

Branes and Quivers

In honor of Philippe Di Francesco's @ 60

Amihay Hanany Jun 2024

Branes and Supersymmetric Gauge Theories

Outline

- Branes on Gorenstein singularities
- Brane Tilings
- Cluster Algebras
- Branes and tropical geometry
- Branes and symplectic singularities

Branes on Gorenstein Singularities

1998

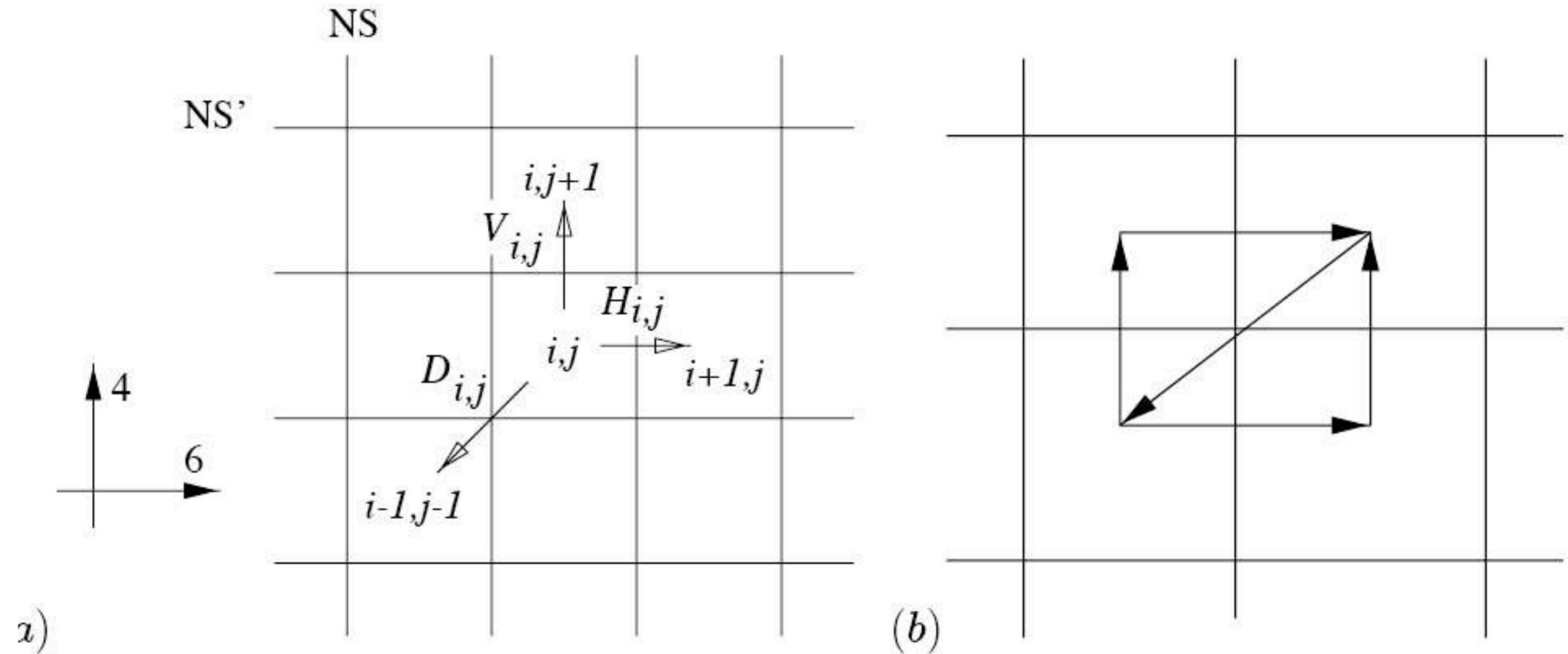
Type IIB superstring (10)

D3 branes (4) on Gorenstein singularities (6)

- N D3 branes on Gorenstein singularities of cplx dim 3
- Fractional branes
- Open strings between the different branes
- Gauge fields, matter fields, interactions — Quiver Q with a superpotential W
- Problem — compute the Q & W for a given singularity

Brane Boxes

1998



$$W = \sum_{i,j} (H_{i,j} V_{i+1,j} D_{i+1,j+1} - V_{i,j} H_{i,j+1} D_{i+1,j+1})$$

Quiver & Superpotential $\mathbb{C}^3 / (\mathbb{Z}_3 \times \mathbb{Z}_3)$

Graph taken from the work of Philippe & Jean-Bernard (1989)

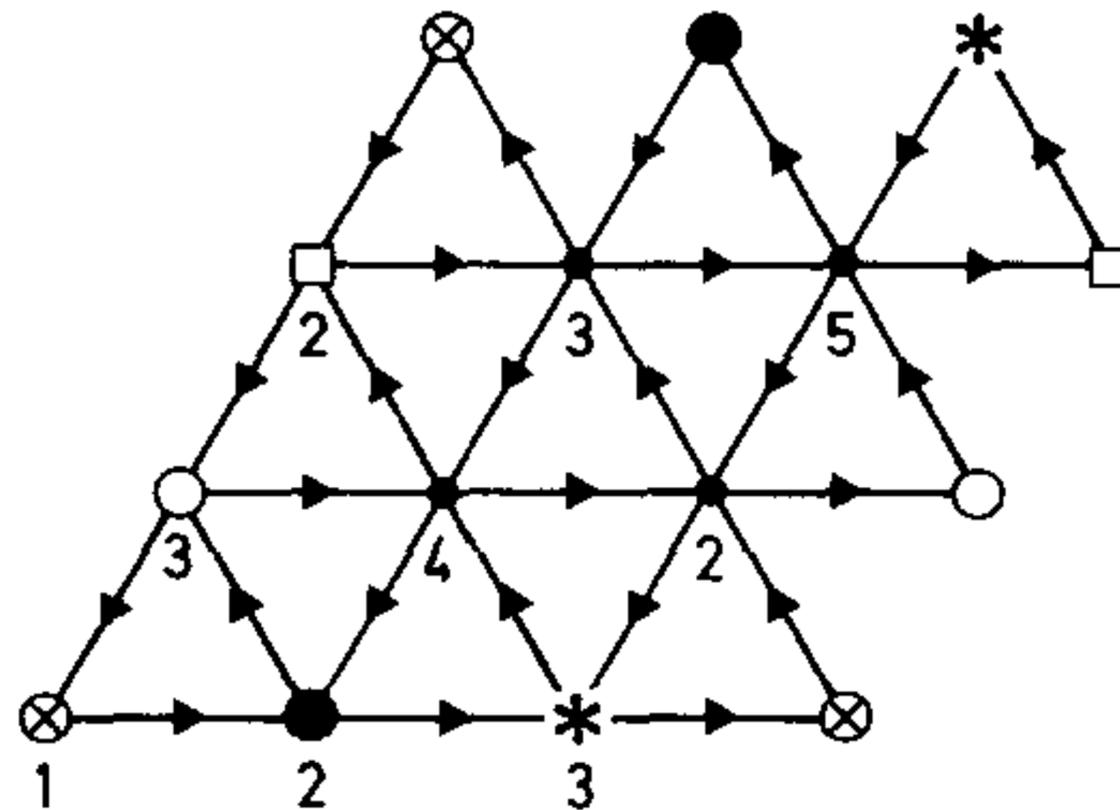


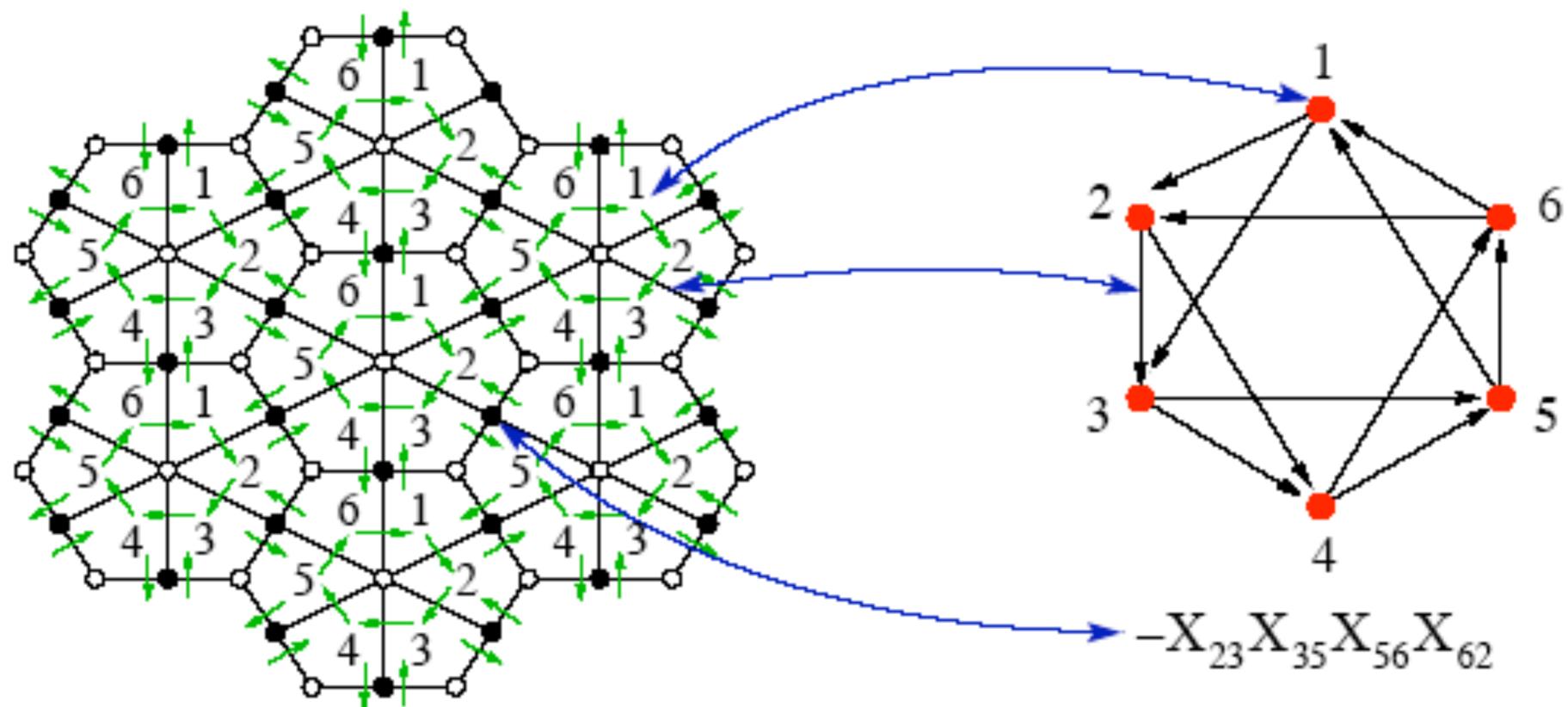
Fig. 9. Graph of the $\mathbb{Z}_3 \times \mathbb{Z}_3$ subgroup of $SU(3)$; points with the same symbol (\bullet , \otimes , etc.) have to be identified.

