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中山大学中法核工程与技术学院  
Institut franco-chinois de l'énergie nucléaire université Sun Yat-sen

# Bayesian model averaging for nuclear symmetry energy from effective proton-neutron chemical potential difference of neutron-rich nuclei

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Mengying QIU

Sun Yat-sen University, China

- Collaborator: Zhen Zhang, Bao-Jun Cai, Lie-Wen Chen, Cen-Xi Yuan

12 September 2024

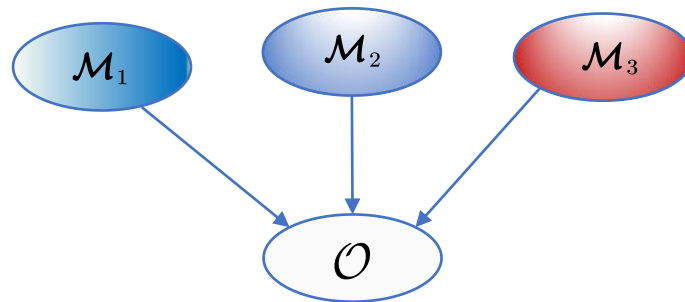
# Why Model Averaging?



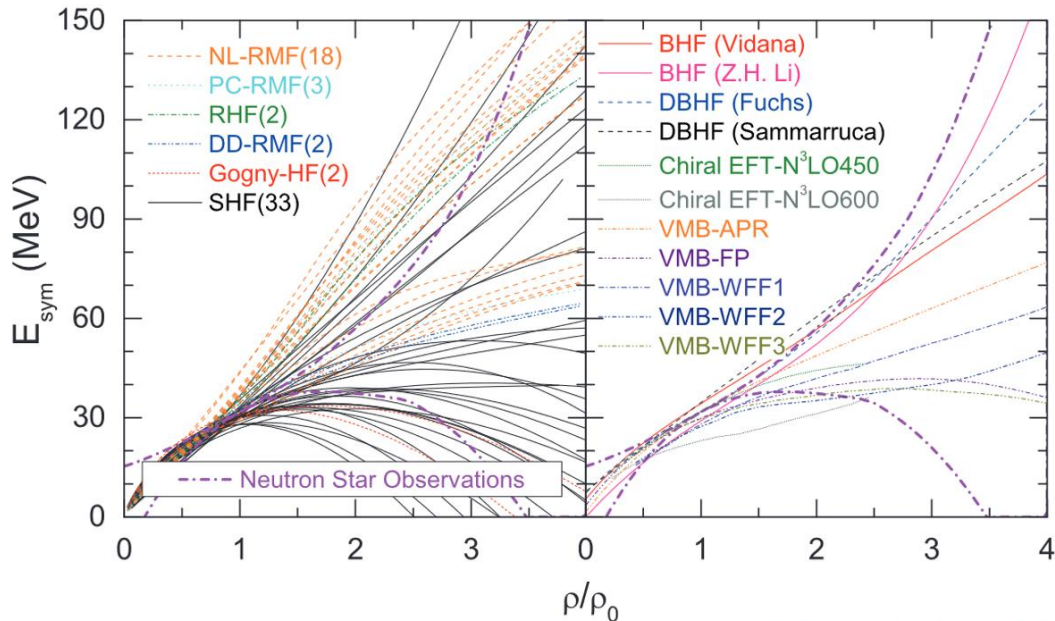
George E.P. Box

“Essentially,  
All models are wrong, but some are useful”

“Which model should we trust?”



# Why Model Averaging for symmetry energy?



L.W.Chen, Nucl.Phys.Rev.34,20 (2017).  
N.B.Zhang, B.A.Li, Eur.Phys.J.A 55,39(2019)

## □ Large uncertainty remains

- Intra-model uncertainty

Experimental data uncertainty  
Theoretical uncertainty  
Correlation between parameters

