

# Subatomic and Astroparticles Physics Master Oral Defence

## **INNOVATIVE METHODS FOR ELECTROMAGNETIC CALORIMETERS BASED ON SCINTILLATING CRYSTALS**

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# **1. Calorimetry background**

- ❏ Radiation-matter interaction
- High energy **electrons** → Bremsstrahlung High energy **photons**  $\rightarrow$  Pair production Radiation length  $X_0$ ⇒ **electromagnetic showers.**



High energy **hadrons**  $\rightarrow$  complex interactions Interaction length λ

⇒ **hadronic showers** made of two components: **electromagnetic and hadronic components**.

## ❏ Response of the calorimeter



➔ **Electromagnetic** (**e.m.**) showers:





**e -**

calorimeter

### ➔ **Hadronic** (**h.**) showers:



with e, h calibration constants and **f em** , the electromagnetic fraction of the shower ( = fraction of hadron energy deposited via e.m. processes).

➔ **Different responses** depending on the showers ⇒ Worsens energy resolution.

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### ➔ **Hadronic** (**h.**) showers:



 $\underline{\text{Signal}_{\text{output}}}$  $f_{em})$  $|e f_{em}| + |$  $|h\left( 1 \right)$ **e.m. h.**

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- ➔ **Different responses** depending on the showers ⇒ Worsens energy resolution.
- ➔ **f em fluctuates** with the energy ⇒ Proportionality lost in hadronic showers.



# **Outlines**

- **1. Calorimeter background**
- **2. Dual-readout calorimetry**
- **3. Simulation**
- **4. Pulse shape discrimination**
- **5. Reconstruction energy**
- **6. Conclusion & Ongoing work**

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## ● **Homogeneous calorimeters based on scintillating crystals**

 $\rightarrow$  High-energy particles interact with the crystals that compose the calorimeter by depositing their energy converted to photons from 2 independent processes:

→ **Scintillation** (isotropic)

→ **Cherenkov** (directional & prompt light, mainly emitted by **e.m.** particles).



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- ➔ Energy detection technique that improves measurement accuracy by using **two types of readings** to correct complex hadron interactions in calorimeters:
	- Scintillation (**S**) signals
	- Cherenkov (**C**) signals
- Both signals can be calibrated with electrons of known E energy.



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- $\triangleright$  Both signals can be calibrated with electrons of known E energy.
- $\triangleright$  The dual-readout method works thanks to the fact that  $(e/h)_{s \neq 0}$ **(e/h)**<sub>C</sub>. The larger the difference between both values, the better the e.m. shower fraction  $f_{\rm em}$  can be extrapolated.

$$
\left\{ \begin{aligned} S & = E \left[ f_{em} + \frac{1}{\left( e/h \right)_S} (1 - f_{em}) \right] \\ C & = E \left[ f_{em} + \frac{1}{\left( e/h \right)_C} (1 - f_{em}) \right] \end{aligned} \right.
$$



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$$



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$$

$$
\boxed{f_{em}=\frac{\left(h/e\right)_C-\left(C/S\right)\left(h/e\right)_S}{\left(C/S\right)\left[1-\left(h/e\right)_S\right]-\left[1-\left(h/e\right)_C\right]}}
$$



 $\Rightarrow$  Use of 2 photodetectors to extract signals and obtain the (C/S) ratio. ⇒ Objective : extract the **f em from the pulse shape** using 1 photodetector

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# **3. Simulation**

- $\triangleright$  **Homogeneous** calorimeter made of 3x3 **PbWO**<sub>4</sub> crystals of dimensions: **2.2 x 2.2 x 25 cm3 (~ CMS ECAL**).
- $\triangleright$   $\pi^+$  and **e** are shot in the center crystal of the calorimeter.



$$
\triangleright \quad f_{em} \text{ implementation:} \quad \left( \quad f_{em,sim} = \frac{E_{em}}{E_{tot}} \quad \right)
$$

- $\triangleright$  Data extracted from the simulation:
	- Numbers of **C** and **S** photons detected at the PMT
	- **Pulse** shape at the PMT
	- **○ fem,sim**

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❏ Some features extracted from the pulse

Ratio of C and S areas of the pulse (a.p.)





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$$
\bigg|\,E^{\text{a. calib.}}_{reco,\pi^+}(f_{em}) = E^{\text{b. calib.}}_{reco,\pi^+} + \langle E_{reco,e^-} \rangle - [a f_{em,sim} + b]\bigg|
$$





### Tables with the results obtained from the different **features** used to do the calibration:





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# **5. Conclusion & Ongoing work**

- ➢ Proportionality shown between the **features extracted from the pulse shape** with the **f em** . ➢ Maximizing the information from the ECAL by extracting proportional features of the **f em** information from the **pulse shape** at the photodetector event by event.
	- ⤷ Information extracted from **only 1 photodetector** instead of using 2 as done normally.
- Possibility of combining features to estimate the  $f_{em}$  using a machine learning algorithm :
	- $\frac{1}{2}$  Reconstruction of the energy with a better resolution.