Branes Tilings

2005

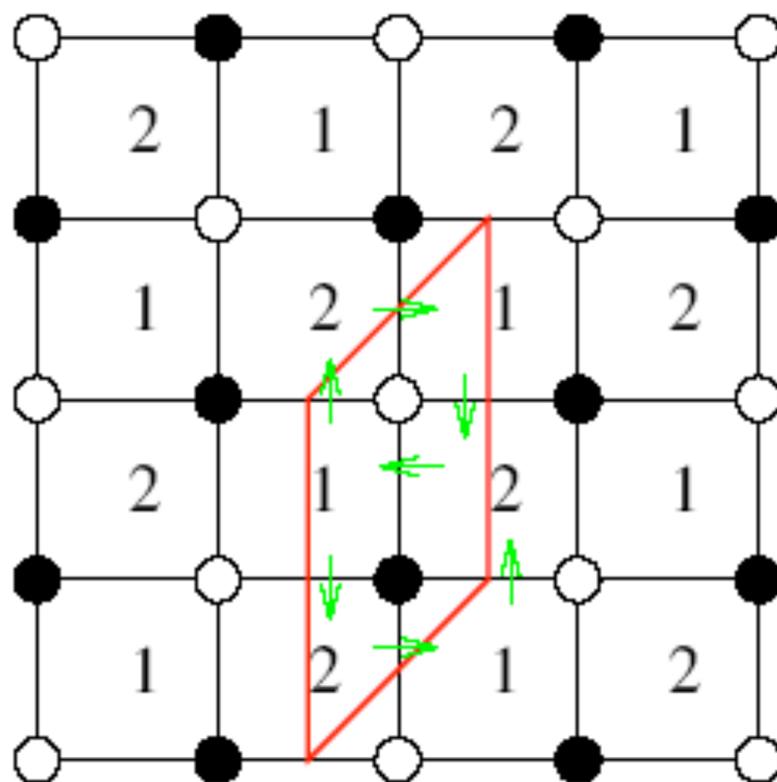
Brane Tilings

2005

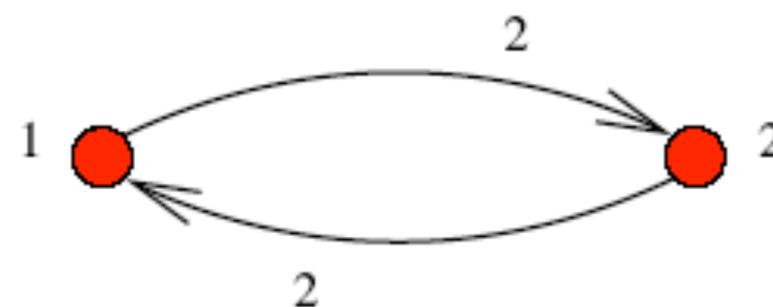


Conifold tiling

2005



brane tiling

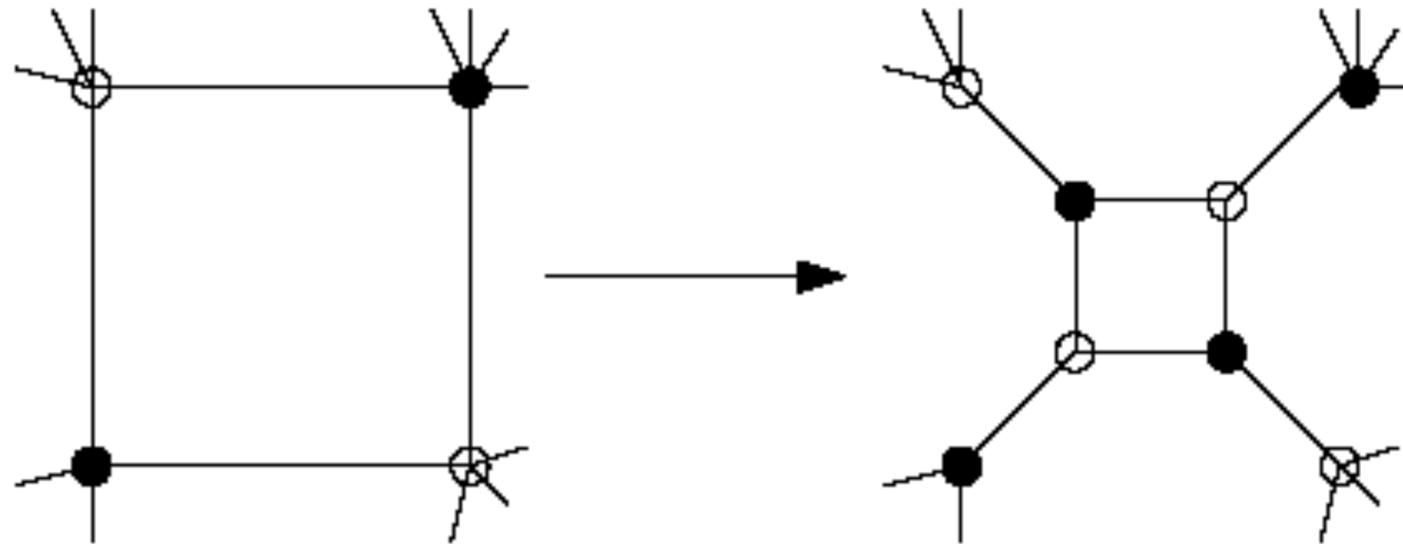


quiver

$$W = X_{12}^{(1)} X_{21}^{(1)} X_{12}^{(2)} X_{21}^{(2)} - X_{12}^{(1)} X_{21}^{(2)} X_{12}^{(2)} X_{21}^{(1)}$$

Seiberg Duality

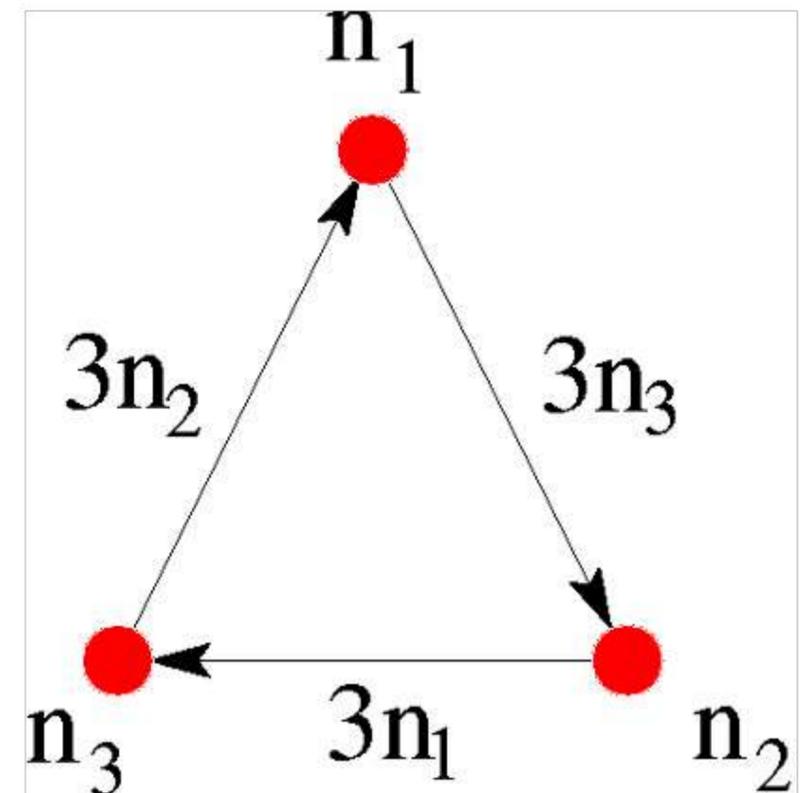
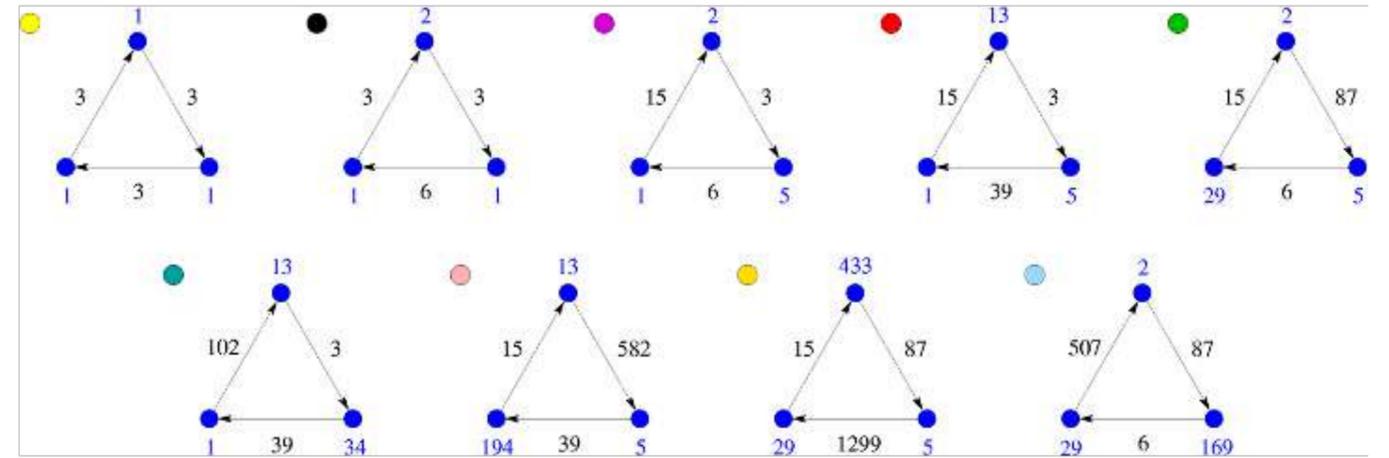
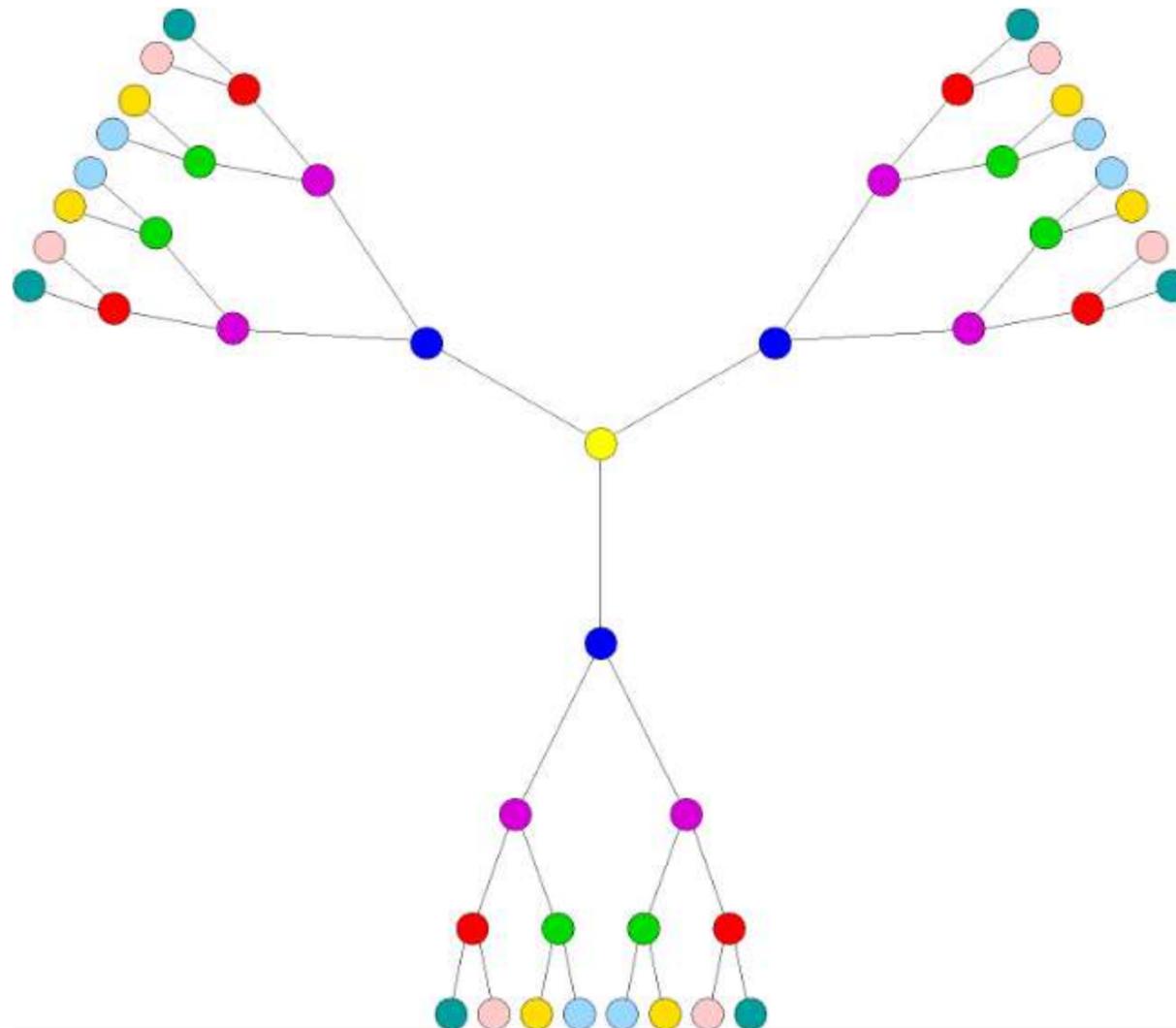
Urban Renewal – quiver mutation



