# Last passage percolation in a strip 

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At the crossroads of physics and mathematics:
The joy of integrable combinatorics
in honor of Philippe Di Francesco's 60th birthday

## Symmetric functions

The Schur functions

$$
s_{\lambda}(x)=\frac{\operatorname{det}\left(x_{i}^{\lambda_{j}+n-j}\right)_{i, j=1}^{n}}{\operatorname{det}\left(x_{i}^{n-j}\right)_{i, j=1}^{n}}
$$

where $\lambda=\left(\lambda_{1} \geqslant \ldots \geqslant \lambda_{n} \geqslant 0\right)$ and $x=\left\{x_{1}, \ldots, x_{n}\right\}$ are symmetric polynomials satisfying the Cauchy summation identity

$$
\sum_{\lambda} s_{\lambda}\left(a_{1}, \ldots, a_{n}\right) s_{\lambda}\left(b_{1}, \ldots, b_{m}\right)=\prod_{i=1}^{n} \prod_{j=1}^{m} \frac{1}{1-a_{i} b_{j}}=: \Pi(a ; b) .
$$

Schur functions satisfy a branching rule that allows to expand them in monomials

$$
s_{\lambda}(x)=\sum_{\emptyset \prec \lambda^{1} \prec \ldots \prec \lambda^{n}=\lambda} \prod_{i=1}^{n} x_{i}^{\left|\lambda^{i}\right|-\left|\lambda^{i-1}\right|}
$$

where $\emptyset \prec \lambda^{1} \prec \cdots \prec \lambda^{n}=\lambda$ is a sequence of interlaced partitions. We write $\mu \prec \lambda$ for $\lambda_{1} \geqslant \mu_{1} \geqslant \lambda_{2} \geqslant \mu_{2} \geqslant \ldots$

## RSK correspondence

Expanding in monomials each side of the Cauchy identity

$$
\sum_{\lambda} s_{\lambda}\left(a_{1}, \ldots, a_{n}\right) s_{\lambda}\left(b_{1}, \ldots, b_{m}\right)=\prod_{i=1}^{n} \prod_{j=1}^{m} \frac{1}{1-a_{i} b_{j}}
$$

the two sides acan be matched using Robinson-Schensted-Knuth correspondence, a bijection

$$
\left\{\begin{array}{c}
\emptyset \prec \lambda^{1} \prec \cdots \prec \lambda^{n}=\lambda \\
\emptyset \prec \mu^{1} \prec \cdots \prec \mu^{m}=\lambda
\end{array}\right\} \quad \longleftrightarrow \quad W=\left(w_{i, j}\right) \in \mathbb{N}^{n \times m}
$$

Greene's theorem implies that $\lambda_{1}=G(n, m)$ where



$$
\begin{equation*}
G(n, m)=\max _{\text {paths }(1,1) \rightarrow(n, m)}\left\{\sum_{(i, j) \in \text { path }} w_{i, j}\right\} \tag{1,1}
\end{equation*}
$$

## Schur measure

Assume that when $w_{i, j} \sim \operatorname{Geom}\left(a_{i} b_{j}\right)$ are independent (we say that $w \sim \operatorname{Geom}(q)$ if $\left.\mathbb{P}(w=k)=(1-q) q^{k}\right)$. Then,

$$
\mathbb{P}(G(n, m) \leqslant r)=\frac{1}{\Pi(a ; b)} \sum_{\lambda: \lambda_{1} \leqslant r} s_{\lambda}(a) s_{\lambda}(b)
$$

In other terms, $G(n, m)$ has the same law as $\lambda_{1}$ when $\lambda$ is a random partition sampled according to the Schur measure

$$
\mathbb{P}(\lambda)=\frac{1}{\Pi(a ; b)} s_{\lambda}(a) s_{\lambda}(b) .
$$

## Asymptotics

[Johansson, 2001] proved that after appropriate rescaling, $G(n, m)$ fluctuates according to the Tracy-Widom GUE distribution (governing fluctuations of the largest eigenvalue of Hermitian random matrices of large size).

