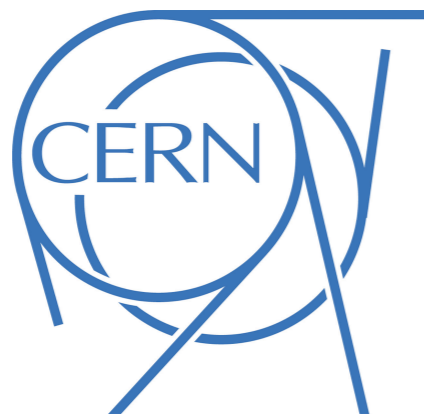

A Swampland Review for Cosmologists



Irene Valenzuela

CERN

IFT UAM-CSIC



String-Cosmo day

Paris, November 2023

Swampland program

What is the space of Effective Field Theories weakly coupled to Einstein gravity that can be consistently UV completed (in quantum gravity)?

Swampland program

What is the space of Effective Field Theories weakly coupled to Einstein gravity that can be consistently UV completed (in quantum gravity)?

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$

Swampland program

What is the space of Effective Field Theories weakly coupled to Einstein gravity that can be consistently UV completed (in quantum gravity)?

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$

What can it go here?

The diagram features two orange circles highlighting the terms \mathcal{L}_{EFT} and $\frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}}$ in the equation above. Two orange arrows originate from these circles and point towards the text "What can it go here?" located below the equation.

Swampland program

What is the space of Effective Field Theories weakly coupled to Einstein gravity that can be consistently UV completed (in quantum gravity)?

$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$

Not every EFT can be UV completed!

What can it go here?

Swampland program

What is the space of Effective Field Theories weakly coupled to Einstein gravity that can be consistently UV completed (in quantum gravity)?

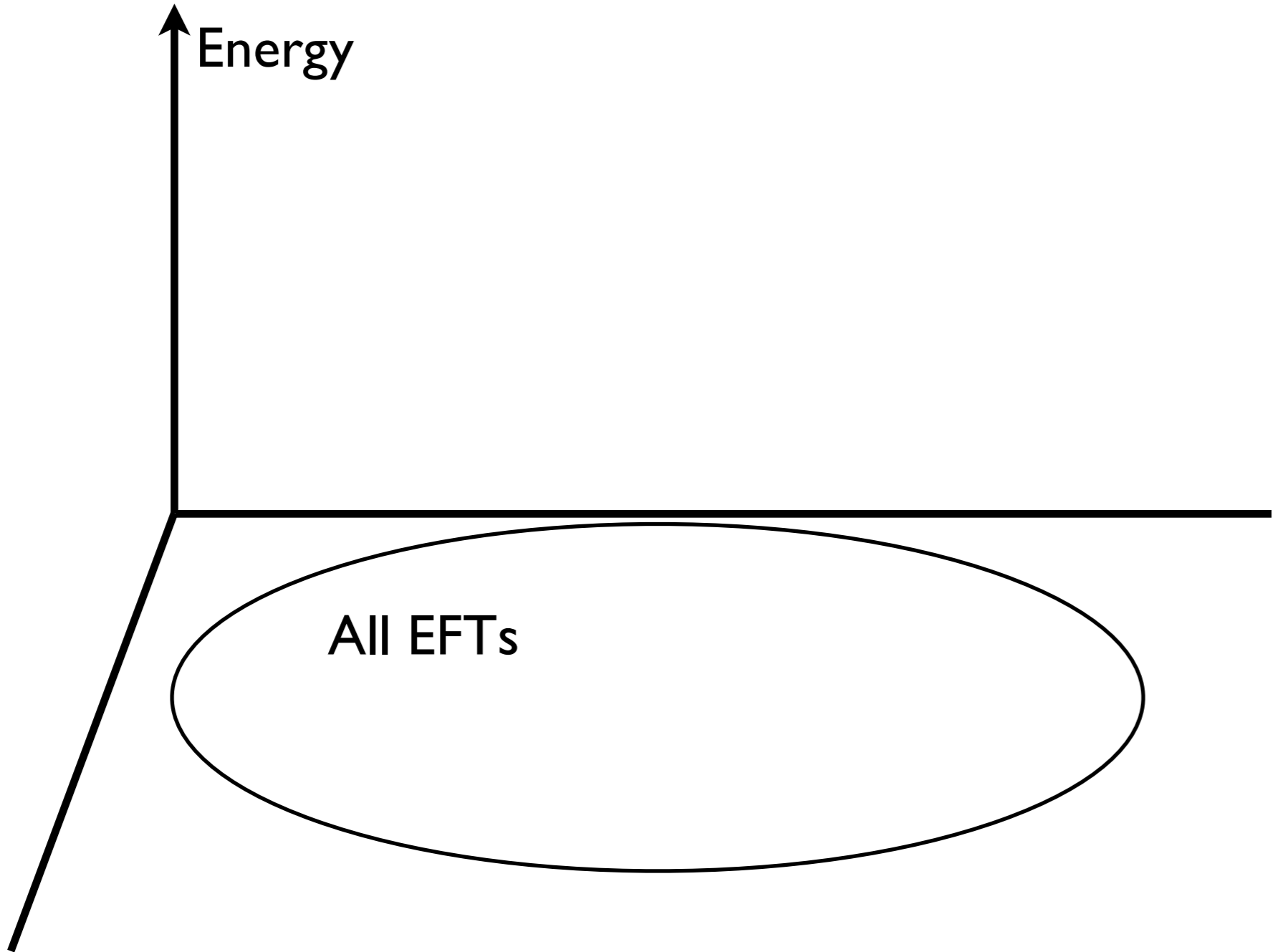
$$S = \int d^4x \left(R + \mathcal{L}_{\text{EFT}} + \sum_n \frac{\mathcal{O}_{n+4}}{\Lambda_{\text{cut-off}}^{n-4}} \right)$$

Not every EFT can be UV completed!

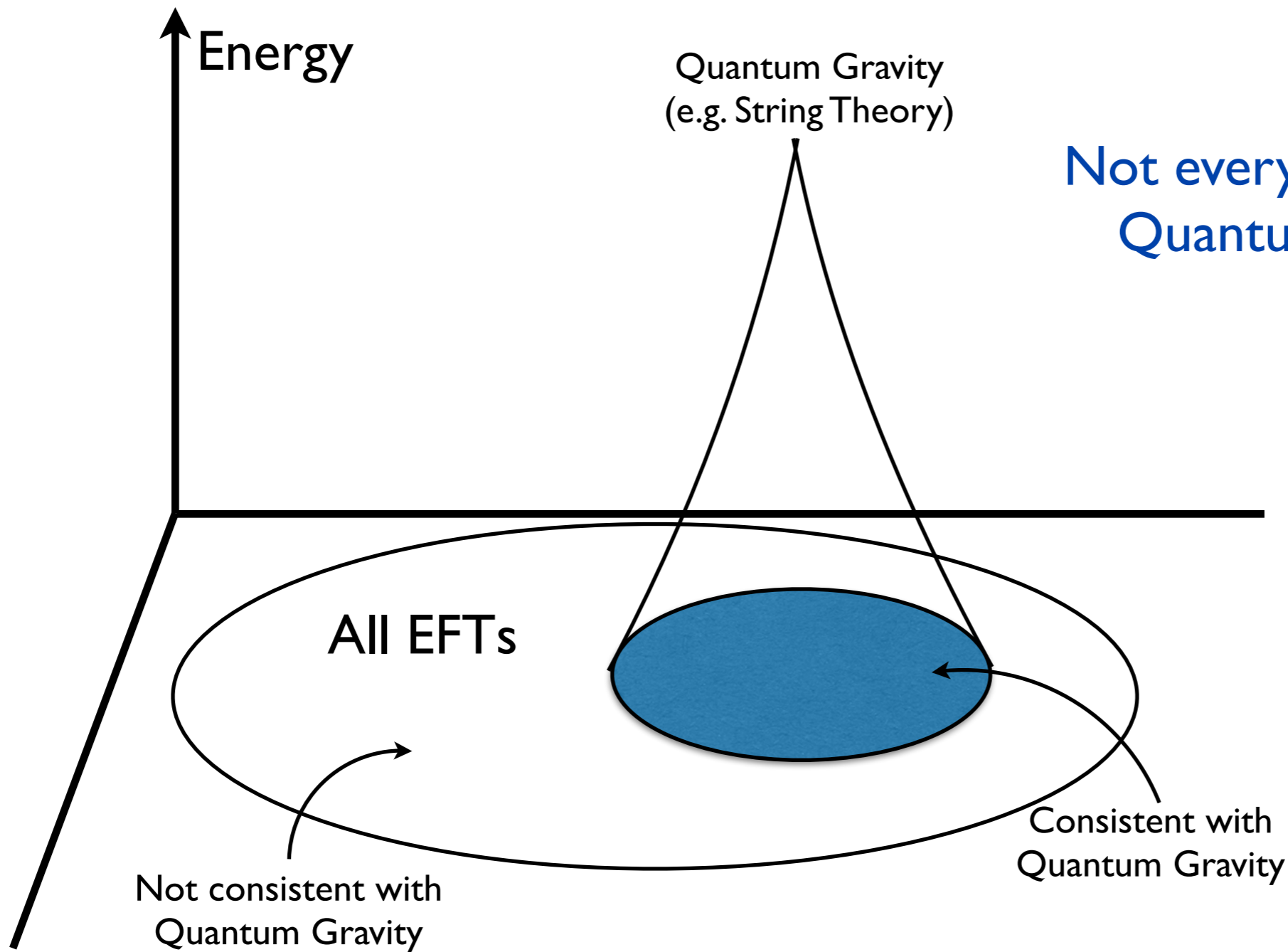
What can it go here?

What is the quantum gravity cut-off?

Energy

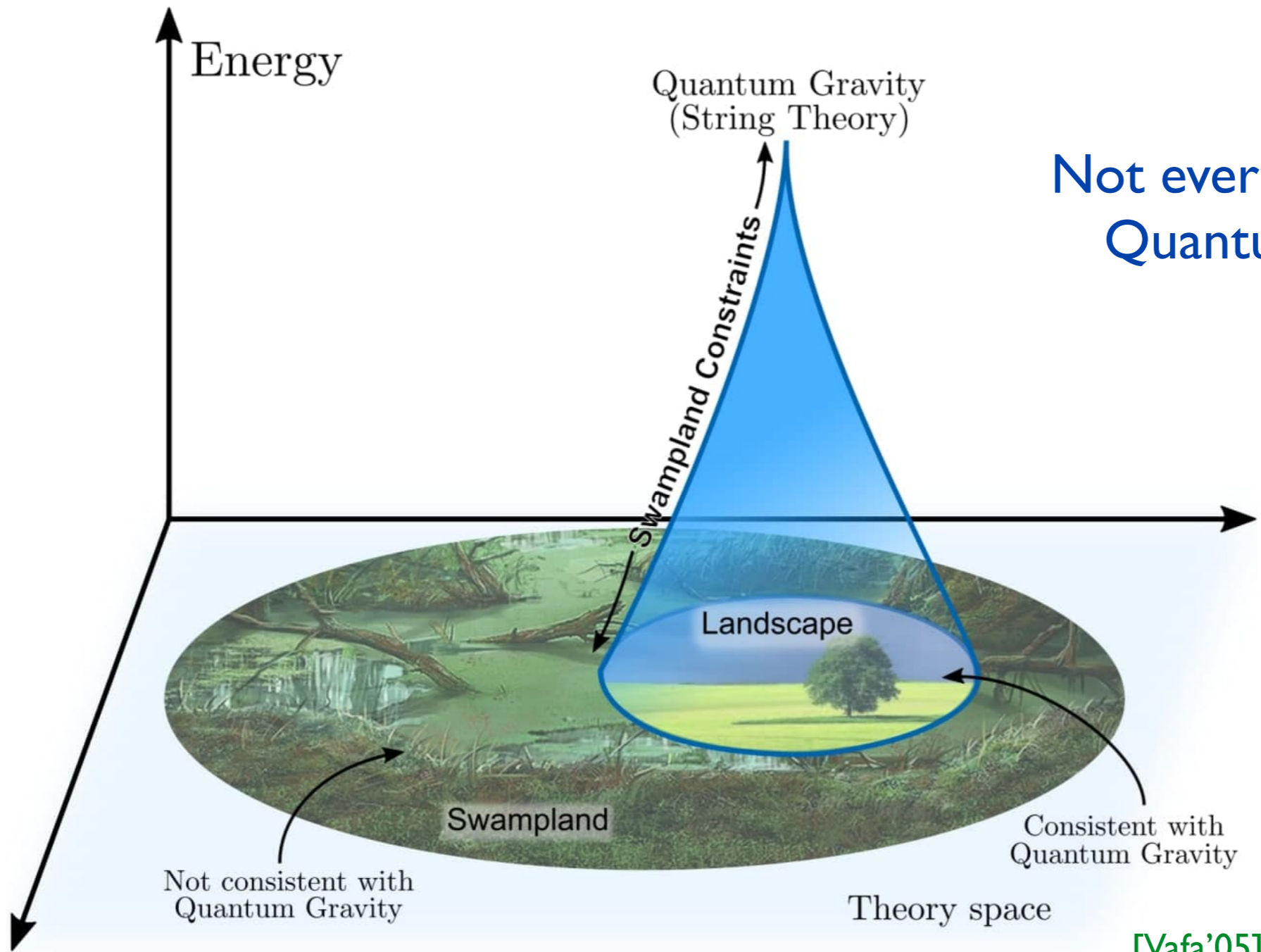


All EFTs



Not everything goes in
Quantum Gravity!

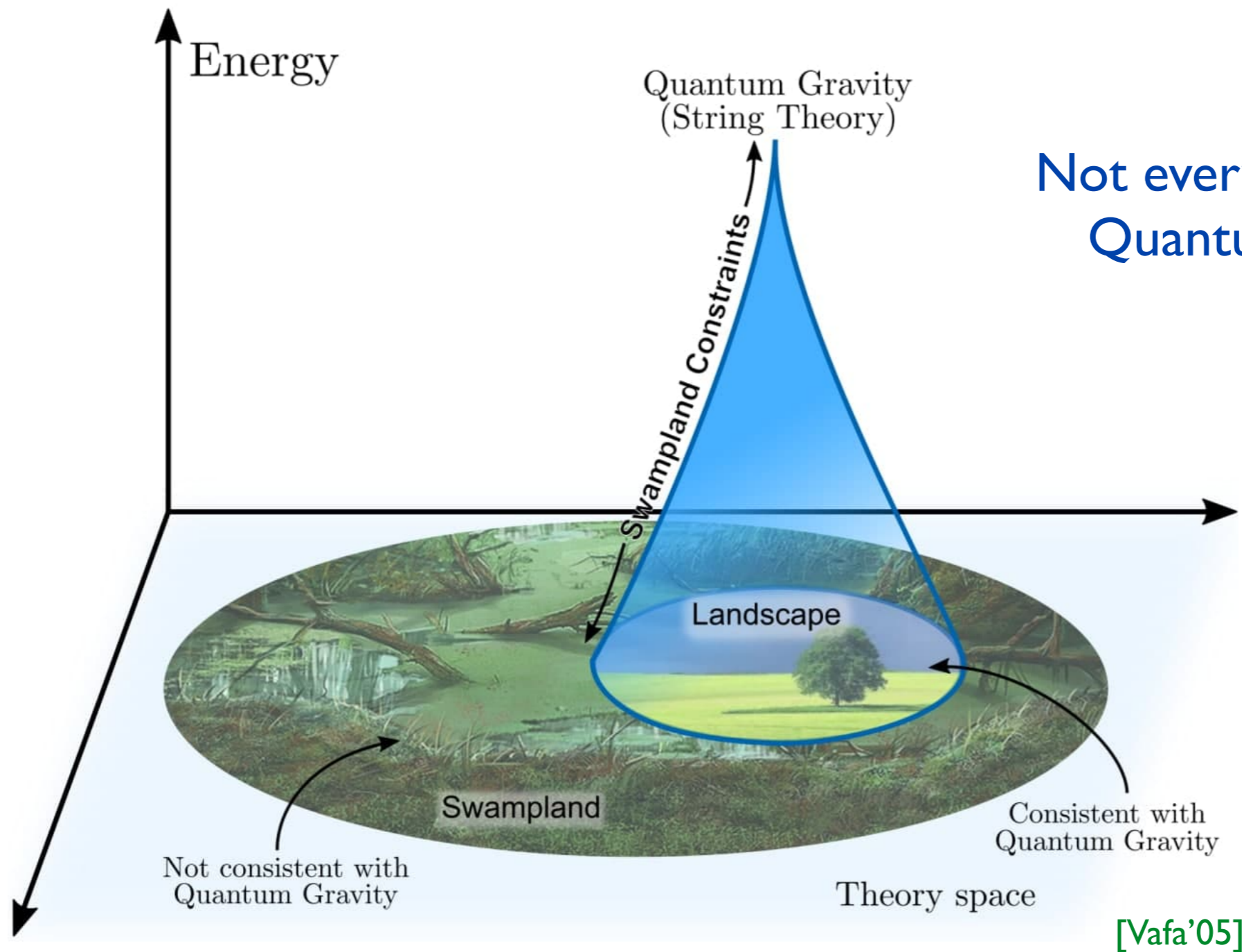
Swampland Lectures/Reviews:
[Brennan, Vafa'17] [Palti'19]
[Van Beest, Calderon-Infante, Mirfendereski, IV'21]
[Graña, Herraez'21]



Not everything goes in
Quantum Gravity!

Swampland Lectures/Reviews:
 [Brennan, Vafa'17] [Palti'19]
 [Van Beest, Calderon-Infante, Mirfendereski, IV'21]
 [Graña, Herraez'21]

[Vafa'05]



Swampland:

Apparently consistent (anomaly-free) quantum **effective field theories** that **cannot** be UV completed in **quantum gravity**

Swampland program

Goal:

Determine the constraints that an effective theory must satisfy to be consistent with quantum gravity

What distinguishes the landscape from the swampland?

Swampland program

Goal:

Determine the constraints that an effective theory must satisfy to be consistent with quantum gravity

What distinguishes the landscape from the swampland?



Universal UV imprint of quantum gravity at low energies

(New approach to connect string theory/quantum gravity to our world)

Swampland program

Goal:

Determine the constraints that an effective theory must satisfy to be consistent with quantum gravity

What distinguishes the landscape from the swampland?



Universal UV imprint of quantum gravity at low energies

(New approach to connect string theory/quantum gravity to our world)

Potential phenomenological implications:

New guiding principles to construct BSM models of Particle Physics and Cosmology

Failure of EFT expectations

These swampland constraints often look surprising from a low energy EFT perspective:

“What seems natural from UV perspective, might look unnatural from the IR perspective”

Failure of EFT expectations

These swampland constraints often look surprising from a low energy EFT perspective:

“What seems natural from UV perspective, might look unnatural from the IR perspective”

Two expectations/assumptions in EFT logic that breakdown in quantum gravity:

- 1) The entire space of parameters (consistent with symmetries) is a priori possible
- 2) Decoupling / Separation of energy scales

Failure of EFT expectations

These swampland constraints often look surprising from a low energy EFT perspective:

“What seems natural from UV perspective, might look unnatural from the IR perspective”

Two expectations/assumptions in EFT logic that breakdown in quantum gravity:

- 1) The entire space of parameters (consistent with symmetries) is a priori possible
- 2) Decoupling / Separation of energy scales

Gravity is different than other interactions! It can induce UV/IR mixing

Failure of EFT expectations

These swampland constraints often look surprising from a low energy EFT perspective:

“What seems natural from UV perspective, might look unnatural from the IR perspective”

Two expectations/assumptions in EFT logic that breakdown in quantum gravity:

- 1) The entire space of parameters (consistent with symmetries) is a priori possible
- 2) Decoupling / Separation of energy scales *golden opportunity!*

Gravity is different than other interactions! It can induce UV/IR mixing

Can we bring new insights to solve naturalness issues in our universe?

Swampland conjectures

Proposals for constraints that EFTs must satisfy to be consistent with QG

They are mainly **motivated by string theory and black hole physics**, but we expect them to be general features of quantum gravity (even beyond string theory)

Swampland conjectures

Proposals for constraints that EFTs must satisfy to be consistent with QG

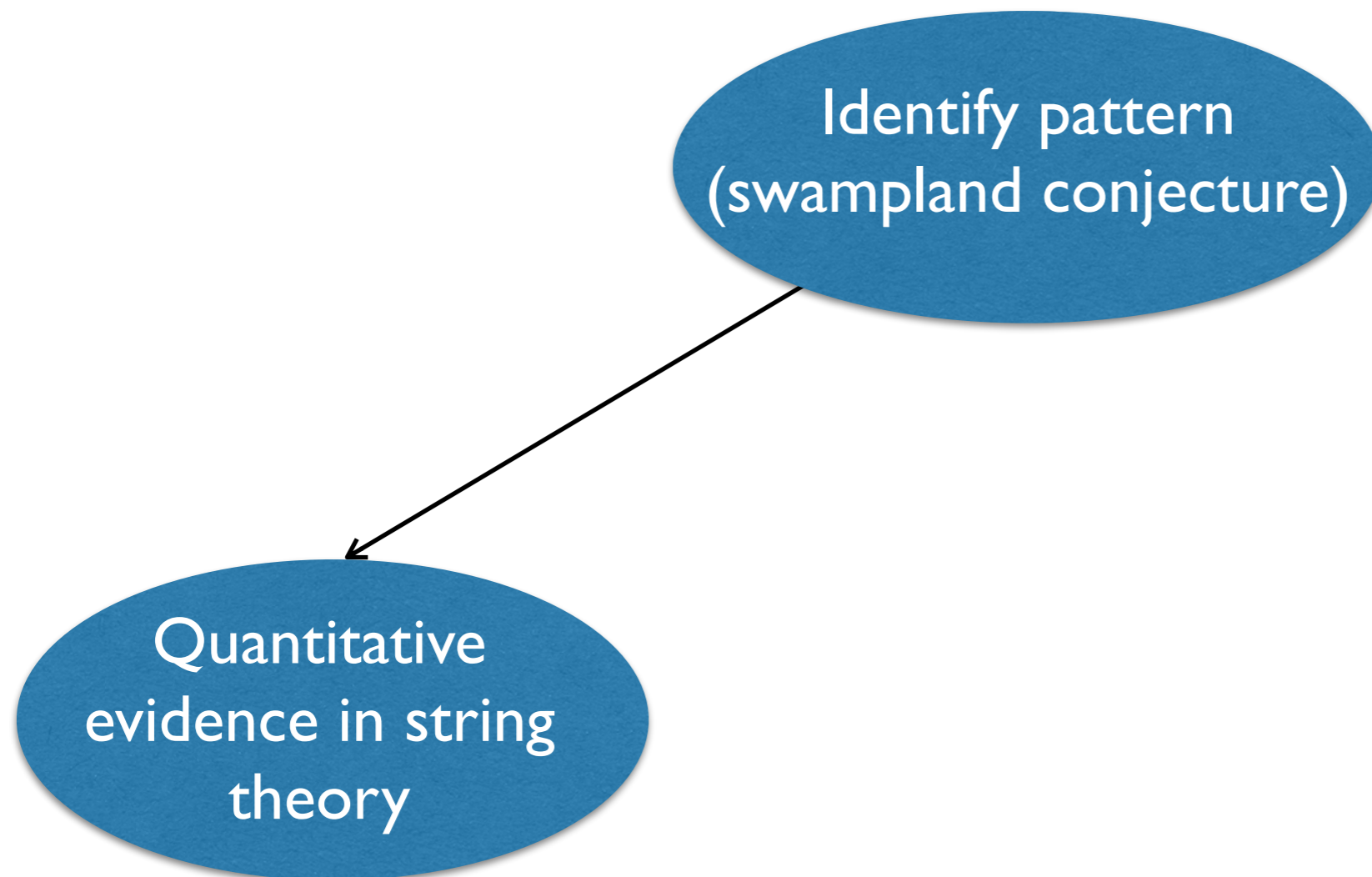
They are mainly **motivated by string theory and black hole physics**, but we expect them to be general features of quantum gravity (even beyond string theory)

Identify pattern
(swampland conjecture)

Swampland conjectures

Proposals for constraints that EFTs must satisfy to be consistent with QG

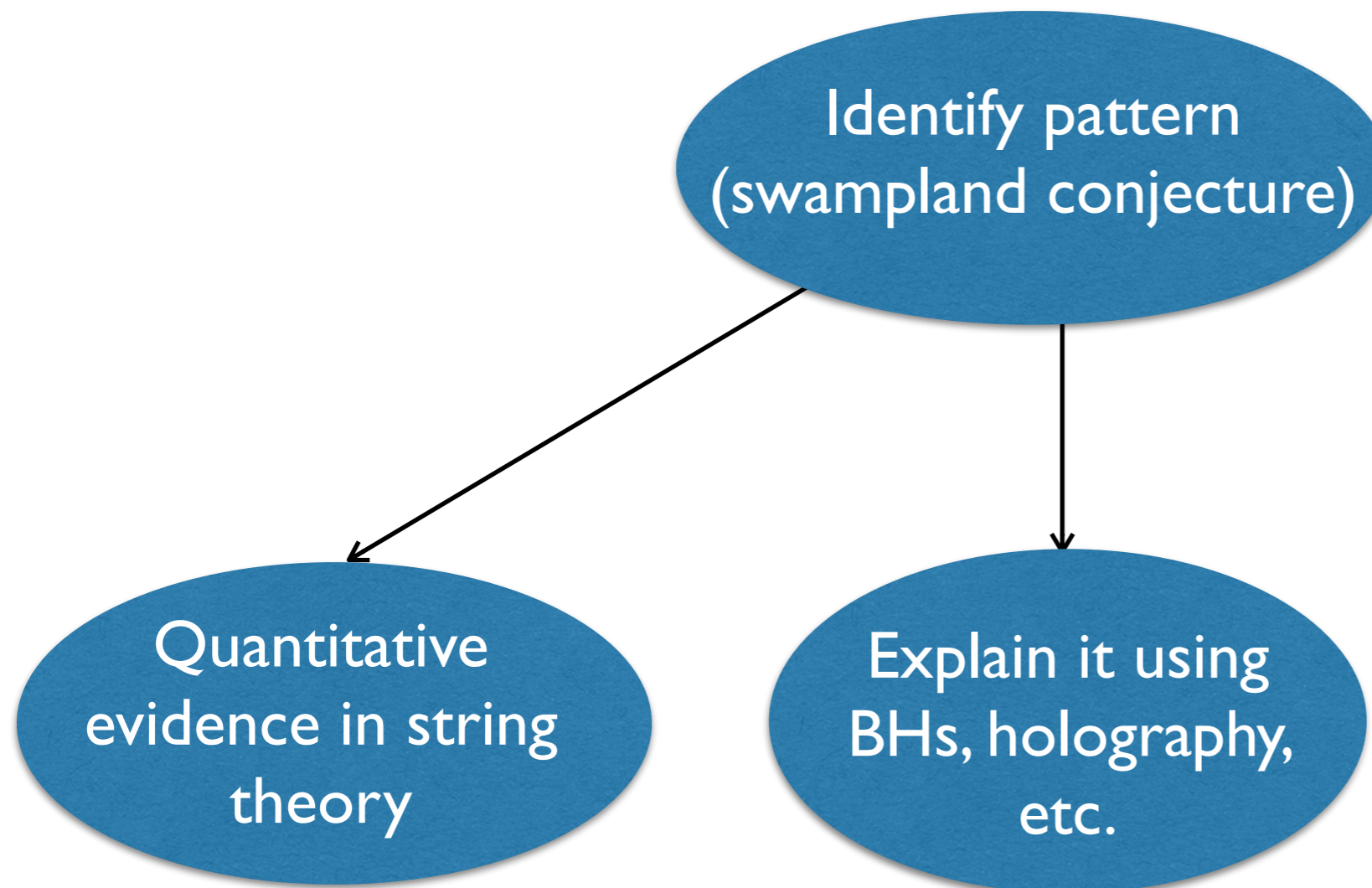
They are mainly **motivated by string theory and black hole physics**, but we expect them to be general features of quantum gravity (even beyond string theory)



Swampland conjectures

Proposals for constraints that EFTs must satisfy to be consistent with QG

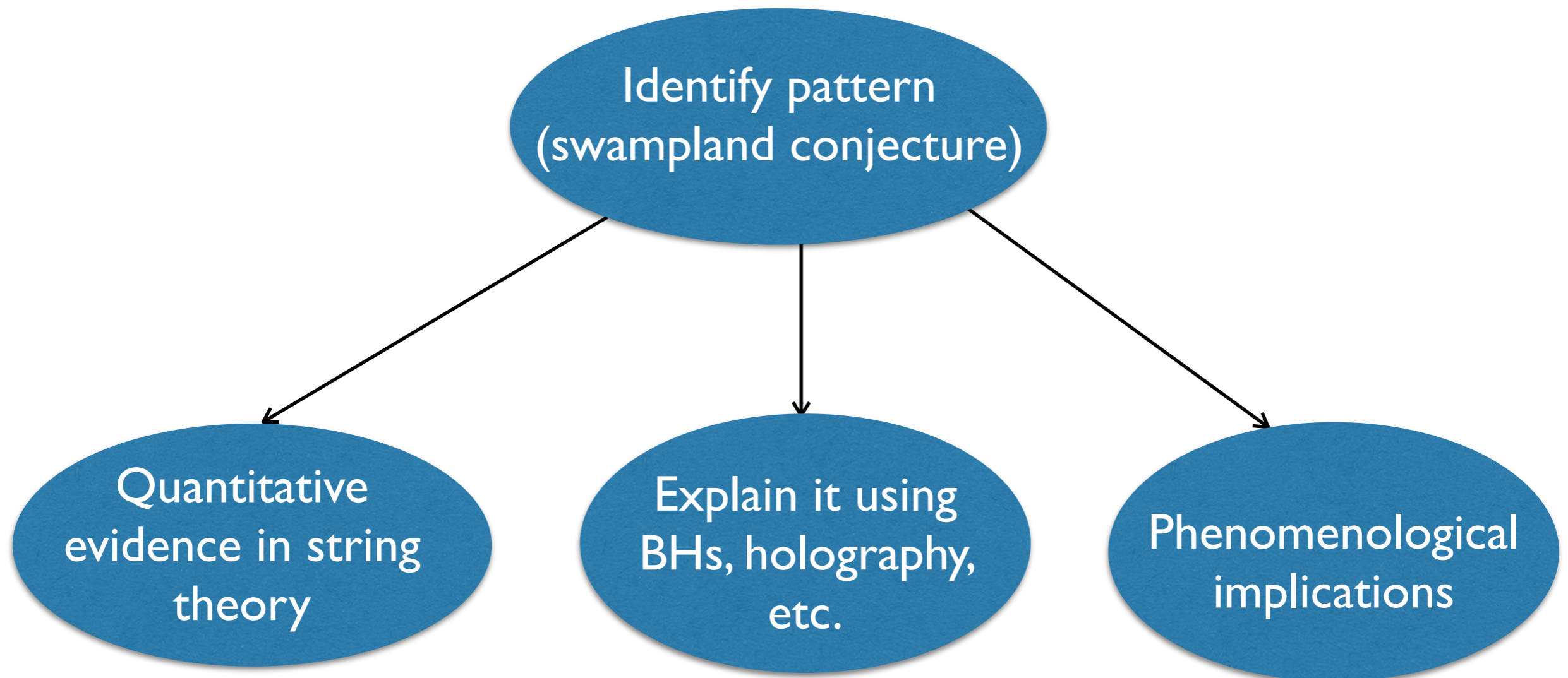
They are mainly **motivated by string theory and black hole physics**, but we expect them to be general features of quantum gravity (even beyond string theory)



Swampland conjectures

Proposals for constraints that EFTs must satisfy to be consistent with QG

They are mainly **motivated by string theory and black hole physics**, but we expect them to be general features of quantum gravity (even beyond string theory)



Swampland conjectures

Completeness
hypothesis

No global
symmetries

Weak Gravity
Conjecture

Distance
Conjecture

Swampland conjectures

Completeness hypothesis

No global symmetries

Cobordism conjecture

Weak Gravity Conjecture

Distance Conjecture

...

