

Neutrino Oscillation Computational Aspects

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APC Laboratory

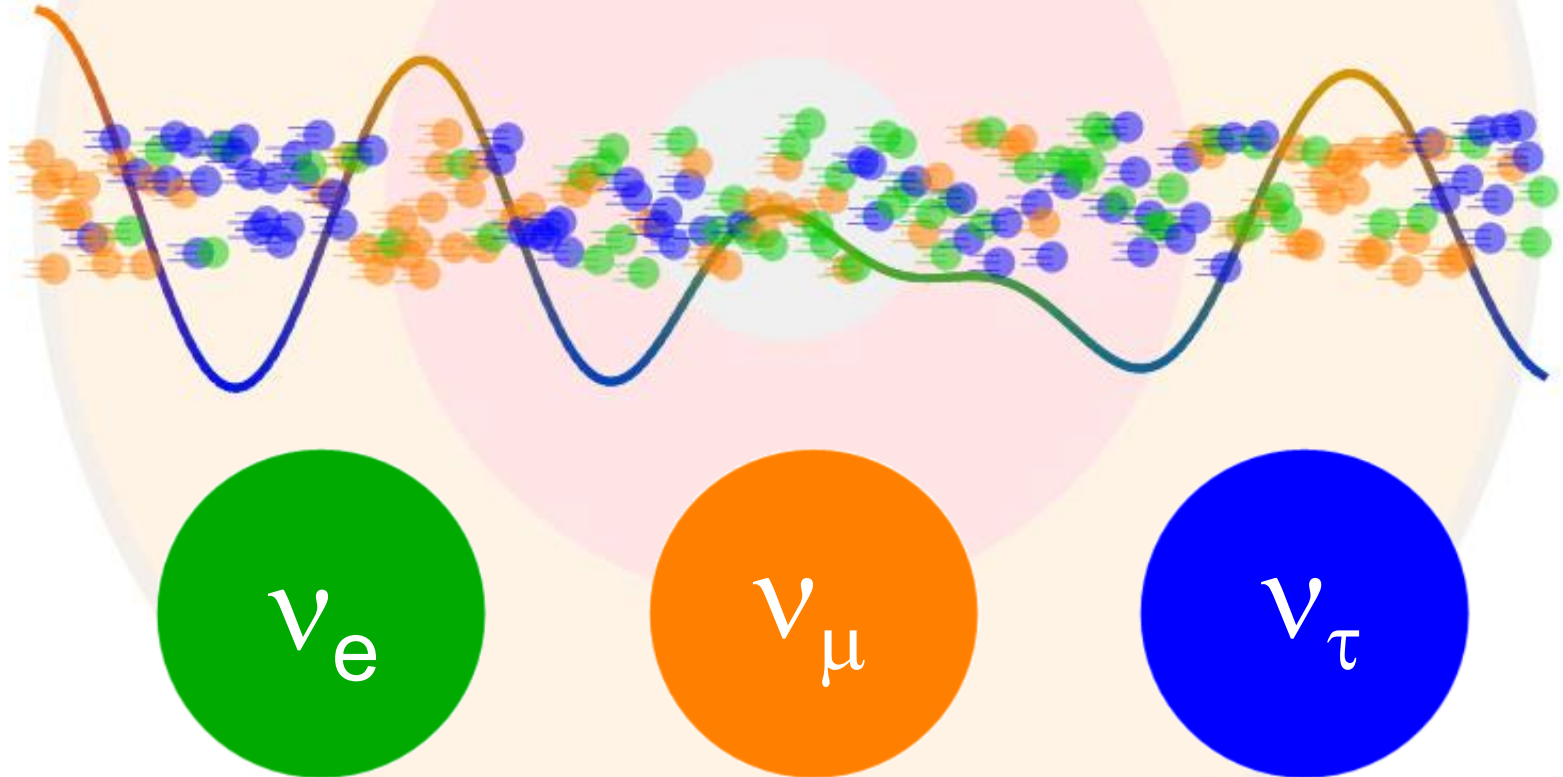
15 November 2023



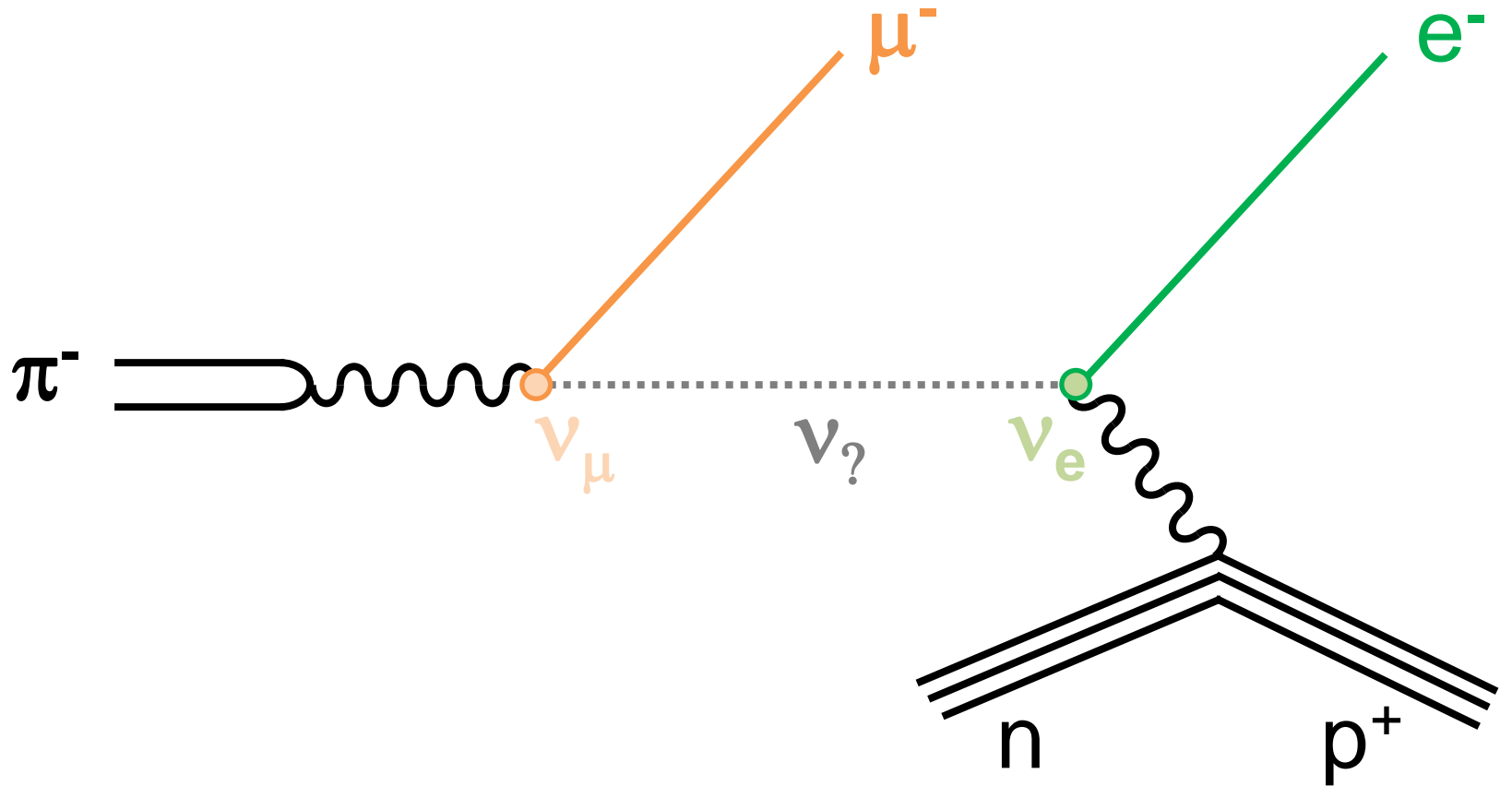
Neutrino Oscillations

- Neutrinos are created in a superposition of mass states
- Time evolution generates flavour oscillations

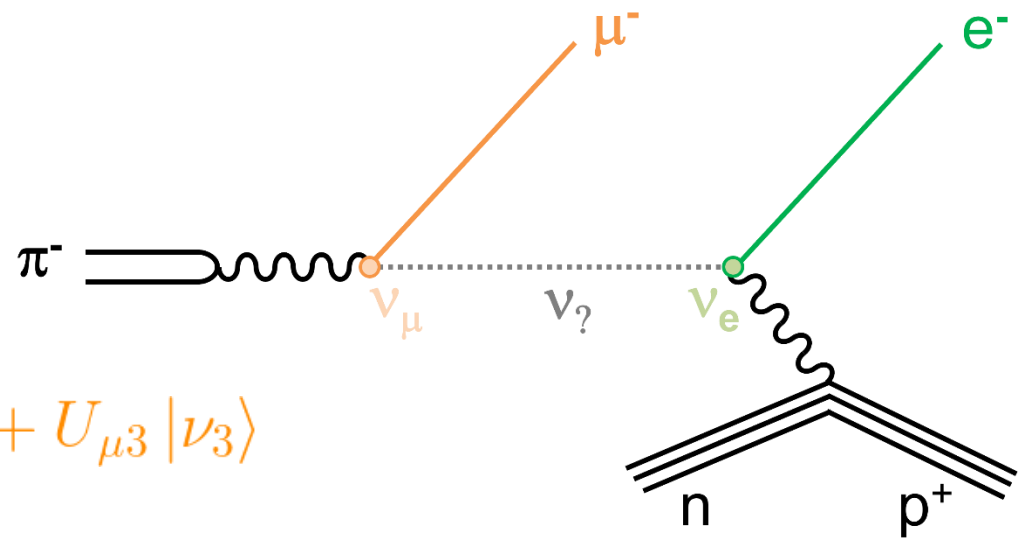
Quantum Mechanics



Neutrino Oscillations



Neutrino Oscillations



$$|\nu_\mu\rangle = U_{\mu 1} |\nu_1\rangle + U_{\mu 2} |\nu_2\rangle + U_{\mu 3} |\nu_3\rangle$$

$$|\nu(t)\rangle = U_{\mu 1} e^{-iE_1 t} |\nu_1\rangle + U_{\mu 2} e^{-iE_2 t} |\nu_2\rangle + U_{\mu 3} e^{-iE_3 t} |\nu_3\rangle$$

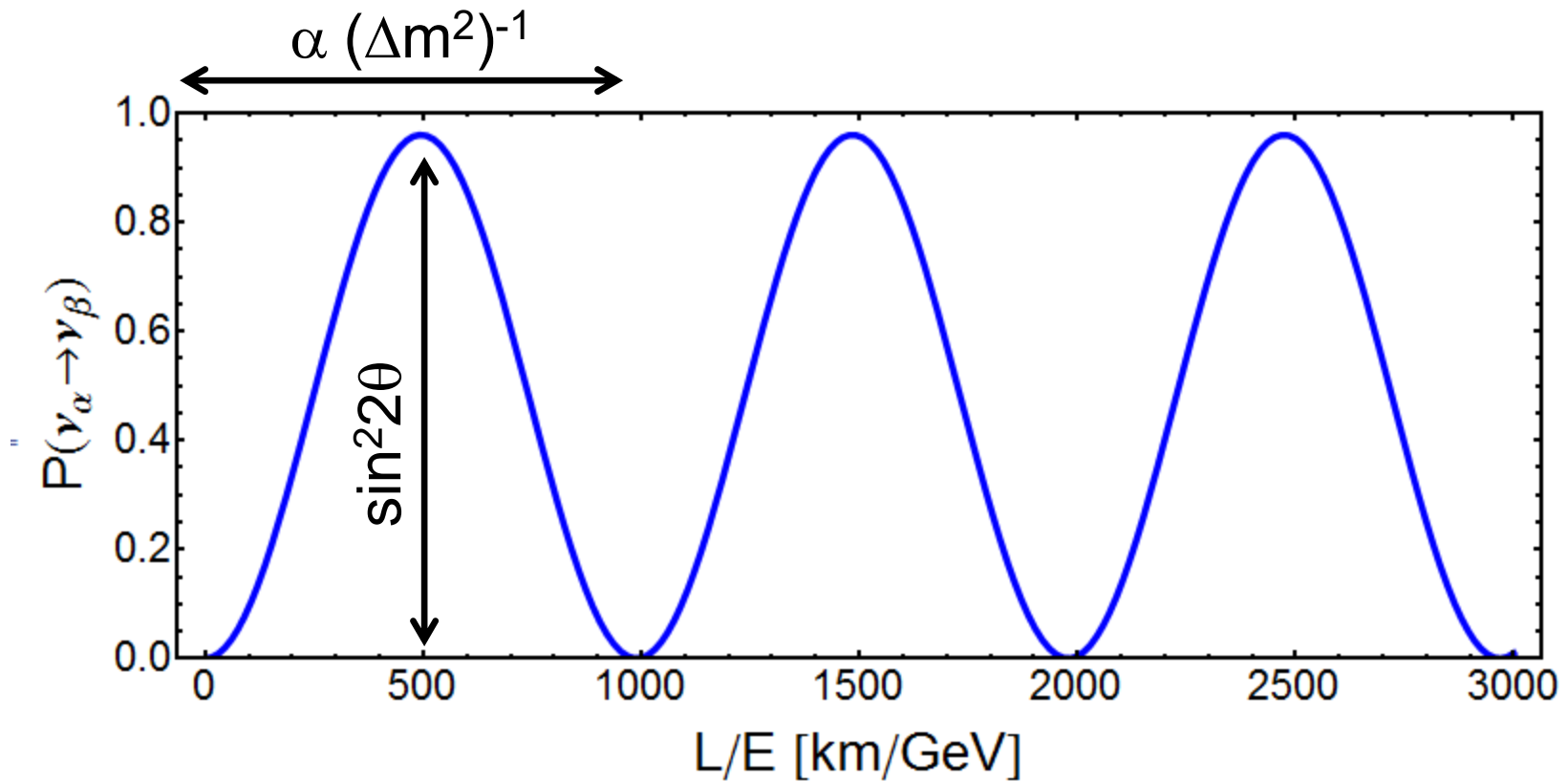
$$P_{\mu e} = |\langle \nu_e | \nu(t) \rangle|^2 = |U_{e 1}^* U_{\mu 1} e^{-iE_1 t} + U_{e 2}^* U_{\mu 2} e^{-iE_2 t} + U_{e 3}^* U_{\mu 3} e^{-iE_3 t}|^2$$

$$E_i \approx E + \frac{m_i^2}{2E} \quad \longrightarrow \quad P_{\alpha \rightarrow \beta} = \delta_{\alpha\beta} - 4 \sum_{j>k} \mathcal{R}_e \left\{ U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* \right\} \sin^2 \left(\frac{\Delta_{jk} m^2 L}{4E} \right) + 2 \sum_{j>k} \mathcal{I}_m \left\{ U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* \right\} \sin \left(\frac{\Delta_{jk} m^2 L}{2E} \right),$$