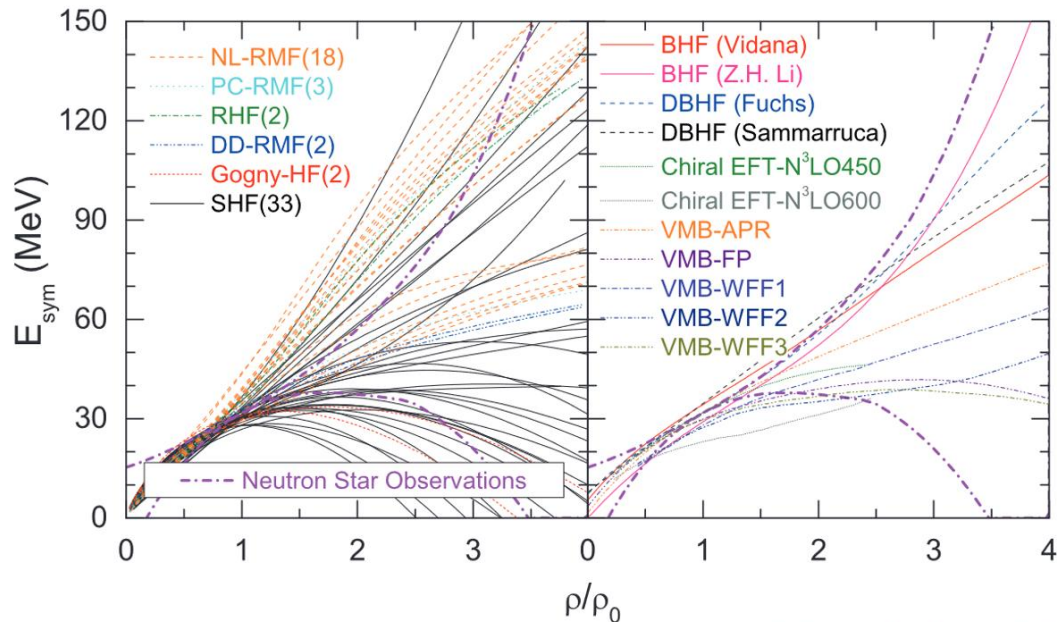
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Variations across different models

→ Possible bias

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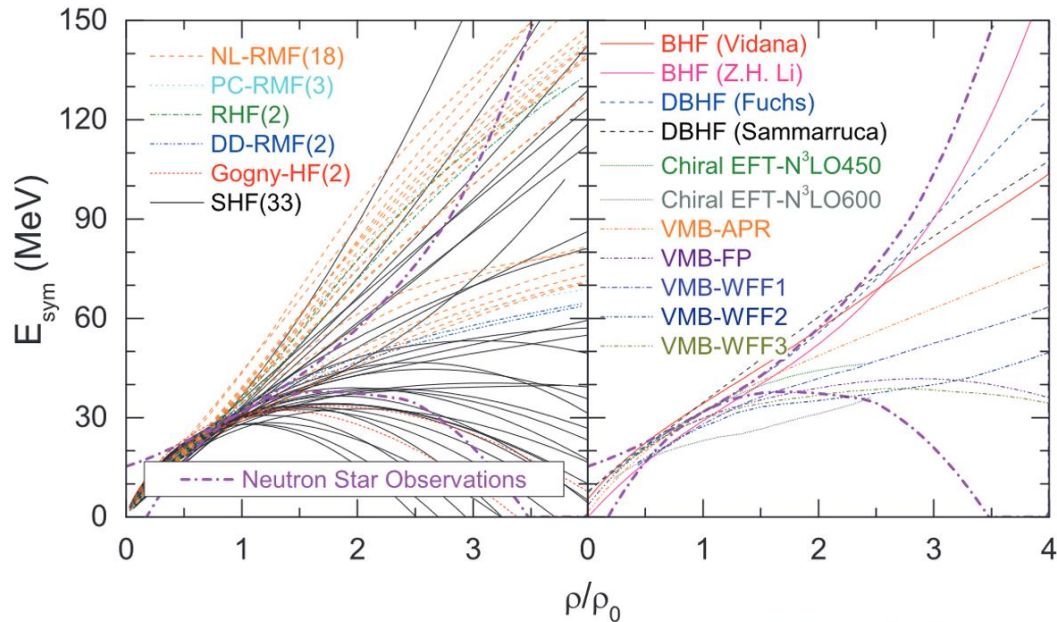
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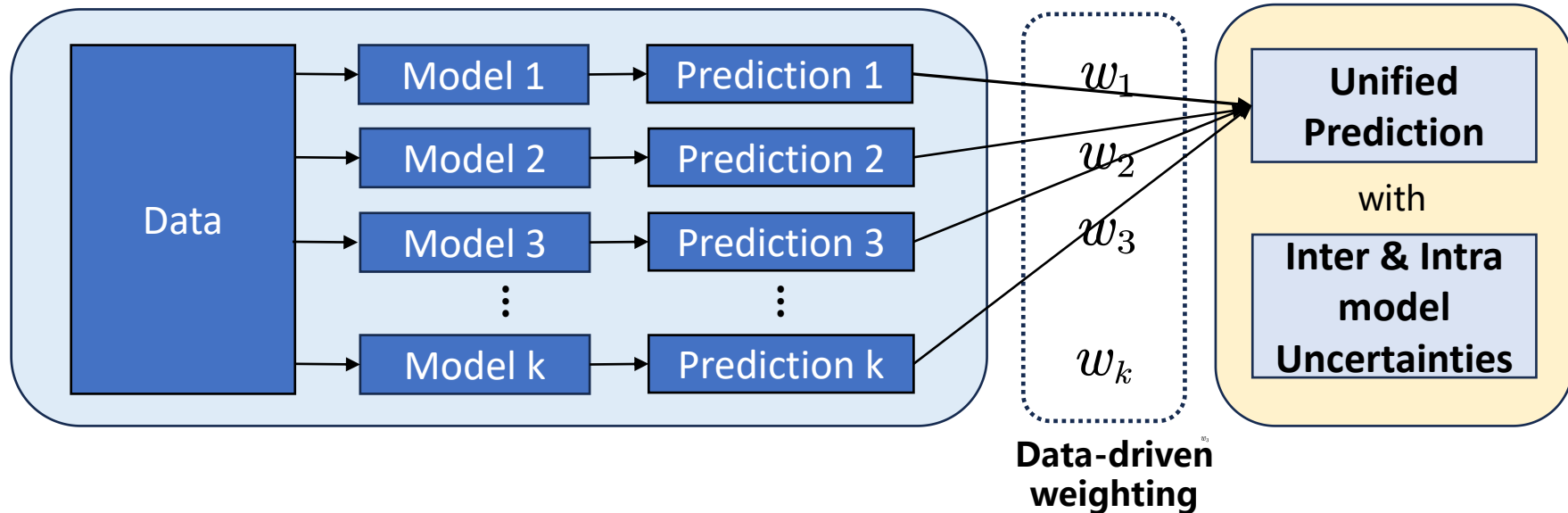
→ Possible bias