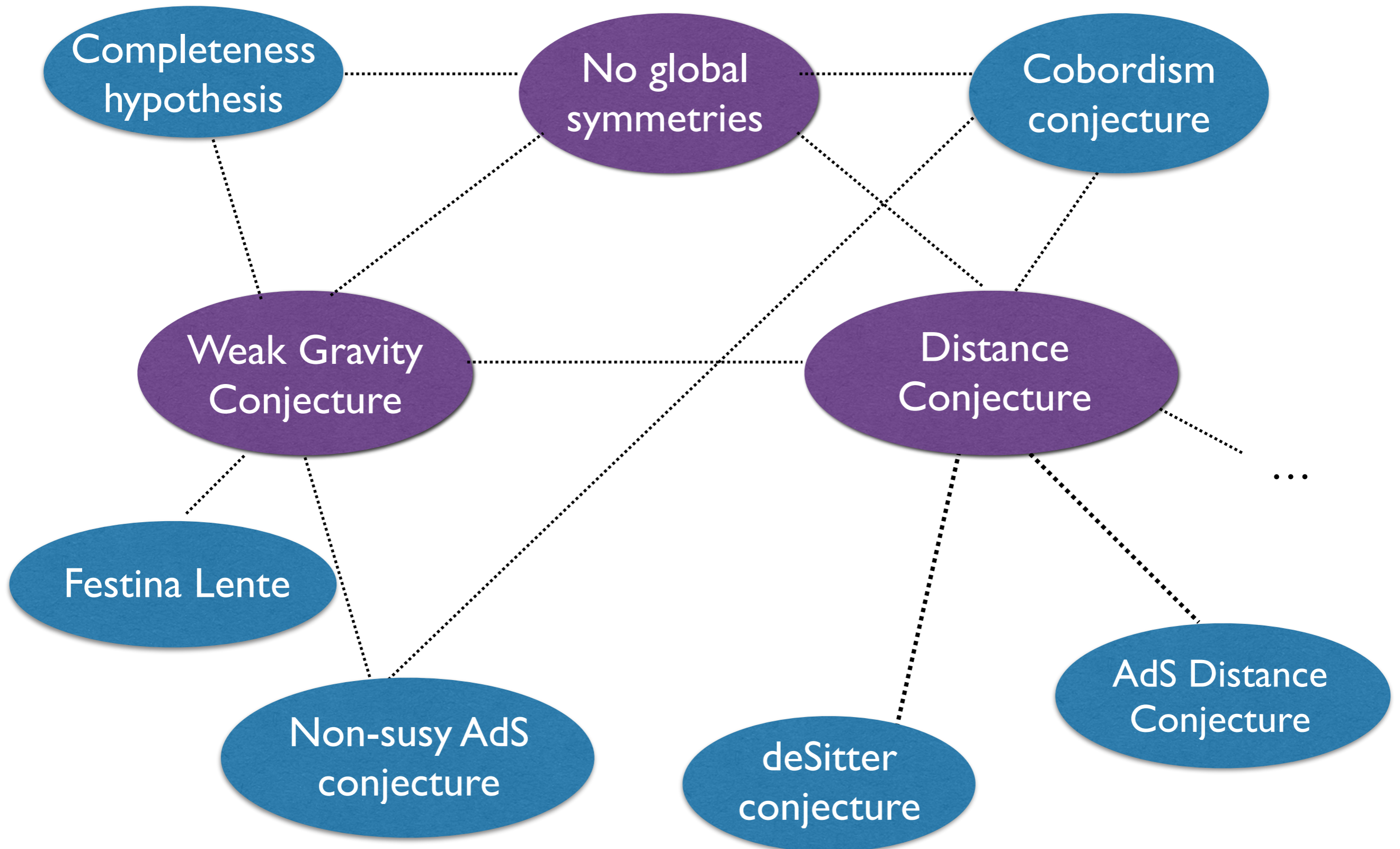
Festina Lente

Non-susy AdS conjecture

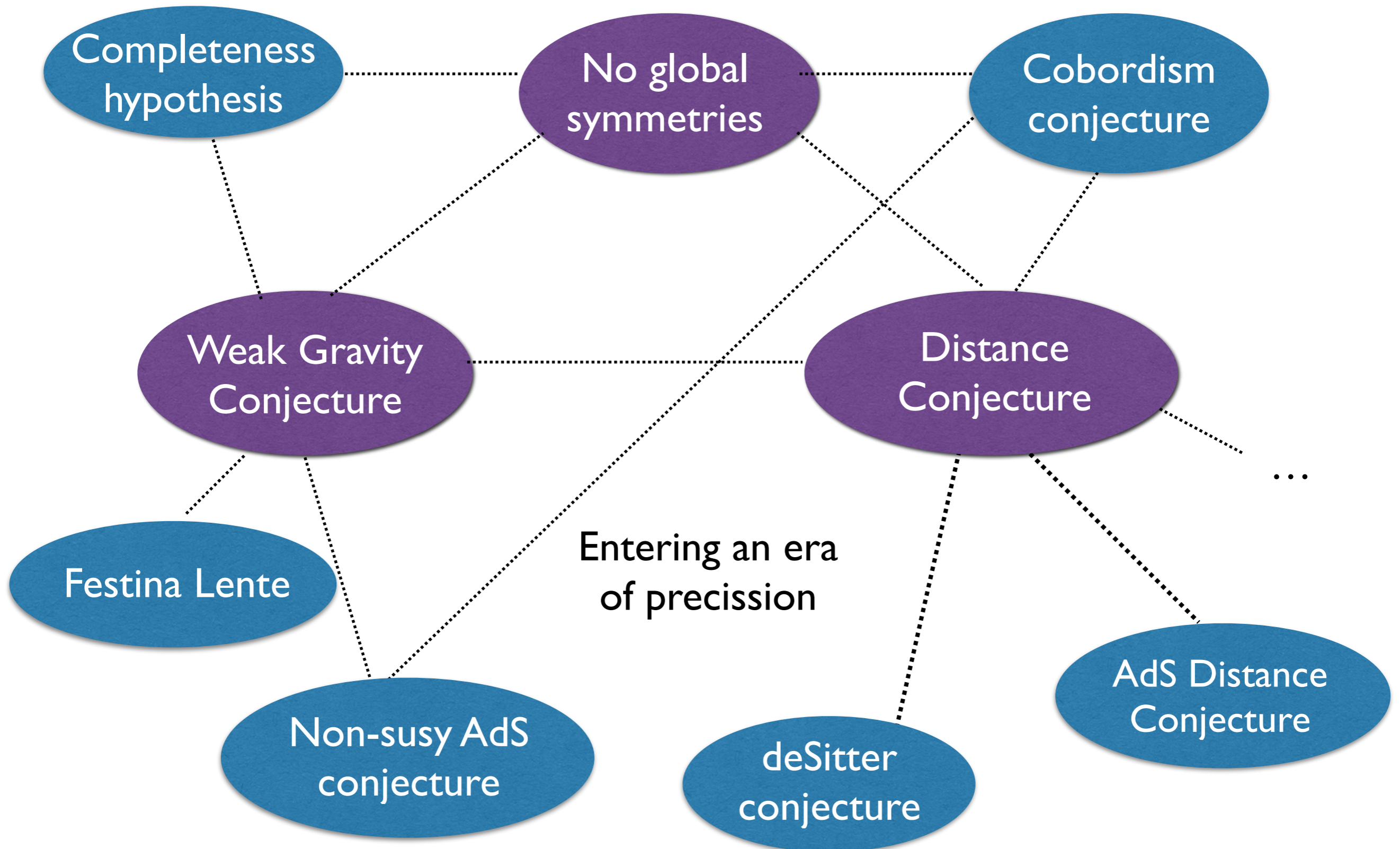
deSitter conjecture

AdS Distance Conjecture

Swampland conjectures



Swampland conjectures

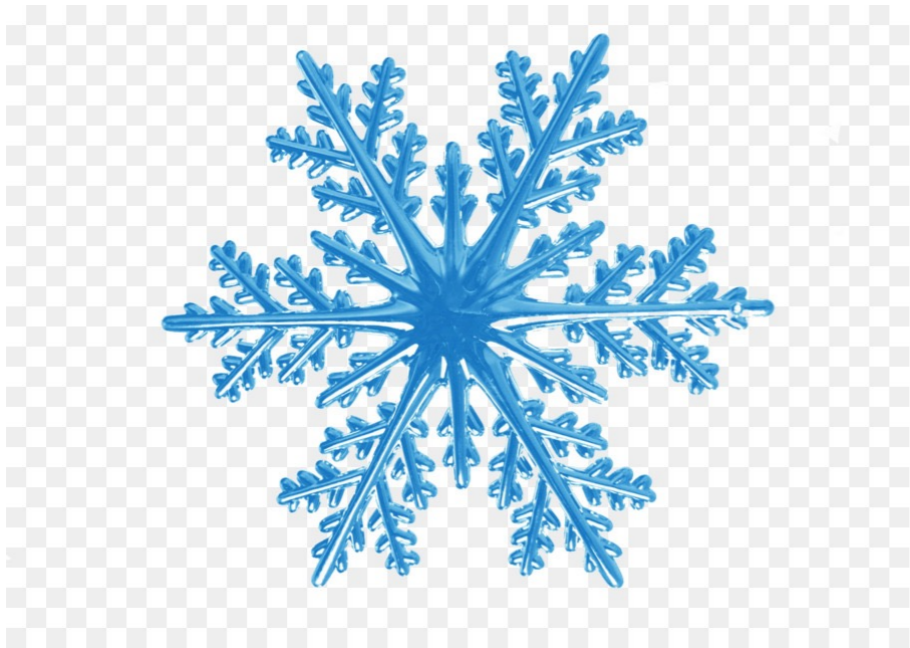


Swampland conjectures

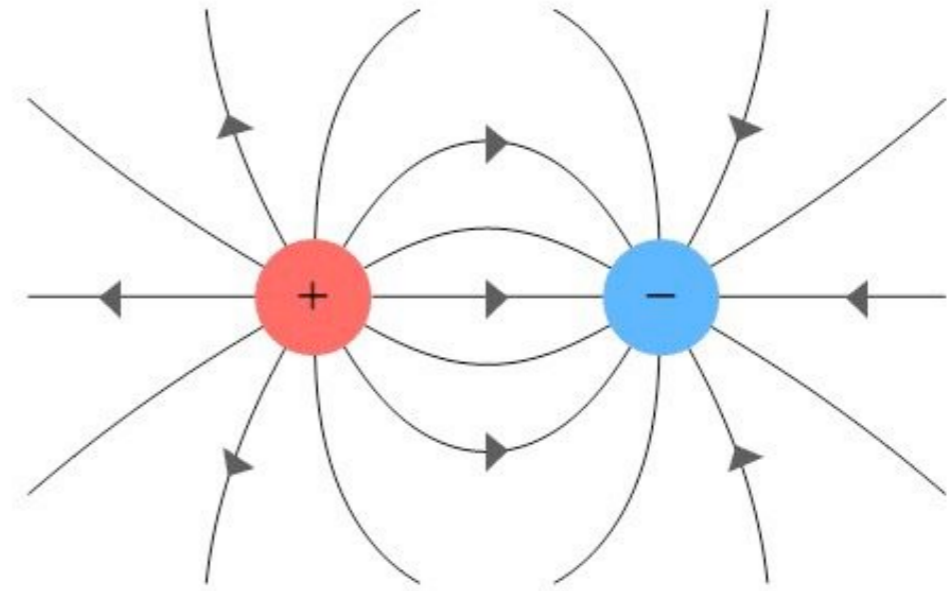
No global
symmetries

Symmetries

Rotational symmetry



Electric charge

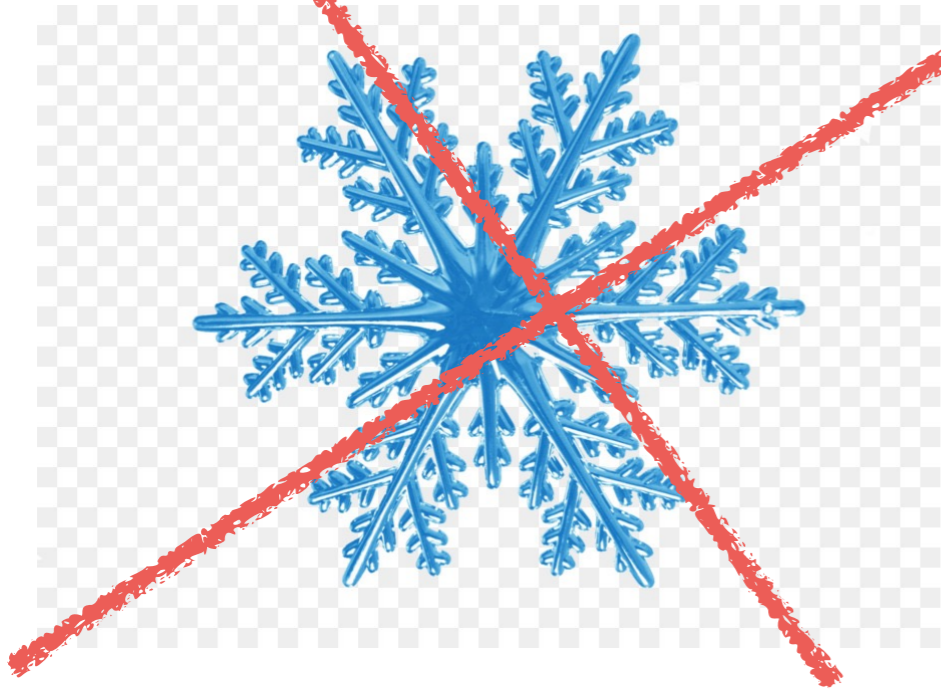


Angular momentum

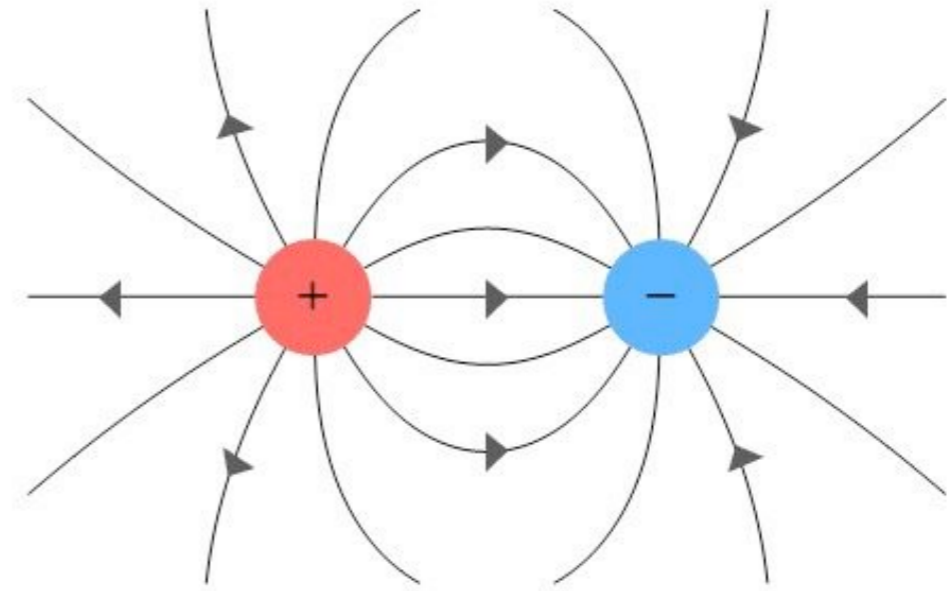


Symmetries

Rotational symmetry



Electric charge



Angular momentum



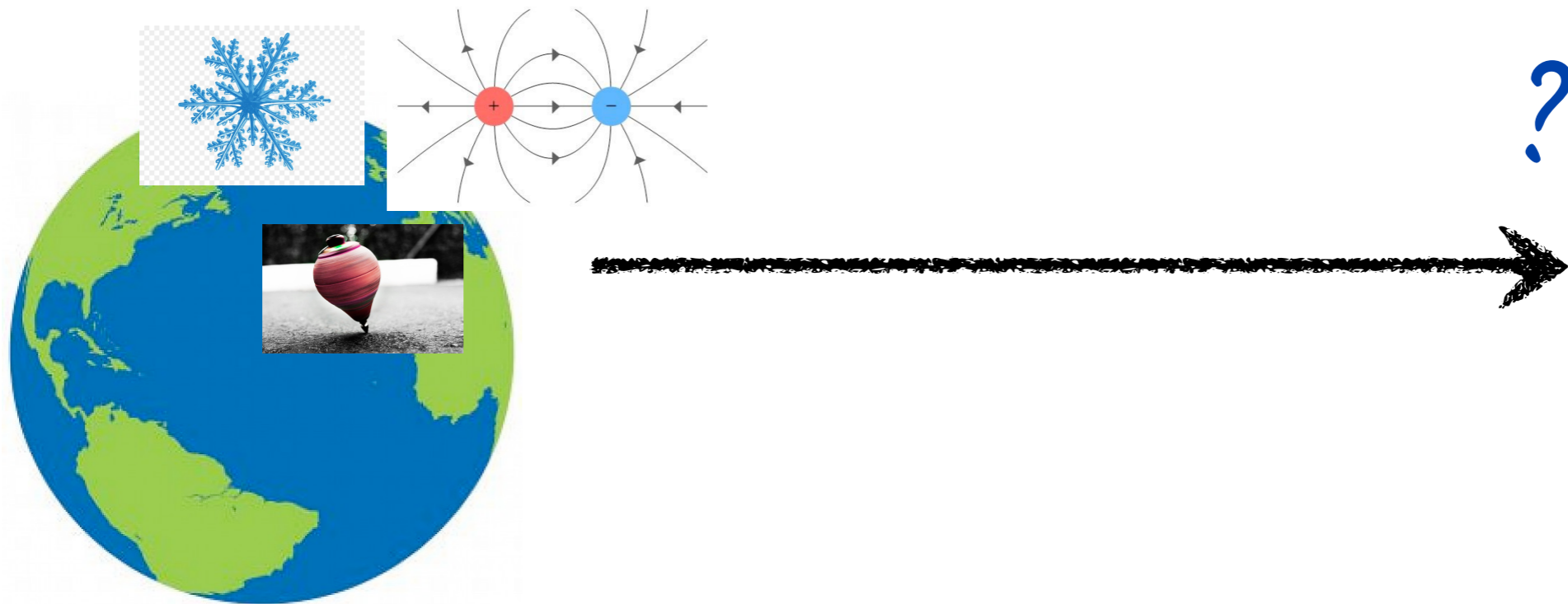
Symmetries

Two types of symmetries: global vs gauged

Symmetries

Two types of symmetries: global vs gauged

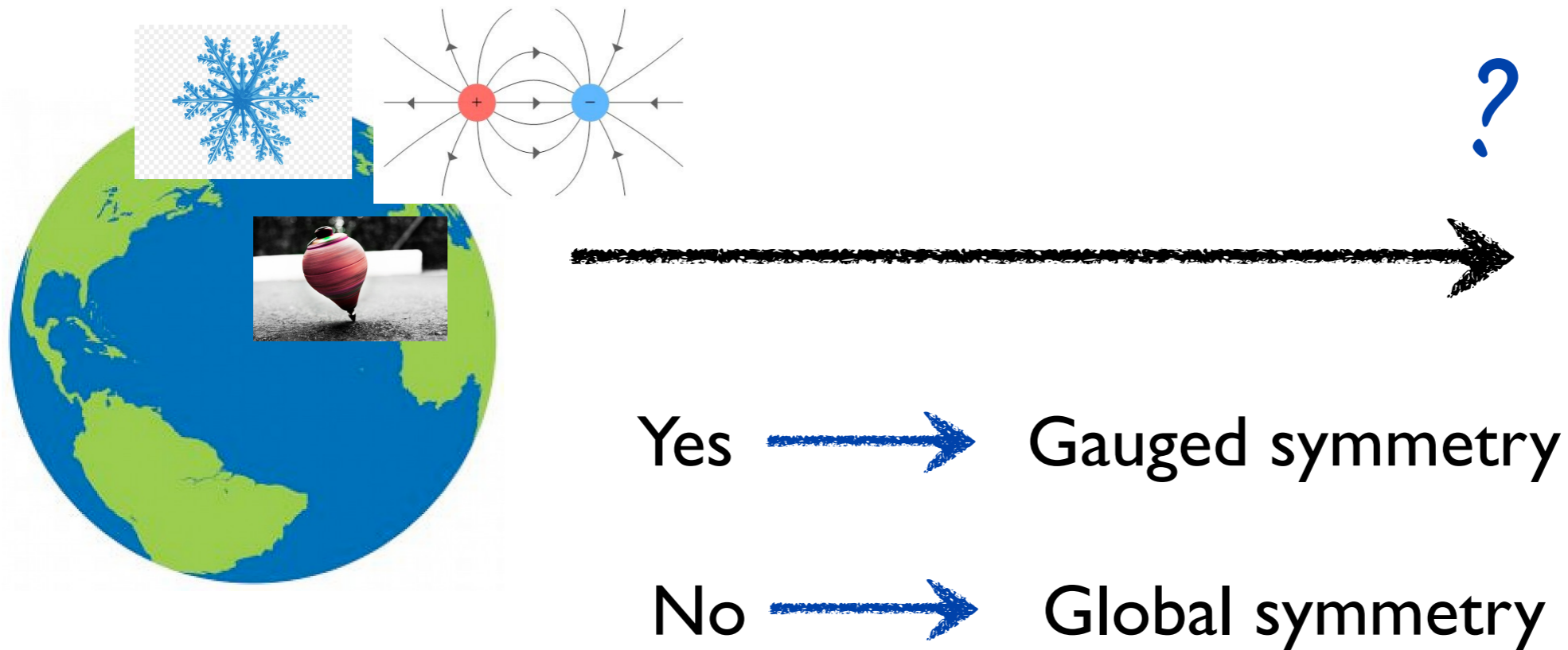
Can we measure the conserved charge from very far away?



Symmetries

Two types of symmetries: global vs gauged

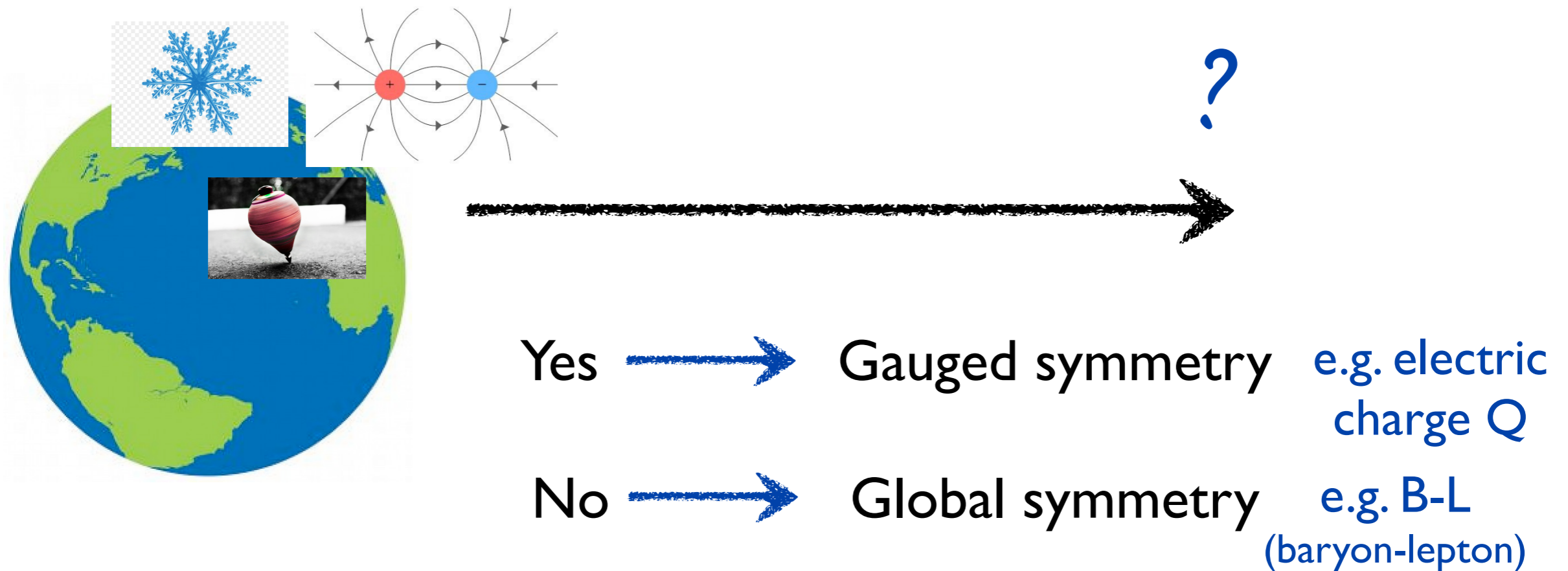
Can we measure the conserved charge from very far away?



Symmetries

Two types of symmetries: global vs gauged

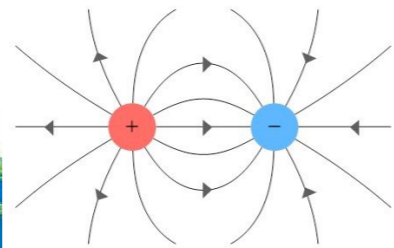
Can we measure the conserved charge from very far away?



Symmetries

Two types of symmetries: global vs gauged

Can we measure the conserved charge from very far away?



?



