Towards ZTF-III

Nao Suzuki (Lawrence Berkeley National Lab / LPNHE)

- Multi SN Spieces (SNIb, SNIc, SNII)
- Probing Dark Galaxies
- RR Lyrae, Cepheid Variables (DESI-II Spectra)
- Mapping the local Universe

Probing "Dark" galaxies (1/3) ZTF × LSST (Subaru/HSC)



Probing "Dark" galaxies (2/3) ZTF × LSST (Subaru/HSC)



Probing "Dark" galaxies (3/3) ZTF × LSST (Subaru/HSC)



Probing IR with Euclid ZTF x Euclid (Image + Spec)



Euclid : IR Spec Schedule is fixed for next 6 years

Euclid : J, H, K

ZTF-III start observing Euclid field 2 weeks in advance

Pre-ZTF Meeting in Japan 2019

TDA-MMS 2019: Time Domain Astronomy in the Era of Massively Multiplexed Spectroscopy February 8-10, 2019, Nikko, Japan



Distribution of Local Matter on BAO Scale Nao Suzuki & Friends

Q Where are we?

Hubble's Law $V = H_0 d$

v = recession velocity in km/sec

d = distance in Mpc

 H_0 = expansion rate today (*Hubble Parameter*) In words:

The more *distant* a galaxy, the *faster* its recession velocity.

General Relativity Metric

Robertson-Walker Metric

$$ds^{2} = dt^{2} - \frac{R^{2}(t)}{c^{2}} \left[\frac{dr^{2}}{1 - kr^{2}} + r^{2}d\theta^{2} + r^{2}sin^{2}\theta d\phi^{2} \right]$$

Schwartzschild Metric

$$ds^{2} = \left(1 - \frac{2m}{r}\right)dt^{2} - \frac{dr^{2}}{1 - \frac{2m}{r}} - r^{2}d\theta^{2} - r^{2}\sin^{2}\theta\,d\varphi^{2}$$

Density Field at z=0



Density Field at z=0



How to quantify density fluctuation?







Einj said: 1

Thanks a lot! Does that simply mean $\sigma_8 = \mathcal{P}(k = 1/8 h/Mpc)$? Where \mathcal{P} is the power spectrum.

Harmonics has to be counted Window Function W(kR)





8 Mpc/h : RMS is Unity (0.8-ish)



150Mpc represents the Universe



Q What is the area to study for 150Mpc?

Q1:z=2

Q2:z=0

CMB Dipole is seen by SNIa

A&A 560, A90 (2013)









Inhomogeneous Expansion of the Universe

The trouble with Hubble: Local versus global expansion rates in inhomogeneous cosmological simulations with numerical relativity

HAYLEY J. MACPHERSON,¹ PAUL D. LASKY,^{1,2} AND DANIEL J. PRICE¹

¹Monash Centre for Astrophysics and School of Physics and Astronomy, Monash University, VIC 3800, Australia ²OzGrav: The ARC Centre of Excellence for Gravitational-wave Discovery, Clayton, VIC 3800, Australia

Macpherson et al $2018^{\text{Submitted to ApJ}}_{\text{ABSTRACT}}$

In a fully inhomogeneous, anisotropic cosmological simulation performed by solving Einstein's equations with numerical relativity, we find a local measurement of the effective Hubble parameter differs by less than 1% compared to the global value. This variance is consistent with predictions from Newtonian gravity. We analyse the averaged local expansion rate on scales comparable to Type 1a supernova surveys, and find that local variance cannot resolve the tension between the Riess et al. (2018a) and Planck Collaboration et al. (2018) measurements.

Density Field at z=0



Expansion Rate at z=0 (normalized by global) Inhomogeneous Expansion of the Universe



 $\mathcal{H}_{\mathcal{D}}/\mathcal{H}_{all}$ Macpherson et al 2018

Hubble Flow vs SNIa Cosmology





Foley et al 2018



Macpherson et al 2018



Figure 3. Local deviations in the Hubble parameter due to inhomogeneities. We show the full distribution of all spheres in the range $75 < r_{\mathcal{D}} < 180 \, h^{-1}$ Mpc in blue. The dashed blue lines represent the 1σ deviation of the inhomogeneous distribution. The blue shaded region represents the 1σ uncertainties on the Planck Collaboration et al. (2016) measurement, while the solid red line and shaded region represent the mean and 1σ deviation in the Riess et al. (2018a) measurement, respectively.



Figure 1. Expansion rate and density of an inhomogeneous, anisotropic universe. Left panel shows the deviation in the Hubble parameter relative to the global mean \mathcal{H}_{all} . Right panel shows the density distribution relative to the global average, $\langle \rho \rangle_{all}$. Both panels show a slice through the midplane of a 256³ resolution simulation with L = 1 Gpc.

