

*Beyond-mean-field approach to shape coexistence phenomena  
in neutron-rich nuclei*

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## *Outline*

- *complex EXCITED VAMPIR beyond-mean-field variational model*
- *shape coexistence and first-forbidden  $\beta^-$  decay of  $^{92}\text{Rb}$  (N=55) to  $^{92}\text{Sr}$  (N=54)*
  - *first-forbidden  $\beta^-$  decay of  $0^-_{gs}$  of  $^{92}\text{Rb}$  to  $0^+_{gs}$  and  $2^+$  states in  $^{92}\text{Sr}$*
- *shape coexistence and first-forbidden  $\beta^-$  decay of  $^{98}\text{Y}$  (N=59) to  $^{98}\text{Zr}$  (N=58)*
  - *low-lying isomeric states in  $^{98}\text{Y}$*
  - *exotic structure and dynamics of low-spin states in  $^{98}\text{Zr}$*
  - *first-forbidden  $\beta^-$  decay of  $0^-_{gs}$  of  $^{98}\text{Y}$  to  $0^+$  and  $2^+$  states in  $^{98}\text{Zr}$*

## *Neutron-rich $A \sim 100$ nuclei manifest dramatic evolution in structure and dynamics :*

- shape transition, shape coexistence, shape mixing*
- drastic changes in structure with particle number, spin, and excitation energy*
- particular variations and even "sudden" changes correlated with the onset of deformation around  $N=58$  in  $Sr$  and  $Zr$  isotopic chains, but a more gradual evolution in the  $Kr$  chain*

## *Challenges for theory*

- beyond-mean-field methods*
- realistic effective Hamiltonians in adequate large model spaces  
aiming at*
- unitary description of structure and  $\beta$ -decay properties*
- comprehensive understanding of multifaceted impact of shape coexistence and mixing*

## *complex EXCITED VAMPIR model*

- *the model space is defined by a finite dimensional set of spherical single particle states*
- *the effective many-body Hamiltonian is represented as a sum of one- and two-body terms*
- *the basic building blocks are Hartree-Fock-Bogoliubov (HFB) vacua*
- *the HFB transformations are essentially complex and allow for proton-neutron, parity, and angular momentum mixing being restricted by time-reversal and axial symmetry*  
*(include unnatural-parity two-body correlations and  $T=1$  and  $T=0$  neutron-proton pairing correlations already at the mean-field level )*
- *the broken symmetries ( $s=N, Z, I, p$ ) are restored by projection before variation*
- \* *The model allows to use rather large model spaces and realistic effective interactions*

# complex EXCITED VAMPIR model - beyond-mean-field variational approach

## VAMPIR

$$E^s[F_1^s] = \frac{\langle F_1^s | \hat{H} \hat{\Theta}_{00}^s | F_1^s \rangle}{\langle F_1^s | \hat{\Theta}_{00}^s | F_1^s \rangle}$$

$\Theta_{00}^s$  - symmetry projector

$|F_1^s\rangle$  - HFB vacuum

$$|\psi(F_1^s); sM\rangle = \frac{\hat{\Theta}_{M0}^s |F_1^s\rangle}{\sqrt{\langle F_1^s | \hat{\Theta}_{00}^s | F_1^s \rangle}}$$

## EXCITED VAMPIR

$$|\psi(F_i^s); sM\rangle = \sum_{j=1}^i |\phi(F_j^s)\rangle \alpha_j^i \quad \text{for } i = 1, \dots, n-1$$

$$|\phi(F_i^s); sM\rangle = \hat{\Theta}_{M0}^s |F_i^s\rangle$$

Allows to identify in a small energy interval spherical, oblate, prolate deformed orthogonal configurations of a given symmetry.

$$|\psi(F_n^s); sM\rangle = \sum_{j=1}^{n-1} |\phi(F_j^s)\rangle \alpha_j^n + |\phi(F_n^s)\rangle \alpha_n^n$$

$$(H - E^{(n)} N) f^{(n)} = 0$$

Projected configurations significantly correlated could become strongly mixed by the final diagonalization.

$$(f^{(n)})^+ N f^{(n)} = 1$$

$$|\Psi_\alpha^{(n)}; sM\rangle = \sum_{i=1}^n |\psi_i; sM\rangle f_{i\alpha}^{(n)}, \quad \alpha = 1, \dots, n$$

