

Non-minimal coupling to gravity in the early universe

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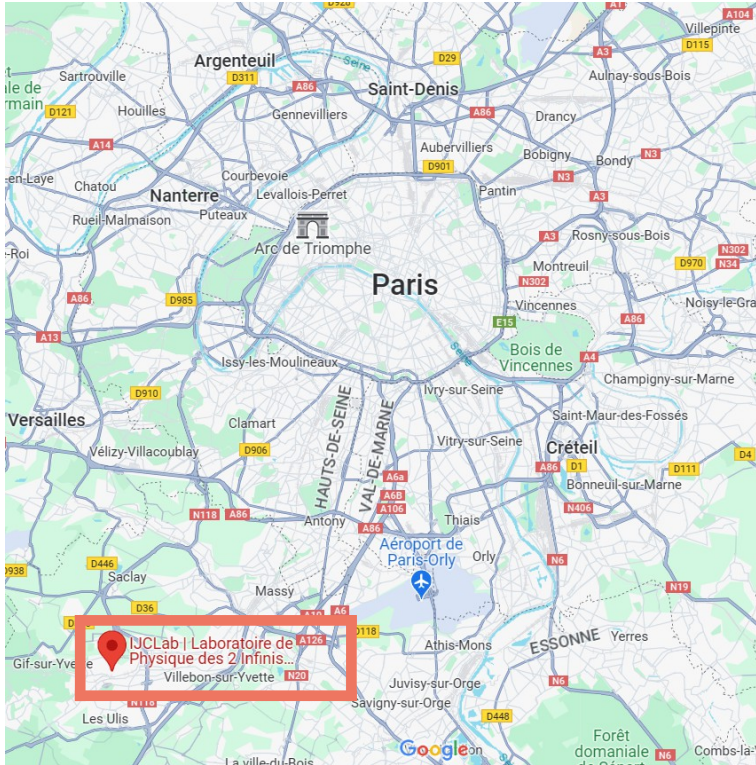
based on “Dark matter production via a non-minimal coupling to gravity” (2211.11773),
and “On unitarity in singlet inflation with a non-minimal coupling to gravity” (2305.05682)

Jong-Hyun Yoon (Jay)
IJCLab CNRS/UPsaclay



RPP, Sorbonne Université
24-26 Jan 2024

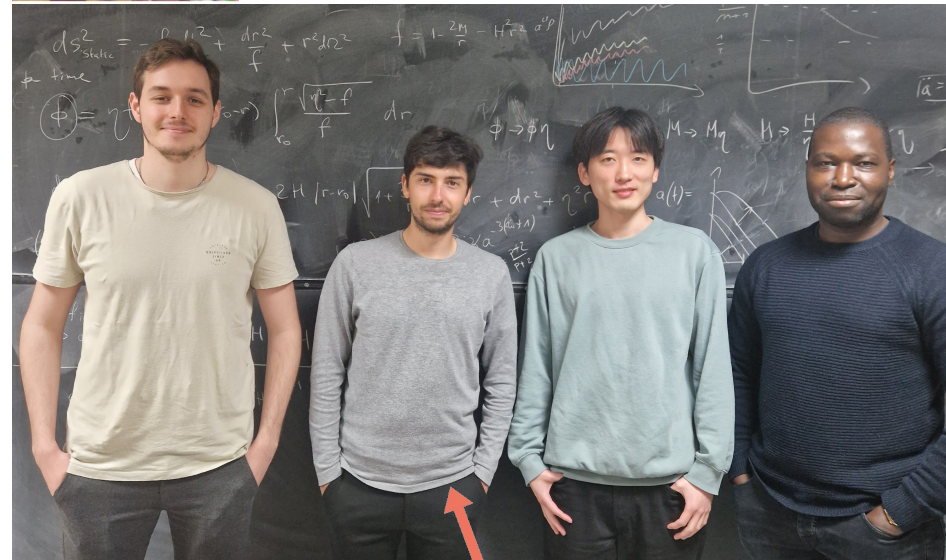
IJCLab @ University Paris-Saclay



Mathieu Gross, “Effects of Fragmentation on Post-Inflationary Reheating” (after Simon)



