Interpretation/Modeling: Prompt/afterglow emission / Population models IAP, 6-7 May, 2024

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GRB Lightcurves: prompt to afterglow

GRB Spectrum: Prompt

GRB Spectrum: Afterglow = non-thermal, probably 2 components

GRB diversity: XRR, XRFs, Low-L GRBs, etc.

- **Short GRBs tend to be harder**, with some exceptions
- Long GRBs show a lot of diversity, with soft or very soft events, usually also weaker: X-ray Rich Bursts, X-Ray Flashes, Low-Luminosity Bursts, etc. **Same physics/progenitors ? Importance of afterglow/host observations**

Initial event & central engine

Two main classes of progenitors:

- **Core-collapse of massive star (collapsar model) – « long » GRB**
- **Merger of binary neutron star system (or NSBH ?) – « short » GRB**

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Relativistic ejection

The GRB prompt emission has to be produced at large distance in a relativistic ejecta.

Relativistic ejection:

- Mechanism?
- Properties of the ejecta: Lorentz factor, geometry, magnetization, etc.

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A few examples of such « simple » diagnostics (not involving a heavy physical modeling)

Exemple 1: constraints on the Lorentz factor & emission radius

- § Fermi-LAT burst GRB 090926A: first observed cutoff at high-energy
- Analysis & interpretation: Yassine et al. [FD] 2017
	- time evolution of the cutoff

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Exemple 1: constraints on the Lorentz factor & emission radius in the emitted spectrum, responsible for both the MeV and the GeV photons. This is not necessarily the case as several *Fermi*-LAT bursts show evidence of an additional high-energy component. We show

• In cases without a detected HE cutoff (usual case): A lower limit on the Lorentz factor of the ejecta can be obtained. than the ejecta can be obtained.

What is needed: Sari & Lithwick 2001, Hascoet et al. [FD] 2012 $\frac{1}{\sqrt{2}}$ min can be obtained from the following formula:

- Variability timescale
- Spectral shape
- Redshift

$$
\Gamma_{\min} \simeq \frac{\left[C_1 2^{1+2\beta} \mathcal{I}(\beta)\right]^{1/2(1-\beta)}}{\left[\frac{1}{2}\left(1+\frac{R_{\text{GeV}}}{R_{\text{MeV}}}\right) \left(\frac{R_{\text{GeV}}}{R_{\text{MeV}}}\right)\right]^{1/2}} \frac{(1+z)^{-(1+\beta)/(1-\beta)}}{\left[\frac{E_{\max}E_{\text{c}}}{(m_{\text{e}}c^2)^2}\right]^{(\beta+1)/2(\beta-1)}},\,
$$
\n
$$
\times \left\{\sigma_{\text{T}}\left[\frac{D_{\text{L}}(z)}{c\Delta t_{\text{var}}}\right]^{2} E_{\text{c}} F(E_{\text{c}})\right\}^{1/2(1-\beta)} \left[\frac{E_{\max}E_{\text{c}}}{(m_{\text{e}}c^2)^2}\right]^{(\beta+1)/2(\beta-1)},\tag{59}
$$

its on the ejecta composition $\mathbf t$ Exemple 2: constraints on the ejecta composition

- photospheric dissipation as in Giannic dispite the change of the change of the change of the change of the change of the change of the change of the change A diagnostic can be applied in cases where the soft gamma-ray spectrum shows s: a non-thermal component and a quasi-thermal the acceleration is component and a quasi-thermal same radius (because *^R*ph [∝] ^Γ−3), additional differences in ar- $\sum_{i=1}^{n}$ evidence for two components: a non-thermal component and a quasi-thermal one.
- Required measurements:
	- r ature of the quasi-thermal component atare or the quasi-thermal component - flux and temperature of the quasi-thermal component
	- flux ratio non-thermal/thermal
- gly m ιουι
ΓΙ **Sari & Lithwick 2007** = strongly model-dependent. § Fireball model: leads to a measurement of the Lorentz factor and initial radius
- General assumptions: Hascoët, Daigne & Mochkovitch 2012
	- When the ejecta is launched: energy = fraction $\epsilon_{\rm th}$ thermal a is iddition on the individual power of a passive magnetic magnetic magnetic magnetic

 ${\bf fraction \ 1} - \epsilon_{\rm th}$ magnetic

- Large distance: magnetization σ Passive field: $\sigma_{\text{passive}} = (1 - \epsilon_{\text{th}})/\epsilon_{\text{th}}$, efficient magnetic acc. $\sigma < \sigma_{\text{passive}}$ Γ $\overline{5}$
- $\,$ Non-thermal emission has an efficiency $\,$ $f_{\rm NTh}$
- $E_{\rm eff} = 0.001$ acceleration leads to σ acceleration leads to σ **• Analysis allows to derive a constraint on some of these parameters.**

Ejecta composition from the thermal/non-thermal ratio

R. Hascoët et al.: Prompt thermal emission in gamma-ray bursts

GRB 090902B: a rare case where the quasi-thermal component could be dominant.

