electron is followed from collision to collision. Provides an accurate simulation of event-by-event fluctuations of the electron signal (very time consuming).

 \rightarrow **AvalancheMC**, similarly as DriftLineRKF, uses the macroscopic drift velocity as function of the electric field for calculating electron or ion drift lines but adds a random diffusion component to each drift line step. It simulates all electrons in an avalanche individually.

• **DriftLineRKF:** there is no random/Monte Carlo element in the drift line simulation. For a given starting point, all electrons/ions will go at the same end point. On the other hand, **microscopic** tracking of electrons is typically used when fluctuations are important (for instance, when an accurate description of diffusion is needed, especially in non-uniform fields, or for studies of gain fluctuations in the avalanche gap to get the ion signal). For the electrons in the drift gap, diffusion will just spread the arrival points of the electrons at the mesh. In this case, we are looking at the signal induced on the mesh plane as a whole.

Both methods are used and compared with each-other

 \rightarrow For a single alpha track: $\Delta T_{RKF} / \Delta T_{MC} \sim 1 / 4000 s$

MC vs RKF



- Generation of 200 vertical 2 MeV α-particles and 2.7 MeV tritons with Garfield++ with both methods
- Collection of simulated data for the deposited charge in the anode electrode (pad) using the RKF method
- The information for the energy loss for both ionizing particles is obtained by using the SRIM code (version: SRIM-2013.00)
- Load a file with the mobility of the ions in argon
- Then, retrieving the "clusters" along the track, the drift is simulated for each of the primary electrons
- The drift lines of the electrons released in the drift gap will stop once they hit the mesh plane
- Use of the MC method for evaluation of the spread of the arrival points of the electrons in the mesh and estimation of the magnitude of the diffusion
- 200 vertical tracks of 2 MeV α-particles and 2.7 MeV tritons were tested using both methods
- Both methods produce the same number of primary electrons that are ~ 152 / alpha track and ~ 24 / triton track, while The secondary electrons produced by the miscroscopic MC method are ~ 75.3 K / alpha track and ~ 11.6 K / triton track
- Are in perfect agreement with the expected quantities of electron-ion pair
- The standard deviations of the distributions of the arrival positions in the XY end plane of the primary electrons as well as of all electrons are almost equal
 - → Using the MC method but tracking only the primary electrons without the tracking of the secondaries

Comparing RKF and MC: results



| Method | RKF | MC no diffusion | MC with diffusion |
|-----------------------------------|------|-----------------|-------------------|
| $\sigma_{\rm alpha} \ ({\rm mm})$ | 0.46 | 0.46 | 0.57 |
| $\sigma_{ m triton} \ ({ m mm})$ | 0.70 | 0.70 | 0.77 |

[0.01 cm / bin]

Simulated total induced charge



The integral of the induced current pulse related to the total induced charge without any convolution yet with a transfer function for a single vertical alpha / triton track

Pulse after convolution with a GET transfer function



0.5

0.4 P(t)

0.3

0.2

0.1

0

-0.1

0

 $h(t) = A \exp(-3t/\tau)(t/\tau)^3 \sin(t/\tau)$

1000

1500

t [ns]

2000

A = 10 $\tau = 232 \text{ ns}$

500



$FLUKA \rightarrow Garfield++$

- Generate a simulated spectrum of α-particle and triton tracks as exited from the ⁶LiF target, having trajectories distributed within the whole detector volume, using Garfield++
- Compare it qualitatively with the simulated total energy deposition histogram as calculated by FLUKA.

Insert E_{kin} and direction of alpha and triton (*LiF* \rightarrow *GAS*) from FLUKA to Garfield++

FLUKA energy deposition spectrum vs RKF deposited charge



Summary

- A solid simulation framework for the XY micromegas detector has been established using the current state-of-the-art MC computer codes FLUKA, Geant4 and Garfield++, SRIM
- Extract at first with a transport MC code (FLUKA, Geant4) the energies and direction cosines of the neutron reaction products and then use Garfield++ to transport them in the drift region and record all the primary electrons that arrive on the micromesh, using the full MC microscopic model enabling only the tracking of primary electrons
- Following this scheme of analysis, we plan to address and carry out studies using other type of converters and fission fragments

Outlook

neBEM / Garfield



From top to bottom: HV1 = -600, HV2 = 0, HV3 = +350 [V] Wire diameter = 0.0018, Hole diameter = 0.0045 [cm] Drift gap: 1 and Amplification gap: 0.0050 [cm]

