



**Abnormal Bifurcation of the Double Binding
Energy Differences and Proton-Neutron Pairing:
Nuclei Close to $N = Z$ line from Ni to Rb**

Yiping Wang 王一平

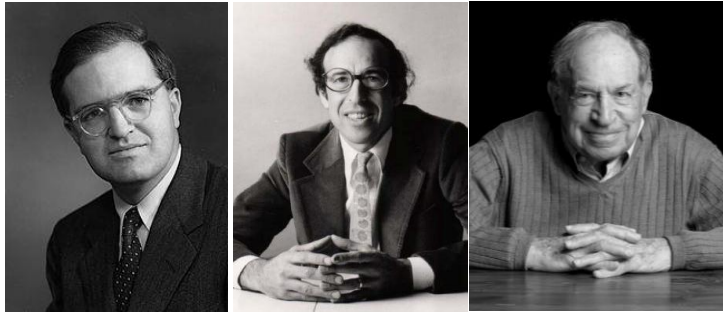
School of Physics, Peking University

Outline

- ◆ Introduction
- ◆ Theoretical Framework
- ◆ Numerical Details
- ◆ Results and
Discussions
- ◆ Summary and Outlook

Pairing correlations

Proposal

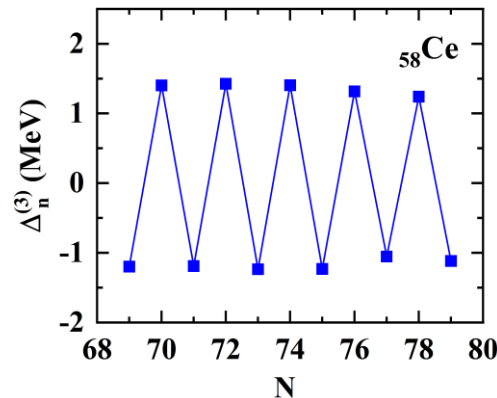
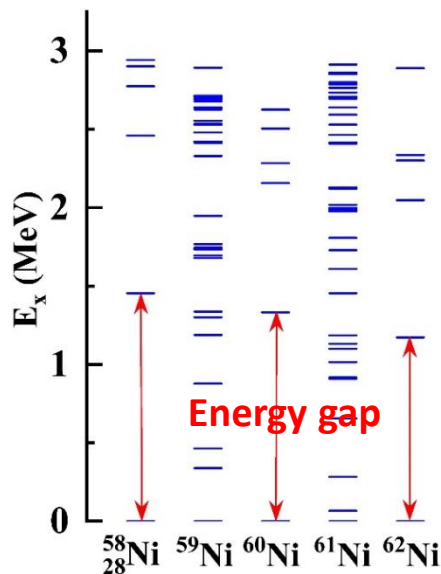


“There is a possible analogy between the excitation spectra of nuclei and those of the superconducting metallic state.”

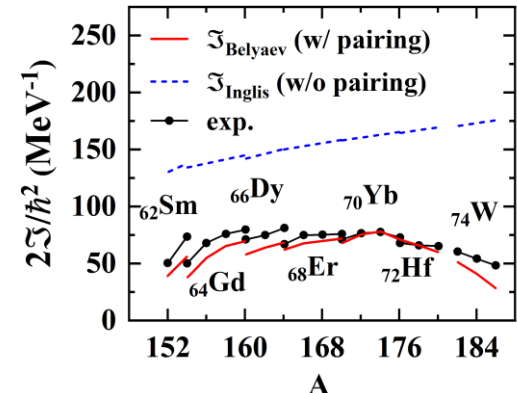
Bohr, Mottelson, Pines, PR 110 (1958) 936

Experimental evidence

neutron-neutron (nn) and proton-proton (pp) pairing



Odd-even mass difference



Moment of inertia

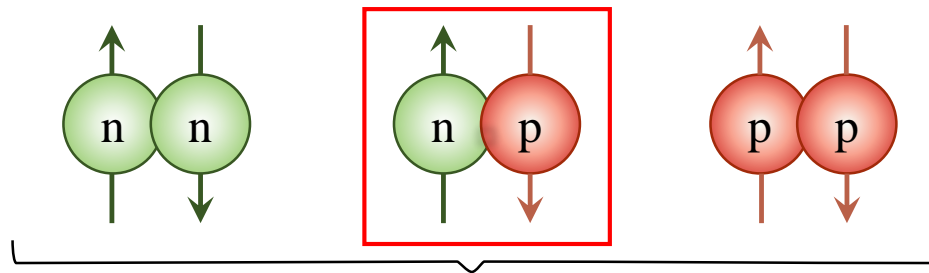
Data from National Nuclear Data Center

Nilsson & Ragnarsson, Shapes and Shells in Nuclear Structure (1995)

Proton-neutron pairing correlations

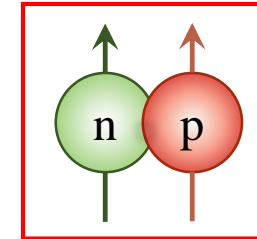
- In atomic nuclei, the **proton-neutron (pn) pairing correlations** may also exist.

Novel pairing phenomenon not found elsewhere in nature



$$T = 1; \quad S = 0; \quad L = 0$$

Isovector Pairing



$$T = 0; \quad S = 1; \quad L = 0$$

Isoscalar Pairing

- Issues of special interest:

- **Existence** of pn pairing condensates;
- **Coexistence** of nn , pp and pn pairing condensates;
- **Different roles** played by isovector and isoscalar pairing correlations.