How to deal with it ?

# Model Averaging



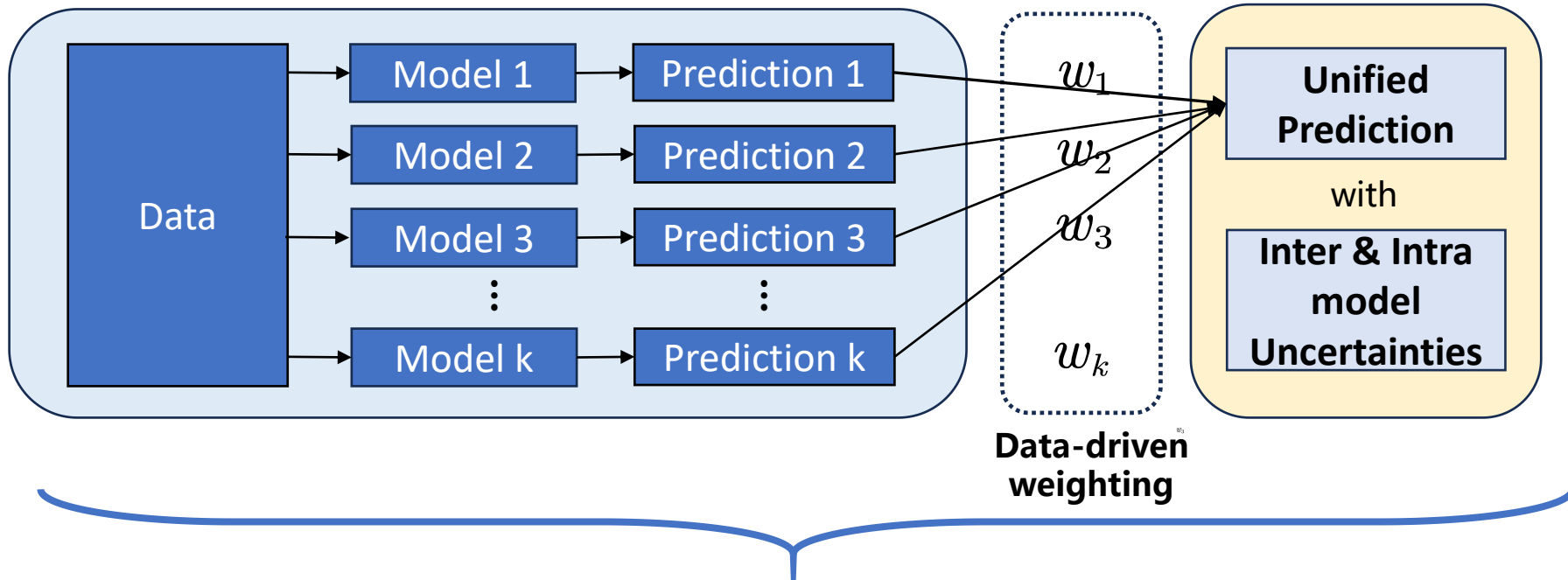
Possible option for combining model predictions