Cluster Algebras

Duality tree for dP_0 , also known as $\mathbb{C}^3/\mathbb{Z}_3$

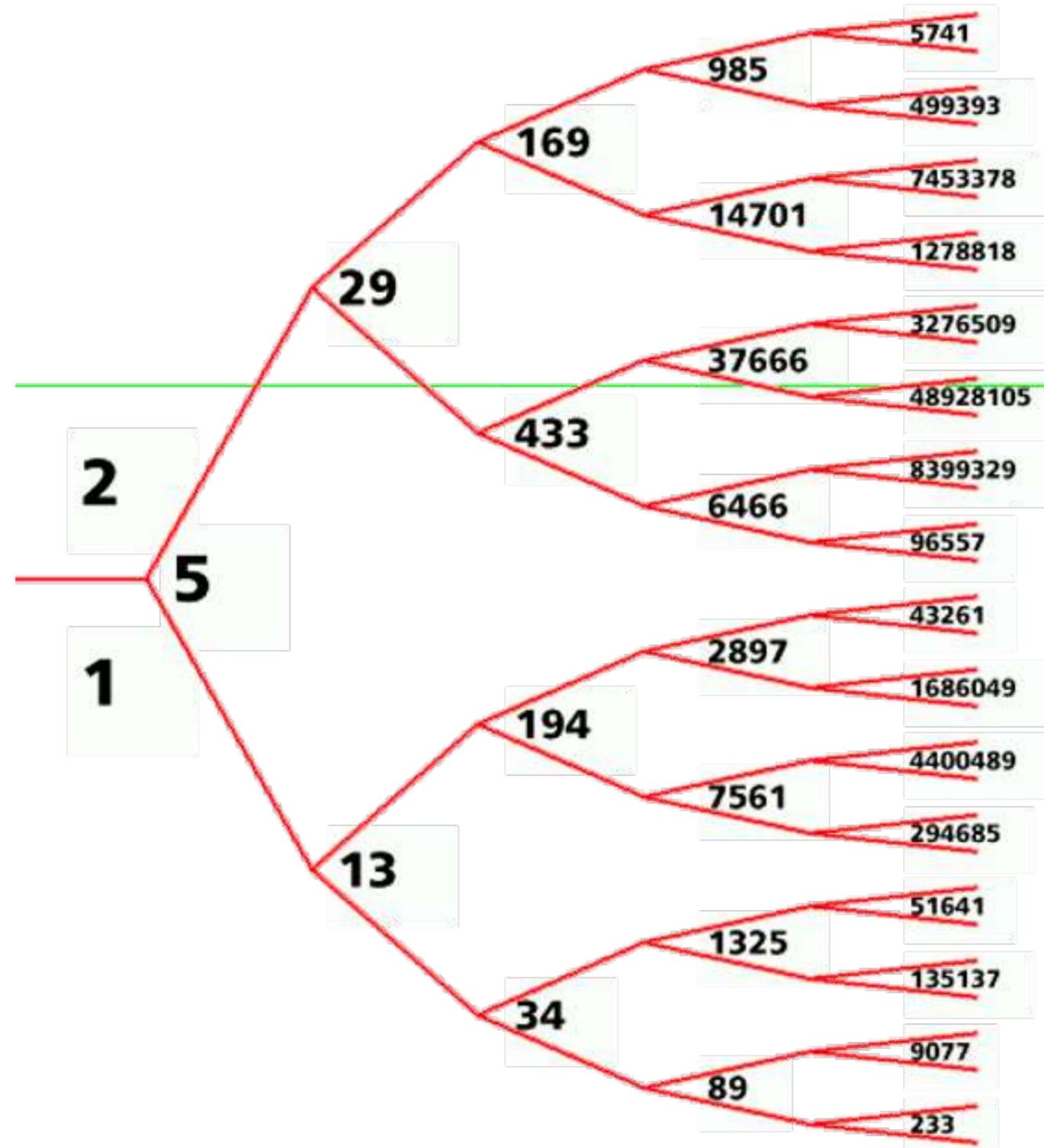
2003



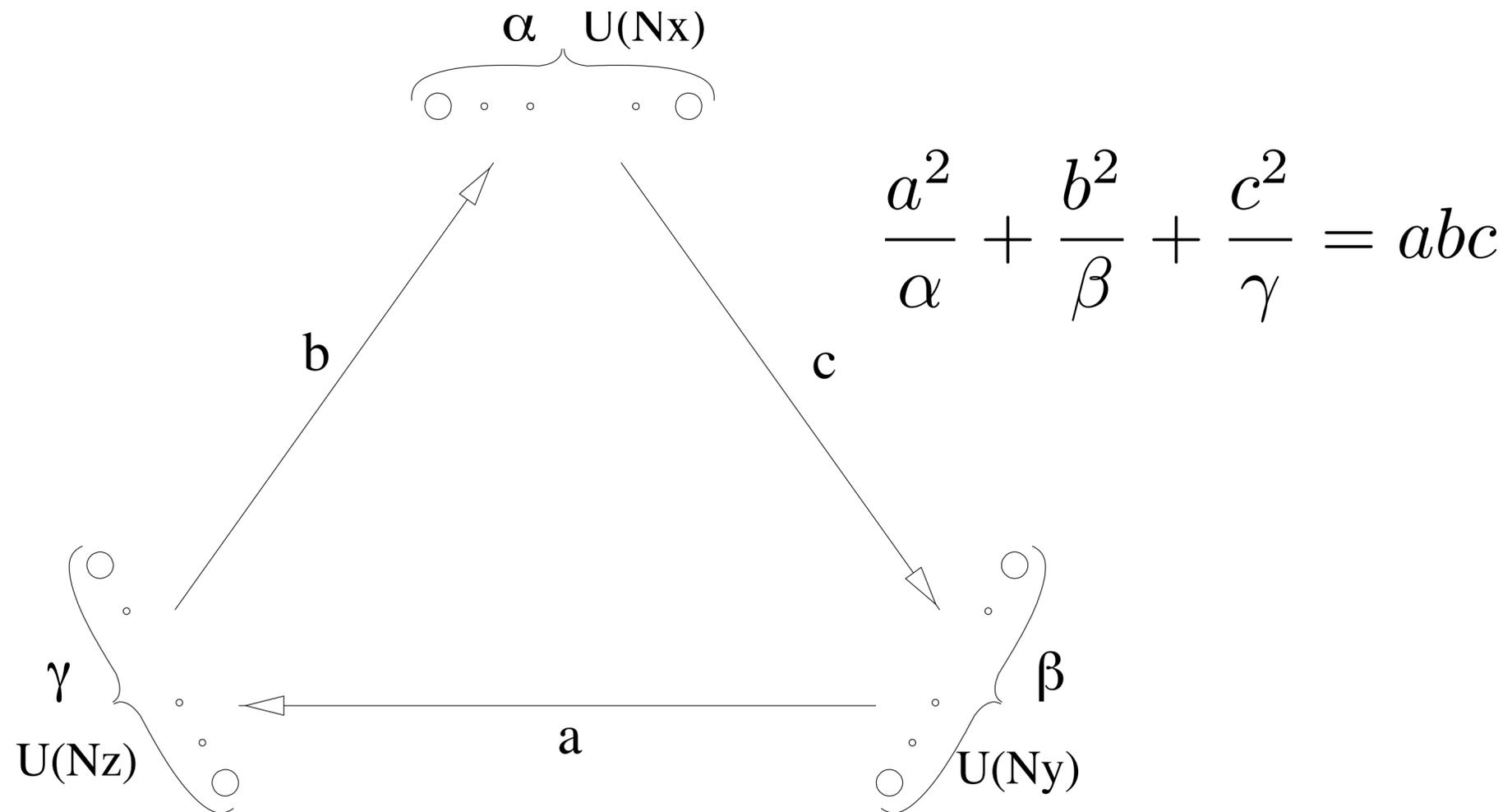
Diophantine (Markov) equation

$$n_1^2 + n_2^2 + n_3^2 = 3n_1n_2n_3$$

Markov Numbers



Quiver 3 block model



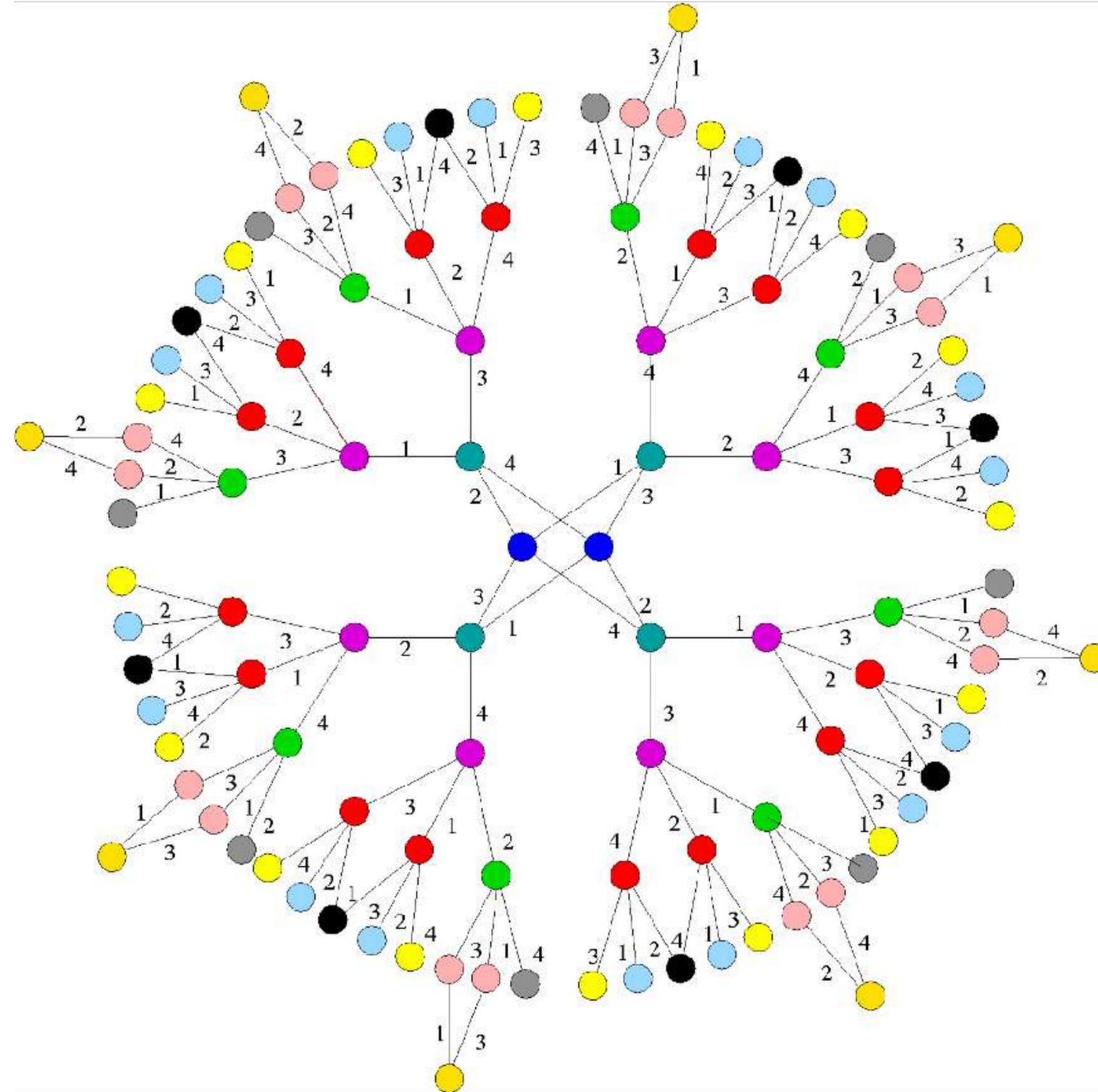
$$R_a = \frac{2}{abc} \frac{a^2}{\alpha}; \quad R_b = \frac{2}{abc} \frac{b^2}{\beta}; \quad R_c = \frac{2}{abc} \frac{c^2}{\gamma}$$