Frauendorf & Macchiavelli, PPNP 78 (2014) 24

Theoretical studies of pn pairing correlations

□ Theoretical approaches for the study of pn pairing correlations

Exactly Solvable Model

Dukelsky et al., PRL 96 (2006) 072503; Lerma et al., PRL 99 (2007) 032501
Pan, Draayer, PRC 66 (2002) 044314; Romero, Dobaczewski, Pastore, APPB 49 (2018) 347

Shell Model

- Cranked Shell Model Monte Carlo Approach Dean et al., PLB 399 (1997) 1
- Lanczos Method Poves & Martinez-Pinedo, PLB 430 (1998) 203
- Projected Shell Model Sun, EPJA 20 (2004) 133

Phenomenological Mean Field Method

- BCS (Woods-Saxon pot.) + Lipkin-Nogami Satuła et al., PLB 393 (1997) 1
- Hartree-Bogoliubov (Nilsson pot.) Bentley et al., PRC 89 (2014) 034302
- Hartree-Fock-Bogoliubov + Excited Vampir Approach A. Petrovici et al., NPA 647 (1999) 197

Density Functional Theory (DFT)

- Cranked Skyrme Hartree-Fock-Bogoliubov Terasaki et al., PLB 437 (1998) 1
- Skyrme Hartree-Fock + Quartet Condensation Model Negrea, Sandulescu, Gambacurta, PRC 105 (2022) 034325
- Cranked Relativistic Hartree-Bogoliubov + Lipkin-Nogami Afanasjev & Frauendorf, PRC 71 (2005) 064318

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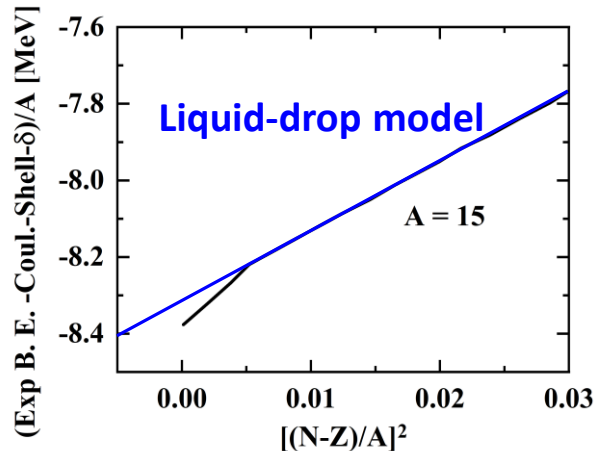
Experimental signals of pn pairing correlations

$N = Z$ Nuclei

ideal place to find pn pairing signals

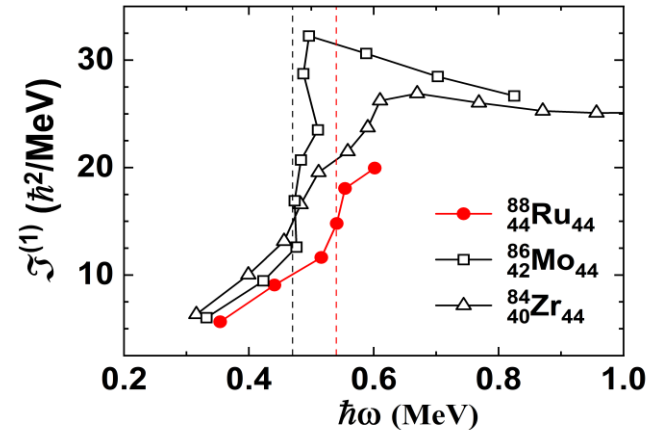
“with enhanced correlations between neutrons and protons that occupy orbitals with the same quantum numbers”

• Additional binding of $N = Z$ nuclei



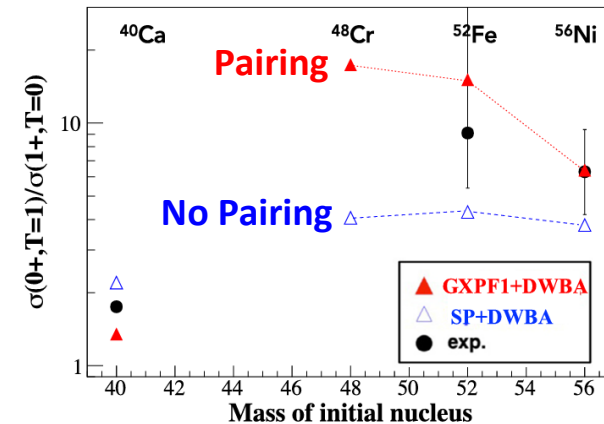
Myers, Swiatecki, Nucl. Phys. 81 (1966) 1

• Delayed alignment in $N = Z$ nuclei



Cederwall et al., Phys. Rev. Lett 124 (2020) 062501

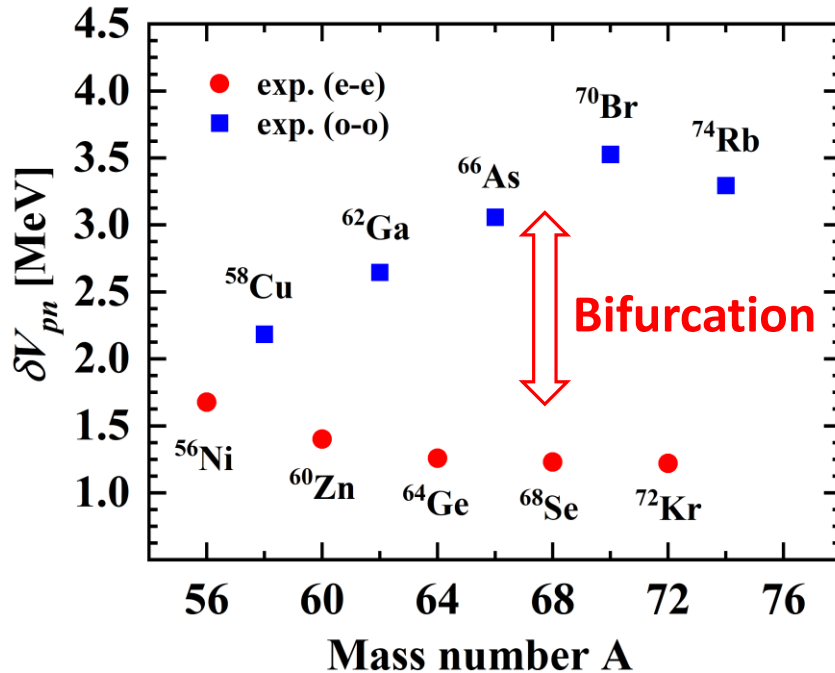
• pn pair transfer cross-section



Le Crom et al., Phys. Lett. B 829 (2022) 137057

Abnormal bifurcation of δV_{pn} for $N = Z$ nuclei

□ State-of-the-art mass measurement reveals an abnormal δV_{pn} bifurcation.



Wang et al., PRL 130 (2023) 192501

Double binding energy difference δV_{pn}

an important mass filter closely related to ...