# Model Averaging



Possible option for combining model predictions



Consistent treatment within Bayesian framework

# Bayesian Analysis



Under model  $\mathcal{M}'$ 's assumption

**The likelihood function**  
of observing  $y$  given the  
model  $\mathcal{M}$  predictions at  $\Theta$

**The posterior probability**  
distribution of quantities of  
interest  $\Theta$  given experimental  
measurements  $y$

$$p(\Theta | y, \mathcal{M}) = \frac{\mathcal{L}(y | \Theta, \mathcal{M}) \pi(\Theta | \mathcal{M})}{p(y | \mathcal{M})}$$

**The prior probability**  
of quantities of interest  $\Theta$  before  
being confronted with the  
experimental measurements  $y$

**The marginal likelihood/Evidence**  
The probability of model  $\mathcal{M}$   
giving experimental measurements  $y$



# Bayesian Model Averaging (BMA)



- Each model's contribution is weighted by its **model posterior probability**

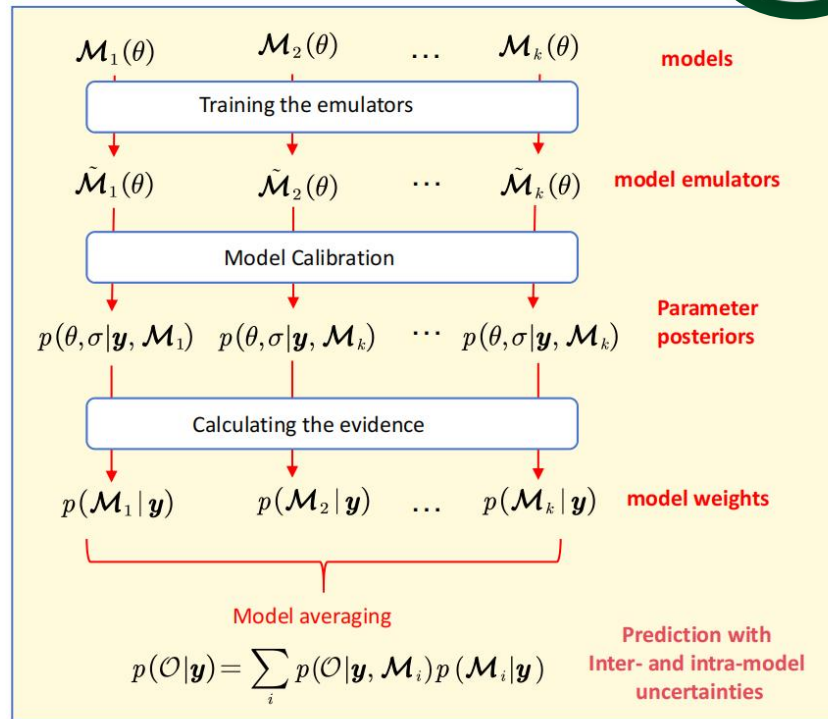
$$p(\mathcal{O}|\mathbf{y}) = \sum_i p(\mathcal{O}|\mathbf{y}, \mathcal{M}_i) p(\mathcal{M}_i|\mathbf{y})$$

- Model posterior probability: a weighting factor

$$p(\mathcal{M}_i|\mathbf{y}) = \frac{p(\mathbf{y}|\mathcal{M}_i)\pi(\mathcal{M}_i)}{\sum_\ell p(\mathbf{y}|\mathcal{M}_\ell)\pi(\mathcal{M}_\ell)}$$

- The model prior  $\pi(\mathcal{M}_i)$  is our preference on  $\mathcal{M}_i$  before seeing the data
- Bayesian evidence/marginal likelihood:** measures the probability that the model reproduces the experimental data

$$p(\mathbf{y}|\mathcal{M}_i) = \int p(\mathbf{y}|\theta_i, \sigma_i, \mathcal{M}_i)\pi(\theta_i, \sigma_i|\mathcal{M}_i)d\theta_i d\sigma_i$$



V. Cirigliano et al, J. Phys. G 49, 120502 (2022)



# Effective Proton-neutron chemical potential difference

## Effective chemical potential

$$\mu_n = \frac{\partial B(N, Z)}{\partial N} \approx \frac{B(N+2, Z) - B(N-2, Z)}{4}, \quad (1)$$

$$\mu_p = \frac{\partial B(N, Z)}{\partial Z} \approx \frac{B(N, Z+2) - B(N, Z-2)}{4}, \quad (2)$$

## Proton-neutron chemical potential differences

$$\Delta\mu_{pn}^* = \frac{1}{4} [B(N, Z+2) - B(N, Z-2) - B(N+2, Z) + B(N-2, Z)]$$