Quiver	α	β	γ	a	b	c	x	y	z	K^2	surface
	1	1	1	3	3	3	1	1	1	9	dP_0
	1	1	2	2	2	4	1	1	1	8	F_0
	1	2	3	1	2	3	1	1	1	6	dP_3
	1	1	5	1	2	5	1	2	1	5	dP_4
	2	2	4	1	1	2	1	1	1	4	dP_5
	3	3	3	1	1	1	1	1	1	3	$dP_6 I$
	2	1	6	1	1	3	1	2	1	3	$dP_6 II$
	4	4	2	1	1	1	1	1	2	2	$dP_7 I$
	1	1	8	1	1	4	2	2	1	2	$dP_7 II$
	3	1	6	1	1	2	1	3	1	2	$dP_7 III$
	1	1	9	1	1	3	3	3	1	1	$dP_8 I$
	8	2	1	2	1	1	1	2	4	1	$dP_8 II$
	2	3	6	1	1	1	3	2	1	1	$dP_8 III$
	5	5	1	1	2	1	1	2	5	1	$dP_8 IV$
	2	2	1	2	2	2	1	1	2	4	$sh dP_5$
	2	1	4	2	1	2	2	2	1	2	$sh dP_7$

Table 1: The complete list of the minimal models for 3-block chiral quivers.

Duality Flower for $\mathbb{F}_0 = \mathbb{P}^1 \times \mathbb{P}^1$

2003

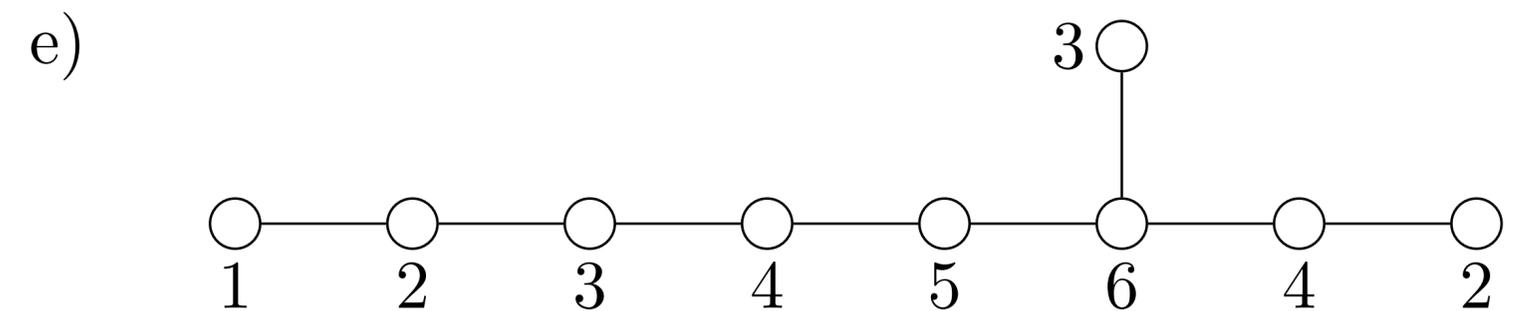
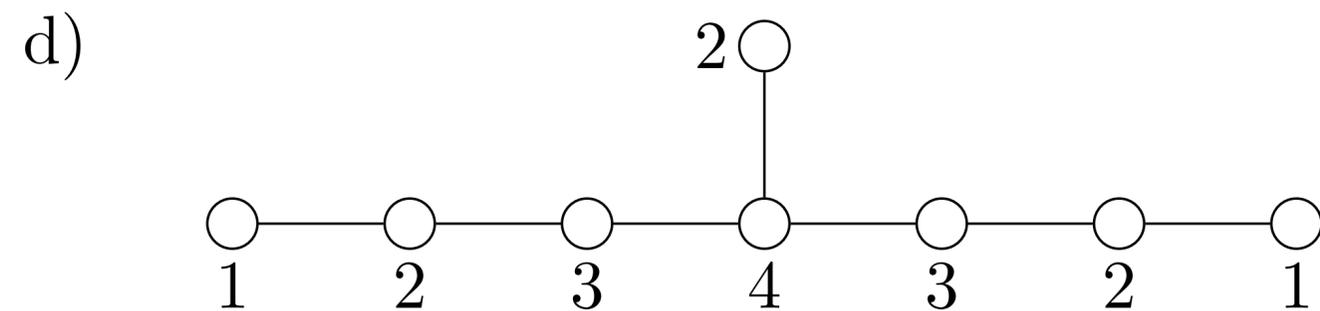
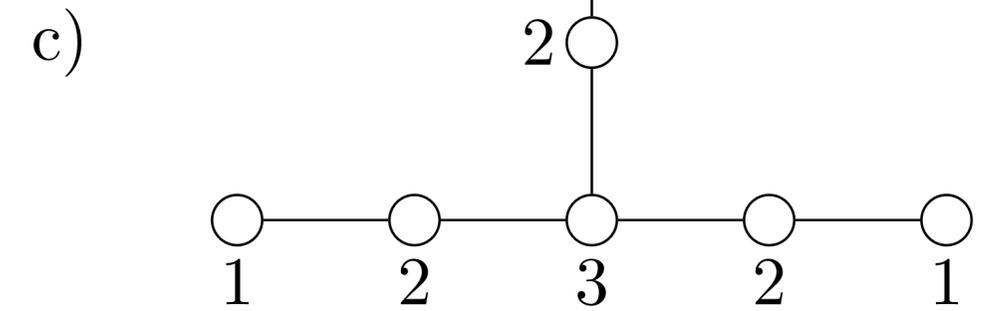
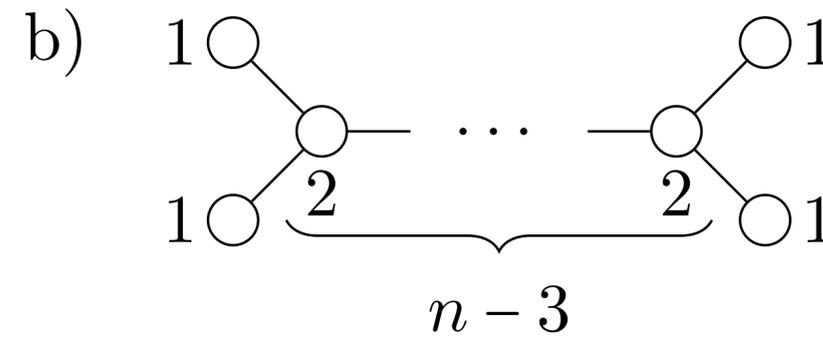
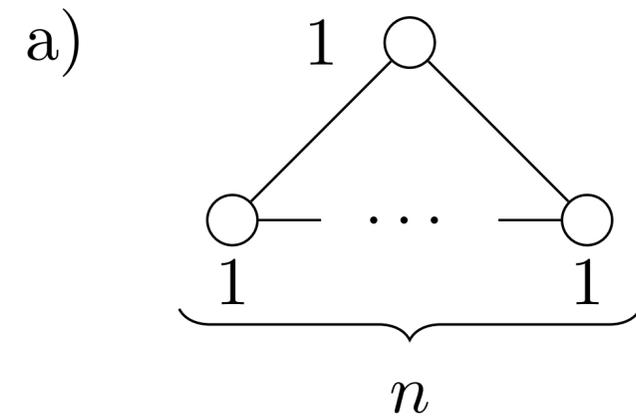


Coulomb Branch

2013

Affine ADE quivers

Coulomb branch – closure of minimal nilpotent orbit



Monopole formula — the ingredients

per each node of label k

- $W = S_k$ — the Weyl group of $GL(k)$
- $\hat{\Lambda}$ — The (Langlands) dual lattice — lattice of coweights
- A set of integer numbers $\hat{\Lambda} = \mathbb{Z}^k \ni m = (m_1, \dots, m_k)$ — magnetic charges (coweights)
- $\hat{\Lambda}/W$ — Principal Weyl chamber $m_1 \leq \dots \leq m_k$
- Boundaries of the Weyl chamber — when some m_i coincide
- H_m — stabilizer of m in $GL(k)$ — a Levi subgroup of $GL(k)$
- d_i^m — degrees of Casimir invariants of H_m

The conformal dimension — $\Delta(m)$

\mathbb{C}^* grading on the Coulomb branch

- Given a quiver with a set of nodes, each with labels k_a
- $\Delta(m)$ is a sum of contributions from nodes and edges:
- For each node with magnetic charges $m_i^a, i = 1 \dots k_a$ there is a negative contribution
 - $-\sum_{1 \leq i < j \leq k_a} |m_i^a - m_j^a|$ (associated with positive roots of $GL(k_a)$)
- For each edge connecting nodes a, b with magnetic charges m_i^a and m_j^b a positive contribution
 - $\frac{1}{2} \sum_{i=1}^{k_a} \sum_{j=1}^{k_b} |m_i^a - m_j^b|$ (associated with bifundamental representation)

The monopole formula

Hilbert series of the Coulomb branch

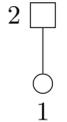
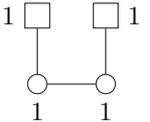
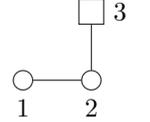
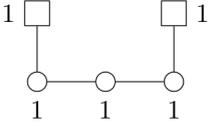
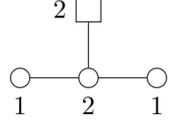
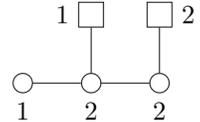
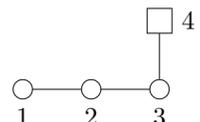
- Given a quiver with all the ingredients defined so far
- Introduce a variable t
- The Hilbert series is given by (flavor nodes have fixed m . Set to 0.)