Yann Mambrini



Essodjolo Kpatcha (Donald)

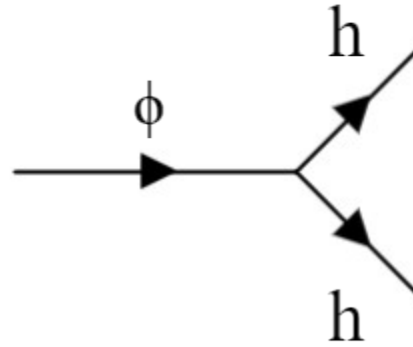
Simon Cléry, “Probing Reheating with Graviton Bremsstrahlung” (right after this talk)

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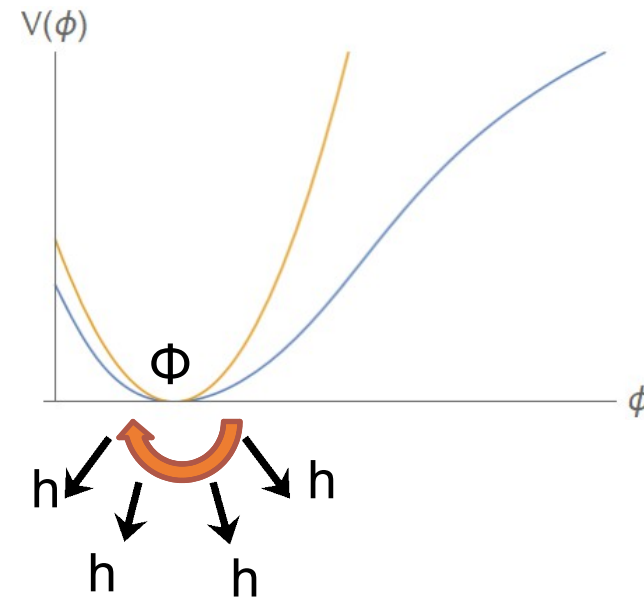
- Reheating and Preheating
- DM production via a Non-Minimal Coupling to gravity
- Singlet inflation with a NMC

Reheating (Inflaton \rightarrow SM bath)

- From Inflaton quanta



- From Inflaton oscillations
 - B.E. enhancement ('preheating')



Minimal Scalar DM models

- While Inflaton \rightarrow SM (reheating the universe),

DM is produced during preheating:

Inflaton=DM

Inflaton-DM scattering

Inflaton F.O., decay to DM

Inflaton-DM non-renormalizable couplings

Inflaton-DM via gravity

Minimal Scalar DM models

- While Inflaton \rightarrow SM (reheating the universe),

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Inflaton F.O., decay to DM

Inflaton-DM non-renormalizable couplings

Inflaton-DM via gravity

Inflaton-DM via gravity

- Non-minimal coupling to gravity

ξ : coefficient
R: Ricci scalar
 Φ : Inflaton field
s: scalar DM

$$\mathcal{S} = \int d^4x \sqrt{-g} \left(\frac{1}{2} M_{\text{Pl}}^2 R - \frac{1}{2} \xi R s^2 - \frac{1}{2} g^{\mu\nu} \partial_\mu s \partial_\nu s - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V \right)$$