## Semi empirical mass formula

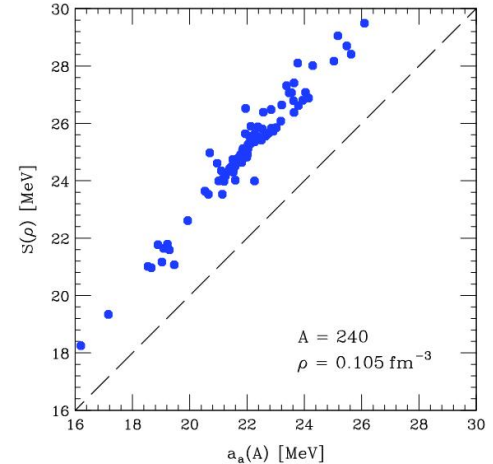
$$B(N, Z) = a_v A - a_s A^{2/3} - a_c \frac{Z^2}{A^{1/3}} - a_{\text{sym}} I^2 A + E_{\text{mic}},$$

## Expected sensitivity

pairing and shell effects

$$\Delta\mu_{pn}^* \simeq a_c \left[ \frac{1-Z}{(A-2)^{1/3}} - \frac{1+Z}{(A+2)^{1/3}} \right] + a_{\text{sym}} \frac{4A^2 I}{A^2 - 4} \simeq -2a_c \frac{Z}{A^{1/3}} + 4a_{\text{sym}} I$$

$$\Delta\mu_{pn}^* \propto a_{\text{sym}} \approx E_{\text{sym}}(\rho_r)$$



Pawel Danielewicz, Jenny Lee, Nuclear Physics A 922 (2014)  
 M. Centelles, Phys. Rev. Lett. 102, 122502 (2009)  
 L.-W. Chen, Phys. Rev. C 83, 044308 (2011)  
 N. Wang, L. Ou, and M. Liu, Phys. Rev. C 87, 034327(2013)

# Non relativistic & covariant EDFs



$$G_s, G_v, W_0 \quad m_{s,0}^* / m, m_{v,0}^* / m$$

$$\rho_0, E_0(\rho_0), K_0 \quad E_{\text{sym}}(\rho_0), L$$

$$m_{\text{Dirac}}^* / m, m_\sigma, c_\omega$$

Analytical transformation with  
pseudo-observables in nuclear matter

## Standard Skyrme Hartree Fock (SHF) Model

$$\begin{aligned} v(\mathbf{r}_1, \mathbf{r}_2) = & t_0(1 + x_0 P_\sigma) \delta(r) \\ & + \frac{1}{2} t_1(1 + x_1 P_\sigma) [\mathbf{k}'^2 \delta(\mathbf{r}) + \delta(\mathbf{r}) \mathbf{k}^2] \\ & + t_2(1 + x_2 P_\sigma) \mathbf{k}' \cdot \delta(\mathbf{r}) \mathbf{k} \\ & + \frac{1}{6} t_3(1 + x_3 P_\sigma) [\rho(R)]^\alpha \delta(\mathbf{r}) \\ & + i W_0 \sigma \cdot [\mathbf{k}' \times \delta(\mathbf{r}) \mathbf{k}] \end{aligned}$$

$$t_0 \sim t_3, x_0 \sim x_3, \alpha, W_0$$

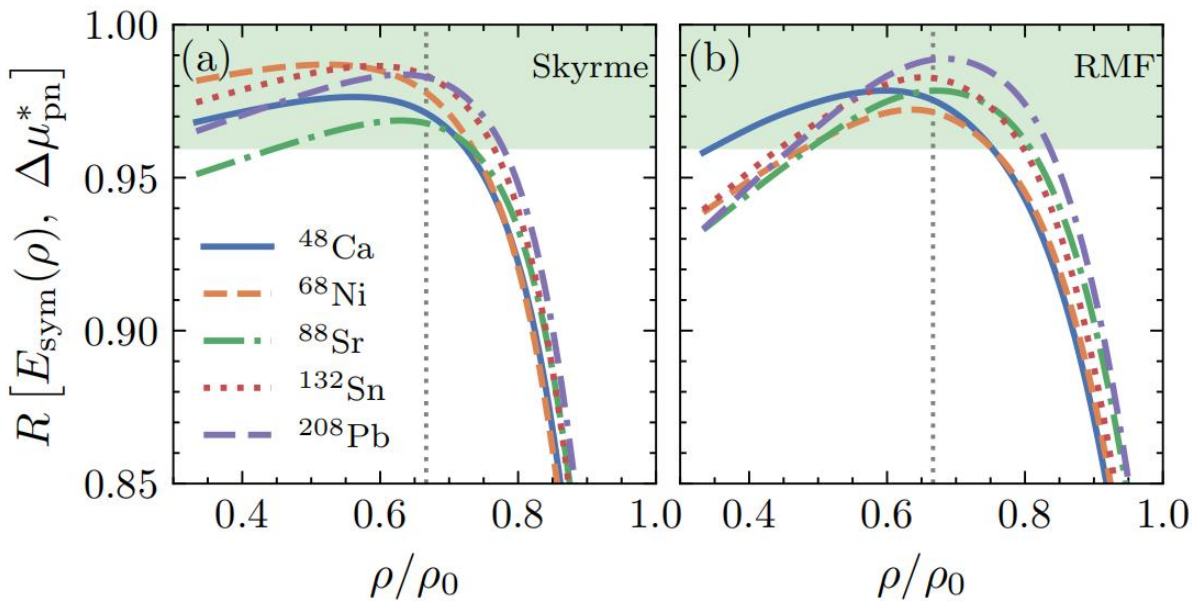
## Non-linear Relativistic Mean-Field (RMF) Model

