

# The extended Lipkin model: implementation in a quantum platform and machine learning analysis of its phase diagram

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Supported by MCIN/AEI/10.13039/501100011033  
PID2022-136228NB-C21

Shapes and Symmetries in Nuclei: from Experiment to Theory (SSNET'24)

Orsay (France), November 2024

# Outline

- 1 Motivation
- 2 Quantum Computing
- 3 The Extended Lipkin Model
- 4 Phase diagram and Machine Learning use
- 5 Conclusions

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# Why the Lipkin/Extended Lipkin Model?

- It is a solvable many-body model that allows **to deal with nuclear, molecular, solid state or quantum optics systems.**
- It can be “mapped” into the Interacting Boson Model of Nuclear Physics.
- Nowadays, it is used to **benchmark many-body approximations** because of its great flexibility and simplicity to be solved for large systems.
- The model **has a rich phase diagram.**
- **It is a model highly used in Quantum Information Science** and, therefore, of great interest.
- **It is directly mapped into Pauli matrices without using the Jordan-Wigner mapping.**

# Why Quantum Computing in Nuclear Physics?

- The U.S. Department of Energy in its document Energy Research Opportunity II (2019) establishes “A broad theory program should be supported, which can, e.g., develop methods to address problems in NP using digital quantum computers and quantum simulators, utilize QIS concepts to better understand nuclear phenomena (such as the nuclear many-body problem and hadronization), and develop new QIS applications of importance to nuclear physics”.
- In the future, quantum computers will allow to outperform present computational (classical) capabilities.
- The applications of Quantum Computing in Nuclear Physics are increasing.
- At present (Noisy Intermediate-Scale Quantum era) is a rather active area, but still it is only dealing with small systems and schematic models: Variational Quantum Eigensolver or Quantum Equation of Motion applied to the Lipkin model, Shell-Model applications, implementation/restoration of symmetries.

# Why Quantum Computing in Nuclear Physics?

- Variational Quantum Eigensolver or Quantum Equation of Motion applied to the Lipkin model, Shell-Model applications, implementation/restoration of symmetries.
  - “Cloud Quantum Computing of an Atomic Nucleus”, PRL **120**, 210501 (2018).
  - “Symmetry-Assisted Preparation of Entangled Many-Body States on a Quantum Computer”, PRL **125**, 230502 (2020).
  - “Lipkin model on a quantum computer”, PRC **104**, 024305 (2021).
  - “Simulating excited states of the Lipkin model on a quantum computer”, PRC **106**, 024319 (2022).
  - “Quantum computing of the  ${}^6\text{Li}$  nucleus via ordered unitary coupled cluster”, PRC **106** (2022).
  - “Nuclear shell-model simulation in digital quantum computers”, Sci. Rep. 13:12291 (2023).

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# What a (digital) Quantum Computer is?

A device composed of:

- $m$  **2-level** quantum systems (qubits),
- a set of quantum gates (acting as unitary operators),
- a set of measurement operators (measuring the state of defined subset of qubits),
- a classical control unit which determines which gate should be applied.

I. M. Georgescu, S. Ashhab, and Franco Nori, *Rev. Mod. Phys.* **86**, 153 (2014).

## Implementations

- **Trapped ions.**
- **Superconducting circuits.**
- Nuclear spins (NMR).
- Photons.
- Neutral atoms.
- Cavity arrays.