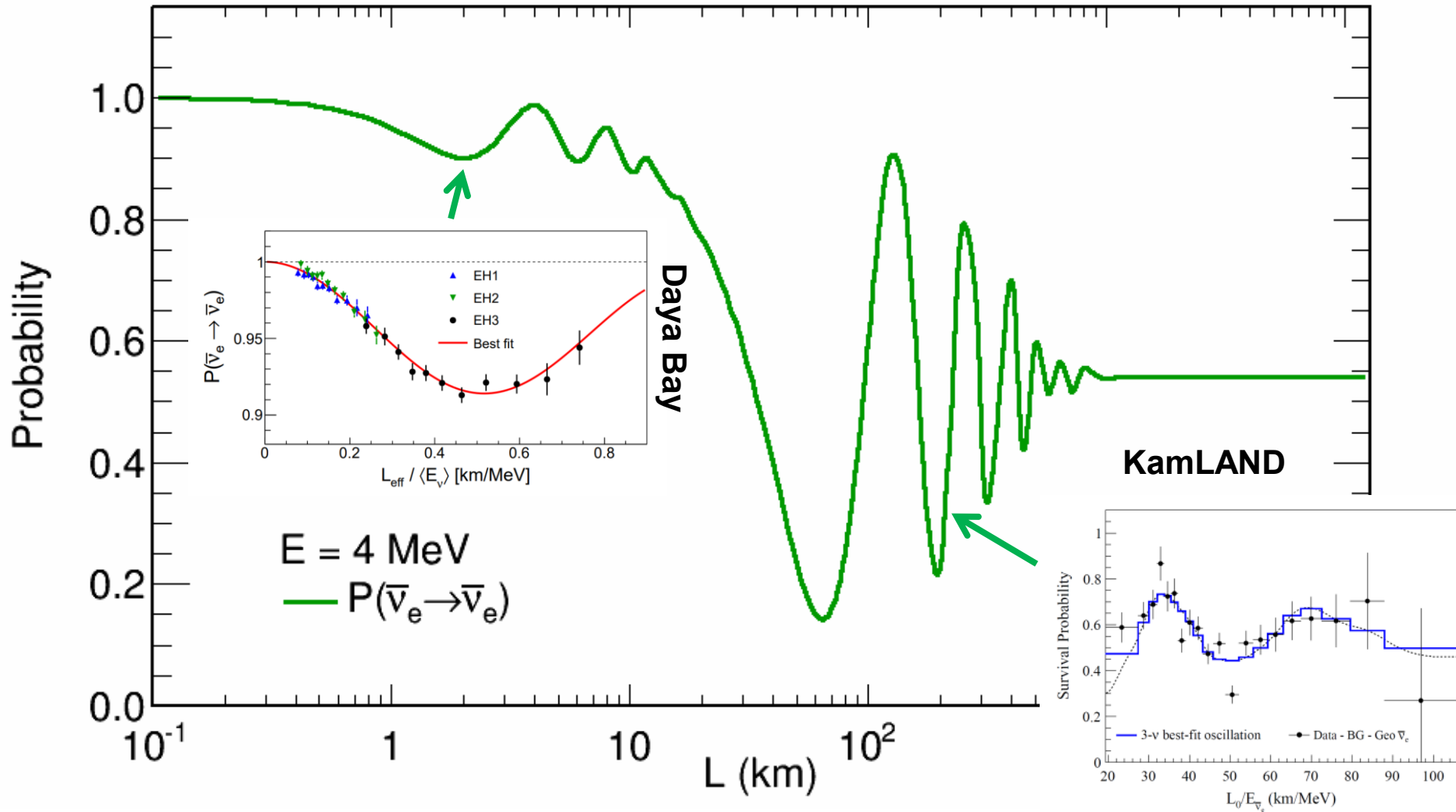
$$t \approx L$$

Neutrino Oscillations

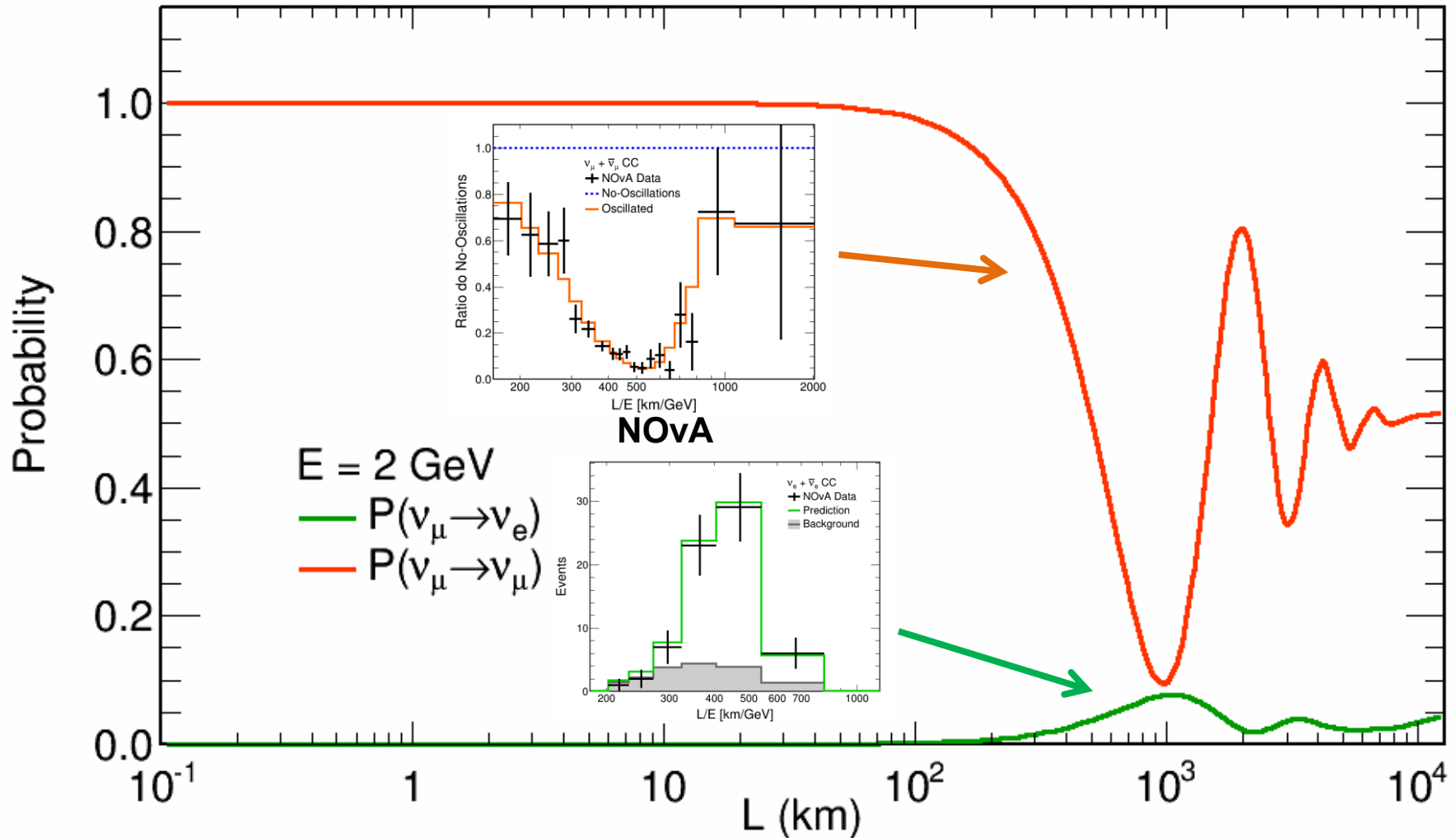
$$P(\nu_\alpha \rightarrow \nu_\beta) \approx \sin^2 2\theta \times \sin^2 \left(1.27 \times \Delta m^2 [\text{eV}^2] \times L/E [\text{km/GeV}] \right)$$



The Data: Reactor Neutrinos

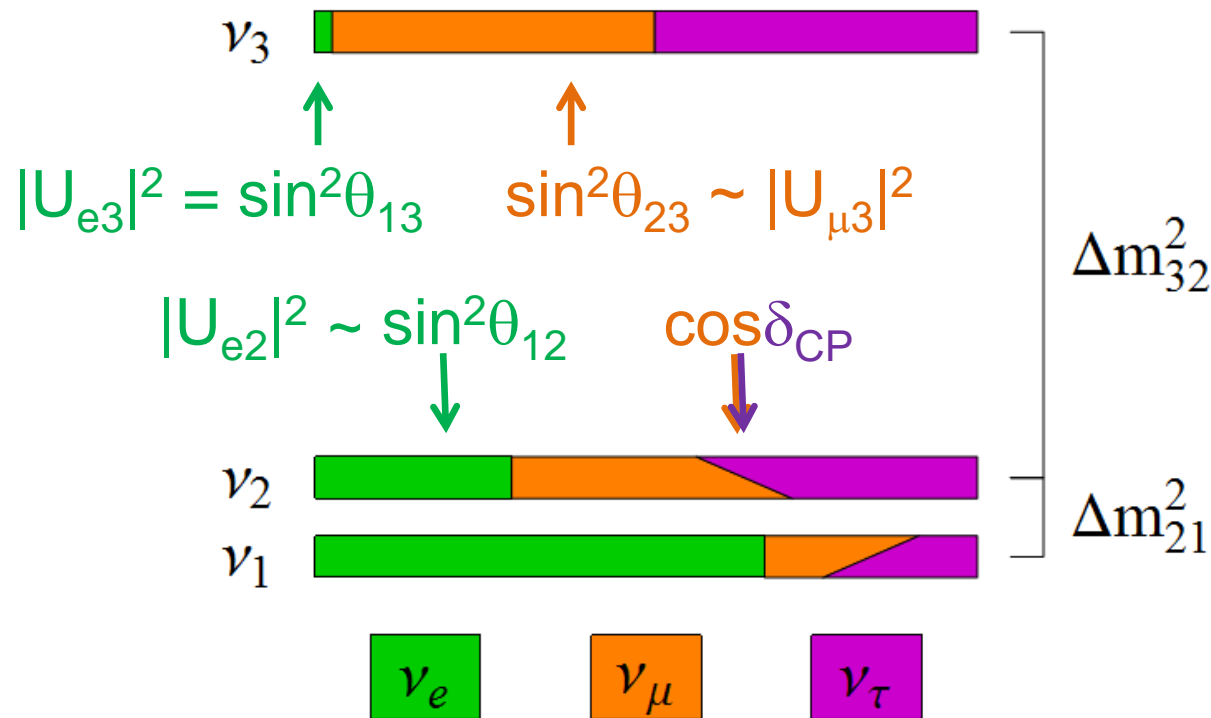


The Data: Accelerator Neutrinos



Neutrino Oscillations

- There are 3 neutrinos, so things are a bit more complicated
- Two independent differences in mass-squared (Δm_{21}^2 , Δm_{32}^2)
- 3 mixing angles (θ_{12} , θ_{13} , θ_{23}) and 1 CPV phase δ_{CP}



Quantum Evolution

Schrödinger: $i \frac{\partial}{\partial t} \mathcal{U} = H \mathcal{U}$

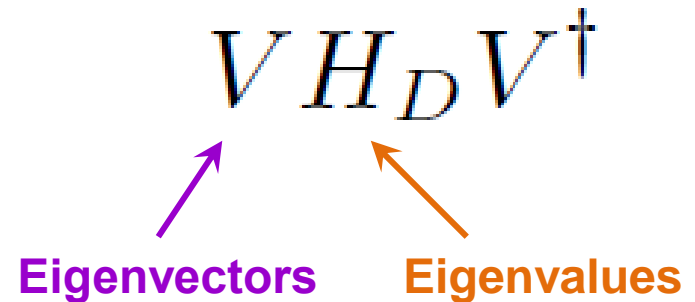
$$\mathcal{P}_{\alpha \rightarrow \beta} = |\langle \beta | \mathcal{U}(t) | \alpha \rangle|^2$$

Time-independent H :

$$\mathcal{U}(t) = e^{-iHt}$$

$$H = V H_D V^\dagger \quad \leftarrow \text{Main problem}$$

$$\mathcal{U}(t) = V e^{-iH_D t} V^\dagger \quad \leftarrow \text{Easy to compute}$$