- R is a function of energy and dominated by Φ , so DM can interact with Φ via

$$R = -\frac{1}{M_{\text{Pl}}^2} T^\mu{}_\mu$$

For $\xi \gg 1$

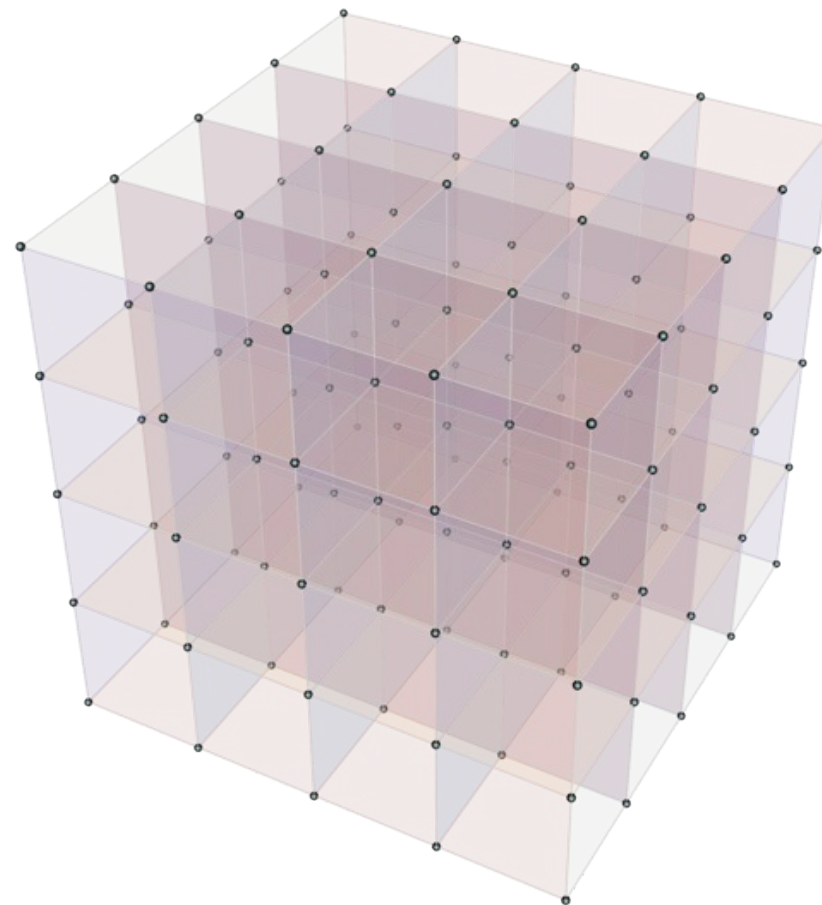
- Previous studies based on perturbative methods (for small ξ)
 - Resonant particle production followed by backreaction and rescattering
 - Non-analytic behavior of curvature
- Numerical approach needed

Lattice simulations

- EOMs + Friedmann Eqs.

