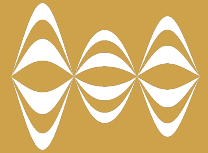


**hexacontatetrapole, E6
 γ decay of ^{53m}Fe**

**AJ Mitchell
Australian National University**



Heavy Ion Accelerators



NCRIS
National Research
Infrastructure for Australia
An Australian Government Initiative

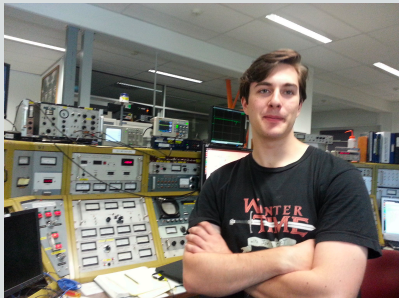


Australian Government
Australian Research Council



**Australian
National
University**

Collaborators



Thomas Palazzo
ANU MPhil student



Greg Lane
Supervisor



AJM

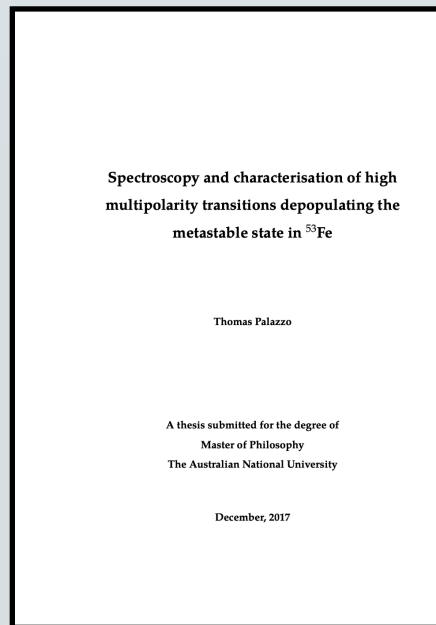


Andrew Stuchbery



Alex Brown
MSU/FRIB (Theory)

... as well as:
A. Akber, B. J. Coombes,
J. T. H. Dowie, M. S. M. Gerathy,
M. W. Reed, T. Kibedi,
and M. O. de Vries.



Featured in Physics Editors' Suggestion PDF HTML

Direct Measurement of Hexacontatetrapole, E6 γ Decay from ^{53}mFe

T. Palazzo, A. J. Mitchell, G. J. Lane, A. E. Stuchbery, B. A. Brown, M. W. Reed, A. Akber, B. J. Coombes, J. T. H. Dowie, T. K. Eriksen, M. S. M. Gerathy, T. Kibédi, T. Tornyi, and M. O. de Vries
Phys. Rev. Lett. **130**, 122503 (2023) – Published 24 March 2023

Physics Synopsis: [Highest-Order Electromagnetic Transition Observed](#)

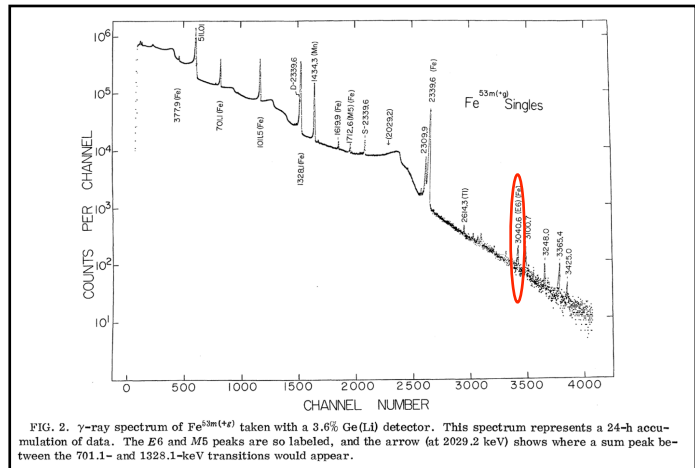
Observations deliver evidence of an exotic “sixth-order” electromagnetic transition in the gamma-ray emission of an iron isotope, a finding that could provide new ways to test nuclear models.

Show Abstract +



Genesis of the project

BA Brown: Several conference presentations drawing attention to $E6$ gamma decay, including Nuclear Data 2013.

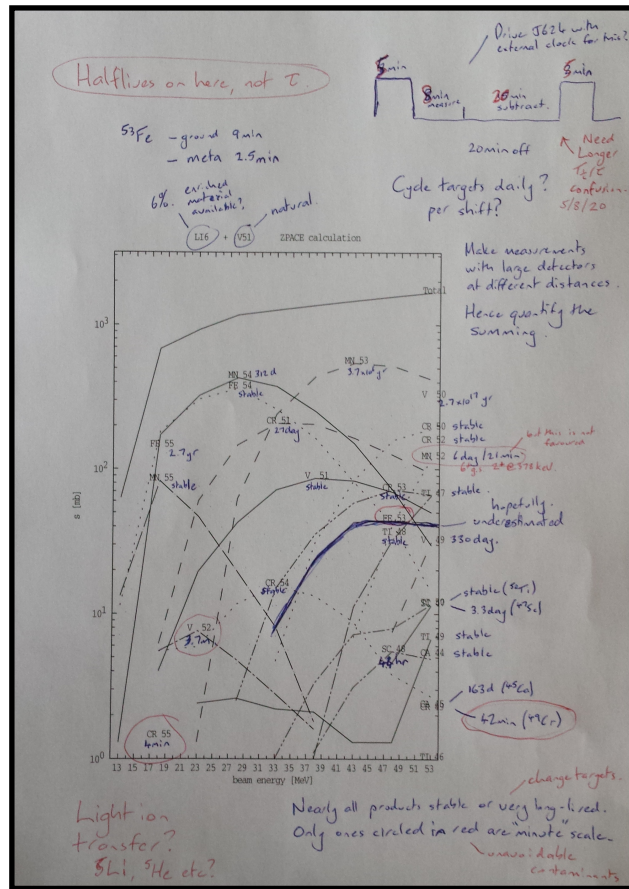


J.N. Black *et al.*,
Phys. Rev. Lett **26**,
451 (1971).

ND2013 audience: “*I bet that’s a sum-peak*”.

AE Stuchbery to T Kibedi: “*I bet we could measure that*”.

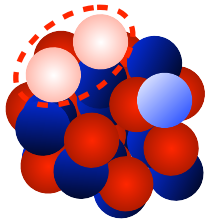
GJ Lane (2015): “*This might work*”.



