



Search for a heavy scalar  $X$  decaying to a scalar  $S$  and a Higgs boson in the  $X \rightarrow SH \rightarrow bb\gamma\gamma$  channel with ATLAS Run-2 data

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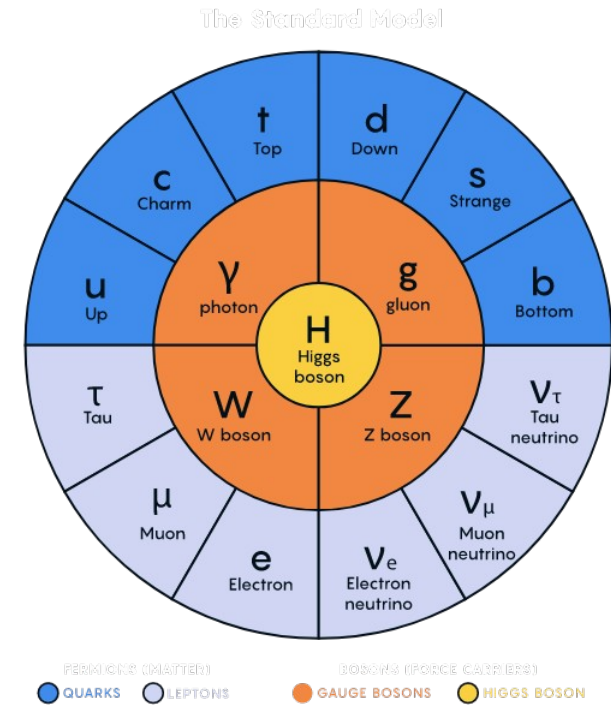
Aix-Marseille University / CNRS / CPPM



# Search for Beyond Standard Model physics

- SM is great but have some caveats :
  - Doesn't explain gravity
  - Doesn't explain neutrino oscillations and mass
  - No dark matter candidate

→ We need physics beyond standard model (BSM)

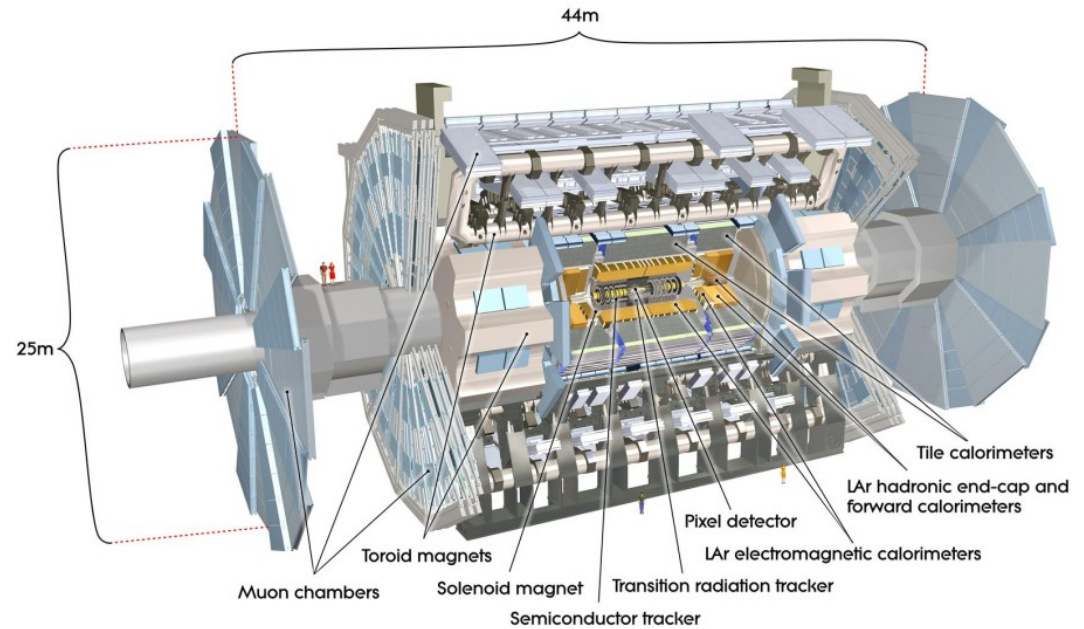
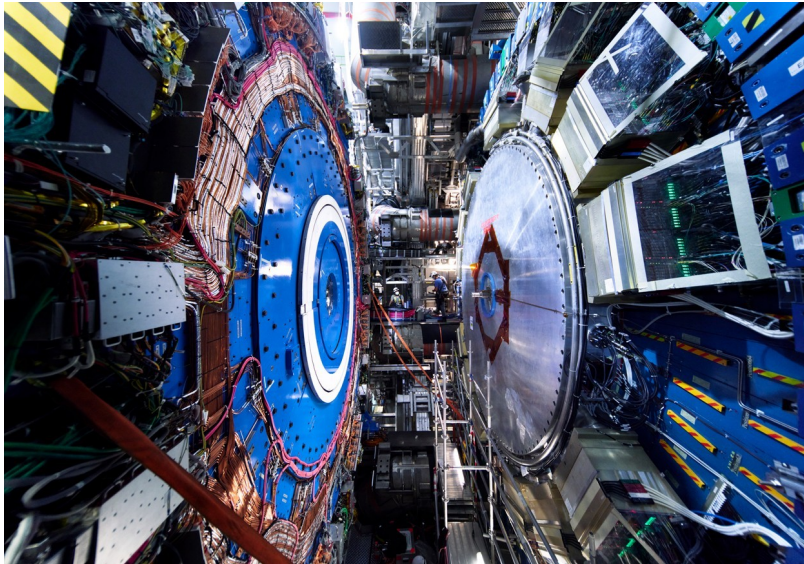


- Some BSM models predict additional scalar particles in different range of masses, notably in the Higgs sector
    - Double Higgs doublet model : 2HDM
    - Next-to-Minimal Supersymmetric Standard Model : NMSSM
    - ... etc
- Could be at reach of LHC and Atlas !



# ATLAS detector

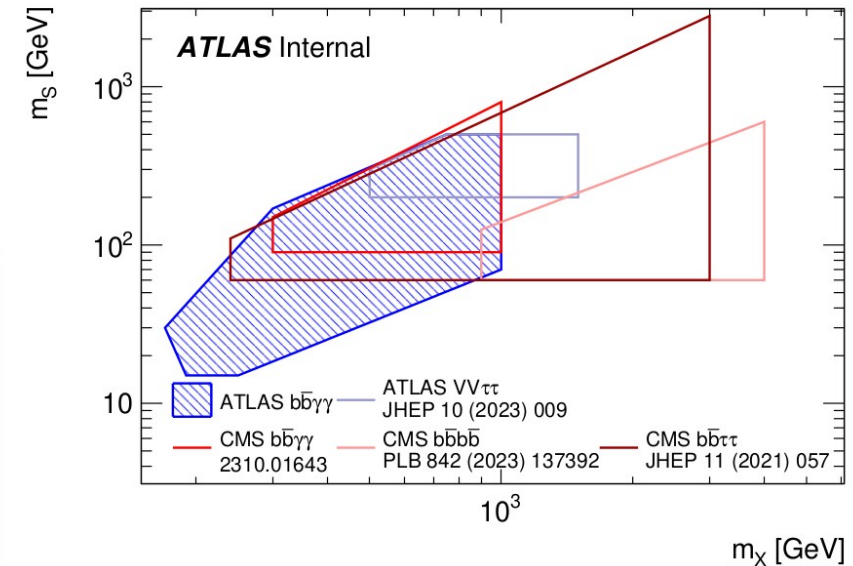
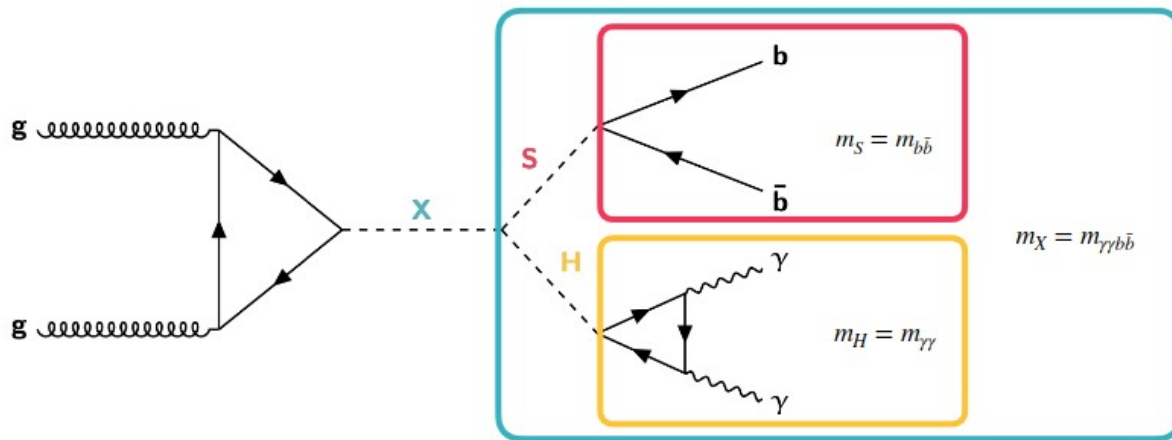
- ATLAS is a multipurposed detector using proton-proton collisions at the Large Hadron Collider (LHC) located at CERN
  - Higgs boson discovered in 2012
  - Make precision measurements to detect deviations from the SM
  - Search for BSM particles, especially in the Higgs sector