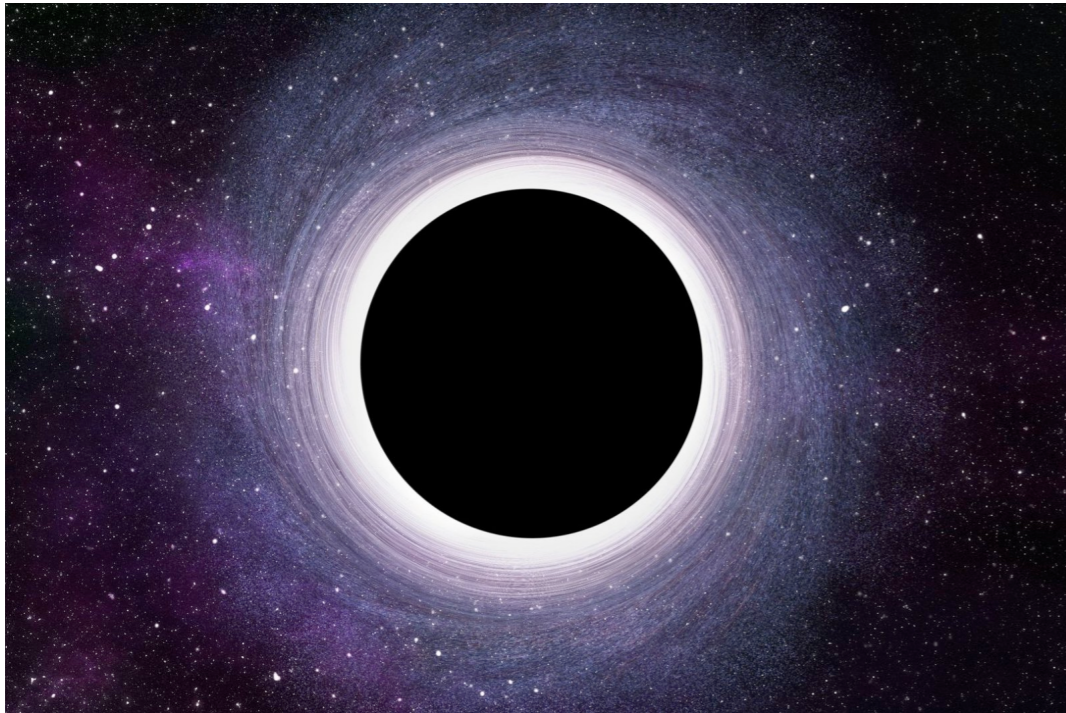
Yes \longrightarrow Gauged symmetry e.g. electric charge Q

No \longrightarrow Global symmetry e.g. B-L
(baryon-lepton)

They are broken in quantum gravity

No global symmetries conjecture

Heuristic motivation from black holes:



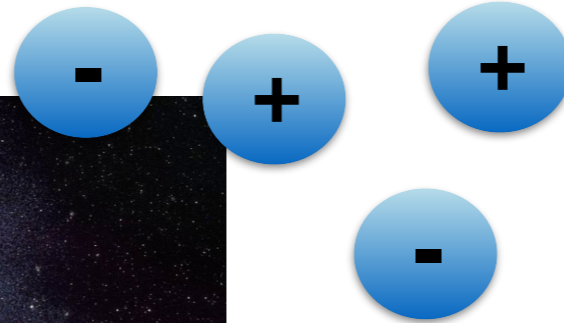
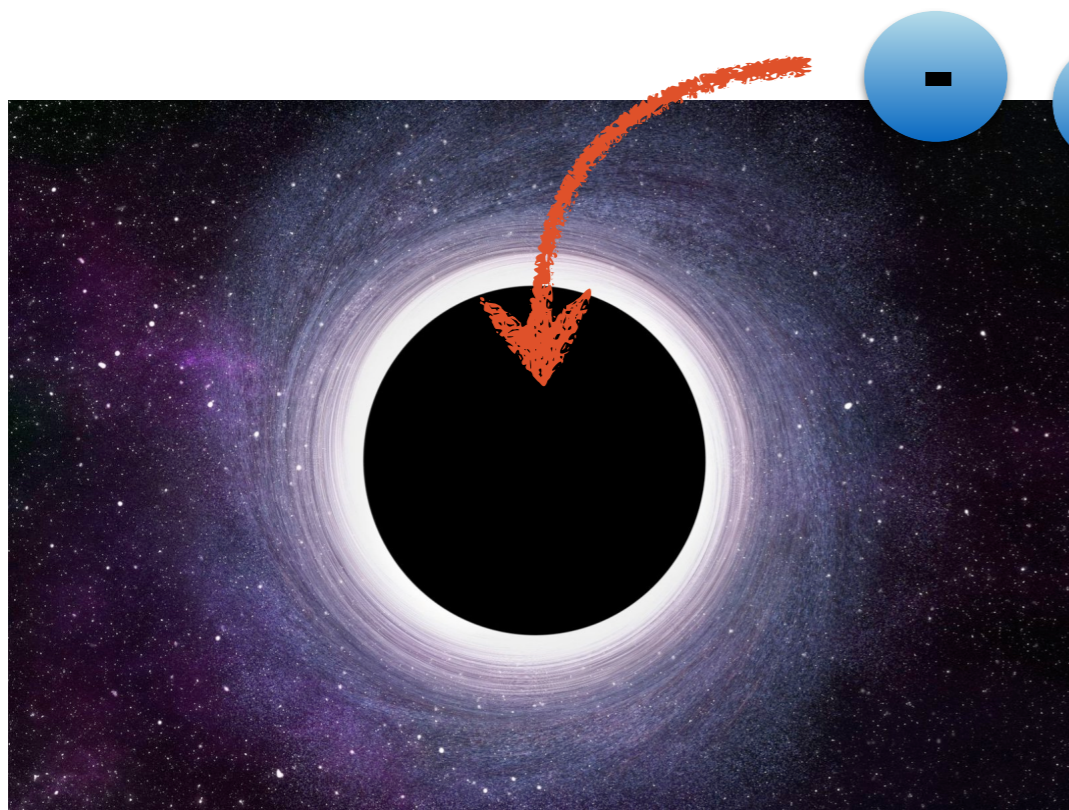
Black holes have
finite entropy



(“how many different ways lead to
the same black hole exterior”)

No global symmetries conjecture

Heuristic motivation from black holes:

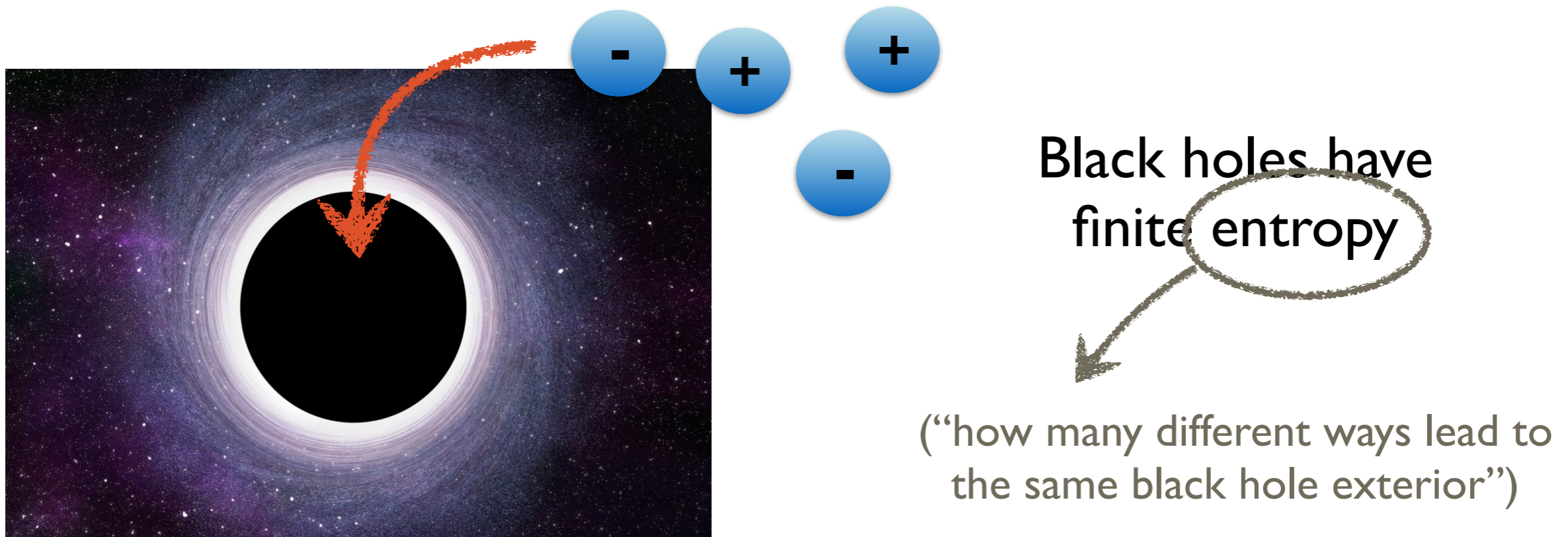


Black holes have
finite entropy

(“how many different ways lead to
the same black hole exterior”)

No global symmetries conjecture

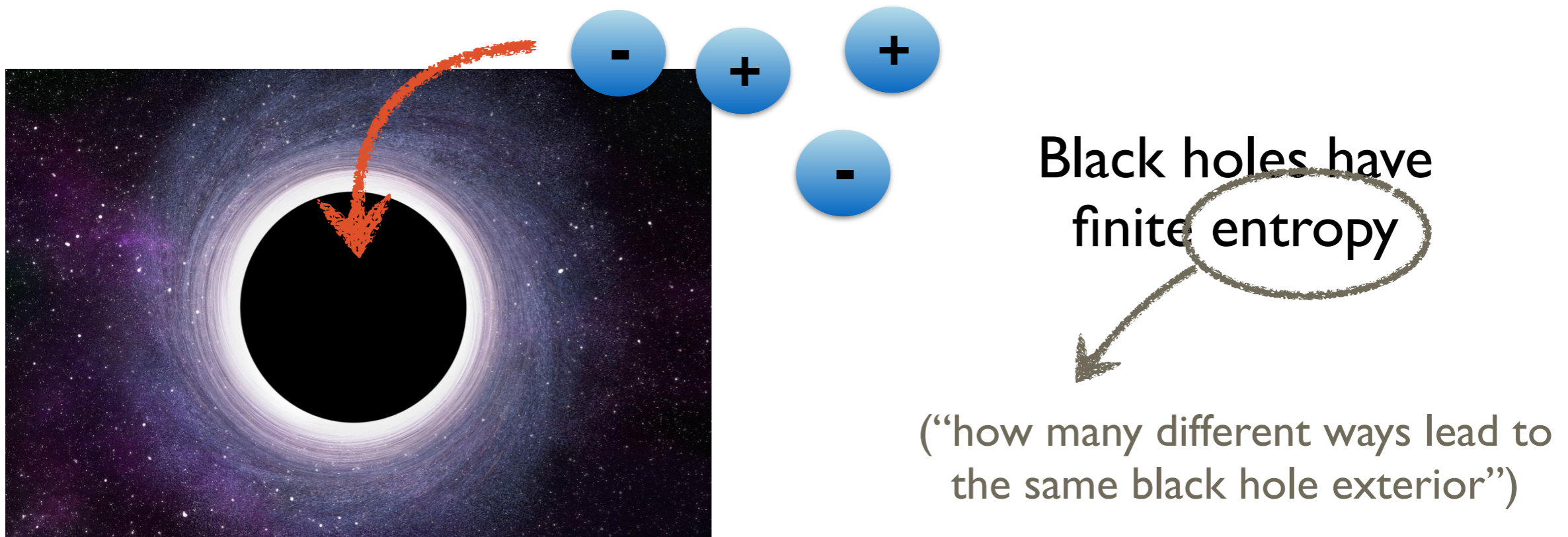
Heuristic motivation from black holes:



Since global symmetries cannot be detected from far away,

No global symmetries conjecture

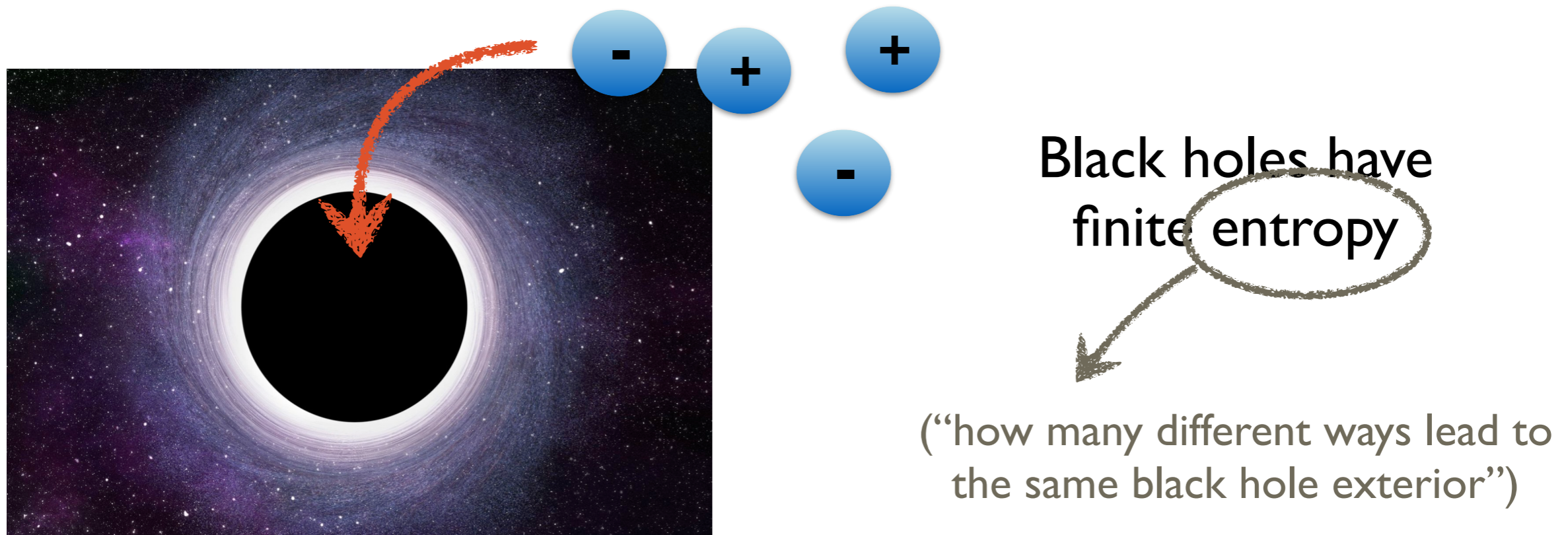
Heuristic motivation from black holes:



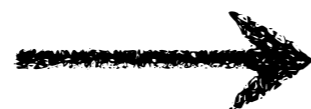
Since global symmetries cannot be detected from far away, we could have infinitely many black holes with different values of the global charge, but that look the same from far way

No global symmetries conjecture

Heuristic motivation from black holes:



Since global symmetries cannot be detected from far away, we could have infinitely many black holes with different values of the global charge, but that look the same from far way



Infinite entropy

contradiction!

No global symmetries conjecture

Global symmetries cannot be exact in quantum gravity
(unless they are gauged)

[Banks-Dixon'88] [Horowitz, Strominger, ...] [Susskind] [Banks, Seiberg'11]

Evidence:

- Proof in perturbative string theory [Polchinski's book]
- Proof in AdS/CFT [Harlow, Ooguri '18]
- Correlation to unitary black hole evaporation (and topology changing processes) [Harlow, Shaghoulian '20] [Chen, Lin '20] [Hsin et al '20] [Yonekura '20] [Bah, Chen, Maldacena'22]

Global symmetries are not well defined in quantum gravity as the topology itself fluctuates

Swampland conjectures

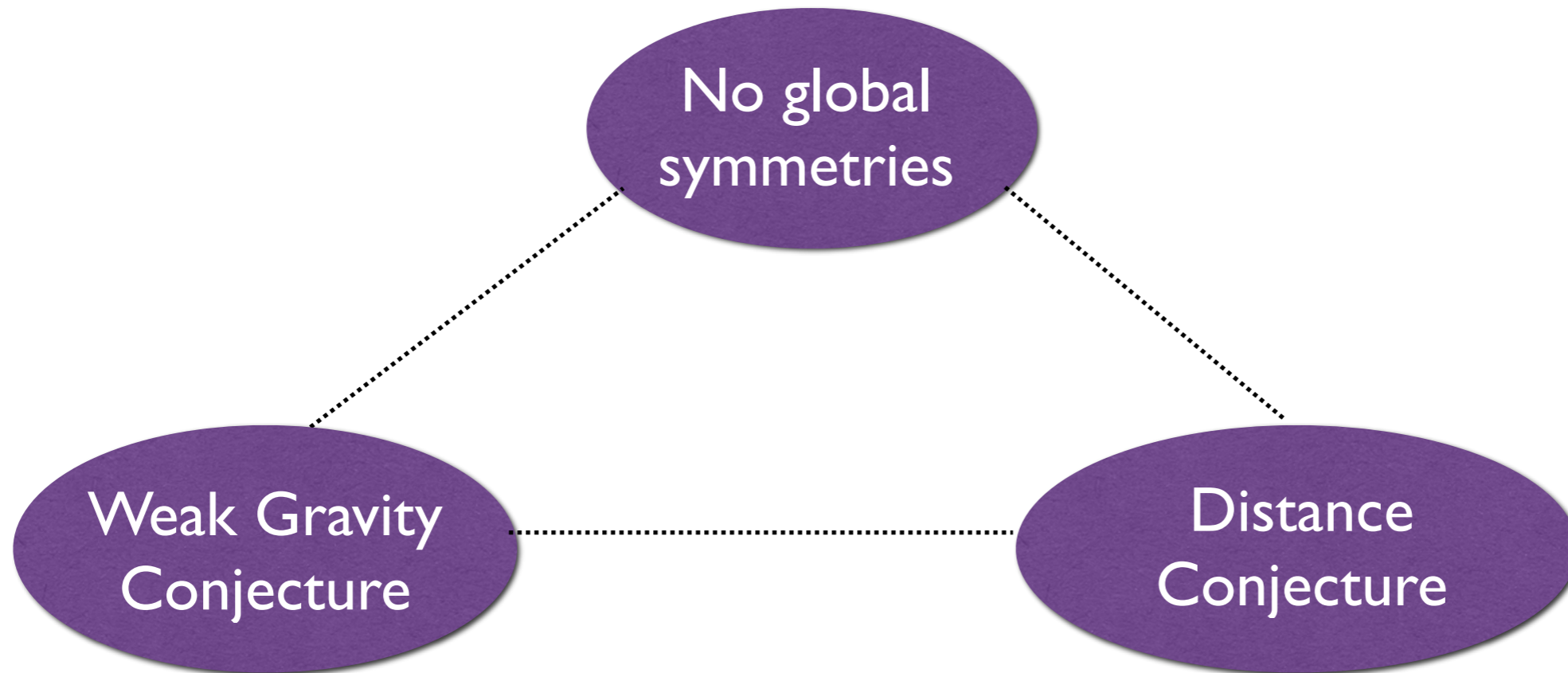
No global
symmetries

Open questions:

How much are the symmetries broken?

Or if they are gauged, how small can the gauge coupling be?

Swampland conjectures



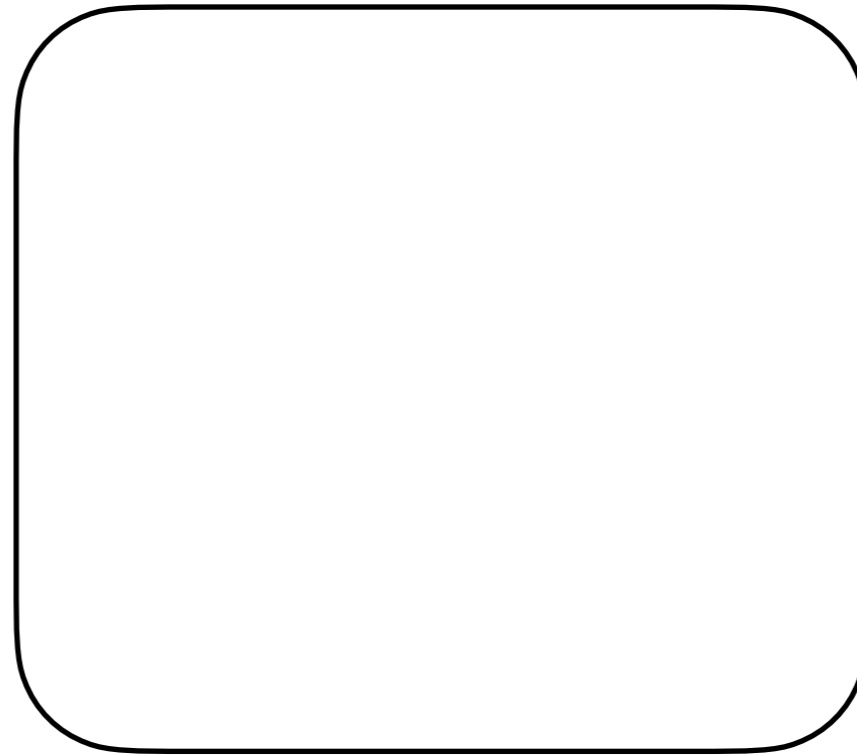
Open questions:

How much are the symmetries broken?

Or if they are gauged, how small can the gauge coupling be?

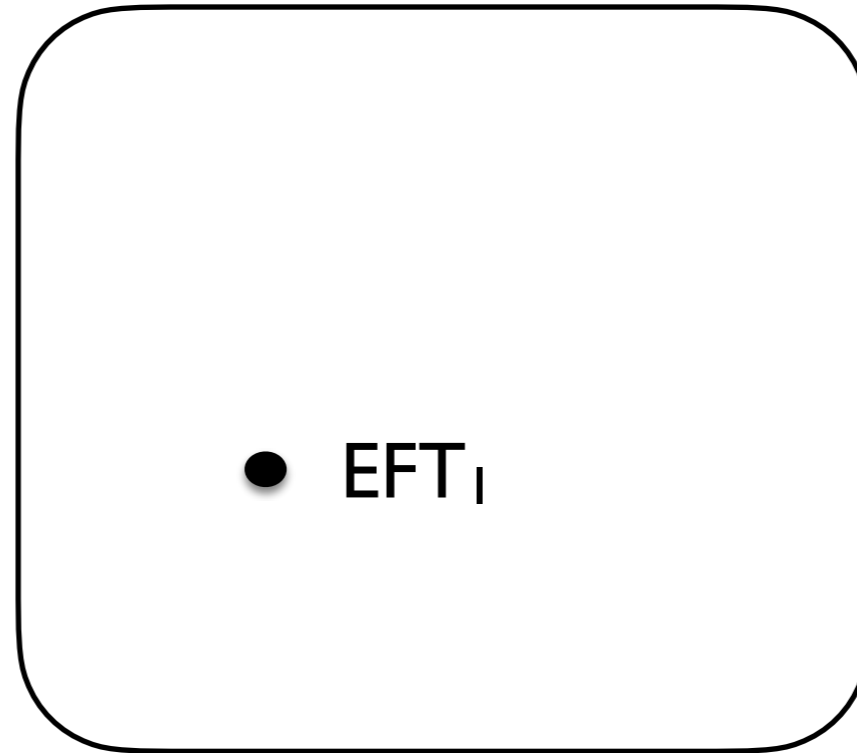
Approximate global symmetries

Parameter space:



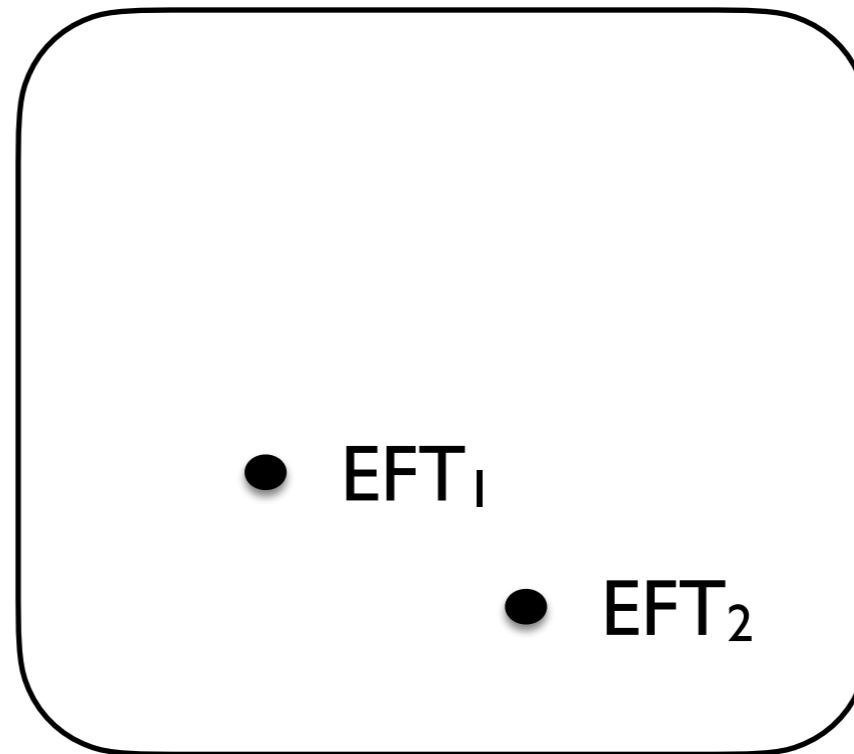
Approximate global symmetries

Parameter space:



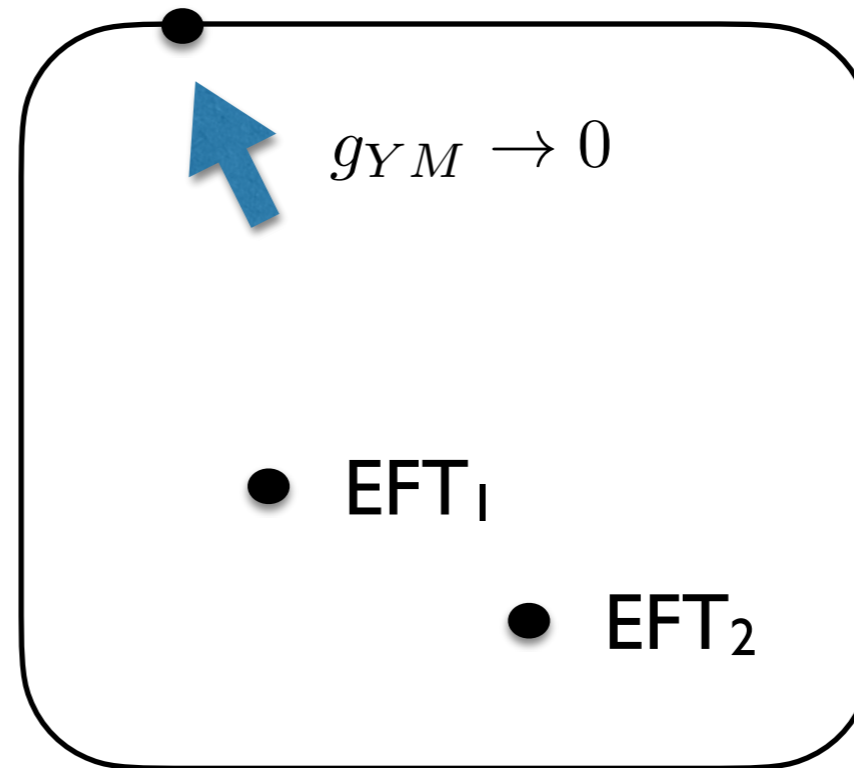
Approximate global symmetries

Parameter space:



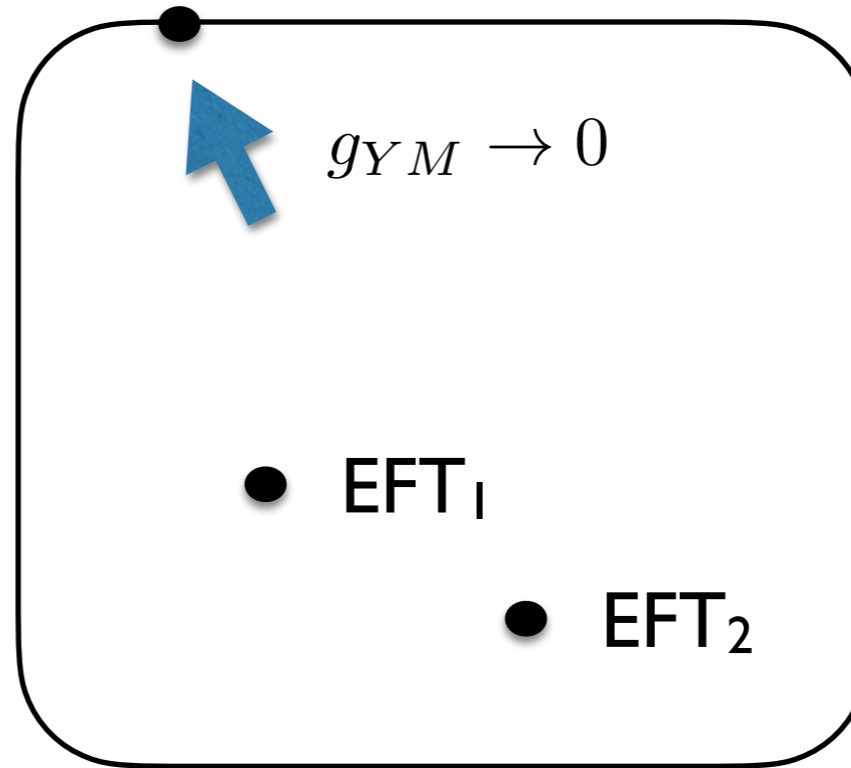
Approximate global symmetries

Parameter space:



Approximate global symmetries

Parameter space:

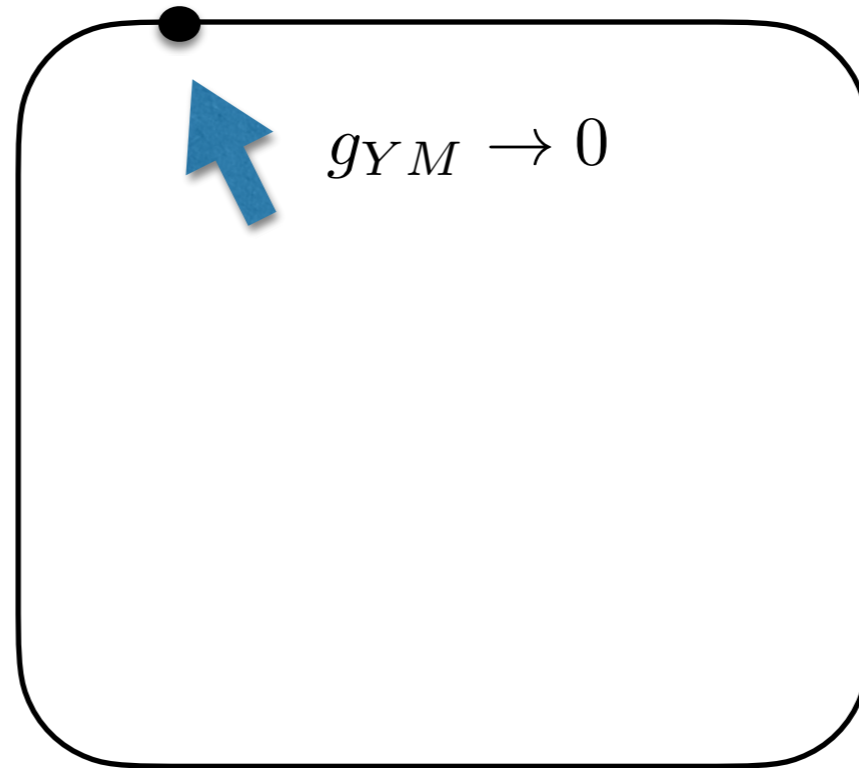


Assume some **global symmetry can be restored in a continuous way** at some special points of the parameter space

e.g. by sending gauge coupling $g_{YM} \rightarrow 0$
we restore a $U(1)$ global symmetry

Approximate global symmetries

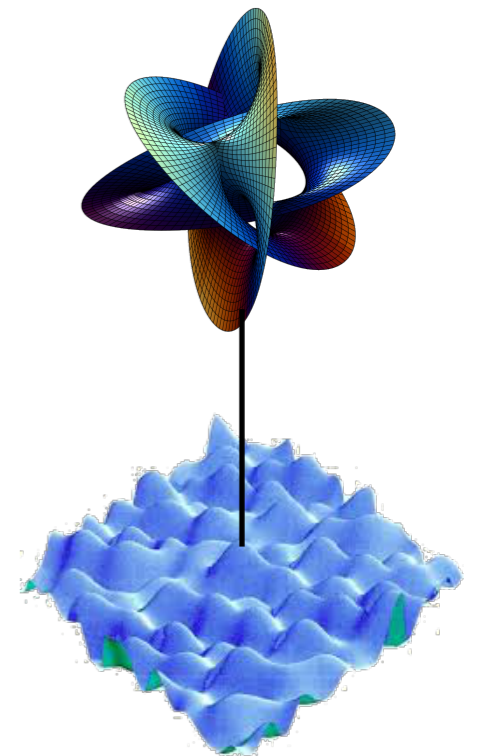
Parameter space:



String Theory has no free parameters:

All masses/couplings are given by vacuum expectation values of scalar fields that fix the size/shape of extra dimensions

e.g. $g_{YM}(\phi)$



Approximate global symmetries

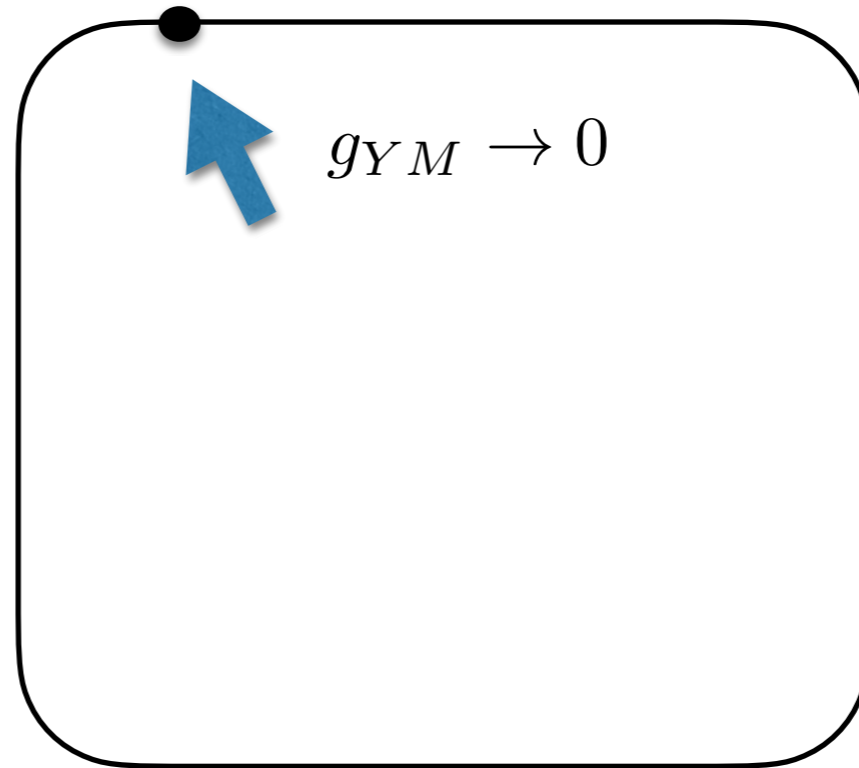
Parameter space:

=

Scalar field space
in String Theory

$$\mathcal{L} = g_{ij}(\phi) \partial\phi^i \partial\phi^j$$

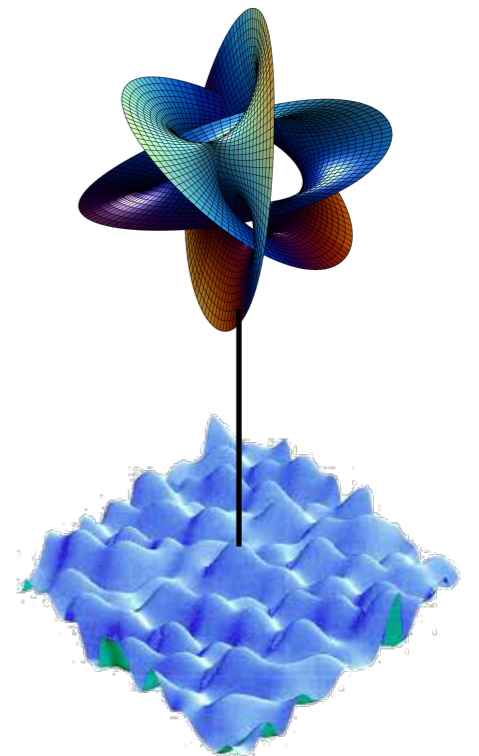
↪ field metric



String Theory has no free parameters:

All masses/couplings are given by vacuum expectation values of scalar fields that fix the size/shape of extra dimensions

e.g. $g_{YM}(\phi)$



Approximate global symmetries

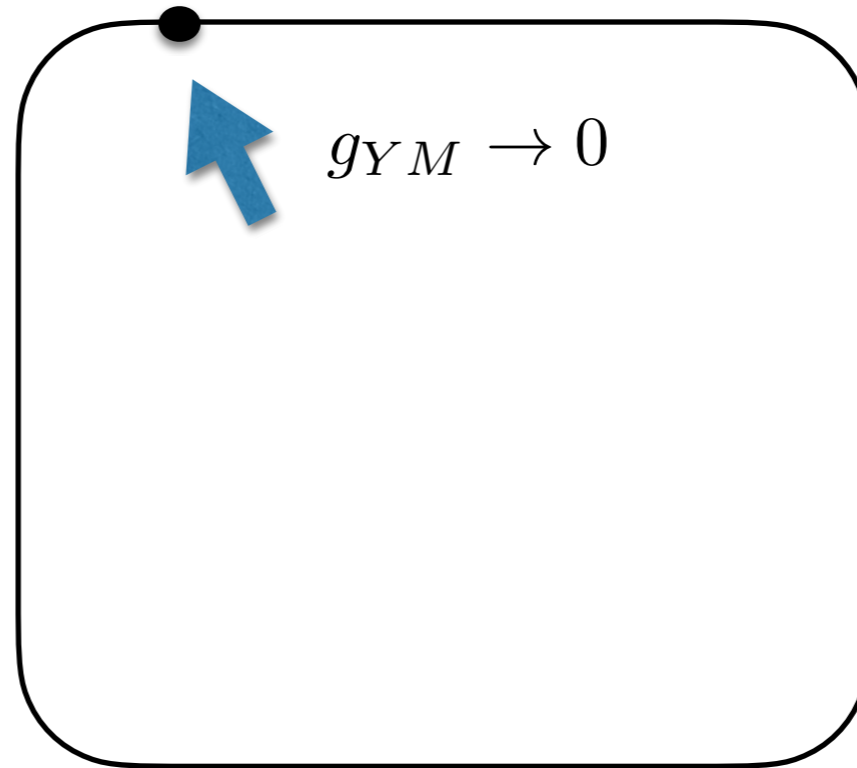
Parameter space:

=

Scalar field space
in String Theory

$$\mathcal{L} = g_{ij}(\phi) \partial\phi^i \partial\phi^j$$

↪ field metric

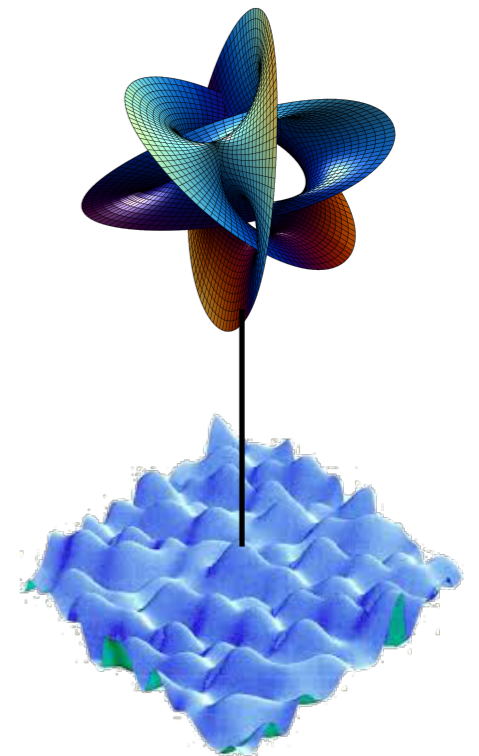


String Theory has no free parameters:

All masses/couplings are given by vacuum expectation values of scalar fields that fix the size/shape of extra dimensions

e.g. $g_{YM}(\phi)$ (like the Higgs boson parametrizes the masses)

$$m(H) = y\langle H \rangle$$



Approximate global symmetries

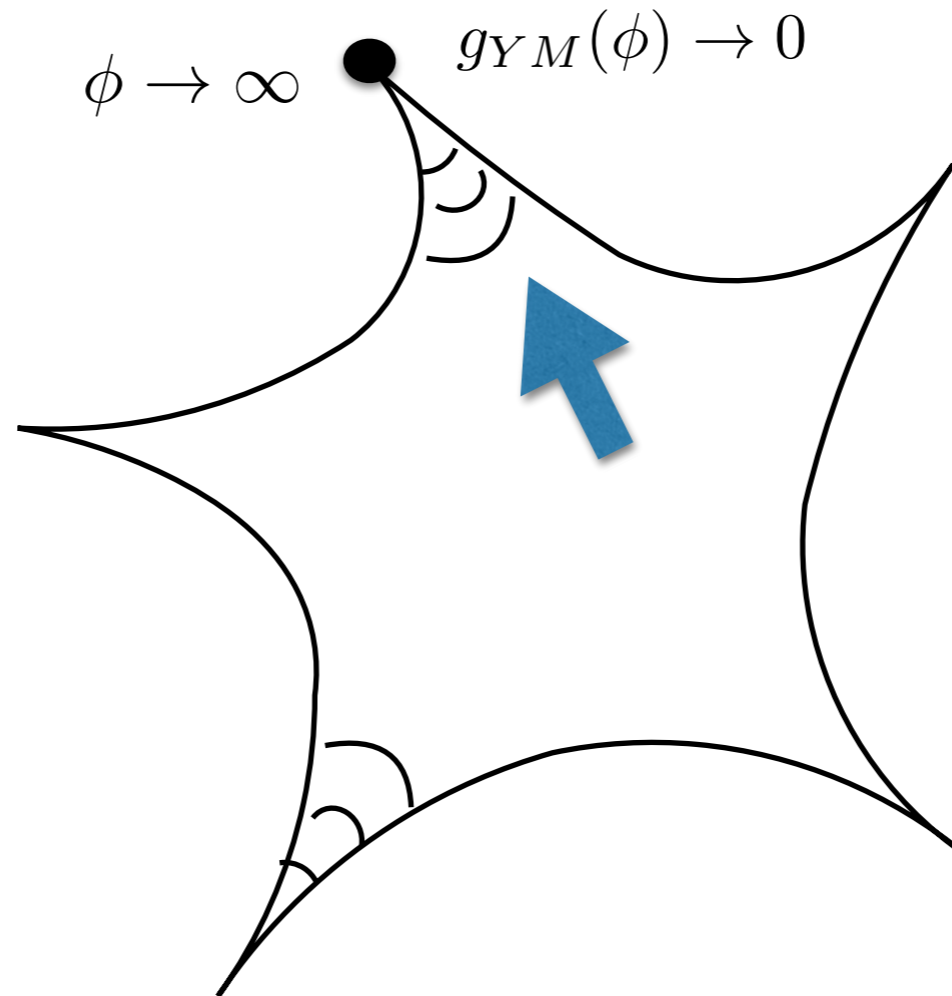
Parameter space:

=

Scalar field space
in String Theory

$$\mathcal{L} = g_{ij}(\phi) \partial\phi^i \partial\phi^j$$

↪ field metric

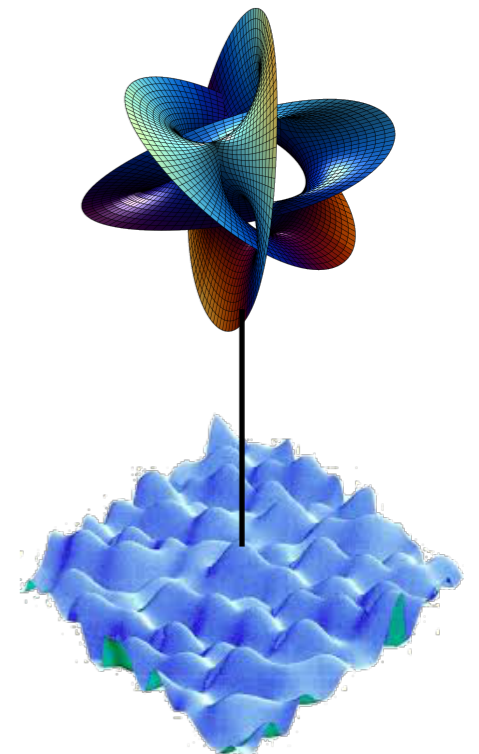


String Theory has no free parameters:

All masses/couplings are given by vacuum expectation values of scalar fields that fix the size/shape of extra dimensions

e.g. $g_{YM}(\phi)$ (like the Higgs boson parametrizes the masses)

$$m(H) = y \langle H \rangle$$



Approximate global symmetries

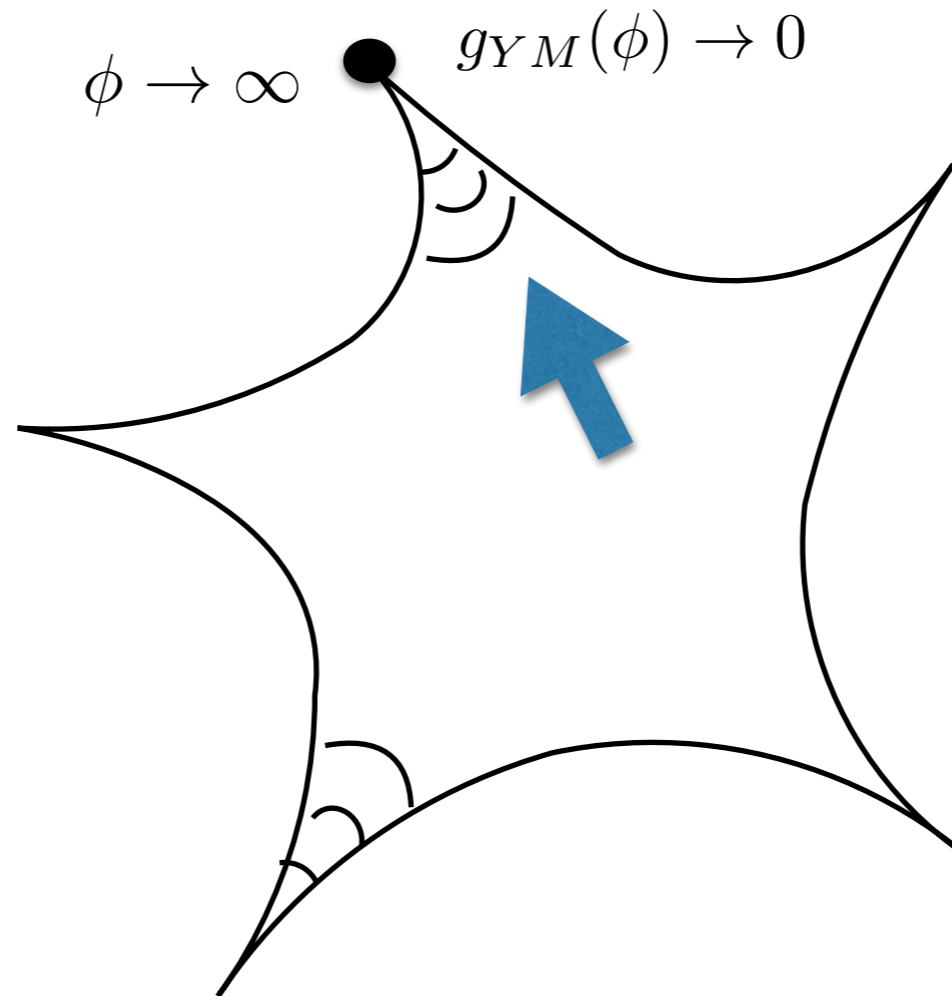
Parameter space:

=

Scalar field space
in String Theory

$$\mathcal{L} = g_{ij}(\phi) \partial\phi^i \partial\phi^j$$

↪ field metric



approx global
symmetry

Global symmetries not
allowed in quantum gravity



They can only be restored at infinite
field distance (asymptotic limits)

Approximate global symmetries

Parameter space:

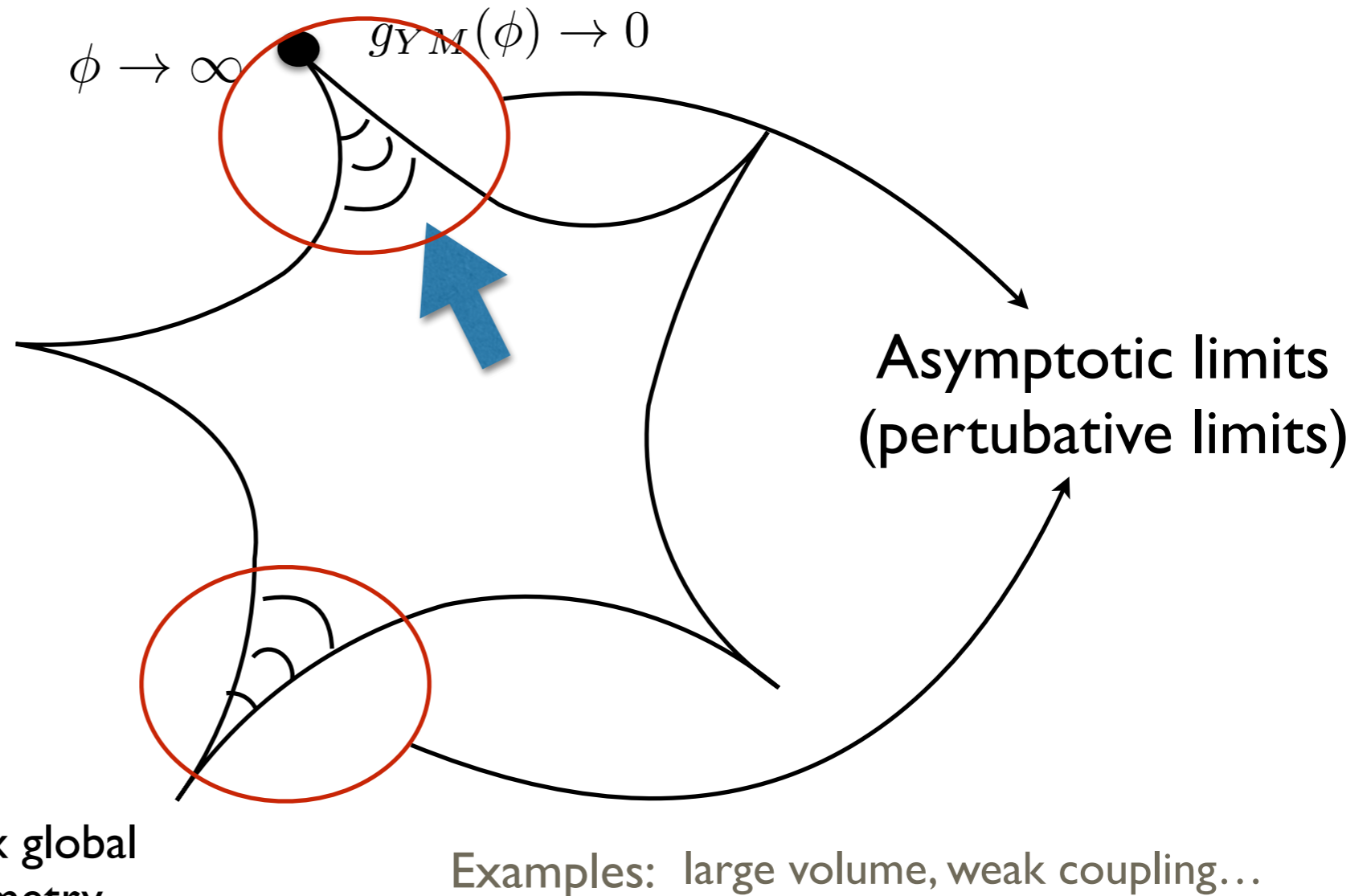
=

Scalar field space
in String Theory

$$\mathcal{L} = g_{ij}(\phi) \partial\phi^i \partial\phi^j$$

↪ field metric

approx global
symmetry



Global symmetries not
allowed in quantum gravity



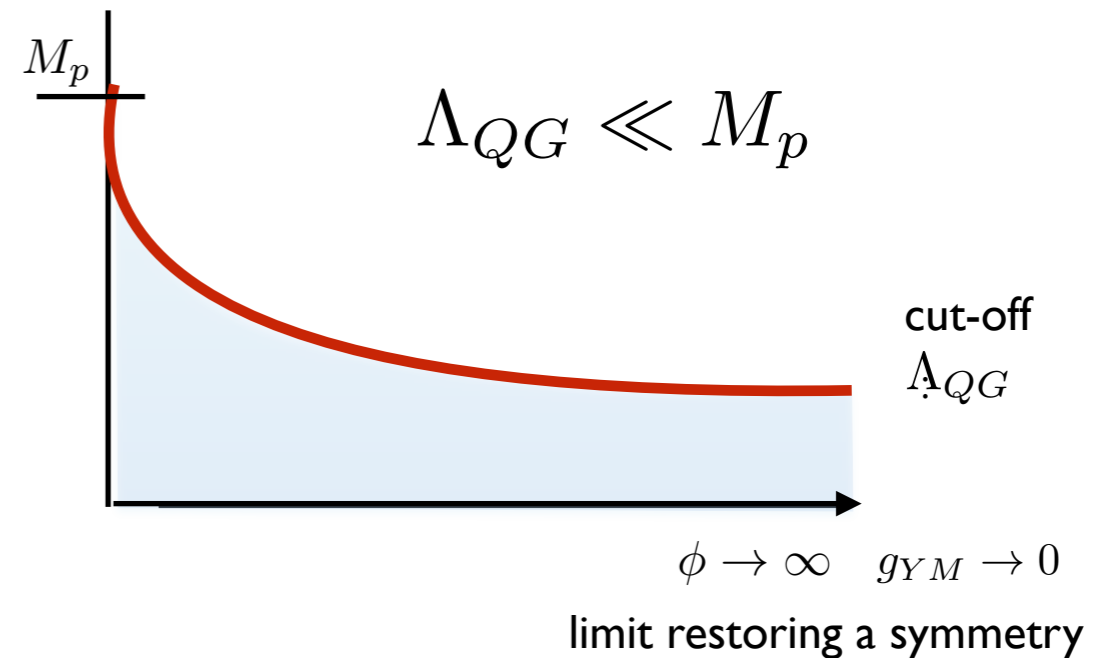
They can only be restored at infinite
field distance (asymptotic limits)

Asymptotic limits

These limits seem under computational control from the point of view of QFT

Asymptotic limits

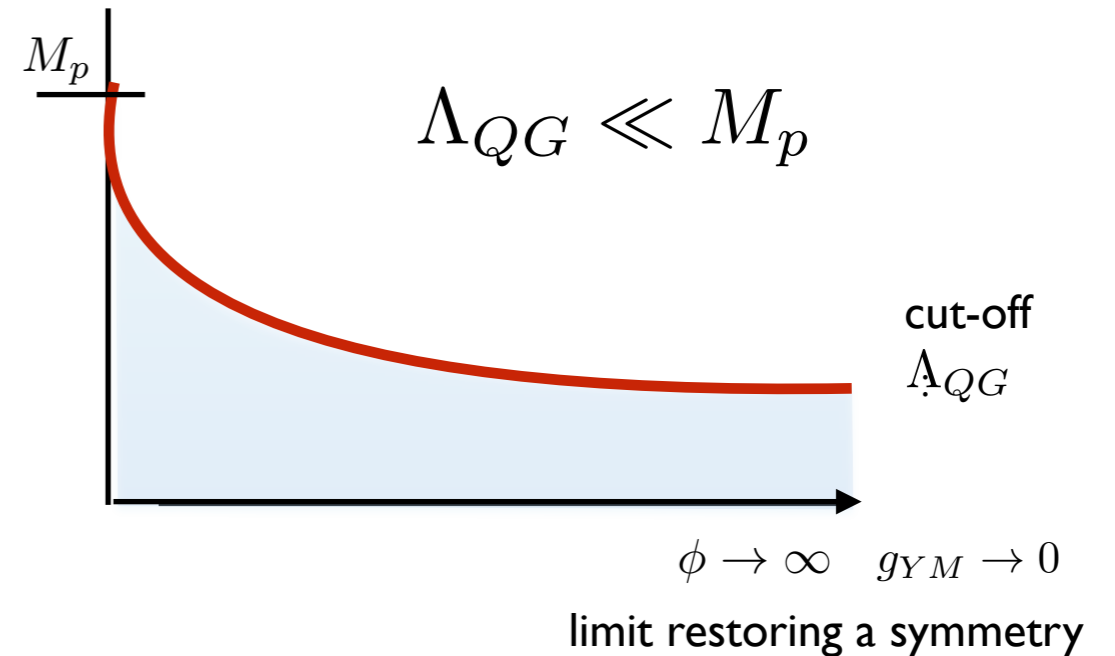
These limits seem under computational control from the point of view of QFT
but the EFT must break down when approaching the boundary by
quantum gravity effects



Asymptotic limits

These limits seem under computational control from the point of view of QFT
but the EFT must break down when approaching the boundary by
quantum gravity effects

How? Does this happens in a
universal way in string theory?

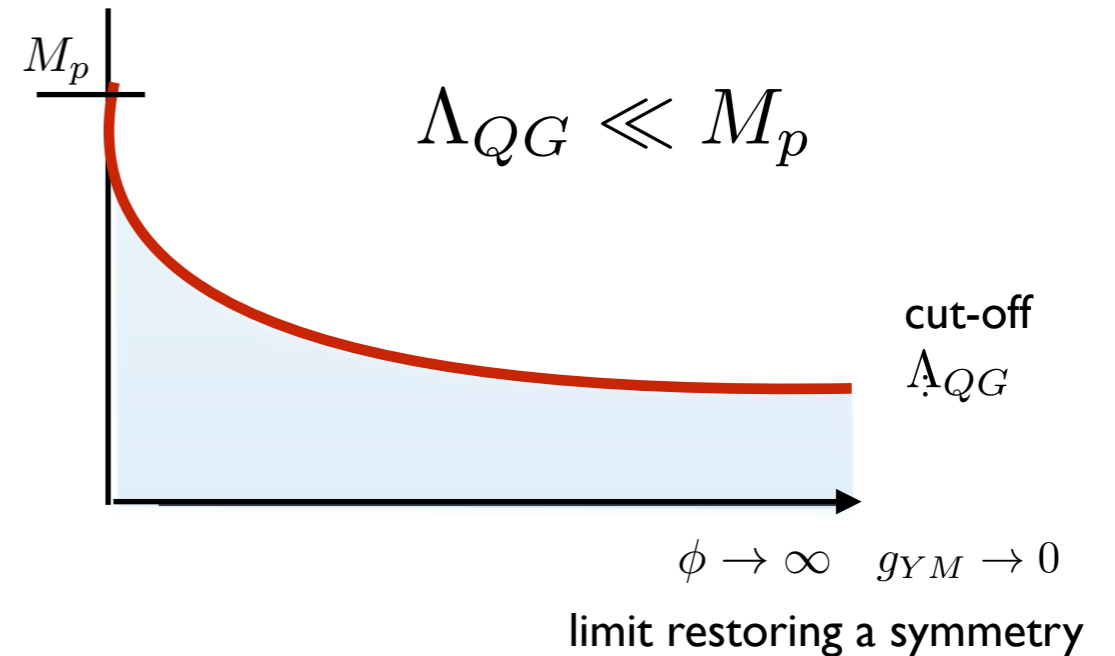


Asymptotic limits

These limits seem under computational control from the point of view of QFT
but the EFT must break down when approaching the boundary by
quantum gravity effects

How? Does this happens in a
universal way in string theory?

YES!

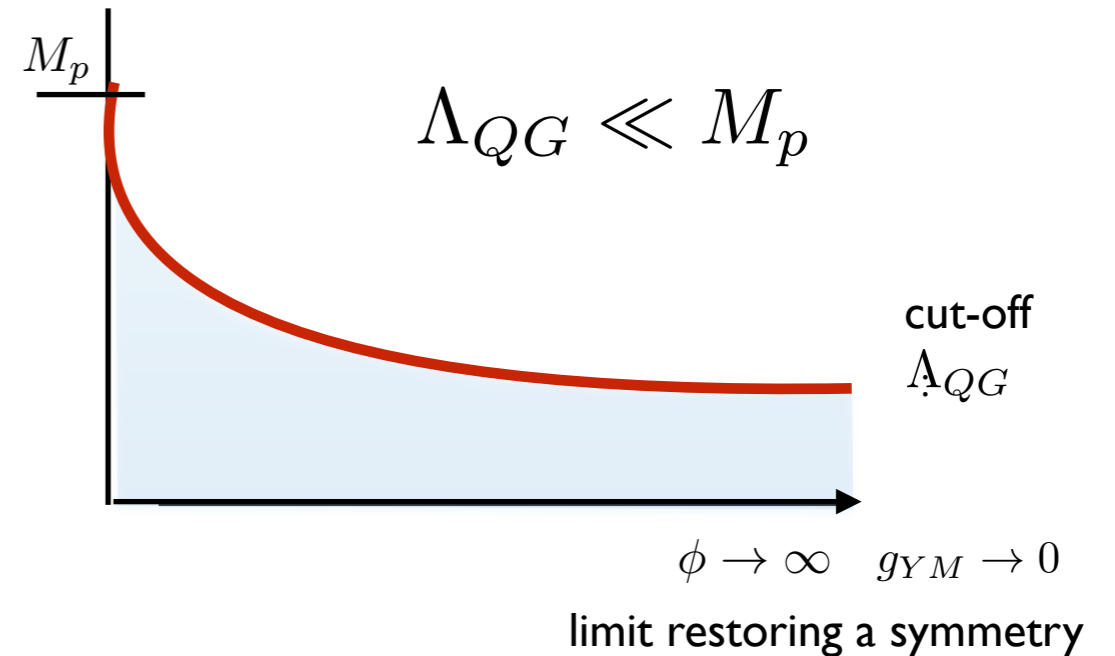


Asymptotic limits

These limits seem under computational control from the point of view of QFT
but the EFT must break down when approaching the boundary by
quantum gravity effects

How? Does this happens in a
universal way in string theory?

YES!

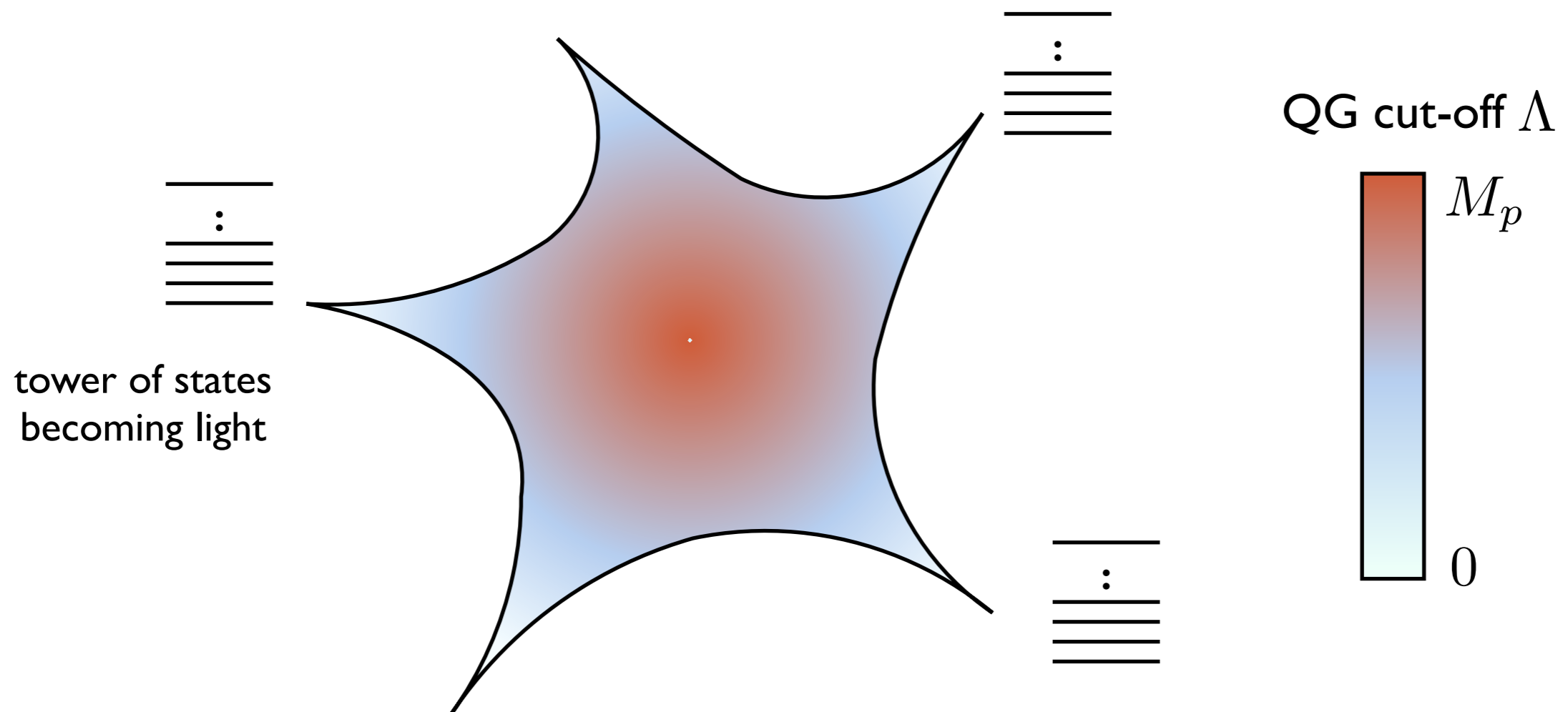


tower of states

There is new light physics that forces the cut-off to go to zero
and acts as a censorship mechanism to restore global symmetries.