- $$H(t) = \sum_{m \in \hat{\Lambda}/W} t^{2\Delta(m)} P_m(t)$$

- $$P_m(t) = \prod_i \frac{1}{1 - t^{2d_i^m}}$$

Examples — from the world of nilpotent orbits

Simple quivers and their Hilbert Series

Nilpotent Orbit	$\text{Dim}_{\mathbb{H}}$	Quiver	HS	HWG
[1, 1]	0	-	1	1
[2]	1		$\frac{(1-t^4)}{(1-t^2)^3}$	$\frac{1}{(1-\mu^2 t^2)}$
[1, 1, 1]	0	-	1	1
[2, 1]	2		$\frac{(1+4t^2+t^4)}{(1-t^2)^4}$	$\frac{1}{(1-\mu_1 \mu_2 t^2)}$
[3]	3		$\frac{(1-t^4)(1-t^6)}{(1-t^2)^8}$	$\frac{(1-\mu_1^3 \mu_2^3 t^{12})}{(1-\mu_1 \mu_2 t^2)(1-\mu_1 \mu_2 t^4)(1-\mu_1^3 t^6)(1-\mu_2^3 t^6)}$
[1, 1, 1, 1]	0	-	1	1
[2, 1, 1]	3		$\frac{(1+t^2)(1+8t^2+t^4)}{(1-t^2)^6}$	$\frac{1}{(1-\mu_1 \mu_3 t^2)}$
[2, 2]	4		$\frac{(1+t^2)^2(1+5t^2+t^4)}{(1-t^2)^8}$	$\frac{1}{(1-\mu_1 \mu_3 t^2)(1-\mu_2^2 t^4)}$
[3, 1]	5		$\frac{(1+t^2)(1+4t^2+10t^4+4t^6+t^8)}{(1-t^2)^{10}}$	$\frac{(1-\mu_1^3 \mu_2^3 \mu_3^3 t^{12})}{(1-\mu_1 \mu_3 t^2)(1-\mu_2^2 t^4)(1-\mu_1 \mu_3 t^4)(1-\mu_1^2 \mu_2 t^6)(1-\mu_2 \mu_3^2 t^6)}$
[4]	6		$\frac{(1-t^4)(1-t^6)(1-t^8)}{(1-t^2)^{15}}$	messy

Quivers for hyper surface symplectic singularities

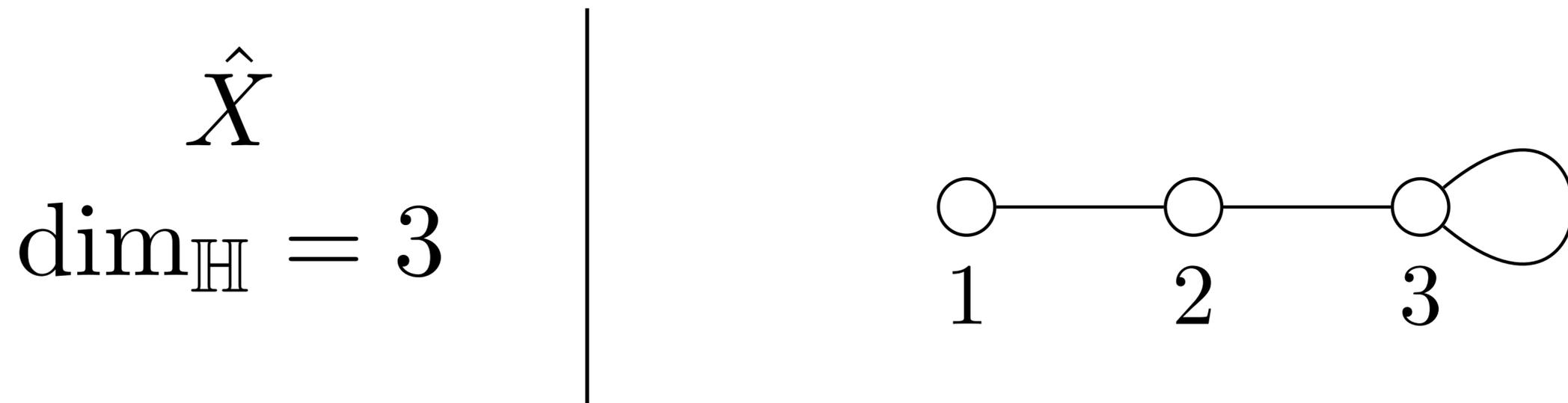
Slices in $\mathrm{Sp}(n)$

$$\begin{array}{c} X_n \\ \dim_{\mathbb{H}} = 2 \end{array} \left| \mathcal{H} \left(\begin{array}{c} \text{---} \text{---} \text{---} \text{---} \text{---} \\ \underbrace{\hspace{10em}} \\ \text{Graph with } n \text{ loops} \end{array} \right) = \mathcal{C} \left(\begin{array}{c} \text{---} \text{---} \text{---} \text{---} \text{---} \\ \underbrace{\hspace{10em}} \\ \text{Graph with } n \text{ loops} \end{array} \right)$$

- $PE \left[\mu^2 t^2 + t^4 + \mu (t^{2n-1} + t^{2n+1}) - \mu^2 t^{4n+2} \right]_{SU(2)}$

Quivers for hyper surface symplectic singularities

Slice in G2



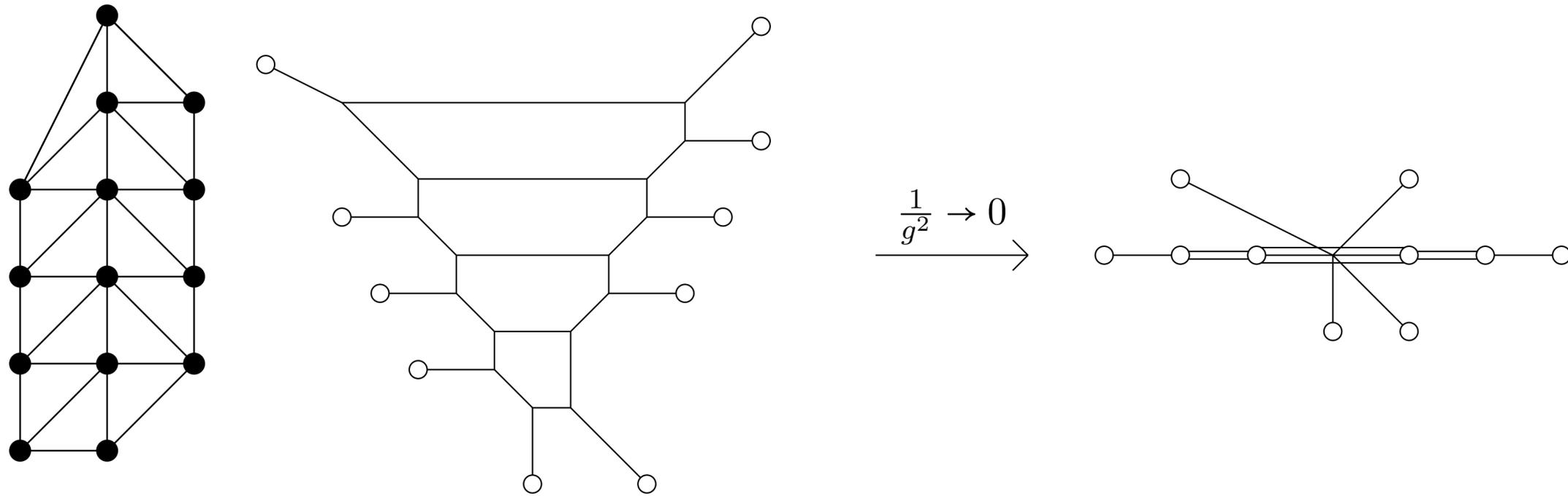
- $H(t, x) = PE \left[[2]t^2 + [3]t^3 - t^{12} \right]$
- Not polynomial PE

Branes

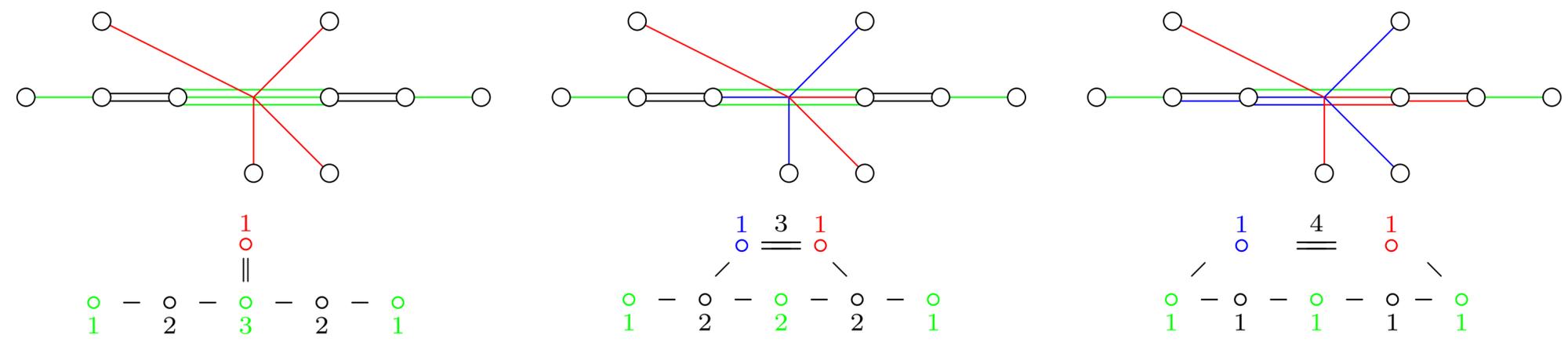
Union of 3 cones

new physics

$$\mathcal{H}_\infty \left(\begin{array}{c} 6 \\ \square \\ \circ \\ SU(5)_1 \end{array} \right) = C_1 \cup C_2 \cup C_3$$



$\xrightarrow{\frac{1}{g^2} \rightarrow 0}$



$$C_1 = \mathcal{C}^{3d} \left(\begin{array}{c} 1 \\ \circ \\ \parallel \\ \circ - \circ - \circ - \circ - \circ \\ 1 \quad 2 \quad 3 \quad 2 \quad 1 \end{array} \right)$$

$$C_2 = \mathcal{C}^{3d} \left(\begin{array}{c} 1 \quad 3 \quad 1 \\ \circ \quad \equiv \quad \circ \\ \diagup \quad \quad \diagdown \\ \circ - \circ - \circ - \circ - \circ \\ 1 \quad 2 \quad 2 \quad 2 \quad 1 \end{array} \right)$$

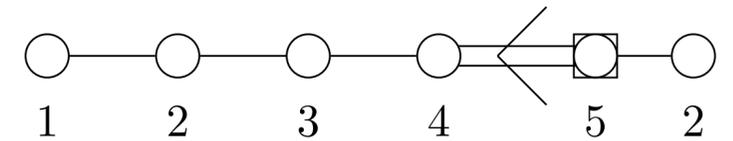
$$C_3 = \mathcal{C}^{3d} \left(\begin{array}{c} 1 \quad 4 \quad 1 \\ \circ \quad \equiv \quad \circ \\ \diagup \quad \quad \diagdown \\ \circ - \circ - \circ - \circ - \circ \\ 1 \quad 1 \quad 1 \quad 1 \quad 1 \end{array} \right)$$

Non simply laced quivers

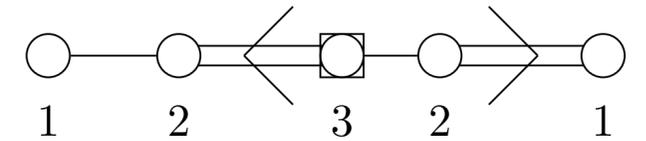
No known Lagrangian or path integral

- A small modification of the monopole formula
- A whole new set of moduli spaces
- A window to exotic moduli spaces
- like rank 1 4d theories