## A variant



Assume that the weight matrix is symmetric:

$$
w_{i, j}=w_{j, i}
$$

$$
w_{i, j} \sim \operatorname{Geom}\left(a_{i} a_{j}\right) \text { for } i>j
$$

$$
w_{i, i} \sim \operatorname{Geom}\left(c a_{i}\right)
$$

$$
G^{\square}(n, m)=\max _{\text {paths }(1,1) \rightarrow(n, m)} \sum_{(i, j) \in \text { path }} w_{i, j}
$$

[Baik-Rains 2003] proved that

$$
\mathbb{P}\left(G^{\square}(n, n) \leqslant r\right)=\frac{1}{\Pi(a, c) \Phi(a)} \sum_{\lambda: \lambda_{1} \leqslant r} c^{\lambda_{1}-\lambda_{2}+\lambda_{3}-\lambda_{4}+\ldots} s_{\lambda}(a)
$$

where

$$
\Phi(x)=\prod_{i<j} \frac{1}{1-x_{i} x_{j}}=\sum_{\lambda^{\prime} \text { even }} s_{\lambda}(x) .
$$

## Asymptotics

As $n \rightarrow \infty, G^{\square}(n, n)$ fluctuates according to the Tracy-Widom GSE or GOE distributions (depending if $c=1$ or $c<1$ ).

## Other variants

[Baik-Rains 2003] also computed $\mathbb{P}\left(G^{\bullet}(n, n) \leqslant r\right)$ in terms of Schur functions for other symmetry types • :


The problem is however open for symmetries such as


Last passage percolation with walls
Imposing two symmetry axis in the diagonal direction is equivalent to assuming that paths are constrained to remain between two walls.

Last passage percolation in a strip
Let $a_{1}, \ldots, a_{N} \in(0,1), c_{1}, c_{2}>0$.

$$
w_{i, j} \sim \operatorname{Geom}\left(a_{i} a_{j}\right) \text { for } j<i<j+N
$$ (indices modulo $N$ )

$c_{2} \quad w_{i, i} \sim \operatorname{Geom}\left(c_{1} a_{i}\right)$

$$
w_{j+N, j} \sim \operatorname{Geom}\left(c_{2} a_{j}\right)
$$

$(0,0)$
$(i, 0) \quad(N, 0)$
We fix an initial condition $G(i, 0)=G_{0}(i)$ for some function $G_{0}$.

$$
G(n, m)=\max _{\text {paths }}^{(i, 0) \rightarrow(n, m)}\left\{G_{0}(i)+\sum_{(i, j) \in \text { path }} w_{i, j}\right\}
$$

## Open Problem

Find the asymptotic distribution of $G(n, m)$ as $n, m \rightarrow \infty$, depending on the width $N$ and the boundary parameters $c_{1}, c_{2}$.

## Conjectural phase diagram

The richest behaviour is when $N=L n^{2 / 3}$. For $a_{1}=\cdots=a_{N}=a$, it is expected that


## Stationary measure (=non-equilibrium steady-state)

However, we can find the asymptotic distribution of

$$
G_{t}(i)=G(t+i, t)-G(t, t)
$$

as $t$ goes to infinity. $G_{t}$ is a Markov process on $\mathbb{Z}^{N}$.

## Problem

Find the law of the initial condition $\left(G_{0}(i)\right)_{1 \leqslant i \leqslant N}$ such that for all $t$,
$\left(G_{t}(i)\right)_{1 \leqslant i \leqslant N} \stackrel{(d)}{=}\left(G_{0}(i)\right)_{1 \leqslant i \leqslant N}$

- For models such as Asymmetric Simple Exclusion Process (ASEP), the standard method is the matrix product ansatz [Derrida-Evans-Hakim-Pasquier 1993].
- We will illustrate another approach based on symmetric functions, taking the example of Last Passage Percolation and Schur functions.
- The method works as well for other models: Log-gamma polymer and KPZ equation (Whittaker functions), stochastic six vertex model (Hall-Littlewood polynomials), other examples.


## Stationary measure

Assume for simplicity that $a_{1}=\cdots=a_{N}=a$.
For $R=(R(j))_{1 \leqslant j \leqslant N}$, let

$$
\mathbb{P}_{q}^{\mathrm{RW}}(R)=\prod_{j=1}^{N}(1-q) q^{R(j)-R(j-1)}
$$

be the probability that $R$ is a random walk with $\operatorname{Geom}(q)$ increments.
Define a Pitman-type operation

$$
R_{1} \star R_{2}(k)=\min _{1 \leqslant j \leqslant k}\left\{R_{1}(j-1)+R_{2}(k)-R_{2}(j)\right\} .
$$

Consider the probability measure

$$
\mathbb{P}_{a, c_{1}, c_{2}}\left(R_{1}, R_{2}\right)=\frac{1}{Z}\left(c_{1} c_{2}\right)^{-R_{1} \star R_{2}(N)} \times \mathbb{P}_{a c_{1}}^{\mathrm{RW}}\left(R_{1}\right) \times \mathbb{P}_{a c_{2}}^{\mathrm{RW}}\left(R_{2}\right)
$$

## Theorem ([B.-Corwin-Yang 2023])

For any parameters a, $c_{1}, c_{2}$, the marginal law of $R_{1}$ under $\mathbb{P}_{a, c_{1}, c_{2}}$ is the unique stationary measure of the Markov process $G_{t}$.