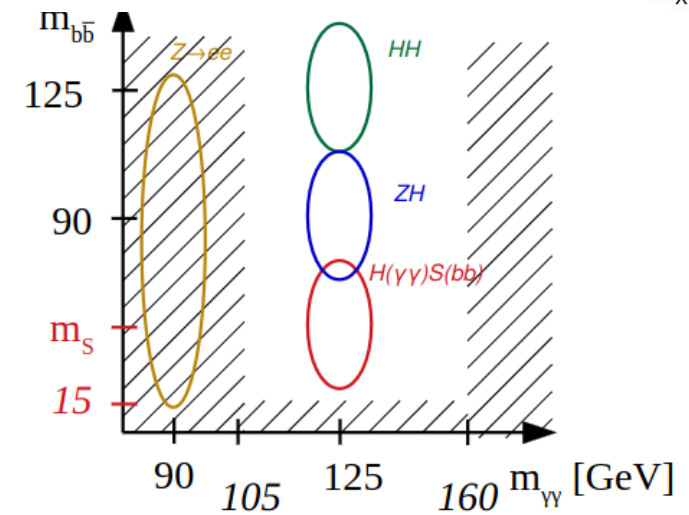


# Targeted signal

- We search for a heavy scalar  $X$  decaying into a light scalar  $S$  and the SM Higgs where  $H \rightarrow \gamma\gamma$  and  $S \rightarrow b\bar{b}$
- Search is **model-independent** and the targeted range of  $X$  and  $S$  masses  $m_X$  and  $m_S$  is :  
 $15 < m_S < 500 \text{ GeV}$  and  $170 \text{ GeV} < m_X < 1 \text{ TeV}$
- Analysis uses  $140 \text{ fb}^{-1}$  Run-2 data at  $\sqrt{s} = 13 \text{ TeV}$



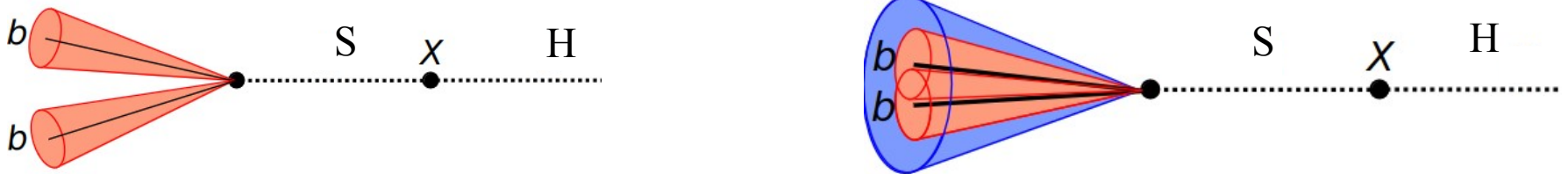
- Analysis is also heavily linked to  $HH \rightarrow b\bar{b}\gamma\gamma$  analysis  
Di-Higgs is a major goal of LHC physics program
  - Can test new strategies with this final state
  - Could also help to remove potential background to  $HH$  signal



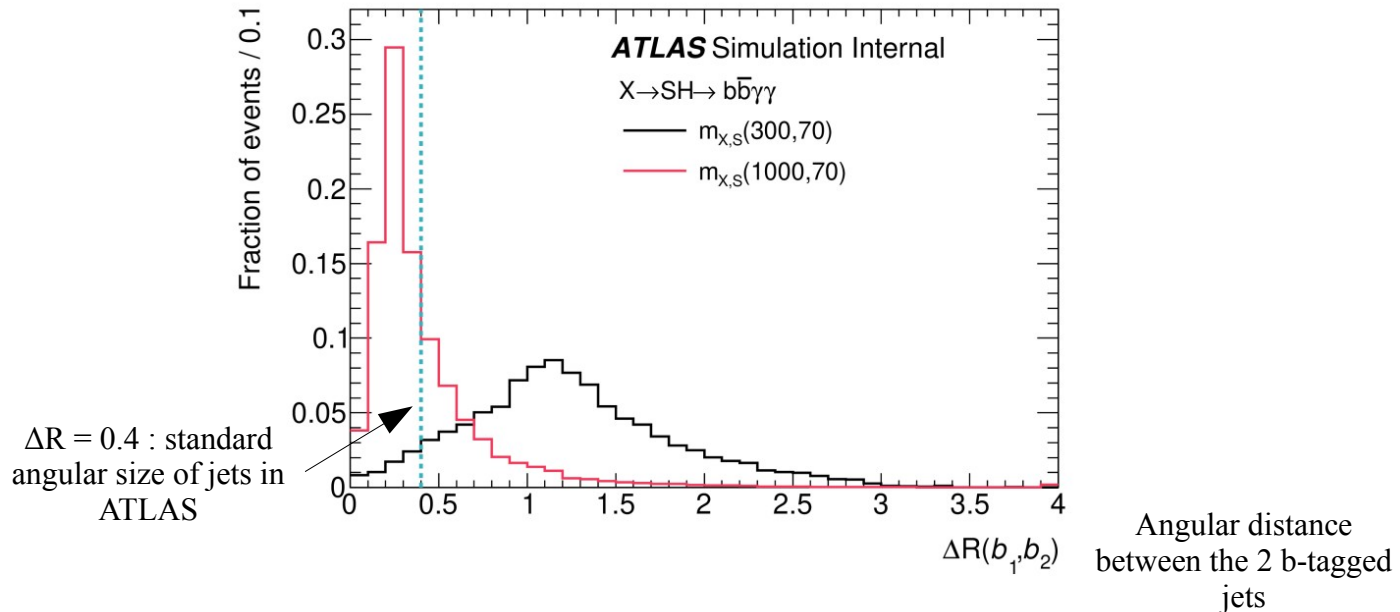


# Different search regions

- A challenging situation arises when  $m_S$  is much smaller than  $m_X$  ( $m_S/m_X < \text{around } 0.1$ ):  
b-jets from the S decay are boosted and reconstructed as one b-tagged jet



→ We separate the search space in a **resolved region** with 2 b-tagged jets and a **merged region** with only one b-tagged jet





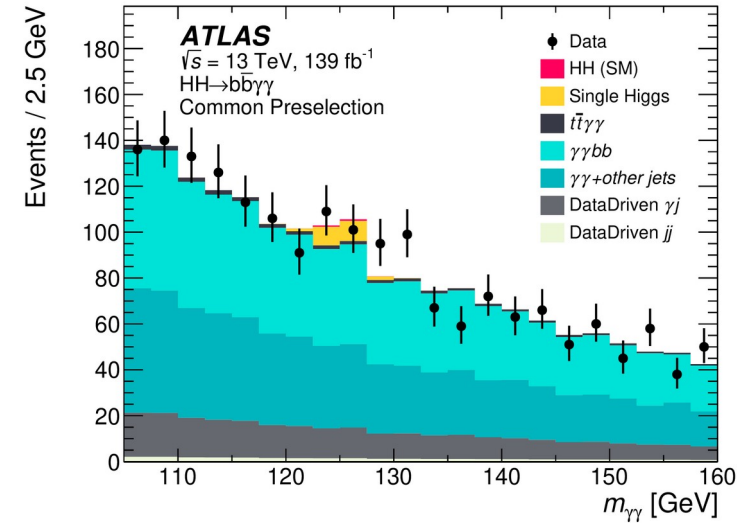
# Event selection

- Selection

Table 4: The definitions of selections used in the analysis.

	2 <i>b</i> -tagged	1 <i>b</i> -tagged
Number of 'tight' and isolated photons		$\geq 2$
$m_{\gamma\gamma}$ [GeV]		$\in [105, 160]$
Number of leptons		$= 0$
Number of central jets		$\in [2, 5]$
Number of <i>b</i> -tagged jets @ 77% WP	$= 2$	$= 1$

## Invariant photon mass distribution



- Predicted background yields (from theoretical cross sections) : main **non-resonant background** is  $\gamma\gamma +$  jets, main **resonant** ones are  $t\bar{t}H$ ,  $ggH$ ,  $ZH$  and also  $VBFH$  for 1 *b*-jet selection

### 2 *b*-tagged jets

	Selection
<i>HH</i> ggF+VBF	$1.691 \pm 0.004$
<i>ZH</i>	$3.691 \pm 0.013$
<i>WH</i>	$0.207 \pm 0.004$
<i>VBFH</i>	$0.685 \pm 0.012$
<i>bbH</i>	$0.62 \pm 0.023$
<i>ggH</i>	$5.453 \pm 0.065$
<i>tHjb</i>	$0.969 \pm 0.029$
<i>tWH</i>	$0.131 \pm 0.005$
<i>ttH</i>	$8.313 \pm 0.014$
$Z(\rightarrow qq)\gamma\gamma$	$20.345 \pm 0.303$
$t\bar{t}\gamma\gamma$	$28.69 \pm 0.109$
$\gamma\gamma$ +jets	$1418.32 \pm 4.596$
Total SM	$1489.116 \pm 4.608$

### 1 *b*-tagged jet

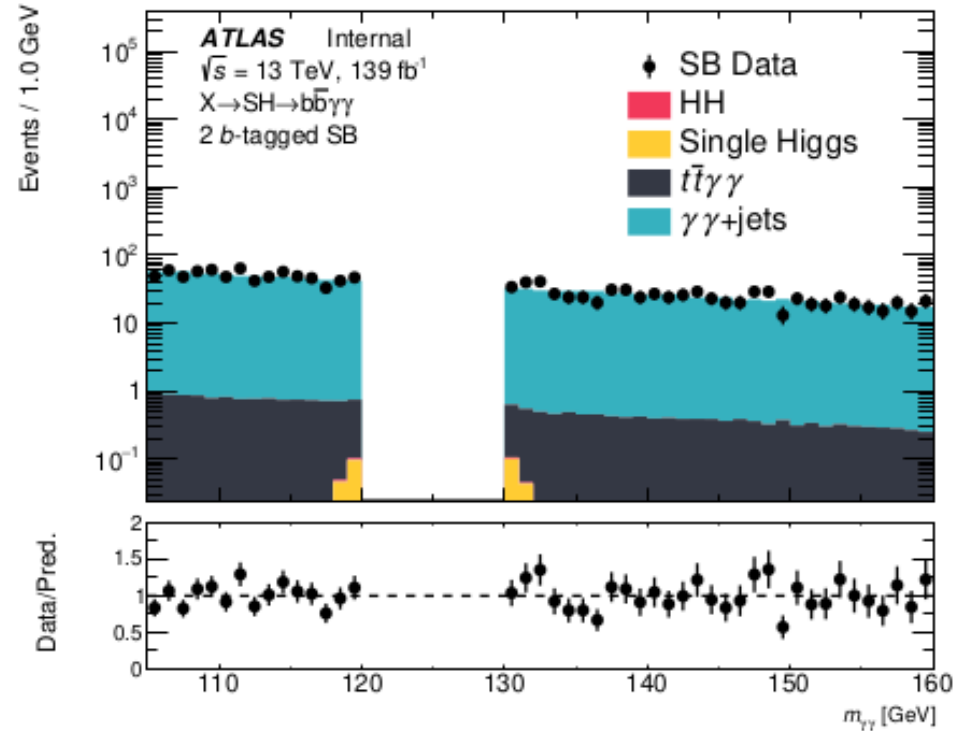
	Selection
<i>HH</i> ggF+VBF	$1.827 \pm 0.004$
<i>WH</i>	$5.97 \pm 0.021$
<i>VBFH</i>	$8.333 \pm 0.042$
<i>ZH</i>	$5.941 \pm 0.015$
<i>bbH</i>	$2.973 \pm 0.047$
<i>ggH</i>	$48.532 \pm 0.202$
<i>ggZH</i>	$1.581 \pm 0.01$
<i>tHjb</i>	$2.681 \pm 0.05$
<i>tWH</i>	$0.572 \pm 0.01$
<i>ttH</i>	$11.681 \pm 0.017$
$Z(\rightarrow qq)\gamma\gamma$	$53.728 \pm 0.8$
$t\bar{t}\gamma\gamma$	$49.78 \pm 0.142$
$\gamma\gamma$ +jets	$16298.8 \pm 16.031$
Total SM	$16492.398 \pm 16.053$