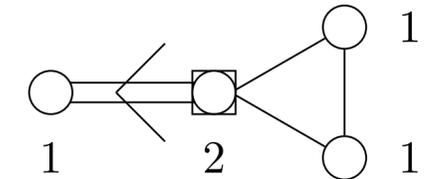
C_5



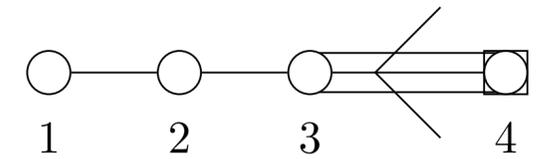
$C_3 \times A_1$



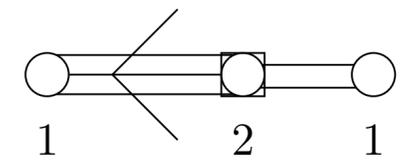
$C_2 \times U_1$



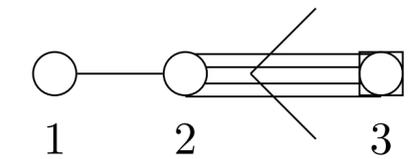
A_3



$A_1 \times U_1$

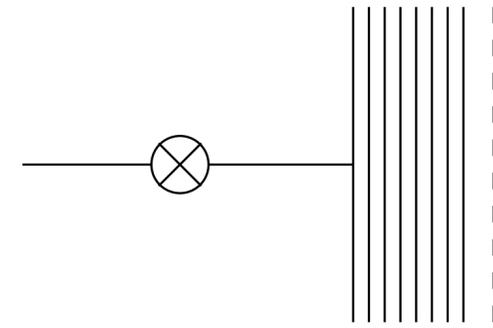
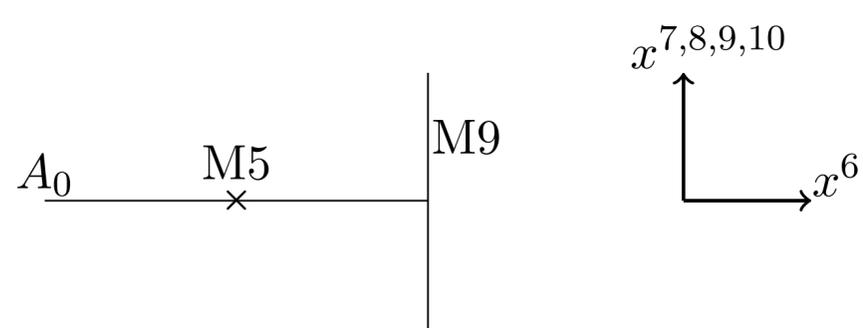


A_2



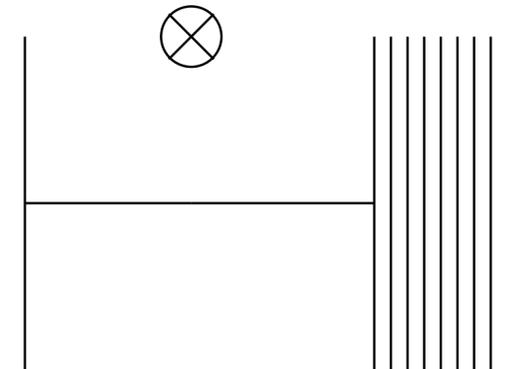
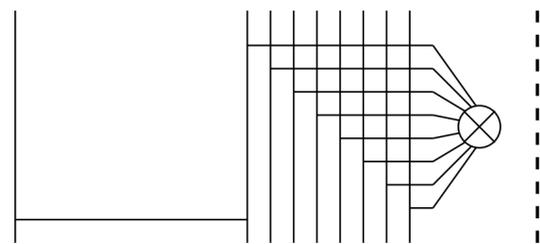
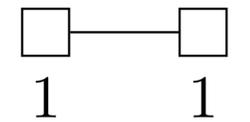
physical effects in 6d

Small instanton transition: $1T \leftrightarrow 29H$



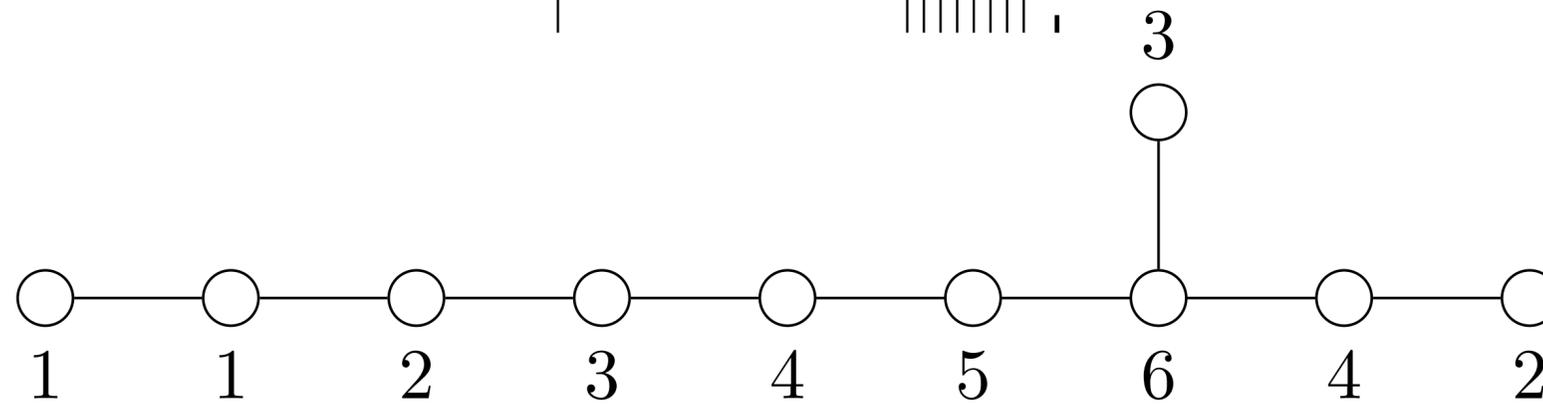
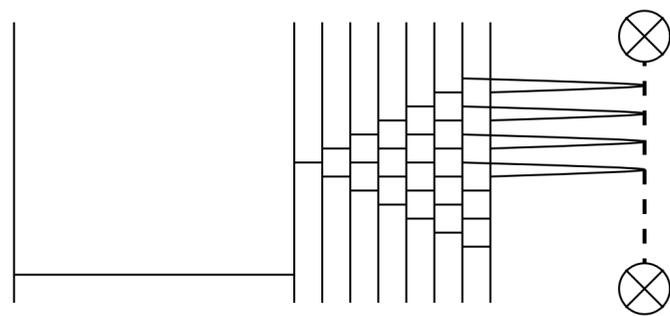
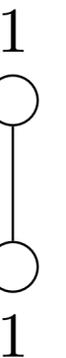
\Leftrightarrow

electric quiver:



\Leftrightarrow

magnetic quiver:



Hasse (phase) diagrams

Hasse diagrams

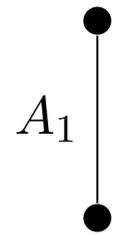
Quiver subtraction

- Recall the work of Kraft and Procesi who classified degenerations in closures of nilpotent orbits
- Minimal degenerations are of two types
- Klein singularity (ADE) — denoted by capital letters
- closure of a minimal nilpotent orbit of some algebra — denoted lower case
- This is reproduced and generalized with the Coulomb branch

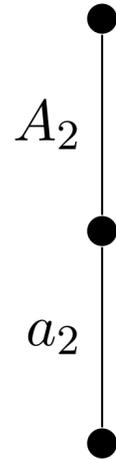
Hasse diagrams for nilpotent orbits

taken from KP

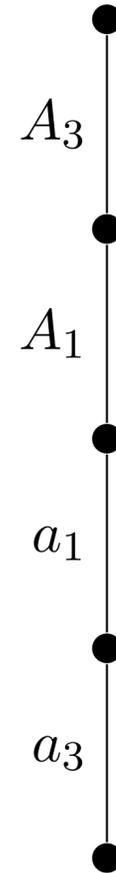
\mathfrak{sl}_2



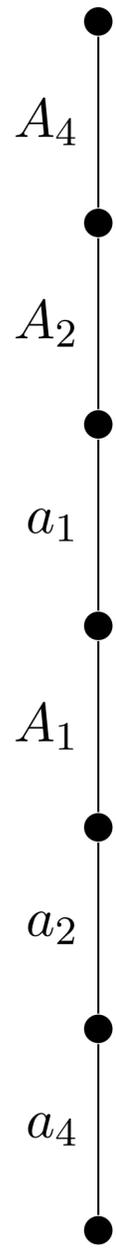
\mathfrak{sl}_3



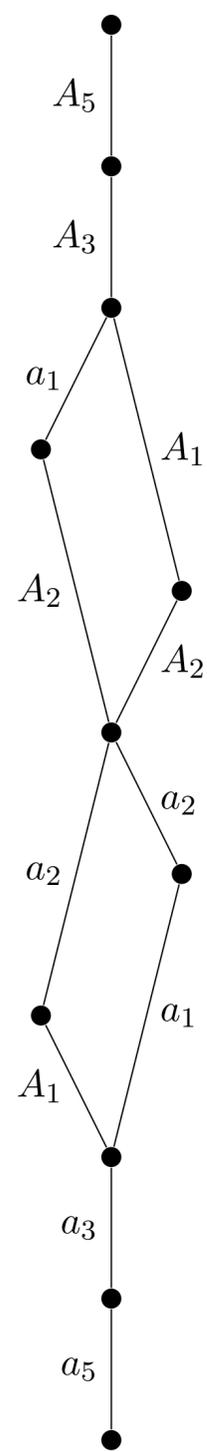
\mathfrak{sl}_4



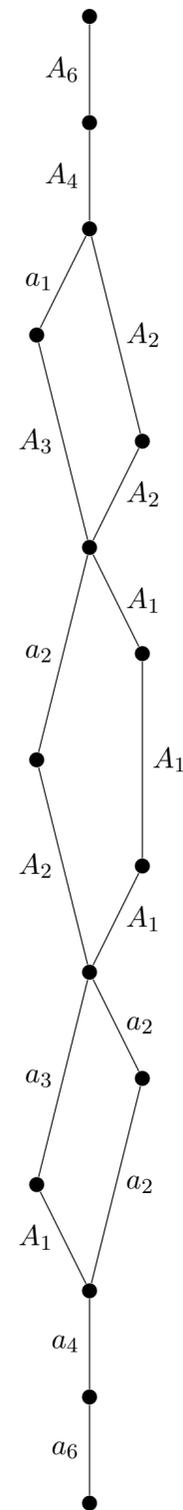
\mathfrak{sl}_5



\mathfrak{sl}_6



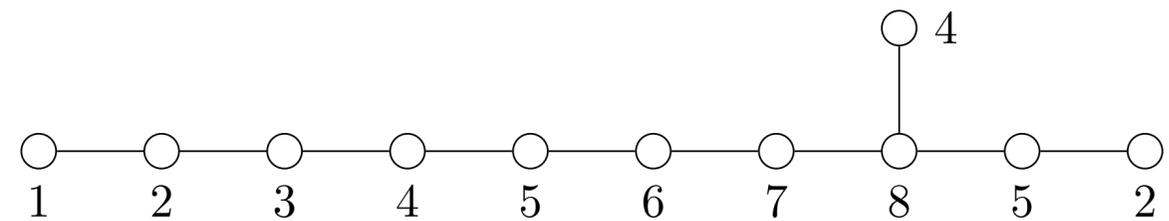
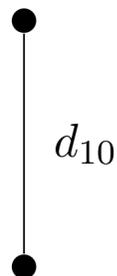
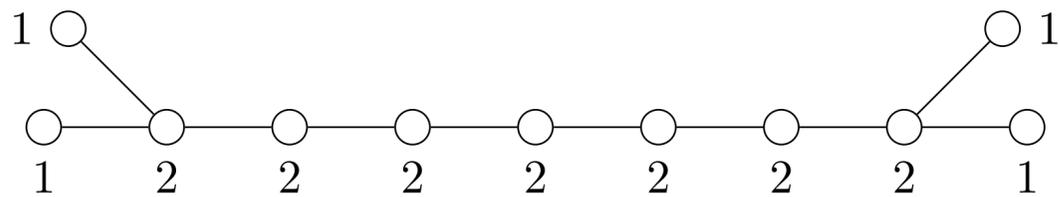
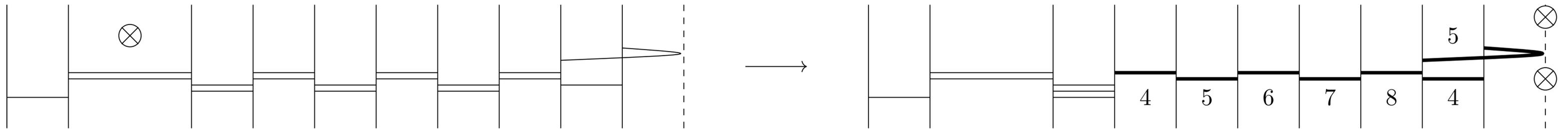
\mathfrak{sl}_7



