Effects of angular momentum projection in RPA and GCM calculations

SSNET'24 - Shapes and Symmetries in Nuclei: from Experiment to Theory

IJCLab, Orsay, November 6th, 2024

Andrea Porro Technische Universität Darmstadt







PHYSICAL REVIEW C 109, 044315 (2024)

Symmetry-restored Skyrme-random-phase-approximation calculations of the monopole strength in deformed nuclei

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 ³ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung GmbH, 64291 Darmstadt, Germany
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 ⁵INFN, Sezione di Milano, Via Celoria 16, 20133 Milano, Italy
 ⁶Instituut voor Kern- en Stralingsfysica, Department of Physics and Astronomy, KU Leuven, 3001 Leuven, Belgium
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(Received 16 December 2023; accepted 12 March 2024; published 9 April 2024)



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arXiV > nucl-th > arXiv:2407.01325

Nuclear Theory

[Submitted on 1 Jul 2024]

Ab initio description of monopole resonances in light- and medium-mass nuclei: IV. Angular momentum projection and rotation-vibration coupling

Andrea Porro, Thomas Duguet, Jean-Paul Ebran, Mikael Frosini, Robert Roth, Vittorio Somà

- I. [EPJA (2024) 60, 133]
 II. [EPJA (2024) 60, 134]
 III. [EPJA (2024) 60, 155]
- IV. [arXiv:2407.01325]



Società Italian di Fisica



Introduction

- Giant Resonances
- GCM and RPA

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- Theoretical introduction
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Schrödinger equation $H |\Psi_n\rangle = E_n |\Psi_n\rangle$

Schrödinger equation

$$H \left| \Psi_n \right\rangle = E_n \left| \Psi_n \right\rangle$$

Open-shell systems

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Symmetry-breaking reference states



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1 Constrained HFB solutions $|\Phi(q)
angle$



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Open-shell systems





1 Constrained HFB solutions

 $|\Phi(q)\rangle$ Generator coordinates (q can be any coordinate)



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 $|\Phi(q)\rangle$

1 Constrained HFB solutions

2 PGCM Ansatz

$$|\Psi_n\rangle = \int \mathrm{d}q \, f_n(q) \, |\Phi(q)\rangle$$

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Initially developed for large-amplitude collective motion



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EHFB (q)

Emin



Linear coefficients

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Linear coefficients



Variational method

$$\delta \frac{\langle \Psi_n | H | \Psi_n \rangle}{\langle \Psi_n | \Psi_n \rangle} = 0$$



Schrödinger equation

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 $\mathcal{H}(p,q) \equiv \langle \Phi(p) | H | \Phi(q) \rangle$ $\mathcal{N}(p,q) \equiv \langle \Phi(p) | \Phi(q) \rangle$

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Thouless theorem

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Introduce the Quasi-Boson approximation (QBA)

>>>> Expand to the quadratic level in $oldsymbol{Z}(q,q_{min})$



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Harmonic approximation



EHFB (q)

No coordinates dependency !

 $|\Phi(q_{min})|$

Emin

All coordinates are explored (differently from GCM)

 $|\Phi(q)\rangle$

 $e^{\mathbf{Z}(q,q_{min})}$

Marmonic approximo

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Harmonic approximation

Eventually rewrites as (Q) RPA equations

$$\begin{pmatrix} A & B \\ -B^* & -A^* \end{pmatrix} \begin{pmatrix} X^n \\ Y^n \end{pmatrix} = E_n \begin{pmatrix} X^n \\ Y^n \end{pmatrix}$$

[Jancovici, Schiff, 1964]



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Nuclei that are **stiff** against deformations (anharmonic effects negligible) $\Phi(q)$

 $_{\rho}\boldsymbol{Z}(q,q_{min})$

Marmonic

Emin
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Intrinsic density is the fundamental variable in EDF

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Much is learnt from symmetry breaking and restoration

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- Common in SR and MR EDF
- Not in linear response on SR EDF (i.e. QRPA)



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SYMMETRY-CONSERVING RANDOM PHASE APPROXIMATION^{\dagger}

C. FEDERSCHMIDT and P. RING*

Physik-Department der Technischen Universität München, D-8046 Garching, West Germany

Received 18 July 1984

Abstract: The projected random phase approximation (PRPA) is derived from a generator coordinate ansatz. It allows the calculation of excited states in the region of phase transitions, where conventional RPA breaks down. The theory is applied for an approximate solution of the R(8) model which shows a pairing collapse at large angular momenta.



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Angular momentum projected mean-field

Deformed

mean-field

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Angular momentum

Projection before solving QRPA VAP QRPA

Deformed

Computationally expensive, no realistic application [Gambacurta, Lacroix, PRC (2012) 86, 064320]

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What about PAV QRPA?