$$\begin{aligned} \mathcal{L} = & \bar{\psi} (i \partial_\mu \gamma^\mu - m) \psi - e \bar{\psi} \gamma_\mu \frac{1 + \tau_3}{2} A^\mu \psi - \frac{1}{4} F^{\mu\nu} F_{\mu\nu} \\ & + g_\sigma \sigma \bar{\psi} \psi - g_\omega \omega_\mu \bar{\psi} \gamma^\mu \psi - g_\rho \bar{\rho}_\mu \bar{\psi} \gamma^\mu \tau \psi \\ & + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - \frac{1}{2} m_\sigma^2 \sigma^2 - \frac{1}{3} b_\sigma M (g_\sigma \sigma)^3 - \frac{1}{4} c_\sigma (g_\sigma \sigma)^4 \\ & - \frac{1}{4} \omega_{\mu\nu} \omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu + \frac{1}{4} c_\omega (g_\omega^2 \omega_\mu \omega^\mu)^2 \\ & - \frac{1}{4} \rho_{\mu\nu} \rho^{\mu\nu} + \frac{1}{2} m_\rho^2 \bar{\rho}_\mu \bar{\rho}^\mu + \frac{1}{2} \Lambda_V (g_\rho^2 \bar{\rho}_\mu \bar{\rho}^\mu) (g_\omega^2 \omega_\mu \omega^\mu), \end{aligned}$$

$$g_\sigma, g_\omega, g_\rho, b_\sigma, c_\sigma, c_\omega, \Lambda_V, m_\sigma$$

$\Delta\mu_{\text{pn}}^*$  for 5 doubly magic nuclei  
 $E_{\text{sym}}(\rho)$  at different densities

# Pearson correlation coefficient

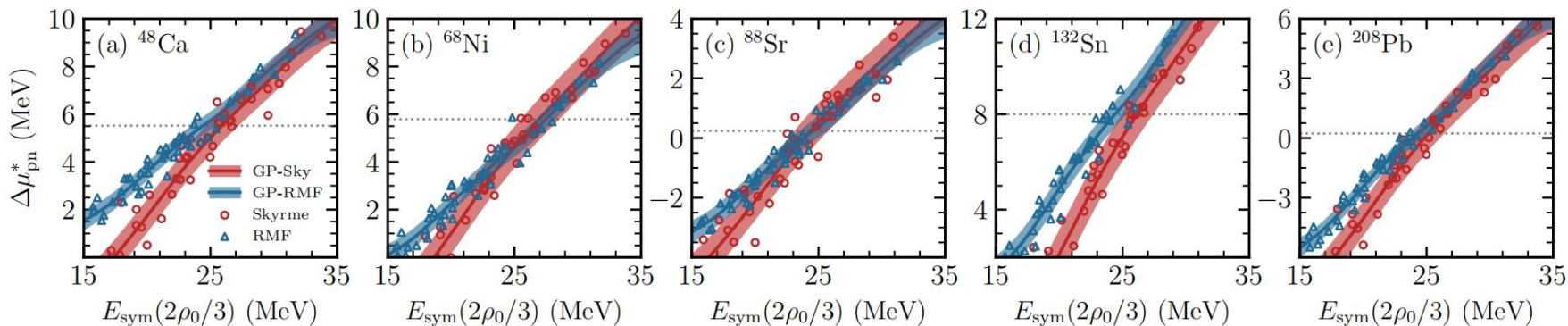


$$\blacklozenge R[A, B] = \frac{\text{Cov}(A, B)}{\sqrt{\text{Var}(A)\text{Var}(B)}}$$

- ◆ A strong linear correlation between the  $\Delta\mu_{pn}^*$  and the symmetry energy at subsaturation densities
- ◆ High sensitivity around  $2\rho_0/3$