GRB100724B: a bright GRB where the presence of a quasi-thermal component can be constrained.

Here: constraints obtained for - a given non-thermal efficiency (x-axis), - and a given initial thermal fraction and large-scale magnetization (lines).

 $F = fireball$

M,is = internal shocks (low sigma) M,rec = reconnection (high sigma)

09092B: $\epsilon_{\rm th} > 0.3 - 0.5$

100724B: is or rec with low $\epsilon_{\rm th} < 0.1$

Fig. 5. Constraints on the thermal and non-thermal emission in GRB 090902B and GRB 100724B. *Top:* for a given thermal fraction $\epsilon_{th} = 10^{-2}$, $10^{-1.5}$, 10^{-1} , $10^{-0.5}$, and 1, the radius $R_0 = \ell/\theta$ at the base of the flow is plotted as a function of the non-thermal efficiency f_{Nth} . The corresponding thermal efficiency f_{th} is also shown (top *x*-axis). *Bottom:* for a given magnetisation $\sigma = 10^{-1}$, $10^{-0.5}$, 1, $10^{0.5}$, 10¹, $10^{1.5}$, and 10^{2} at the end of acceleration phase, the Lorentz factor of the flow is plotted as a function of f_{Nth} (the unmagnetised case $\sigma = 0$ cannot be distinguished from the case $\sigma = 0.1$). Sets of parameters representative of the different classes of scenarios discussed in the paper are indicated: F ($\epsilon_{\text{th}} = 1$, $\sigma = 0$) (standard fireball), M,is₁ (log $\epsilon_{\text{th}} = -0.5$, $\sigma = 0$) and M,is₂ (log $\epsilon_{\text{th}} = -1.5$, $\sigma = 0$) (efficient magnetic acceleration: magnetisation is low above the photosphere and the dominant non-thermal mechanism is internal shocks), M,rec₁ (log $\epsilon_{th} = -0.5$, $\sigma = 10$), and M,rec₂ (log $\epsilon_{th} = -1.5$, $\sigma = 10$) (magnetised flow at large distance, the dominant non-thermal mechanism is magnetic reconnection). The initial radius is fixed to $R_0 = 300$ km, a typical value for long GRBs. The observational data (thermal flux, temperature, ratio of the thermal over the total flux) used for the calculation (see text) are taken from Abdo et al. (2009); Pe'er et al. (2012) for GRB 090902B (*left column*), and from Guiriec et al. (2011) for GRB 100724B (*right column*).

Exemple 3: Lorentz factor from the peak of the optical afterglow we are not aware of any selection effects that could lead to factor of the blast wave at *R*dec; ρdec is the external density at ambient density at the deceleration radius, ρdec, is normalized to ρ0*/m*^p = <u>trom the neak of the ontical afterglo</u> **LORENTZ FROM BE SUBSTANTIAL CAN BE SUBSTANTIALLY**

E Model dependent: standard afterglow model (forward external shock) with some assumptions (efficiency of the prompt emission, ambient density, ...) <u>non-detection</u> of Detection of Late peaks *T there is no intrinsic correlation between Tp and L*^γ *(or E*^γ *). R*dec, and *s* describes the slope of the external density profile, **standard and grow moder from ward external points of the series of the series of the series of the series of the** at the deceleration time, factor of the blast wave at θ the blast wave at θ the external density at θ the external density at θ **Reducity Profile Shows Reducity Profile, and Shows Reducity Profile, and a** *s* = d ln ρ*/*d ln *R*. A uniform medium is described by *s* = 0 and with some assumptions (efficiency of the prompt emission, ambient density, .. (A color version of this figure is available in this figure is available in the online \mathcal{L} (2) The estimate gives the Lorentz factor of the *blast wave*, Γbw,

peak = deceleration time
$$
T_p \simeq T_{\text{dec}} \simeq \frac{R_{\text{dec}}}{2\Gamma_{\text{bw}}^2 c}
$$
. (OK if $T_p \gg T_{\gamma}$: NR RS)

Early afterglow with a detection of the peak: blast wave Lorentz factor Lorentz factor,

> 1 8

$$
\Gamma_{\rm bw} = \left[\frac{3-s}{32\pi c^5} \frac{1-\eta}{\eta} \frac{E_{\gamma}}{\rho_{\rm dec} T_{\rm p}^3}\right]^{\frac{1}{8}}
$$

See e.g. Molinari, Vergani et al. 2007, Ghirlanda et al. 2012 $\frac{1}{2}$ $\frac{1}{2}$ shock propagating back into the relativistic ejecta. As discussed that is converted into prompt radiation.