Intrinsic density is the fundamental variable in EDF

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Projection before solving QRPA VAP QRPA

Computationally expensive, no realistic application [Gambacurta, Lacroix, PRC (2012) 86, 064320]

What about PAV QRPA? Can we treat projection a posteriori?



- Needle approximation for AMP
- RPA reinstates the missing symmetries to some extent



- Needle approximation for AMP
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The Random Phase Approximation: Its Role in Restoring Symmetries Lacking in the Hartree–Fock Approximation

A. M. LANE

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Present work

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- Focus on K=0 (monopole and quadrupole)

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$$P^{J}_{MK} = \frac{2J+1}{2} \int_{-1}^{+1} d\left(\cos\beta\right) \ d^{J}_{MK}(\beta) e^{-i\beta \hat{J}_{y}}$$

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Hartree-Fock wave-functions often lack symmetries possessed by the Hamiltonian. It is often said that the Random Phase Approximation (RPA) restores the missing symmetries. Since the RPA does not readily lead to explicit wave-functions, it is not a trivial matter to verify this assertion. We analyse the situation, and show that, while RPA restores symmetry in some respects, it does not do so completely. Besides the normal RPA, we discuss the generalisation of RPA that describes modes in isobars of the given nucleus. This is needed to enable us to discuss the case of isospin symmetry, which is analysed in detail.

Rotation operator



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Wigner small-d matrices

Rotation operator



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Remark

For J=0 projection is a pure rotation

 $|\text{ROT}\rangle \equiv N_{\text{ROT}}P_{00}^{0}|\text{HF}\rangle$

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QRPA code based on HFBTHO v1.66 SkM* parametrisation Systematic study for ²⁴Mg





[PRC (2024) 109, 044315]



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[PRC (2024) 109, 044315]

AMP RPA results



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AMP RPA results

Intrinsic frame (deformed) Laboratory frame (projected)



AMP RPA results

Intrinsic frame (deformed) Laboratory frame (projected)



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Laboratory frame (projected)


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Observation in other systems







GCM : symmetry-breaking solutions $|\text{GS}\rangle_{\text{def}}$ $|\omega\rangle_{\text{def}}$ PGCM : symmetry-conserving solutions $|\text{GS}\rangle_{\text{sym}}$ $|\omega\rangle_{\text{sym}}$



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Variational treatment of rotations in PGCM !



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Variational treatment of rotations in PGCM !

Projection effects

- Not too dissimilar
- Increased fragmentation (e.g. ²⁴Mg)
- More quantitative agreement



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Can we treat projection a posteriori?



 $\begin{array}{ll} {\rm GCM: symmetry-breaking solutions} & |{\rm GS}\rangle_{\rm def} & |\omega\rangle_{\rm def} \\ \\ {\rm PGCM: symmetry-conserving solutions} & |{\rm GS}\rangle_{\rm sym} & |\omega\rangle_{\rm sym} \\ \\ {\rm Variational treatment of rotations in PGCM!} \end{array}$

PAV GCM: projection of symmetry-breaking solution

- Anomalous spectrum
- Zero-frequency rotations (Goldstone modes)
- Born-Oppenheimer-like approximation



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Rotational state $|\text{ROT}\rangle = \hat{R}(\Omega) |\text{GS}\rangle_{\text{def}}$



 $|\mathrm{GS}\rangle_{\mathrm{def}}$ $|\omega\rangle_{\rm def}$ GCM : symmetry-breaking solutions $|\mathrm{GS}\rangle_{\mathrm{sym}} \quad |\omega\rangle_{\mathrm{sym}}$ PGCM : symmetry-conserving solutions Variational treatment of rotations in PGCM ! PAV GCM: projection of symmetry-breaking solution Anomalous spectrum • Zero-frequency rotations (Goldstone modes) ٠ Born-Oppenheimer-like approximation ٠ $|\text{ROT}\rangle = \hat{R}(\Omega) |\text{GS}\rangle_{\text{def}}$ Rotational state Non-vanishing Coupling $a_{\omega} = \langle \text{ROT} | \omega \rangle_{\text{def}} \quad \langle \text{ROT} | \omega \rangle_{\text{sym}} = 0$



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Similar results in GCM and RPA

- Does not depend on the many-body method
- Consequence of deformed ground state



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Similar results in GCM and RPA

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Rotations must be treated variationally!