Neutrino Hamiltonian in Vacuum

PMNS Matrix = Vacuum Eigenvectors

$$U = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{23} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Mass basis

$$H_0 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} \end{bmatrix}$$

Flavour basis

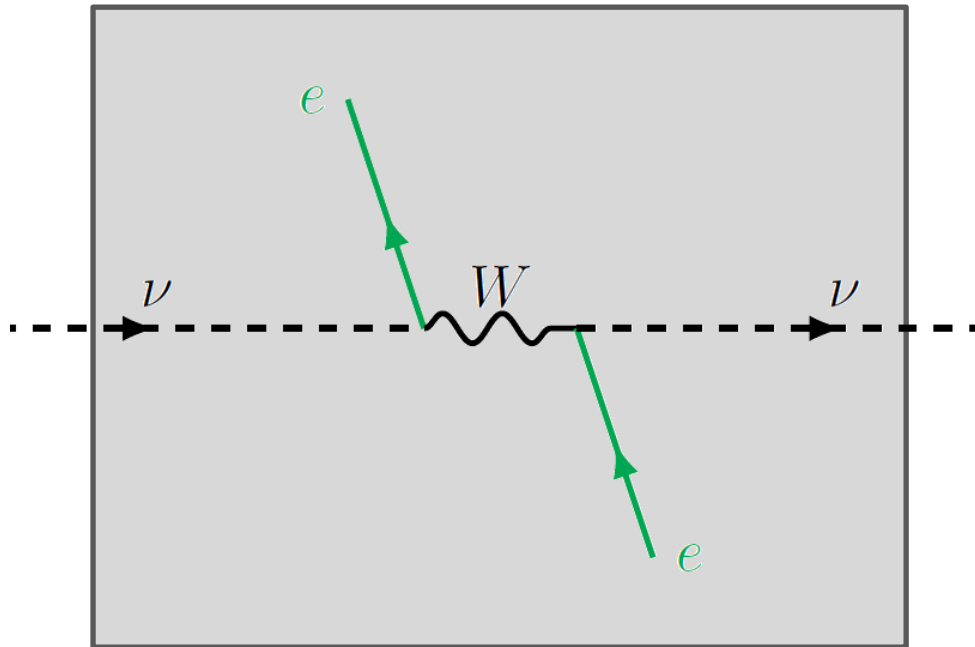
$$H = UH_0U^\dagger$$



Easy to solve

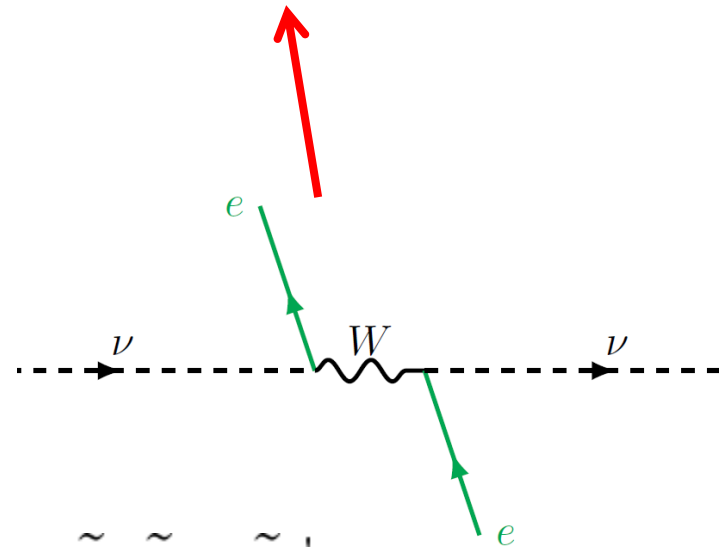
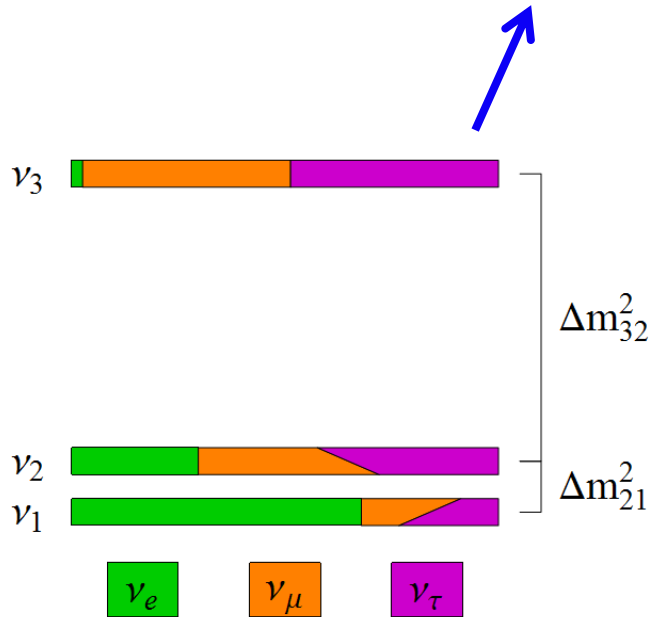
Already know diagonalisation
(Eigenvectors and Eigenvalues)

Neutrinos in Matter



Matter Effects

$$H_{eff} = U \underbrace{\begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} \end{bmatrix}}_{H_0} U^\dagger + V_e \underbrace{\begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}}_V$$