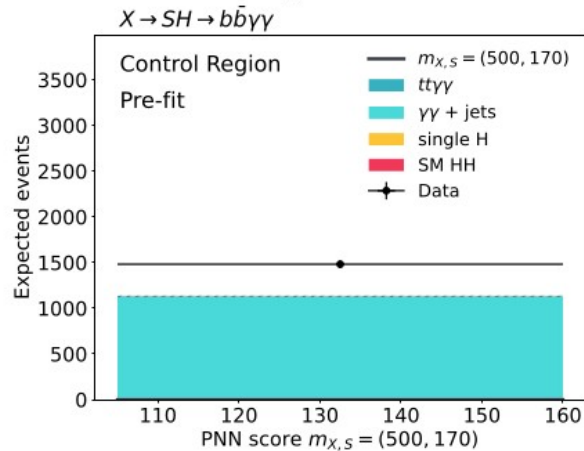
# Analysis recap – Strategy

- The  $m_{\gamma\gamma}$  distribution is used to split the events into a SR with  $120 < m_{\gamma\gamma} < 130$  GeV and a sideband control region (CR)

Here for the  
2 b-tagged jets selection



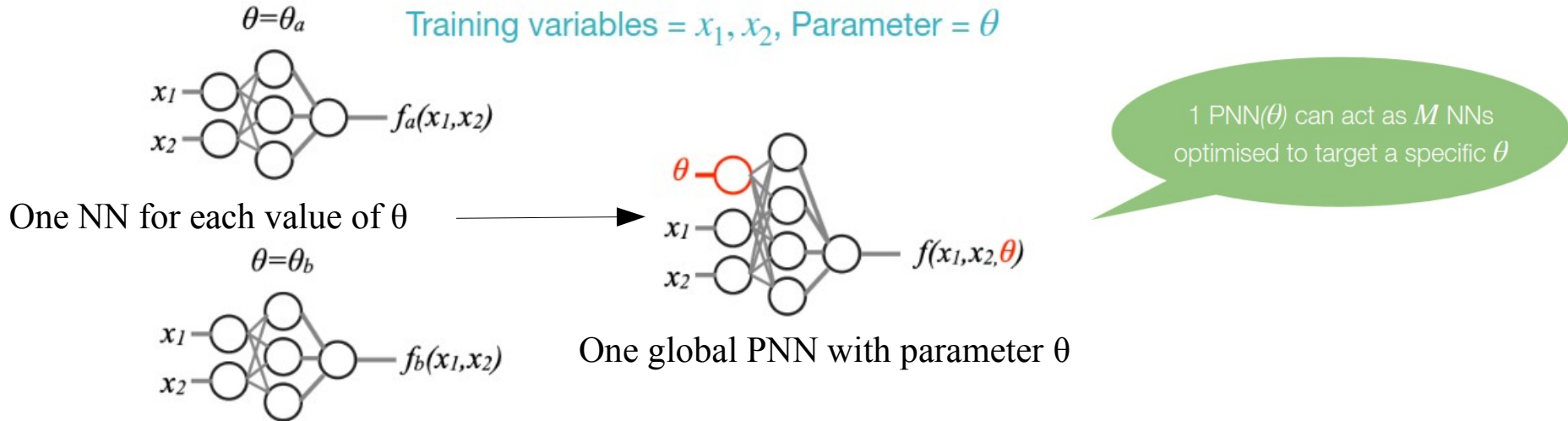
- The CR allows to correct the normalisation of the  $\gamma\gamma$  + jets events using true data





# PNN discriminant

- A parameterised Neural Network (PNN) is used as discriminant in the SR. It is trained with simulated events that pass each selection (SR + SB)



- Two separate PNNs :

- **2 b-tagged :**

Parameter  $\theta = \mathbf{mX, mS}$

Training samples : signal, ttH, ggH, ZH and  $\gamma\gamma$  + jets backgrounds (no HH – too signal-like and confuses the network)

Training variables :  $m_{jj}$  and  $m_{\gamma\gamma jj}^* = m_{\gamma\gamma jj} - (m_{\gamma\gamma} - 125 \text{ GeV})$

- **1 b-tagged :** target low mS/mX values

Parameter  $\theta = \mathbf{mX}$

Training samples : signal, VBFH, HH, ttH, ggH, ZH and  $\gamma\gamma$  + jets backgrounds

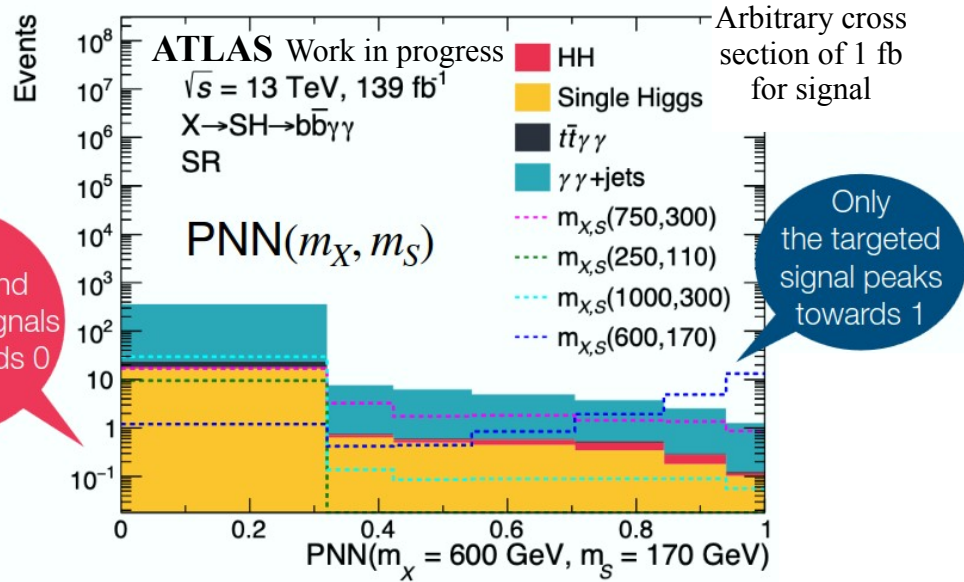
Training variables : b-jet  $p_T$  and  $m_{\gamma j}^* = m_{\gamma j} - (m_{\gamma} - 125 \text{ GeV})$





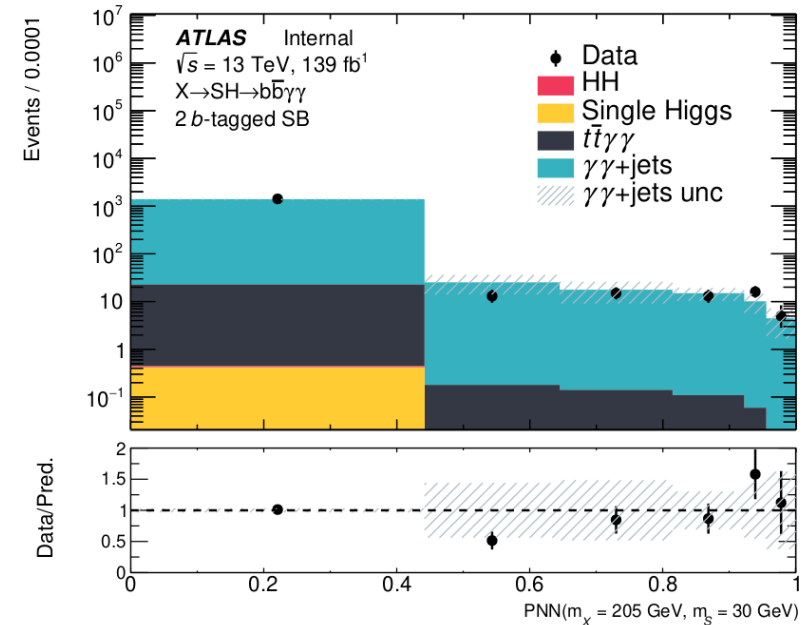
# PNN distribution

- Example of PNN distribution



- PNN shapes of backgrounds and signal comes from MC samples
- The binning is constrained to have at least 1 background event in every bin

- Final results are computed with a binned log-likelihood fit on the PNN distribution
- Consistency between data and MC is checked in the SB