## *Neutron-rich nuclei in the $A=100$ mass region*

*$^{40}\text{Ca}$  - core*

*model space for both protons and neutrons :*

*$1p_{1/2}$   $1p_{3/2}$   $0f_{5/2}$   $0f_{7/2}$   $2s_{1/2}$   $1d_{3/2}$   $1d_{5/2}$   $0g_{7/2}$   $0g_{9/2}$   $0h_{11/2}$*

## *Effective interaction - renormalized $G$ -matrix (Bonn CD potential)*

- pairing properties enhanced by short range Gaussians for:*

$T = 1$  pp, nn channels

$T = 0$  and 1 np channels

- onset of deformation influenced by monopole shifts:*

$\langle 0g_{9/2} 0f; T=0 | G | 0g_{9/2} 0f; T=0 \rangle$

$\langle 0f_{5/2} 1d; T=0 | G | 0f_{5/2} 1d; T=0 \rangle$

$\langle 0g_{9/2} 0h_{11/2}; T=0 | G | 0g_{9/2} 0h_{11/2}; T=0 \rangle$

- Coulomb interaction between valence protons added*

## First-forbidden beta decay formalism

$$\frac{1}{t_{1/2}} = \frac{f}{K} \quad f = \int_1^{W_0} C(W)W(W^2 - 1)^{1/2}(W_0 - W)^2 F(Z, W) dW,$$

*first-forbidden*  $0^- \rightarrow 0^+$   $C(W) = k + kb/W$

$$O(0^-) = \sum_{ab} o^{(0)}(0^-)(ab) \langle \xi_f J_f || [c_a^\dagger \tilde{c}_b]_0 || \xi_i J_i \rangle$$

$$w = -g_A \sqrt{3} \langle a || r [\mathbf{C}_1 \times \boldsymbol{\sigma}]^{(0)} || b \rangle$$

$$w' = -g_A \sqrt{3} \langle a || \frac{2}{3} r I(1, 1, 1, 1, r) [\mathbf{C}_1 \times \boldsymbol{\sigma}]^{(0)} || b \rangle$$

$$v = \frac{\epsilon_{mec} g_A \sqrt{3}}{M} \langle a || [\boldsymbol{\sigma} \times \boldsymbol{\nabla}]^{(0)} || b \rangle$$

*first-forbidden*  $0^- \rightarrow 2^+$   $C(W) = k + kaW + kcW^2$

$$O(0^-) = \sum_{ab} o^{(2)}(0^-)(ab) \langle \xi_f J_f || [c_a^\dagger \tilde{c}_b]_2 || \xi_i J_i \rangle$$

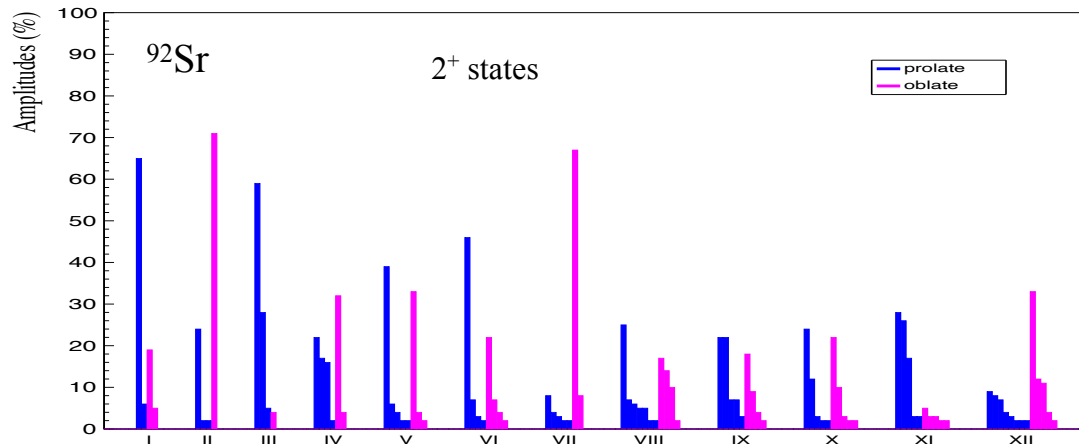
$$z = 2g_A \langle a || r [\mathbf{C}_1 \times \boldsymbol{\sigma}]^2 || b \rangle$$

# Shape coexistence and first-forbidden $\beta^-$ decay of $^{92}\text{Rb}$ to $^{92}\text{Sr}$

A. Petrovici, Phys. Rev. C 109, 024303 (2024)

$^{92}\text{Rb}$  :  $0^-_{\text{gs}}$  - spherical configuration  
(spherical occupation  $(d^{\nu}_{5/2})^{4.81}$ )