# Bits and qubits

## Classical vs quantum

Bits: 0 or 1

Qubits:

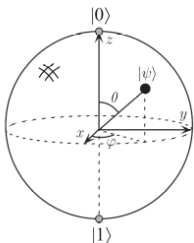
$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$$

$$|\psi\rangle = \cos \frac{\theta}{2} |0\rangle + e^{i\varphi} \sin \frac{\theta}{2} |1\rangle$$

Multiple Qubits:

$$|\psi\rangle = \alpha_{00} |00\rangle + \alpha_{01} |01\rangle + \alpha_{10} |10\rangle + \alpha_{11} |11\rangle$$

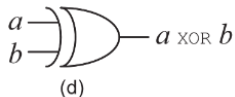
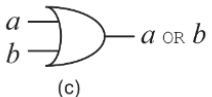
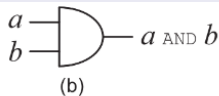
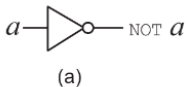
Bell state:  $\frac{|00\rangle + |11\rangle}{\sqrt{2}}$



Bloch sphere

# Gates







## Classical gates



From M. A. Nielsen, I. L. Chuang, Quantum Computation and Quantum Information

# Quantum gates

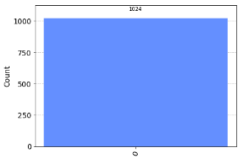
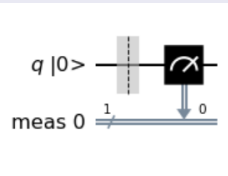
## Single qubit gates

Hadamard		$\frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$
Pauli- $X$		$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$
Pauli- $Y$		$\begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$
Pauli- $Z$		$\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$
Phase		$\begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$
$\pi/8$		$\begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix}$

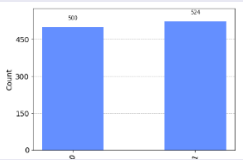
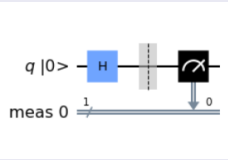
From M. A. Nielsen, I. L. Chuang, Quantum Computation and Quantum Information

# Quantum circuits

## The simplest case



## A step further



Results from Qiskit, <https://www.ibm.com/quantum/qiskit>



# The (Extended) Lipkin Model

## The first appearance

H.J. Lipkin, N. Meshkov, and A.J. Glick, "Validity of many-body approximation methods for a solvable model: (I). Exact solutions and perturbation theory", Nucl. Phys. **62**, 188 (1965).

## The original Hamiltonian

It describes a set of  $N$  particles interacting with a long-range interaction.  $\vec{S} = \sum_{i=1}^N \vec{s}_i$  and  $S = N/2$ .

$$H_L = \varepsilon S_z - \frac{V}{2}(S_+^2 + S_-^2) - \frac{W}{2}(S_+ S_- + S_- S_+)$$

## A convenient way of rewriting the Hamiltonian

$$H_L = (1 - \lambda)(S + S_z) - \frac{\lambda}{N}(S_+ + S_-)^2 = (1 - \lambda)(S + S_z) - 4\frac{\lambda}{N}S_x^2$$

# The ELM as an approximation of the Interacting Boson Model

## The ELM Hamiltonian *à la* CQF

$$H = (1 - \lambda) n_t - \frac{\lambda}{N} Q^{(\alpha)} \cdot Q^{(\alpha)},$$

$$n_t = t^\dagger t \text{ and } Q^{(\alpha)} = (s^\dagger t + t^\dagger s) + \alpha(t^\dagger t)$$

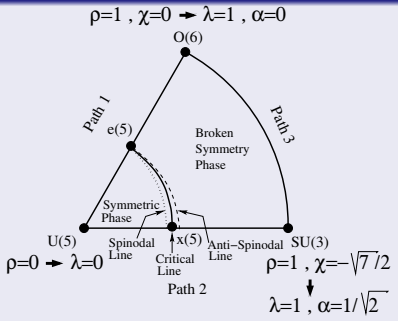
## Equivalence with the Interacting Boson Model

$$S_+ = t^\dagger s$$

$$S_- = s^\dagger t,$$

$$S_z = \frac{1}{2}(t^\dagger t - s^\dagger s).$$

$$H_{EL} = H_L - \frac{\lambda}{N} \left[ \alpha^2 (S_x + S_z)^2 - 2\alpha (S_x (S_x + S_z) + (S_x + S_z) S_x) \right],$$