M. Qiu, B. J. Cai, L.-W. Chen et al. Phys. Lett. B 849 (2024) 138435

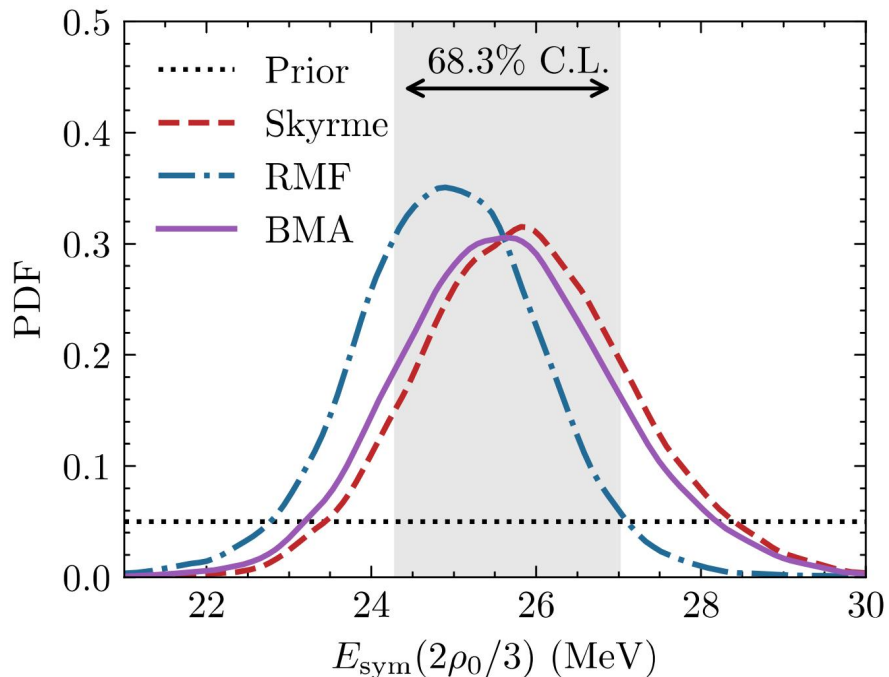
# Gaussian Process(GP)-Sky & GP-RMF



- Tune Gaussian processes using the results of 50 Skyrme EDFs and 50 covariant EDFs.
  - *Surmise* python package by BAND collaboration
- GP predictions with uncertainties.

M. Plumlee, O. Surer, S. M. Wild, and M. Y.-H. Chan, *surmise* 0.2.0,  
<https://surmise.readthedocs.io/en/latest/>

# Symmetry energy at $2\rho_0/3$



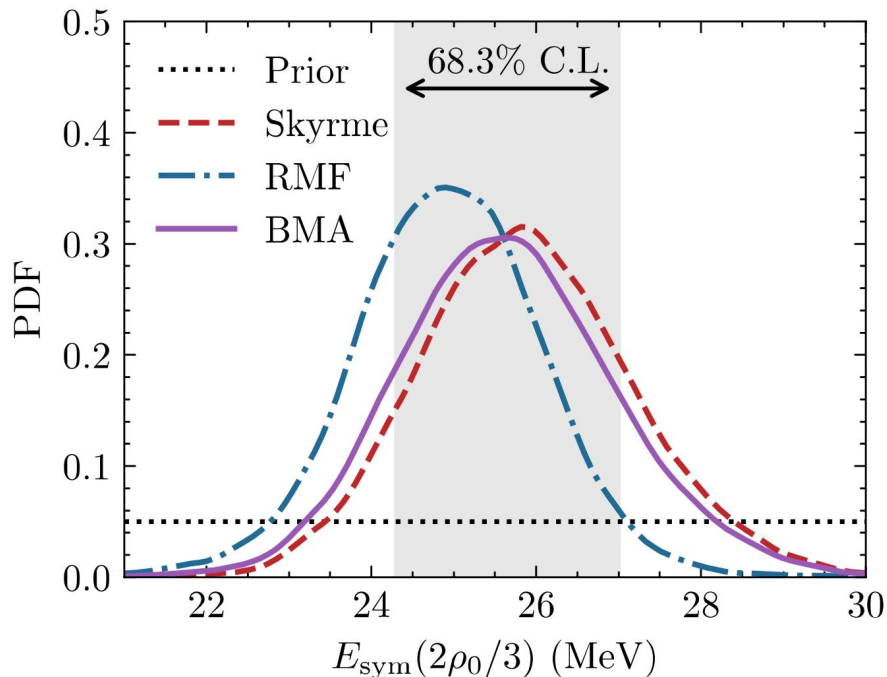
- Posterior by Sequential Monte Carlo algorithm from PyMCv4.0

O. Abril-Pla, et al., PeerJ Computer Science 9,e1516 (2023)

Skyrme EDFs  $E_{\text{sym}}(2/3\rho_0) = 25.8^{+1.3}_{-1.2}$  MeV  
 $\sigma = 0.4^{+0.4}_{-0.2}$  MeV