Asymptotic Towers of States

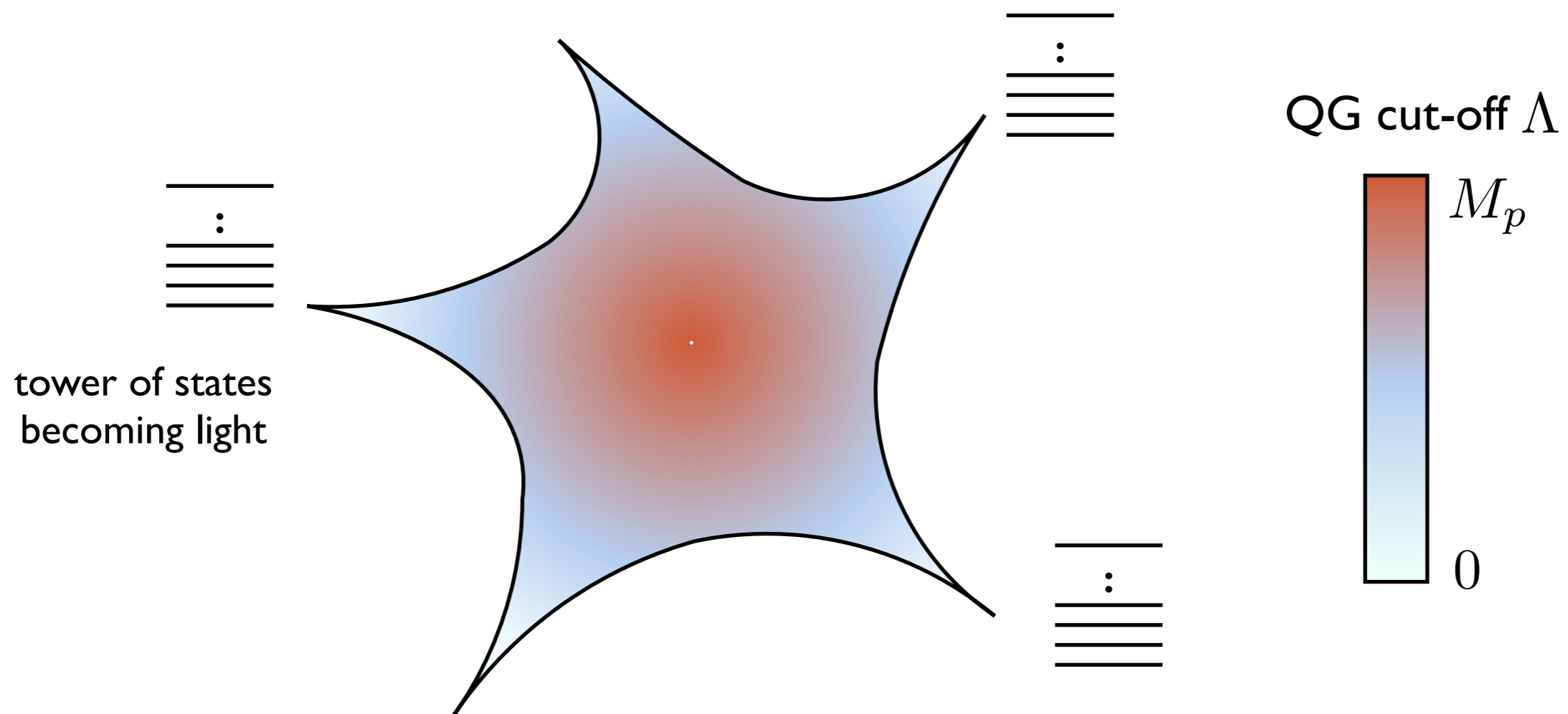
Quantum gravitational effects become significant at a scale $\Lambda_{QG} \ll M_p$ due to towers of states becoming light as approaching the boundaries



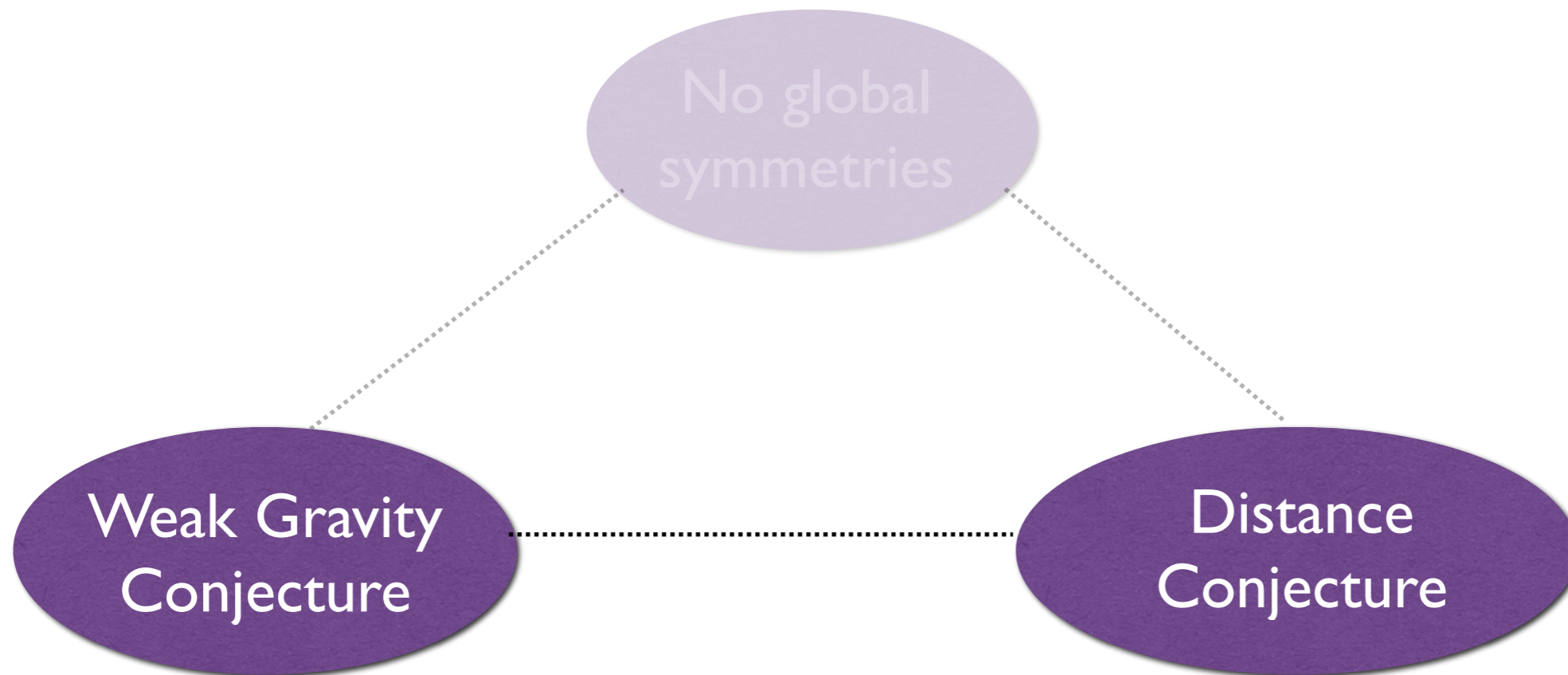
Asymptotic Towers of States

Quantum gravitational effects become significant at a scale $\Lambda_{QG} \ll M_p$ due to towers of states becoming light as approaching the boundaries

Goal: Determine the **quantum gravity cut-off** at which the EFT breaks down in terms of EFT data that quantifies how close we are to the boundaries (e.g. the value of a gauge coupling, the scalar field range...)



Swampland conjectures



They quantify how approximate global symmetries can be by specifying the concrete behaviour of this tower of states

Distance Conjecture

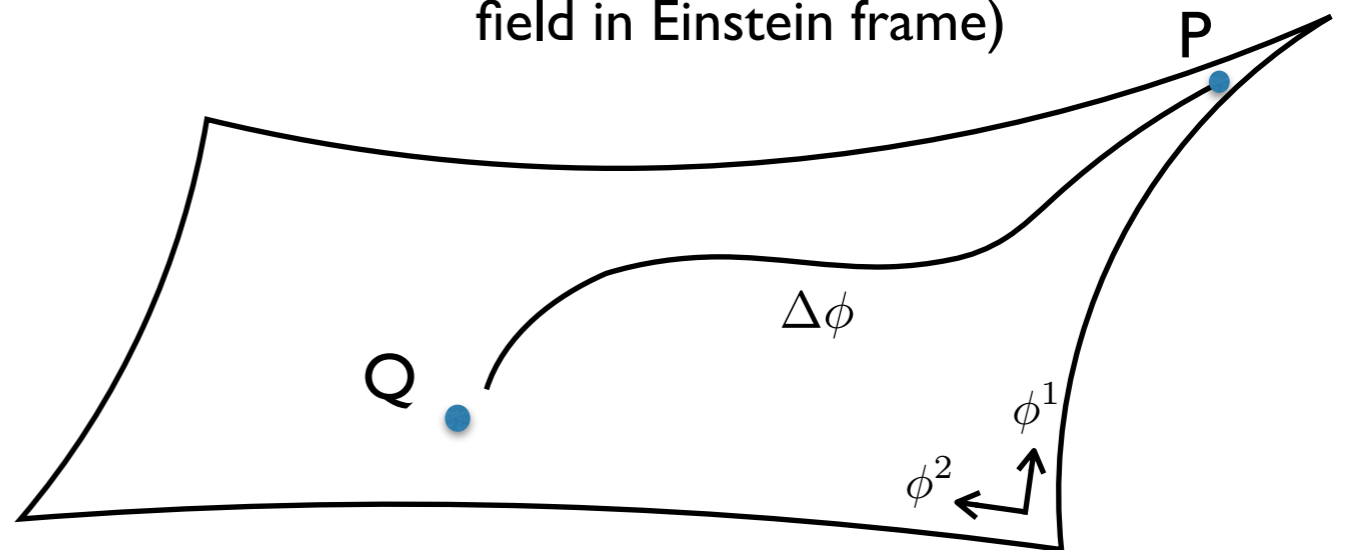
Given an EFT coupled to gravity, with a moduli space parametrized by the vacuum expectation value of some scalar fields:

There is an **infinite tower of states** becoming **exponentially light** at every **infinite field distance** limit of the moduli space

$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \text{when} \quad \Delta\phi \rightarrow \infty \quad [\text{Ooguri-Vafa'06}]$$

$\mathcal{L} = g_{ij}(\phi) \partial\phi^i \partial\phi^j$ \rightarrow scalar manifold (moduli space)

$$\Delta\phi = \int_Q^P \sqrt{g_{ij} \frac{d\phi^i}{ds} \frac{d\phi^j}{ds}} ds \equiv \text{geodesic distance (canonically normalised scalar field in Einstein frame)}$$



Distance Conjecture

Given an EFT coupled to gravity, with a moduli space parametrized by the vacuum expectation value of some scalar fields:

There is an **infinite tower of states** becoming **exponentially light** at every **infinite field distance** limit of the moduli space

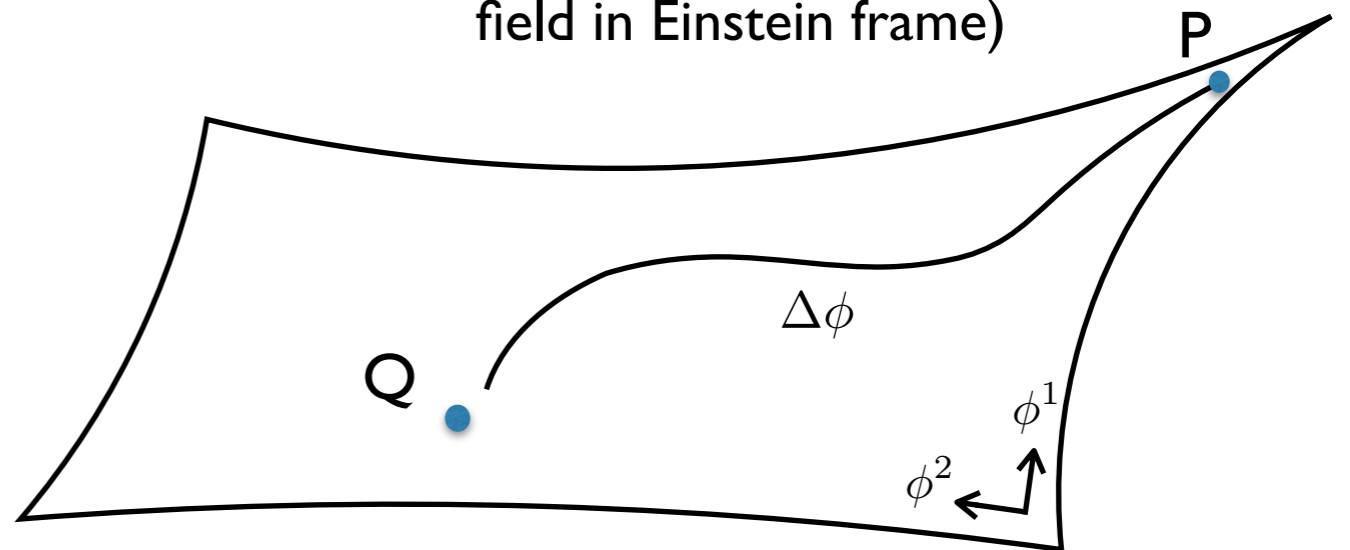
$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \text{when} \quad \Delta\phi \rightarrow \infty \quad [\text{Ooguri-Vafa'06}]$$

$\mathcal{L} = g_{ij}(\phi) \partial\phi^i \partial\phi^j$ \rightarrow scalar manifold (moduli space)

$$\Delta\phi = \int_Q^P \sqrt{g_{ij} \frac{d\phi^i}{ds} \frac{d\phi^j}{ds}} ds \equiv \text{geodesic distance (canonically normalised scalar field in Einstein frame)}$$

For example:

- Kaluza-Klein towers as $R \rightarrow \infty$
- winding modes as $R \rightarrow 0$
- string modes as $g_s \rightarrow 0$



Distance Conjecture

Given an EFT coupled to gravity, with a moduli space parametrized by the vacuum expectation value of some scalar fields:

There is an **infinite tower of states** becoming **exponentially light** at every **infinite field distance** limit of the moduli space

$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \text{when} \quad \Delta\phi \rightarrow \infty \quad [\text{Ooguri-Vafa'06}]$$

Evidence:

- Plethora of quantitative tests in string theory

[Grimm, Palti, IV'18] [Grimm, Palti, Li'18] [Gendler, IV'20] [Lee, Lerche, Weigand'18-21] [Corvilain, Grimm, IV'18]
[Baume, Marchesano, Wiesner'19] [Lanza, Marchesano, Martucci, IV'20-21] [Klaewer, Lee, Weigand, Wiesner'22] ...

- Bottom-up arguments based on black hole physics / entropy bounds

[Hamada, Montero, Vafa, IV'21] [Cribiori et al'22]

- Tests in AdS/CFT [Perlmutter, Rastelli, Vafa, IV'21]

[Calderon-Infante et al'23]

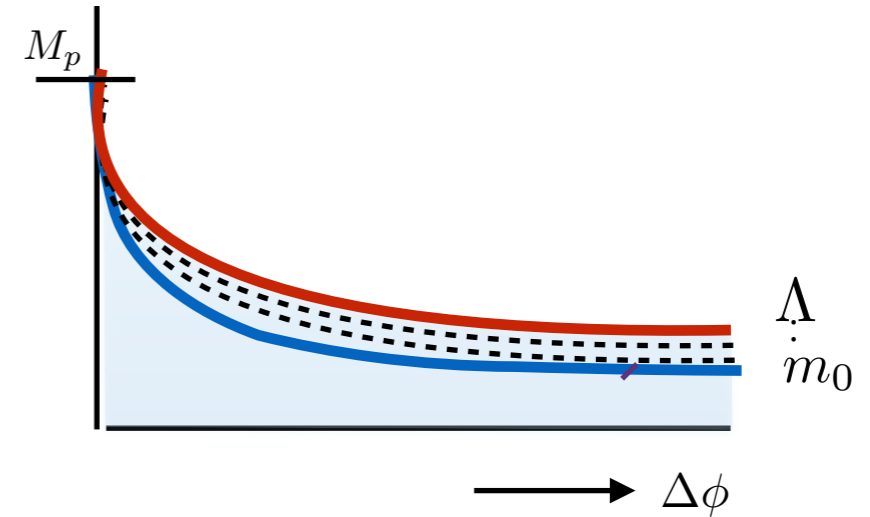
[Baume, Calderon-Infante'21-23]

Phenomenological Implications

This tower signals the quantum gravity **breakdown** of the effective theory:

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi)$$

$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$

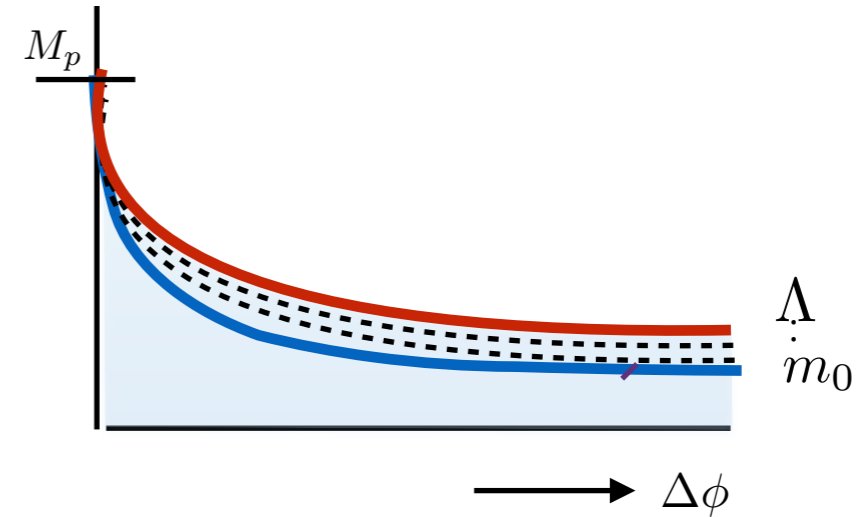


Phenomenological Implications

This tower signals the quantum gravity **breakdown** of the effective theory:

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi)$$

$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$



Maximum scalar field range* that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off

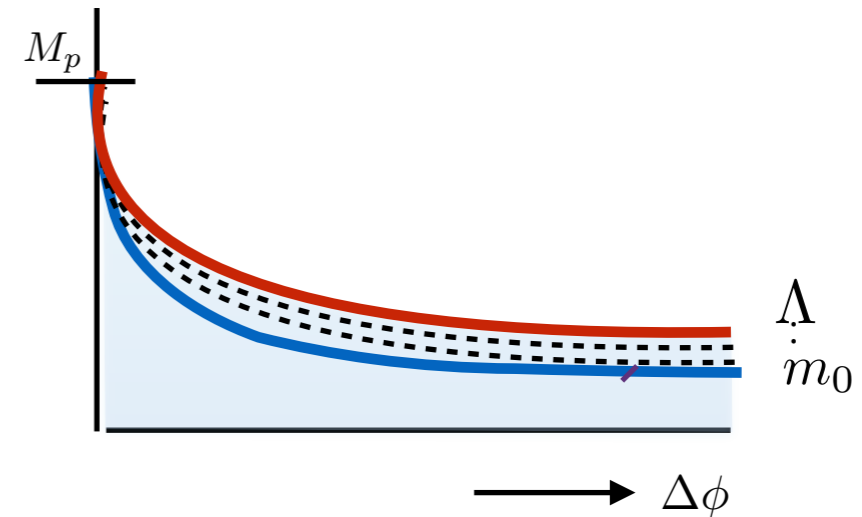
*Caveats to be discussed later

Phenomenological Implications

This tower signals the quantum gravity **breakdown** of the effective theory:

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi)$$

$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log\left(\frac{M_p}{\Lambda}\right)$$



Maximum scalar field range* that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off

Useful to constrain

Inflation

Quintessence

Cosmological solutions to the EW hierarchy problem (e.g. relaxion)

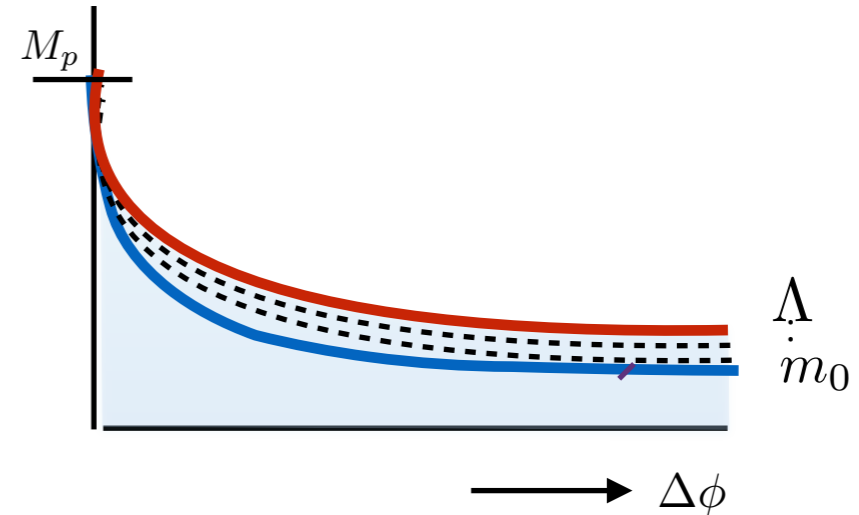
*Caveats to be discussed later

Phenomenological Implications

This tower signals the quantum gravity **breakdown** of the effective theory:

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi)$$

$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log\left(\frac{M_p}{\Lambda}\right)$$



Maximum scalar field range* that can be accommodated in a given EFT as a function of the Quantum Gravity cut-off

Useful to constrain

Inflation

Quintessence

Cosmological solutions to the EW hierarchy problem (e.g. relaxion)

*Caveats to be discussed later

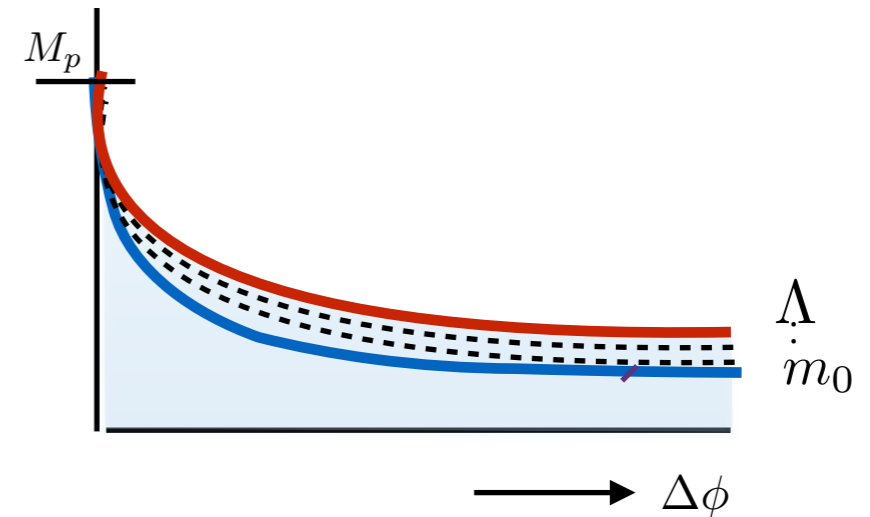
Large field ranges are problematic since the cut-off gets reduced

Phenomenological Implications

This tower signals the quantum gravity **breakdown** of the effective theory:

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi)$$

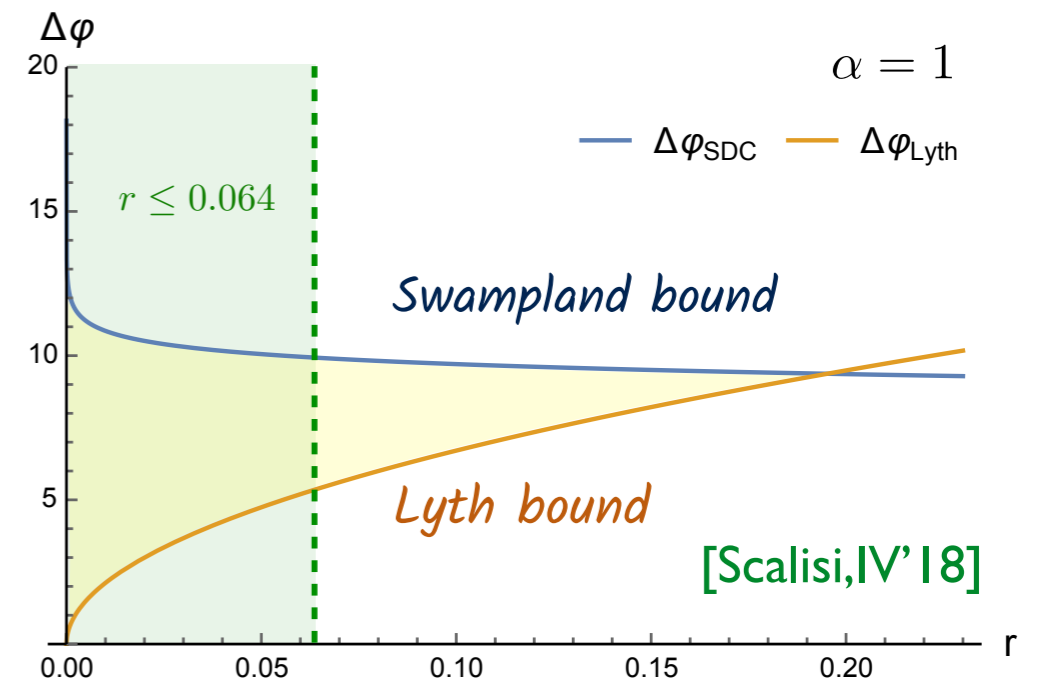
$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$



Example: Constraints on single field inflation

$$\Delta\phi \leq \frac{1}{\alpha} \log \frac{M_p}{H} = \frac{1}{\alpha} \log \sqrt{\frac{2}{\pi^2 A_s r}}$$

$$H \leq \Lambda$$

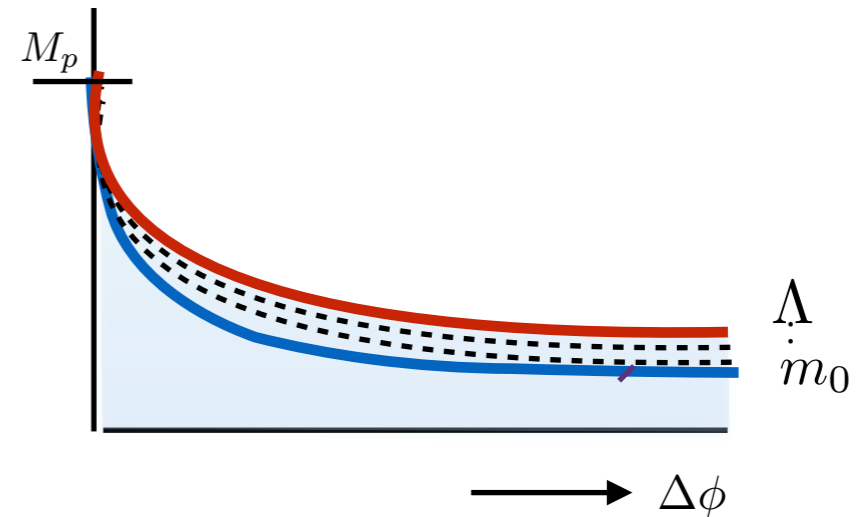


Phenomenological Implications

This tower signals the quantum gravity **breakdown** of the effective theory:

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi)$$

$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$

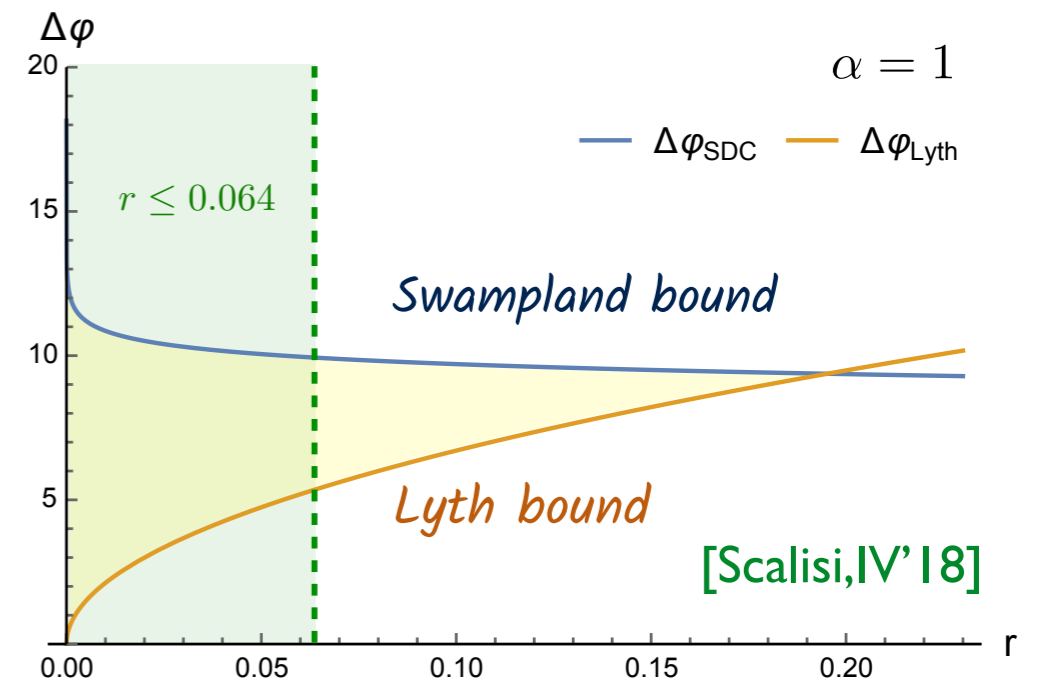


Example: Constraints on single field inflation

$$\Delta\phi \leq \frac{1}{\alpha} \log \frac{M_p}{H} = \frac{1}{\alpha} \log \sqrt{\frac{2}{\pi^2 A_s r}}$$

$$H \leq \Lambda$$

Large field inflation is not ruled out, but can be highly constrained



Cosmological signatures of the tower?