$^{92}\text{Sr}$  :  $0^+_{\text{gs}}$  - multiple shape coexistence : spherical (57%),  
prolate (28%), oblate (15%) configurations

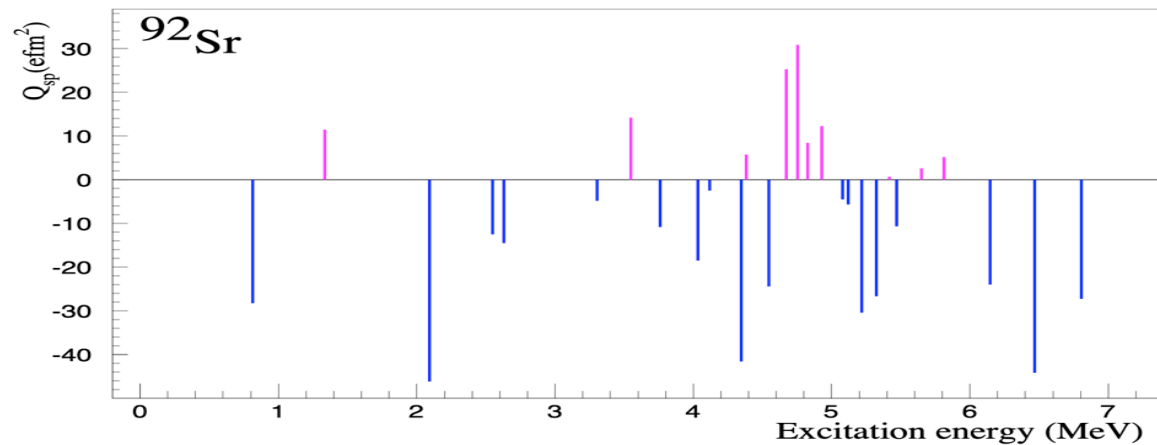


$$\beta_2 = 0.06 \div 0.31$$

$$\beta_2 = -0.18 \div -0.25$$

$$B^{\text{EXP}}(E2; 2^+_1 \rightarrow 0^+_{\text{gs}}) = 197(74) e^2 \text{fm}^4$$

$$B^{\text{EXVAM}}(E2; 2^+_1 \rightarrow 0^+_{\text{gs}}) = 277 e^2 \text{fm}^4$$





*First-forbidden  $\beta^-$  decay of the  $0^-$  ground state of  $^{92}\text{Rb}$  to  $0^+$  ground state of  $^{92}\text{Sr}$*

$$v = \frac{\epsilon_{mec} g_A \sqrt{3}}{M} \langle a || [\boldsymbol{\sigma} \times \boldsymbol{\nabla}]^{(0)} || b \rangle$$

*Log ft values for the decay of  $0^-_{gs}$  of  $^{92}\text{Rb}$  to  $0^+_{gs}$  of  $^{92}\text{Sr}$*

$\epsilon_{mec} = 1.3$	$\epsilon_{mec} = 1.4$	$\epsilon_{mec} = 1.5$	Expt.
5.92	5.78	5.66	5.75