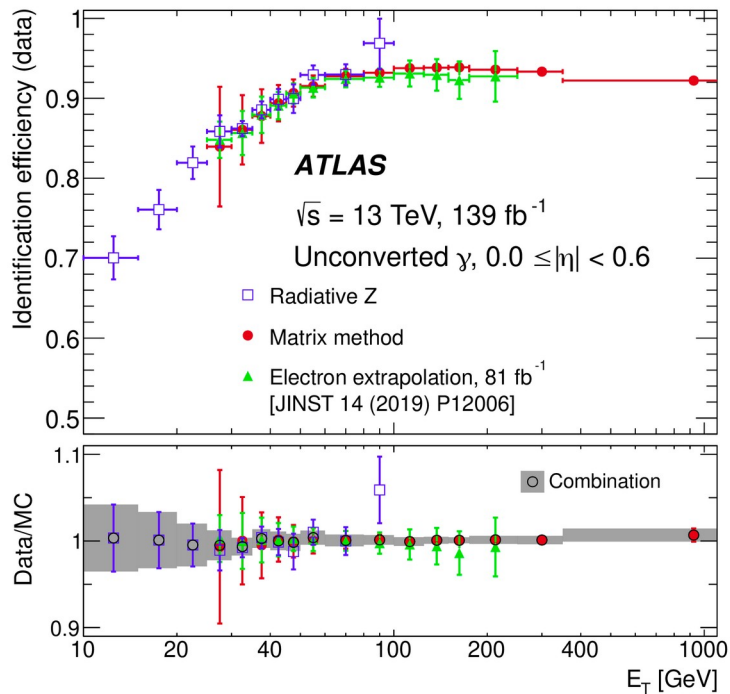




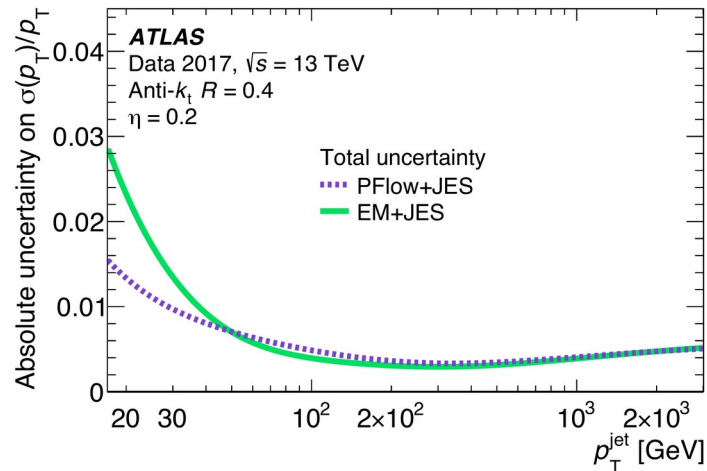
# Experimental Systematics

- Physics is an experimental science → we have uncertainties affecting the measures
- Eventually at our analysis level it can have various impacts :
  - Particle identification can change the number of events in the CR and SR
  - Flavour tagging change the number of b-tagged jets
  - $p_T$  and energy resolution can change the position and width of the peak in the  $m_{\gamma\gamma}$ ,  $m_{bb}$  and  $m_{bb\gamma\gamma}$  distributions and eventually the shape of the PNN distribution

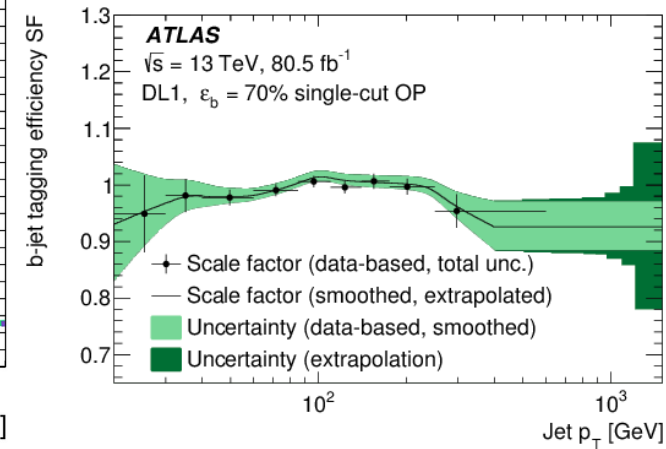
Photon identification efficiency as a function of  $E_T$



Photon  $p_T$  resolution as a function of  $p_T$



b-jet tagging efficiency as a function of  $p_T$





# Experimental Systematics – Yield change

- Systematics change both the yields and the shape of the PNN distribution of the samples
- Here we look at the yields change first

Main uncertainties

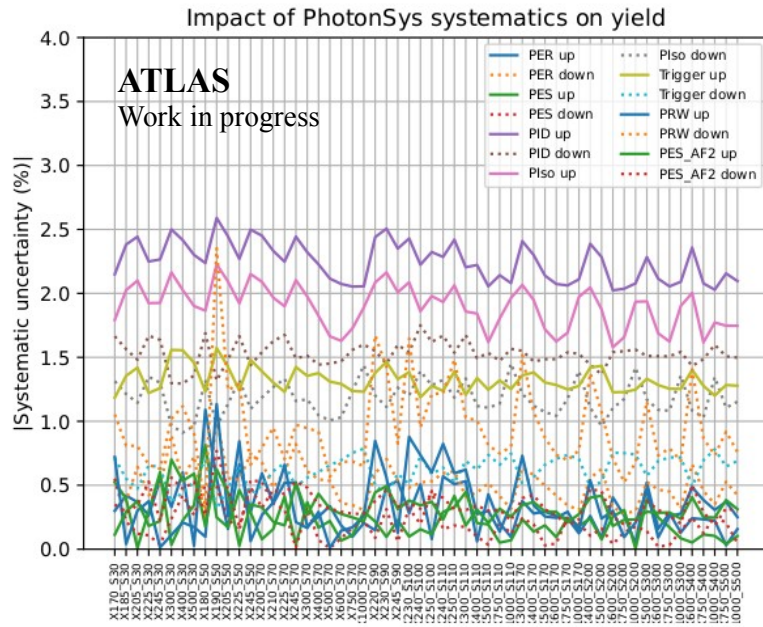
Table 18: Uncertainty [in %] on the yield for backgrounds in the 2  $b$ -tagged category.

	Source	Yield uncertainty [%]								
		ttH	ZH	ggHH	ggH	ggZH	tHjb	VBFHH	Zqqyy	Zbbyy
Event-based	Photon Trigger	1.01	1.02	1.00	1.04	0.98	1.03	1.02	1.00	1.25
	Pile-up reweighting	0.88	0.76	0.56	0.43	0.60	0.99	0.60	3.69	1.15
Photon	Photon Energy Resolution	0.42	0.42	0.34	0.42	0.43	0.63	0.43	2.79	0.78
	Photon Energy Scale	0.17	0.18	0.12	0.07	0.11	0.18	0.24	14.53	0.90
	Photon ID	1.59	1.61	1.44	1.64	1.49	1.60	1.59	1.72	1.96
	Photon Isolation	1.55	1.57	1.45	1.60	1.46	1.59	1.59	1.27	1.88
Jet	Jet Energy Scale	1.36	0.94	0.55	1.81	0.74	0.76	0.72	5.30	1.09
	Jet Energy Resolution	7.33	4.60	2.91	7.50	3.36	4.88	3.08	0.68	5.37
Flavour-tagging	b-jet efficiency	2.07	2.99	2.51	3.05	2.55	2.30	2.83	0.10	3.36
	c-jet efficiency	0.40	0.71	0.06	1.68	0.60	0.92	0.07	13.12	0.22
	light-jet efficiency	0.79	0.38	0.40	2.72	0.51	0.90	0.42	1.85	0.48

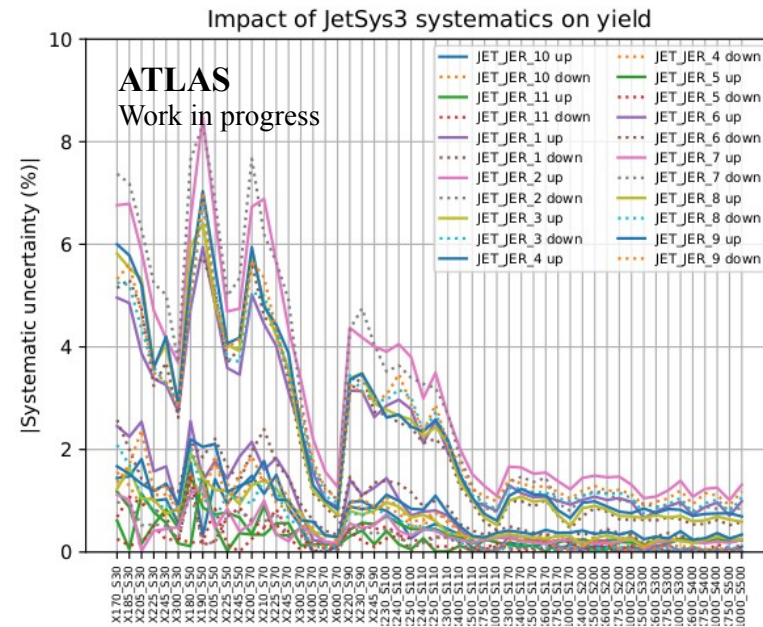


# Experimental Systematics – Yield change

- For signal, yields changes are dependant on  $m_S$  and  $m_X$

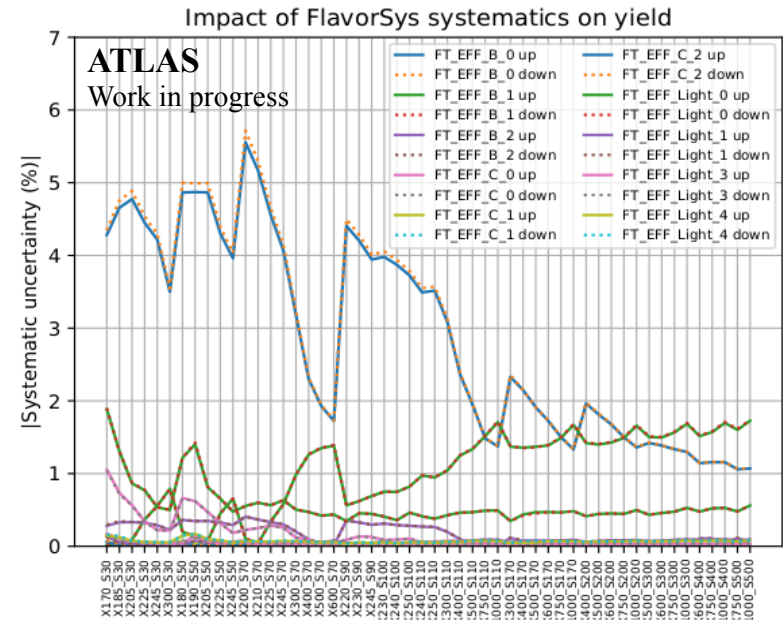


(b) 2  $b$ -tagged



(b) 2  $b$ -tagged

For instance here are the systematics NP related to photon, jet and flavour tagging as a function of  $m_S$  and  $m_X$