- **Collectivity and deformation**

Talmi, Rev. Mod. Phys. 34 (1962) 704

Casten, Phys. Rev. Lett. 54 (1985) 1991

Cakirli, Casten, Phys. Rev. Lett. 96 (2006) 132501

- **Underlying shell structure**

Heyde et al., Phys. Lett. B 155 (1985) 303

- **Phase transition behavior**

Federman, Pittel, Phys. Lett. B 69 (1977) 385

Federman, Pittel, Phys. Lett. B 77 (1978) 29

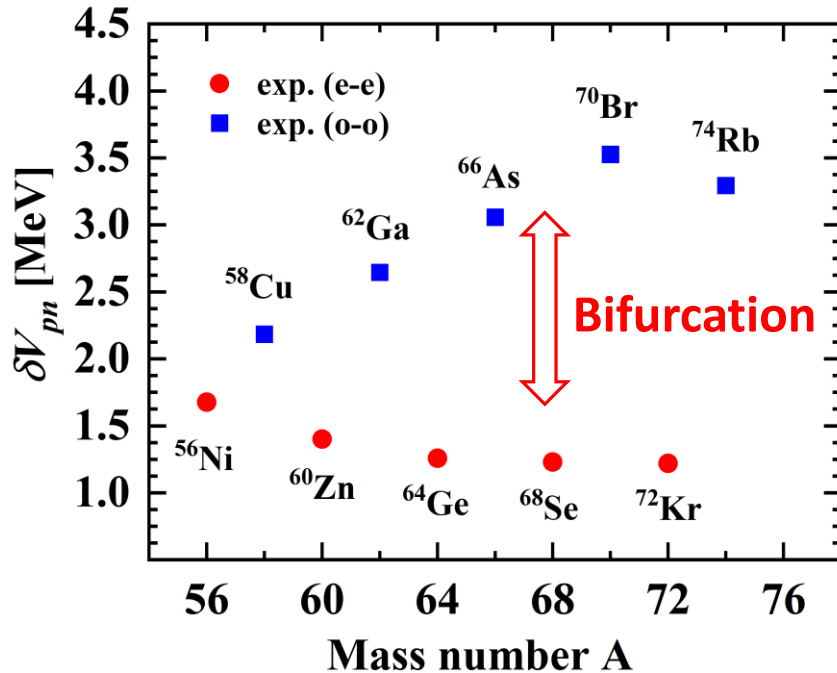
$$\delta V_{pn}^{o-o}(N, Z) = \frac{\partial^2 B}{\partial N \partial Z} = B(N, Z) - B(N-1, Z) - B(N, Z-1) + B(N-1, Z-1)$$

$$\delta V_{pn}^{e-e}(N, Z) = \frac{\partial^2 B}{\partial N \partial Z} = \frac{1}{4} [B(N, Z) - B(N-2, Z) - B(N, Z-2) + B(N-2, Z-2)]$$

Van Isacker et al., PRL 74 (1995) 4607

Abnormal bifurcation of δV_{pn} for $N = Z$ nuclei

State-of-the-art mass measurement reveals an abnormal δV_{pn} bifurcation.



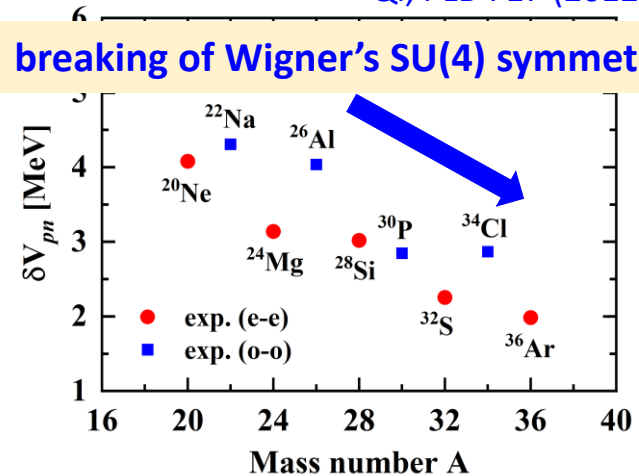
Wang et al., PRL 130 (2023) 192501

Study of δV_{pn} along $N = Z$ line deepens our understanding of nuclear force.

Warner et al., Nat. Phys. 2 (2006) 311

Qi, PLB 717 (2012) 436

breaking of Wigner's SU(4) symmetry



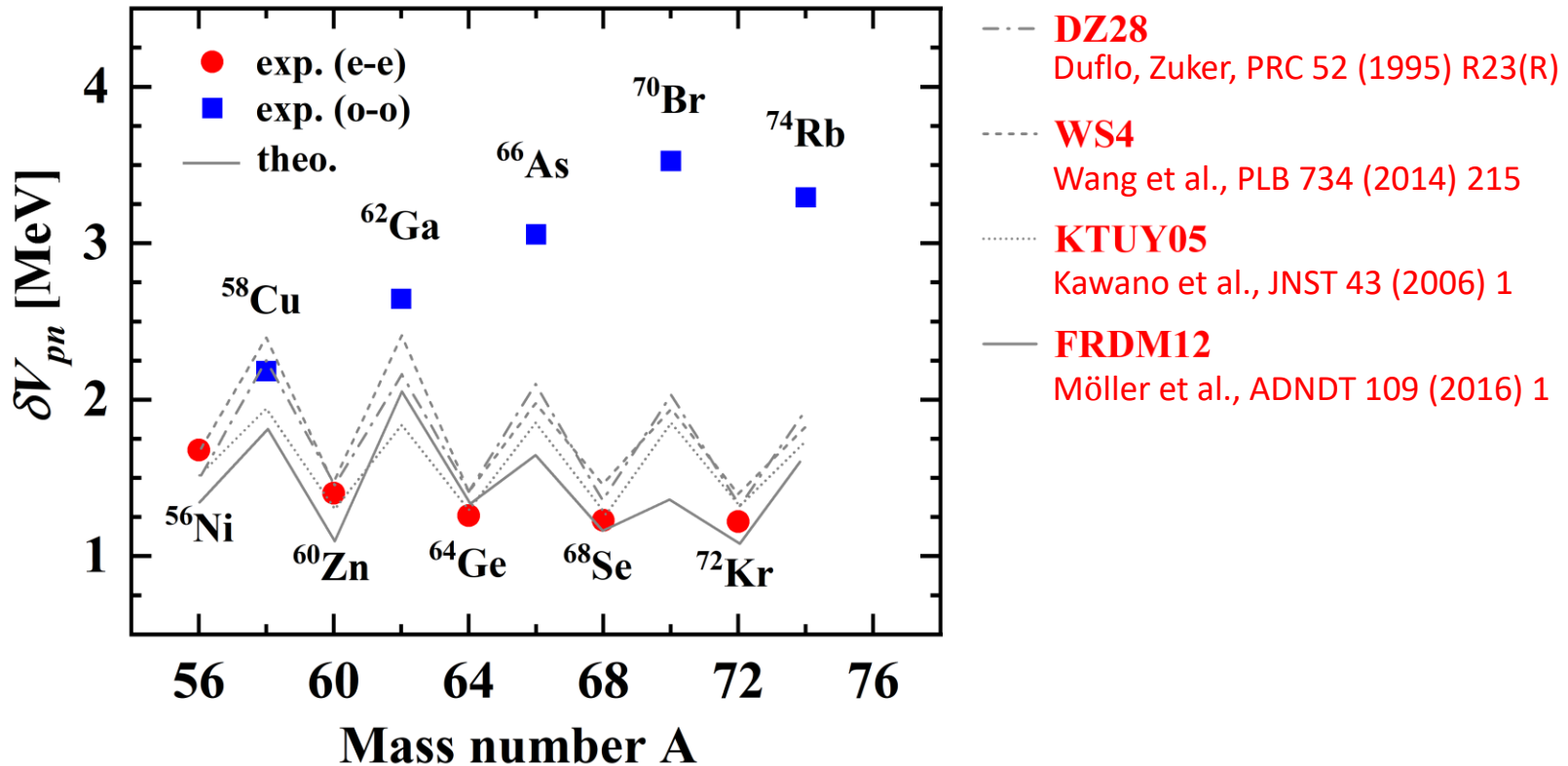
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Van Isacker et al., PRL 74 (1995) 4607