Phenomenological Implications

Caveats to be better understood at the theoretical level in order to give precise phenomenological implications:

I) What is the value of the exponential rate? Is there a lower bound? 

$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \rightarrow \quad \Delta\phi \lesssim \frac{1}{\lambda} \log\left(\frac{M_p}{\Lambda}\right) \quad \alpha_{\min} = \left| \frac{\vec{\nabla} m}{m} \right| \geq \frac{1}{\sqrt{d-2}} \quad \text{[Etheredge et al'22]}$$

??

$$\frac{\vec{\nabla} m}{m} \cdot \frac{\vec{\nabla} \Lambda}{\Lambda} = \frac{1}{d-2} \quad \text{[Lee et al'19]}$$
$$\frac{\vec{\nabla} m}{m} \cdot \frac{\vec{\nabla} \Lambda}{\Lambda} = \frac{1}{d-2} \quad \text{[Castellano et al'23]}$$

Phenomenological Implications

Caveats to be better understood at the theoretical level in order to give precise phenomenological implications:

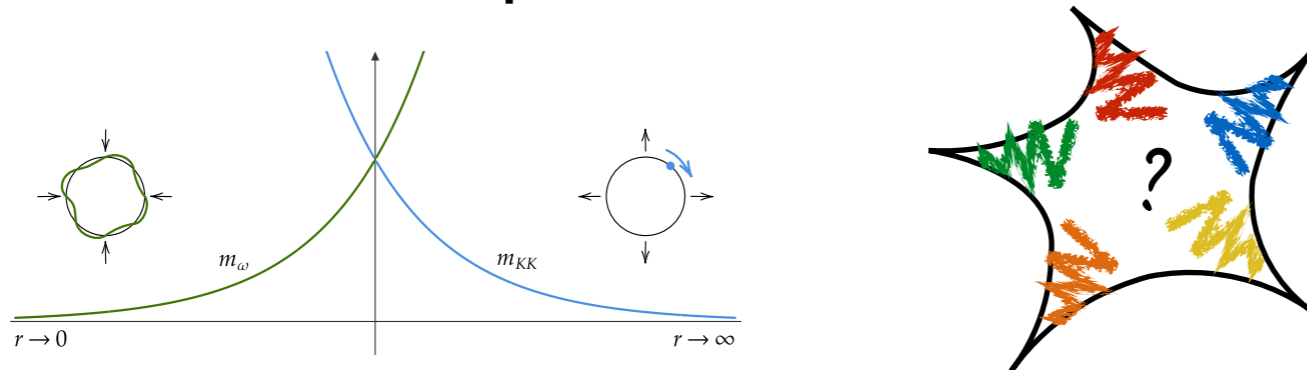
1) What is the value of the exponential rate? Is there a lower bound? 

$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \rightarrow \quad \Delta\phi \lesssim \frac{1}{\lambda} \log\left(\frac{M_p}{\Lambda}\right) \quad \alpha_{\min} = \left| \frac{\vec{\nabla} m}{m} \right| \geq \frac{1}{\sqrt{d-2}} \quad \begin{array}{l} \text{[Etheredge et al'22]} \\ \text{[Lee et al'19]} \end{array}$$

??

$$\frac{\vec{\nabla} m}{m} \cdot \frac{\vec{\nabla} \Lambda}{\Lambda} = \frac{1}{d-2} \quad \text{[Castellano et al'23]}$$

2) When does the exponential behaviour of the masses kick in?



[Long et al'21]

[van de Heisteeg et al'22-23]

Phenomenological Implications

Caveats to be better understood at the theoretical level in order to give precise phenomenological implications:

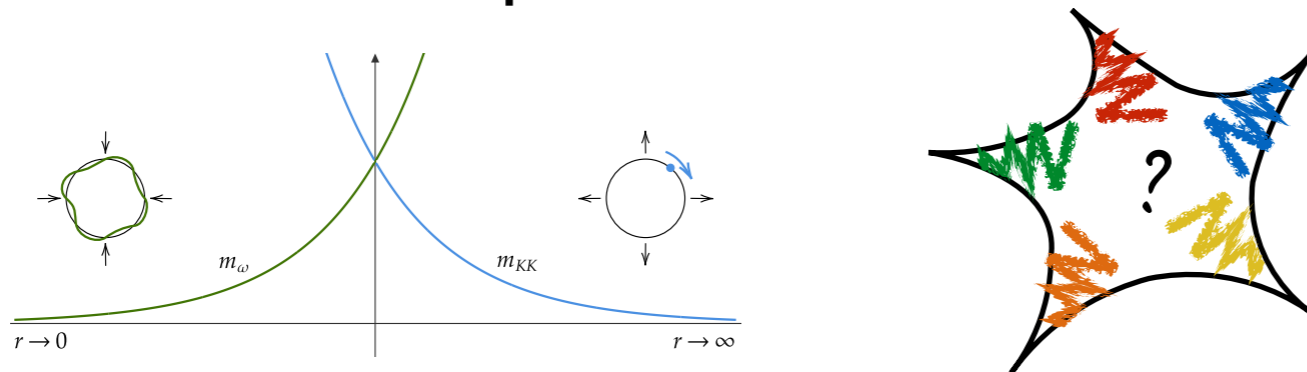
1) What is the value of the exponential rate? Is there a lower bound? 

$$m \sim m_0 e^{-\alpha \Delta\phi} \quad \rightarrow \quad \Delta\phi \lesssim \frac{1}{\lambda} \log\left(\frac{M_p}{\Lambda}\right) \quad \alpha_{\min} = \left| \frac{\vec{\nabla} m}{m} \right| \geq \frac{1}{\sqrt{d-2}} \quad \begin{array}{l} \text{[Etheredge et al'22]} \\ \text{[Lee et al'19]} \end{array}$$

??

$$\frac{\vec{\nabla} m}{m} \cdot \frac{\vec{\nabla} \Lambda}{\Lambda} = \frac{1}{d-2} \quad \text{[Castellano et al'23]}$$

2) When does the exponential behaviour of the masses kick in?

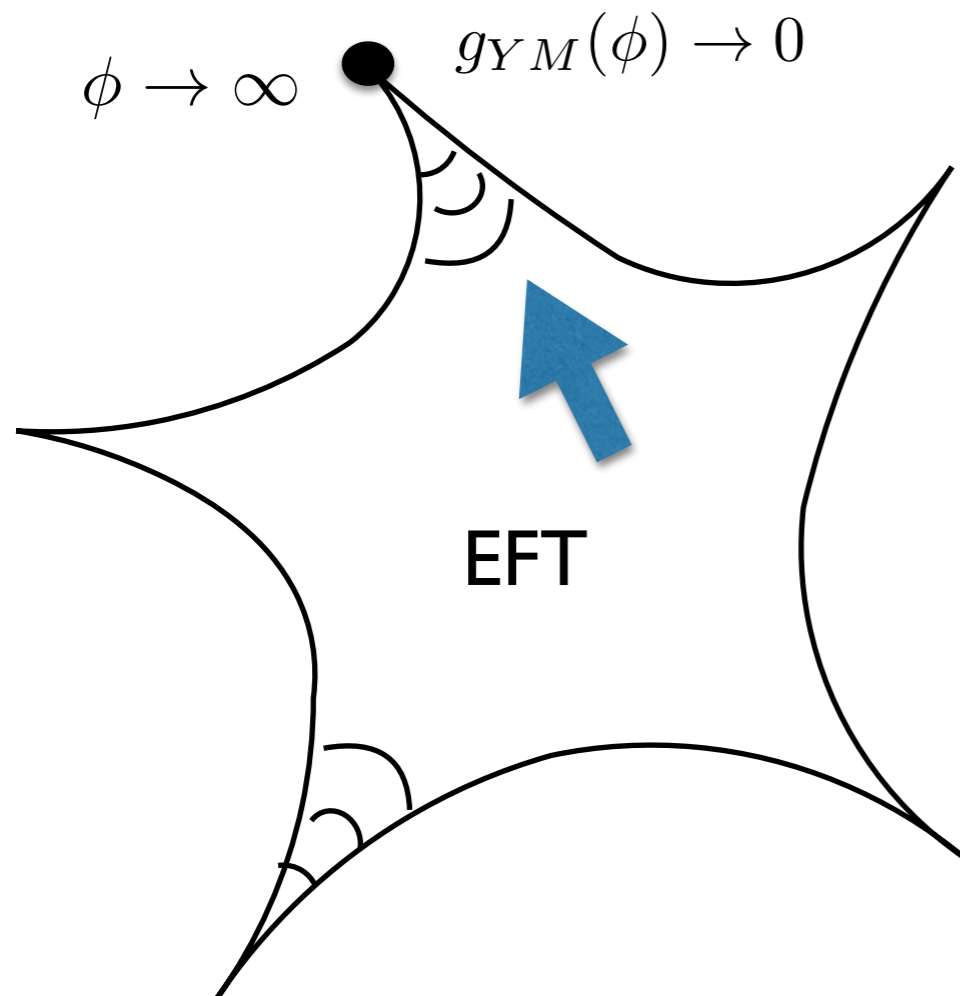


[Long et al'21]
[van de Heisteeg et al'22-23]

3) This is a bound on geodesic distances, that must be read from the kinetic term of the scalars. What about non-geodesic trajectories?
(related to constraints on the potential)

[Calderon-Infante et al'20] [Freigang et al'23]...

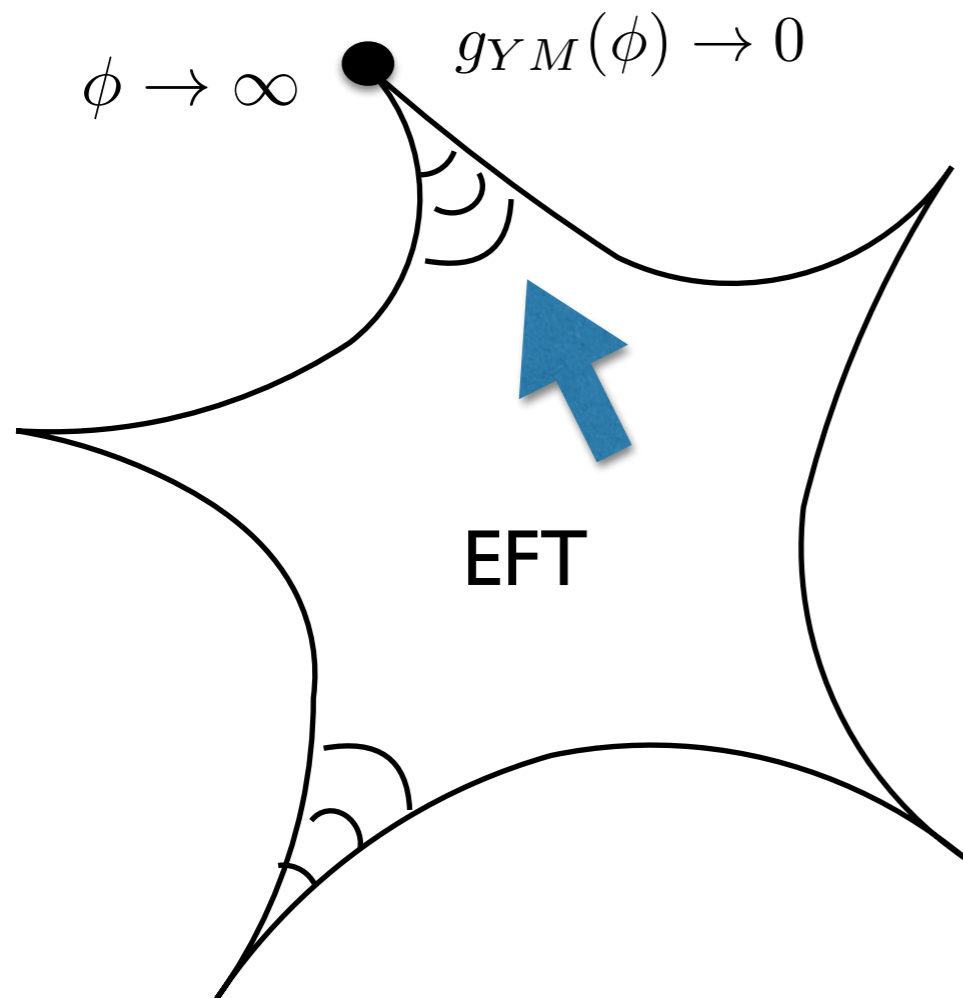
How does the tower/cut-off behaves in terms of EFT data?



For weak coupling limits:

\exists tower of states satisfying the Weak Gravity Conjecture

How does the tower/cut-off behaves in terms of EFT data?



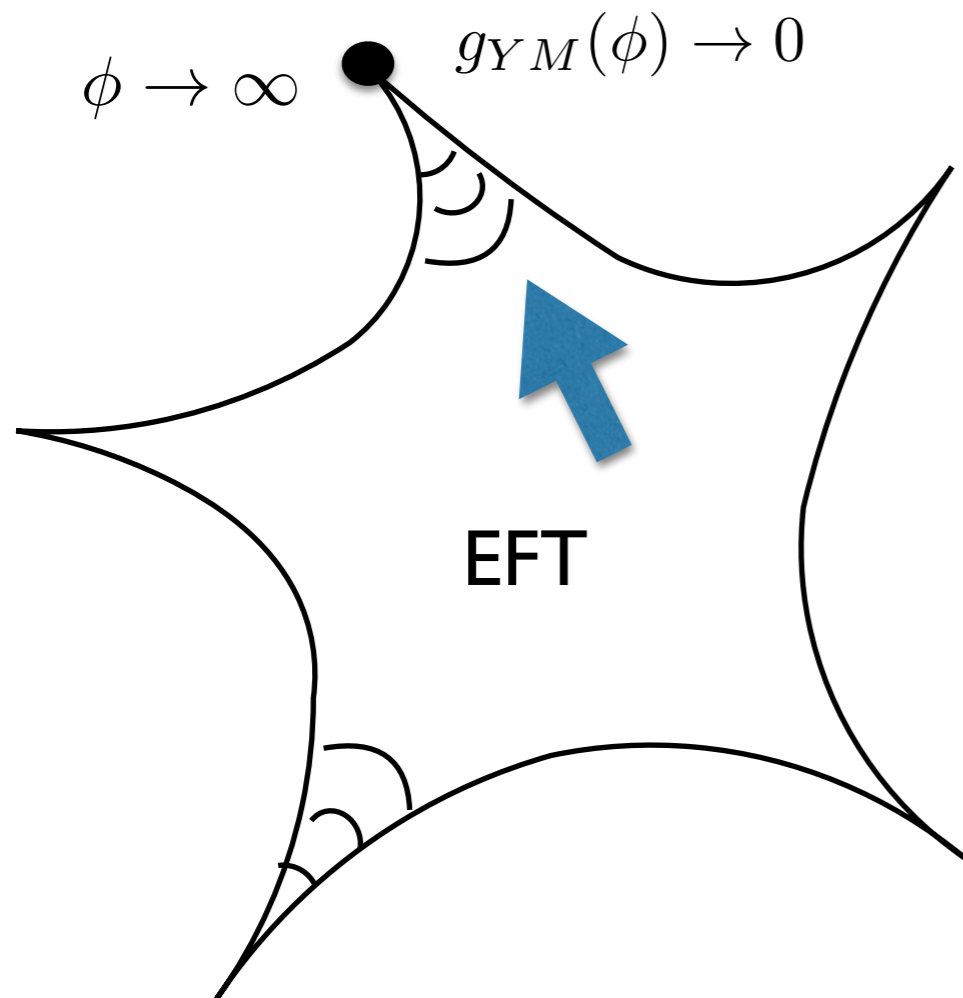
For weak coupling limits:

\exists tower of states satisfying the Weak Gravity Conjecture

\rightarrow EFT breaks down at $\Lambda \lesssim g_{YM}^p M_p$
with $1/3 \leq p \leq 1$

Lower bound on gauge coupling!

How does the tower/cut-off behaves in terms of EFT data?



For weak coupling limits:

\exists tower of states satisfying the Weak Gravity Conjecture

\rightarrow EFT breaks down at $\Lambda \lesssim g_{YM}^p M_p$
with $1/3 \leq p \leq 1$

Lower bound on gauge coupling!

Small gauge couplings are problematic since the cut-off gets reduced

Weak Gravity Conjecture

Weak Gravity conjecture: [Arkani-Hamed et al'06]

Given a gauge theory coupled to gravity, there must exist an electrically charged state with:

$$m \leq \gamma_{\text{BH}} Q M_p$$

mass ← *electric charge* → $\mathcal{O}(1)$ factor (extremality bound of the black holes)

$$Q = q g_{\text{YM}}$$

← *quantized charge* ← *gauge coupling*



Weak Gravity Conjecture

Weak Gravity conjecture: [Arkani-Hamed et al'06]

Given a gauge theory coupled to gravity, there must exist an electrically charged state with:

$$m \leq \gamma_{\text{BH}} Q M_p$$

mass ← *electric charge* → $\mathcal{O}(1)$ factor (extremality bound of the black holes)

$$Q = q g_{\text{YM}}$$

← *quantized charge* ← *gauge coupling*

- Independent motivation based on black hole physics
- Evidence based on string theory, AdS/CFT, scattering amplitudes,...

review: [Harlow et al'22]



Weak Gravity Conjecture

Weak Gravity conjecture: [Arkani-Hamed et al'06]

Given a gauge theory coupled to gravity, there must exist an electrically charged state with:

$$m \leq \gamma_{\text{BH}} Q M_p$$

mass ← *electric charge* → $\mathcal{O}(1)$ factor (extremality bound of the black holes)

$$Q = q g_{\text{YM}}$$

← *quantized charge* ← *gauge coupling*

- Independent motivation based on black hole physics
- Evidence based on string theory, AdS/CFT, scattering amplitudes,...

review: [Harlow et al'22]



If $g_{\text{YM}} \ll 1$ there is a tower of states satisfying the WGC

[Montero et al.'16][Heidenreich et al.'15-16][Andriolo et al'18][Gendler,IV'20]

Weak Gravity Conjecture

Weak Gravity conjecture: [Arkani-Hamed et al'06]

Given a gauge theory coupled to gravity, there must exist an electrically charged state with:

$$m \leq \gamma_{\text{BH}} Q M_p$$

mass ← *electric charge* → $\mathcal{O}(1)$ factor (extremality bound of the black holes)

$$Q = q g_{\text{YM}}$$

← *quantized charge* ← *gauge coupling*

Festina Lente: [Montero et al'19]

Given a gauge theory coupled to gravity in dS space, all particles must satisfy:

$$m^2 \geq \sqrt{6} g H M_p \quad \text{to allow BHs to evaporate back to dS}$$

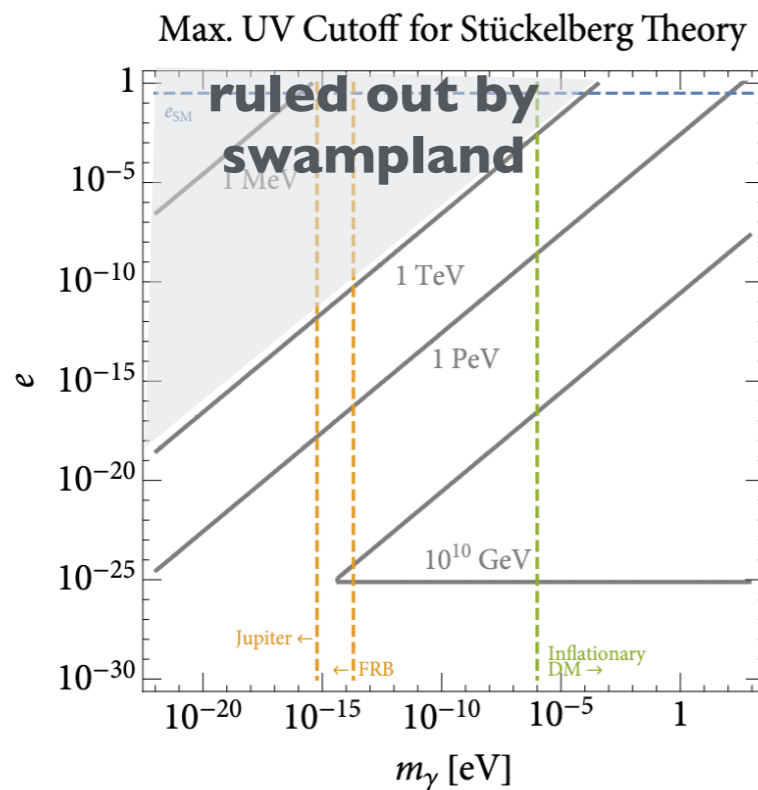
Motivated by a similar reasoning than the WGC for BHs in dS space

Phenomenological implications

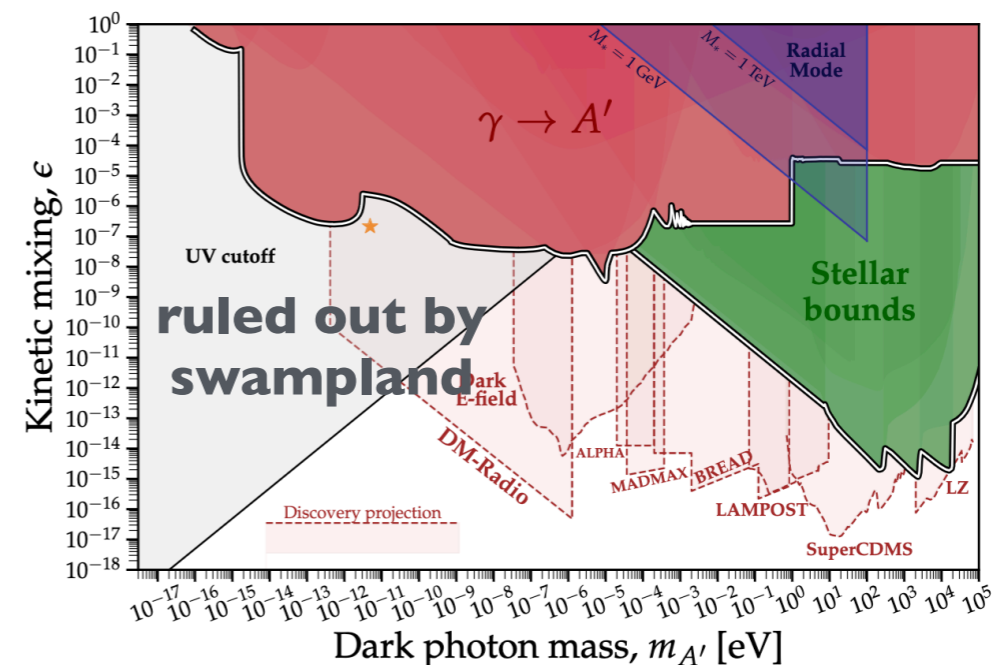
Constraints on dark photons for dark matter:

The quantum gravity cut-off becomes $\Lambda_{QG} \ll M_p$

for weakly coupled or very light dark photons

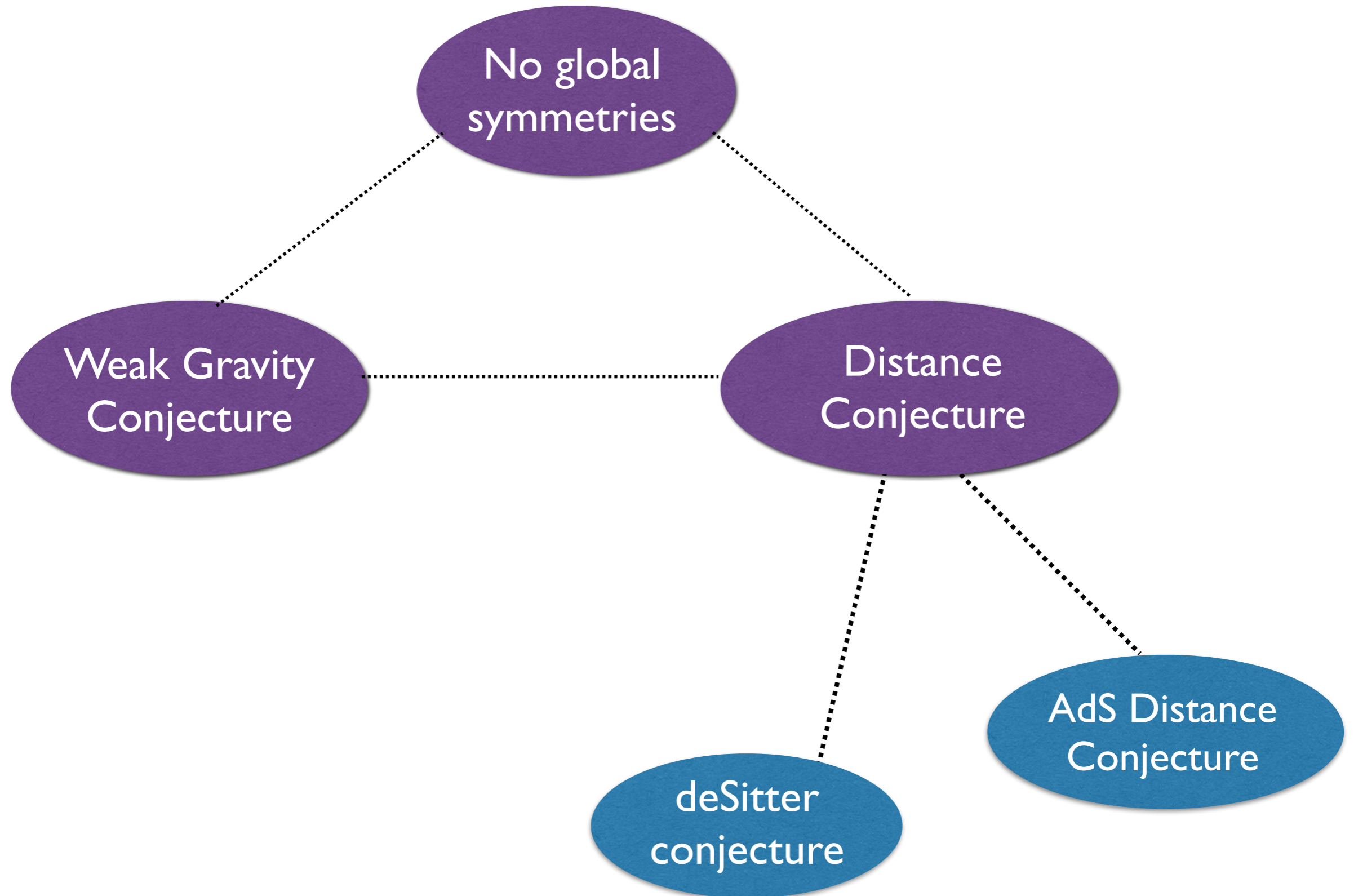


[Reece'18]

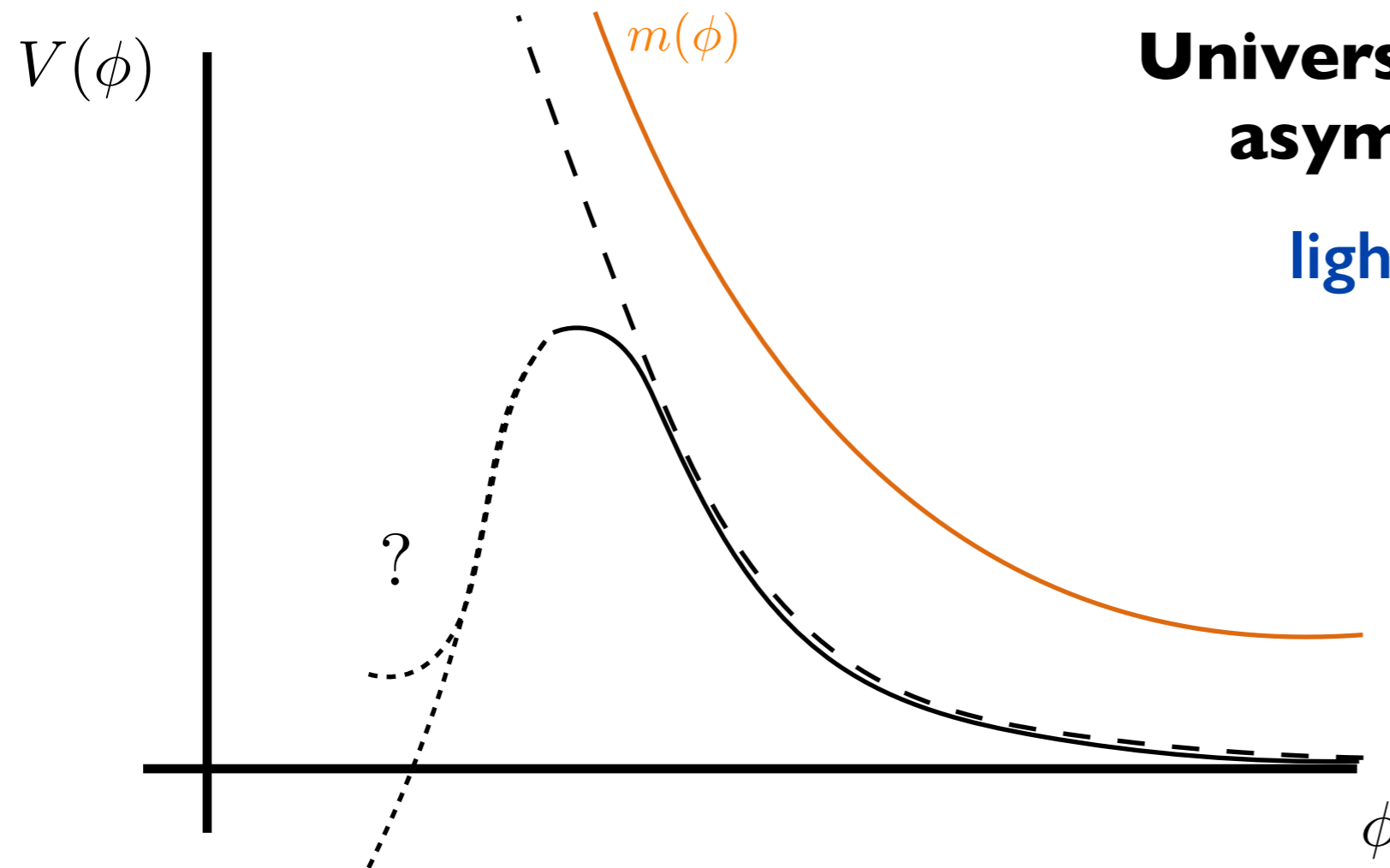


[Montero, Muñoz, Obied'22]