6d – small instanton transition

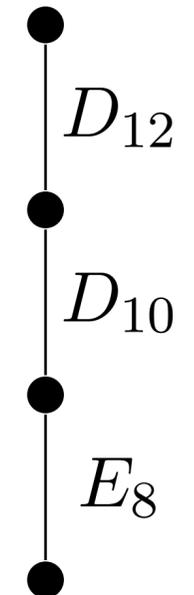
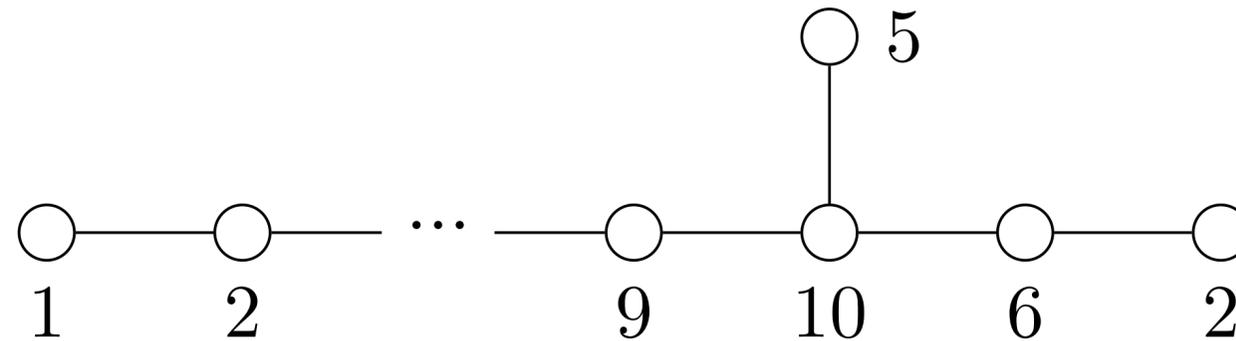
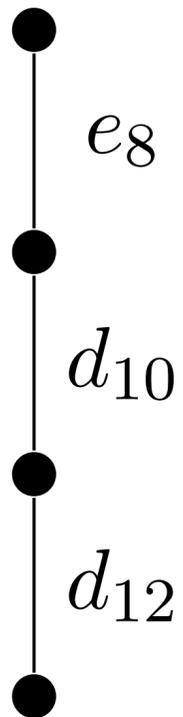
SU(2) with 10 flavors

- The Classical Higgs branch – minimal nilpotent orbit of SO(20)
- The moduli space of 1 SO(20) instanton on \mathbb{C}^2



E_8 family $k = 2$ – shows up in 5d, 6d

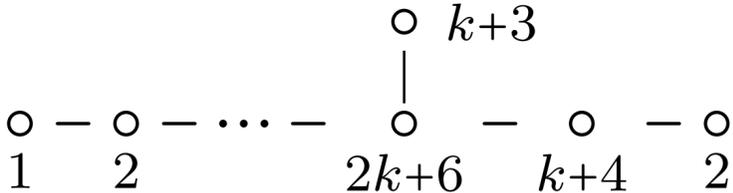
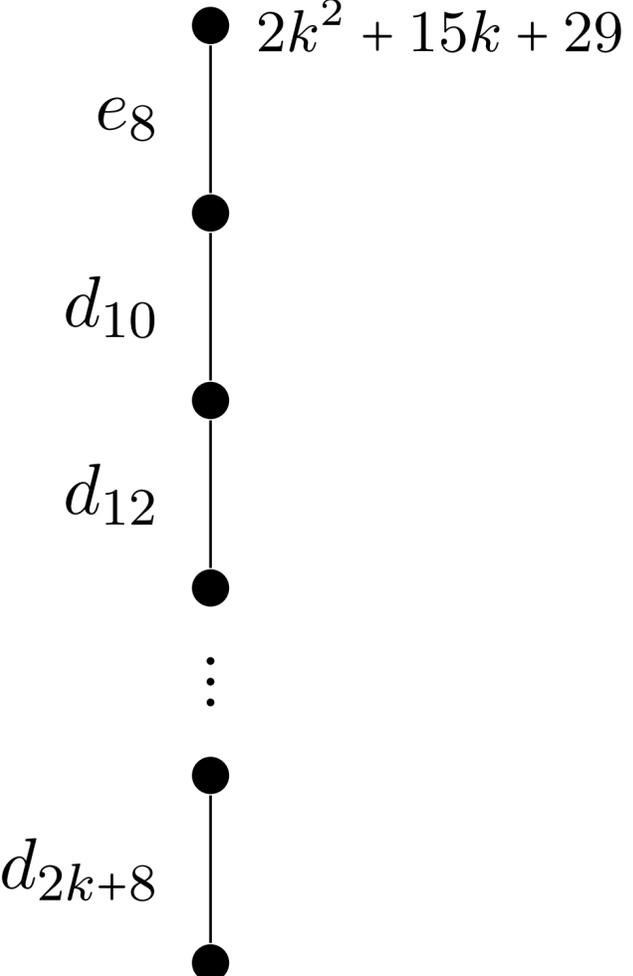
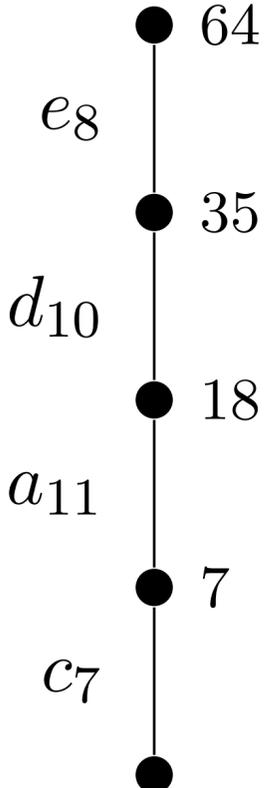
Structure of symplectic leaves – left: Coulomb branch; right: Higgs branch



$$\bullet \quad PE \left[\sum_{i=1}^{k+3} \mu_{2i} t^{2i} + t^4 + \mu_{2k+8} (t^{k+2} + t^{k+4}) \right]_{SO(4k+16)}$$

small instanton transition

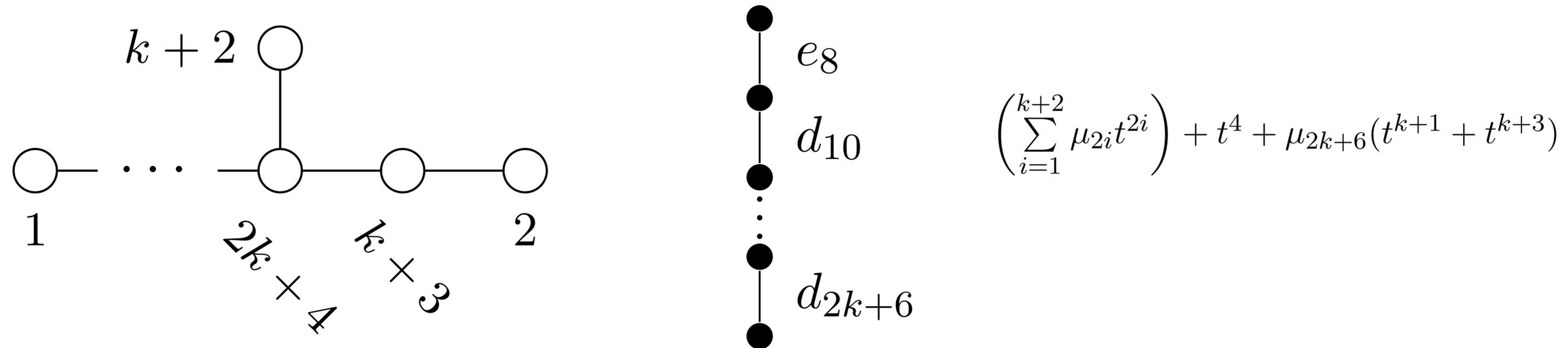
finite - infinite coupling

6d SCFT	$\mathrm{Sp}(k)$ with $N = 4k + 16$ flavours	G_2 with 7 flavours
Magnetic quiver		Not known
Hasse diagram		

6d — small instanton transition

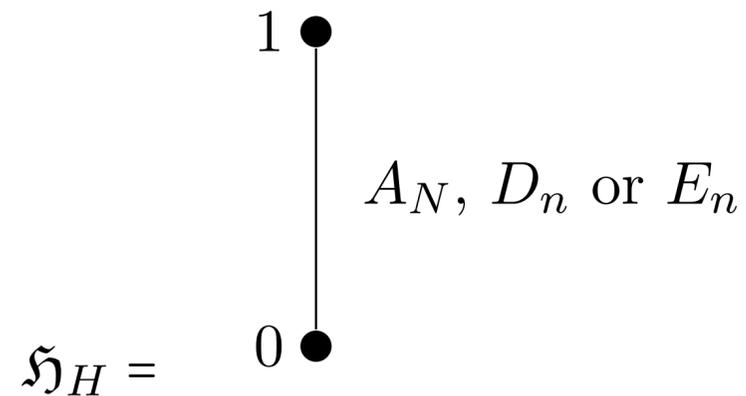
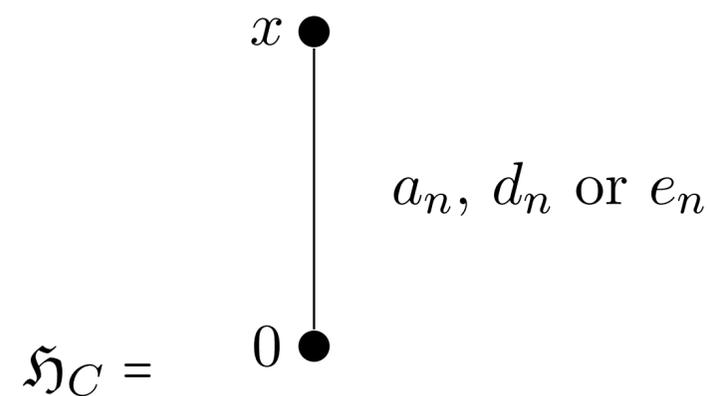
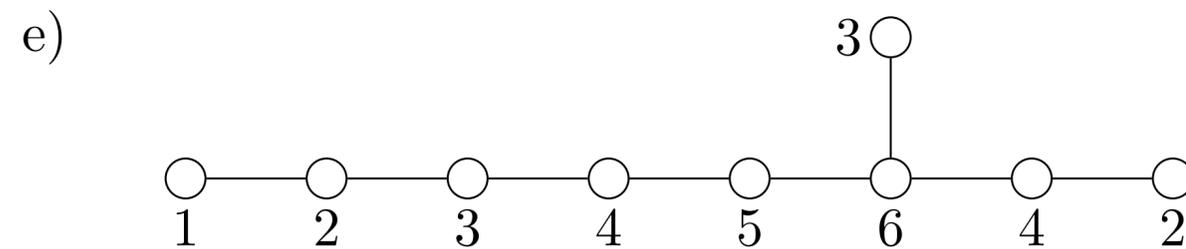
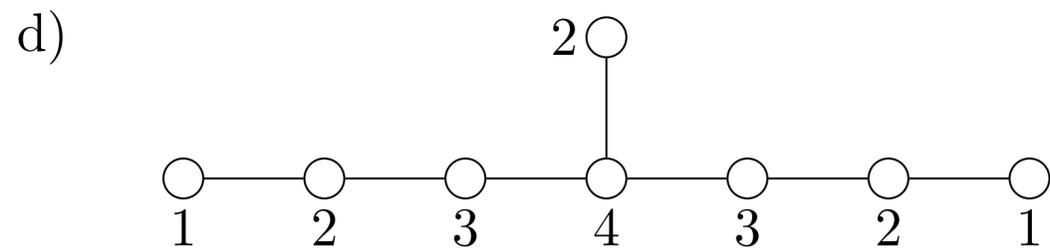
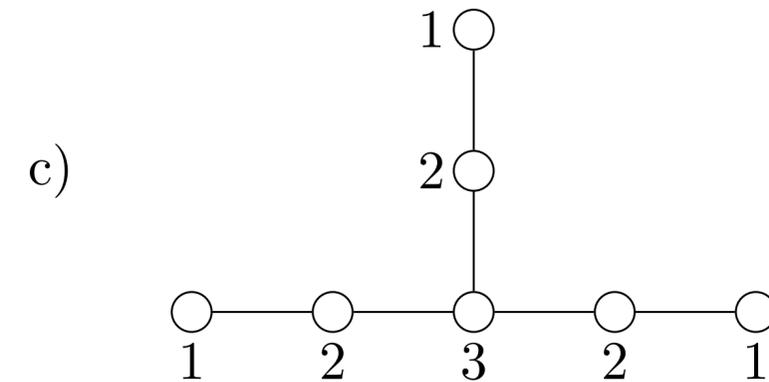
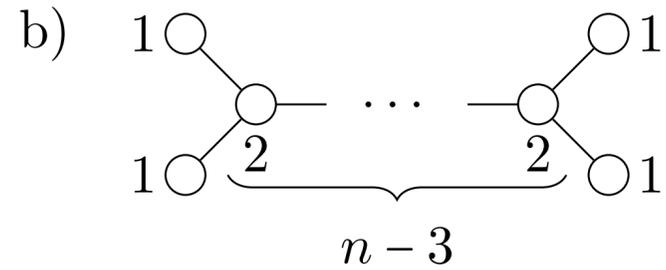
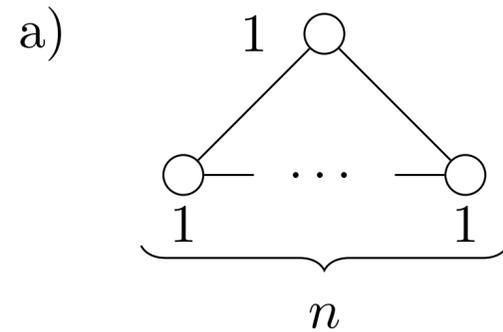
a family of $Sp(k-1)$ theories with $2k+6$ flavors