Macpherson et al 2018

Today (500 SNeIa z <0.1)



ZTFII (3000 SNeIa z <0.1)



ZTFIII (30,000 SNeIa)



Beginning to see the sign



çiile ROTAL ASTRONOMICAL SOCT MNRAS \$10, 216-229 (2022)

The DES view of the Eridanus supervoid and the CMB cold spot

A. Kovács⁶,^{1,2}*† N. Jeffrey⁶,^{3,4} M. Gatti,^{3,6} C. Chang,^{7,8} L. Whiteway,⁴ N. Hamaus,⁹ O. Lahav,⁴ G. Pollina,⁹ D. Bacon,¹⁰ T. Kacprzak,¹¹ B. Mawdsley,¹⁰ S. Nadathur,⁴ D. Zeurcher,¹¹ J. García-Bellido,¹² A. Alarcon,¹³ A. Amon,¹⁴ K. Bechtol,¹⁵ G. M. Bernstein,⁶ A. Campos,¹⁶ A. Carnero Rosell,^{1,2,17} M. Carrasco Kind,^{18,19} R. Cawthon,¹⁵ R. Chen,²⁰ A. Choi,²¹ J. Cordero,²² C. Davis,¹⁴ J. DeRose,^{23,24} C. Doux,⁶ A. Drlica-Wagner,^{7,5,25} K. Eckert,⁶ F. Elsner,⁴ J. Elvin-Poole,^{21,26} S. Everett,²⁴ A. Ferté,²⁷ G. Giannini,⁵ D. Gruen,^{9,28,29} R. A. Gruendl,^{18,19} I. Harrison,^{22,30} W. G. Hartley,³¹ K. Herner,²⁵ E. M. Huff,²⁷ D. Huterer,³² N. Kuropatkin,²⁵ M. Jarvis,⁵ P. F. Leget,¹⁴ N. MacCrann,³³ J. McCullough,¹⁴ J. Muir,³⁴ J. Myles,^{14,28,29} A. Navarro-Alsina,³⁵ S. Pandey,⁶ J. Prat,⁷ M. Raveri,⁸ R. P. Rollins,²² A. J. Ross,²¹ E. S. Rykoff,^{14,29} C. Sánchez,⁶ L. F. Secco,⁶ I. Sevilla-Noarbe,³⁶ E. Sheldon,³⁷ T. Shin,⁶ M. A. Troxel,²⁰ I. Tutusaus,^{38,39} T. N. Varga,^{9,40} B. Yanny,²⁵ B. Yin,¹⁶ Y. Zhang,²⁵ J. Zuntz,⁴¹ M. Aguena,^{17,42} S. Allam,²⁵ F. Andrade-Oliveira,^{17,43} J. Annis,²⁵ E. Bertin,^{44,45} D. Brooks,⁴ D. Burke,^{14,29} J. Carretero,⁵ M. Costanzi,^{45,47,43} L. N. da Costa,^{17,49} M. E. S. Pereira,³² T. Davis,⁵⁰ J. De Vicente,³⁵ S. Desai,⁵¹ H. T. Diehl,²⁵ I. Ferrero,⁵² B. Flaugher,²⁵ P. Fosalba,^{35,39} J. Frieman,^{8,25} E. Gaztañaga,^{38,39} D. Gerdes,^{32,53} T. Giannantonio,^{54,55} J. Gschwend,^{17,49} G. Gutierrez,²⁵ S. Hinton,⁵⁰ D. L. Hollowood,²⁴ K. Honscheid,²¹ D. James,⁵⁶ K. Kuchn,^{57,58} M. Lima,^{17,42} M. A. G. Maia,^{17,49} J. L. Marshall,⁵⁹ P. Melchior,⁵⁰ F. Menanteau,^{18,19} R. Miquel,^{5,61} R. Morgan,¹⁵ R. Ogando,^{17,79} F. Paz-Chinchon,^{18,54} A. Pieres,^{17,49} A. A. Plazas,⁶⁰ M. Rodriguez Monroy,³⁶ K. Romer,⁶² A. Roodman,^{14,29} E. Sanchez,³⁶ M. Schubnell,³² S. Serrano,^{33,39} M. Smith,⁶³ M. Soares-Santos,³² E. Suchyta,⁶⁴ M. E. C. Swanson,¹⁸ G. Tarle,³² D. Thomas,¹⁰ C.-H. To^{14,28,29} and J. Weller^{9,40}

Affiliations are listed at the end of the paper

 3×10^{-1}

 4×10^{-1}

DES and the cold spot 217



Brout et al. 2022

Summary : Fact Sheet

- <u>150Mpc (BAO scale)</u> represents the local universe
- All Sky Survey is needed
- $\sigma 8$ is the measure of fluctuation amplitude, $\sigma 100$ is about 2%
- Redshift is a sum of peculiar velocity + Hubble Flow
- By using SNIa, we can extract peculiar velocity to prove local density fluctuations
- With 100 SNeIa, CMB Dipole is seen
- z<0.05 SNeIa were threw away in the past for SNIa cosmology business