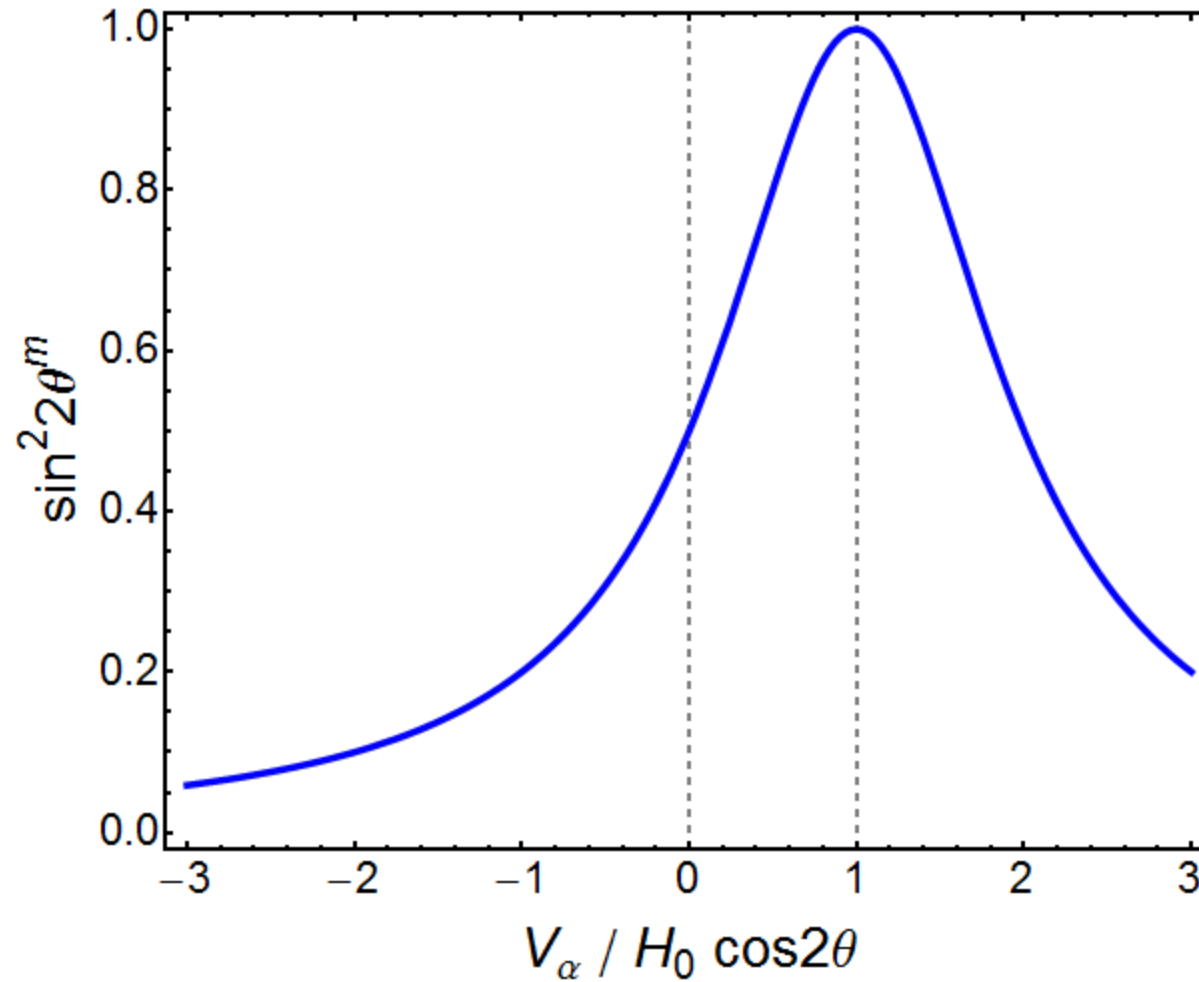


$$H_{eff} = \tilde{U} \tilde{H}_D \tilde{U}^\dagger$$

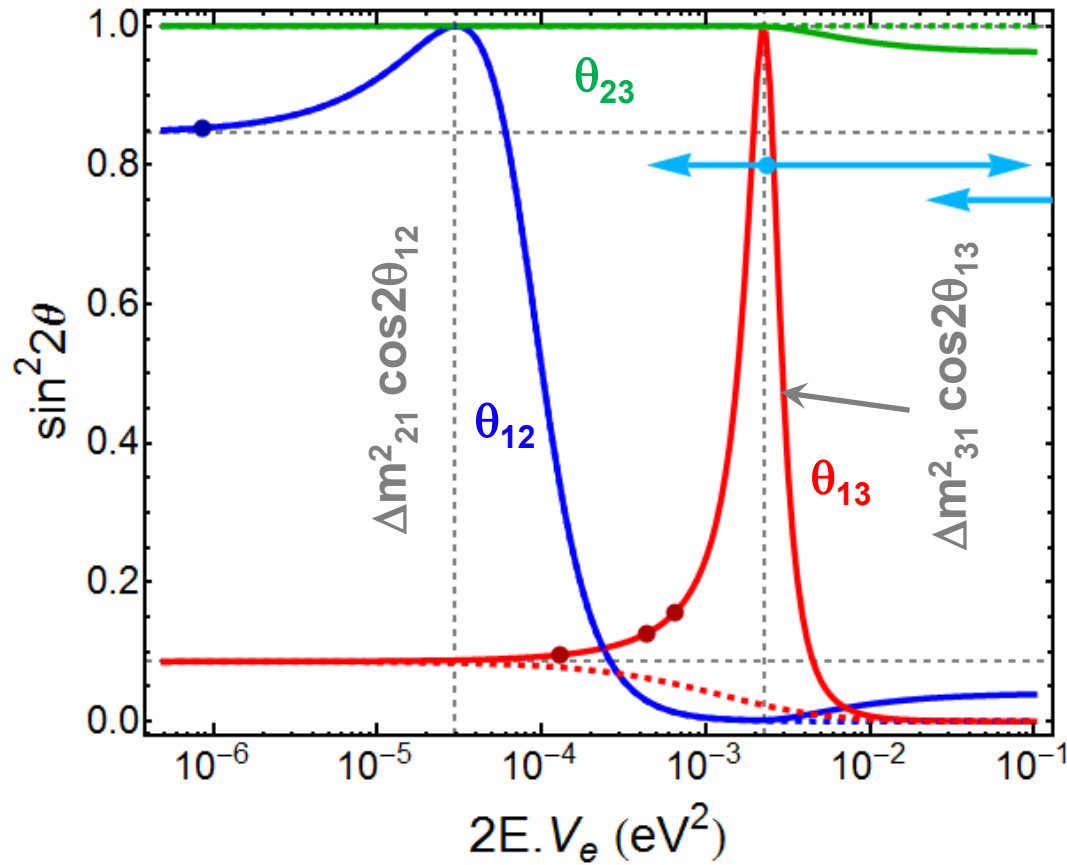
Effective Mixing

Effective Masses

Resonance

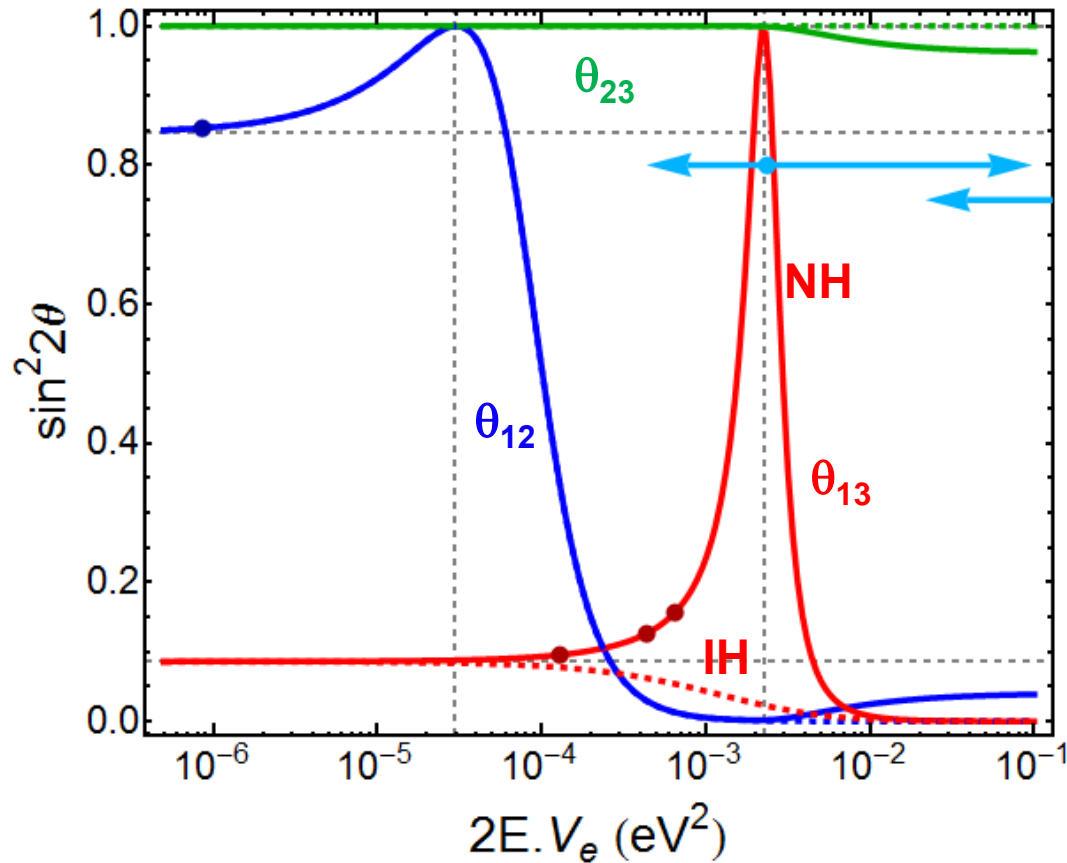


Resonances



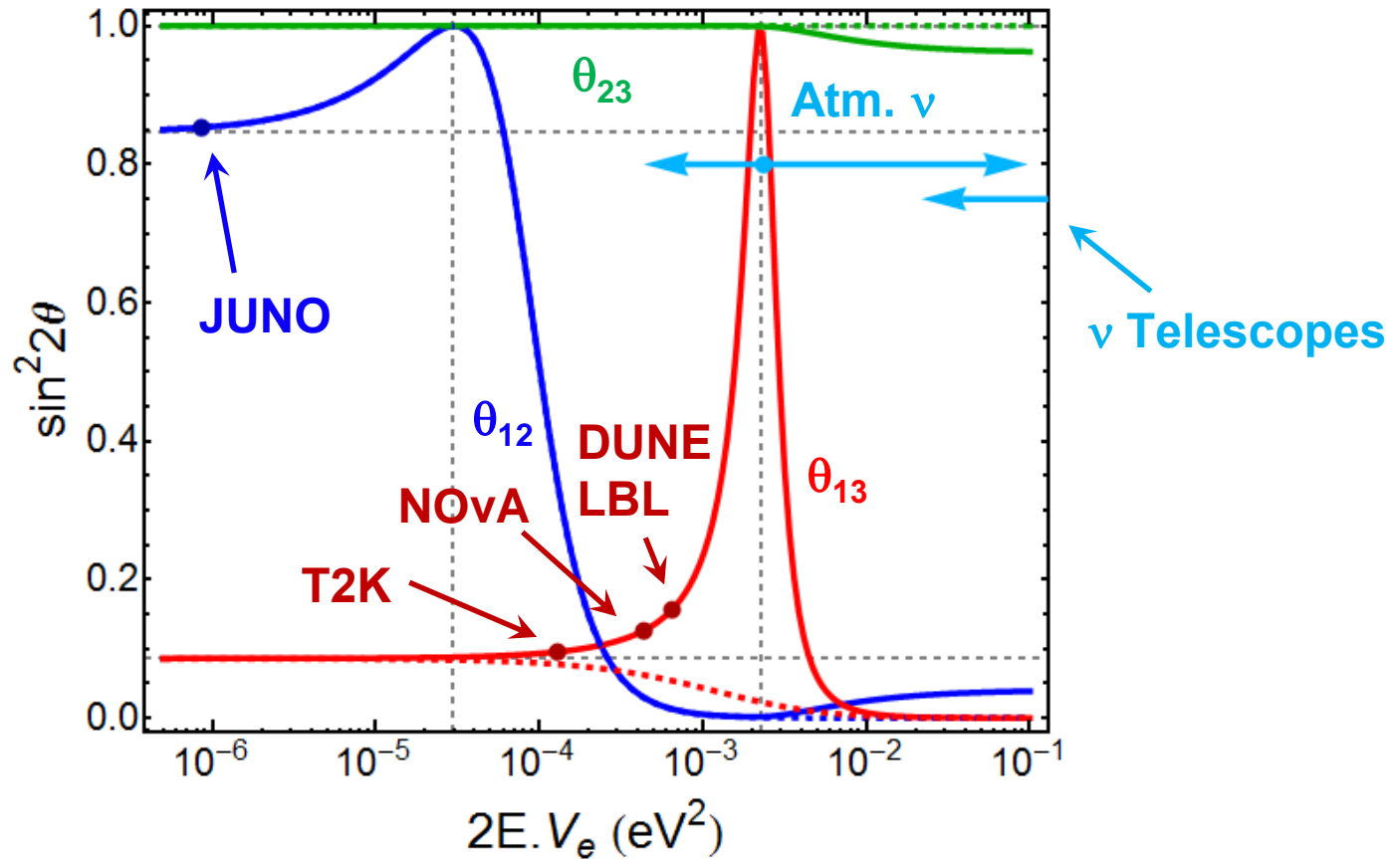
$$H_{eff} = U \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m^2_{21}}{2E} & 0 \\ 0 & 0 & \frac{\Delta m^2_{31}}{2E} \end{bmatrix} U^\dagger + V_e \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Resonances



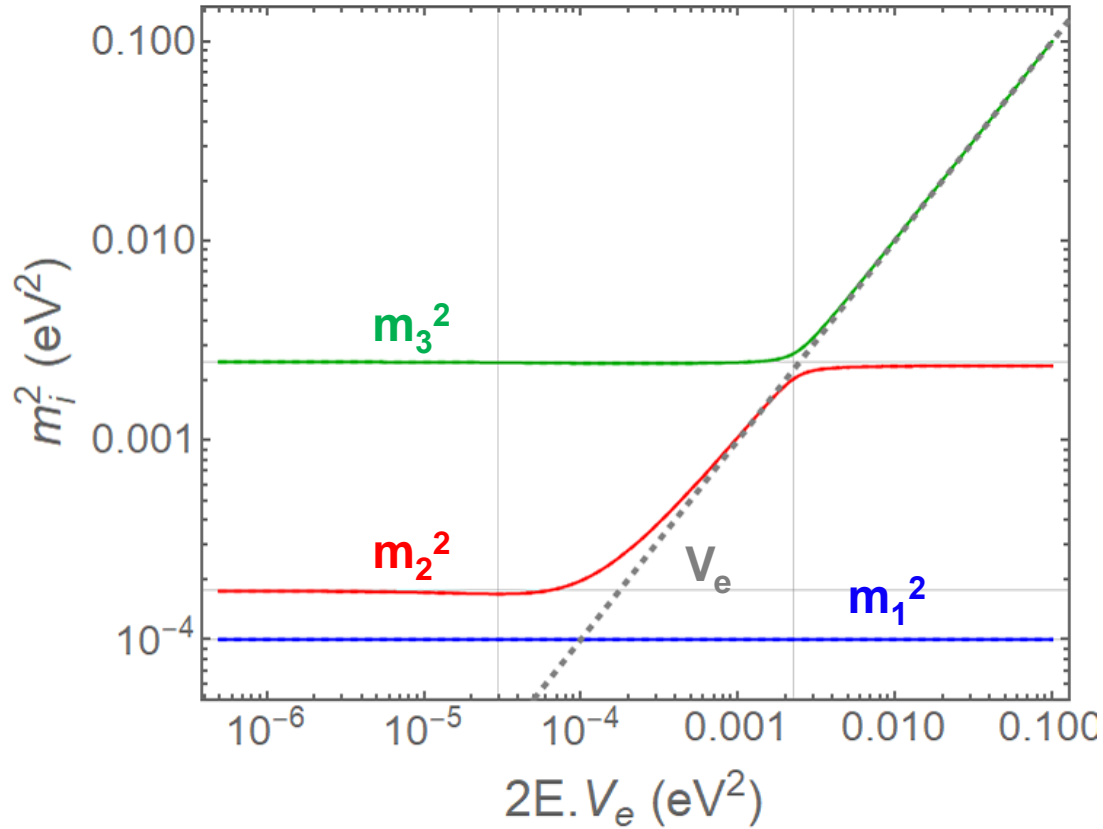
$$H_{eff} = U \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} \end{bmatrix} U^\dagger + V_e \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Resonances



$$H_{eff} = U \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} \end{bmatrix} U^\dagger + V_e \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

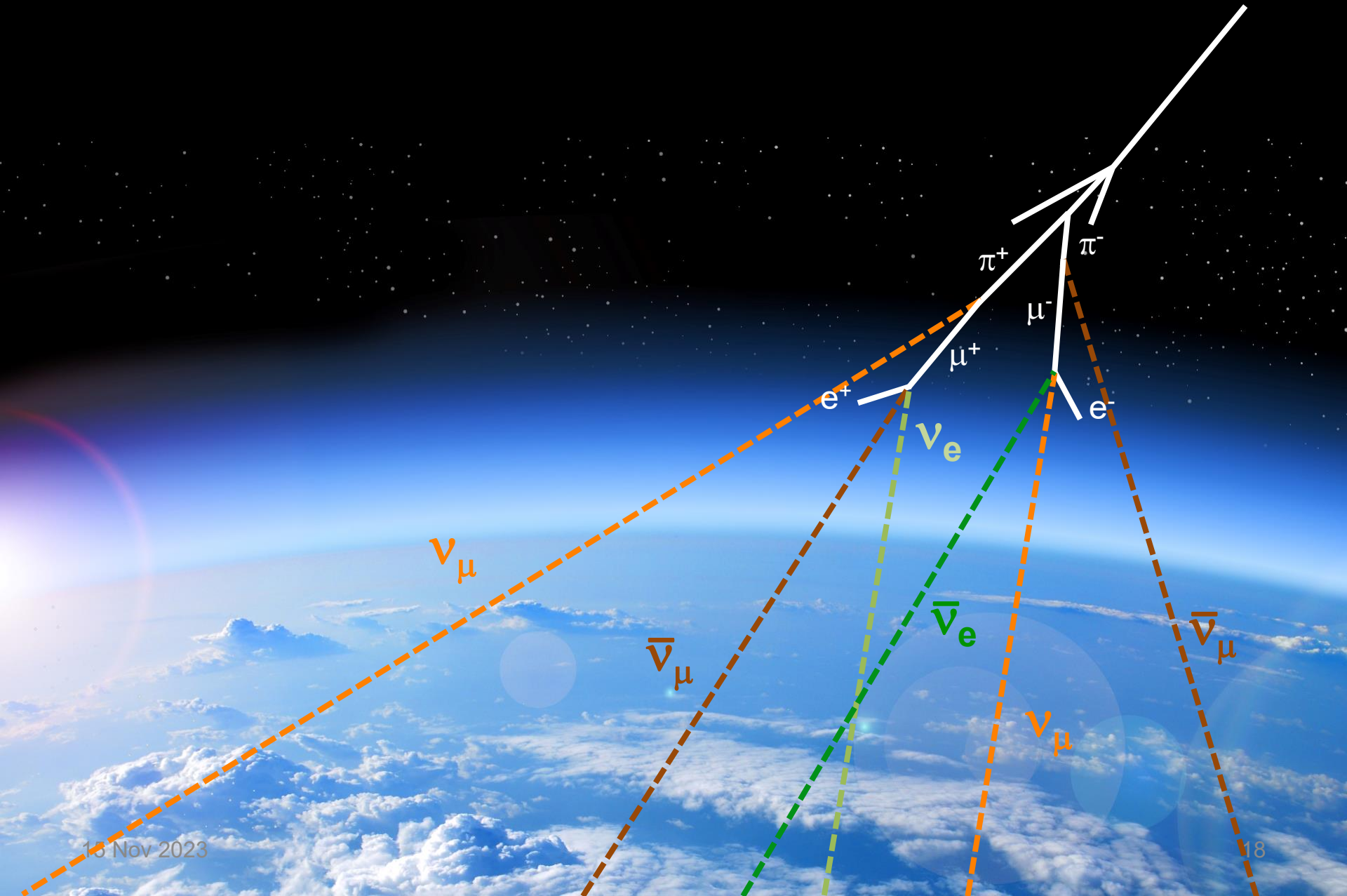
Resonances



$$H_{eff} = U \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} \end{bmatrix} U^\dagger + V_e \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Atmospheric Neutrinos

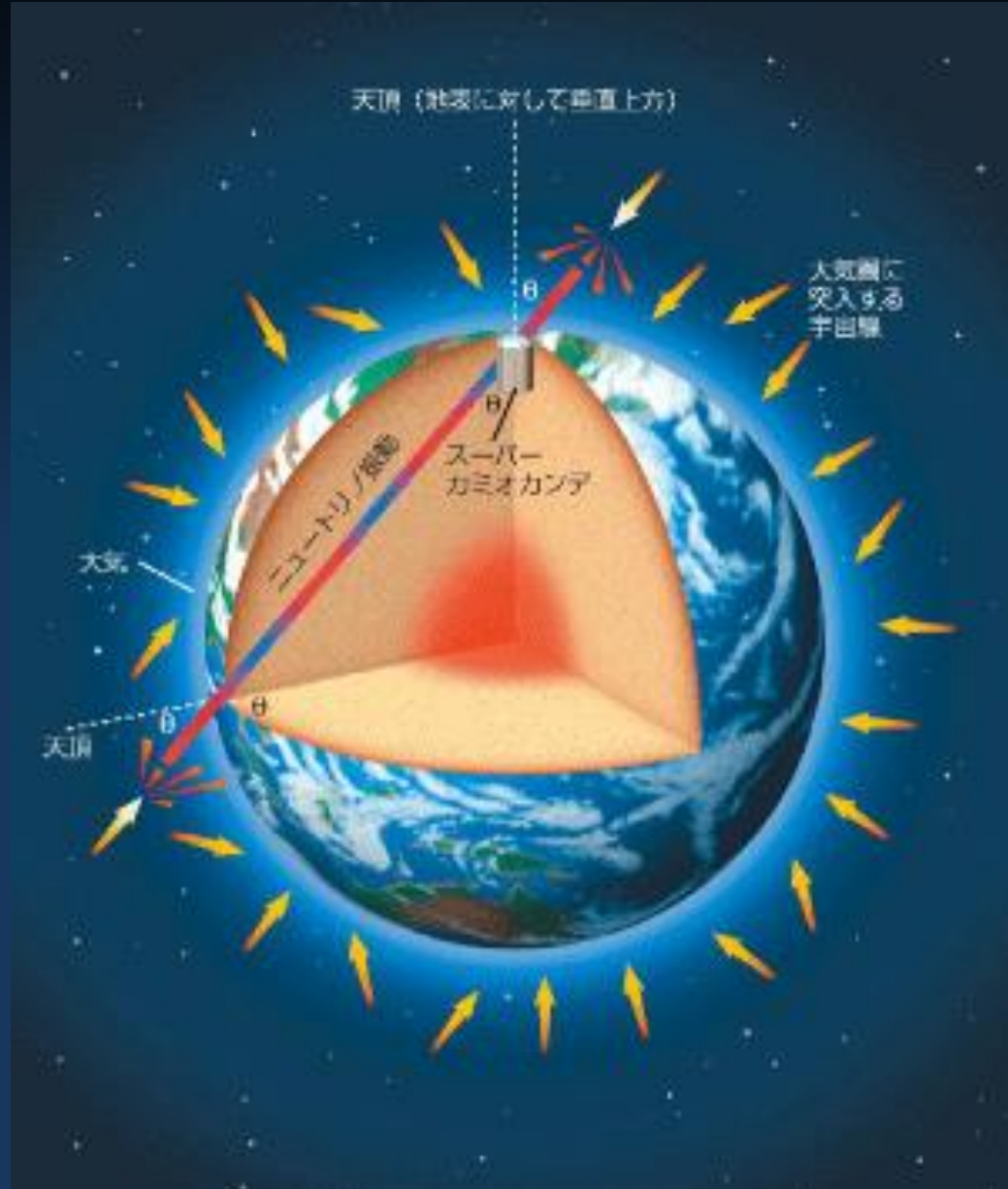
Cosmic Ray



16 Nov 2023

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Atmospheric Neutrinos



Neutrino Hamiltonian in Matter

$$H_{eff} = U \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} \end{bmatrix} U^\dagger + V_e \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