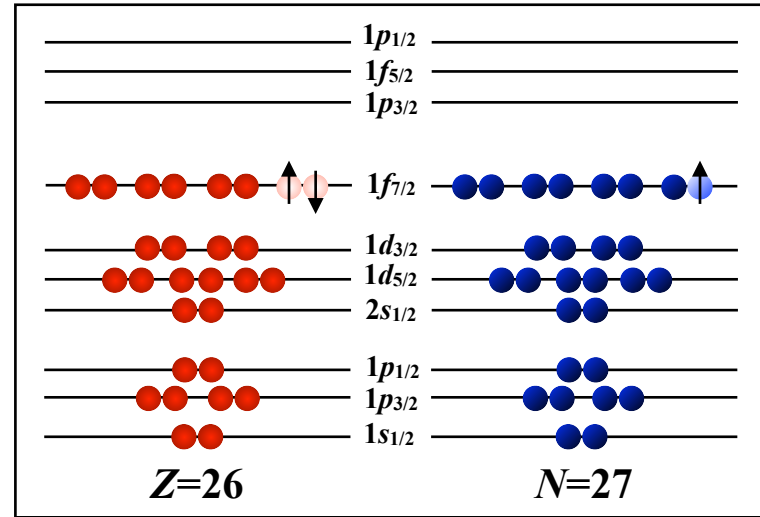
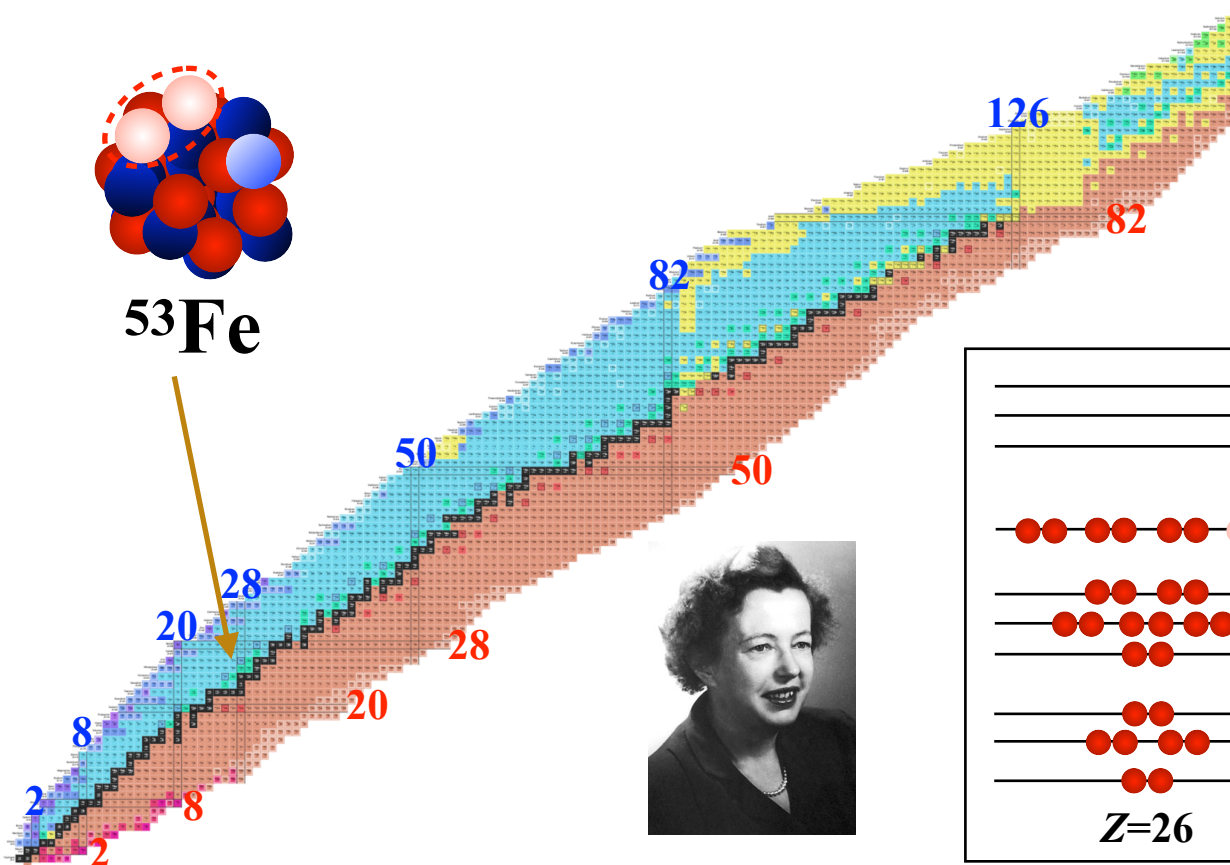
G.J. Lane, private communication (2015)