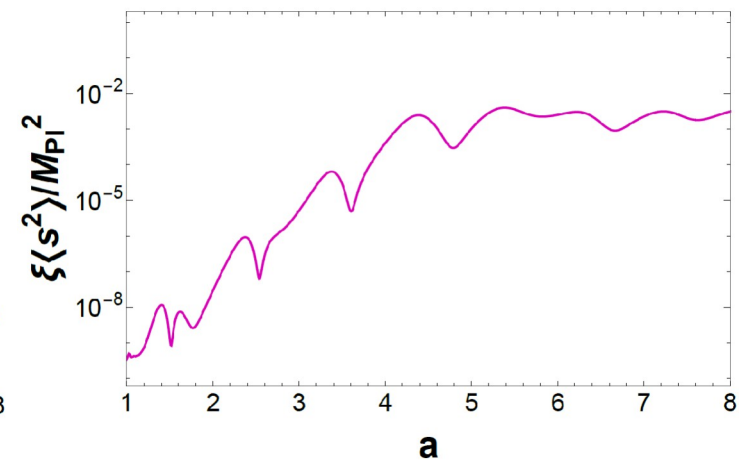
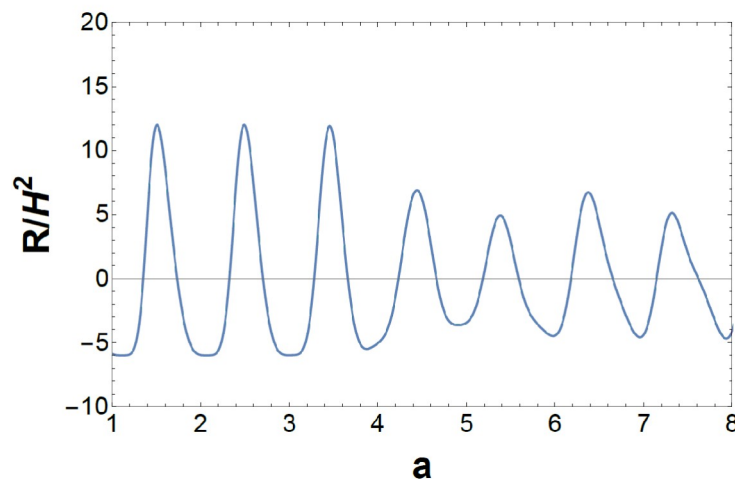
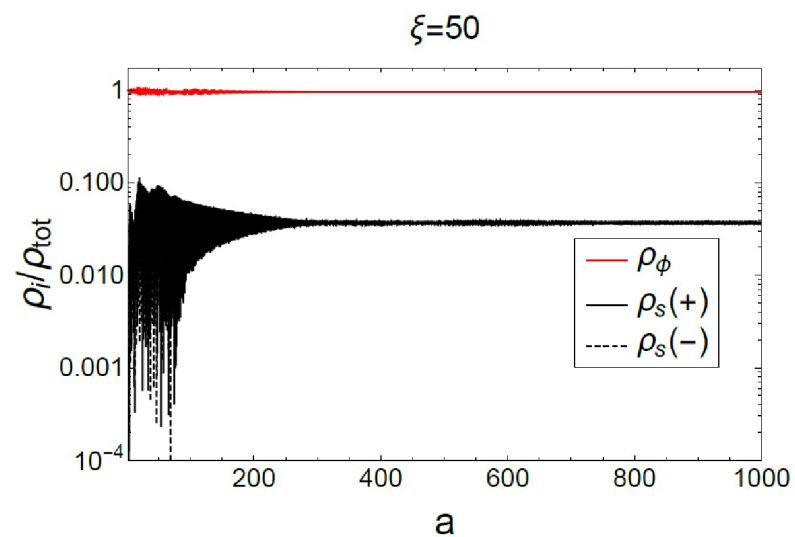
$$\ddot{f} + 3\frac{\dot{a}}{a}\dot{f} - \frac{1}{a^2}\nabla^2 f + \frac{\partial V}{\partial f} = 0$$
$$\ddot{a} = -\frac{4\pi a}{3}(\rho + 3p)$$
$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi}{3}\rho$$

$$\rho = T + G + V ; p = T - \frac{1}{3}G - V$$
$$T = \frac{1}{2}\dot{f}^2 ; G = \frac{1}{2a^2}|\nabla f|^2 .$$



Simulation outcome

- CosmoLattice customized for NMC
- Energy distribution, R breakdown, resonant production, etc.



Simulations provide intuitive insights into events in the early universe

Exp. constraints on DM

- DM relic abundance (conserved since reheating)

$$Y = \frac{n}{s_{\text{SM}}} \quad , \quad s_{\text{SM}} = \frac{2\pi^2}{45} g_{*s} T^3$$