Cez

100

120



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neBEM / Garfield



neBEM / Garfield++

rest-framework

root [0] REST_Detector_ViewReadout("readouts.root", "pixel")





(int) 0 root [1] .q

Readout: pixelReadout --- plane: 0



back-up



Cer



Neutron detection with Micromegas

Due to the so-called 3He shortage crisis, many detection techniques used nowadays for thermal neutrons are based on alternative converters. Thin films of ¹⁰B or ¹⁰B₄C are used to convert neutrons into ionizing particles which are subsequently detected in gas proportional counters, but only for small or medium sensitive areas so far.

Neutron detection \rightarrow neutron to charge converter

- Solid converter: thin layers deposited on the drift or mesh electrode (¹⁰B, ¹⁰B₄C, ⁶Li, ⁶LiF, U, actinides...)
 - ✓ Sample availability & handling
 - \checkmark Efficiency estimation
 - Limitation on sample thickness from fragment range

 flimited efficiency
 - × Not easy to record all fragments
- $\blacktriangleright \quad Detector gas (^{3}He, BF_{3}...)$
 - ✓ Record all fragments
 - ✓ No energy loss for fragments [▲] reaction kinematics
 - \checkmark No limitation on the size $^{\diamond}$ high efficiency
 - ★ Gas availability
 - ★ Handling (highly toxic or radioactive gasses)
- Neutron elastic scattering
 - ▶ gas (H, He)
 - ▹ solid (paraffin etc.)
 - Availability
 - High energies
 - * Efficiency estimation & reaction kinematics



Neutron detection with high efficiency (~50%):

- ³He crisis
- Increased demand for neutron detectors
 - → Science
 - ➔ Homeland security
 - → Industry

Micromegas for neutrons

- Micro-Pattern Gaseous Detector (gain, fast timing, high rate, granularity, radiation hardness, simplicity...)
- Low mass budget
- Transparent to neutrons
- Large area detectors cheap & robust



ELECTRONICS



Challenge: No global trigger signal =>

AGET electronics* + Reduced CoBo configuration

Self triggering mode / timing difference between strips.

- 64 analog channels /chip.
- Auto trigger: discriminator and threshold
- Multiplicity signal: analog OR of 6discriminators
- Address of the hitted channels
- SCA readout mode (all/hitted/selected channels)
- Max sampling rate: 100 MHz.
- 16 peaking time values: 50 ns-1us.
- 4 charge ranges/channel: 120fC/ 240fC/ 1pC/ 10 pC.



*GET, General electronics for TPC, ANR proposal / GET-QA-000-0005, AGET Data Sheet.

22 NEUTRON BEAM MONITOR + PROFILER



Accurate neutron cross section measurements require:

Neutron fluence/Beam interception factor

Number/fraction of neutrons hitting the area covered by the sample.

Shape of the beam profile

Beam optics misalignment => Beam fluence variations.

For non-monoenergetic neutron sources:



sample



=>Dependence of profile on the neutron energy

Requirements:

- Quasi-online neutron flux + beam profiler as well
- Minimal perturbation of the neutron beam / Minimal induced background
- · Stay permanently in the beam

PRINCIPLE OF NEUTRON BEAM PROFILER + MONITOR OPERATION





The neutron flux will be simultaneously extracted from the SUM SIGNAL.

cea

THE n_TOF FACILITY (CERN)





Neutron beam characteristics:

- 1. White neutron beam (0.025 eV-1 GeV).
- 2. Neutron energy defined by the Time-Of-Flight.
- 3. High instantaneous flux.
- 4. 7ns width proton pulse every > 1s
- 5. Excellent energy resolution



MOUNTING THE DETECTOR IN EAR-2





Test Beams with Resistive - Micromegas prototypes



VMM3a

- Setup of of 2x MMFE1s on 2x Resistive Micromegas chambers (Ar+7%CO₂ 400µm pitch, 5mm drift)
- Custom made firmware and software was

developed allowing to trigger with scintillator system

- Mode to control the CKBC externally
- High data rate ~20KHz/channel (VMM can reach 4MHz), arrived at the limit of Gbps UDP connection
- Noise levels of 300 e- ENC at gain 9mV/fC, 200ns





BROOKHAVEN

VMM3

VMM ASIC - G. lakovidis

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