Nuclear shell model



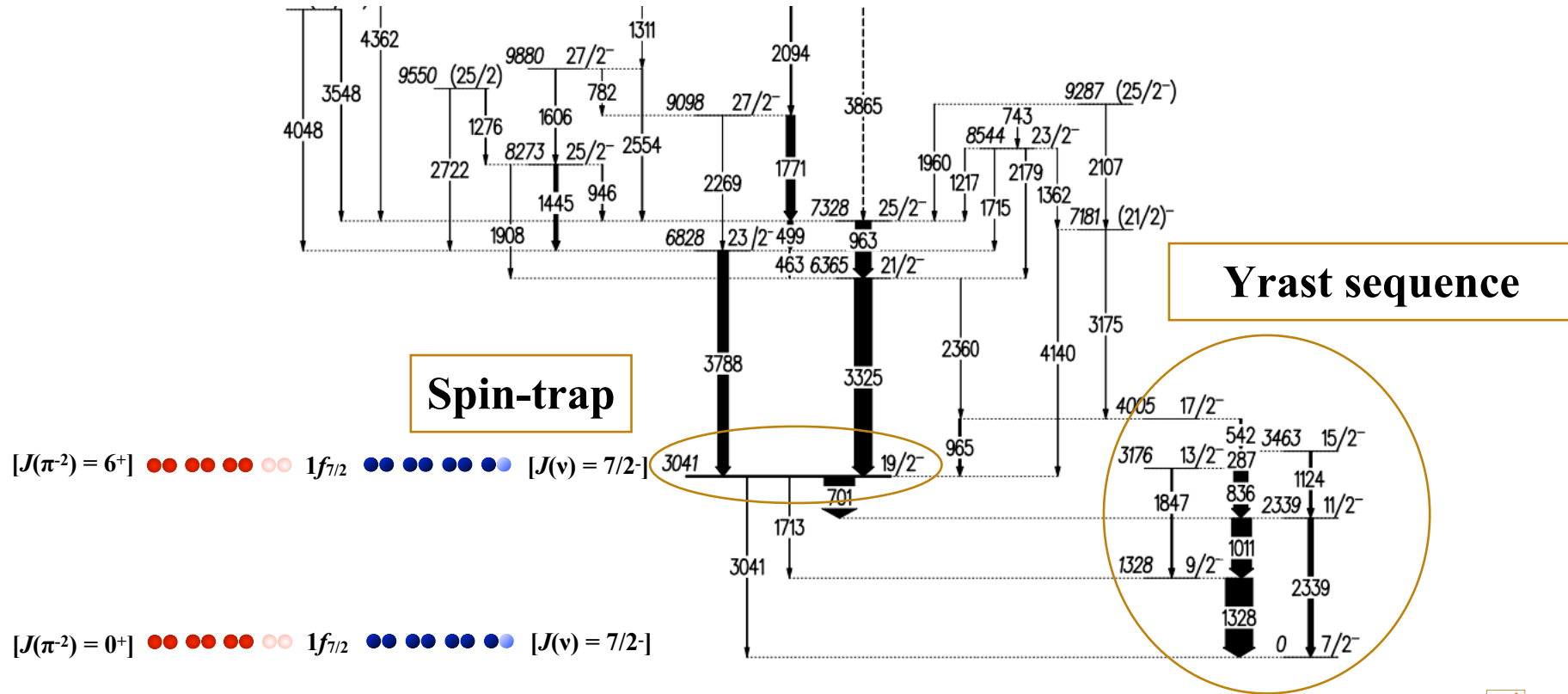
^{53}Fe



<https://people.physics.anu.edu.au/~ecs103/chart/>



High-spin states in ^{53}Fe



^{53m}Fe decay

Selection rules:

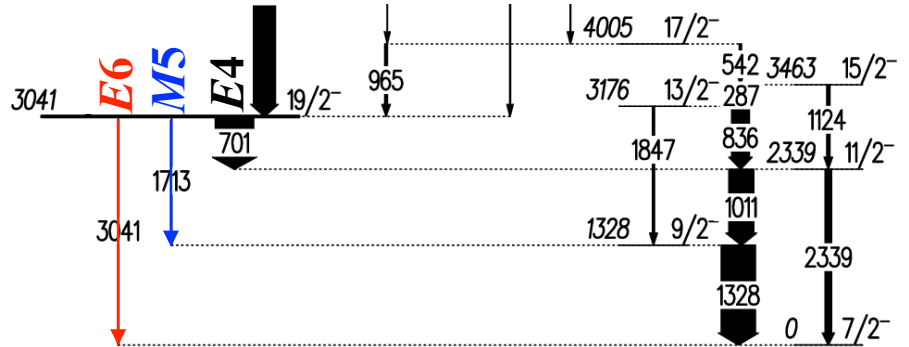
- $|I_i - I_f| \leq L \leq |I_i + I_f|$
- $\Delta P = (-1)^L$ or $(-1)^{L-1}$

In general, gamma decay is dominated by the lowest multipole order permitted:

$$\frac{\lambda(E(L + 1))}{\lambda(EL)} \approx 10^{-5}$$

$$\frac{\lambda(M(L + 1))}{\lambda(ML)} \approx 10^{-5}$$

$$\frac{\lambda(EL)}{\lambda(ML)} \approx 10^2$$



- $L = 1, 2$ are prevalent in atomic and nuclear systems
- $L = 3$ is rare (around 1100 known)
- $L = 4$ is very rare (around 170 known)
- $L = 5$ is very, very rare (around 25 known)
- $L = 6$ is unique (one claim so far)



Australian Heavy Ion Accelerator Facility (HIAF)

Accelerators:

- 14UD tandem accelerator (~ 14 MV)
- LINAC (6 MV)

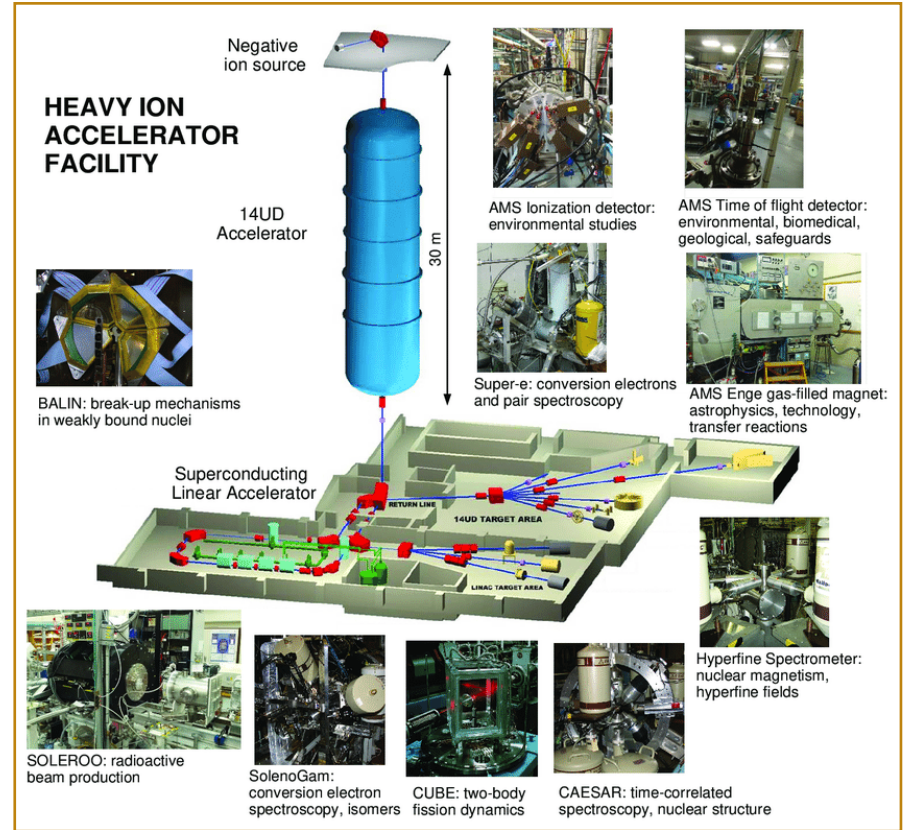
10 beam lines:

- Fundamental research
- Space Irradiation Beam Line

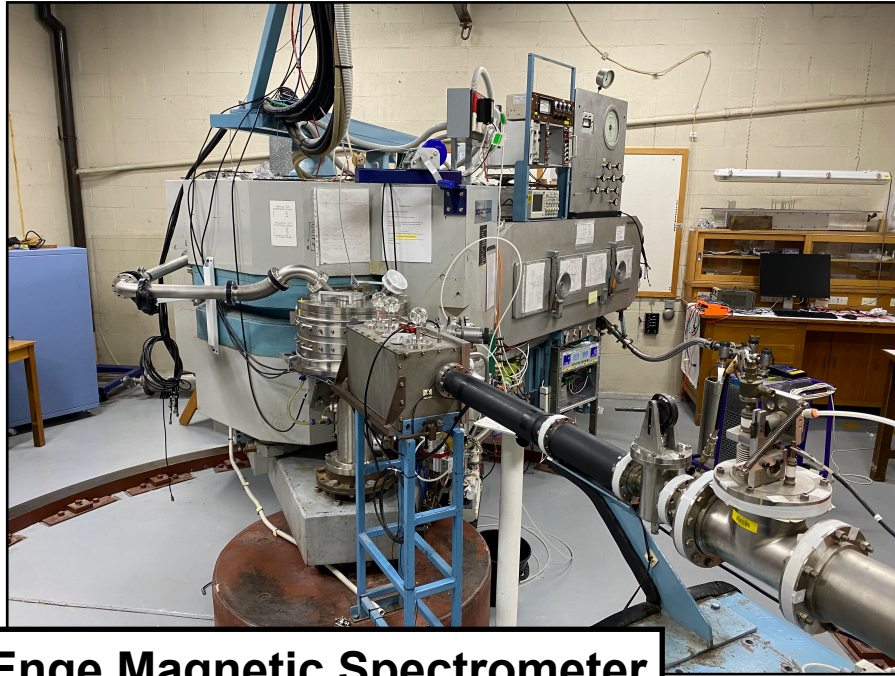
Research:

- Nuclear structure
- Nuclear reaction dynamics
- AMS
- Dark Matter / Astroparticle physics

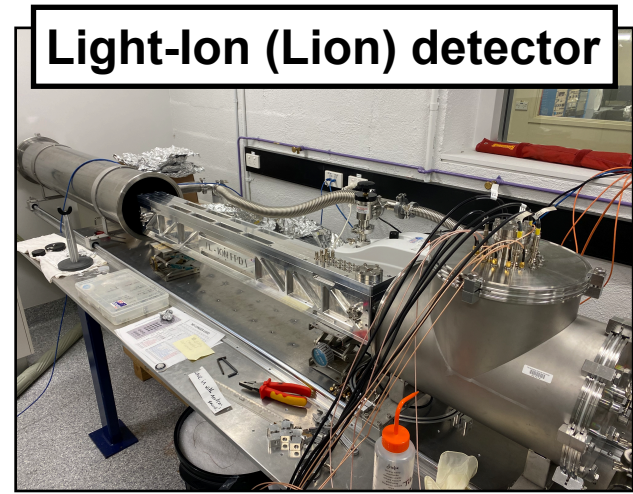
<https://physics.anu.edu.au/tour/nuclear/>



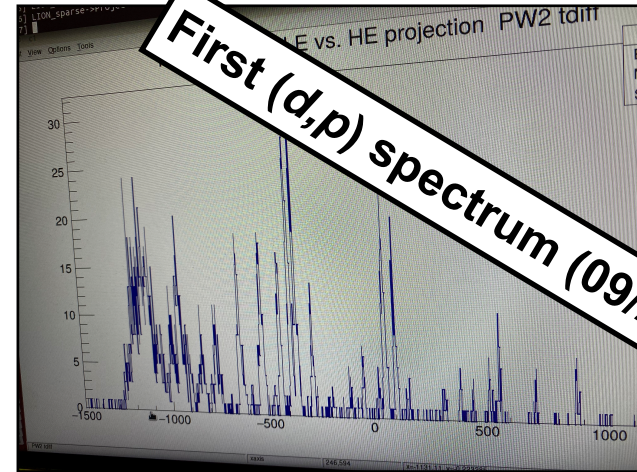
Side-step



Enge Magnetic Spectrometer



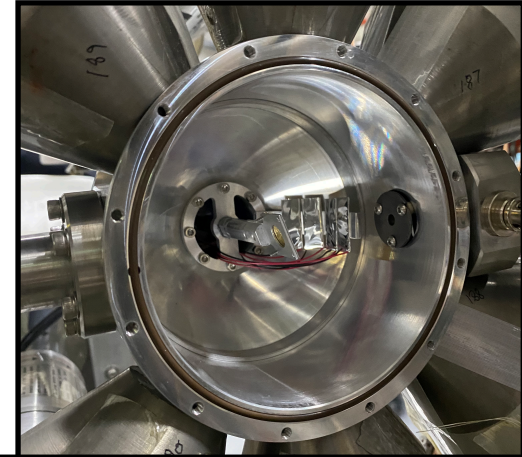
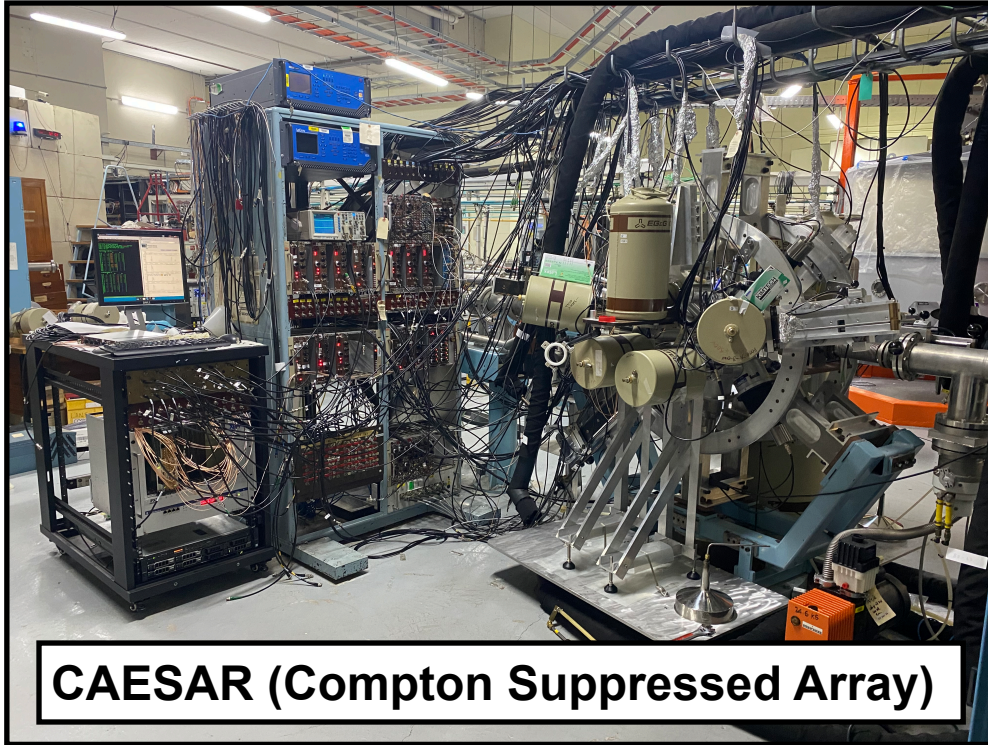
Light-Ion (Lion) detector