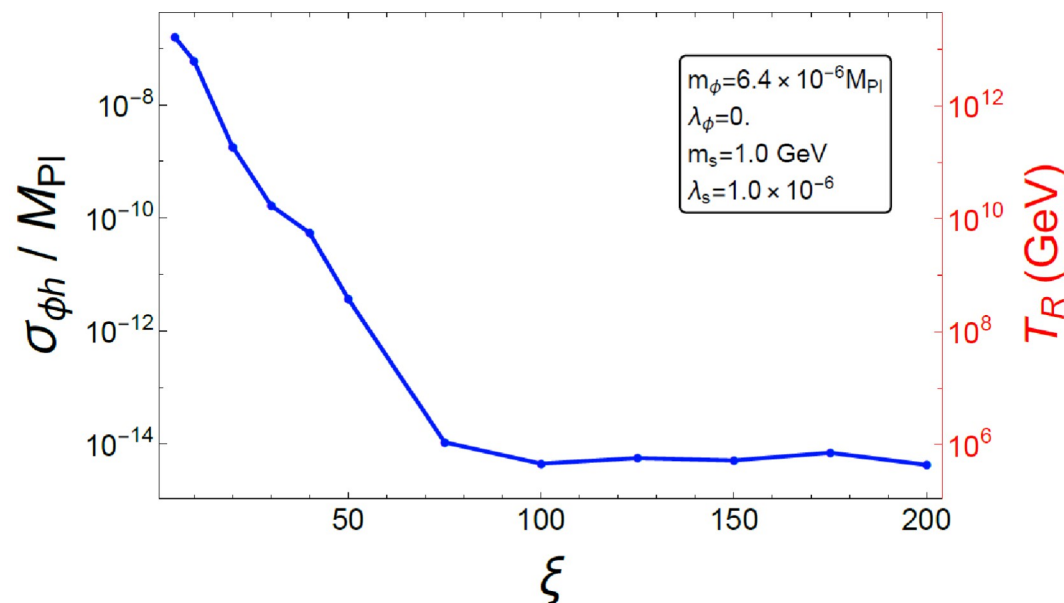
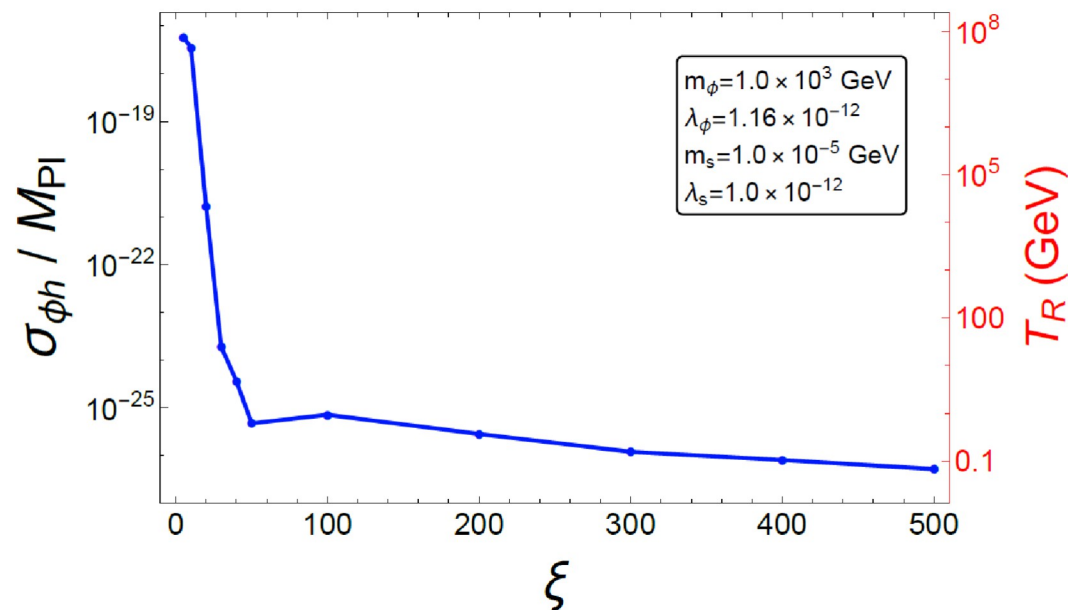
$$Y_{\infty} = 4.4 \times 10^{-10} \left(\frac{\text{GeV}}{m_s} \right)$$

- Reheating via inflaton decay into Higgs $V_{\phi h} = \sigma_{\phi h} \phi H^{\dagger} H$

$$H_R \simeq \Gamma_{\phi \rightarrow hh} \quad , \quad \Gamma_{\phi \rightarrow hh} = \frac{\sigma_{\phi h}^2}{8\pi m_{\phi}} \quad H_R = \sqrt{\frac{\pi^2 g_*}{90}} \frac{T_R^2}{M_{\text{Pl}}}$$