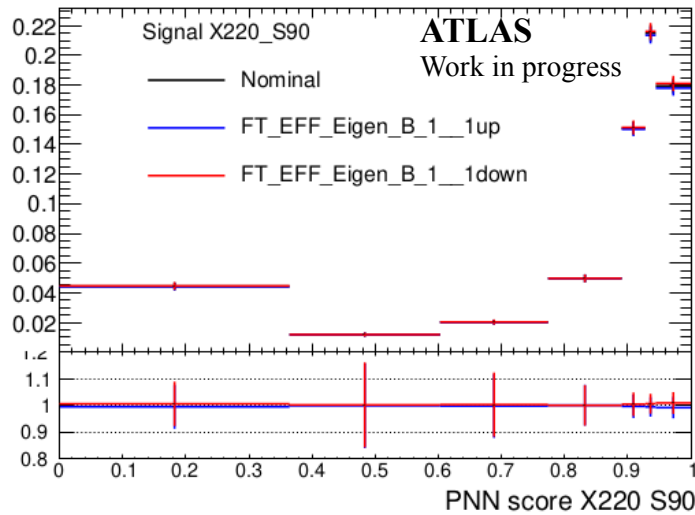


(b) 2  $b$ -tagged

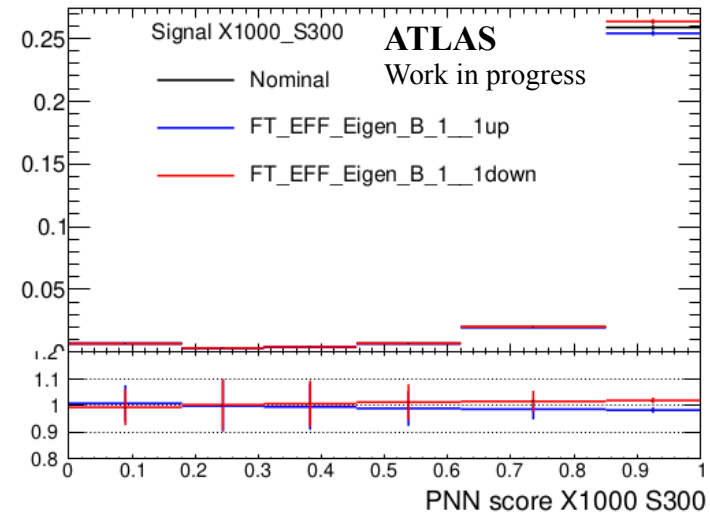
Jet energy resolution systematics are the most important but remain below 10%



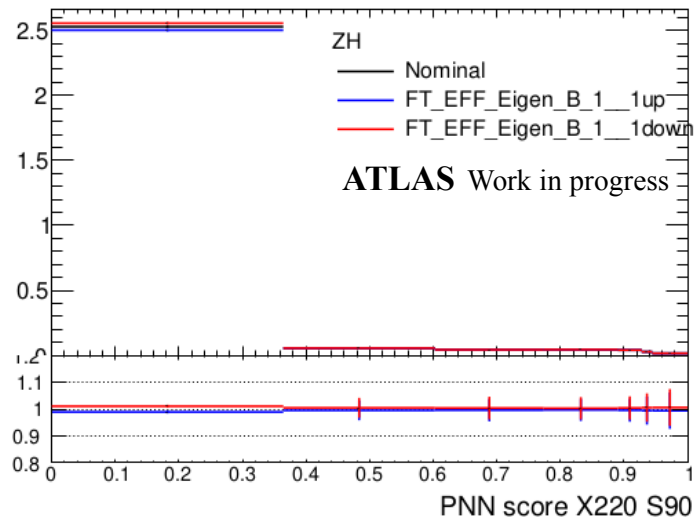
# Experimental Systematics – Shape change



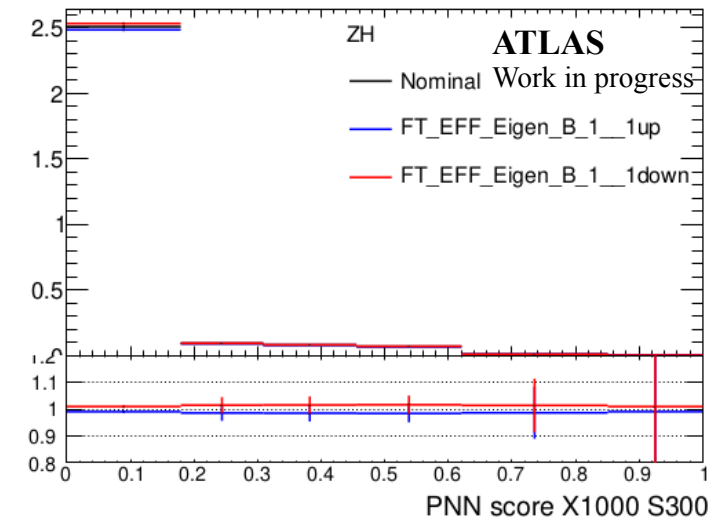
(a) SH signal,  $m_X = 220$  GeV,  $m_S = 90$  GeV



(b) SH signal,  $m_X = 1000$  GeV,  $m_S = 300$  GeV



(e) ZH,  $m_X = 220$  GeV,  $m_S = 90$  GeV



(f) ZH,  $m_X = 1000$  GeV,  $m_S = 300$  GeV

- NB : for  $\gamma\gamma + \text{jets}$ , only shape changes are used as normalisation is imposed by the sideband

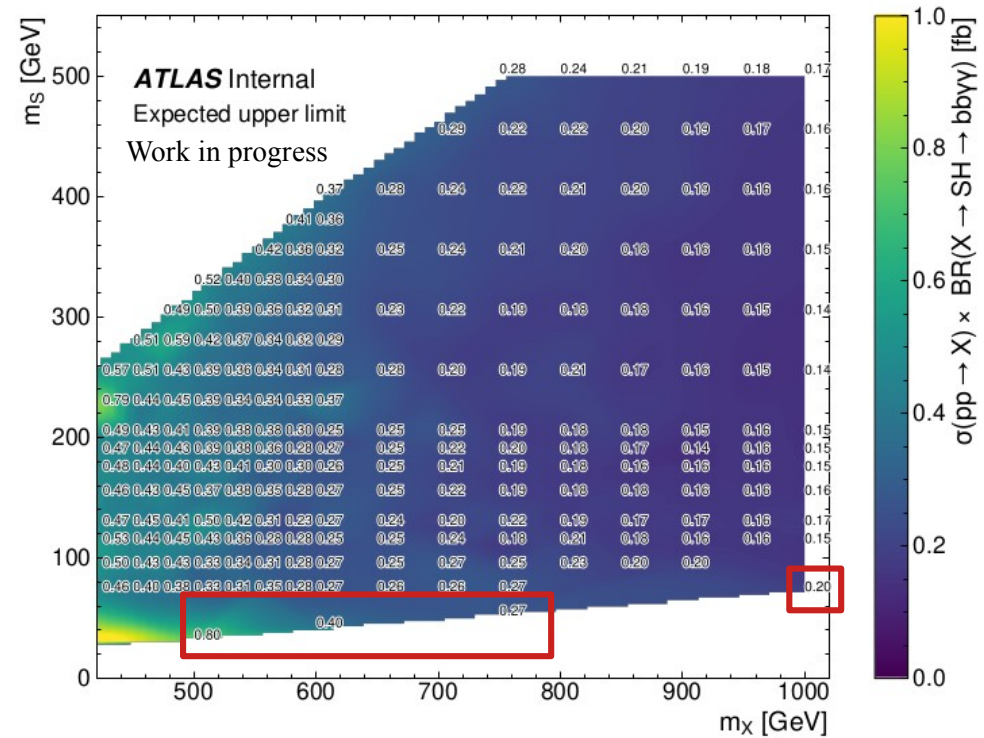
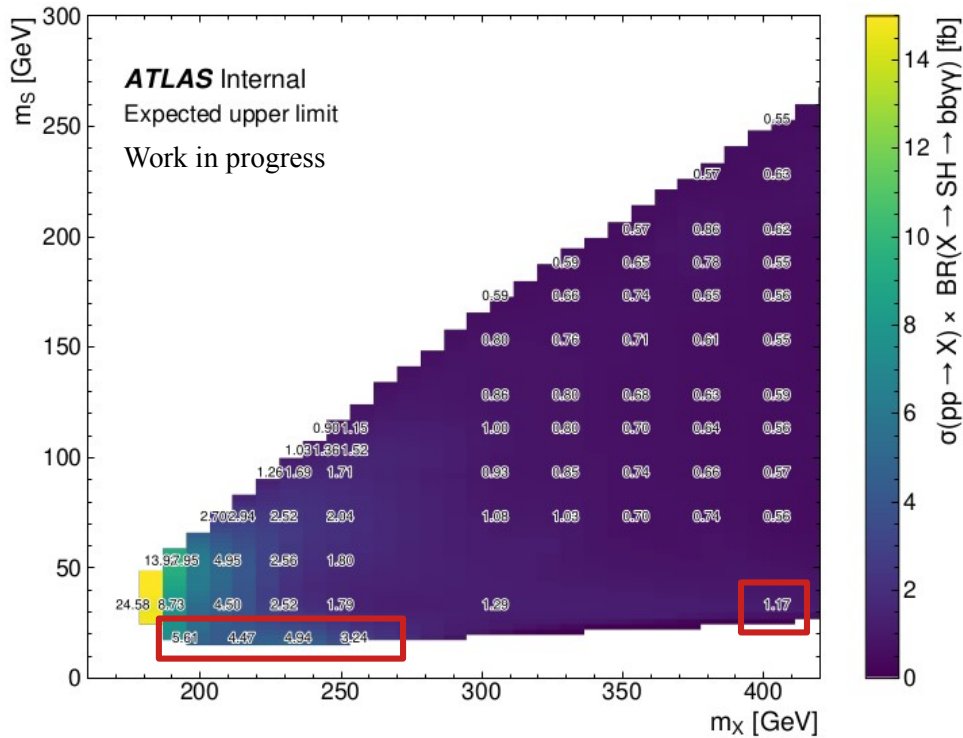


# Blinded expected limits

- If no signal is observed, upper limits on the cross section of the  $X \rightarrow SH$  signal in the  $b\bar{b}\gamma\gamma$  final state can be set
- I can only show blinded results (i.e using Monte-Carlo and not true ATLAS data)  
Not final results but gives an idea of the analysis sensibility