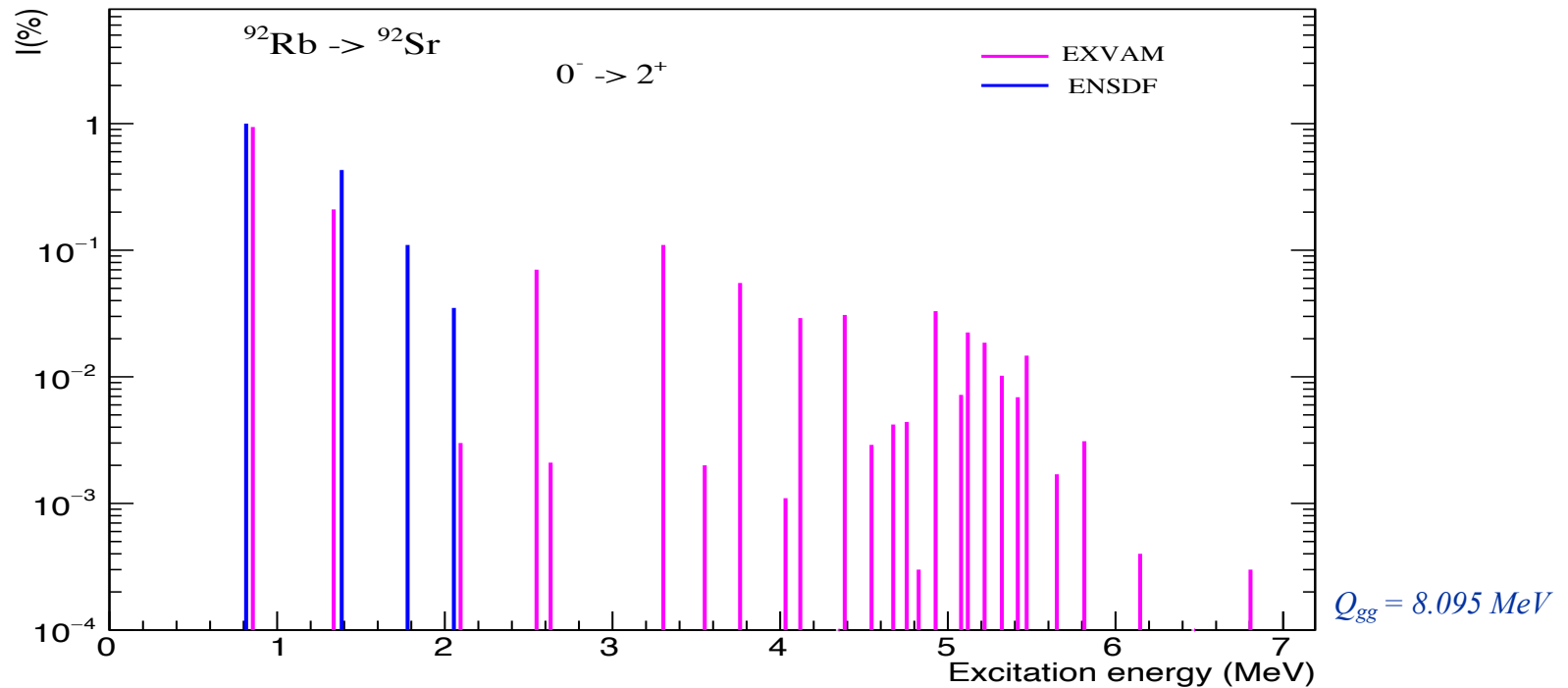
*Contributions: - main :  $d^v_{5/2} f^{\pi}_{5/2}$  matrix elements*

*- very small effects:  $s^v_{1/2} p^{\pi}_{1/2}$ ,  $d^v_{3/2} p^{\pi}_{3/2}$ ,  $g^v_{7/2} f^{\pi}_{7/2}$  matrix elements*

*Log ft values for the decay of  $0^-_{gs}$  of  $^{92}\text{Rb}$  to  $0^+_{gs}$  of  $^{92}\text{Sr}$  constraining the final state to 100% spherical, prolate or oblate content*

	$\epsilon_{mec} = 1.3$	$\epsilon_{mec} = 1.4$	$\epsilon_{mec} = 1.5$
$0^+_{\text{spherical}}$	5.62	5.49	5.37
$0^+_{\text{prolate}}$	7.57	7.45	7.34
$0^+_{\text{oblate}}$	9.15	9.06	9.02

# First-forbidden $\beta^-$ decay of $0^-$ ground state of $^{92}\text{Rb}$ to $2^+$ states in $^{92}\text{Sr}$

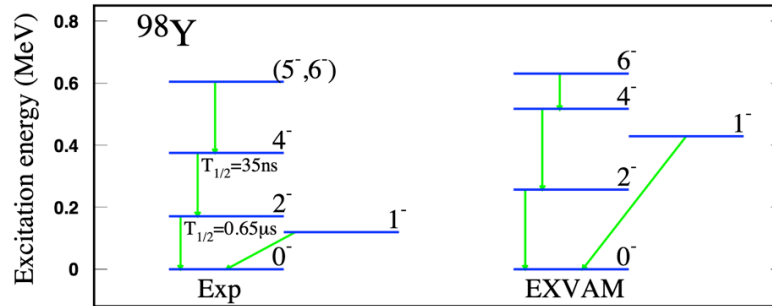


*Significant first-forbidden unique transitions in the high-energy region –  
relevant input for the investigation of the reactor antineutrino spectra*

