# ATHENA X-IFU'S TES/SQUID DETECTION CHAIN

**DRTBT 2024** 

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#### Outline

- Athena and X-IFU instrument overview
  - Scientific context
  - Athena mission
  - > X-IFU instrument
- TES detectors for X-IFU
  - > TES small signal model and main resolution drivers
  - Influence of count rate
- Evolution of TES design across project lifetime
- The X-IFU TDM detection chain
  - Functioning principle and overview
  - Dimensioning rules
    - Case study: how to optimize the pixel speed wrt. performance targets
    - Current dimensioning point in large signal

- Going beyond signal to noise ratio optimization
  - TDM bandwidths constraints
  - Crosstalk
  - Gain drifts and associated correction algorithms
- What is next for X-IFU?
  - Long term timeline
  - Short term demonstration activities

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# ATHENA X-IFU SCIENTIFIC CONTEXT AND INSTRUMENT OVERVIEW







#### **The Universe seen in X-rays**





#### The X-ray data in astronomy

- The final output of most X-ray instruments is an « event list »
- Contrary to longer wavelength experiments where fluxes are measured, each individual X-ray photon is measured. For each one of them, we have:
- Its arrival time
- Its energy
  - The impacted pixel (= position)
    - Auxiliary data
- From this we can construct
  - Energy spectra
  - Images
  - Light curves and Fourier spectra
    - o Oscillation of the emission of compact objects
    - o Pulsars
    - o ...











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#### X-ray spectroscopy: what we learn from X-ray lines

- Rest frame energy: gas composition (atomic element, ionization state)
  - The most famous X-ray line in astronomy: neutral iron = 6.4 keV in rest frame
- **Energy** shift: gas velocity from Doppler effect (combined with relativistic redshift from source distance)
  - > 20km/s Doppler shift at 6 keV =  $E \times \frac{\Delta v}{c}$  = 0.4 eV
- Line integral: abundance (relative density of the element in the gas)
- Line broadening: state of the gas (temperature, turbulence, …)



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#### High resolution spectroscopy (microcalorimeters vs. CCDs)



Note: Gratings based CCD spectrometers typically have a better resolution than microcalorimeters at low energy but the spatial information is blurred with the spectral one



#### True spatially resolved spectroscopy: from Hitomi to Athena



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# **Athena science objectives**

- Athena is the next generation large X-ray observatory
- Initially selected to address the hot and energetic Universe science theme
- How did matter assemble to form the large scale structures that we see today?
  - Formation and evolution of the largest structures of the Universe
  - Detection of the « missing baryons »
  - Mainly through the observation of faint extended sources
  - How did black holes grow and shape the Universe?
    - o Black hole formation
    - Feedback on galaxy and clusters formation
    - o Accretion physics
    - Mainly through the observation of point sources, some among the brightest of the X-ray sky
  - A large part dedicated to "observatory science" (open observation calls)







# **The Athena Observatory**

- ATHENA: Advanced TElescope for High ENergy Astrophysics
  - Launch scheduled for 203X on Ariane 6
  - Few tons and few kW ~ 14 m long satellite
  - Halo orbit around the L1 Lagrange point
    - Preferred to L2 for stability of soft protons flux (create instrument background)







### **The Athena mirror**

- Silicon pore optics mirror (Wolter I) working in grazing incidence
  - Millions of pores stacked into so-called mirror modules
    - Stacks of Si wafers assembled in a mostly fully automated process
- Effective area an order of magnitude larger than current X-ray missions
- Accommodated on an hexapod platform. Allows:
- Switching between instruments (WFI, X-IFU)
- Defocusing for bright point sources observations





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# **Payload Compartment (PLC) and Wide Field Imager (WFI)**

- The Payload Compartment (PLC)
  - Accommodates both Athena instruments (X-IFU + WFI)
  - ~ 2000 kg, few kW
- Major overhaul (aka rescope) of the mission architecture in 2023 following ESA budget constraints
  - Addition of passive cooling down to 50 K for X-IFU thanks to V-groove system
  - o Reduction of WFI allocated volume and sky access angle
- Wide Field Imager (WFI)
  - ➢ 40'x40' FoV
  - DEPFET detectors (depleted p-channel field-effect transistor)
  - Moderate energy resolution (170 eV at 7 keV)
  - Small detector mounted out of focus to observe the brightest objects of the X-ray sky







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# **Main X-IFU performance requirements**

Parameter	Requirement
Energy range	0.2 – 12 keV
Energy resolution	3 eV at 7 keV design target (4 eV requirement)
Field of view	4' equivalent diameter
Pixel size	~ 5" (1504 pixels)
Instrument efficiency	>12%, >56%, >63%, >42% @ 0.35, 1.0, 7.0, 10.0 keV
Count rate capability	80% high res throughput at 1 mCrab (goal at 10 mCrab)
	80% high res throughput on Perseus cluster (goal at CasA)
	50 % 10 eV throughput at 1 Crab (5- 8 keV)







#### The X-IFU in the PLC







#### The X-IFU in the PLC





Netherlands

France

Italy

Belgium

Finland





Spain











#### The X-IFU in the PLC



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#### The new cryogenic architecture

- New architecture following mission rescope
- Dewar now under consortium responsibility (AVS). Passively cooled down to 50K via successive radiator panels
- Active cooling down to 4.5K via remove US cooler
  - Compressors and electronics in the « basement »
  - o Competitive process in the US
- Last stage cooling via multi-stage ADR (CEA SBT)
  - > 5 ADR pills
    - $\circ$  ~ 1 mW at 2K
    - $\circ~~$  ~ 10  $\mu W$  at 325 mK
    - ~ 1µW at 50 mK (~ 60 nW without margin from main detection chain cold electronics, ~ half in TES shunts, half in SQUIDs)