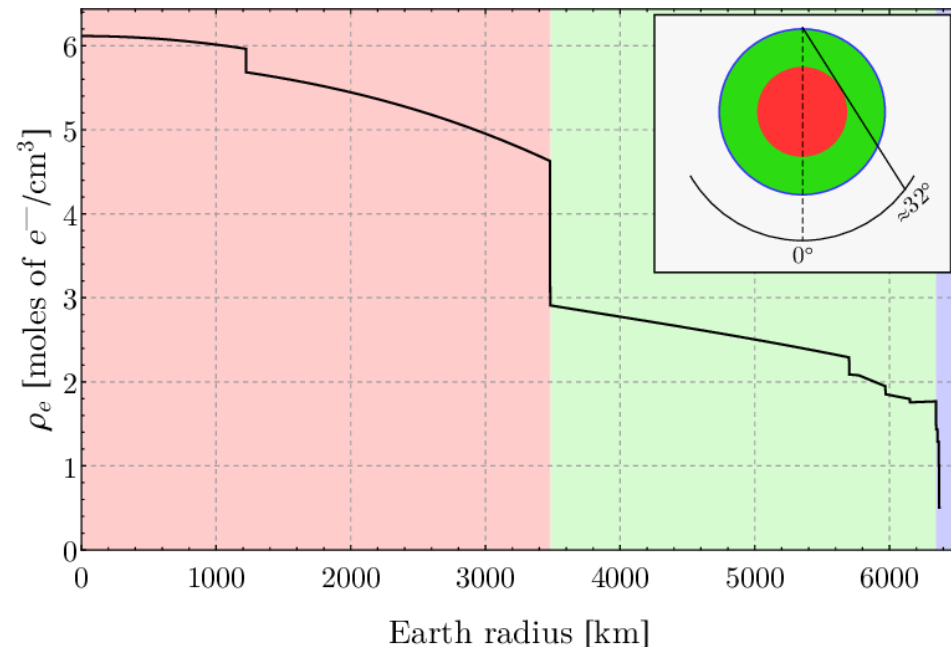
Time Dependent

$$V_e \equiv \pm \sqrt{2} G_F n_e \quad (+ \text{ for } \nu, - \text{ for } \bar{\nu})$$

$$\text{Relevant when } V_e \gtrsim \frac{\Delta m^2}{2E} \sim 1/L$$

$$\text{In general, } 1/V_e \sim 1700 \text{ km} \times \left(\frac{3 \text{ g/cm}^3}{\rho} \right)$$

$$\text{Comparable to } \frac{\Delta m_{31}^2}{2E} \text{ for } E \sim 10 \text{ GeV}$$



Quantum Evolution

$$\text{Schrödinger: } i \frac{\partial}{\partial t} \mathcal{U} = H \mathcal{U}$$

$$\mathcal{P}_{\alpha \rightarrow \beta} = |\langle \beta | \mathcal{U}(t) | \alpha \rangle|^2$$

Time-independent H :

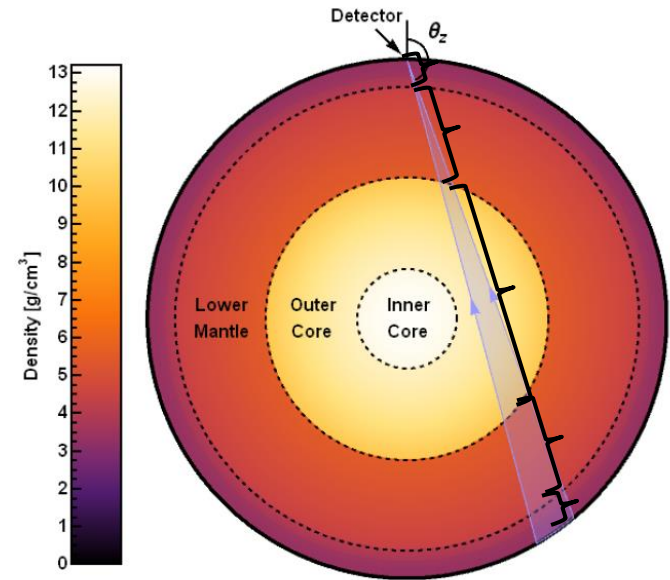
$$\mathcal{U}(t) = e^{-iHt}$$

$$H = V H_D V^\dagger$$

$$\mathcal{U}(t) = V e^{-iH_D t} V^\dagger$$

Time-dependent H :

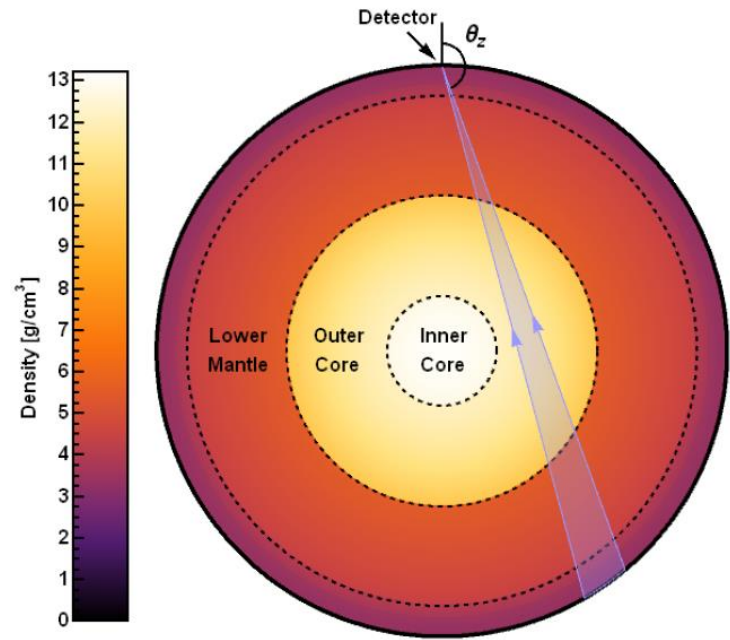
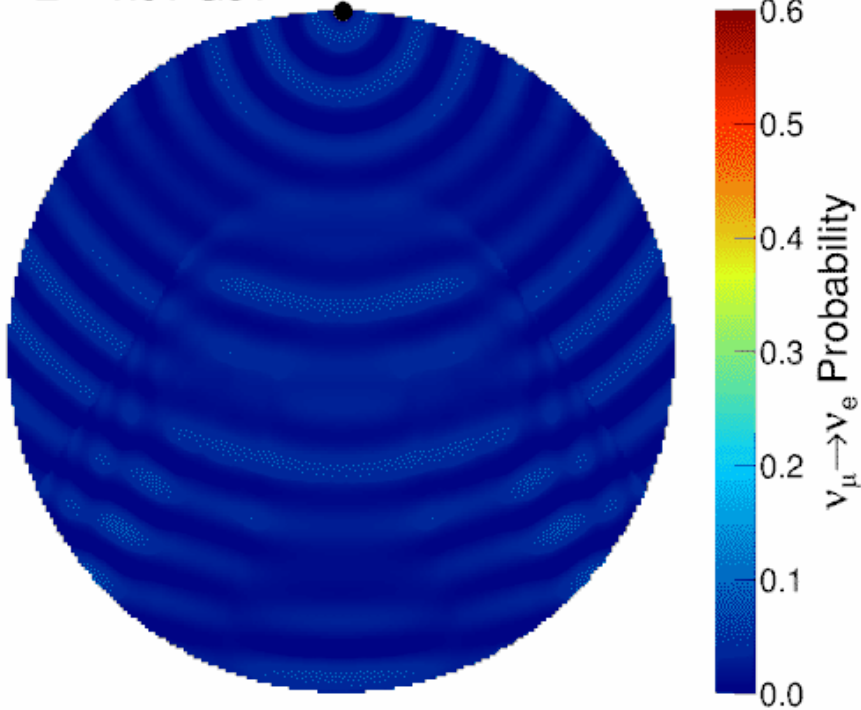
$$\mathcal{U}(t) = \mathcal{T} e^{-i \int_0^t H(t') dt'} \approx \prod_k e^{-iH(t_k) \Delta t}$$



- Trace neutrino path through the Earth
- Break path into N segments of similar electron density
- Compute evolution through each segment with constant density assumption

Atmospheric Neutrinos

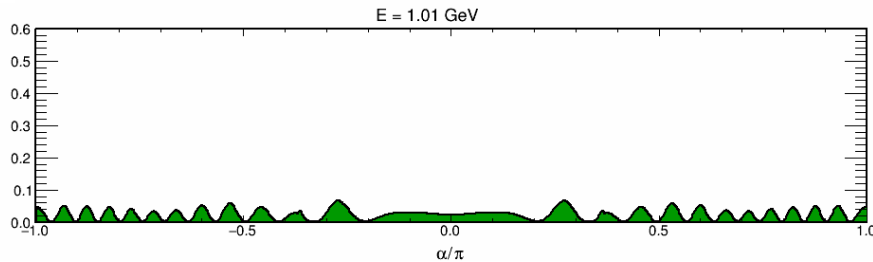
$E = 1.01 \text{ GeV}$



Oscillations are **resonant** at certain energies

$$E_{\text{res}} \sim 7 \text{ GeV in Mantle}$$

$$E_{\text{res}} \sim 3 \text{ GeV in Core}$$



ν_e appearance at the surface

Extended Models

Non-Standard Interactions (NSI)

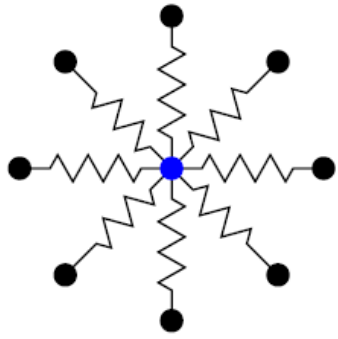
$$H_{eff} = U \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} \end{bmatrix} U^\dagger + V_e \begin{bmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{bmatrix}$$

Sterile Neutrinos (3+N Flavours)

$$H_{eff} = U_S \begin{bmatrix} 0 & 0 & 0 & 0 & \cdots \\ 0 & \frac{\Delta m_{21}^2}{2E} & 0 & 0 & \cdots \\ 0 & 0 & \frac{\Delta m_{31}^2}{2E} & 0 & \cdots \\ 0 & 0 & 0 & \frac{\Delta m_{41}^2}{2E} & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix} U_S^\dagger + \begin{bmatrix} V_e & 0 & 0 & 0 & \cdots \\ 0 & 0 & 0 & 0 & \cdots \\ 0 & 0 & 0 & 0 & \cdots \\ 0 & 0 & 0 & V_n/2 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix}$$

$$U_S = U_{N-1,N} \cdots U_{34} U_{24}^{(c)} U_{14}^{(c)} U_{23} U_{13}^{(c)} U_{12}$$

Extended Models



Decoherence

**HELP
NEEDED**

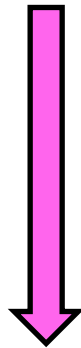
$$\partial_t \rho = \underbrace{-i[H, \rho]}_{\text{Unitary}} + \underbrace{\frac{1}{2} \sum_j 2A_j \rho A_j^\dagger - \{A_j^\dagger A_j, \rho\}}_{\text{Non-Unitary}}$$

Operators as sum
SU(3) of generators

$$\mathcal{O} = \sum_j \text{tr} [\mathcal{O} F_j] F_j$$

In general* 36 parameters!