Bring challenges to the nuclear theories

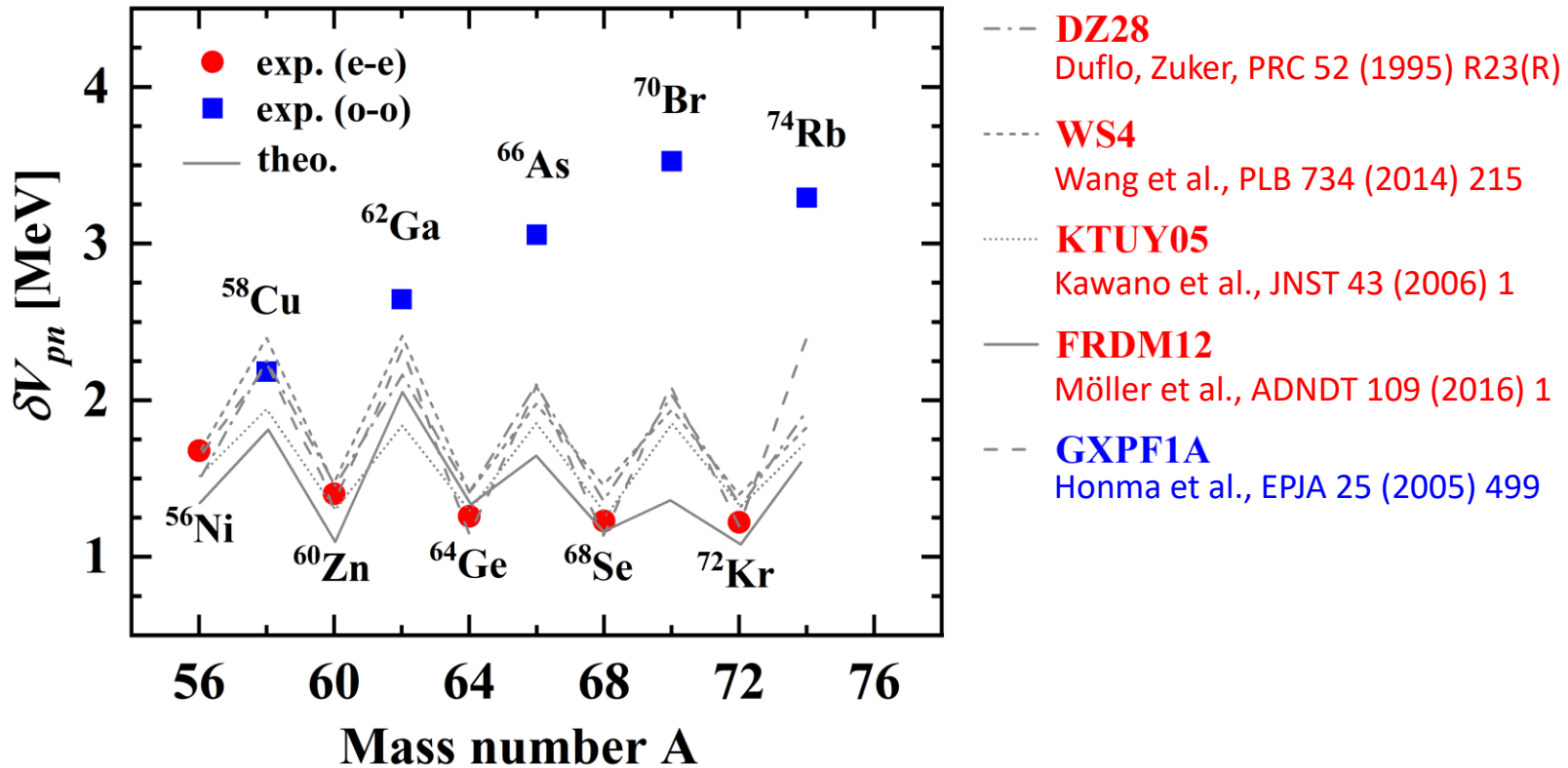
- Frequently used mass models (**macroscopic-microscopic models**, **shell model**, and **DFTs**) fail to describe the δV_{pn} bifurcation.