Discussion (z<0.05 Science)

- Peculiar Velocity = Local Density Field
- Local Density Field can be modeled
- Supergalactic Plane can be mapped
- Are we in a void?
- Local Density Map can be drawn from 10,000 SNeIa
- How do we get SNIa Spectra?
- Southern Hemisphere Facilities are needed
- Inhomogeneous Expansion of the Universe should be detectable

Backup Slides

(8)

Local gravity versus local velocity: solutions for β and non-linear bias

Marc Davis,^{1*} Adi Nusser,² Karen L. Masters,³ Christopher Springob,⁴ John P. Huchra⁵ and Gerard Lemson⁶

¹Departments of Astronomy & Physics, University of California, Berkeley, CA 94720, USA

²Physics Department and the Asher Space Science Institute-Technion, Haifa 32000, Israel

³Institute for Cosmology and Gravitation, University of Portsmouth, Dennis Sciama Building, Burnaby Road, Portsmouth PO1 3FX

⁴Anglo-Australian Observatory, PO Box 296, Epping, NSW 1710, Australia

⁵Harvard–Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138, USA

⁶Max-Planck Institute of Astrophysics, Karl-Schwarzschild-Str. 1, 85741 Garching, Germany

2.2 TF sample

20 yr ago, the mis-calibration of full sky Tully–Fisher (TF) data was the problem that led to very discrepant results for the determination of $\beta \equiv \Omega/b$, with $\beta = 0.5 \pm 0.2$ (DNW96) and $\beta = 1.0 \pm 0.2$ (e.g. Dekel et al. 1993). The mistaken TF calibration led to a largescale flow that confused both analyses, but in the end, it was a calibration error in the southern sky which made a false large-scale flow (Willick et al. 1997). In one analysis, this led to a higher χ^2 than was acceptable, and in the other it led to a biased result.

$$\boldsymbol{v}_{\mathrm{g}}(\boldsymbol{r}) = \frac{2f(\Omega)}{3H_0\Omega}\boldsymbol{g}(\boldsymbol{r}).$$

3.2 Peculiar velocities from the inverse Tully-Fisher relation

Given a sample of galaxies with measured circular velocity parameters, $\eta_i = \log \omega_i$, linewidth ω_i , apparent magnitudes m_i and redshifts z_i , the goal is to derive an estimate for the smooth underlying peculiar velocity field. We assume that the circular velocity parameter, η , of a galaxy is, up to a random scatter, related to its absolute magnitude, M, by means of a linear ITF relation, i.e.

$$\eta = \gamma M + \eta_0.$$

One of the main advantages of inverse TF methods is that samples selected by magnitude, as most are, will be minimally plagued by Malmquist bias effects when analysed in the inverse direction (Schechter 1980; Aaronson et al. 1982). We write the absolute magnitude of a galaxy.

$$A_i = M_{0i} + P_i \tag{9}$$

where

$$M_{0i} = m_i + 5\log(z_i) - 15$$
(10)

and

 $P_i = 5\log(1 - u_i/z_i),$ (11)

where m_i is the apparent magnitude of the galaxy, z_i is its redshift in units of km s⁻¹ and u_i is its radial peculiar velocity in the LG frame.



Figure 5. The derived peculiar velocities v_{ITF} , v_g and $v_{\text{ITF}} - v_g$ of galaxies on aitoff projections on the sky in galactic coordinates. The rows correspond to galaxies with cz < 2000, 2000 < cz < 4000, $4000 < cz < 6000 \text{ km s}^{-1}$ and $6000 < cz < 10\,000 \text{ km s}^{-1}$, respectively. The size of the symbols is linearly proportional to the velocity amplitude (see key to the size of the symbols given at the bottom of the figure). In order to better see the differences, a 400 km s^{-1} dipole, in the direction of the CMB dipole, has been subtracted from the v_{ITF} and v_g velocities. Note that $v_{\text{ITF}} - v_g$ is considerably smaller than v_{ITF} or v_g , even for the most distant galaxies.

SNIa Supernova (z<0.05) Feindt et al 2013 : SN Factory

U. Feindt et al.: Measuring cosmic bulk flows



Fig. 1. Peculiar velocities of individual SNe determined from their distance moduli, μ_i , by solving Eq. (2) for v_{DF} . The plots show the Union2 (*top row*) and SNFACTORY (*bottom row*) datasets in the redshift bins 0.015 < z < 0.035 (*left column*), 0.035 < z < 0.045 (*middle column*) and 0.045 < z < 0.06 (*right column*). The marker diameter of each SN is proportional to the absolute value of the velocity plus an offset (see the scale at the top right), with red circles corresponding to positive velocities and blue squares corresponding to negative ones. For reference, the directions of the CMB dipole and the Shapley supercluster (SSC) are shown.