*Contributions: - main :  $d^{\nu}_{5/2}f^{\pi}_{5/2}$  matrix elements  
- small effects:  $d^{\nu}_{5/2}p^{\pi}_{1/2}$ ,  $d^{\nu}_{5/2}p^{\pi}_{3/2}$ ,  $g^{\nu}_{7/2}f^{\pi}_{5/2}$  matrix elements*

# Multiple shape coexistence in $^{98}\text{Zr}$ by first-forbidden $\beta^-$ decay of $^{98}\text{Y}$

## Low-lying isomeric states in $^{98}\text{Y}$



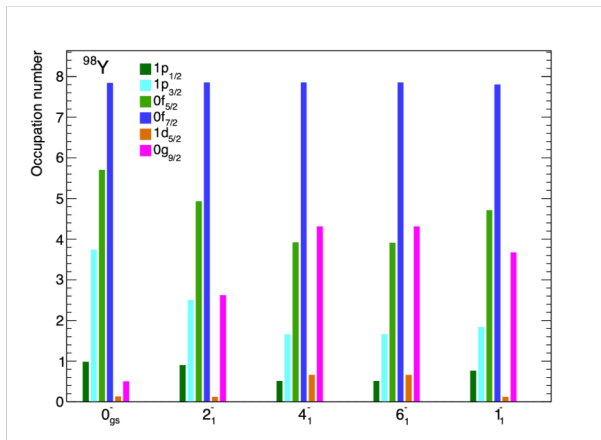
$0^-_{gs}$  : spherical configuration

$2^-_1$  : mixing of oblate deformed configurations

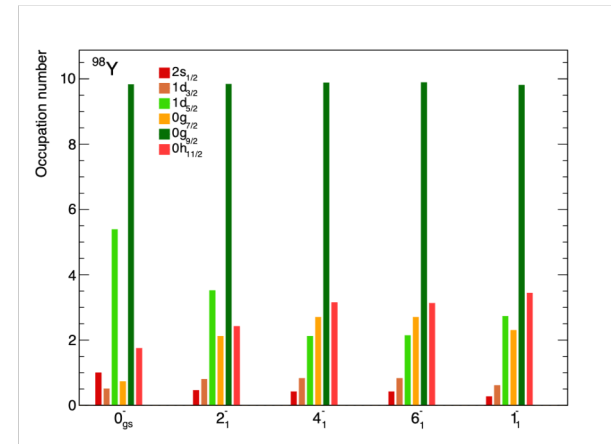
$4^-_1$  : prolate deformed configuration

$6^-_1$  : prolate deformed configuration

$1^-_1$  : mixing of oblate deformed configurations



Proton occupation of valence spherical orbitals

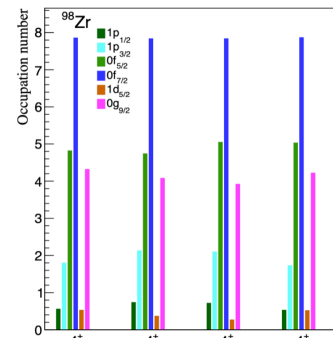
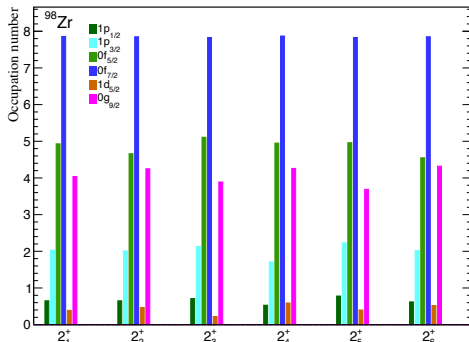
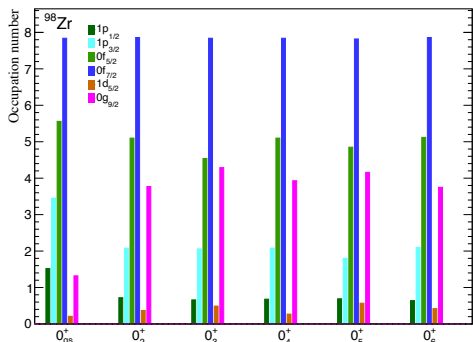