Wang et al., PRL 130 (2023) 192501

Bring challenges to the nuclear theories

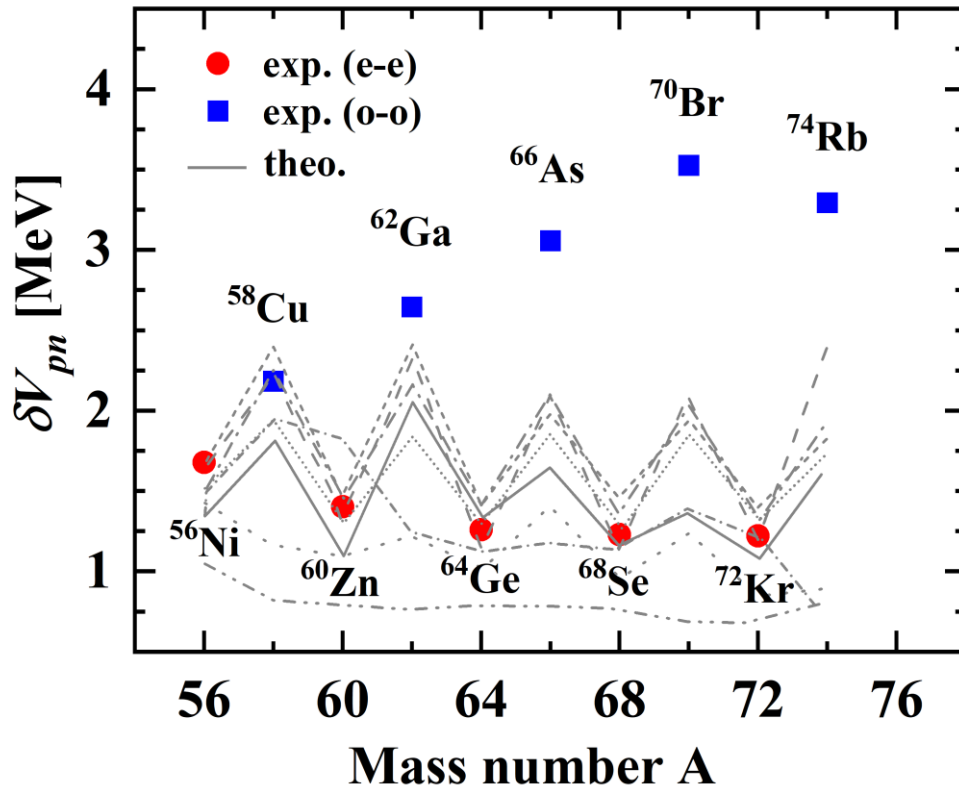
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Wang et al., PRL 130 (2023) 192501

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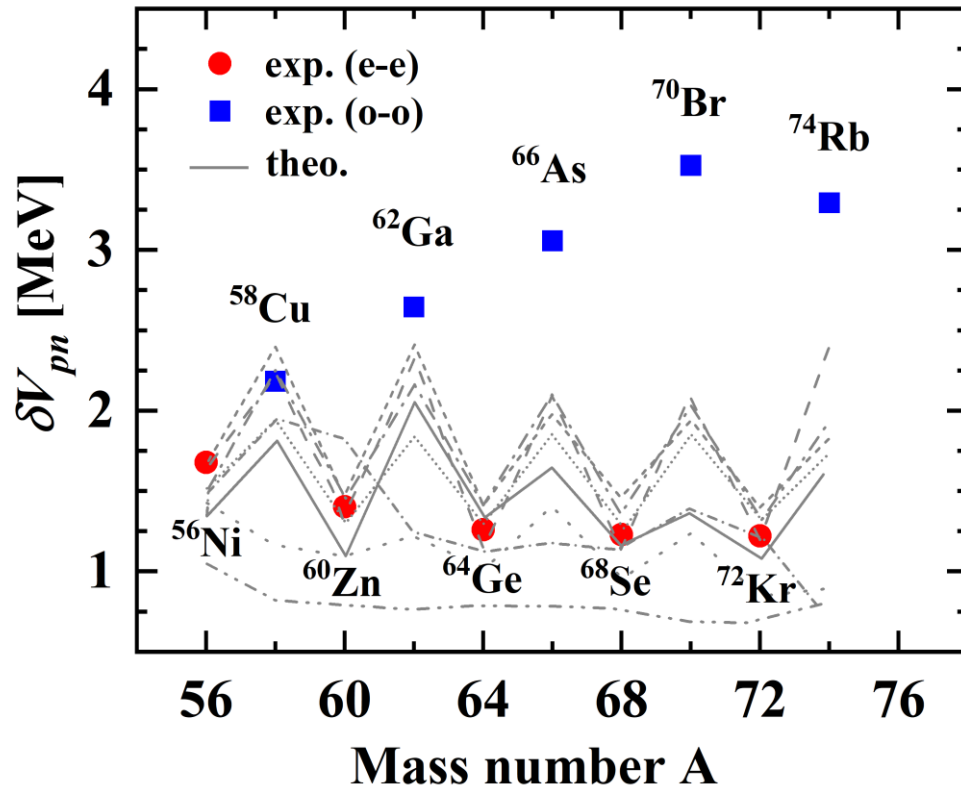


Wang et al., PRL 130 (2023) 192501

- **DZ28**
Duflo, Zuker, PRC 52 (1995) R23(R)
- .-.- **WS4**
Wang et al., PLB 734 (2014) 215
- **KTUY05**
Kawano et al., JNST 43 (2006) 1
- **FRDM12**
Möller et al., ADNDT 109 (2016) 1
- - - **GXPF1A**
Honma et al., EPJA 25 (2005) 499
- **HFB27**
Kortelainen et al., PRC 85 (2012) 024304
- .-.- **UNEDF1**
Goriely et al., PRC 88 (2013) 061302(R)
- .-.- **RMF**
Geng et al., PTP 113 (2005) 785

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Wang et al., PRL 130 (2023) 192501

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Goriely et al., PRC 88 (2013) 061302(R)
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Geng et al., PTP 113 (2005) 785

- What is the missing piece in these models?

maybe pn pairing correlations.

DFT + Shell-Model-Like Approach (SLAP)

- The nuclear DFT starts from a universal density functional and has achieved great successes in describing many nuclear phenomena.

a promising framework to microscopically consider pn pairing correlations

Vautherin, Brink, PRC 5 (1972) 626

Meng (Ed.), *Relativistic Density Functional for Nuclear Structure* (2016)

- Shell-Model-Like Approach *an important tool to treat pairing correlations*

- Advantages:

Yang, Zeng, Acta Physica Sinica 20 (1964) 846; Zeng, Cheng, NPA 405 (1983) 1

Volya et al., PLB 509 (2001) 37

- good particle number;
- treating the blocking effects exactly.

Particle-Number-Conserving Method

Exact Pairing Method

- Different Versions:

only nn and pp pairing correlations considered

- Cranking Nilsson Model + SLAP

Zeng et al., PRC 50 (1994) 1388; He et al., NPA 817 (2009) 45; Zhang et al., NPA 816 (2009) 19

- Nonrelativistic DFTs + SLAP

Pillet et al., NPA 697 (2002) 141; Liang et al., PRC 92 (2015) 064325

- Relativistic DFTs (RDFT) + SLAP

Meng et al., FPC 1 (2006) 38; WYP, Meng, PLB 841 (2023) 137923; Xu et al., PRL 133 (2024) 022501

- ...