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#### The 4K core

- The « heart » of the instrument is organized in a 4K core
  - ~ 45 kg, suspended inside 50 K dewar (likely via struts) with an interface at the green plate level
  - Faraday cage (on top of the 50 K dewar one) to protect scientific signals

### Accommodates two main subsystems

- Focal Plane Assembly suspended at 2K by CFRP struts
- Multistage ADR providing 2K, 325 mK (thermal intercept) and 50 mK cooling
- Superconducting looms between 4K and 2K
- (« Tubes » are principle representations of harnesses to show interfaces)





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# **TES DETECTORS FOR X-IFU**





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### **Detection principle: Transition Edge Sensors (TES)**





- **TES thermal equilibrium and electro-thermal feedback**
- The TES is maintained inside its transition by the Joule power provided by its bias
- Under (quasi) constant bias voltage, we have a negative electro-thermal feedback loop:

$$\triangleright \quad P_{Joule} = \frac{V^2}{R_{TES}(T,I)}$$

$$\succ T \uparrow \Rightarrow R \uparrow \Rightarrow P_{Joule} \downarrow \Rightarrow T \downarrow$$

- Allows a stable operation of the TES in transition and speeds up the cooling of the TES after an X-ray impact => pushes the information towards higher frequency
- Note that devices with negative dR/dT (e.g. doped semiconductor detectors) in turn require being biased via a current





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# Main TES parameters and governing differential equations

The response of TES detectors is governed by two coupled differential equations

$$L\frac{dI}{dt} = V_0 - IR_L - IR_{TES} + Noise$$

$$\frac{dT}{dt} = -P_{sink} + R_{TES}I^2 + P_{Xray}$$
  
with  $P_{sink} = K(T^n - T^n_{bath}) \approx P_{sink,0} + G\delta T$   
 $P_{Xray} = E_X/dt$ 

The TES transition itself is characterized locally by two main <u>unitless</u> parameters:

$$\alpha = \frac{\partial \log R}{\partial \log T} \bigg|_{I_0} \qquad \beta = \frac{\partial \log R}{\partial \log I} \bigg|_{T_0}$$
$$R_{TES} \approx R_0 + \alpha \frac{R_0}{T_0} + \beta \frac{R_0}{I_0}$$

Pixel parameters definitions		
$R_n$ (m $\Omega$ ) – TES normal resistance		
$T_0$ (mK) – TES bias point temperature		
$G_b$ (pW/K) – Thermal conductance		
C (pJ/K) – Total heat capacity		
n – Thermal exponent		
$T_{bath}$ (mK) – Bath temperature		
$R_0$ (m $\Omega$ ) – TES bias point resistance		
$R_s$ ( $\mu\Omega$ ) – Shunt resistance		
$I_0$ (µA) – TES current at bias point		
$\alpha (T/R \partial R/\partial T)$		
$\beta (I/R \partial R/\partial I)$		
$\mathcal{L}_{I}(P_{0}\alpha/GT_{0})$ – ETF loop gain		
M – Unexplained Johnson noise factor		
L (nH) – Damping inductance		
$\tau_+/\tau$ (µs) – Detector time constants at L		
$f_{eff}$ (Hz) – Effective frequency		



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 $P_{Xray} = E_X/dt$ 







Plot from M. Gonzalez's LTD intro

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WARNING: Different convention from M. Gonzalez's Itd intro (see above)



#### **Small signal TES impulse response: an X-ray pulse**

In the small signal limit (see Lindeman+00, Irwin&Hilton+05), the TES response follows

 $\delta I \propto (e^{-t/\tau_+} - e^{-t/\tau_-})$ 

with  $1/\tau_+$  and  $1/\tau_-$  the eigen values of the differential system

- \* The response is stable (critically or overdamped) if  $\tau_+$  and  $\tau_-$  are real. This reduces to a constraint on the circuit inductance
  - > The circuit is underdamped (there are oscillations) if  $L_{crit} \leq L \leq L_{crit+}$
  - > Typically, TES circuits are dimensioned such that  $L \leq L_{crit-}$
- In the low inductance and stiff bias voltage limit ( $R_{shunt} \ll R_0$ ):

 $\tau_{+} \rightarrow \frac{L}{R_{shunt} + R_{0}(1 + \beta)} \quad \text{electrical time constant}$  $\tau_{-} \rightarrow \tau_{eff} = \frac{C}{G} \frac{1 + \beta}{1 + \beta + \mathcal{L}_{I}} \quad \text{thermal time constant,}$ accelerated by ETF



- The readout circuit has to cope with the initial sharp electrical rise whereas the pulse information is rather contained in the slower thermal response
  - For a large part of the project, the circuit was thus optimized to operate near critical damping ( $\tau_+ = \tau_-$ )



# Main TES noise contributors

- Phonon noise (aka thermal fluctuation noise)
  - ►  $S_{P_{TEN}} = 4k_B T_0^2 GF(T_0, T_{bath})$  (fluctuations across the thermal link to the bath)
- Johnson noise from the load resistor
- >  $S_{V_{IN,L}} = 4k_BT_0R_L$  (electrical fluctuations in the shunt resistor)
- TES Johnson noise
  - ►  $S_{V_{JN,TES}} = 4k_BT_0R_0\zeta(I_0)$  (electrical fluctuations in the TES resistance, but amplified by the TES non-linear response)
  - > Without any other fluctuation mechanism than the impact of the TES current sensitivity,  $\zeta(I_0) = 1 + 2\beta$  (I&H+00)
  - ► In practice, only really works at low  $\beta$  values. Most recent TESs for X-ray applications exhibit an "unexplained" noise in the TES electrical band:  $S_{VTES} = 4k_BT_0R_0(1 + 2\beta + M^2)$ 
    - Hard to distinguish from internal thermal fluctuation noise (thermal fluctuations inside the TES, e.g. due to more than one heat capacity, imperfect thermalization, ...). It has the same frequency response
    - Each white thermodynamical fluctuations is transformed by the TES system and gets a distinct frequency response