Neutron occupation of valence spherical orbitals

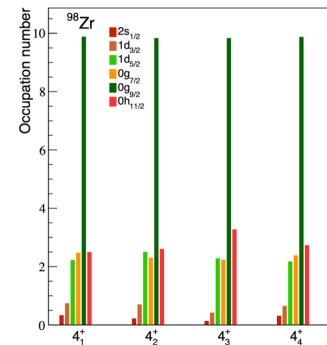
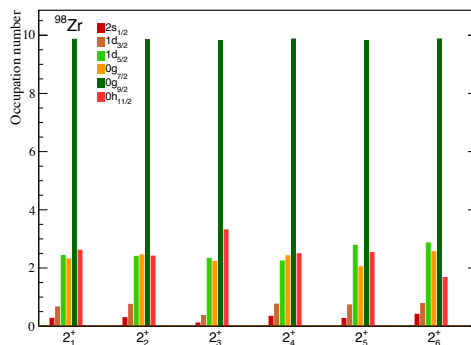
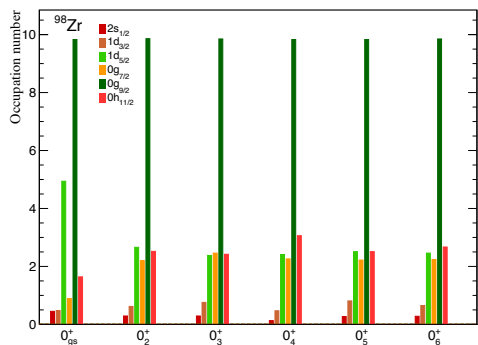


## Occupation of valence spherical orbitals

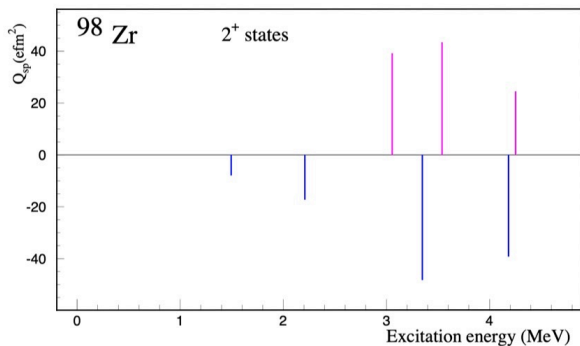
### Protons



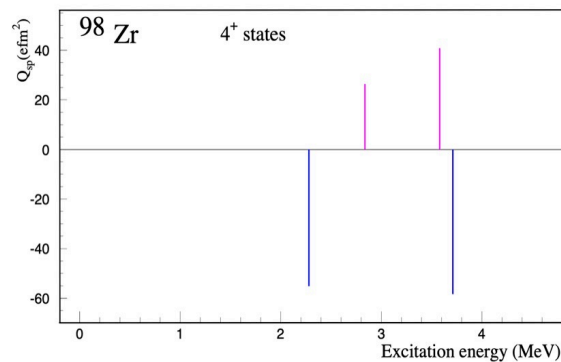
### Neutrons



## Spectroscopic quadrupole moments



$$g(2^+_{1}) = 0.48 \quad g(2^+_{2}) = 0.48 \quad g(2^+_{3}) = 0.435$$



$$g(4^+_{1}) = 0.49 \quad g(4^+_{2}) = 0.43 \quad g(4^+_{3}) = 0.33$$

## *E0 transitions*

$\rho^2(E0)$  values for the lowest six  $0^+$  states in  $^{98}\text{Zr}$

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$I[\hbar]$	EXVAM $0_2^+$	EXP $0_2^+$	EXVAM $0_3^+$	EXP $0_3^+$	EXVAM $0_4^+$	EXP $0_4^+$	EXVAM $0_5^+$	EXVAM $0_6^+$
$0_1^+$	0.033	0.0112(12)						
$0_2^+$			0.020	0.076(6)			0.005	0.007
$0_3^+$					0.022	0.061(8)	0.013	0.011
$0_4^+$								0.006

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$\rho^2(E0)$  values for  $2^+$  and  $4^+$  states in  $^{98}\text{Zr}$

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Transition	EXVAM	Transition	EXVAM
$2_2^+ \rightarrow 2_1^+$	0.013	$4_2^+ \rightarrow 4_1^+$	0.009
$2_3^+ \rightarrow 2_1^+$	0.004	$4_4^+ \rightarrow 4_1^+$	0.016
$2_4^+ \rightarrow 2_1^+$	0.012	$4_3^+ \rightarrow 4_2^+$	0.008
$2_5^+ \rightarrow 2_1^+$	0.006		
$2_3^+ \rightarrow 2_2^+$	0.011		
$2_4^+ \rightarrow 2_2^+$	0.006		
$2_5^+ \rightarrow 2_2^+$	0.020		
$2_5^+ \rightarrow 2_3^+$	0.005		
$2_7^+ \rightarrow 2_5^+$	0.012		