$$\tilde{L}_{jk} = \frac{1}{2} \sum_{lmn} (\vec{a}_l \cdot \vec{a}_m) f_{lkn} f_{nmj}$$



$$\partial_t \vec{\rho} = (\tilde{H} - \tilde{L}) \vec{\rho}$$

System of 9
coupled equations

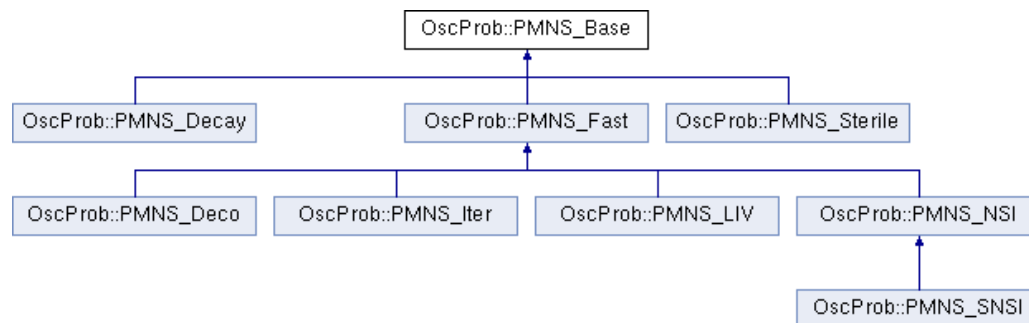
Diagonal w/ energy conserv.

$$\tilde{L} = \begin{bmatrix} 0_{3 \times 3} & 0 & 0 & 0 \\ 0 & I_2 \Gamma_{21} & 0 & 0 \\ 0 & 0 & I_2 \Gamma_{31} & 0 \\ 0 & 0 & 0 & I_2 \Gamma_{32} \end{bmatrix}$$

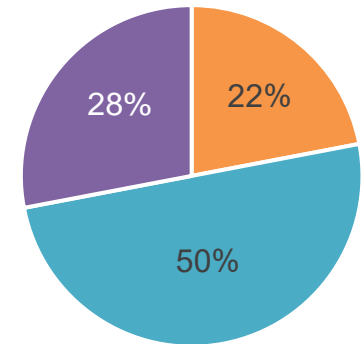
OscProb Package

- Diagonalises Hamiltonian to obtain **exact probabilities**
- Three step process:
 - **Build Hamiltonian** from parameters
 - **Solve Hamiltonian**
 - Fast algorithm from GLOBES for 3 neutrinos*
 - **Propagate neutrino** state
- Repeat for each step of constant matter in neutrino path
- PremModel class has built-in Earth layers model

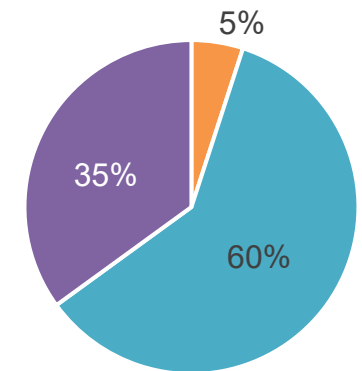
<https://github.com/joaobcoelho/OscProb>



Single Step (2.2 μ s)



85 steps (110 μ s) †

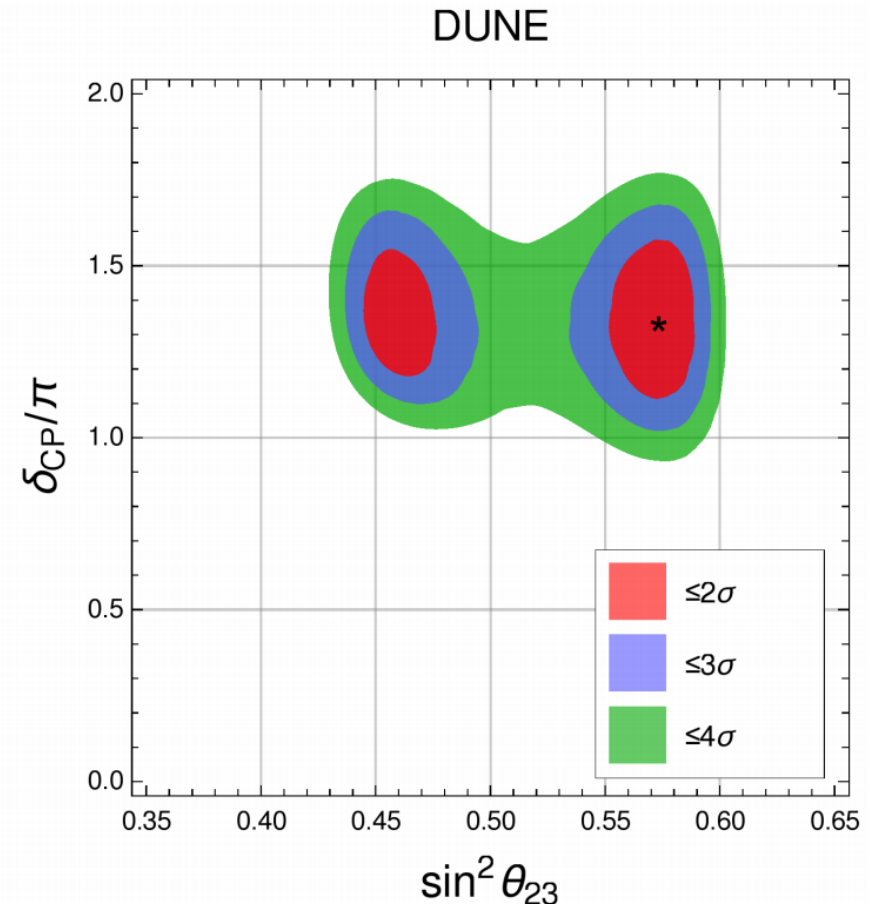


■ Build H ■ Solve H ■ Propagate

† Up-going (42+2 layers)

Why Performance Matters

- Computing oscillation probabilities is a big part of CPU time
- Total computations: $F_p \times F_d \times CP \times E_b \times \Theta_b \times L \times C \times P_n \times M$
 - F_p : Initial flavours produced (2)
 - F_d : Flavours detected (3+NC)
 - CP: Nu and nubar (2)
 - E_b : Energy bins (~ 100)
 - Θ_b : Direction bins (~ 100)
 - L: Earth layers to cross (~ 40)
 - C: $\Delta\chi^2$ surface points (e.g. contour: $\sim 50 \times 50$)
 - P_n : Nuisance parameters (e.g. syst.: ~ 20)
 - M: Minimisation steps (~ 100)
- Typically $\sim 10^{13}$ computations to obtain a full likelihood surface
- At $1\mu\text{s}$ / comp.: **~ 1 CPU-year**
- M typically grows with P_n
- **Feldman-Cousins corrections would require $\sim 10\text{k}$ contours**



Main Optimisations

- Use a fast algorithm for solving **eigensystem**
 - J. Kopp, Int. J. Mod. Phys. C, **19**, 523 (2008)
 - Uses analytical solutions for 3x3 matrix
 - Developed for GLoBES
- Some BSM models need different methods: Eigen library
 - PMNS_Sterile class solves NxN matrices
 - PMNS_Decay class solves non-Hermitian matrices
- Reduce number of operations when **computing Hamiltonian**
 - Take into account known form of PMNS matrix
- Caching of eigensystem for reusing known solutions

Hamiltonian Optimisation

$$U_S = U_{N-1,N} \cdots U_{34} U_{24}^{(c)} U_{14}^{(c)} U_{23} U_{13}^{(c)} U_{12}$$

$$H = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 10 & 0 & 0 \\ 0 & 0 & 20 & 0 \\ 0 & 0 & 0 & 40 \end{bmatrix} \xrightarrow{U_{12}} \begin{bmatrix} 3.2 & 4.7 & 0 & 0 \\ 4.7 & 6.8 & 0 & 0 \\ 0 & 0 & 20 & 0 \\ 0 & 0 & 0 & 40 \end{bmatrix}$$

$$\xrightarrow{U_{13}} \begin{bmatrix} 3.9 & 4.6 & 3.3 & 0 \\ 4.6 & 6.8 & -0.9 & 0 \\ 3.3 & -0.9 & 19 & 0 \\ 0 & 0 & 0 & 40 \end{bmatrix} \xrightarrow{U_{23}} \begin{bmatrix} 3.9 & 5.6 & -0.7 & 0 \\ 5.6 & 12 & 6.2 & 0 \\ -0.7 & 6.2 & 14 & 0 \\ 0 & 0 & 0 & 40 \end{bmatrix}$$

$$\xrightarrow{U_{14}} \begin{bmatrix} 5.3 & 5.5 & -0.7 & 7 \\ 5.5 & 12 & 6.2 & -1.1 \\ -0.7 & 6.2 & 14 & 0.1 \\ 7 & -1.1 & 0.1 & 39 \end{bmatrix} \xrightarrow{U_{24}} \begin{bmatrix} 5.3 & 6.8 & -0.7 & 5.8 \\ 6.8 & 12 & 6.1 & 4.2 \\ -0.7 & 6.1 & 14 & -1.1 \\ 5.8 & 4.2 & -1.1 & 38 \end{bmatrix}$$

$$\xrightarrow{U_{34}} \begin{bmatrix} 5.3 & 6.8 & 2.2 & 5.4 \\ 6.8 & 12 & 7.4 & 0.8 \\ 2.2 & 7.4 & 19 & 9.3 \\ 5.4 & 0.8 & 9.3 & 34 \end{bmatrix}$$

of Operations

	Opt.	Std.
2x2	3	8
3x3	13	54
4x4	34	192
5x5	70	500
6x6	125	1080

- Each rotation only affects some columns and rows
- Hermitian, so only upper triangle
- Total of **$O(n^3)$ operations**
- Std. matrix mult. $\rightarrow O(n^4)$

Extra Considerations

1. Caching of eigensystem solutions
 - Often we have to solve the same eigensystem multiple times
 - E.g.: Neutrinos cross same Earth layer from different angles
 - OscProb can save these to avoid repeated computations
 - Balance between hashing overhead and eigensystem
2. Earth models can be too detailed for our needs
 - Default is 44 layers of constant density
 - Can easily go down to 15 layers
 - Even better, OscProb can dynamically merge similar layers
3. Parallelise oscillation propagations
 - Usually we need to compute 2x3 transitions ($\nu_{e,\mu} \rightarrow \nu_{e,\mu,\tau}$)
 - OscProb can compute all of these in parallel
 - Avoids some extra propagation costs

Iterative Approximation

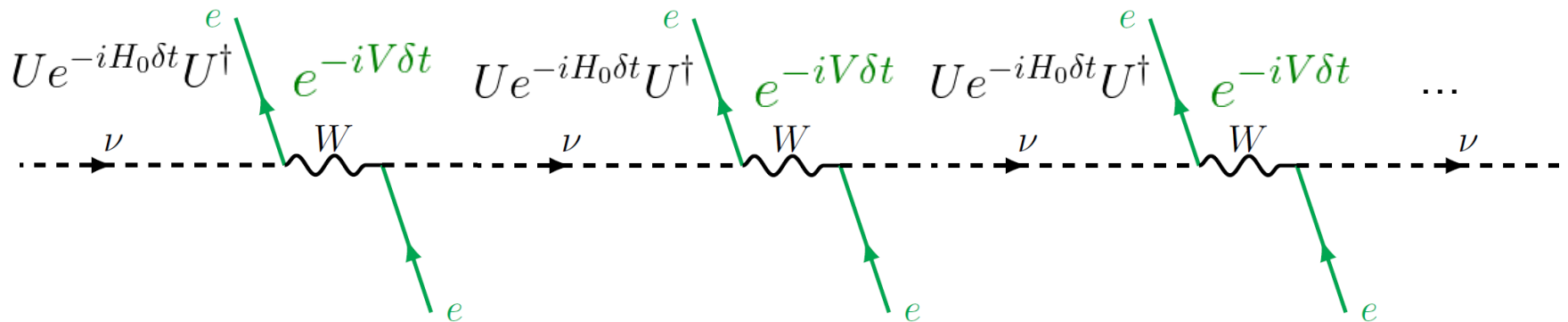
Zassenhaus Formula

$$e^{(X+Y)t} = e^{Xt} e^{Yt} e^{-[X,Y]\frac{t^2}{2}} e^{(2[Y,[X,Y]]+[X,[X,Y]])\frac{t^3}{6}} \dots$$