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The TES energy resolution is given by the integrated signal to noise density:

$$\Delta E_{FWHM} = \sqrt{8 \log(2)} \left( \int_0^\infty 4SNRD^2(f) df \right)^{-1/2}$$
  
$$= \sqrt{8 \log(2)} \left( \int_0^\infty \frac{4}{NEP^2(f)} df \right)^{-1/2}$$
 TES responsivity  
(= FT of normalized pulse)  
$$= \sqrt{8 \log(2)} \left( \int_0^\infty \frac{4|s_I(f)|^2}{S_{ITFN}(f) + S_{IJN,TES}(f) + S_{IJN,L}(f)} df \right)^{-1/2}$$

In the small signal limit, without amplifier noise and in strong ETF case:

$$\Delta E_{FWHM} \propto \sqrt{\frac{4k_B T_0^2 C \sqrt{\zeta(I_0)}}{\alpha}}$$

Simplified TES model at high ETF and (very) low inductance => Information limited by Johnson noise





# **Impact of different TES parameters on performances: Temperature**

$$\Delta E_{FWHM} \approx \sqrt{\frac{4k_B T_0^2 C \sqrt{\zeta(I_0)}}{\alpha} \sqrt{\frac{nF(T_0, T_{bath})}{1 - (T_{bath}/T_0)^n}}} \sim 0.5$$
 for X-IFU TES

• Usual « the colder the better » ( $4k_BT_0^2$  noise scaling shared by all thermal detectors)

- But need lower Tc and colder thermal bath (fluctuations become much stronger when trying to operate a TES close to the thermal bath temperature)
- For X-IFU,  $T_{bath} = 55 \ mK$  (including  $\Delta T$  in thermal link between last stage cooler and detector),  $T_c \approx 90 \ mK$
# Impact of different TES parameters on performances: Heat capacity

$$\Delta E_{FWHM} \approx \sqrt{\frac{4k_B T_0^2 C \sqrt{\zeta(I_0)}}{\alpha}} \sqrt{\frac{nF(T_0, T_{bath})}{1 - (T_{bath}/T_0)^n}}$$

- **\Rightarrow** Reducing C (but keeping  $\tau$  fixed):
  - Both sensitivity and transduced thermal noise increase linearly
  - > Translates into a bandwidth increase =>  $\sqrt{C}$  impact on  $\Delta E$
- Limited by intrinsic heat capacity of TES, heat capacity of absorber required for the X-ray stopping power and required energy range
  - X-IFU absorber made of a main bilayer of gold (good lateral thermal conductivity) and electroplated bismuth (better stopping power to heat capacity ratio)
    - Small gold cap on top of absorber to increase IR reflectivity and protect absorber (Bi) from humidity

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# Impact of different TES parameters on performances: alpha

$$\Delta E_{FWHM} \approx \sqrt{\frac{4k_B T_0^2 C \sqrt{\zeta(I_0)}}{\alpha}} \sqrt{\frac{nF(T_0, T_{bath})}{1 - (T_{bath}/T_0)^n}}$$

- Increasing alpha:
- > In strong ETF limit, manifests itself as a bandwidth increase =>  $1/\sqrt{\alpha}$  scaling



Simplified TES model at high ETF and (very) low inductance



# What about pixel speed/bath conductivity?

$$\Delta E_{FWHM} \approx \sqrt{\frac{4k_B T_0^2 C \sqrt{\zeta(I_0)}}{\alpha}} \sqrt{\frac{nF(T_0, T_{bath})}{1 - (T_{bath}/T_0)^n}}$$

- What happens when we change G? Warning: « rough arguments »
  - > The low frequency response density scales as 1/G
  - > Phonon noise amplitude scales as  $\sqrt{G}$
- **BUT...** to recover same TES setpoint temperature, need to decrease Joule power and thus current scales as  $1/\sqrt{G}$  => overall signal range scales as  $1/\sqrt{G}$ . In total:
- **Response scales as**  $1/\sqrt{G}$
- Phonon noise does not change
- Johnson noise does not change
- And bandwidth scales as *G*
- Overall  $SNR \times \sqrt{BW}$  is conserved =>  $\Delta E$  does not depend on *G* at first order

Note Fourier transform scaling: if  $x(t) \rightarrow X(w)$ , then  $x(at) \rightarrow X(w/a)/|a|$ 

Simplified TES model at high ETF and (very) low inductance => Information limited by Johnson noise





# **Optimal filtering (matched filter)**

- The energy of each X-ray pulse needs to be measured on board
  - At 2 cps/pxl and ~ 50 ms pulse records sampled on 16 bits at 130 kHz, the current useful X-IFU data rate ~ 300 Mbits/s...
- Most X-ray calorimeters use Optimal Filtering. It assumes (as an approximation both are wrong! But it works):
  - > Pulses are linear with energy :  $d(t) = E \times s(t)$
  - Noise is stationary: noise PSD N<sup>2</sup>(f) is sufficient to describe the full system noise
- \*  $\chi^2$  minimization problem in the Fourier Space:



 $\min\left(\sum \frac{|D(f) - E \times S(f)|^2}{|N^2(f)|}\right) \Rightarrow \hat{E} = \frac{\langle D(f)(S^*(f)/N^2(f))}{\sum |S(f)|^2/N^2(f)} \qquad \text{Weighting by Noise Spectral Density}$ Equivalent in time domain:  $\hat{E} = \frac{\left\langle d(t) \mathcal{FFT}^{-1}\left\{\frac{S(f)}{N^2(f)}\right\}\right)}{\sum |S(f)|^2/N^2(f)} \qquad \text{The time domain optimal filter}$ 



# **Optimal filtering (matched filter)**





# **Record length degradation**

- The 0Hz bin of the OF is usually removed (the filter becomes zero-summed)
  - Energy reconstruction becomes insensitive to offsets and in general less sensitive to low frequency noise (e.g. 1/f)
    - Not to low frequency gain variation though (e.g. temperature, bias, drifts)!
- Information is lost. The shorter the filter (record), the larger the degradation (Doriese+09):

$$\Delta E = \sqrt{8\log(2)} \left( \int_{1/2t_{rec}}^{\infty} 4SNRD^2(f)df \right)^{-1/2} = \frac{\lim_{t_{rec} \to \infty} \Delta E}{\sqrt{1 - 1/(2t_{rec}f_{eff})}}$$
  
with  $f_{eff} = \frac{\int_0^{\infty} SNRD^2(f)df}{\lim_{f \to 0} SNRD^2(f)}$ 

Slower (easier to readout) pixels lead to larger degradation









# **Count rate capability**

- X-ray photons arrival times are governed by Poisson statistics
- Depending on time separation, the quality of the energy reconstruction can be degraded
- Grading scheme used to identify high (and lower grade) resolution events. Creates branching ratios as a function of count rate
- An X-ray calorimeter count rate capability is thus usually defined as a fraction of high resolution events to be met at a given count rate
- This C.R. requirement, yields a maximum total time period  $\Delta t_{tot}$  that is permitted to process an X-ray event:

 $\Delta t_{tot} = 10 \ \tau_{fall} + t_{rec} \quad \blacktriangleleft$ 

For uniform count rate:  $P(HR) = P(no \ evt \ in \ \Delta t_{tot}) = e^{-c\Delta t_{tot}} \Rightarrow \Delta t_{tot} = -\frac{\log P(HR)}{\log P(HR)}$ 

For non-uniform count rate (e.g. point sources):

$$P(HR) = \sum_{pixels} f_i P(HR)_i = \sum_{pixels} f_i e^{-c_i \Delta t_{tot}} \Rightarrow \Delta t_{tot} = \cdots$$

Available time for pulse processing

Simplified secondaries exclusion





# What about X-IFU? On focus constraint

- Two main types of count rate capability requirements:
- Point source: 80% throughput at 1 mCrab, goal at 10 mCrab
- Extended source: 80% throughput on Perseus cluster, goal at CasA
- Without defocusing, in the original X-IFU concept (250µm pixel pitch), 1mCrab lead to ~ 42 cps on central pixel (also larger mirror)!
  - > 80% throughput criterion yields  $\Delta T_{tot} = 11 ms$
  - > Much tighter constraint than extended source goal (2.3 cps/pxl =>  $\Delta T_{tot} = 97ms$ )
  - Which pixel speed is required to make this work?
    - ► As a simplification,  $\tau_{fall} \approx \tau_{eff}$  and  $f_{eff}$  scaled together with respect to original « Large Pixel Array » pixel parameters (Smith+16)  $\Leftrightarrow$  scaling G without changing other parameters
    - > For a 2 eV intrinsic detector, need  $f_{eff}$  in the range of 1 kHz





# What about X-IFU? Introduction of defocusing

- Early in the program, ability to defocus the mirror was introduced (same mechanism as for instrument selection)
  - Spreads point source flux over several hundreds of pixels
  - $\label{eq:constraint} \begin{array}{ll} & 1 & mCrab \\ & \Delta T_{tot,ext\;goal} \end{array} \quad constraint \quad becomes \quad \Delta T_{tot} = 1.45\;s \gg \\ \end{array}$



- Instrument can simply be dimensioned on extended source count rate goal
  - Count rate constraint becomes much less stringent
    - Can "spend" associated margin to either:
      - Decrease the pixel speed => easier to read out (see later)
      - Increase the pixel size => lower number of required pixels to populate the field of view
        - Of course, still need to "properly" sample the PSF
        - Still reasonable imaging with pixel size ~ telescope HEW
        - Athena HEW = 5" (now 10")  $\Leftrightarrow$  291 µm with 12 m focal length









### **Evolution of TES array design over X-IFU study**



Preserved instrument performance for most of the study while significantly decreasing demands on the cryogenic system