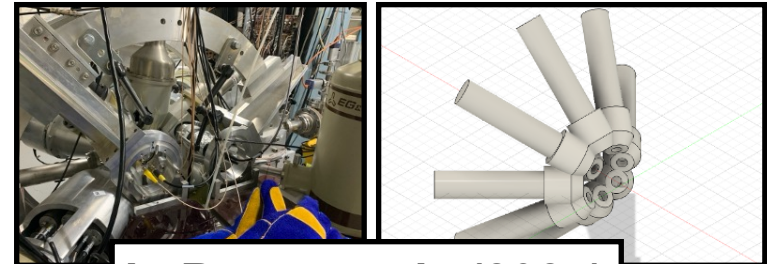
First (d,p) spectrum (09/24)



γ -ray spectroscopy



New target chamber (2022)



LaBr₃ upgrade (2025)



^{53m}Fe γ -ray data

$^{51}\text{V}(^6\text{Li},4n)^{53m}\text{Fe}$, 2 pnA, 50 MeV
10 mg/cm² targets

Repeating irradiation cycle:

7.5 minutes beam on (production)
20 minutes beam off (isomer decay)

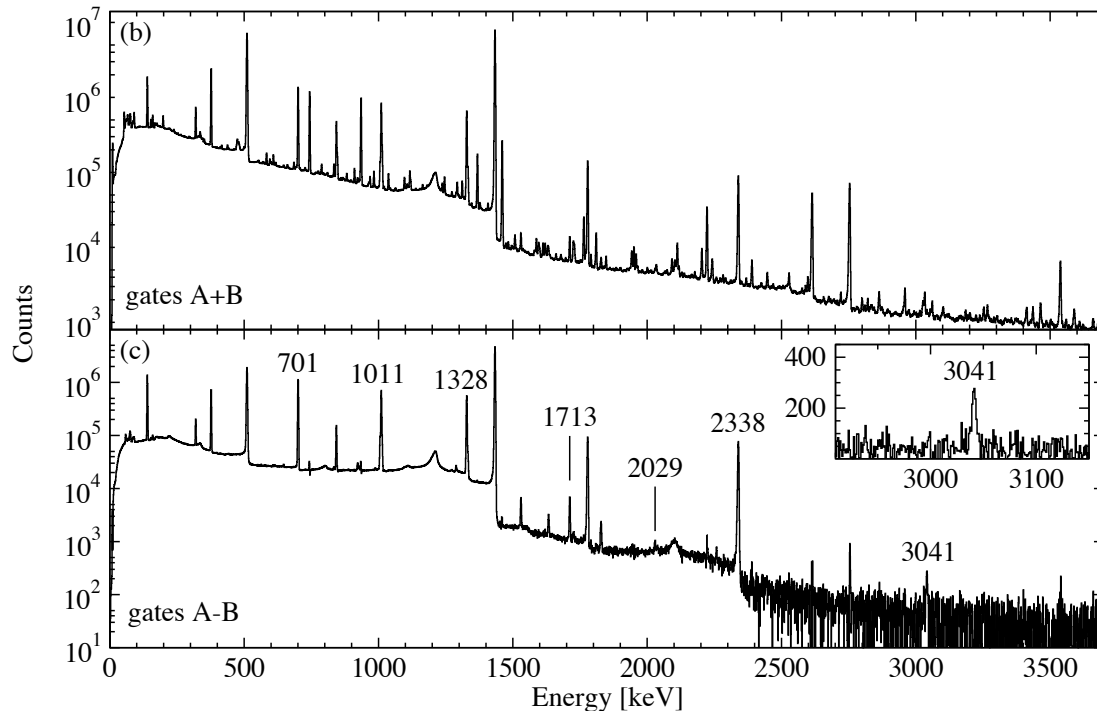
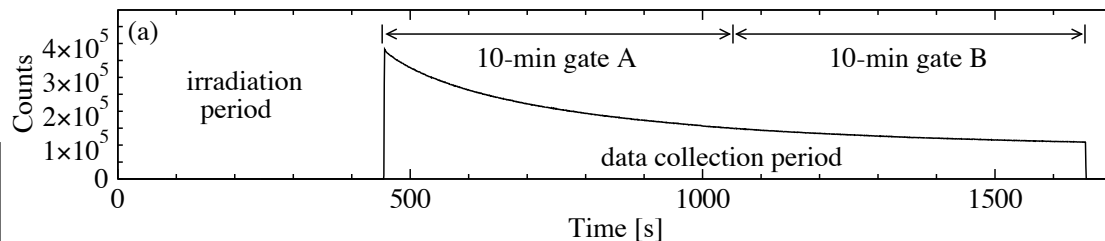
Gates A + B:

~ 10 different nuclides.

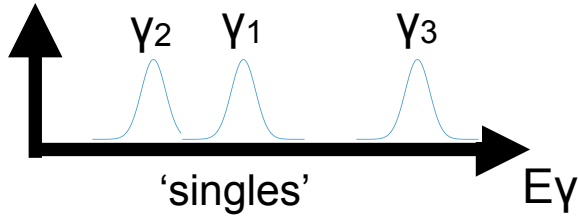
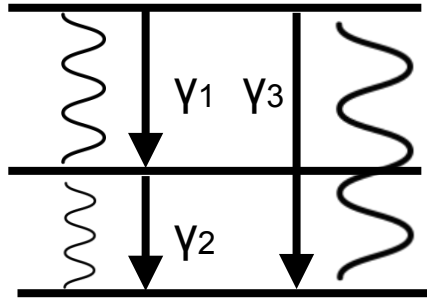
Gates A - B:

Isolate ^{53m}Fe decay.