- PGCM already does
- Projected QRPA needed

<u>Outline</u>

Introduction

- Giant Resonances
- GCM and RPA

Random Phase Approximation

- Theoretical introduction
- Angular Momentum Projection

Results

- Rotation-Vibration coupling
- Comparison to ab initio PGCM

Conclusions



(Q) R P A

Symmetry breaking

GCM

Symmetry conserving





Symmetry conserving





(1) [Erler, PhD Thesis, TUD, 2012]

Symmetry conserving





Symmetry conserving

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Symmetry conserving

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(1) [Erler, PhD Thesis, TUD, 2012]

Symmetry conserving

Thank you for the attention

<u>ces</u>

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Gianluca Colò Danilo Gambacurta



Robert Roth Achim Schwenk Alexander Tichai Backup slides





Extended derivation

Standard RPA derivation

$$\begin{aligned} \langle n|T|\text{RPA} \rangle &= \langle \text{RPA}|TQ_n^{\dagger}|\text{RPA} \rangle \\ &= \langle \text{RPA}|\left[T, Q_n^{\dagger}\right]|\text{RPA} \rangle \\ &\approx \langle \text{HF}|\left[T, Q_n^{\dagger}\right]|\text{HF} \rangle \\ &= \sum_{ph} \left\{ X_{ph}^n \left\langle h|T_{\lambda\mu}|p \right\rangle + Y_{ph}^n \left\langle p|T_{\lambda\mu}|h \right\rangle \right\} \end{aligned}$$

Projected RPA derivation

$$\begin{split} \langle \operatorname{RPA}|T_{\lambda\mu}P^{J}_{K_{0}-\mu,K}|n\rangle &= \langle \operatorname{RPA}|T_{\lambda\mu}P^{J}_{K_{0}-\mu,K}Q^{\dagger}_{n}|\operatorname{RPA}\rangle \\ &= \langle \operatorname{RPA}|T_{\lambda\mu}P^{J}_{K_{0}-\mu,K}Q^{\dagger}_{n} - Q^{\dagger}_{n}P^{J}_{K_{0}-\mu,K}T_{\lambda\mu}|\operatorname{RPA}\rangle \\ &\approx \langle \operatorname{HF}|T_{\lambda\mu}P^{J}_{K_{0}-\mu,K}Q^{\dagger}_{n} - Q^{\dagger}_{n}P^{J}_{K_{0}-\mu,K}T_{\lambda\mu}|\operatorname{HF}\rangle \\ &= \sum_{\mathrm{ph}} X^{\mathrm{ph}}\langle \operatorname{HF}|T_{\lambda\mu}P^{J}_{K_{0}-\mu,K}a^{\dagger}_{p}a_{h}|\operatorname{HF}\rangle + Y^{\mathrm{ph}}\langle \operatorname{HF}|a^{\dagger}_{h}a_{p}P^{J}_{K_{0}-\mu,K}T_{\lambda\mu}|\operatorname{HF}\rangle \end{split}$$

Reduced transition amplitudes

$$\langle \text{RPA} || T_{\lambda} || n \rangle = (2J_0 + 1) N_0 N_n (-1)^{J_0 - K_0} \sum_{\text{ph}} \sum_{\mu} \left[X_n^{\text{ph}} + (-1)^{\mu} Y_n^{\text{ph}} \right] \begin{pmatrix} J_0 & \lambda & J \\ -K_0 & \mu & K_0 - \mu \end{pmatrix} \int_{-1}^{1} d(\cos\beta) \, d_{K_0 - \mu, K}^J(\beta) \langle \text{HF} | T_{\lambda \mu} e^{i\beta J_y} a_p^{\dagger} a_h | \text{HF} \rangle$$

$$N_{0} = \left[\int_{-1}^{1} d(\cos\beta) \ d_{K_{0},K_{0}}^{J_{0}}(\beta) \langle \mathrm{HF}|e^{i\beta J_{y}}|\mathrm{HF} \rangle \right]^{-1/2} \qquad N_{n} = \left[\sum_{\mathrm{p,h};\mathrm{p}',\mathrm{h}'} \left(X_{ph}^{n} X_{p'h'}^{n} - Y_{ph}^{n} Y_{p'h'}^{n} \right) \int_{-1}^{1} d(\cos\beta) \ d_{K,K}^{J}(\beta) \langle \mathrm{HF}|a_{h'}^{\dagger}a_{p'}e^{i\beta J_{y}}a_{p}^{\dagger}a_{h}|\mathrm{HF} \rangle \right]^{-1/2}$$

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<u>Comparison to rotational transition density</u>



Overlap maximised at $\alpha \sim 24^{\circ}$

$$\rho_{\alpha} = \left(\frac{\hat{R}_y(\alpha) + \hat{R}_y(-\alpha)}{2}\right)\rho_0$$

<u>AMP identity resolutions</u>

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Rotational overlap



PAV RPA convergence



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PAV RPA convergence