1 b-taged jet points

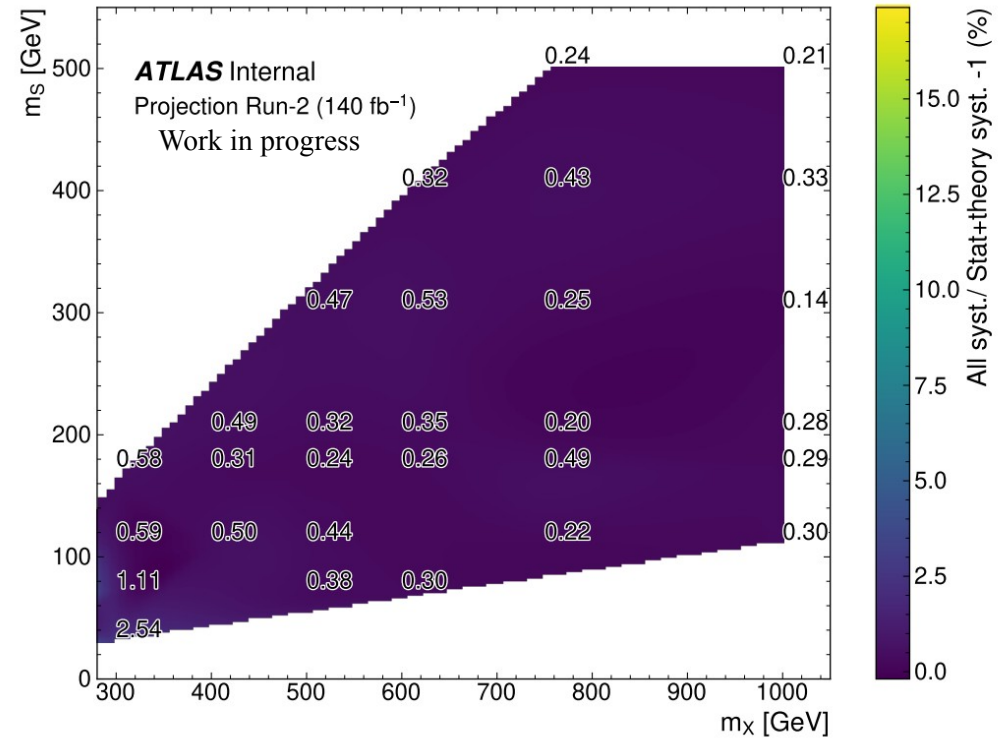
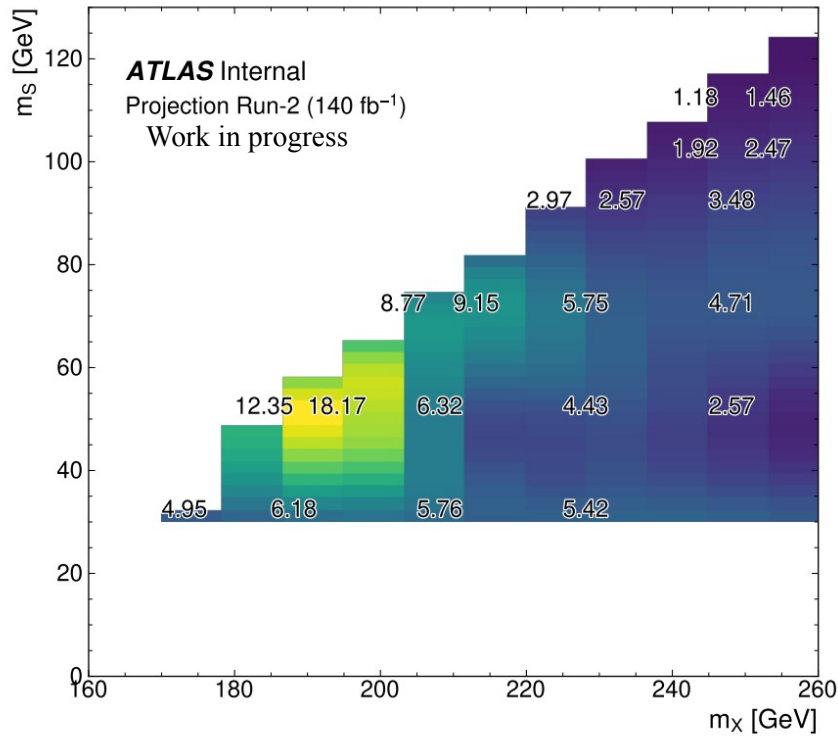


- Limits range from 0.15 to 25 fb  
Sensibility is better in high mass region



# Impact of experimental systematics uncertainty

- Here we plot the ratio between blinded expected limits and the limits obtained without taking into account experimental systematics to check their impact :

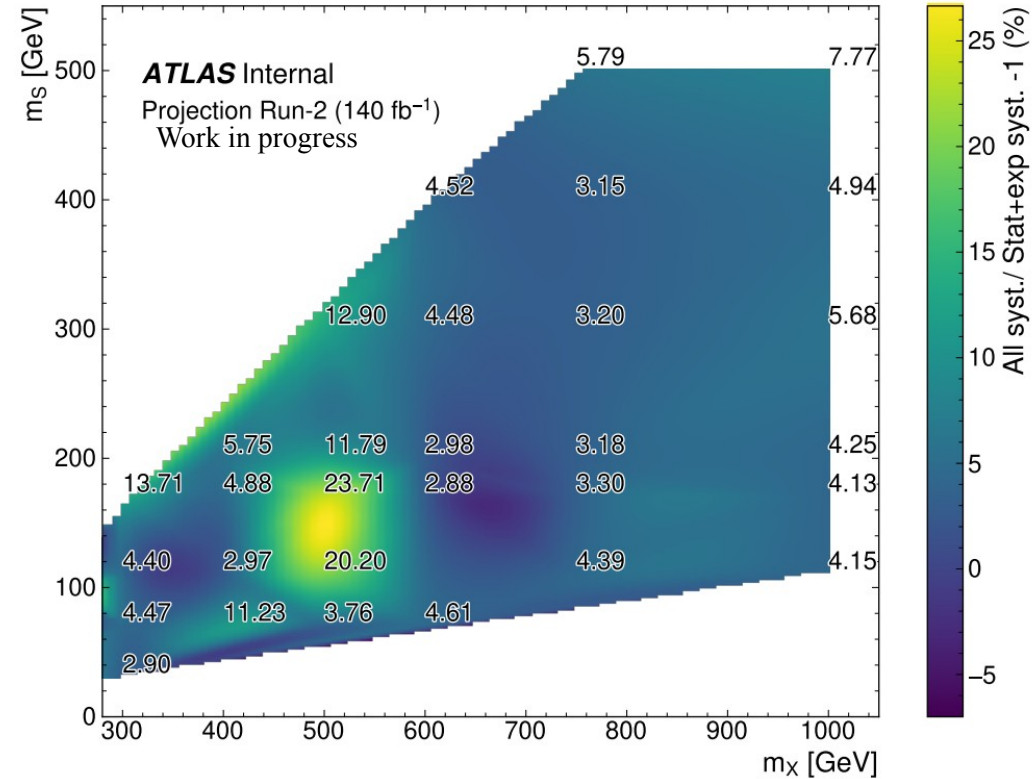
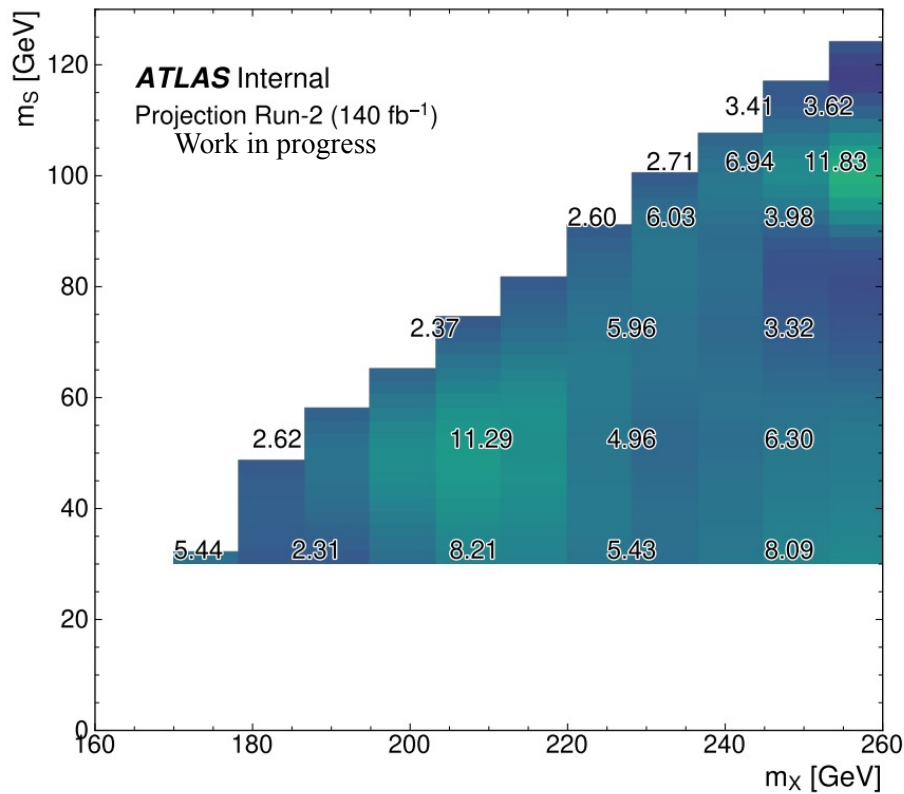


- Their impact on limit can reach 18% but are mostly between 5-10% at low mass and below 1% at high mass