This work

a) A new version of RDFT + SLAP is developed, which allows a microscopic treatment of the nn , pp and pn pairing correlations simultaneously.

b) The developed approach is applied to investigate the abnormal δV_{pn} bifurcation for the $N = Z$ nuclei from Ni to Rb.

RDFT-SLAP: Many-body Hamiltonian

In RDFT-SLAP, the many-body Hamiltonian reads

$$\hat{H} = \hat{H}_0 + \hat{H}_{\text{pair}}.$$

□ One-body part:
$$\hat{H}_0 = \sum_{k>0} \left[\varepsilon_k^\pi (a_k^\dagger a_k + a_{\bar{k}}^\dagger a_{\bar{k}}) + \varepsilon_k^\nu (b_k^\dagger b_k + b_{\bar{k}}^\dagger b_{\bar{k}}) \right].$$

$\varepsilon_k^{\pi(\nu)}$: single-proton (neutron) energies obtained from the Dirac equation,

$$[-i\boldsymbol{\alpha} \cdot \boldsymbol{\nabla} + \beta(m + S) + V]\psi_k = \varepsilon_k \psi_k.$$

$$S(\mathbf{r}) = \alpha_S \rho_S + \beta_S \rho_S^2 + \gamma_S \rho_S^3 + \delta_S \Delta \rho_S,$$

$$V(\mathbf{r}) = \alpha_V \rho_V + \gamma_V \rho_V^3 + \delta_V \Delta \rho_V + \tau_3 \alpha_{TV} \rho_{TV} + \tau_3 \delta_{TV} \Delta \rho_{TV} + e \frac{1 - \tau_3}{2} A^0.$$

□ Pairing part:
$$\hat{H}_{\text{pair}} = \sum_{T_z=0,\pm 1} \hat{H}_{\text{pair}}^{T_z}, \quad \hat{H}_{\text{pair}}^{T_z} = -G \sum_{k,k'>0}^{k \neq k'} P_{k,T_z}^\dagger P_{k',T_z}.$$

The pp , pn and nn pair creation operators are respectively:

$$P_{k,-1}^\dagger = a_k^\dagger a_{\bar{k}}^\dagger, \quad P_{k,0}^\dagger = \frac{1}{\sqrt{2}} (b_k^\dagger a_k^\dagger + a_k^\dagger b_{\bar{k}}^\dagger), \quad P_{k,1}^\dagger = b_k^\dagger b_{\bar{k}}^\dagger.$$

RDFT-SLAP: Many-body Wave function

- Nuclear many-body wave functions:

$$|\Psi\rangle = \sum_{i, \{s_k\}} C_i^{\{s_k\}} |\text{MPC}_i^{\{s_k\}}\rangle.$$

$C_i^{\{s_k\}}$: expansion coefficients

$$|\text{MPC}_i^{\{s_k\}}\rangle = |l_1 l_2 \cdots l_N m_1 m_2 \cdots m_Z\rangle = b_{l_1}^\dagger b_{l_2}^\dagger \cdots b_{l_N}^\dagger a_{m_1}^\dagger a_{m_2}^\dagger \cdots a_{m_Z}^\dagger |0\rangle$$

- Occupation probabilities for the single-particle states:

$$n_k^{\pi(\nu)} = \sum_{i, \{s_k\}} |C_i^{\{s_k\}}|^2 P_i^{k, \pi(\nu)}, \quad P_i^{k, \pi(\nu)} = \begin{cases} 1, & \psi_k^{\pi(\nu)} \text{ is occupied in } \text{MPC}_i^{\{s_k\}}, \\ 0, & \text{otherwise.} \end{cases}$$

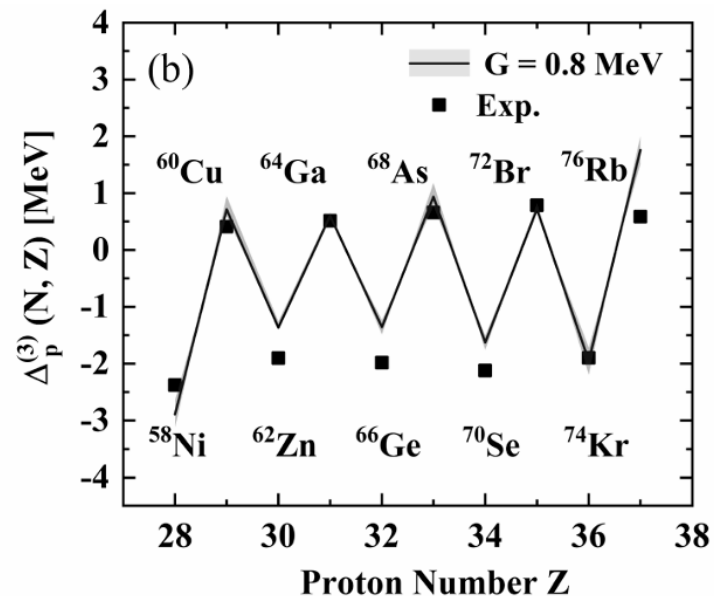
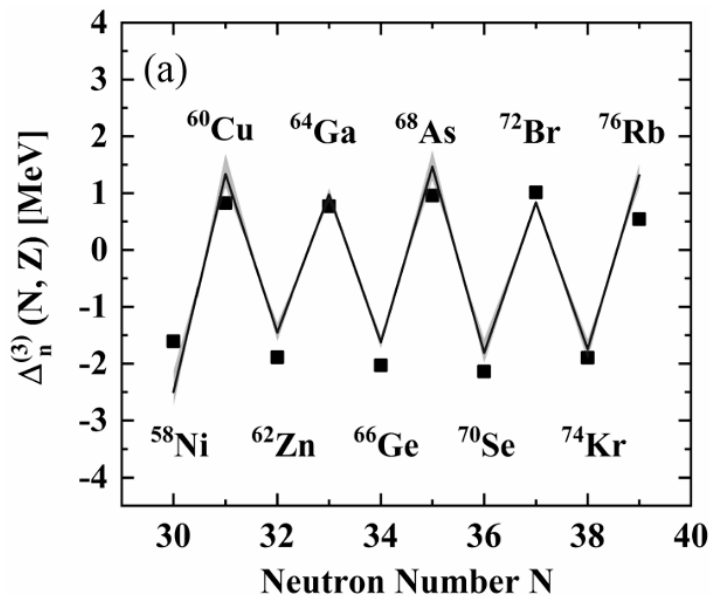
- Nucleon densities calculated by the occupation probabilities:

$$\begin{aligned} \rho_S(\mathbf{r}) &= \sum_k n_k \bar{\psi}_k(\mathbf{r}) \psi_k(\mathbf{r}), & \rho_V(\mathbf{r}) &= \sum_k n_k \psi_k^\dagger(\mathbf{r}) \psi_k(\mathbf{r}), \\ \rho_{TV}(\mathbf{r}) &= \sum_k n_k \psi_k^\dagger(\mathbf{r}) \tau_3 \psi_k(\mathbf{r}), & \rho_c(\mathbf{r}) &= \sum_k n_k \psi_k^\dagger(\mathbf{r}) \frac{1 - \tau_3}{2} \psi_k(\mathbf{r}). \end{aligned}$$

which in turn determines S and V in the Dirac equation.

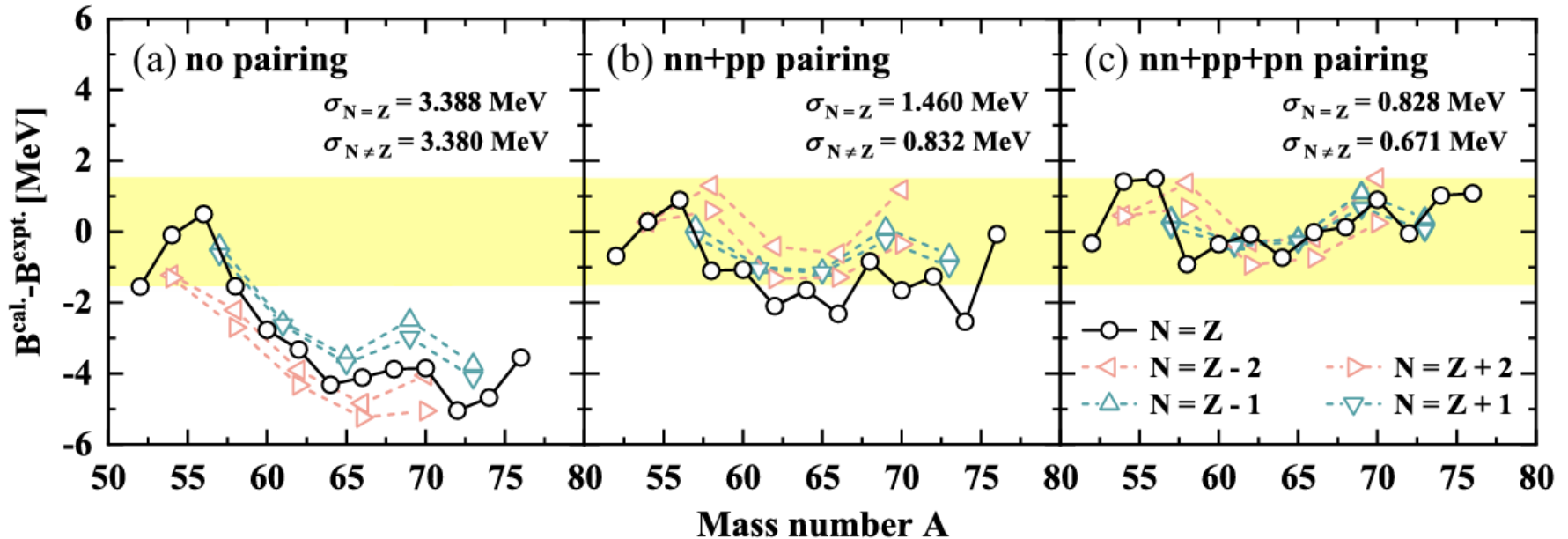
Numerical Details

- Density functional: PC-PK1 [P. W. Zhao et al., PRC 82 \(2010\) 054319](#)
- Major shells of the 3DHO bases: 10
- Energy truncation for MPC space: 16 MeV
- Pairing strength: $G = 0.8$ MeV



Data from M. Wang *et al.*, CPC 45 (2021) 030003 ; F. G. Kondev *et al.*, CPC 45 (2021) 030001

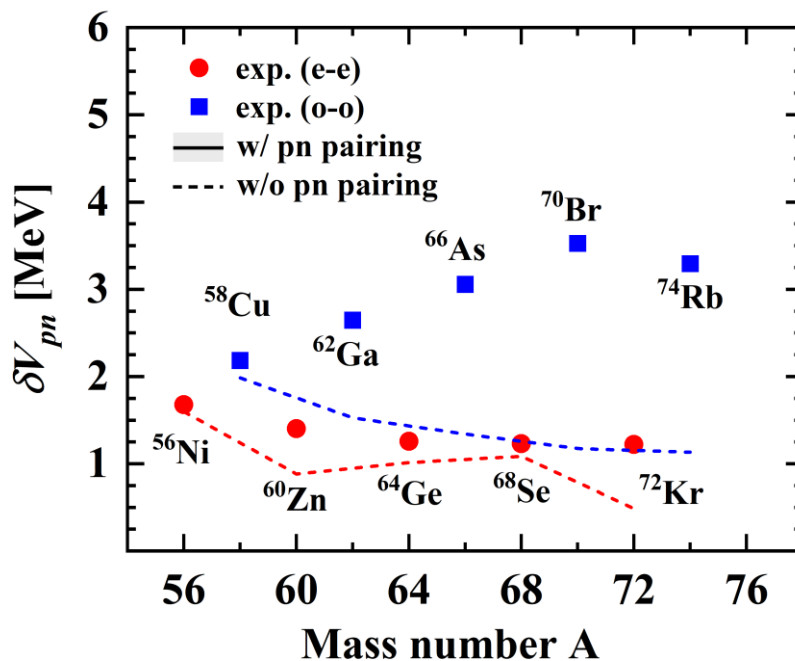
Descriptions of Binding Energies



Data from M. Wang *et al.*, CPC 45 (2021) 030003 ; F. G. Kondev *et al.*, CPC 45 (2021) 030001
M. Wang *et al.*, PRL 130 (2023) 192501

The pairing correlations are important for these nuclei near the $N = Z$ line, in particular the pn pairing correlations.