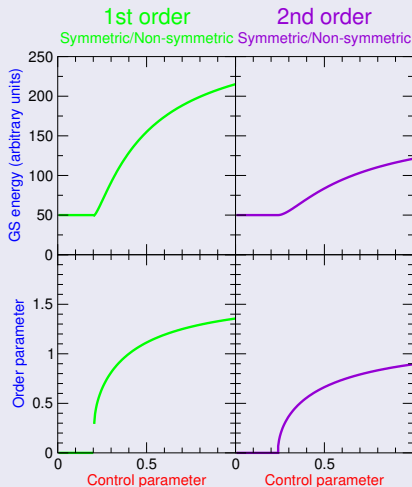




# What a Quantum Phase Transition (QPT) is?

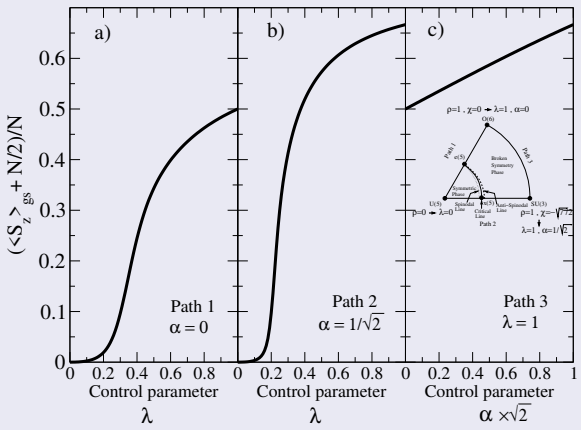
A QPT appears when the ground state a quantum system experiences a sudden change in its structure (order parameter) when a parameter that affects the Hamiltonian (control parameter) slightly changes around its critical value. This transitions are assumed to occurs at zero temperature.

$$\hat{H} = (1 - \xi)\hat{H}_1 + \xi\hat{H}_2$$



# Precursors of a quantum phase transition

## Order parameter



Expectation value of  $S_z + N/2$ , normalized to  $N$ . Calculations are done for  $N = 6$ .

# How to determine the shape/phase of the system?

## The obvious things

- Shape is not really an observable.
- The shape of the system is a property of its **ground state** (it is true that it can be also defined for a excited state).
- It is well defined at the mean-field level.

## A different view

- The shape of the system characterizes its spectrum.
- **An observable depending on the spectrum could encode the shape of the system. That, in general, will happen for the time evolution of the matrix element of a non-eigenstate.**
- Most probably the results will depend of the state and on the used operator. Difficult to determine a priori the best state and operator.
- These types of measurements are the easiest ones in Quantum Computing.

# Our Ansatz

The shape of the system is encoded in the time evolution of an appropriately chosen operator.

## The correlation operator

$$\begin{aligned} C_\nu(i, j, t) &= \langle \Phi_0 | U(t)^\dagger \sigma_\nu^i \sigma_\nu^j U(t) | \Phi_0 \rangle \\ &\quad - \langle \Phi_0 | U(t)^\dagger \sigma_\nu^i U(t) | \Phi_0 \rangle \langle \Phi_0 | U(t)^\dagger \sigma_\nu^j U(t) | \Phi_0 \rangle, \end{aligned}$$

$\sigma_\nu^i$  is a Pauli matrix,  $|\Phi_0\rangle$  is an arbitrary state, not an eigenstate. Also can be used  $S_z$ ,

$$S_z(t) = \langle \Phi_0 | U(t)^\dagger S_z U(t) | \Phi_0 \rangle,$$

## Time evolution as a tool determine the phase diagram of a system

### The evolution operator

$$U(t) = \exp(-i H t)$$

In practice, it is implemented through the **Lie-Trotter-Suzuki decomposition** (Trotter in short)

$$U(t) \approx \left( \prod_k e^{-iH_k t/n_T} \right)^{n_T} = U_T(t, n_T),$$

where the error produced will depend on the commutator  $[H_i, H_j]$  and scale as  $1/n_T$ , where  $n_T$  denotes the number of Trotter steps.