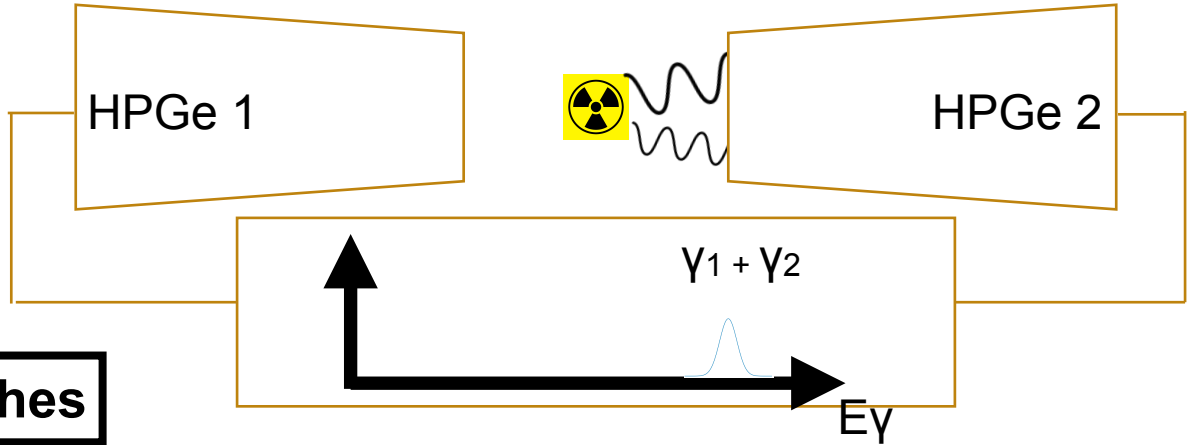
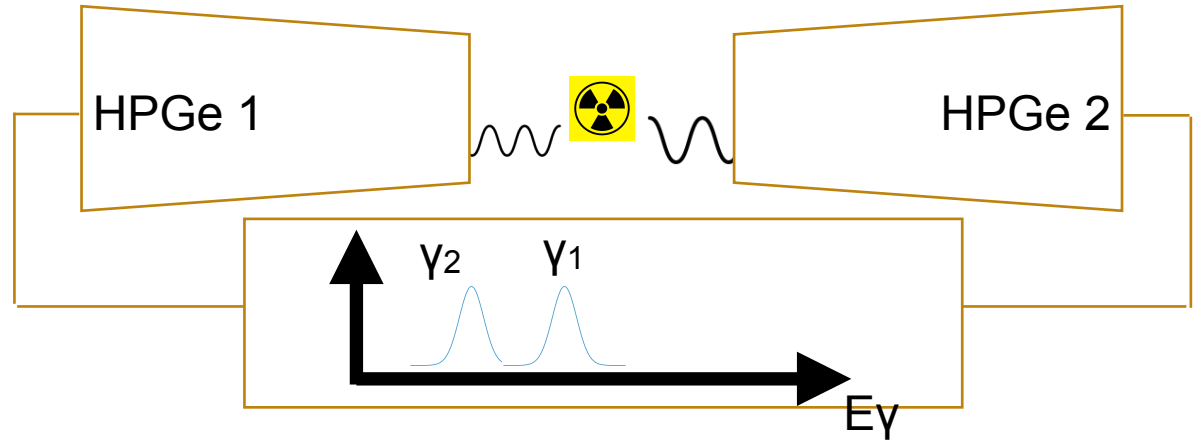
- Known γ rays from ^{53m}Fe
- Including (weak) peak at 3041 keV
- And a feature at 2029 keV



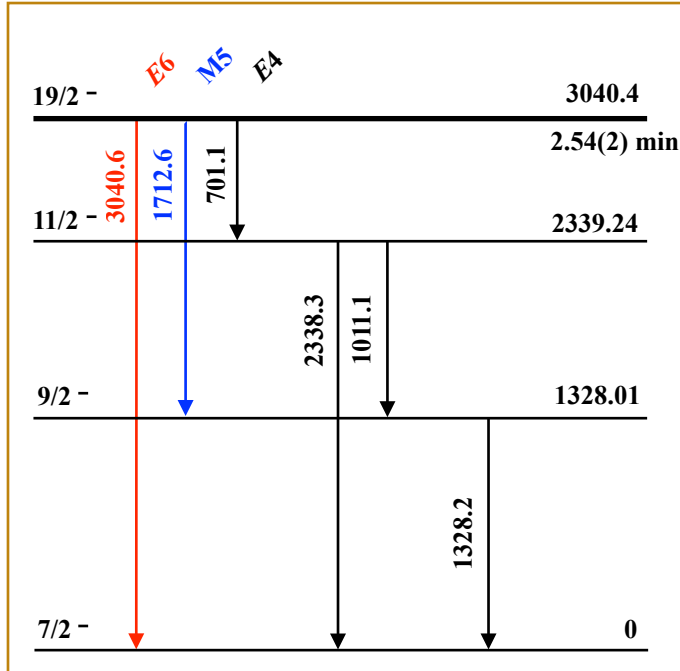
γ -ray summing



Problem for weak branches



γ-ray summing



$$Y_\gamma = I_\gamma + \sum S_i,$$

Method 1: 'Experimental'

$$Y_{2029} = S_{2029} = I_{701} \cdot \varepsilon_{701} \cdot b_{1011} \cdot b_{1328} \cdot \varepsilon_{1328} \times \overline{W}_{701,1328}(0).$$

Expressions that connect sum components to S_{2029} , I_i , b_i , ε_i , $\overline{W}_{i,j}(0)$.

Method 2: 'Geometric'

Considering the change in counting efficiency with moving the radial detectors.

Method 3: 'Computational'

$$Y_{3041} = I_{3041} \cdot \varepsilon_{3041} + I_{701} \cdot b_{2338} \cdot \varepsilon_{701} \cdot \varepsilon_{2338} \cdot \overline{W}_{701,2338}(\theta) + I_{1713} \cdot b_{1328} \cdot \varepsilon_{1713} \cdot \varepsilon_{1328} \cdot \overline{W}_{1713,1328}(\theta) + I_{701} \cdot b_{1011} \cdot b_{1328} \cdot \varepsilon_{701} \cdot \varepsilon_{1011} \cdot \varepsilon_{1328} \times \overline{W}_{701,1011,1328}(\theta).$$

Single expression that combines quantities that were measured in the experiment.

Method 4: 'Monte Carlo'

Decay of ^{53}mFe proceeds via randomised pathways that are weighted by the measured transition branching ratios.

Sum-component is $\approx 50\%$ of the total yield of the 3041-keV γ ray



Results

E_{Level}	E_{γ}	σL	I_{γ}			$B(\sigma\lambda)$ (W.u)		$B(\sigma\lambda)$ ($e^2\text{fm}^{2\lambda}$, $\mu_N^2 \text{fm}^{2\lambda-2}$)	
			This work	Ref. [1]	Ref. [2]	This work	I_{γ} [2]	This work	I_{γ} [2]
3040.4	701.1(1)	$E4$	$\cong 100$	$\cong 100$	$\cong 100$	0.2593(21)	0.2587(21)	$6.46(5) \times 10^2$	$6.44(6) \times 10^2$
	1712.6(3)	$M5$	1.05(5)	0.7(1)	1.3(1)	4.34(21)	5.4(4)	$3.31(16) \times 10^5$	$4.1(3) \times 10^5$
	3040.6(5)	$E6$	0.056(17)	0.020(5)	0.06(1)	0.42(12)	0.45(8)	$2.61(81) \times 10^5$	$2.8(5) \times 10^5$

[1] 10.1103/PhysRevLett.26.451

[2] 10.1103/PhysRevC.11.939

- Sum contributions to $E4$, $M5$, $E6$ accounted for.
- Reduced transition strengths deduced.
- Consistent with the 1975 value of Black *et al*, and an unpublished result in D. Geesaman's PhD thesis.