Parameter space (ξ , T_R)

- (locally) quartic and quadratic inflaton potential



- BBN, Cold DM

$$T_R > 4 \text{ MeV}$$

$$\frac{T_{\text{SM}}}{\langle E(s) \rangle} \sim \left(\frac{1}{\lambda_\phi} \frac{a_R}{a_*} \right)^{1/4} \gg 1$$

Singlet inflation w/ NMC

- Singlet scalar Φ with a quartic potential + a non-minimal coupling to gravity

$$\mathcal{S} = \int d^4x \sqrt{-g} \left(\frac{1}{2} M_{\text{Pl}}^2 R + \frac{1}{2} \xi R \phi^2 - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V \right) \quad V = \frac{1}{4} \lambda_\phi \phi^4$$

- Well justified from a phenomenological perspective
- Higgs inflation? Unitarity problem ($E \ll \Lambda$?)

Inflaton dynamics with $\xi (=10000)$

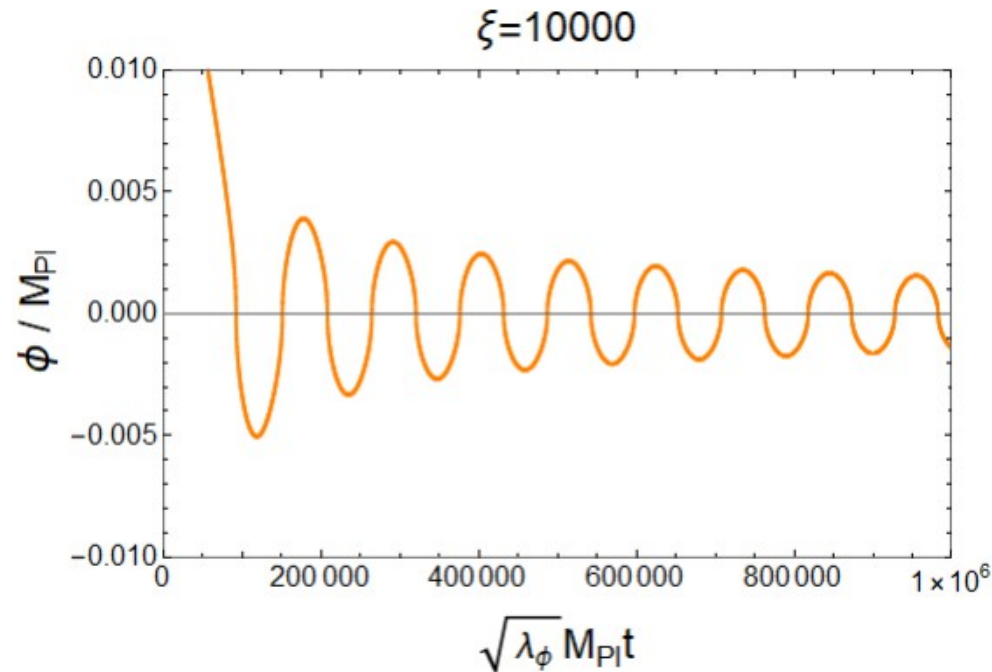
- Solve EOM (linear analysis)

$$\ddot{\phi} + 3H \dot{\phi} - \frac{1}{a^2} \nabla^2 \phi - \xi R \phi + \frac{\partial V}{\partial \phi} = 0 ,$$

$$3H^2 = \rho(\phi) .$$

$$R = \frac{1}{1 + (6\xi + 1)\xi \phi^2} [(1 + 6\xi) \partial^\mu \phi \partial_\mu \phi + 4V + 6\xi \phi V'_\phi]$$