Swampland conjectures

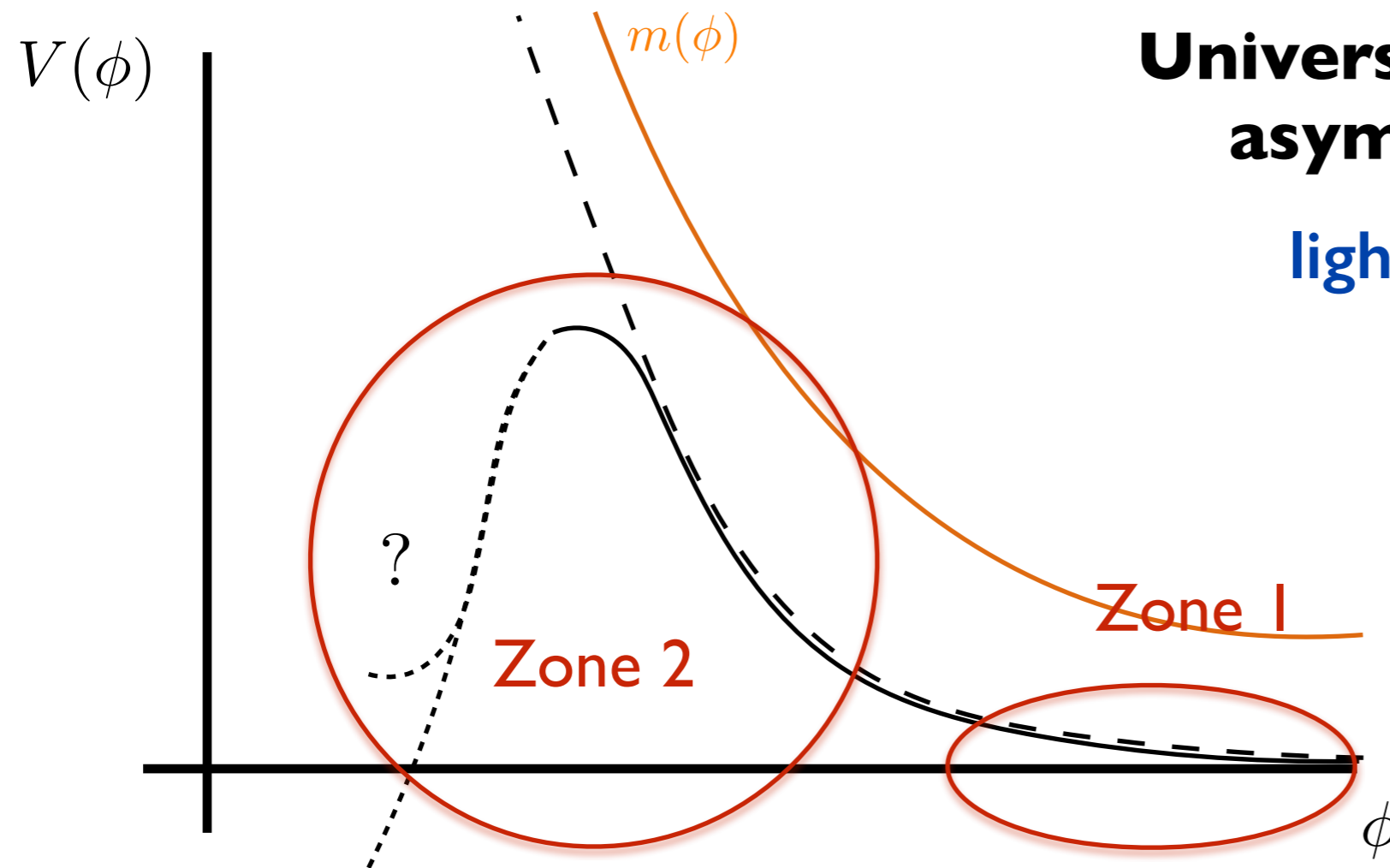


Asymptotic vacuum energy



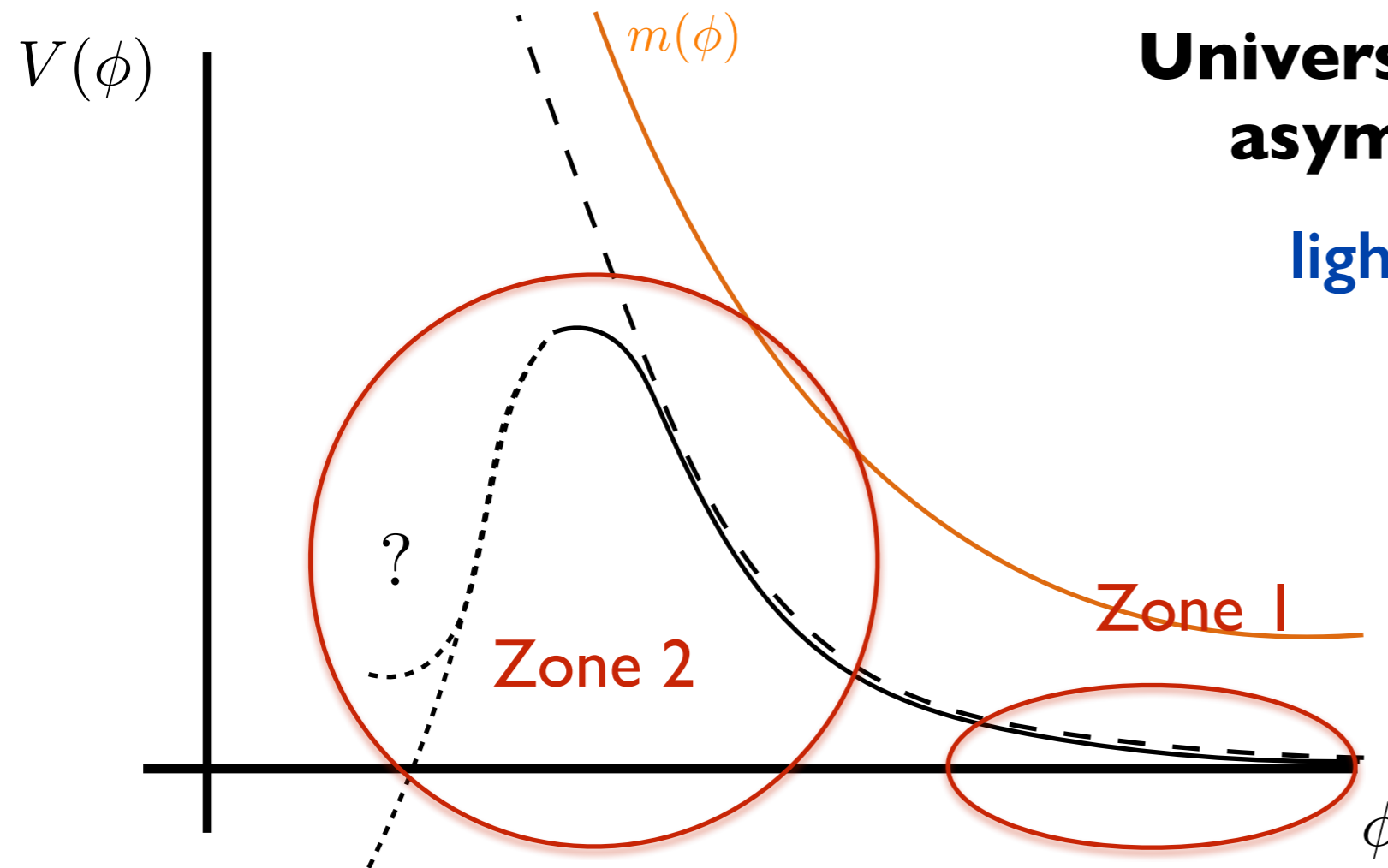
Universal consequence of asymptotic regimes:
light tower of states

Asymptotic vacuum energy



Universal consequence of asymptotic regimes:
light tower of states

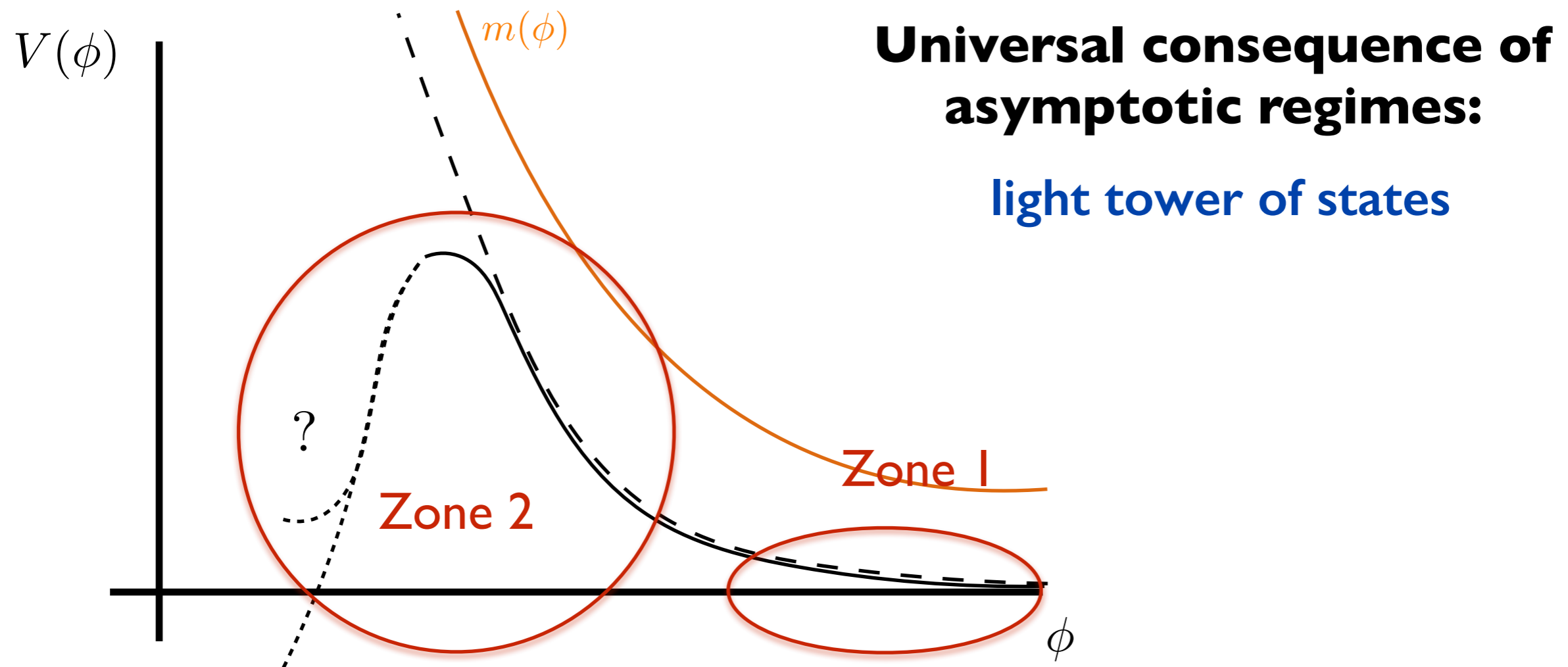
Asymptotic vacuum energy



Universal consequence of asymptotic regimes:
light tower of states

How does the tower of states relates to the vacuum energy?

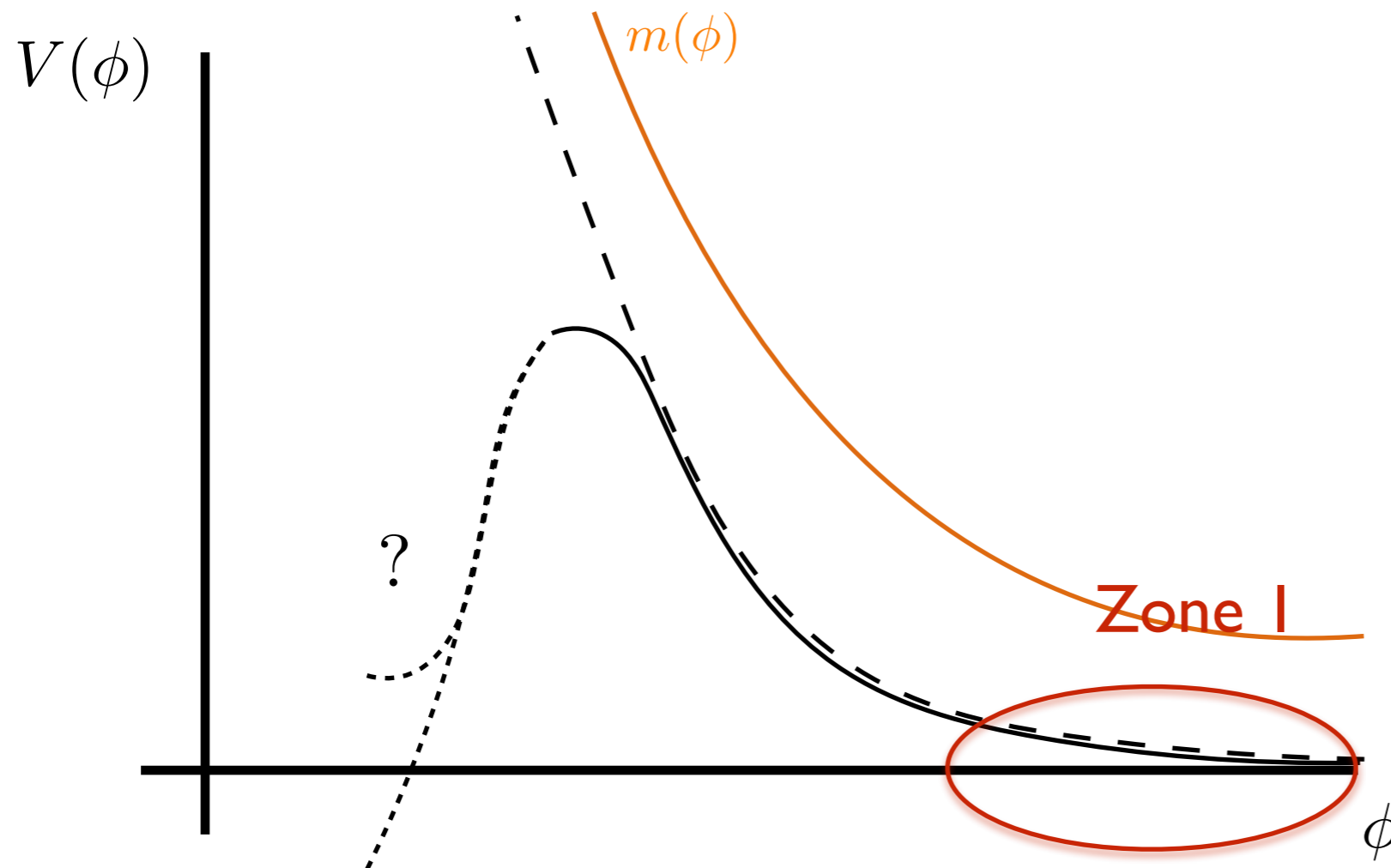
Asymptotic vacuum energy



How does the tower of states relates to the vacuum energy?

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units, as } V_0 \rightarrow 0$$

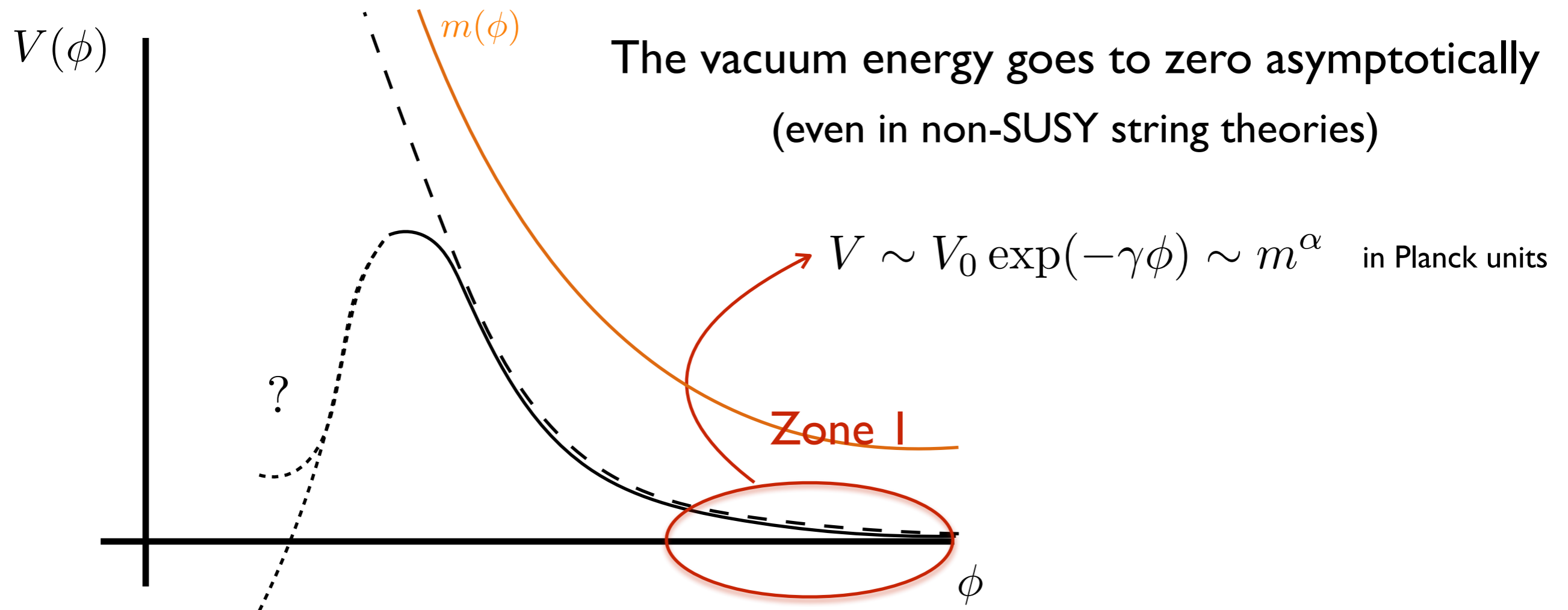
Asymptotic vacuum energy



What is the asymptotic behaviour of the potential?

Could the universe be accelerating forever?

Asymptotic vacuum energy



What is the asymptotic behaviour of the potential?

Could the universe be accelerating forever?

Asymptotic potential at parametric control

Asymptotically: $V \sim V_0 \exp(-\gamma\phi) \sim m^\alpha$ in Planck units

Asymptotic potential at parametric control

Asymptotically: $V \sim V_0 \exp(-\gamma\phi) \sim m^\alpha$ in Planck units

Asymptotic deSitter conjecture: $\gamma = \frac{\nabla V}{V} \geq \mathcal{O}(1)$

[Ooguri,Palti,Shiu,Vafa'18]

- Swampland motivation: consequence of the light tower of states

Asymptotic potential at parametric control

Asymptotically: $V \sim V_0 \exp(-\gamma\phi) \sim m^\alpha$ in Planck units

Asymptotic deSitter conjecture: $\gamma = \frac{\nabla V}{V} \geq \mathcal{O}(1)$

[Ooguri,Palti,Shiu,Vafa'18]

- Swampland motivation: consequence of the light tower of states
- Plethora of string theory evidence: [Wrase,Junghans,Andriot... '18-19] [Grimm et al'19]
(new no-go's for dS at parametric control)

Asymptotic potential at parametric control

Asymptotically: $V \sim V_0 \exp(-\gamma\phi) \sim m^\alpha$ in Planck units

Asymptotic deSitter conjecture: $\gamma = \frac{\nabla V}{V} \geq \mathcal{O}(1)$

[Ooguri,Palti,Shiu,Vafa'18]

- Swampland motivation: consequence of the light tower of states
- Plethora of string theory evidence: [Wrase,Junghans,Andriot... '18-19] [Grimm et al'19]
(new no-go's for dS at parametric control)

What is γ ? **Could it describe an accelerating cosmology?**

Notice! If $\gamma \leq \frac{2}{\sqrt{d-2}}$ accelerated expansion à la quintessence

Asymptotic potential at parametric control

Asymptotically: $V \sim V_0 \exp(-\gamma\phi) \sim m^\alpha$ in Planck units

Asymptotic deSitter conjecture: $\gamma = \frac{\nabla V}{V} \geq \mathcal{O}(1)$

[Ooguri,Palti,Shiu,Vafa'18]

- Swampland motivation: consequence of the light tower of states
- Plethora of string theory evidence: [Wrase,Junghans,Andriot... '18-19] [Grimm et al'19]
(new no-go's for dS at parametric control)

What is γ ? **Could it describe an accelerating cosmology?**

Notice! If $\gamma \leq \frac{2}{\sqrt{d-2}}$ accelerated expansion à la quintessence

Many recent checks in string theory (no clear no-go yet) [Li,Grimm,IV'19][Valeixo et al'20] [Andriot et al'20-22] [Cicoli et al'21-22] [Calderon-Infante et al'22] [Shiu,Tonioni'23] [Cremonini et al'23][Hebecker et al'23]...

Asymptotic potential at parametric control

Asymptotically: $V \sim V_0 \exp(-\gamma\phi) \sim m^\alpha$ in Planck units

Asymptotic deSitter conjecture: $\gamma = \frac{\nabla V}{V} \geq \mathcal{O}(1)$


[Ooguri,Palti,Shiu,Vafa'18]

- Swampland motivation: consequence of the light tower of states
- Plethora of string theory evidence: [Wrase,Junghans,Andriot... '18-19] [Grimm et al'19]
(new no-go's for dS at parametric control)

What is γ ? **Could it describe an accelerating cosmology?**

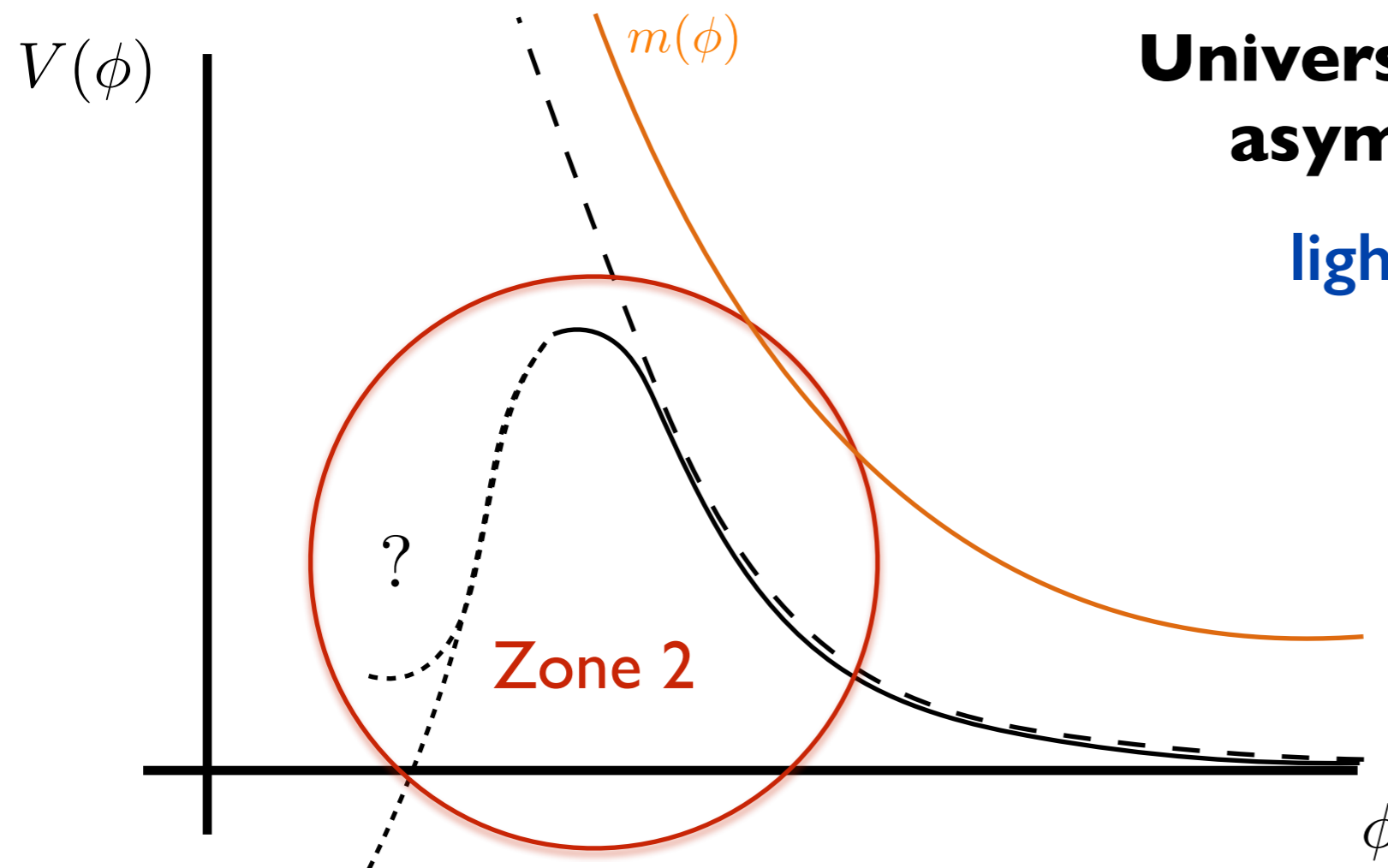
Notice! If $\gamma \leq \frac{2}{\sqrt{d-2}}$ accelerated expansion à la quintessence

If the tower contains higher spin fields: $\alpha \geq 2$ (Higuchi bound)

$\left| \frac{\vec{\nabla} m}{m} \right| \geq \frac{1}{\sqrt{d-2}}$ implies $\left| \frac{\vec{\nabla} V_0}{V_0} \right| \geq \frac{2}{\sqrt{d-2}}$  no accelerated expansion at parametrically late times

[Bedroya,Vafa'19] [Rudelius'22]

Asymptotic vacuum energy



Universal consequence of asymptotic regimes:
light tower of states

The behaviour of the potential becomes model-dependent

How does the tower of states relates to the vacuum energy?

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units, as } V_0 \rightarrow 0$$

Relation with vacuum energy

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units,} \quad \text{as } V_0 \rightarrow 0$$

From string theory perspective: $V = V_{\text{tree}} + V_{\text{loop}} + \dots$

If the different terms compete to generate a vacuum, at the minimum one naturally still has: $V_0 = \lambda m^\alpha$

Relation with vacuum energy

$$V_0 \sim m_{\text{tower}}^\alpha \text{ in Planck units, as } V_0 \rightarrow 0$$

From string theory perspective: $V = V_{\text{tree}} + V_{\text{loop}} + \dots$

If the different terms compete to generate a vacuum, at the minimum one naturally still has: $V_0 = \lambda m^\alpha$



Caveat: One could try to fine-tune λ to decouple them,

but naturally V_0 remains small if the tower is light

Relation with vacuum energy

$$V_0 \sim m_{\text{tower}}^\alpha \text{ in Planck units, as } V_0 \rightarrow 0$$

From string theory perspective: $V = V_{\text{tree}} + V_{\text{loop}} + \dots$

If the different terms compete to generate a vacuum, at the minimum one naturally still has: $V_0 = \lambda m^\alpha$



Caveat: One could try to fine-tune λ to decouple them,

but naturally V_0 remains small if the tower is light

- ☑ All known families of holographic AdS vacua (even DGKT)
- ☑ KKLT-like proposals for dS in string theory
- ☑ AdS/dS proposals using Casimir energies: $V_0 \sim m^d$

Relation with vacuum energy

$$V_0 \sim m_{\text{tower}}^\alpha \text{ in Planck units, as } V_0 \rightarrow 0$$

From swampland perspective:

It can be motivated by generalising the conjecture to distances in the space of metric configurations (rather than only in moduli space):

$$\textit{Distance} \sim \log |V_0| \quad \textit{Flat space limit } V_0 \rightarrow 0 \textit{ is at infinite distance}$$

$$m \sim \exp(-\alpha \text{ distance}) \sim \Lambda^\alpha \text{ as } \Lambda \rightarrow 0$$

(AdS Distance Conjecture) [Luest,Palti,Vafa'19]

Relation with vacuum energy

$$V_0 \sim m_{\text{tower}}^\alpha \text{ in Planck units, as } V_0 \rightarrow 0$$

From swampland perspective:

It can be motivated by generalising the conjecture to distances in the space of metric configurations (rather than only in moduli space):

$$\text{Distance} \sim \log |V_0| \quad \text{Flat space limit } V_0 \rightarrow 0 \text{ is at infinite distance}$$

$$m \sim \exp(-\alpha \text{ distance}) \sim \Lambda^\alpha \text{ as } \Lambda \rightarrow 0$$

(AdS Distance Conjecture) [Luest,Palti,Vafa'19]

No clear bottom-up explanation (indep. of string theory) yet

Our universe

Let me assume this relation and study its consequences

$$m \sim V_0^{1/\alpha} \rightarrow 0 \quad \text{Could it be the case of our universe?}$$

[Montero, Vafa, IV'22]

Our universe

Let me assume this relation and study its consequences

$$m \sim V_0^{1/\alpha} \rightarrow 0 \quad \text{Could it be the case of our universe?}$$

[Montero, Vafa, IV'22]

Possible scenario:

The **smallness of our vacuum energy** is not due to a huge fine-tuning of contributions in a landscape, but **is a signal of being near an asymptotic limit** where it naturally goes to zero

Our universe

Let me assume this relation and study its consequences

$$m \sim V_0^{1/\alpha} \rightarrow 0 \quad \text{Could it be the case of our universe?}$$

[Montero, Vafa, IV'22]

Possible scenario:

The **smallness of our vacuum energy** is not due to a huge fine-tuning of contributions in a landscape, but **is a signal of being near an asymptotic limit** where it naturally goes to zero

Consequence:

But then, there should be a **light tower of states** whose mass is correlated to the cosmological constant

$$\begin{array}{ccc} 2 \leq \alpha \leq d & \longrightarrow & V^{1/2} \lesssim m \lesssim V^{1/4} \\ \text{Higuchi bound} \longleftarrow & & \\ & \longleftarrow & \text{no fine-tuning} \\ & & d = \text{space-time dimension} \end{array}$$

Experimental constraints

Is a tower with $V^{1/2} \lesssim m \lesssim V^{1/4}$ compatible with
experimental constraints?

In our universe: $V^{1/4} \sim 2.31 \text{ meV}$

Experimental constraints

Is a tower with $V^{1/2} \lesssim m \lesssim V^{1/4}$ compatible with experimental constraints?