# **Evolution of TES array design over X-IFU study**

<ul> <li>[1] Smith+16</li> <li>[2] No published reference, design principles similar to Nagayoshi+20</li> <li>[3] Smith+21</li> <li>[4] Wakeham+23</li> </ul>	SiN membrane Au stripes		Absorber supports 50 µm SIN → 50 µm membrane Nb bias leads Absorber outline	Silicon nitride membrane L Au Absorber stem attachment
Parameter	LPA1 [1] Start of phase A baseline	LPAa 75 AR:0.5 [2] End of phase A baseline (FDM)	LPA2.5a [3] First TDM baseline pixel	LPA 50x30 µm [4] Current optimization
Pixel pitch (µm)	249	275	317	317
$T_0$ (mK) – TES bias point T $^\circ$	90	88	89	89
$G_b$ (pW/K) – Thermal cond.	200	60	72	66
<i>C</i> (pJ/K) – Total heat capacity	0.8	1.0	0.73	0.73
$R_0 (m\Omega) - TES$ bias point resistance	1	7.1	0.97	1.4
$P_0$ (pW) – TES bias power	4.63	1.31	1.55	1.41
$\alpha \left( T/R \ \partial R/\partial T \right)$	75	532	619	619
$\beta(I/R \partial R/\partial I)$	1.25	8	22	22
M – Unexplained Johnson noise factor	0	1.73	1.88	1.88
L/L <sub>crit</sub>	1	1	0.64	0.74
$f_{eff}(Hz)$ – Effective frequency	962	421	477	435
$\Delta E_{smallsig}$ – Small sig resolution	1.7	1.68	1.72	1.72

COPS



# **TES optimization for AC bias**

- Original X-IFU concept used FDM multiplexing (see D. Prêle's presentation tomorrow) with TESs operated under AC bias
- Two main degradations observed wrt. DC bias \*\*
  - AC losses from Eddy currents generated in the normal metal structures
    - Prevents biasing low in the transition, spoils alpha, adds an 0 additional noise source
    - AC Josephson effects
      - Relatively large non-linear Josephson inductance in parallel with the TES resistance
      - Generates step-like structures in the TES transition 0
- Joint NASA/GSFC and SRON effort to optimize TES design
  - Remove normal metal features as much as possible (notably stripes)
  - Reduce non microstip area loop formed by TES and leads
  - Increase Rn, notably via increased aspect ratio
    - Higher Rn at the same power reduces the weak link effect
  - Reduction of pixel speed to mitigate carrier leakage









f = 1.48 MHz



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# THE X-IFU TDM DETECTION CHAIN: PRINCIPLE AND DIMENSIONING RULES





#### **SQUID** readout



- TES current is readout by SQUIDs (Superconducting Quantum Interference Devices)
  - > The current flows through a coil to create a magnetic field
  - The SQUID transduces the magnetic field into a measurable voltage signal



- ✤ A SQUID characteristic is very non-linear (almost a sine wave)
  - Its linear readout requires a feedback signal generated in the warm electronics
  - Operated in a so-called « Flux Locked Loop » aka FLL









*t<sub>row</sub>*: time between two rows/pixels (160 ns for X-IFU)

 t<sub>frame</sub> = N<sub>MUX</sub> × t<sub>row</sub>: time between two reading of the same row/pixel (=48x160 ns =7.68 µs for X-IFU)

 X-IFU operated in 4 independent (but synchronized) 8x48 TDM schemes















Credit: NASA/GSFC

 Array of 1504 TES detectors (NASA/GSFC)









Durkin+23

#### Fully differential MUX SQUID chip (NIST)

- Two level FAS to limit wiring for RA signals
- Also contains shunt and Nyquist L for TES











 Differential AMP SQUID (VTT)

- Most of the cold gain (~600 Ω transresistance)
- Located at 2K





#### **The overall X-IFU readout architecture**



- Warm Front-End Electronics (APC)
  - Located in PLC basement at 300 K
  - Differential LNA and current sources for cold electronics bias implemented in dedicated ASIC
  - Implements buffers on DRE dynamic lines



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#### The overall X-IFU readout architecture



#### Digital Readout Electronics (IRAP)

- Digitalization and demultiplexing of TDM signals
- Feedbacks generation for both SQUID stages
  - Not a closed loop for AMP SQUIDs
  - Regular pattern to compensate varying offset from one MUX SQUID to another
- Row address signals generation
- On-board processing









WFEE 3-4

Number of channels	32	
MUX	48	
Pixels	1504 (+32)	
FoV	3.98' ~ 4'	
RA signals	6 + 8	
FPA T0 panels	4	
FPA T2 panels	2	
Science harness	2	
Wires per science harness	~ 500 (with	
	isolation pairs)	
WFEE physical box	2	
WFEE physical box WFEE functional units	2 4	
WFEE physical box WFEE functional units Columns per ASIC	2 4 2	
WFEE physical box WFEE functional units Columns per ASIC ASICs per WFEE unit	2 4 2 4 2 4	
WFEE physical box WFEE functional units Columns per ASIC ASICs per WFEE unit DRE physical boxes	2 4 2 2 4 4 4 4	
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WFEE physical boxWFEE functional unitsColumns per ASICASICs per WFEE unitDRE physical boxesDRE functional unitsColumns per FPGA	1solation pairs) 2 4 2 2 4 4 4 4 4 4 4	



# **TDM dimensioning rules: FLL operation**

- FLL algorithm:
  - $\succ I_{FB}[n] = I_{FB}[n-1] + K_{I}(I_{TES}[n] I_{FB}[n-1]) + I_{offset}$
- Feedback signal generated at the previous frame
- TES signal change between two frames should not exceed the SQUID monotonous/linear range for the most energetic event in the band (12 keV)
  - In practice, pulse slew rate decay makes it so that one could tolerate an overshoot on the pulse onset without loosing stability
  - Not used for now in dimensioning

$$\Delta \phi_{max} = \frac{dI}{dt} \bigg|_{max} M_{in} N_{MUX} t_{row} \le \phi_{range} \approx 0.3 \ \phi_0$$



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# **TDM dimensioning rules: aliasing**