$$B(XL; J_i \rightarrow J_f) = \frac{L[(2L+1)!!]^2}{8\pi(L+1)} \left(\frac{\hbar c}{E_{\gamma}}\right)^{2L+1} P_{\gamma}(XL; I_i \rightarrow I_f)$$

$$B_w(EL; J_i \rightarrow J_f) = \frac{1}{4\pi} \left(\frac{3}{L_{\gamma} + 3}\right)^2 r^{2L} \{e^2(fm)^{2L}\}$$

$$B_w(ML; J_i \rightarrow J_f) = \frac{10}{\pi} \left(\frac{3}{L_{\gamma} + 3}\right)^2 r^{2L-2} \{\mu_n^2(fm)^{2L-2}\}$$



Interpretation

Reduced transition strength

$$B(XL; J_i \rightarrow J_f) = \frac{\mathcal{M}^2}{2J_i + 1}$$

Reduced matrix element

$\mathcal{A}_{p,n}$
calculated

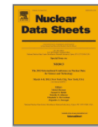
$$\mathcal{M} = \mathcal{A}_p \cdot \epsilon_p + \mathcal{A}_n \cdot \epsilon_n$$

Proton Neutron Effective nucleon charge

$$\epsilon_{p,n} = e_{p,n} + \delta_{p,n}$$

Bare nucleon charge Core-polarisation charge





The Shell-Model Code NuShellX@MSU

B.A. Brown ^a , W.D.M. Rae ^b 

- Shell-model calculations: **restricted** ($f_{7/2}$)¹³ and **full *fp* shell**
- GFPX1A and KB3G Hamiltonians used
 - Restricted model space similar to historical work.
 - Full model space reduced by roughly half.

σL	$\mathcal{A}_p \times 10^3$	$\mathcal{A}_n \times 10^3$	$\mathcal{M} \times 10^3$	$\mathcal{M}_p^{\text{expt.}} \times 10^3$
<i>E4</i>	0.142(17)	0.045(7)	-	0.1137(5)
<i>M5</i>	5.09(76)	-0.11(2)	4.98(76)	2.57(6)
<i>E6</i>	3.52(63)	0.22(4)	-	2.29(35)

Two ‘observations’:

- *E2* transitions are generally enhanced in the full *fp*-shell model space.
- Dominated by the proton component (\mathcal{A}_p and \mathcal{A}_n similar in strong *B(E2)*s in the region).



Proton effective charges

$$\epsilon_p + \epsilon_n \approx 2.0$$

E2 effective charge: $\epsilon_p \approx 1.12$

E4 effective charge: $\epsilon_p = 0.64(6)$

E6 effective charge: $\epsilon_p = 0.62(13)$

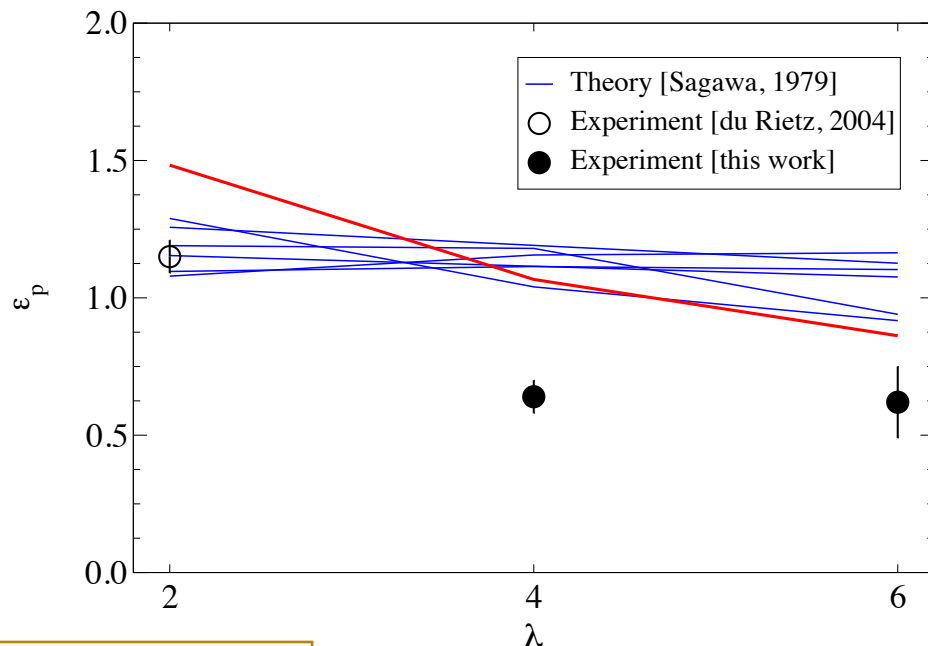
These can be evaluated by considering coupling of valence nucleons to (core) particle-hole excitations.

Choice of—and sensitivity to—the residual particle-hole interaction adopted in the calculation.

Considered for seven interactions by Sagawa.
- Wigner-type (red) interaction closest match.

Excellent agreement for $\lambda = 2$

All of the theoretical results are too large for $\lambda=4$ and $\lambda=6$.



[10.1103/PhysRevC.19.506](https://arxiv.org/abs/10.1103/PhysRevC.19.506)



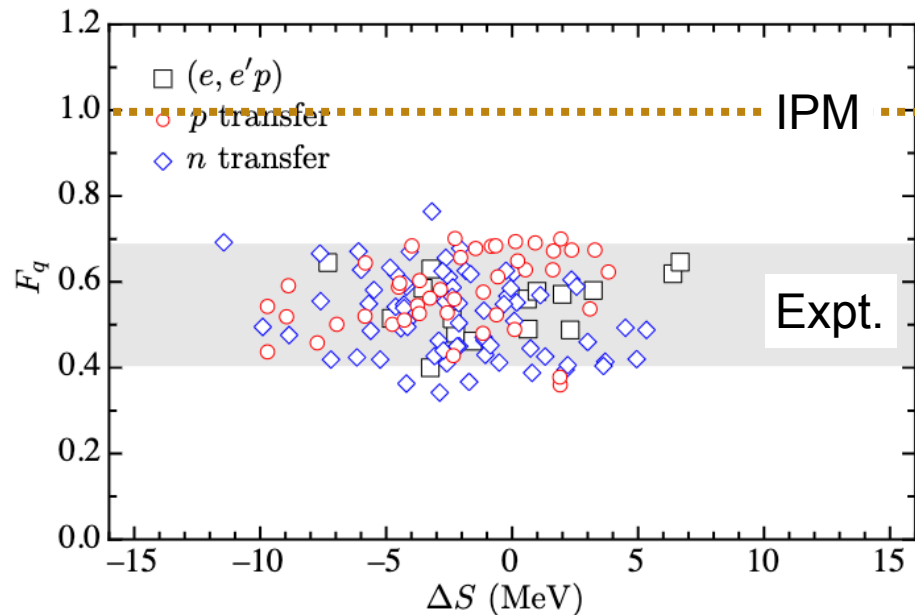
Connection to single-particle behaviour?