Nonlinear RMF  $E_{\text{sym}}(2/3\rho_0) = 24.9 \pm 1.1$  MeV  
 $\sigma = 0.8^{+0.6}_{-0.3}$  MeV

# Symmetry energy at $2\rho_0/3$



- Model posterior probability

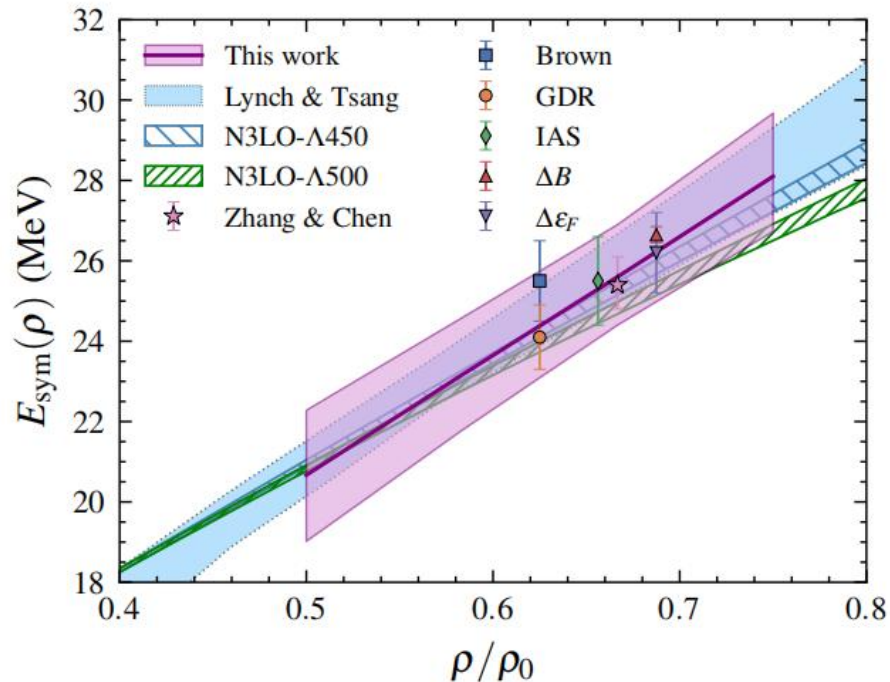
$$p(\mathcal{M}_i | \mathbf{y}) = \frac{p(\mathbf{y} | \mathcal{M}_i) \pi(\mathcal{M}_i)}{\sum_{\ell} p(\mathbf{y} | \mathcal{M}_{\ell}) \pi(\mathcal{M}_{\ell})}$$

- Equal prior preference
- Evidence ratio Sky/RMF  $\approx 3.3$



- $E_{\text{sym}}(2/3\rho_0)$  is inferred to be  $25.6^{+1.4}_{-1.3}$  MeV

# Symmetry energy at subsaturation densities



- **Brown:** Doubly magic nuclei  
B.A.Brown, *Phys. Rev. Lett.* 111, 232502 (2013)
- **Lynch & Tsang:** various terrestrial and astrophysical constraints  
W.G.Lynch and M.B.Tsang, *Phys. Lett. B.* 830, 137098 (2022)
- **Zhang & Chen:** Doubly magic nuclei+PREX+CREX  
Z. Zhang and L.-W. Chen, *Phys. Rev. C.* 108, 024317 (2023)
- **GDR:** Giant dipole resonance  
L. Trippa *et al*, *Phys. Rev. C* 77, 061304 (2008)
- **IAS:** Isobaric analog states  
Pawel Danielewicz, Jenny Lee, *Nuclear Physics A* 922 (2014)
- **$\Delta B$ :** Isotope binding energy difference  
Z. Zhang and L.-W. Chen, *Phys. Lett. B* 726, 234 (2013)
- **$\Delta \epsilon_F$ :** Fermi energy difference  
N. Wang, L. Ou, and M. Liu, *Phys. Rev. C* 87, 034327(2013)



# Summary



- Within both the non-relativistic Skyrme EDFs and the nonlinear RMF model, the effective proton-neutron chemical potential difference  $\Delta\mu_{pn}^*$  of neutron-rich nuclei is found to be strongly sensitive to the symmetry energy  $E_{sym}(\rho)$  around  $2\rho_0/3$ ,
  - We carried out a Bayesian model averaging analysis based on Gaussian process emulators to extract the symmetry energy around  $2\rho_0/3$   $25.6_{-1.3}^{+1.4}$  MeV
  - Since both the intra- and inter-model uncertainties are taken into account in our BMA analyses, the present results are statistically more reliable.
- Inclusion of more experimental observables and more theoretical models

**Thank you for your attention!**