- The Classical Higgs branch — a nilpotent orbit of $SO(4k+12)$
- The Higgs branch at infinite coupling with magnetic quiver and Hasse diagram. The HWG is PE of the polynomial below



Basic Hasse diagrams - affine ADE quivers

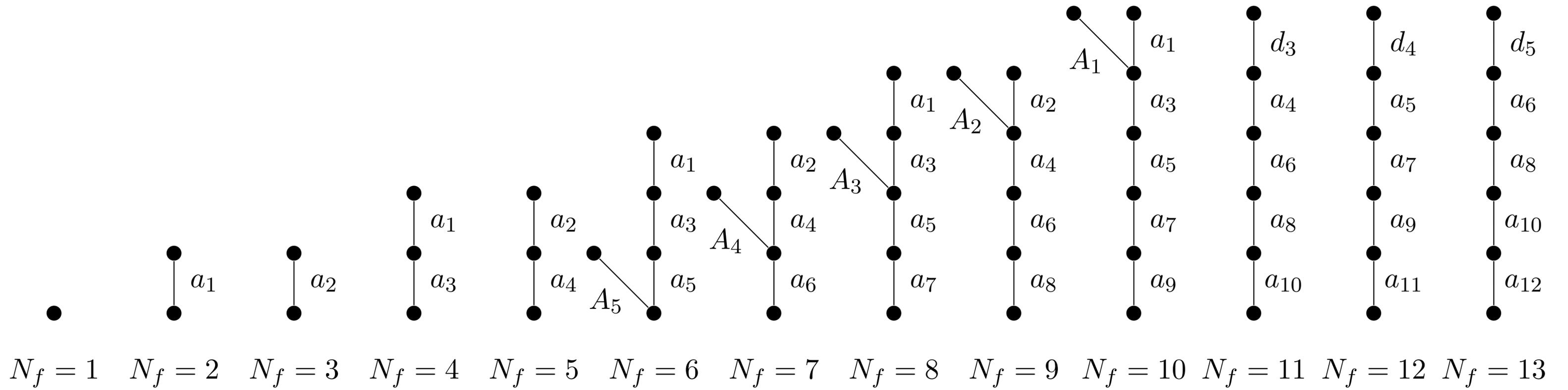
2 symplectic leaves, minimal slices



$$x = \begin{cases} n & \text{for } a_n \\ 2n - 3 & \text{for } d_n \\ 11 & \text{for } e_6 \\ 17 & \text{for } e_7 \\ 29 & \text{for } e_8 \end{cases}$$

4d $\mathcal{N} = 2$ SU(6) with fundamental matter

union of 2 cones



Minimal transverse slices

Symplectic singularities

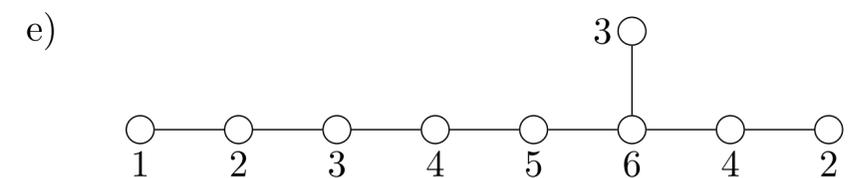
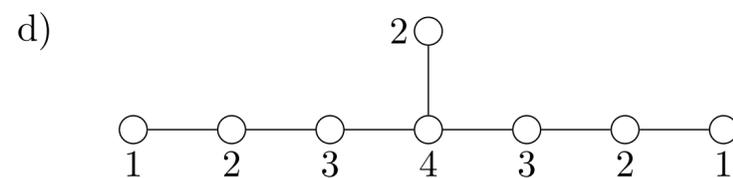
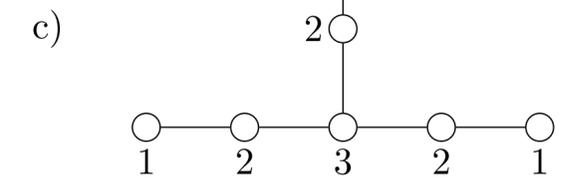
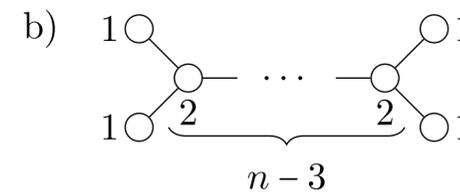
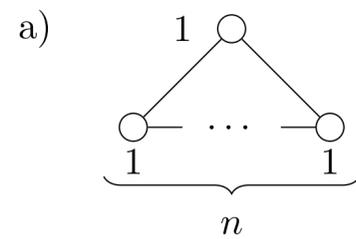
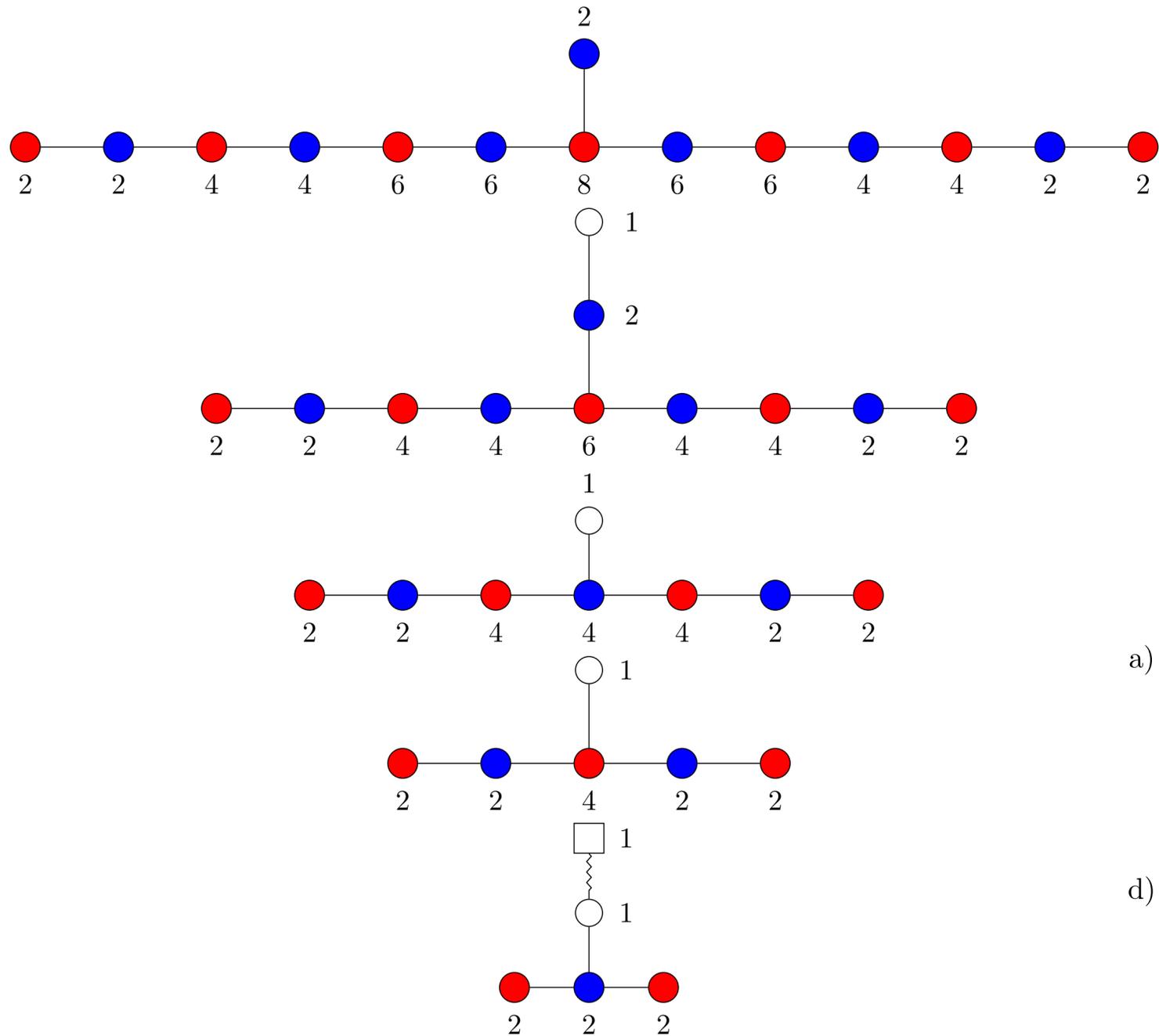
- A list of minimal transitions

Slice	Framed quiver	Unframed quiver
a_n		
b_n		
c_n		
d_n		
e_6		
e_7		
e_8		

Slice	Framed quiver	Unframed quiver
f_4		
g_2		
ac_n		
ag_2		
cg_2		
$h_{n,k}$		
$\bar{h}_{n,k}$		
A_n		

Orthogonal Symplectic

closures of minimal nilpotent orbits of type E_n for $n = 8, 7, 6, 5, 4$



**For a future of fruitful
discussions!**