# Impact of theoretical systematics uncertainties

- Same plot as before but with theoretical uncertainties



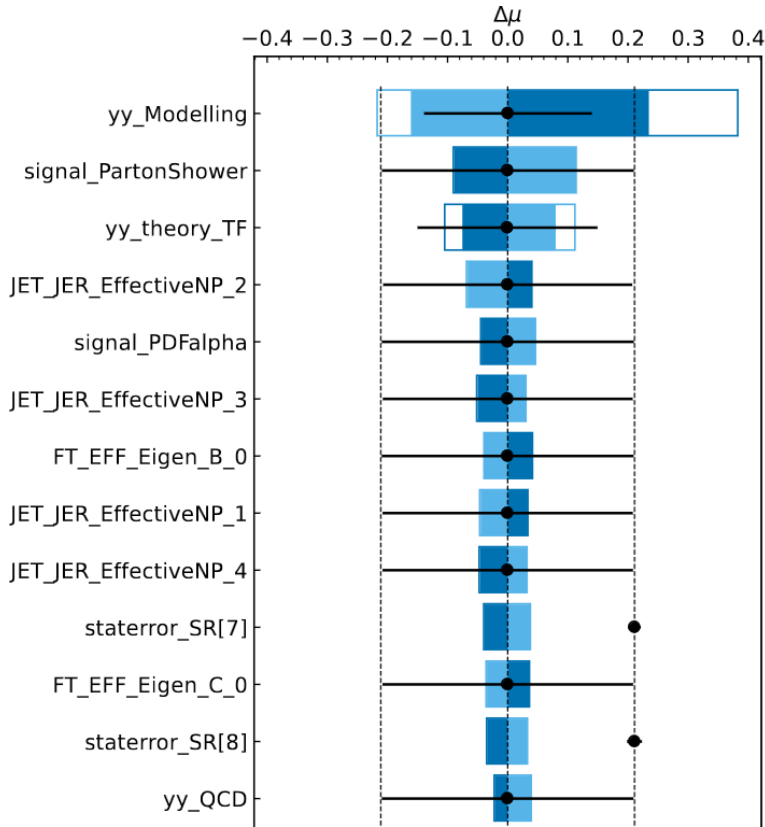
- Theory systematics impact depends a lot on  $m_X$  and  $m_S$  and can reach 20%
- They are dominated by  $\gamma\gamma$  modelling  
Other theoretical systematics account for 3-4%



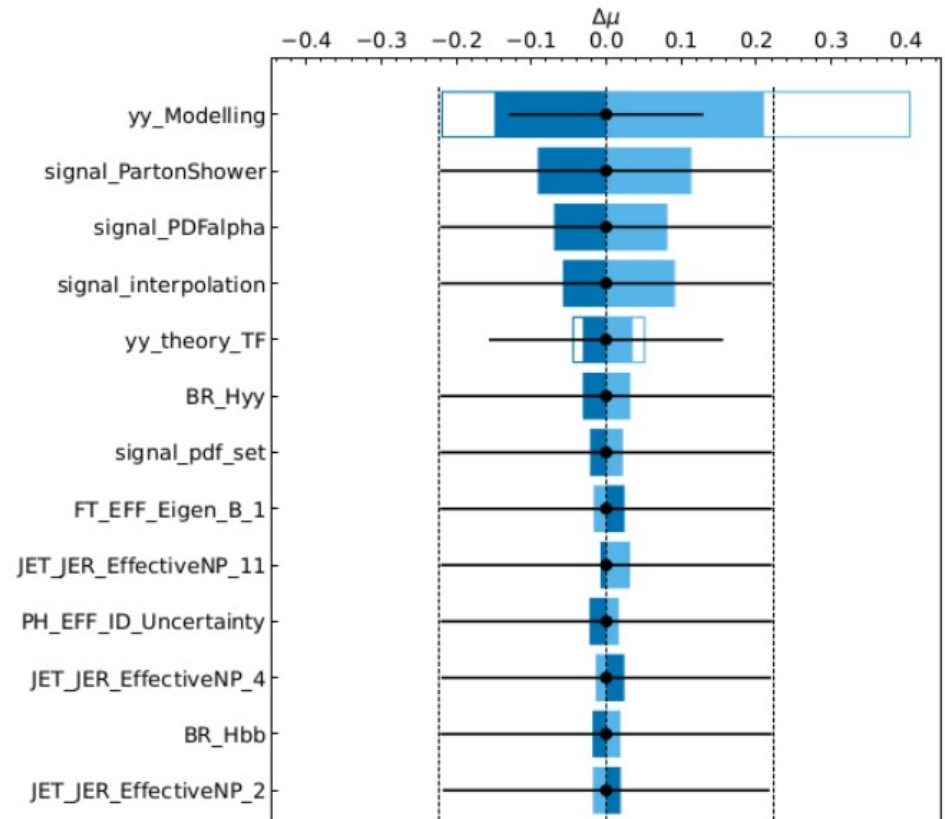


# Impact of systematics uncertainty – ranking plots

- Ranking plots : class the uncertainties with the impact they have on the fit POI



mX = 250 GeV, mS = 100 GeV



mX = 575 GeV, mS = 200 GeV

● pulls  
 □ pre-fit impact:  $\theta = \hat{\theta} + \Delta\theta$     ■ post-fit impact:  $\theta = \hat{\theta} + \Delta\hat{\theta}$   
 □ pre-fit impact:  $\theta = \hat{\theta} - \Delta\theta$     ■ post-fit impact:  $\theta = \hat{\theta} - \Delta\hat{\theta}$

- Largest systematics is the modelling of the  $\gamma\gamma$  + jets background
- Largest experimental systematics are flavour tagging and jet energy resolution



# Summary

- A search for a resonant scalar particle  $X$  decaying into a scalar  $S$  and SM Higgs is performed on the  $X \rightarrow SH \rightarrow b\bar{b}\gamma\gamma$  channel with the ATLAS Run-2 data
- Most interesting point in this analysis is the parameterised neural network (PNN) that has been developed to target the signal for any values of  $m_X$  and  $m_S$
- Analysis currently in ATLAS internal review  
New analyses will come with Run-3 data, for both HH and SH !



Back-up



# All systematics uncertainties

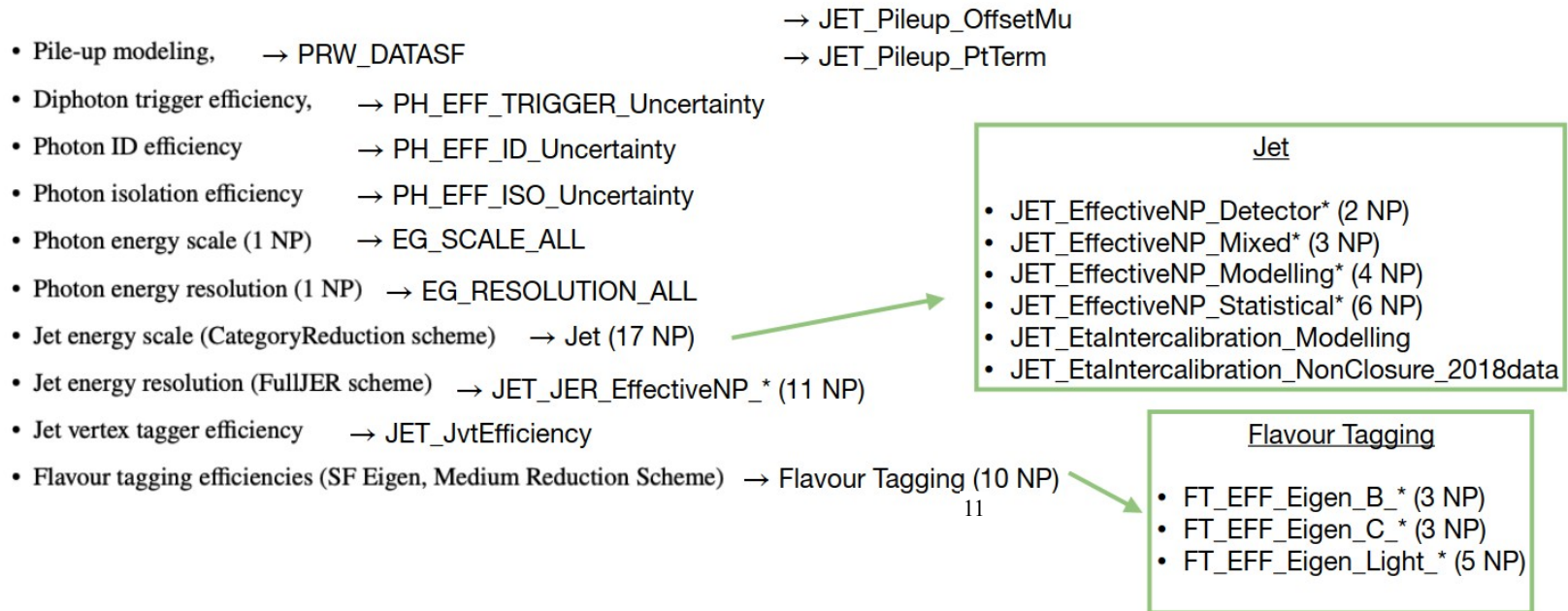
## Main uncertainties

		Signal	HH ggF	HH VBF	ttH & ZH	Other Single Higgs	Continuum $\gamma\gamma$ +jets
Theory	Normalisation	$BR(H \rightarrow \gamma\gamma)$	$BR(H \rightarrow \gamma\gamma)$ $BR(H \rightarrow b\bar{b})$ PDF+ $\alpha_S$ Scales + $m_t$	$BR(H \rightarrow \gamma\gamma)$ $BR(H \rightarrow b\bar{b})$ PDF+ $\alpha_S$ Scales	$BR(H \rightarrow \gamma\gamma)$	$BR(H \rightarrow \gamma\gamma)$  PDF+ $\alpha_S$ Scales	$\gamma\gamma$ transfer factor
	Shape+Norm.	Scales, PDF+ $\alpha_S$ Parton shower Interpolation	Parton Shower		Scales, PDF+ $\alpha_S$ Parton Shower		Scales, PDF+ $\alpha_S$ Modelling
Exp.	Shape+Norm.	Pile-up modelling Diphoton trigger efficiency Photon identification and isolation efficiency Photon energy scale and resolution (all exp. systematics are neglected for $bbH$ , $tH$ and $VBF H$ ) Jet energy scale and resolution Jet vertex tagger efficiency Flavour tagging efficiency					



# Experimental Systematics

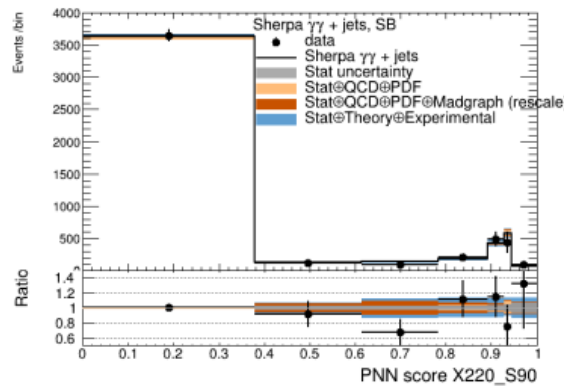
- In the analysis framework experimental systematics are studied through MC samples at  $\pm 1\sigma$  away from nominal values from each effect
- They have been produced for major backgrounds only : ttH, ZH, ggH, ggHH and VBFHH, ggZH,  $\gamma\gamma$  + jets and  $Z(bb/qq)\gamma\gamma$
- Systematics are treated as nuisance parameters (NP) in the fit  $\rightarrow$  47 in total !