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## *E2 transition strengths*

$B(E2; \Delta I = 2)$  (W.u.) in  $^{98}\text{Zr}$

Transition	EXVAM	Cologne(2020)	ENSDF
$2_1^+ \rightarrow 0_1^+$	3.5	1.1 $+0.3-0.2$	2.9 $+8-5$
$2_1^+ \rightarrow 0_2^+$	28.1	11 $+3-2$	29 $+8-6$
$0_3^+ \rightarrow 2_1^+$	3.7		58 8
$0_4^+ \rightarrow 2_1^+$	0.45		0.103 8
$2_2^+ \rightarrow 0_1^+$	0.04	0.26 $+0.20-0.08$	
$2_2^+ \rightarrow 0_2^+$	1.24	1.8 $+1.4-0.6$	
$2_2^+ \rightarrow 0_3^+$	34.6	40 $18^{\text{Triumf}}$	
$0_4^+ \rightarrow 2_2^+$	0.75		42 3
$2_3^+ \rightarrow 0_1^+$	0.13	0.14 $+0.12-0.04$	
$2_3^+ \rightarrow 0_2^+$	0.25	1.7 $+1.5-0.5$	
$2_3^+ \rightarrow 0_3^+$	0.25		
$2_3^+ \rightarrow 0_4^+$	23.8		
$2_4^+ \rightarrow 0_3^+$	2.3		
$2_4^+ \rightarrow 0_4^+$	0.64		
$2_4^+ \rightarrow 0_5^+$	24.9		
$0_6^+ \rightarrow 2_4^+$	53.5		
$2_5^+ \rightarrow 0_5^+$	6.9		
$2_5^+ \rightarrow 0_6^+$	18.7		
$4_1^+ \rightarrow 2_1^+$	43.6	25 $+15-7$	42 $+10-7$
$4_1^+ \rightarrow 2_2^+$	10.6	38 $+26-13$	54 $+18-16$
$4_2^+ \rightarrow 2_1^+$	1.56	0.6 $+0.17-0.12$	
$4_2^+ \rightarrow 2_2^+$	40.4	4.6 $+1.7-1.3$	
$4_3^+ \rightarrow 2_3^+$	30.7		
$4_3^+ \rightarrow 2_4^+$	4.7		
$4_4^+ \rightarrow 2_4^+$	47.7		
$4_4^+ \rightarrow 2_5^+$	3.9		

$B(E2; \Delta I = 0)$  (W.u.) in  $^{98}\text{Zr}$

Transition	EXVAM	Cologne(2020)	Triumf(2023)
$2_2^+ \rightarrow 2_1^+$	41.6	46 $+35-14$	33 17
$2_3^+ \rightarrow 2_1^+$	0.3	7.6 $+6.5-2.3$	
$2_3^+ \rightarrow 2_2^+$	0.5		
$2_4^+ \rightarrow 2_1^+$	2.3		
$2_4^+ \rightarrow 2_2^+$	5.9		
$2_4^+ \rightarrow 2_3^+$	4.6		
$2_5^+ \rightarrow 2_3^+$	0.4		
$2_5^+ \rightarrow 2_4^+$	13		
$2_6^+ \rightarrow 2_3^+$	4.6		
$4_2^+ \rightarrow 4_1^+$	28.2		
$4_3^+ \rightarrow 4_2^+$	1.2		
$4_4^+ \rightarrow 4_2^+$	5.5		
$4_4^+ \rightarrow 4_3^+$	16.8		

## First-forbidden $\beta^-$ decay of $0^-$ ground state of $^{98}\text{Y}$ to $0^+$ states in $^{98}\text{Zr}$

$^{98}\text{Y}$ :  $0^-_{gs}$  - spherical configuration

$$0^-_{gs}: (s^v_{1/2})^1 (d^v_{5/2})^{5.39}$$

$^{98}\text{Zr}$ :  $0^+_{gs}$ : spherical (87%), prolate (3%), oblate (10%) configurations

$$0^+_{gs}: (d^v_{5/2})^{4.94}$$

$$v = \frac{\epsilon_{mec} g_A \sqrt{3}}{M} \langle a || [\boldsymbol{\sigma} \times \boldsymbol{\nabla}]^{(0)} || b \rangle$$