Interpretation of δV_{pn} bifurcation



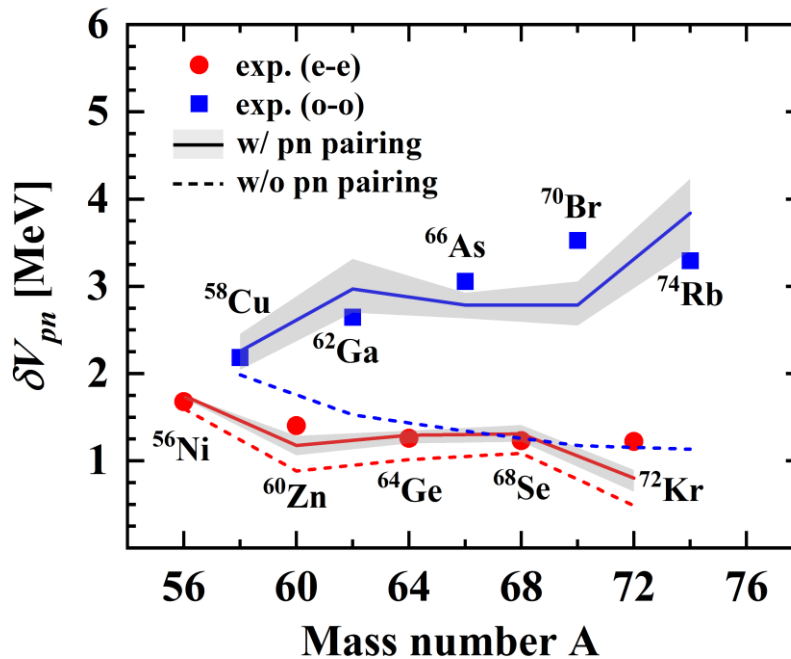
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Van Isacker et al., PRL 74 (1995) 4607

The pn pairing correlations would significantly enhance the δV_{pn} for odd-odd $N = Z$ nuclei, and thus result in the δV_{pn} bifurcation.

Interpretation of δV_{pn} bifurcation



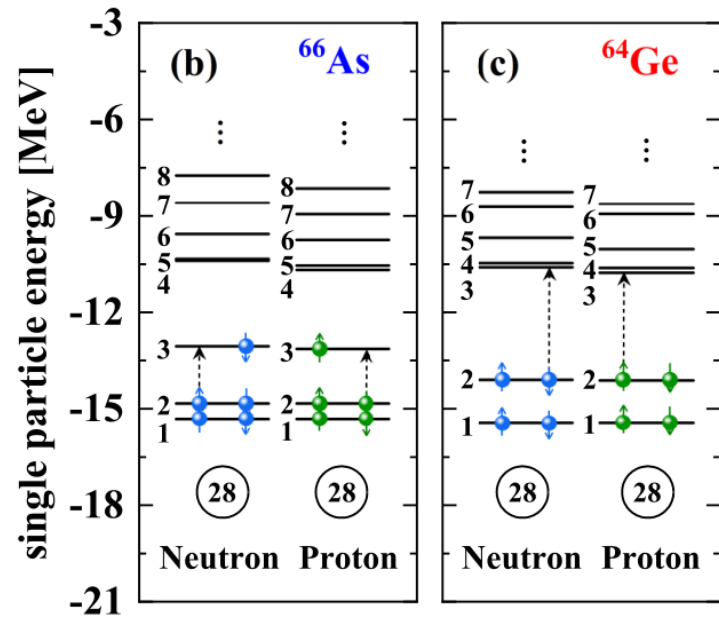
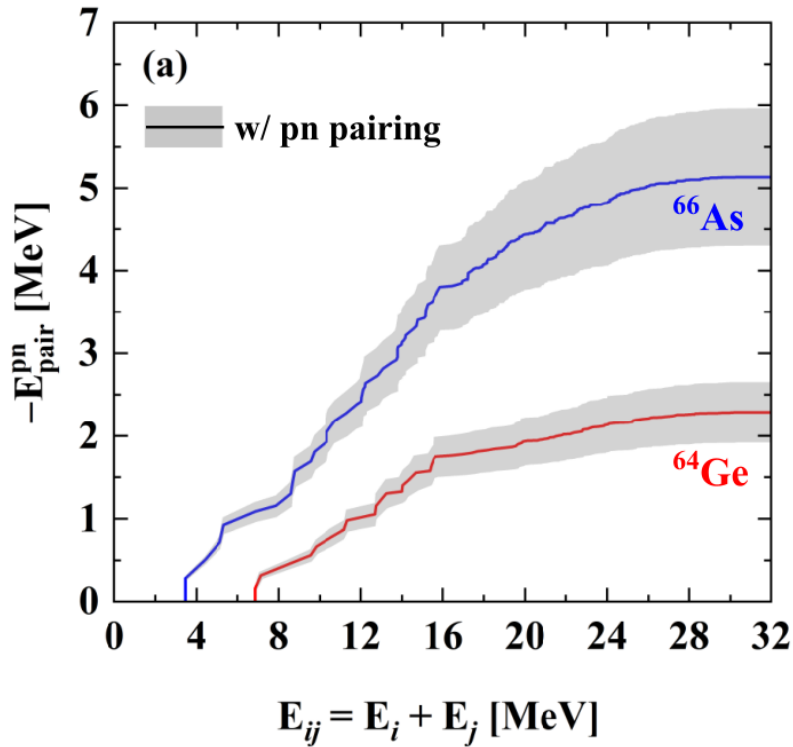
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How to understand the pn pairing effects on δV_{pn}



$$E_{\text{pair}}^{\text{pn}} = \langle \Psi | \hat{H}_{\text{pair}}^{T_z=0} | \Psi \rangle = \sum_{ij} C_i^* C_j \langle \text{MPC}_i | \hat{H}_{\text{pair}}^{T_z=0} | \text{MPC}_j \rangle$$

Compared with ^{64}Ge , there are more MPCs with low excitation energies contributing nonzero pn pairing energy $E_{\text{pair}}^{\text{pn}}$ for ^{66}As .

Summary and Outlook

Summary

RDFT + SLAP calculations with nn , pp and pn pairing correlations:

- provide a good description of the masses for nuclei near $N = Z$ line
- provide an excellent interpretation of the δV_{pn} bifurcation
- reveal a clear signal for the existence of pn pairing correlations

Outlook

The newly developed method and its extensions can also be used for:

- δV_{pn} for nuclei in other mass regions
- Rotational properties for $N = Z$ nuclei
- Possible signals of isoscalar pairing correlations

Thank you for your attention!