## Next steps:

- ❖ Improvement on the ML algorithm.
- ❖ Applying the techniques developed to a realistic configuration, as in a testbeam or a full detectors.
- ❖ Including a HCAL in simulation to allow full containment of the shower and evaluate the influence of the ECAL with dual readout capability on the performance, the results will be evaluated with experimental measurements in testbeam. From: Lucchini M.T., et al.,, JINST 15 P11005, 2020 [4]



# **Thank you for your attention !**

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# **References**

[1] - Roldan, Pablo & Lecoq, Paul, Quality control and preparation of the PWO crystals for the electromagnetic calorimeter of CMS. (2011)

[2] - Auguste Besson, extract of chapter 6 "High Energy" of "Particle interaction with matter In nuclear and particle physics" lecture, Strasbourg's University. (2022)

[3] - Sehwook Lee, Michele Livan, Richard Wigmans, Dual-readout calorimetry, Reviews of Modern Physics, volume 90 , https://doi.org/10.1103/RevModPhys.90.025002. (2018)

[4] - Lucchini M.T., et al., New perspectives on segmented crystal calorimeters for future colliders, JINST 15 P11005. (2020)

[5] - Particle Data Group. "Review of Particle Physics". Progress of Theoretical and Experimental Physics (PTEP) 083C01. (2022)

# **Backup slides**

### ❏ Radiation-matter interaction



Figure 1 : Fractional energy loss per radiation length in lead as a function of electron or positron energy



Figure 2 : Photon total cross sections as a function of energy in carbon and lead, showing the contributions of different processes



## ❏ HCAL dual-readout capabilities

S-C diagram





 $\circ$  50 x 50 x 200 cm<sup>3</sup> ( $\sim$  HCAL) to test the dual-readout capabilities of the simulated calorimeter.

## ❏ ECAL dual-readout capabilities







$$
4.37~{\rm GeV} \leqslant E_{depo} \leqslant 4.81~{\rm GeV}
$$

## ❏ Linearity for e.m. showers



 $\mathsf{f}_{\mathsf{em}}$ 

 $-0.9$ 

 $0.8$ 

 $\times 10^{7}$ 

 $280$ otorla be

260

240

 $\Box$  Features -  $(C/S)_{a.p.}$ 



 $\Box$  Features -  $(C/S<sub>late</sub>)<sub>a.p.</sub>$ 



❏ Features - ToT@5%



❏ Features - ToT@10%



❏ Features - ToT@40%



❏ Features - Rise time





### ❏ Gradient Boost Decision Tree

 $\rightarrow$  an ensemble technique that **combines the predictions of multiple decision trees to produce a more accurate and robust model**.

**1. An initial model**: often a simple one like the mean of the target values for regression.

**2. Building Trees Iteratively**: new decision trees added one at a time, each new tree trained to correct the errors or residuals of the combined ensemble of all previous trees. **3. Gradient Descent** to minimize a loss function. **4. Updating the mode**l: predictions from each new tree added to the ensemble model, weighted by a learning rate. **5. Repeat** 



From: Deng, Haowen, et al., BMC Medical Informatics and Decision Making, 2021

### ❏ Preliminary results with ML



### $\rightarrow$  33% of data for testing  $\rightarrow$  score = 0.82

### ❏ Preliminary results with ML



 $\rightarrow$  33% of data for testing  $\rightarrow$  score = 0.96

Influence of the incident and deposited energies in the linear  $\mathbf{A}$ fit to reconstruct the deposited energy inside an ECAL



Table 1: Values of the linear fit to reconstruct the deposited energy for different incident energies and sections of deposited energy. The incident and deposited energy are expressed in GeV. In the cell, the first value corresponds to the slope of the fit and the second below to the constant coefficient. They are also expressed in GeV.

### ❏ Hadronic shower



### Hadronic shower

- . Charged Hadronic particle are slowed down by dEdx  $\sqrt{\text{Small effect}}$
- . In addition collision with nuclei (strong interaction)
	- $\checkmark$  Inelastic scattering: intial state  $\neq$  final state
	- $\checkmark$  At high energy:  $\sigma_r \sim 40$  mb  $\times$  A<sup>0.71</sup>
- $\cdot$  Interaction length  $\lambda$ <sub>r</sub>:

 $\checkmark$  e.g. Carbon:  $\rho.\lambda \sim 86$  g/cm<sup>2</sup>  $\checkmark$  e.g. Lead:  $\rho.\lambda \sim 194$  g/cm<sup>2</sup>

$$
= \frac{A}{\mathcal{N}_{A\nu}\rho\sigma_l} \qquad \text{which yields:} \qquad N(x) = N_0 \exp\left(-\frac{1}{2}\int_{-\infty}^{\infty} \frac{1}{N(x)}\right)
$$

 $\lambda_{\tau} \gg X_0$ 

- ✔ Hadronic showers longer than e.m. showers
- ~1/3 of the shower produces  $\pi^0 \Rightarrow$  quickly stopped by e.m. shower

 $\lambda_I$ 

- Energy threshold to produce a pion (lightest hadrons)  $\bullet$ 
	- ✔ Play a similar role of E, for electrons
	- ✔ Hadronic shower stops below this energy
- · Large variety of process
	- √ A fraction of the energy loss can not be converted in signal
		- " (nuclear binding energy, target recoil, etc.
	- ✔ Large fluctuations in hadronic showers
	- $\checkmark$  Energy measurement less precise

2022-23

A.Besson, Université de Strasbourg

### From: A. Besson, lectures in M2PSA 2022

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 $-x/\lambda_{\rm int}$