- Each pixel is only read out every N<sub>MUX</sub> row
- Mismatch between the open-loop bandwidth...
  - > Needs to allow the system to respond on the timescale of a row:  $f_{0L} > 1/t_{row}$  (6.25 MHz for X-IFU)
- … and the readout frequency

$$\succ f_{sample} = \frac{1}{t_{frame}} = \frac{1}{N_{MUX}t_{row}}$$
 (130 kHz for X-IFU)

Readout noise gets aliased 2N<sub>MUX</sub> times:

 $M_{in} \sqrt{S_{I_{MUX}}} = \sqrt{\pi N_{MUX} S_{\phi}}$ 

Assumes open loop has an equivalent 1st order roll-off at  $1/t_{row} => factor \pi/2$ 

with  $\sqrt{S_{I_{MUX}}}$  multiplexed readout noise referred to TES (unit: A/VHz)  $M_{in}$  mutual inductance coupling  $I_{TES}$  to MUX SQUID (unit: H)

- $\sqrt{S_{\phi}}$  non-multiplexed readout noise, referred to MUX SQUID (unit:  $\phi_0$ /VHz)
- $N_{MUX}$  multiplexing factor = number of TDM rows







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# **Dimensioning exercise (small signal model)**

#### **Question: which TDM performance is required**

## as a function of pixel speed?

#### Target performance:

- > 2 eV (for a 3 eV instrument) at detector + readout chain level
  - Actual pixel + readout allocation ~ 2.25 eV, but would lead to unrealistic constraints. 2 eV taken to somewhat compensate the fact that the small signal model underpredicts resolutions
- Target throughput before rescope: 90% at 2.5 cps/pxl (extended source count rate goal, setting aside 10 % for crosstalk rejection see later):  $\Delta T_{tot} = 42ms$
- Post rescope focal plane population: 4' equivalent diameter, populated by 317 µm pitch pixels in 32 TDM columns => MUX 47(+1: one pure resistive pixel per column to monitor gain drifts)

#### First done with LPA1 pixel (phase A design). Then pixel speed modified only via G

- > Setpoint current rescaled to keep same working temperature
- L modified to remain at critical damping
- All other parameters kept the same

#### Disclaimers

- Use simplified dimensioning rules here for ease of description and small signal model
- Will highlight after how TES non-linearity (among others) modifies the dimensioning for the latest X-IFU design point





$$t_{rec} = \Delta t_{tot} - 10\tau_{eff} = 37.8 ms$$

$$\Delta E_{FWHM} = \sqrt{8 \log(2)} \left( \int_{1/2 t_{rec}}^{\infty} \frac{4|s_I(f)|^2}{S_{ITFN}(f) + S_{I_{JN,TES}}(f) + S_{I_{JN,L}}(f) + S_{I_{MUX}}} df \right)^{-1/2} \le 2 eV$$

$$=> \text{Gives a limit on the maximum allowed (multiplexed) readout noise at TES level } S_{I_{MUX},max}$$

$$S_{I_{MUX},max} = 43.9 \text{pA}/\sqrt{Hz}$$

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# **Optimization of SQUID input mutual inductance**

- Two main TDM dimensioning relations
  - > SQUID linear flux range :  $\frac{dI}{dt}\Big|_{max} M_{in} N_{MUX} t_{row} \le 0.3 \phi_0$
- > Maximum TES ref'd noise to reach performance :  $\frac{\sqrt{\pi N_{MUX}S_{\phi}}}{M_{in}} \leq \sqrt{S_{I_{MUX,max}}}$
- Min is a free parameter that can (needs to) be optimized for each array model
  - > Min ↑
    - Readout noise ↓ wrt. TES noise
    - Required linear flux range ↑
  - ➤ Min↓
    - Readout noise ↑ wrt. TES noise
    - $\circ \quad \text{Required linear flux range} \downarrow$





## **TDM allowed parameter space**

- Two main TDM dimensioning relations
  - > SQUID linear flux range :  $\frac{dI}{dt}\Big|_{max} M_{in} N_{MUX} t_{row} \le 0.3 \phi_0$
- > Maximum TES ref'd noise to reach performance :  $\frac{\sqrt{\pi N_{MUX}S_{\phi}}}{M_{in}} \leq \sqrt{S_{I_{MUX,max}}}$
- **\therefore** Combine with Min (free parameter): provides the allowed parameter space for the multiplexing scheme ( $t_{row}$  and  $N_{MUX}$ )







#### **TDM allowed parameter space**







### **TDM allowed parameter space**



- Would not be able to implement the X-IFU with the original pixel design by almost a factor 2 on either (or a combination)
  - Noise amplitude spectral density
  - Squid dynamic
  - Row time
- But... We know that original pixel concept was largely too fast. What happens when we slow down the pixel?



# Dimensioning exercise (small signal model): Slowing down the pixel

#### **Reminder: simple scaling of G as a toy exercise**



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### Dimensioning exercise (small signal model): Slowing down the pixel





#### **Impact of non-linearity**

- Whereas useful to describe trends, the small signal is certainly not reliable to dimension an X-IFU like readout
- Especially true for stripe-less high alpha TESs as those foreseen for X-IFU
- Intrinsic energy resolution is largely underestimated
  - Need to rather dimension the impact of readout noise as a quadrature term on top of the pixel intrinsic resolution
  - Also allows having a less ambiguous share of requirements between subsystems...
- Maximum slew rate is overestimated by a factor of a few
- Note that operation slightly below L<sub>crit</sub> may also be favored when taking into account nonlinearity



Courtesy data from same pixels as Smith+21

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# **Application to latest target pixel design**

- Current target pixel is an adaptation of the 8x32
   TDM demo pixel
- Aim at higher aspect ratios (Wakeham+23) to
- Further decrease the pixel speed
- Reduce the pixel sensitivity to magnetic field (among the main source of gain drifts)
- Latest measurements predict
  - A slew rate at 12 keV of ~ 0.17 A/s
  - An acceptable readout noise of 25 pA/√Hz at TES level for the budgeted ~ 0.8 eV quadratic degradation from readout
  - $\blacktriangleright = DSR = 6.8 \times 10^9 \sqrt{Hz}$





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## GOING BEYOND SIGNAL TO NOISE RATIO OPTIMIZATION





.....