If t is small, we can ignore higher order terms

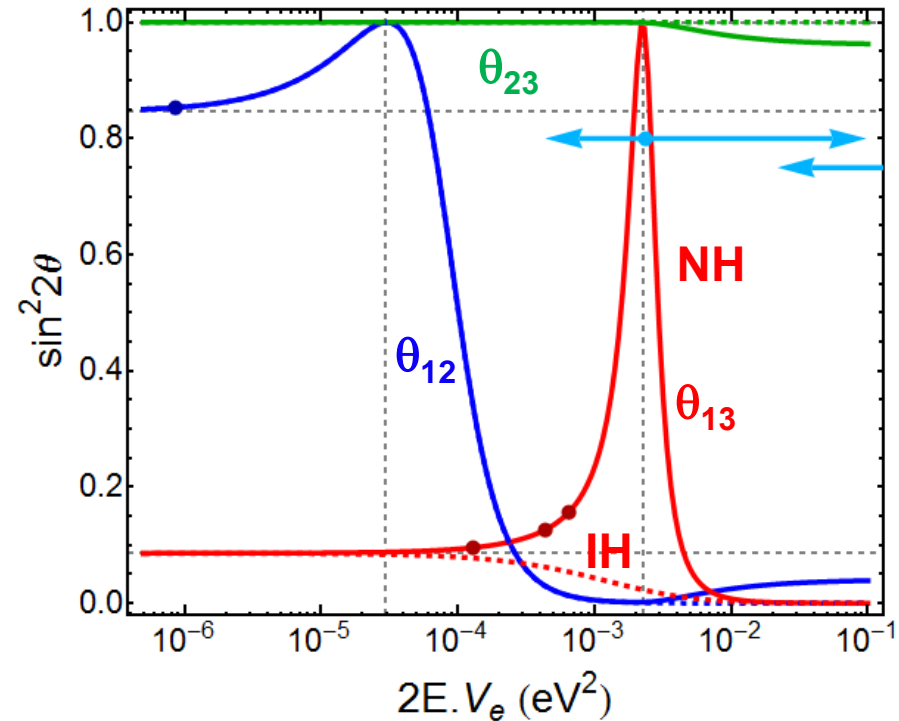
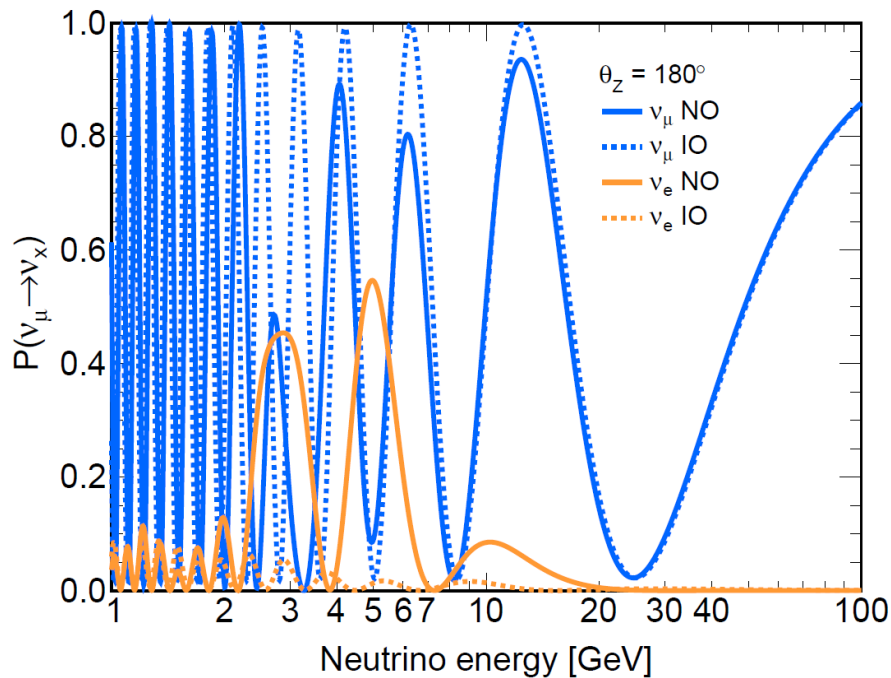
$$e^{-i(UH_0U^\dagger+V)t} \approx e^{-iUH_0U^\dagger t} e^{-iVt} = U e^{-iH_0 t} U^\dagger e^{-iVt}$$

Everything is already diagonal. No need for solving eigensystems



Interpolation?

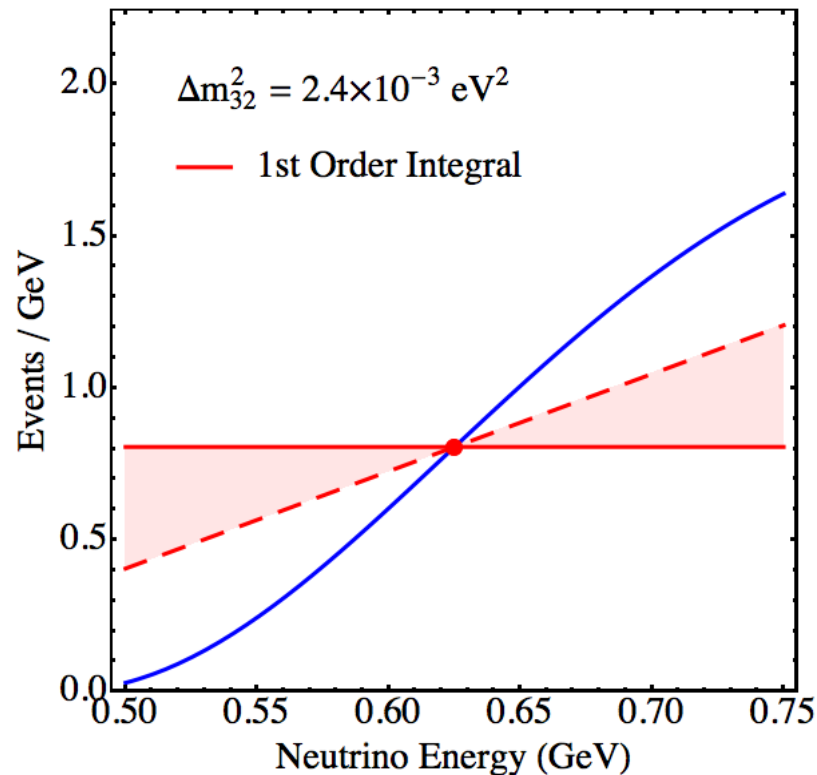
**HELP
NEEDED**



- Interpolating oscillation probabilities can be difficult due to fast oscillations at some energies
- One idea would be to instead interpolate the eigensystem solutions
- Equivalent to interpolating the effective mixing parameters

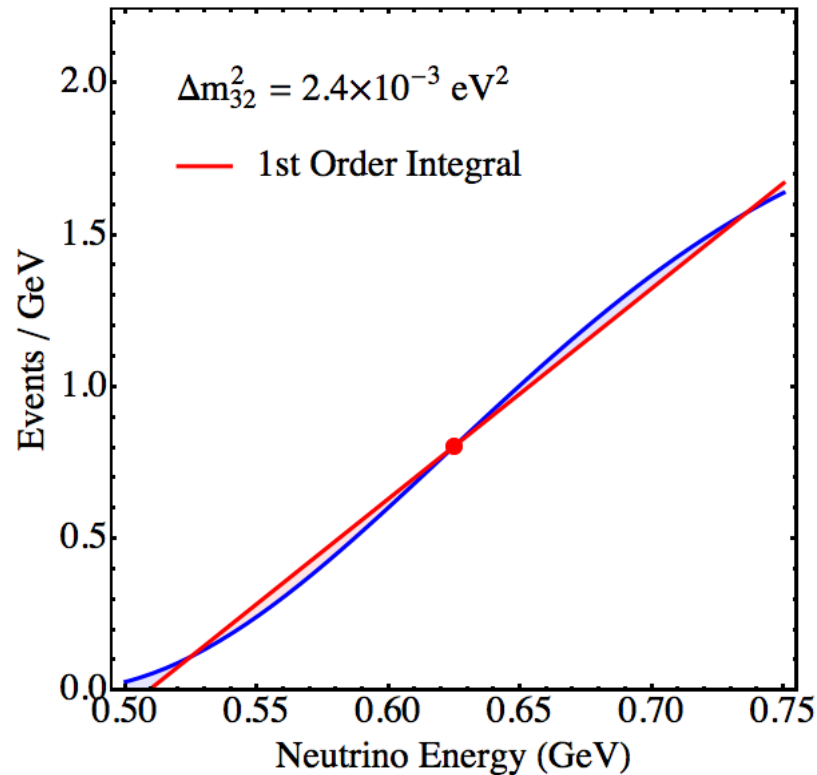
Oscillation Averaging

- Another problem with fast oscillations is computing the average oscillation over a bin
- Simply taking the bin center value is equivalent to a linear approximation of the function



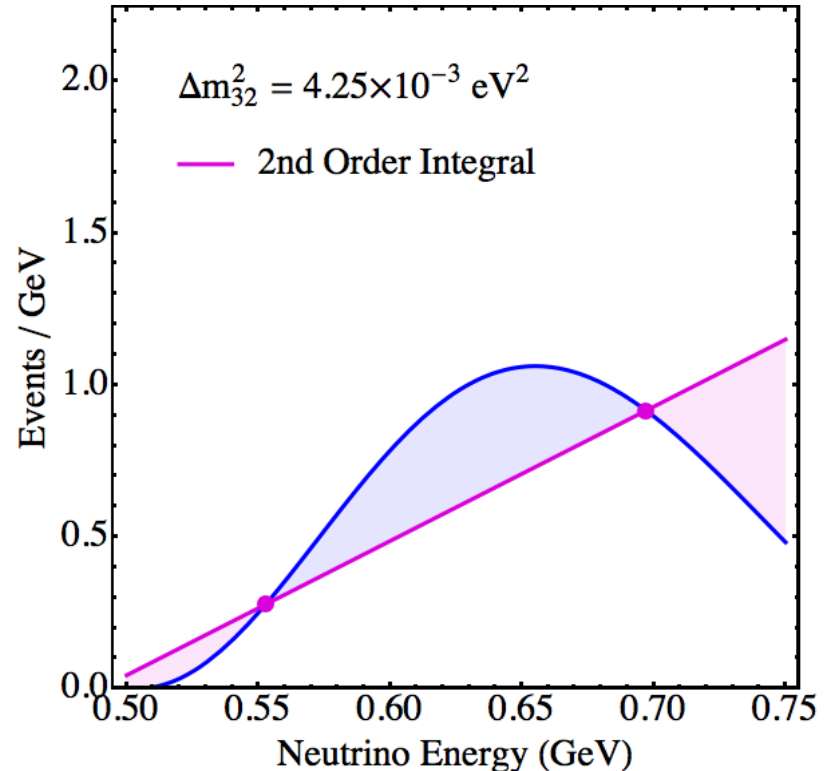
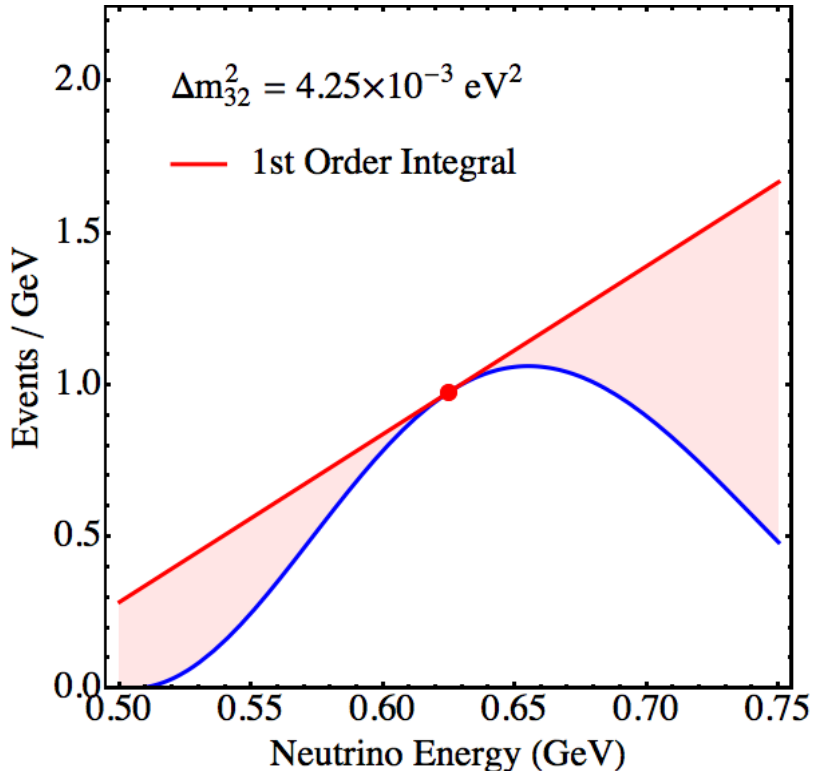
Oscillation Averaging

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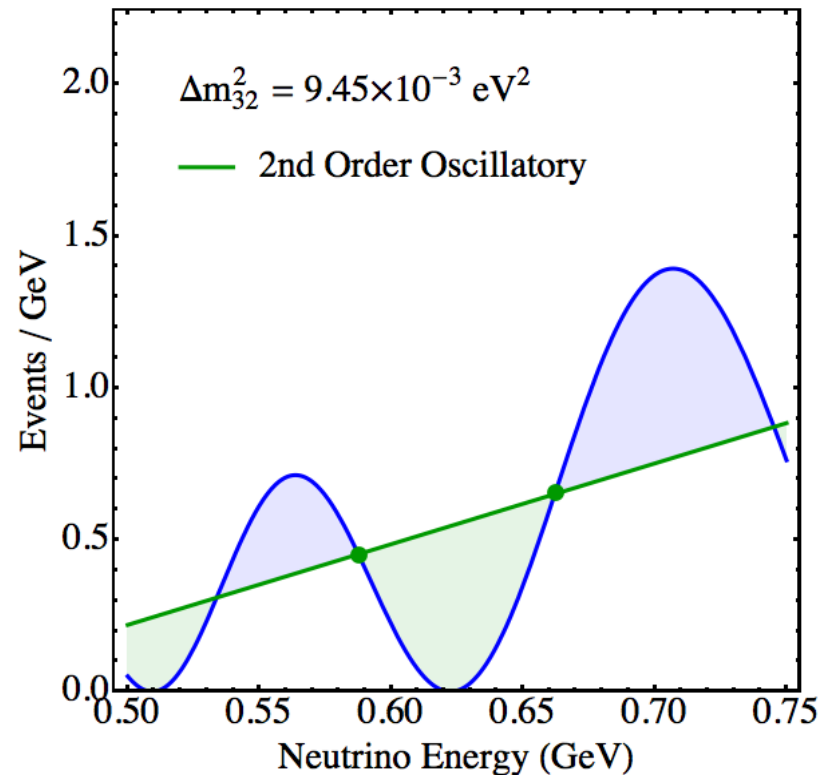
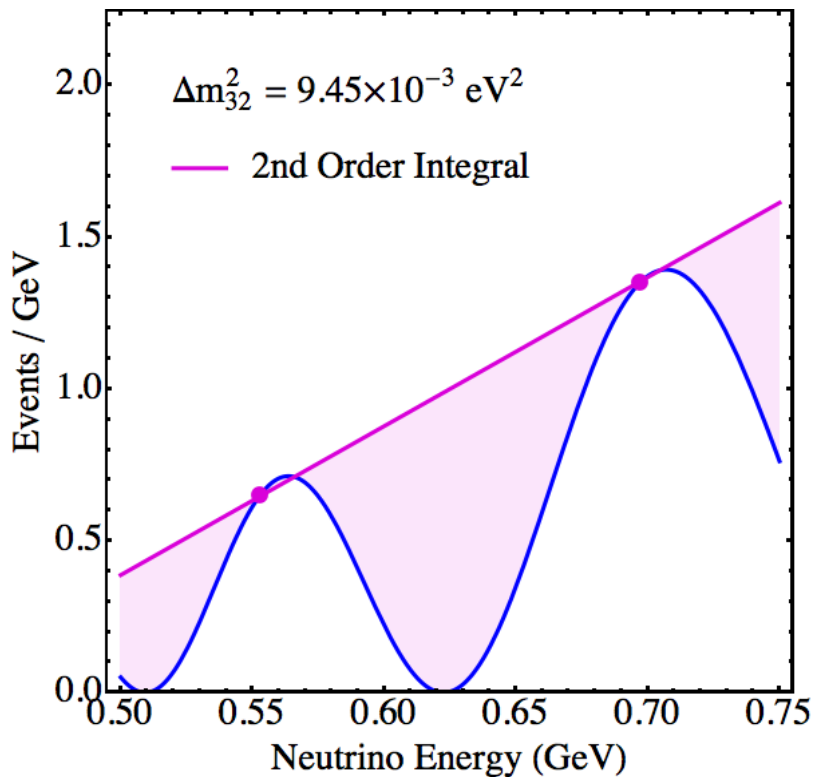
Oscillation Averaging