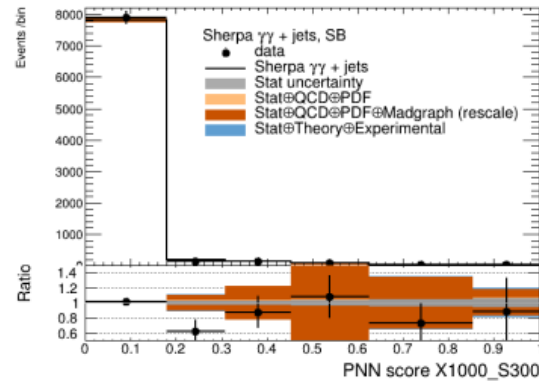


# Theoretical systematics uncertainties

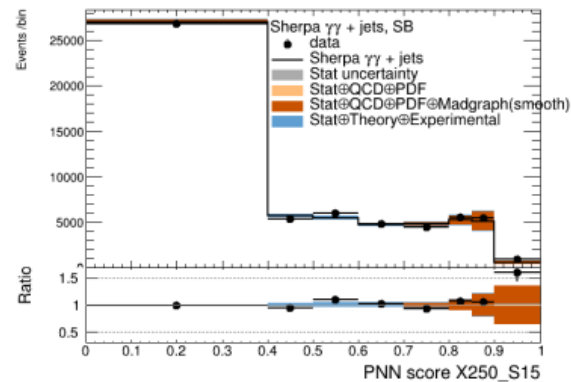
- Largest theoretical uncertainty is the modelling of  $\gamma\gamma + \text{jets}$  events which is difficult to handle
- Normalisation of the  $\gamma\gamma + \text{jets}$  events is determined by a normalisation factor from the sideband distribution
- The uncertainty regarding the modelling is evaluated by comparing simulated events from two different MC generators : Sherpa and MadGraph



(a)  $m_X = 220$  GeV,  $m_S = 90$  GeV



(b)  $m_X = 1000$  GeV,  $m_S = 300$  GeV

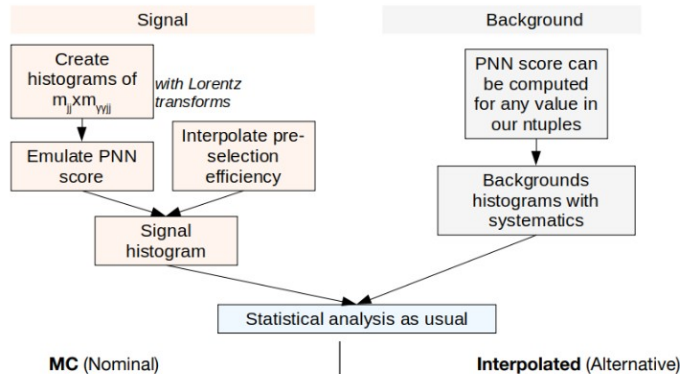


(c)  $m_X = 250$  GeV,  $m_S = 15$  GeV



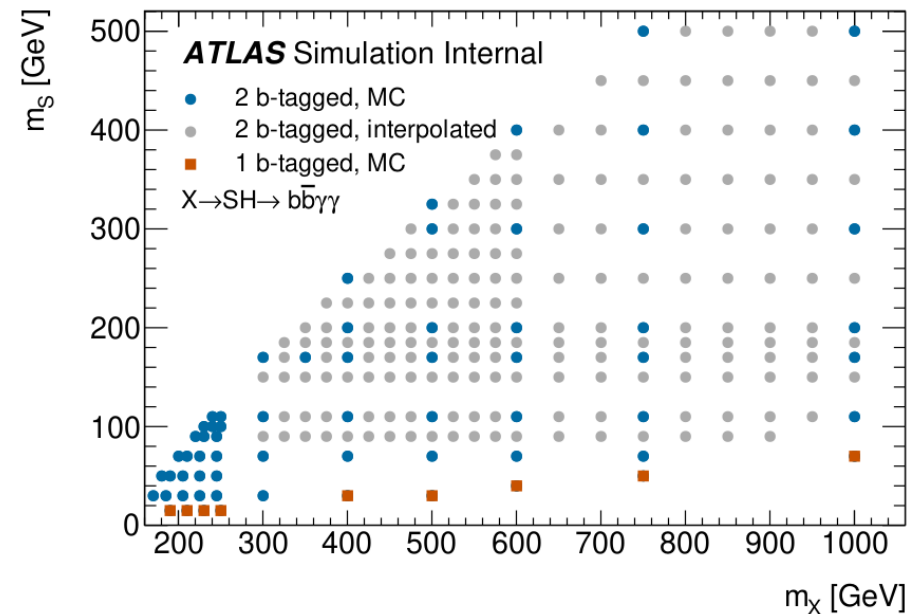
# Interpolation strategy

- Why interpolate ?
  - We need to be able to look for any signal in the region and granularity is not precise enough with MC samples
- Interpolation works separately for signal and background
  - PNN score can be computed for any  $m_X$ ,  $m_S$  values in background samples
  - For signal we need to interpolate both the yields and the PNN shape
    - The shape is obtained with Lorentz transforms
    - The yields are obtained using Delaunay triangulation from the available MC samples



- Signal from MC
- **Shape** signal PDF uncertainties from MC
- **Shape** signal exp systematics from MC

- Signal from Lorentz interpolation
- **Flat** signal PDF uncertainties from the signal-most like bin
- **Flat** signal exp systematics from the signal-most like bin
- Additional shape interpolation uncertainty from the Lorentz interpolation fit parameters

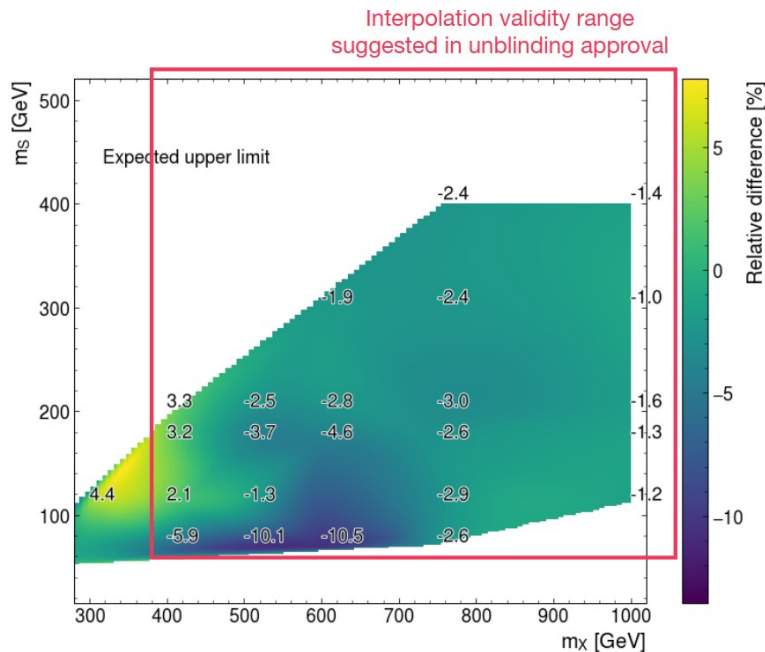
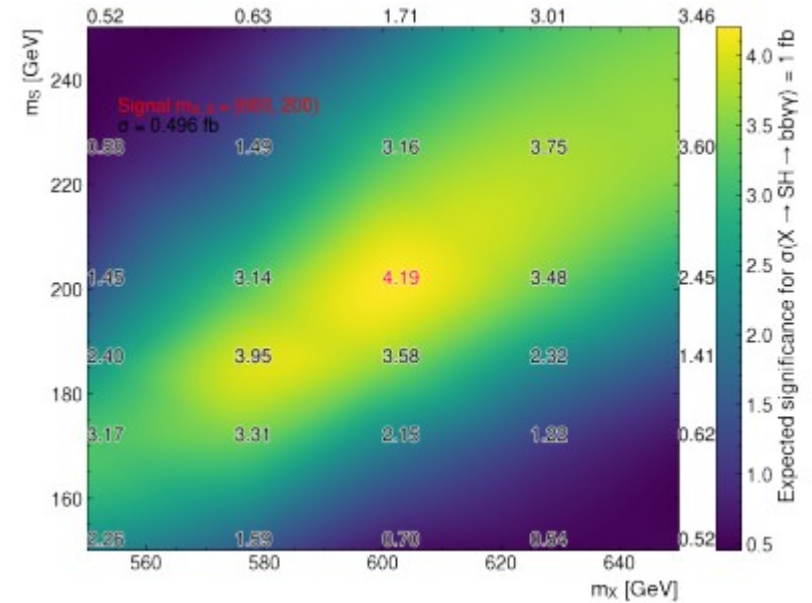




# Interpolation range and validity

- Where should we interpolate ?
- Injection tests are made to ensure granularity is enough to allow us to be sensible to any signal in the probed region

→ we inject signal at  $\sigma = 2 \times$  expected limit and want at least one neighbouring point to have an expected significance  $\geq 3$



- Validity of interpolation is evaluated by comparing interpolated and MC signal limits
  - Difference is below 5% for most points with a maximum of  $\sim 10\%$
  - Interpolation is more difficult in low  $m_S$  regions where there are some jets overlap
- Interpolation is made for points with  $m_X \geq 300 \text{ GeV}$  and  $m_S \geq 70 \text{ GeV}$