E Upper limit on the peak time: lower limit on the Lorentz factor at the decement of the decement of the decement of the decement of the set of the energy has a set of the energy has σ estimate depends on the poorly known and \downarrow 3 **Example 2** Upper limit on the peak time: \blacksquare estimate depends on the poorly known ambient density as

■ With a rich data set starting early, a full AG **Tenacle 12**, [erg s⁻¹] modelling is possible: constraints on Lorentz deceleration radius. The product Γ_{bw} ρ_{dec} is estimated using Equation (5) factor, density, prompt efficiency = a good science case for SVOM

T
Figure 2. Estimated Lorentz factor of the GRB blast wave, Γ_{bw}, at the state of smaller than *T*_γ . In the smaller than *T*_γ (which $\eta = 0.5$) and shown vs. the burst funniosity L_{γ} (Equation (2)). The ambient density at the deceleration radius, ρ_{dec} , is normalized to $\rho_0/m_{\text{p}} =$ **b**
1 cm⁻³. Bursts with $T_p < T_\gamma$ are highlighted in red; for these bursts the ejecta Lorentz factor Γ can be substantially higher than Γ_{bw} (see text). the bursts with bursts with low \mathbb{R} Figure 3 illustrates how the GRB sample is transformed $P_{p} =$ between ^Γbw and *^T*p, ^Γbw [∝] *^E*¹*/*⁸ (with $\eta = 0.5$) and shown vs. the burst luminosity L_V (Equation (2)). The

 \mathbf{H}^{H} (A color version of this figure is available in the online journal.)

A word of caution:

- These diagnostics are never 100% model-independent
- Lower limit on the Lorentz factor from gamma gamma: robust
- § Measurement of the Lorentz factor from the HE cutoff: less robust (e.g. the cutoff could be some curvature due to the natural shape of the HE component)
- Constrain on the initial thermal content of the ejecta: partially model dependent (valid only in scenarios with a non-dissipative photosphere)
- Measurement of the Lorentz factor from the peak of the early visible afterglow (standard afterglow model $+$ some assumptions on the prompt efficiency, the external density)

Prompt emission

Observed short timescale/non-evolving variability in GRB lightcurves imply an internal dissipation in the ejecta (Sari & Piran 1997a,b).

Internal dissipation / radiation processes:

- (Dissipative) Photosphere? (thermal + comptonization)
- Internal shocks? (synchrotron + IC)
- Reconnection? (synchrotron + IC)

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Photosphere

At the photospheric radius, the ejecta becomes transparent to its own radiation.

Photospheric emission:

- non-dissipative photosphere: thermal (Paczynski, Peer, Beloborodov, ...)
- dissipative photosphere: non-thermal (Rees & Meszaros, Beloborodov, ...) dissipation? shocks (radiation mediated shocks, see e.g. Samuelsson), reconnection (see e.g. Giannios), other?

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Dissipation in the optically thin regime

Reference model: internal shocks (electron acceleration in mildly relativistic collisionless shocks?) [low magnetization at large distance] **Alternative: reconnection** (electron acceleration ?) [low magnetization at large distance]

Radiation: synchrotron + Inverse Compton Scatterings in both cases

Dissipation in the optically thin regime

Reference model: internal shocks (electron acceleration in mildly relativistic collisionless shocks?) [low magnetization at large distance] **Alternative: reconnection** (electron acceleration ?) [interfactor magnetization at large $\frac{1}{2}$ distance] **Reference models internal shocks** (electron acceleration in milal evolution.
 Collisionless shocks?) [low magnetization at large distance]
 Alternative: reconnection (electron acceleration at large temporal evolution. Internal shocks: models available @IAP, difficult to directly apply to fit data
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Just an example: modelling the synchrotron spectrum in the optically thin regime

 $=$ the dominant process expected above the photosphere (standard prediction in fast cooling regime: alpha $= -3/2$)

Band function

§ Fermi/GBM spectral catalog: (10 years: Poolakkil et al. 2021)

- 2297 GRBs
- Time-integrated spectra / Spectrum at peak flux

Robustness of the spectral analysis: see Fred Piron's presentation...