- Another problem with fast oscillations is computing the average oscillation over a bin
- Will fail if function is not approximately linear
- Gaussian quadrature improves this to 3rd order polynomial



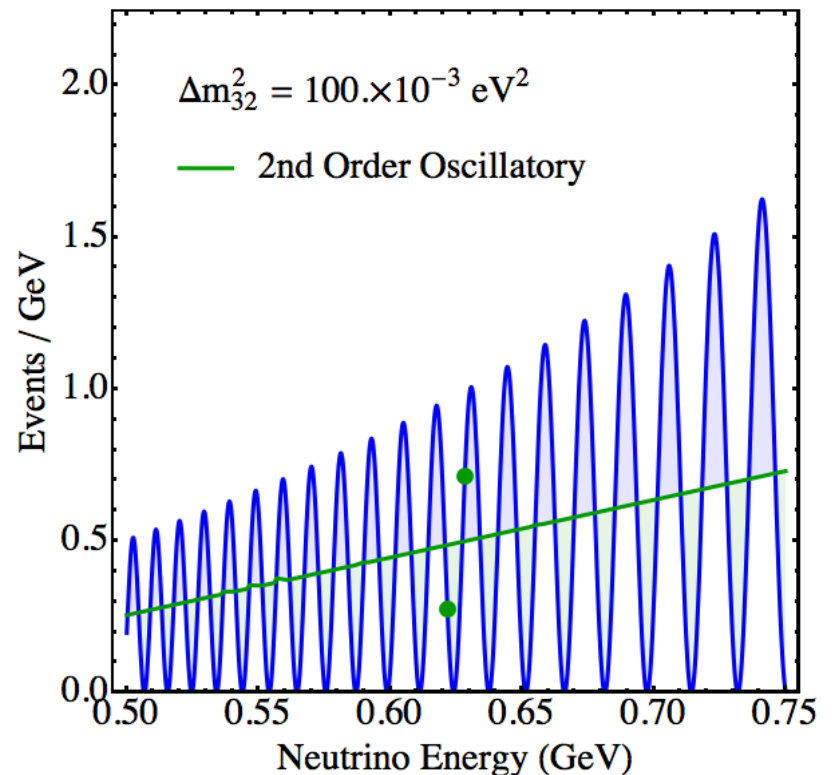
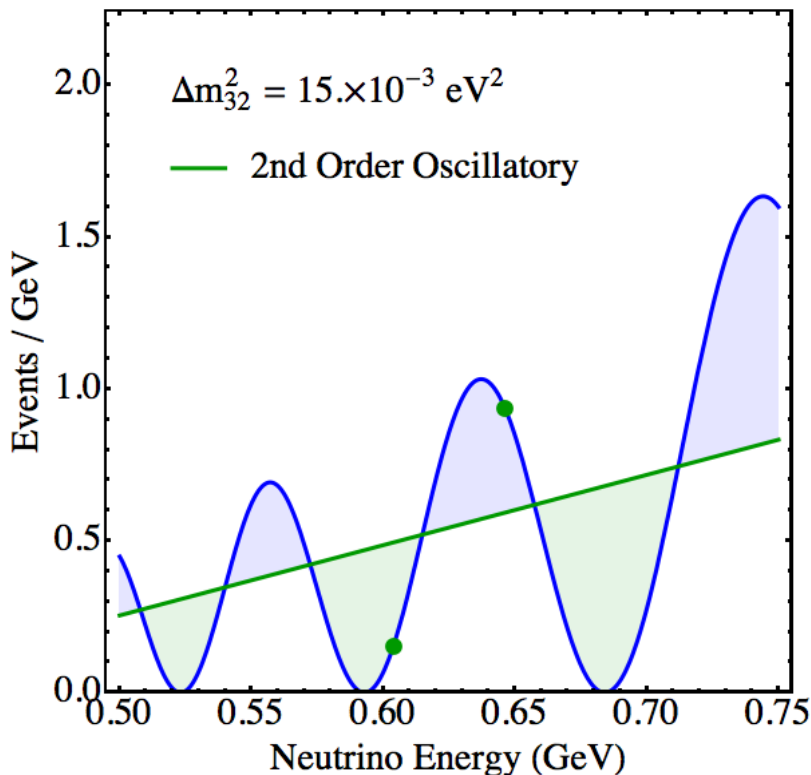
Oscillation Averaging

- Another problem with fast oscillations is computing the average oscillation over a bin
- But also fails for oscillating functions
- Can be extended by moving from polynomial to trig functions



Oscillation Averaging

- Another problem with fast oscillations is computing the average oscillation over a bin
- Works even at high frequencies, but requires known freqs
- Also, only valid for single frequency





Oscillation Averaging

HELP
NEEDED

- Solution can be generalized for multiple frequencies
- However, no analytical solution and hard to solve numerically
- Current implementation approximates by assuming frequencies are hierarchical and can be solved independently
- Numerical solution improves on this, but doesn't work reliably

perf: Use improved Gaussian quadrature rule in AvgProb #43

 Open joaoabcoelho wants to merge 1 commit into `master` from `dev-avgprob` 

 Conversation 2  Commits 1  Checks 1  Files changed 1



joaoabcoelho commented 2 weeks ago • edited ▾

Owner ...

Given that oscillation probabilities typically are of the form:

$$f(x) = a + b x + \sum_{i=1}^n c_i \cos(k_i x + \phi_i)$$

We can compute the integral of $f(x)$ via a quadrature rule solving for:

$$\int_{x_0 - \Delta x}^{x_0 + \Delta x} f(x) dx = \frac{1}{2^n} \sum_{j=1}^{2^n} f(x_0) + \sum_{i=1}^n (-1)^{[(j-1)/2^{i-1}]} \delta x_i$$

The solution satisfies:

$$\forall j \in \{1 \dots n\} \prod_{i=1}^n \cos(k_j \delta x_i) = \text{sinc}(k_j \Delta x)$$

Here I'm implementing solutions to this system of equations by Newton's method.



Alternative Approach

HELP
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- Use perturbation theory to approximate evolution:

$$S(\bar{E} + \xi_E) \approx \bar{S} e^{-i\mathbf{K}_E \xi_E} \quad \text{with} \quad \bar{S} = e^{-i\bar{H}L}$$

- Based on very interesting new paper:
 - <https://arxiv.org/abs/2308.00037>

From ray to spray: augmenting amplitudes and taming fast oscillations in fully numerical neutrino codes

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- MaCh3 does something similar adapted from SK

AvgProb Precision

**HELP
NEEDED**

- Whatever the solution, we need to be able to estimate the approximation error
- Current estimates are too conservative
- Default precision set to 0.01% but samples too many points
- At analysis level, one needs to optimize the balance between precision and speed in these calculations

Benchmarking

**HELP
NEEDED**

- Compute 100x20 oscillograms for $\nu_{\mu} \rightarrow \nu_e$
- Using AvgProb to remove fast oscillations

```
32 Processing StressTest.C...
33 PMNS_Fast: Performance = 142 μs/iteration
34 PMNS_Iter: Performance = 72 μs/iteration
35 PMNS_Sterile: Performance = 940 μs/iteration
36 PMNS_NSI: Performance = 155 μs/iteration
37 PMNS_Deco: Performance = 267 μs/iteration
38 PMNS_Decay: Performance = 241 μs/iteration
39 PMNS_LIV: Performance = 328 μs/iteration
40 PMNS_SNSI: Performance = 217 μs/iteration
```

- Would be great to have some comparisons with other tools

Conclusion

**HELP
NEEDED**

- Oscillation calculations are the core of many analyses
- Computations are fast, but need to be done trillions of times
- OscProb is an open source option for computing oscillations
- Goal is to be fast and cover many BSM models
- If you have ideas or would like to help with current issues, you are very welcome to contribute to the project in Github
- Integrating OscProb into DUNE tools like MaCh3 would also be very much appreciated

Joao Coelho, Rebekah Pestes, Alba Domi, Simon Bourret, Ushakrhm, Imaderer, & vicacuen. (2023). joaoabcoelho/OscProb: v2.0.12 (v2.0.12). Zenodo. <https://doi.org/10.5281/zenodo.10104847>

Backup

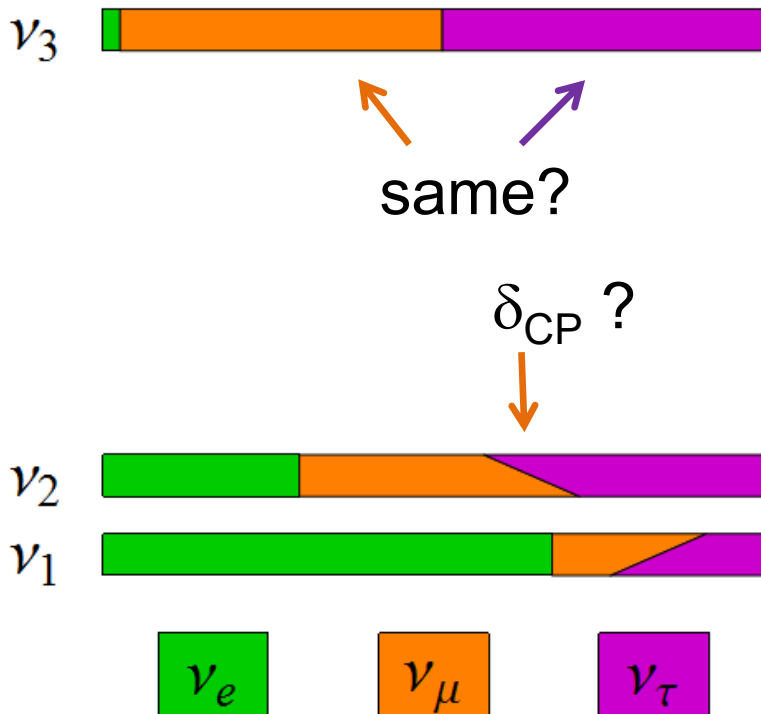
Missing Pieces

symmetries

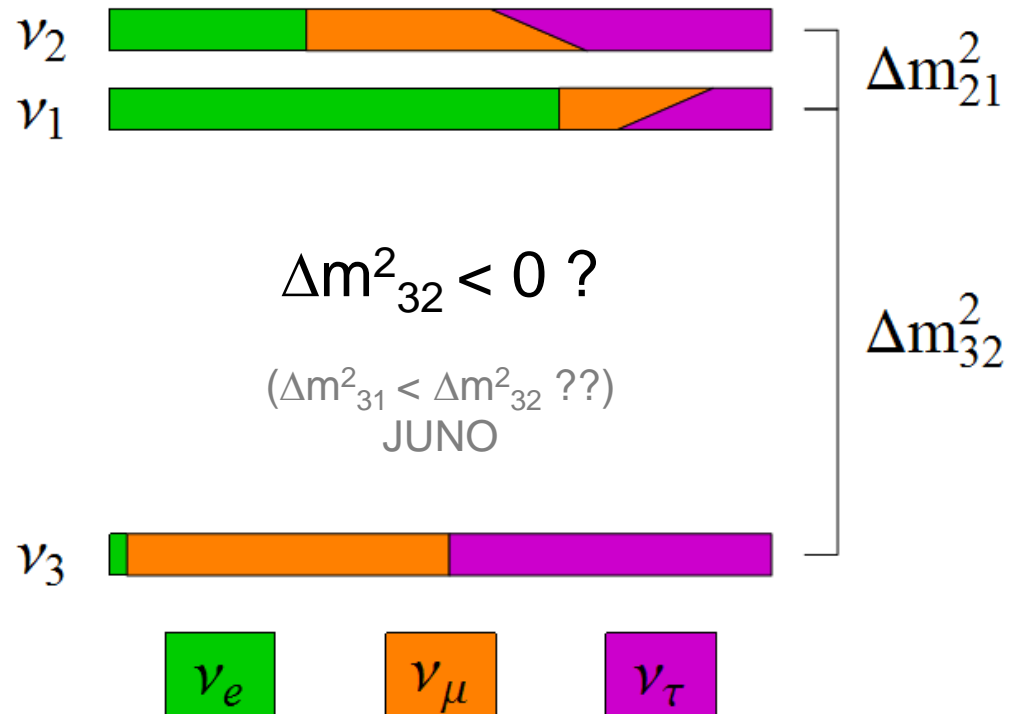
$$\sin^2 2\theta \times \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$$

- Is $\theta_{23} = \pi/4$? Underlying symmetry?
- Do neutrinos violate CP? (δ_{CP})
- **What is the mass ordering? (Mass Hierarchy)**

Normal Hierarchy



Inverted Hierarchy



Matter Effects