#### TDM bandwidths optimization

- At each TDM row, a settling transient occurs when the system switches from one signal level (pixel) to the other
- Its shape depends on:
- The physics of the MUX SQUID (including FAS)
- The relative timing between the different TDM signals (FB and RA) and their filtering by the system, therefore on the system bandwidths
  - It creates:
    - Crosstalk
    - Gain scale distortions





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#### TDM bandwidths optimization: crosstalk vs. bandwidth

- The amplitude of the end of row residual depends on the signal from the previous row
- Different frequency cut-offs lead to different crosstalk types. At first order:
- A cut-off in the pre-modulation bandwidth will slow the settling of the FB signal (~ pulse) => creates proportional crosstalk
- A cut-off in the post-modulation bandwidth will slow the settling of the error signal (~ pulse derivative) => creates derivative crosstalk
  - Simulations of this effect for X-IFU showed that most bandwidths need to be in the range of at least 10 MHz, i.e. quite higher than the "canonical" 6.25 MHz given by 1/trow



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#### **TDM** bandwidths optimization: gain scale effect

- \* As the settling shape is non-linear (notably because full FB signal spans several  $\phi_0$ ), the end of row residual, even in absence of a pulse in the previous row, depends on the signal level
- This creates gain scale distortions that can be significant with respect to the "main gain scale", which is dominated by the TES response
- Simulations show that BWs need to be more open than to respect the previous crosstalk criterion
  - SQUID-level model shows very good match to measurements
- End-to-end measurement performed using linear pulse injections show similar distortions as simulations (quantitative comparison in progress)
- Will increase the Noise Equivalent Bandwidth of the readout noise
  - > In practice, combination of bandwidths leads to similar NEB as  $\frac{\pi}{2} f_{row}$







- Crosstalk = small fraction of a perpetrator pulse leaking onto a victim pulse
- Creates energy offsets as a function of time separation and the match between the crosstalk pulse shape and the optimal filter
- E.g. proportional crosstalk will create a much larger effect than derivative crosstalk
- Main sources of crosstalk in X-IFU
  - **TDM transients**
  - Thermal crosstalk (largely mitigated by backside gold coating of TES wafer few µm thickness)
  - Inductive coupling (in MUX SQUID chip, harnesses, between TES wires, ..)
- Advantage = it is a fully deterministic process
  - Can create exclusion rules to reject from the science the most impacted events
  - In principle, could try to correct for it (even if not very accurately)
- Overall, in pre-rescope design, expected ~ 10% additional (on top of grading) throughput loss from crosstalk







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#### **Environmental drifts**

Parameter	Assumed SSR sensitivity requirements	$\Delta E = 0.1 \text{ eV}$ equivalent
Temperature	$\partial E/\partial T = 150 \text{ meV}/\mu\text{K}$	0.28 μK rms
Bias voltage	$\partial E/(\partial V/V) = 15 \text{ meV/ppm}$	2.8 ppm rms
Magnetic field	$\partial E / \partial B = 3 \text{ eV/nT}$	15 pT rms

- TESs are very performant but sensitive detectors
- Drifts of the environment will modulate the X-ray response
  - Degrades the resolution on short time scales
  - Creates non-linear modifications of the instrument gain scale on longer timescales
- Requires
  - ~ 1µK rms thermal stability
  - 10<sup>5</sup> 10<sup>6</sup> shielding towards external magnetic fields
  - Electronics gain stability in the TES bias and FB lines of the order of few 10 ppm/K









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FPA DM1.0, credit SRON



#### Long timescale drift: gain drift correction algorithms

- Similar to XRISM, the X-IFU will implement a modulated X-ray source to track the pixels' gains in a set of fiducial lines
  - > Will shine for 1-1.5 s every few tens of s
- Due to the non-linear nature of the TES gain scale (and of its modification by an external parameter), linear correction is not sufficient
  - Rather interpolate in a non-linear manner between calibrated gain scales
  - Higher order corrections also being tested using the pre-pulse baseline as an additional fiducial





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#### Long timescale drift: gain drift correction algorithms performance







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# WHAT IS NEXT FOR X-IFU?







#### Long term X-IFU timeline





- Detection chain « early verification » (ongoing, up to 2025)
  - Validation of differential detection chain design at warm and then coupled to a cold differential focal plane
- Cold harness breadboarding (ongoing)
  - Crosstalk characterization
- Thermalization performance
- Connectorization
  - > ...
- FPA DM1.1 (build 2024, test up to ~ mid 2026)
  - > At scale FPA demonstration model with differential TDM cold electronics
  - Re-use of FPA DM1.0 structure
  - High fidelity prototypes of cold electronics
- 4K Core DM (AIT and test mid 2026 to mid 2027)
  - Demonstration model of 4K core, with ADR DM and FPADM1.1
  - Objectives
    - o Validation of ability to integrate full 4K core
    - o Compatibility between ADR and FPA
    - o 4K core thermal regulation optimization and thermal model correlation
    - $\circ$  Assessment of sensitivity to external perturbations (microvibrations, thermal stability, mag field TBC)