### The Hamiltonian terms

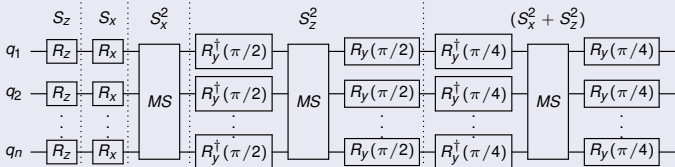
$$H_1 = g_z S_z + g_{zz} S_z^2,$$

$$H_2 = g_x S_x + g_{xx} S_x^2,$$

$$H_3 = g_{xz} (S_x + S_z)^2.$$

# The implementation of the time evolution

## Quantum circuit



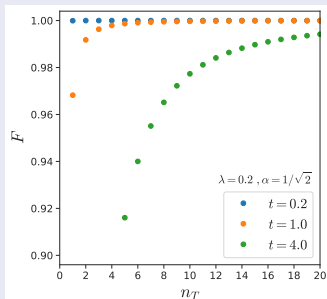
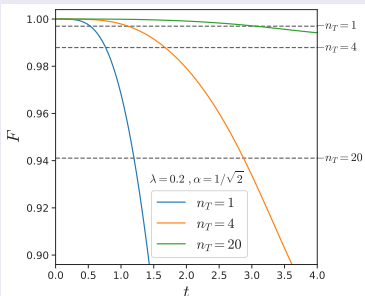
Fidelities of up to 99.9999% for single-qubit gates and up to 99.9% for MS gates.

$$f = \left[ (f_s)^{6N} (f_{MS})^3 \right]^{n_T}.$$

The depth of the circuit depends on  $n_T$ .

# The fidelity

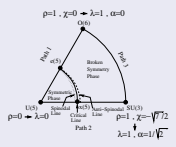
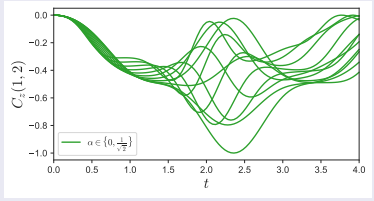
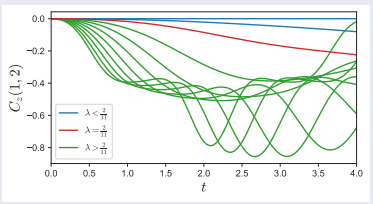
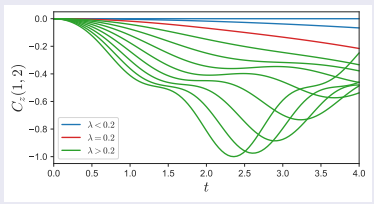
$$F = |\langle \phi(0) | U_T(t, n_T) U(t) | \phi(0) \rangle|^2.$$



Fidelity of the trotterized evolution.  $N = 6$ , initial state  $|\Phi_0\rangle = |\uparrow\downarrow\uparrow\downarrow\uparrow\downarrow\rangle$ .

# Really useful information from time evolution?

## Time evolution of the correlation operator



Time evolution in selected cases for path1, path2, and path3



# Machine learning to recognize the shape of the system

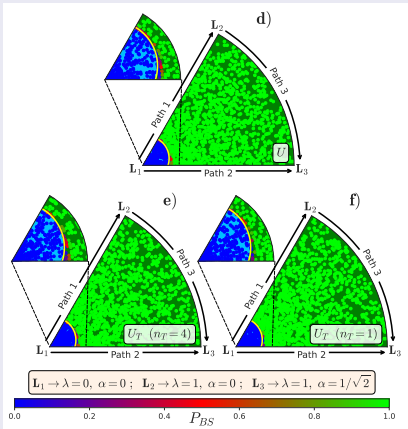
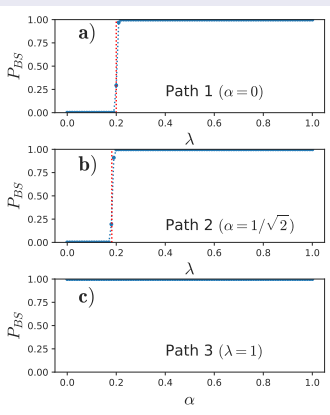
## Machine learning in a classical computer