Effects of SSC and of a B decay on syn+SSC radiation

A&A proofs: manuscript no. output

- $f(t)$ (fast cooling electron) and $f(t)$ otherwise (slowers) and $f(t)$ of the synchrotron spectrum: In presence of a decaying magnetic field with *t* **E** Low-energy photon index
- appears a SSC in Klein-Nishina regime: T therefore, electrons radiate exists r $\alpha = -3/2$ to -1 - SSC in Klein-Nishina regime:
- c,e↵ ' *^t* - B decay on an intermediate scale between the electron radiative dyn/*t* 0 B mescale and the dynamical decay, i.e. *t* $time$ cole: $\alpha = 2/2$ to $2/2$ timescale and the dynamical timescale: $\alpha = -3/2$ to $-2/3$

Full modelling: coupling the radiative model to the dynamics of the ejecta

Much more complex: not shown here. Advantage: spectral and temporal properties.

- IAP: tools available for internal shocks, can not be used to directly fit data (too heavy) $=$ intermediate step $=$ ISSM, see Fred Piron's presentation
- IAP: towards a hybrid model, with internal shocks starting below the photosphere (radiation mediated shocks) and still propagating above it (standard internal shocks): long term project with Filip Samuelsson.
- § Reconnection models: Bing Zhang et al.
- § Shock breakouts? (may be useful to study low-L/soft GRBs)
- § Other models?

Afferglow

The afterglow is associated to the deceleration of the relativistic ejecta by the external medium. (Rees & Meszaros, Piran & Sari, …)

- Ultra-relativistic forward shock in the external medium (electron acceleration in UR collisionless shocks?)
- Low magnetization: reverse shock in the ejecta (NR / UR)
- Synchrotron + Inverse Compton scatterings

Afterglow

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external medium. (Rees & Meszaros, Piran & Sari accelerated electrons

in the forward external shock.
 « Standard » model: $\frac{SVD}{N}$ and the stat standard a model: synchrotron radiation from shock-accelering relations and
standard a model: synchrotron radiation are available (« closure relations »).
Simple analytical implementation are available for data fitting)
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Still many open issues, especially related to the models for some scenarions of the property of the discussed, some capital many open issues, plated to the model models for some scenarion and principle of the some complem

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Data fitting with the forward external shock model

An exemple of an afterglow fit: GW170817

Posterior distribution on model parameters not shown here: should always be checked (for instance when using afterglowpy). One can find a good fit with strange parameters: says probably something (note the case here!).

Code still under development:

- synchrotron self-absorption
- reverse shock
- etc.

(structured jet, FS: syn+SSC)

Pellouin & Daigne, submitted

Spectrum at the peak:

Fig. 9. Posterior distribution of the afterglow spectrum around its peak ($t_{\rm obs} = 110$ days) for the "SSC (with KN)" fit of the afterglow of GW 170817. Data points show the multi-wavelength observations at $t_{\text{obs}} \pm 4$ days. The upper limit from H.E.S.S. is also indicated (Abdalla et al. 2020). The low-energy component (solid line) is produced by synchrotron radiation, while the high-energy emission (dashed line) is powered by SSC diffusions. Thick lines represent the median value at each observing frequency, dark contours the 68% confidence interval and light contours the 97.5% confidence interval. Some instrument observing spectral ranges are shown in colors.

One word on population models

- Expertise available: Jesse Palmerio's PhD work (Palmerio & Daigne 2021)
- **Population is described in terms of rate(z), luminosity function, etc.**
- **Main difficulties: correcting for selection effects.**
	- **= building a large & complete sample (e.g. flux-limited) to apply the model**
- Including the afterglow in the model is complex.
- Present status: Long GRBs: OK ; Short GRBs : more uncertain ; other kind of events ?

The example of the population of classical long GRBs

Long GRB Population Model

Blue: model with evolution (here: comoving rate)

 \mathcal{G} and \mathcal{G} and \mathcal{G} and \mathcal{G} and \mathcal{G} are defined by the distribution \mathcal{G} Here we use flux-limited samples to avoid the detailed modeling of the detection efficiency.

[yr¹ Gpc3] log *L*⇤ *p k*evol log *Ep*⁰ *^E*^p *zm a b* an cyolution is hecueu (comoving rate anu/or iuminosity function) $\lambda = a$ constraint on progenitor models 0.1 1.34 + 0. Results: an evolution is needed (comoving rate and/or luminosity function) $=$ a constraint on progenitor models

Palmerio & Daigne 2021

& Daigne 2021

Long GRB Population Model

Distinguising between rate and luminosity evolution?

Luminosity distribution in four redshift bins

Would need a sample allowing to build the luminosity distribution in a different redshift bins...

Long GRB Population Model

rate:

(violet)

Prediction for high-z GRBs? Interest of low-energy threshold

Pop. model: all-sky rate above z=6 as a function of the peak flux limit

Palmerio & Daigne 2021 & Daigne 2021

Questions, discussion?