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Geoffray+24



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    - $\circ$   $% \ensuremath{\mathsf{Assessment}}$  of sensitivity to external perturbations (microvibrations, thermal stability, mag field TBC)



Credit: Absolut System





#### Conclusion



- Athena and X-IFU restarting at full speed following mission rescope
- Several similarities with other low temperature detection experiment
- Happy to discuss this with you further
  - Good ideas welcome!
- Numerous demonstration activities on-going/planned in the coming years
  - > Major step forward for the instrument
  - Excited to share these results soon with the community

## BIBLIOGRAPHY





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#### **Bibliography**

- **Reviews for TES detectors and associated readout chains particularly relevant for X-ray astronomy** 
  - McCammon D., 2005, cpd..book, 1. doi:10.1007/10933596\\_1
  - Irwin K. D., Hilton G. C., 2005, cpd..book, 63. doi:10.1007/10933596\\_3
  - **Gottardi, L., Nagayashi, K., 2021, Appl. Sci., 11, 3793. doi:10.3390/app11093793**
  - Gottardi L., Smith S., 2022, hxga.book, 103. doi:10.1007/978-981-16-4544-0\\_22-1
- **t** Latest full X-IFU instrument description paper (pre-rescope configuration, new architecture to be published this year)
  - Barret D., Albouys V., Herder J.-W. den ., Piro L., Cappi M., Huovelin J., Kelley R., et al., 2023, ExA, 55, 373. doi:10.1007/s10686-022-09880-7



### 

#### **Bibliography**

- Other articles cited in this presentation (in order of citation)
  - Wakeham+23: Wakeham N. A., Adams J. S., Bandler S. R., Beaumont S., Chervenak J. A., Cumbee R. S., Finkbeiner F. M., et al., 2023, ITAS, 33, 3253067. doi:10.1109/TASC.2023.3253067
- Lindeman+00: Lindeman, M. A., "Microcalorimetry and the transition-edge sensor", PhDT, 2000
- Gottardi+21: Gottardi L., de Wit M., Taralli E., Nagayashi K., Kozorezov A., 2021, PhRvL, 126, 217001. doi:10.1103/PhysRevLett.126.217001
- Doriese+09: Doriese, W. B., "Optimal filtering, record length, and count rate in transition-edge-sensor microcalorimeters", in *The Thirteenth International Workshop on Low Temperature Detectors LTD13*, 2009, vol. 1185, no. 1, pp. 450–453. doi:10.1063/1.3292375
- Smith+16: Smith S. J., Adams J. S., Bandler S. R., Betancourt-Martinez G. L., Chervenak J. A., Chiao M. P., Eckart M. E., et al., 2016, SPIE, 9905, 99052H. doi:10.1117/12.2231749
- Nagayoshi+20: Nagayoshi K., Ridder M. L., Bruijn M. P., Gottardi L., Taralli E., Khosropanah P., Akamatsu H., et al., 2020, JLTP, 199, 943. doi:10.1007/s10909-019-02282-8
- Smith+21: Smith S. J., Adams J. S., Bandler S. R., Beaumont S., Chervenak J. A., Denison E. V., Doriese W. B., et al., 2021, ITAS, 31, 3061918. doi:10.1109/TASC.2021.3061918
- Gottardi+17: Gottardi L., Akamatsu H., van der Kuur J., Smith S. J., Kozorezov A., Chervenak J., 2017, ITAS, 27, 2655500.
- de Wit+20: de Wit M., Gottardi L., Taralli E., Nagayoshi K., Ridder M. L., Akamatsu H., Bruijn M. P., et al., 2020, JAP, 128, 224501. doi:10.1063/5.0029669
- Sonzalez+24: Gonzalez M., Parot Y., Kirsch C., Cucchetti E., Peille P., Murat D., et al. 2024, JLTP, in press
- Durkin+23: Durkin M., Backhaus S., Bandler S. R., Chervenak J. A., Denison E. V., Doriese W. B., Gard J. D., et al., 2023, ITAS, 33, 3264175. doi:10.1109/TASC.2023.3264175





#### **Bibliography**

- Other articles cited in this presentation (in order of citation)
  - Doriese+12: Doriese W. B., Alpert B. K., Fowler J. W., Hilton G. C., Hojem A. S., Irwin K. D., Reintsema C. D., et al., 2012, JLTP, 167, 595. doi:10.1007/s10909-012-0509-7
- Durkin+21: Durkin M., Adams J. S., Bandler S. R., Chervenak J. A., Denison E. V., Doriese W. B., Duff S. M., et al., 2021, ITAS, 31, 3065279. doi:10.1109/TASC.2021.3065279
- Kirsch+22: Kirsch C., Lorenz M., Peille P., Dauser T., Ceballos M. T., Cobo B., Merino-Alonso P. E., et al., 2022, JLTP, 209, 988. doi:10.1007/s10909-022-02700-4
- Peille+18: Peille P., Dauser T., Kirsch C., den Hartog R., Cucchetti E., Wilms J., Barret D., et al., 2018, JLTP, 193, 940. doi:10.1007/s10909-018-1964-6
- Smith+23: Smith S. J., Witthoeft M. C., Adams J. S., Bandler S. R., Beaumont S., Chervenak J. A., Cumbee R. S., et al., 2023, ITAS, 33, 3258908. doi:10.1109/TASC.2023.3258908
- Geoffray+24: Geoffray H., Adams J., Fossecave H., Cucchetti E., Bellouard E., Peille P., Daniel C., et al. 2024, JLTP, in press