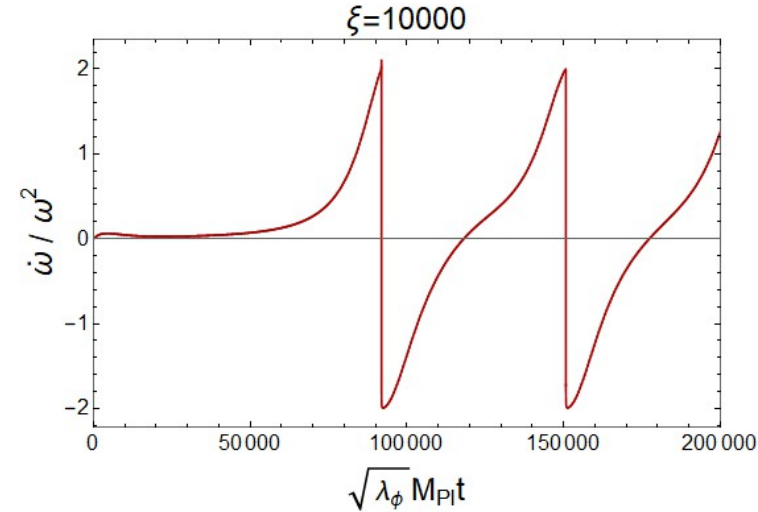
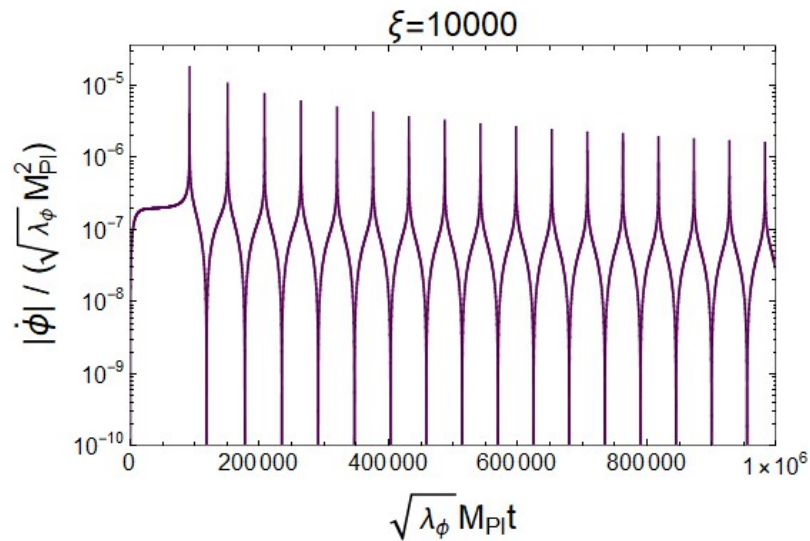
$$\rho(\phi) = \frac{1}{2} \dot{\phi}^2 + V(\phi) - 3\xi H^2 \phi^2 - 6\xi H \phi \dot{\phi} ,$$