In our universe: $V^{1/4} \sim 2.31 \text{ meV}$

Nature of the tower (according to string theory): [Lee, Lerche, Weigand '19]

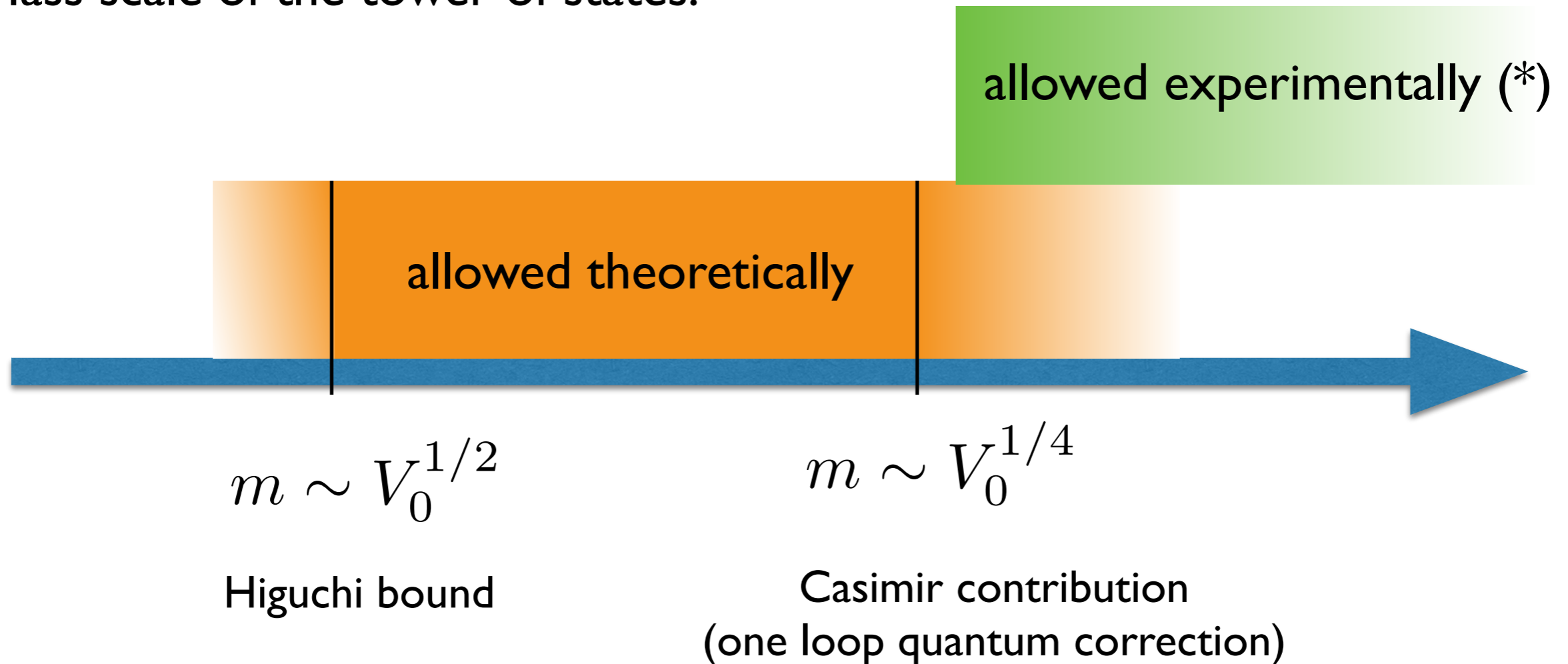
- ❖ ~~String perturbative limit~~ *ruled out exp.*
- ❖ Decompactification of n extra dimensions

Experimental constraints:

- ❖ Astrophysical bounds: $m^{-1} \leq 10^{-4} \mu\text{m}$ (~~$n = 2$~~) *ruled out*
[Hannestad and Raffelt '03] $m^{-1} \leq 44 \mu\text{m}$ ($n = 1$)
- ❖ Dev. from Newton's laws ($n=1$): $m^{-1} \leq 30 \mu\text{m}$ [Lee et al '21]

Experimental constraints

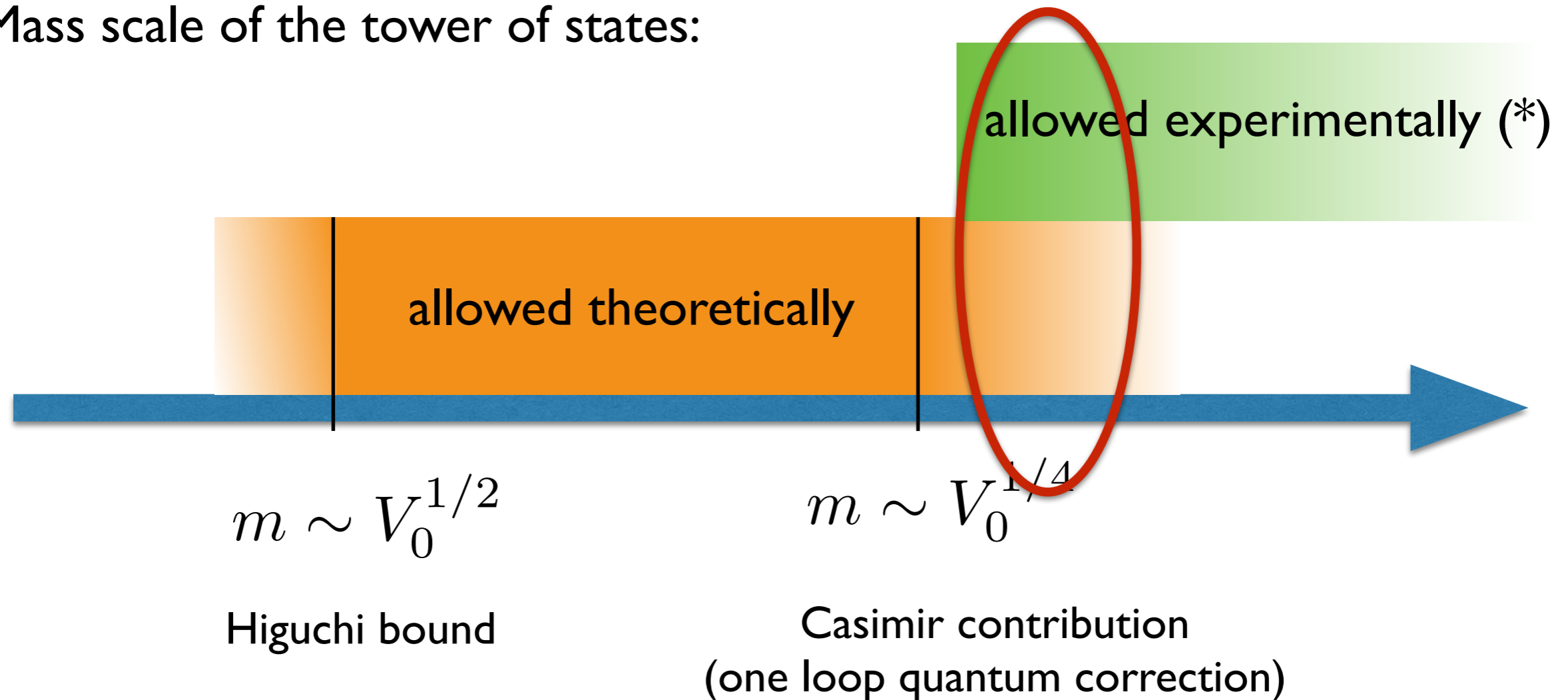
Mass scale of the tower of states:



(*) astrophysical bounds and deviations from Newton's law

Experimental constraints

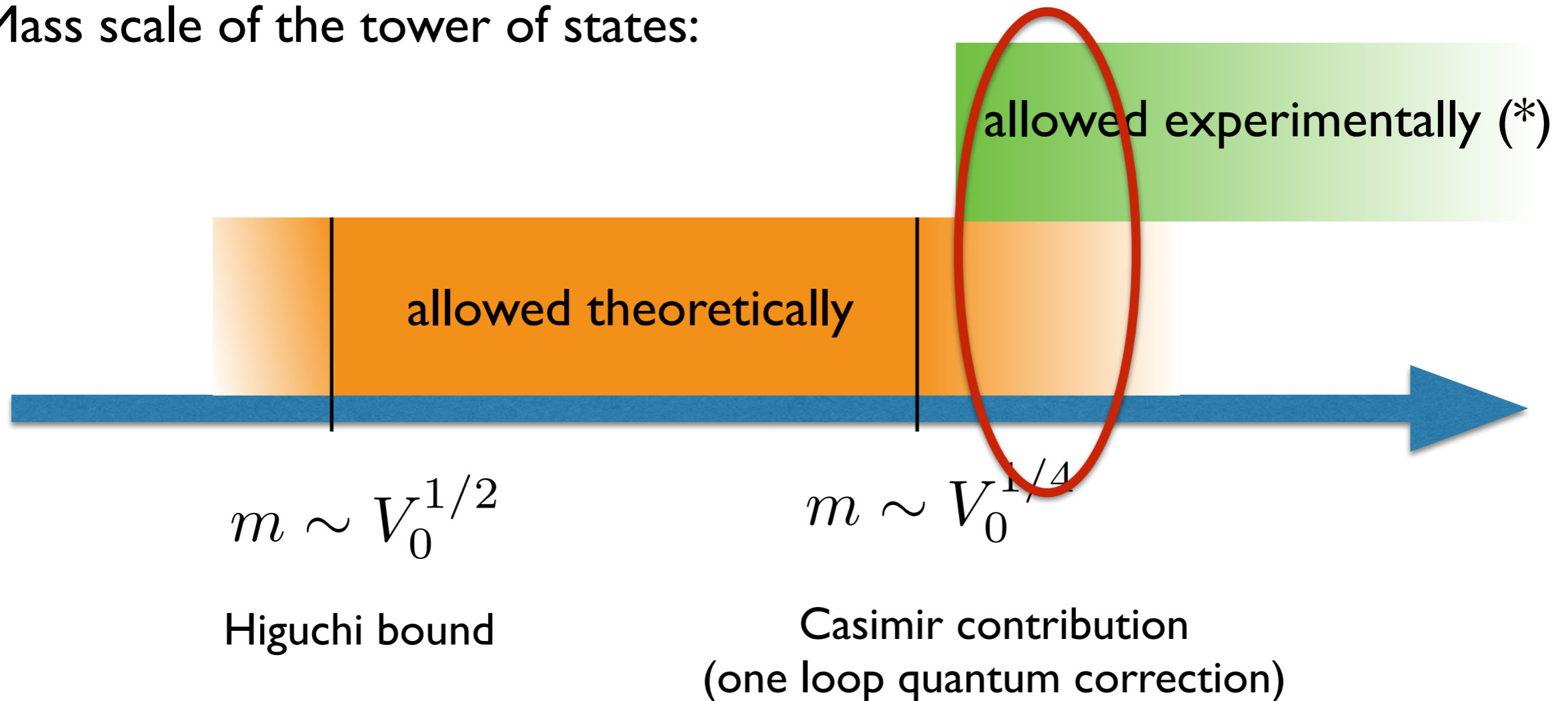
Mass scale of the tower of states:



(*) astrophysical bounds and deviations from Newton's law

Experimental constraints

Mass scale of the tower of states:



(*) astrophysical bounds and deviations from Newton's law

Only $n=1$ (one large extra dimension) is marginally compatible!

Dark Dimension Scenario

If our universe lives near an infinite distance limit $V_0 \rightarrow 0$,

there should be a **light tower of states** of mass: (mod extra fine-tunings)

$$m \sim V_0^{1/4} \sim \mathcal{O}(meV)$$

Dark Dimension Scenario

If our universe lives near an infinite distance limit $V_0 \rightarrow 0$,

there should be a **light tower of states** of mass: (mod extra fine-tunings)

$$m \sim V_0^{1/4} \sim \mathcal{O}(meV)$$



neutrino scale!

Tower of right handed neutrinos?

Dark Dimension Scenario

If our universe lives near an infinite distance limit $V_0 \rightarrow 0$,

there should be a **light tower of states** of mass: (mod extra fine-tunings)

$$m \sim V_0^{1/4} \sim \mathcal{O}(meV)$$



neutrino scale!

Tower of right handed neutrinos?

(it could explain coincidence between neutrino masses and cosmological constant)

Dark Dimension Scenario

If our universe lives near an infinite distance limit $V_0 \rightarrow 0$,

there should be a **light tower of states** of mass: (mod extra fine-tunings)

$$m \sim V_0^{1/4} \sim \mathcal{O}(meV)$$



neutrino scale!

Tower of right handed neutrinos?

(it could explain coincidence between neutrino masses and cosmological constant)

implying one large extra dimension $l \sim 0.1 - 10\mu m$

The Dark Dimension

[Montero, Vafa, IV'22]

Dark Dimension Scenario

If our universe lives near an infinite distance limit $V_0 \rightarrow 0$,

there should be a **light tower of states** of mass: (mod extra fine-tunings)

$$m \sim V_0^{1/4} \sim \mathcal{O}(meV)$$



neutrino scale!

Tower of right handed neutrinos?

(it could explain coincidence between neutrino masses and cosmological constant)

implying one large extra dimension $l \sim 0.1 - 10\mu m$

The Dark Dimension

[Montero,Vafa,IV'22]

(This tower also helps to avoid violation of the AdS Distance conjecture upon compactification of the Standard Model) [Gonzalo,Ibanez,IV'21]

Dark Dimension Scenario

[Arkani-Hamed, Dimopoulos, Dvali'98]

This scenario is an example of the **Large Extra Dimension models (ADD)**

although the scale is different than usual, since it was motivated by the smallness of the cosmological constant and not by the EW hierarchy problem:

$$\text{QG cut-off: } \hat{M} \sim m^{1/3} M_P^{2/3} \sim 10^{10} \text{ GeV}$$

Dark Dimension Scenario

[Arkani-Hamed, Dimopoulos, Dvali'98]

This scenario is an example of the **Large Extra Dimension models (ADD)**

although the scale is different than usual, since it was motivated by the smallness of the cosmological constant and not by the EW hierarchy problem:

$$\text{QG cut-off: } \hat{M} \sim m^{1/3} M_P^{2/3} \sim 10^{10} \text{ GeV}$$

Proposals for dark matter in this scenario: [Gonzalo et al'22] [Law-Smith et al'23] [Anchordoqui et al'23]

Dark Dimension Scenario

[Arkani-Hamed, Dimopoulos, Dvali'98]

This scenario is an example of the **Large Extra Dimension models (ADD)**

although the scale is different than usual, since it was motivated by the smallness of the cosmological constant and not by the EW hierarchy problem:

$$\text{QG cut-off: } \hat{M} \sim m^{1/3} M_P^{2/3} \sim 10^{10} \text{ GeV}$$

Proposals for dark matter in this scenario: [Gonzalo et al'22] [Law-Smith et al'23] [Anchordoqui et al'23]

Open challenges: We do not have a concrete string theory embedding including the SM!

Dark Dimension Scenario

[Arkani-Hamed, Dimopoulos, Dvali'98]

This scenario is an example of the **Large Extra Dimension models (ADD)**

although the scale is different than usual, since it was motivated by the smallness of the cosmological constant and not by the EW hierarchy problem:

$$\text{QG cut-off: } \hat{M} \sim m^{1/3} M_P^{2/3} \sim 10^{10} \text{ GeV}$$

Proposals for dark matter in this scenario: [Gonzalo et al'22] [Law-Smith et al'23] [Anchordoqui et al'23]

Open challenges: We do not have a concrete string theory embedding including the SM!

It will be **tested in future experiments** that will improve the precision measurements on deviations from Newton's law

New ISLE at the Conrad Observatory

[Aspelmeyer, Adelberger, Shayeghi, Zito...]

Conclusions

- ❖ Consistency with Quantum Gravity can have important implications for our universe at energies much below the Planck scale.

Conclusions

- ❖ Consistency with Quantum Gravity can have important implications for our universe at energies much below the Planck scale.
- ❖ Not every EFT is consistent with UV completion in Quantum Gravity, unless it satisfies the swampland constraints.

Conclusions

- ❖ Consistency with Quantum Gravity can have important implications for our universe at energies much below the Planck scale.
- ❖ Not every EFT is consistent with UV completion in Quantum Gravity, unless it satisfies the swampland constraints.
- ❖ Approximate global symmetries, weakly coupled gauge theories and large field ranges are disfavoured in Quantum Gravity
 - ➔ new towers of states become light yielding $\Lambda_{QG} \ll M_p$

Conclusions

- ❖ Consistency with Quantum Gravity can have important implications for our universe at energies much below the Planck scale.
- ❖ Not every EFT is consistent with UV completion in Quantum Gravity, unless it satisfies the swampland constraints.
- ❖ Approximate global symmetries, weakly coupled gauge theories and large field ranges are disfavoured in Quantum Gravity
 - ➔ new towers of states become light yielding $\Lambda_{QG} \ll M_p$
- ❖ Swampland constraints motivated by string theory motivate an scenario in which the smallness of our vacuum energy is tied to the existence of one mesoscopic extra dimension of

$$l \sim 0.1 - 10\mu m \quad \text{in our universe.}$$

Thank you!

If you want to hear more about the Swampland program:

Online series of Swampland seminars / open mic discussions

on Mondays at 11:00 am ET (5:00 pm CET)

You can subscribe here: <https://sites.google.com/view/swamplandseminars/>

Everybody is welcome! :)

back-up slides

Approximate global symmetries,
Weakly coupled gauge theories,
Large field ranges...

...come at a price.

(Swampland) Distance Conjecture (SDC):

There is an infinite tower of states becoming exponentially light at every infinite field distance limit of the moduli space

$$m(P) \sim m(Q)e^{-\alpha\Delta\phi} \quad \text{when} \\ \Delta\phi \rightarrow \infty$$

(geodesic distance)

[Ooguri-Vafa'06]

[Arkani-Hamed et al'06]

Weak Gravity Conjecture (WGC):

Given a gauge theory, there must exist an electrically charged state with

$$\frac{Q}{M} \geq \left(\frac{Q}{M}\right)_{\text{extremal}} = \mathcal{O}(1) \quad \begin{array}{l} Q=qg : \text{charge} \\ m : \text{mass in} \\ \text{Planck units} \end{array}$$

Strong version: there is a sublattice/tower of superextremal states

[Montero et al.'16][Heidenreich et al.'15-16][Andriolo et al.'18]

UV cut-off goes to zero
due to new light states

$$\Lambda \sim gM_p$$

$$\Lambda \sim M_p \exp(-\alpha\Delta\phi)$$

Evidence for WGC and SDC

❖ String theory compactifications: Plethora of quantitative tests!

- Systematic approach according to the level of supersymmetry
- Interesting connections to mathematics

[Grimm, Palti, IV'18]

[Grimm, Palti, Li'18]

[Lee, Lerche, Weigand'18-19]

...

❖ AdS/CFT:

- WGC proven for AdS3 using modular invariance of the CFT
- WGC from QI theorems and entanglement entropy
- SDC formulated in terms of a CFT Distance conjecture

[Heidenreich et al'16]

[Montero et al'16]

[Montero'18]

[Perlmutter et al'20]

❖ Black hole arguments:

- WGC follows from requiring black holes to decay
- WGC/SDC follows from entropy bounds associated to small BHs
- Connection between WGC and weak cosmic censorship

[Arkani-Hamed et al'06]

[Hamada et al'21]

[Crisford et al'17]

❖ Using positivity/unitarity bounds: lead to mild versions of the WGC

[Cheung et al'18][Hamada et al'18]...

WGC and SDC from Entropy Bounds

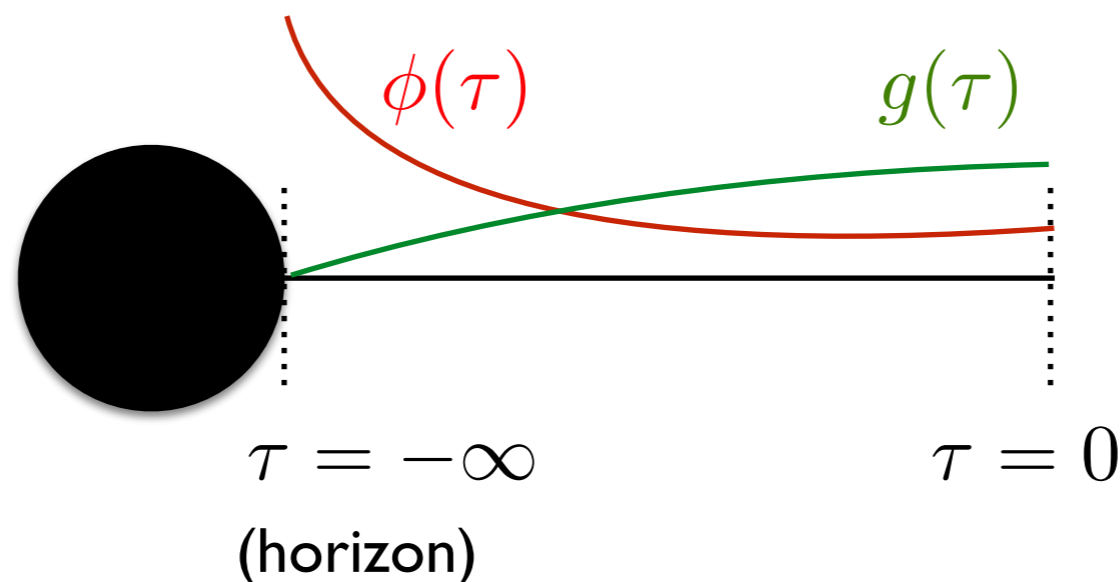
Take Einstein-Maxwell-Dilaton theory:

$$S = \int d^4x \sqrt{-g} \left[R + 2|d\phi|^2 + \frac{1}{2g(\phi)^2} |F|^2 \right] \quad \text{s.t.} \quad g(\phi) \rightarrow 0 \quad \text{as} \quad \phi \rightarrow \infty$$

There are electrically charged BH solutions with classical zero area (small BHs)

If $g(-\infty) \rightarrow 0$ then $A(-\infty) \rightarrow 0$: **Small BH**

BH induces a running of the scalar field and gauge coupling as approaching the horizon leading to:



large field range!
small gauge coupling!

WGC and SDC from Entropy Bounds

Small BHs lead to a violation of the Bekenstein bound, unless the EFT cutoff decreases as dictated by the SDC / WGC

Entropy Bound:

A region of size L cannot have more entropy than a Schwarzschild black hole of the same area $A = L^2$

$$N_{\text{species}} = Q_{\text{max}} \lesssim L^2 = A$$

Using extremality condition and that EFT breaks down at $|d\phi|^2 \sim \Lambda^2$



$$\Lambda \lesssim g \quad \text{in Planck units}$$

due to an infinite tower of states

Cobordism conjecture

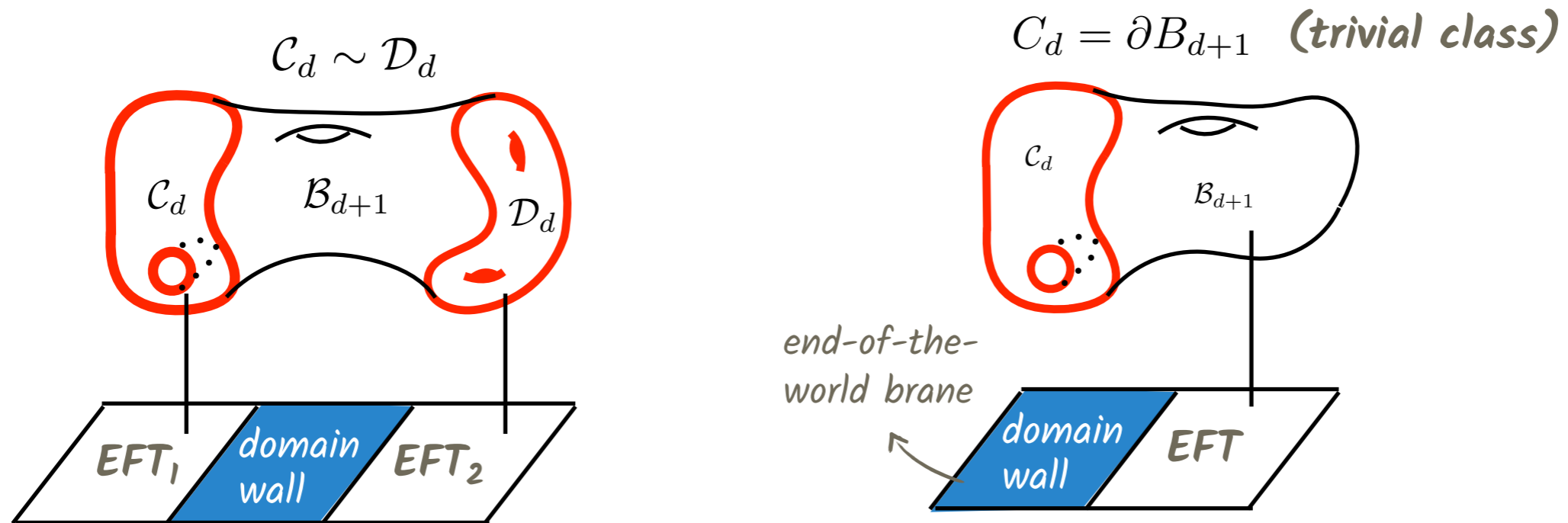
The cobordism group of a quantum gravity theory must be trivial:

$$\Omega_k^{QG} = 0 \quad [\text{McNamara, Vafa'19}]$$

k : internal dimension

D : total dimension

to avoid a $(D-k-1)$ -form global symmetry with charges $[M] \in \Omega_k^{QG}$

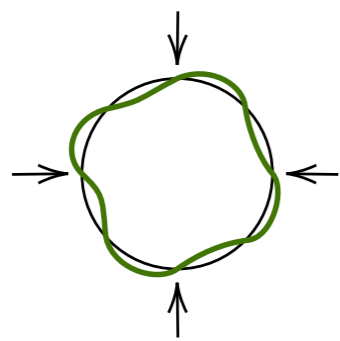


It implies all theories of same dimension are connected by finite energy domain walls, and predicts the existence of new defects in string theory!

Asymptotic Towers of States

Simplest example: Circle compactification of a string theory

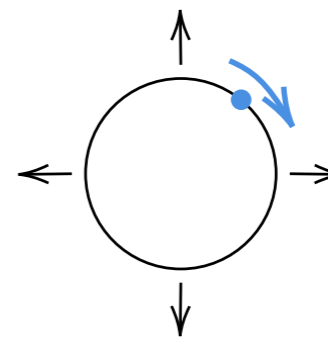
Tower of winding modes



m_ω

$r \rightarrow 0$

Tower of Kaluza-Klein states



m_{KK}

$r \rightarrow \infty$

Quantum gravity cut-off: $\Lambda \equiv M_{\text{pl},d+1} = \frac{M_{\text{pl},d}}{r} \rightarrow 0$ as $r \rightarrow \infty$

(same for T-dual theory as $r \rightarrow 0$)

Pattern

In all known string theory examples so far, it occurs that

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units,} \quad \text{as } V_0 \rightarrow 0$$

We can bound $2 \leq \alpha \leq d$ in quasi-dS space [Montero,Vafa,IV'22]

$\alpha \geq 2$: Higuchi bound: $m_{\text{tower}} \geq H$ since the tower contains higher spin fields

$\alpha \leq d$: Even if tree level is small, there is at least a **one-loop contribution**:

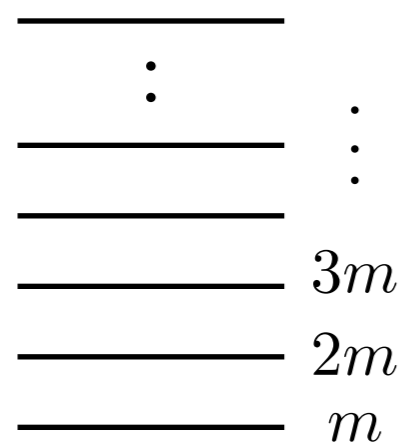
$$V \sim m^d \quad (\text{if non-susy})$$

$d =$ space-time dimension

Failure of IR EFT expectations

In all known string theory examples so far, it occurs that

$$V_0 \sim m_{\text{tower}}^\alpha \quad \text{in Planck units,} \quad \text{as } V_0 \rightarrow 0$$



It is important to consider the contribution of the entire tower of states (vs a finite number of fields):

- If integrating out the infinite tower: (e.g. Casimir energies)

$$V_0 \sim m_1^d \quad \text{first light state of the tower!}$$

- If integrating a finite number of fields below a cut-off:

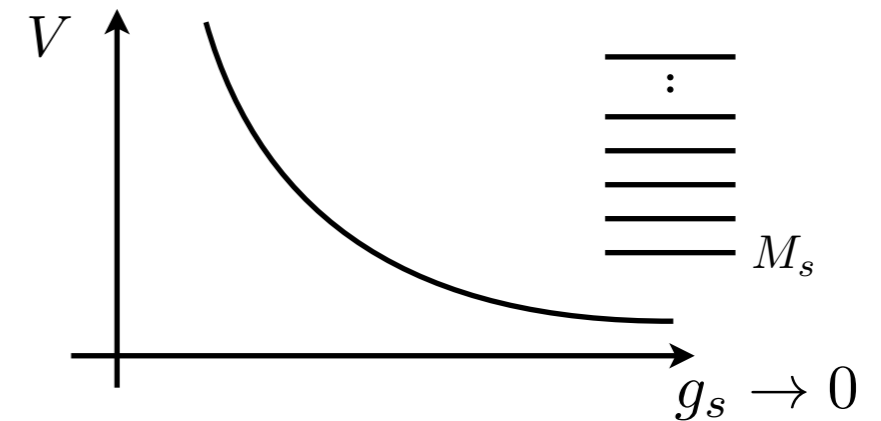
$$V_0 \sim m_{\text{heavy}}^d \quad \text{the heavy states dominate}$$

Non-SUSY string theory example

**SO(16)xSO(16) non-SUSY (tachyon-free)
heterotic string theory:**

Tower of string modes becoming light
in the weak coupling limit, starting at

$$m_{\text{tower}} = M_s$$



Positive runaway on the dilaton

$V_{\text{tree}} = 0$ by conformal invariance

$$V_{1\text{-loop}} \sim - \sum_i (-1)^{F_i} \int_{\Lambda_{UV}^{-2}}^{\infty} \frac{ds}{s^6} \exp\left(-\frac{m_i^2 s}{2}\right) \quad \rightarrow \quad V \sim m_{\text{tower}}^{10}$$

Contribution of massive string excitations is “cut-off” at M_s due to modular invariance (to avoid redundancy of counting the same states more than once)

More massive NS5-branes do not contribute either, they are composite objects