$E6$ matrix element and $(e,e'p)$ cross sections:

- both expressed in spectroscopic amplitudes.
- both 'quenched' by a similar magnitude.

Attributed to short- and long-range correlations.

Similarities suggest these are connected.

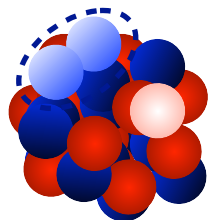


Any model developed to understand quenching of reaction cross sections should be extended to calculations of electromagnetic matrix elements.

[10.1103/PhysRevLett.111.042502](https://doi.org/10.1103/PhysRevLett.111.042502)



Where to from here?

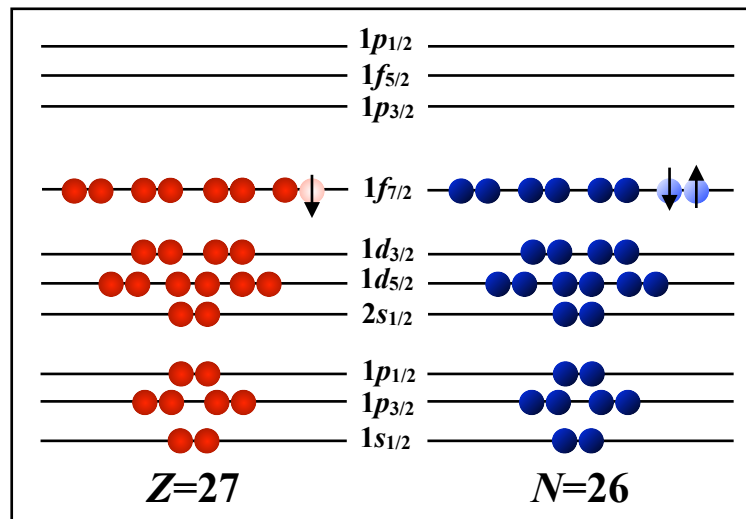
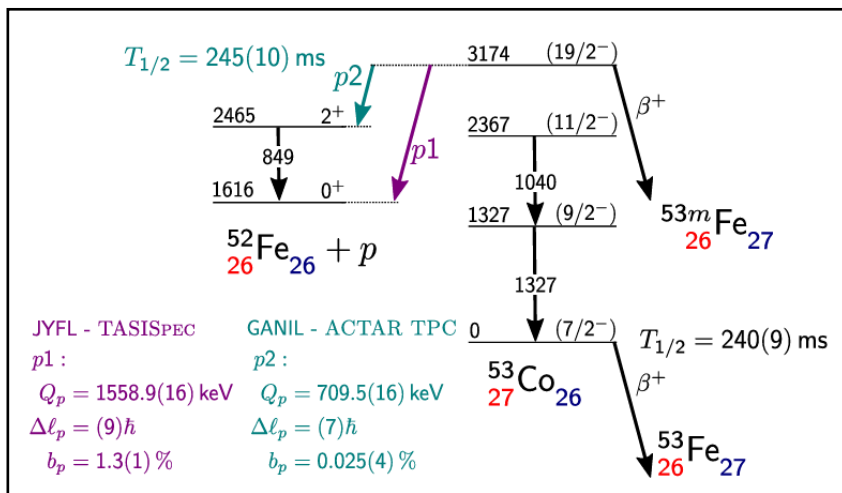


^{53}Co

L. G. Sarmiento, *et al.* Nat. Commun. **14**, 5961 (2023).

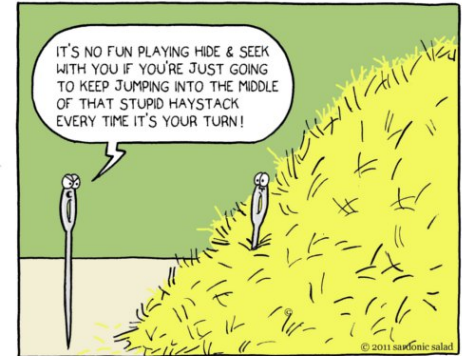
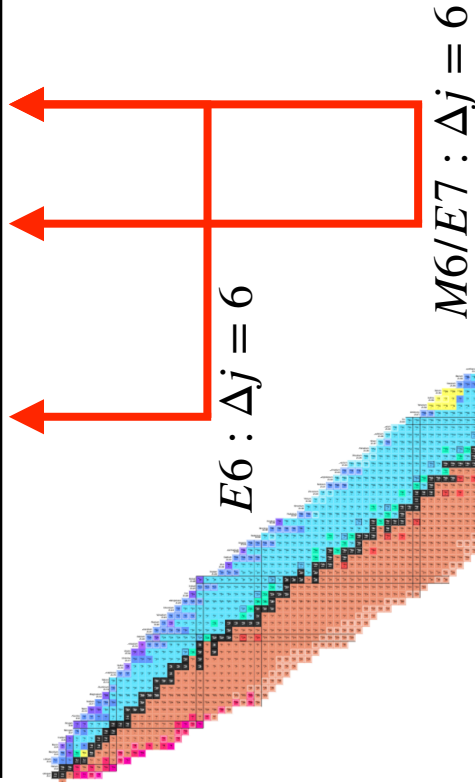


Elucidating the nature of the proton radioactivity and branching ratio on the first proton emitter discovered ^{53m}Co



Where to from here?

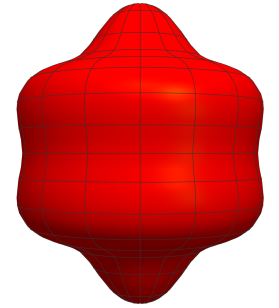
6	126	$0i_{13/2}$
3		$1f_{5/2}$
5		$0h_{9/2}$
1		$2p_{1/2}$
1		$2p_{3/2}$
3		$1f_{7/2}$
5	82	$0h_{11/2}$
0		$2s_{1/2}$
2		$1d_{3/2}$
4		$0g_{7/2}$
2		$1d_{5/2}$
4	50	$0g_{9/2}$
4	(40)	$1p_{1/2}$
1		$0f_{5/2}$
3		$1p_{3/2}$
1	28	



Summary

- **Unambiguous confirmation** of the highest-known transition multipolarity in nature ($E6$).
- **Transition strengths** for the high-multipolarity transitions from the 2.54-minute, $J=19$ -isomer in ^{53}Fe have been determined.
- **Shell-model calculations** highlight the need for cross-shell mixing to explain the experimentally observed strengths.
- **Proton effective charges** are suppressed in high-multipolarity, electric transitions, which are fundamentally different in nature from collective $E2$ transitions.
- **Deeper theoretical investigation required to fully understand the difference.**





THANK YOU

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Australian
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