NMC makes 'steep' zero-crossings

Inflaton dynamics with $\xi (=10000)$

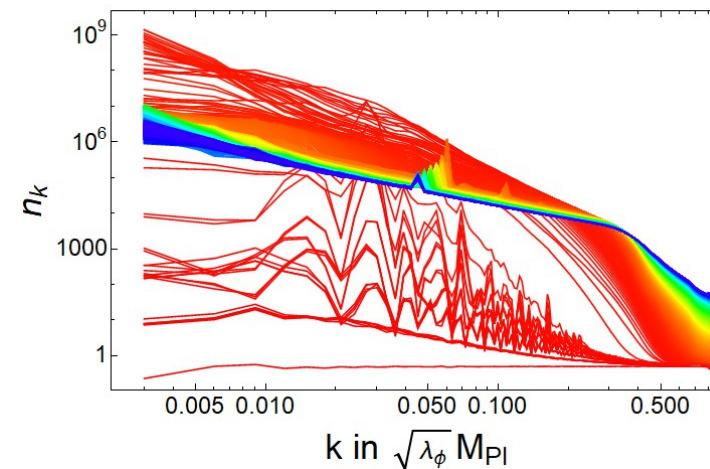
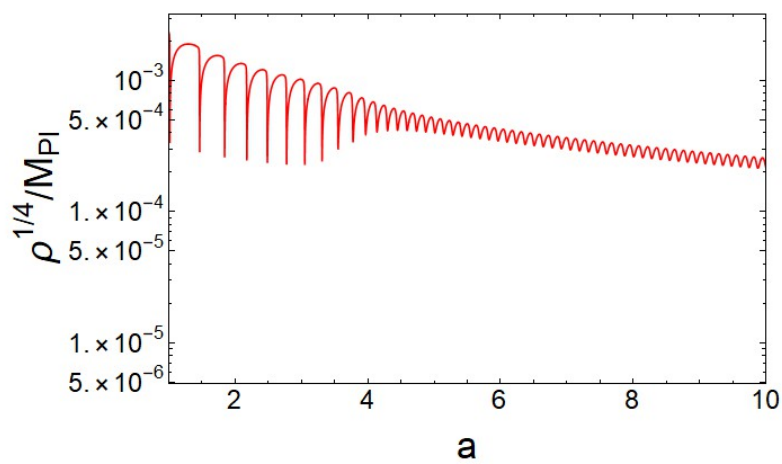
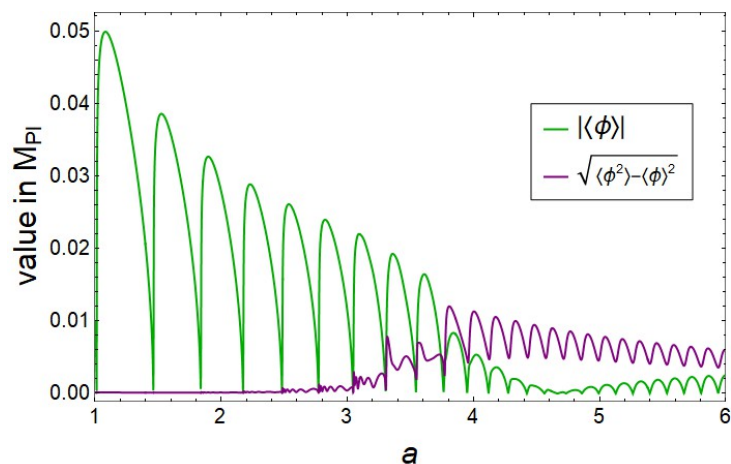
- Non-adiabaticity ($\sim dw/dt \gg w^2$) \rightarrow Particle production



- Efficient particle production can spoil the linear analysis

Simulation outcome for $\xi=100$

- Non-instantaneous but very-early fragmentation
- Non-trivial scaling law in energy ($0 < w < 1/3$)
- Non-perturbative effects (resonance, backreaction, and rescattering)



Collective effects

- Possible Higgs couplings $\Delta V = \sigma_{\phi h} \phi H^\dagger H$

In Einstein frame $\longrightarrow \Delta V / [4(1 + \xi\phi^2)^2]$

χ : canonical
normalized inflaton $\sigma_{\phi h} H^\dagger H \frac{\chi^n}{\Lambda^{n-1}}$

- Unitarity in n-particles states (transition prob. < 1)

$$\rho(n \rightarrow 2) \propto \left(c_n \frac{p_{\max} \sqrt{f}}{\Lambda} \right)^{2n} \quad \Lambda \equiv \frac{1}{\xi}$$

$$\kappa \sim \frac{p_{\max} \sqrt{f}}{\Lambda} \lesssim 1 \quad \longrightarrow \quad \xi_{\max} \sim \text{few} \times 100$$

Conclusion

- Studied DM production during preheating via a NMC (for $\xi \gg 1$)
 - Obtained viable parameter space in terms of ξ and T_R
- Implemented singlet inflation model (w/ NMC) on the lattice and learned about its nontrivial dynamics (for $\xi \gg 1$)
 - Considering unitarity with collective effects within EFT, ξ can be extended to a few 100