- Regression
- **Clustering**
- Decision Trees
- Reinforced Learning
- Genetic Algorithms
- **Neural Networks**

## The options

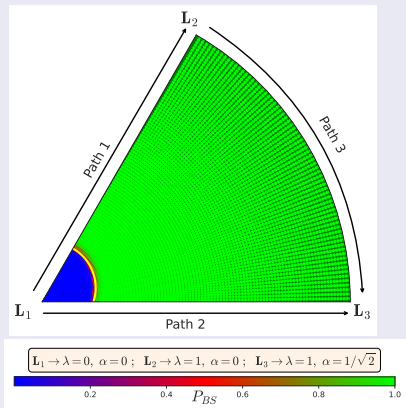
- To use supervised learning.
- To use supervised learning with partial information.
- To use unsupervised learning.

# Results from a Convolutional Neural Network



Using  $C_z(1, 2, t)$  for a system of size  $N = 6$ , with initial state  $|\Phi_0\rangle = |\uparrow\downarrow\uparrow\downarrow\uparrow\downarrow\rangle$ .

# Results from a CNN with partial information



Training the CNN only with points within path 1 ( $\overline{L_1 L_2}$ ), i.e.  $\alpha = 0$ .

# The unsupervised case

## C-means

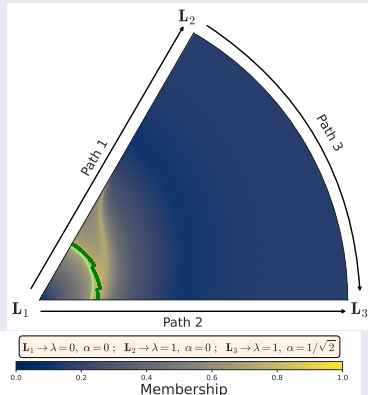
Objective function to be minimized

$$f(\vec{C}) = \sum_i^n \sum_j^k w_{i,j} (\|\vec{x}_i - \vec{c}_j\|)^2$$

where

$$w_{i,j} = \frac{1}{\sum_{j'}^k \left( \frac{\|\vec{x}_i - \vec{c}_j\|}{\|\vec{x}_i - \vec{c}_{j'}\|} \right)^{\frac{2}{m-1}}}$$

In K-means  $w_{i,j} \in \{0, 1\}$ .



Cluster 1 membership for the phase diagram of the ELM



# Conclusions and outlook

- It has been presented the implementation of the Extended Lipkin model in a quantum simulator.
- The phase diagram of the model has been determined calculating the time evolution and using a Convolutional Neural Network.
- **The time evolution method does not need the knowledge of the ground state.**
- Supervised and unsupervised flavours have been presented.

## Further reading

- P. Pérez-Fernández, J.M. Arias, JEGR, and L. Lamata, "A digital quantum simulation of the Agassi model", Phys. Lett. B **829**, 137133 (2022).
- A. Sáiz, JEGR, J.M. Arias, L. Lamata, and P. Pérez-Fernández, "Quantum Simulations of an Extended Agassi Model in Trapped Ions using Machine Learning", Phys. Rev. C **106**, 064322 (2022).
- S. Baid, A. Sáiz, L. Lamata, P. Pérez-Fernández, A. Rios, A.M. Romero, J.M. Arias, and JEGR, "Extended Lipkin model: Proposal for implementation in a quantum platform and machine learning analysis of its phase diagram", Phys. Rev. C **110**, 044318 (2024).

Thank you for your attention