Log ft values for the decay of  $0^-_{gs}$  of  $^{98}\text{Y}$  to  $0^+$  states in  $^{98}\text{Zr}$

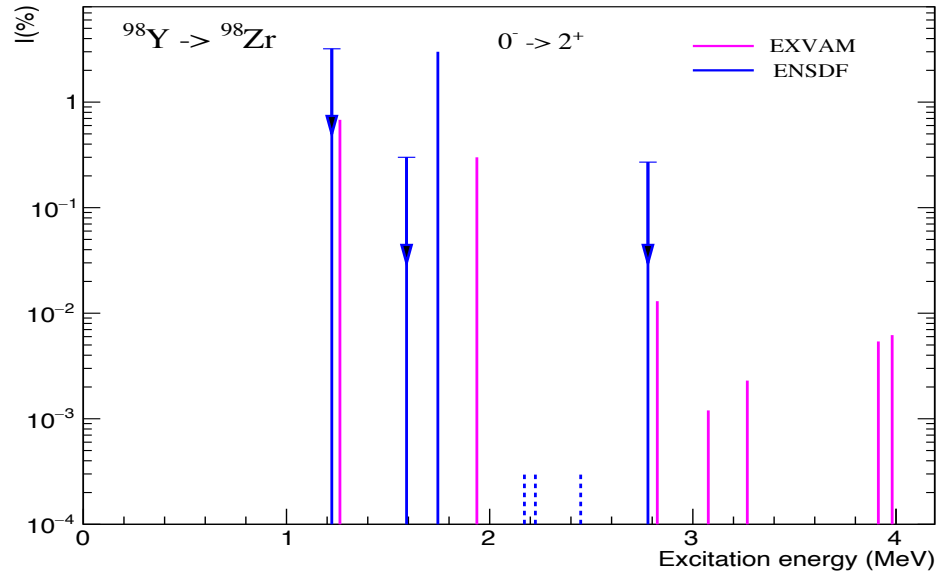
$I[\hbar]$	EXVAM ( $\epsilon_{mec}=1.15$ )	Exp.
$0^+_{gs}$	5.75	5.8 (2)
$0^+_2$	6.56	> 7.2
$0^+_3$	7.53	6.0 (1)
$0^+_4$	7.76	5.6 (1)
$0^+_5$	7.34	
$0^+_6$	7.66	

**Contributions:** - main :  $s^v_{1/2} p^\pi_{1/2}$  matrix elements

- very small (or cancelling effects):  $d^v_{3/2} p^\pi_{3/2}$  and  $d^v_{5/2} f^\pi_{5/2}$  matrix elements



# First-forbidden $\beta^-$ decay of $0^-$ ground state of $^{98}\text{Y}$ to $2^+$ states in $^{98}\text{Zr}$



$$Q_{gg} = 8.992 \text{ MeV}$$

Contributions:  $s_{1/2}^{\nu} p_{3/2}^{\pi}$ ,  $d_{5/2}^{\nu} p_{1/2}^{\pi}$ , and  $d_{5/2}^{\nu} f_{5/2}^{\pi}$  matrix elements

## Summary

*Comprehensive understanding of multifaceted impact of shape coexistence in neutron-rich nuclei within complex Excited Vampir beyond-mean field variational model*

- *Shape coexistence and first-forbidden  $\beta^-$  decay of  $^{92}\text{Rb}$  to  $^{92}\text{Sr}$* 
  - $^{92}\text{Rb}$  - spherical  $0^-$  ground state
  - $^{92}\text{Sr}$  - spherical-prolate-oblate mixing in  $0^+$  ground state
    - complex prolate-oblate mixing in  $2^+$  states ( signature - spectroscopic quadrupole moments )

\* significant first-forbidden unique transitions in high-energy region - relevant for reactor antineutrino anomalies
- *Multiple shape coexistence and first-forbidden  $\beta^-$  decay of  $^{98}\text{Y}$  to  $^{98}\text{Zr}$* 
  - $^{98}\text{Y}$  - spherical  $0^-$  ground state
    - variable mixing of prolate and oblate configurations creates low-spin isomers
  - $^{98}\text{Zr}$  - spherical-prolate-oblate mixing in  $0^+$  states induces significant  $E0$  transitions
    - variable prolate-oblate mixing in  $2^+$  and  $4^+$  states induces complex decay pattern

\* unitary description of exotic structure and first-forbidden  $\beta^-$  decay induced by multiple shape coexistence