## A variant of the Schur process

The random walks $R_{1}, R_{2}$ are related to a sequence of partitions signatures:

$$
R_{1}(j)=\lambda_{1}^{(j)}-\lambda_{1}^{(0)}, \quad R_{2}(j)=\lambda_{2}^{(j)}-\lambda_{2}^{(0)}
$$

where $\boldsymbol{\lambda}=\lambda^{(0)} \prec \lambda^{(1)} \prec \cdots \prec \lambda^{(N)}$ is a sequence of interlaced signatures $\lambda^{(j)}=\left(\lambda_{1}^{(j)} \geqslant \lambda_{2}^{(j)}\right) \in \mathbb{Z}^{2}$ distributed as

$$
\mathbb{P}(\boldsymbol{\lambda})=\frac{1}{Z_{a, c_{1}, c_{2}}(N)} c_{1}^{\lambda_{1}^{(0)}-\lambda_{2}^{(0)}} c_{2}^{\lambda_{1}^{(N)}-\lambda_{2}^{(N)}} \prod_{j=1}^{N} s_{\lambda}(j) / \lambda^{(j-1)}\left(a_{i}\right)
$$

and $s_{\lambda / \mu}$ denote skew Schur functions

$$
s_{\lambda / \mu}(x)=\mathbb{1}_{\lambda_{1} \geqslant \mu_{1} \geqslant \lambda_{2} \geqslant \mu_{2} x^{\lambda_{1}+\lambda_{2}-\mu_{1}-\mu_{2}} .} .
$$

The construction is similar to the free boundary Schur process [Betea-Bouttier-Nejjar-Vuletic 2017] except that the $\lambda^{(j)}$ are no longer integer partitions:

- parts can be negative,
- all signatures have length 2.
- the measure $\mathbb{P}$ is infinite (but becomes a probability measure if we fix $\lambda_{2}^{(0)}=0$ ).

Properties of Schur functions yield explicit formulas:

$$
\mathbb{E}\left[t^{2 R_{1}(N)}\right]=\frac{1}{Z_{a, c_{1}, c_{2}}(N)} \oint \frac{d z}{2 \dot{i} \pi z}\left|\frac{1-z^{2}}{(1-t a z)^{N}\left(1-z c_{1} / t\right)\left(1-z c_{2} t\right)}\right|^{2}
$$

More generally, for $0=x_{0}<\cdots<x_{k}=N$, there is a simple formula for

$$
\mathbb{E}\left[\prod_{i=1}^{k} t_{i}^{2\left(R_{1}\left(x_{i}\right)-R_{1}\left(x_{i-1}\right)\right)}\right]
$$

In particular, one can deduce that starting from the stationary initial data, i.e. $G_{0} \stackrel{(d)}{=} R_{1}$, we have

$$
\mathbb{E}[G(n, n)]=n \times v\left(c_{1}, c_{2}, N\right)
$$

where

$$
v\left(c_{1}, c_{2}, N\right)=\frac{\left(1-a^{2}\right) Z_{a, c_{1}, c_{2}}(N+1)-Z_{a, c_{1}, c_{2}}(N)}{Z_{a, c_{1}, c_{2}}(N)}
$$

with

$$
Z_{a, c_{1}, c_{2}}(N)=\oint \frac{d z}{2 \mathbf{i} \pi z}\left|\frac{1-z^{2}}{(1-a z)^{N}\left(1-z c_{1}\right)\left(1-z c_{2}\right)}\right|^{2}
$$

## More general two-layer Schur process



Vertices on the path are decorated by signatures $\lambda=\left(\lambda_{1} \geqslant \lambda_{2}\right) \in \mathbb{Z}^{2}$. We define a probability measure on $\boldsymbol{\lambda}$ by taking the product of Boltzmann weights

$$
\begin{aligned}
& \mathrm{wt}\left(\mathrm{a}^{\lambda}\right)=\mathrm{wt}\left({ }^{\lambda}{ }_{\mu}^{\mathrm{a}}\right)=s_{\lambda / \mu}(a) \\
& \mathbb{P}(\boldsymbol{\lambda})=\frac{1}{Z_{a, c_{1}, c_{2}}(N)} c_{1}^{\lambda_{1}^{(0)}-\lambda_{2}^{(0)}} c_{2}^{\lambda_{1}^{(N)}-\lambda_{2}^{(N)} \prod_{\text {edges } e} \mathrm{wt}(e)}
\end{aligned}
$$

When $c_{1} c_{2}<1$, this is a well-defined probability measure.

## Dynamics on the two-layer Schur process



$$
\boldsymbol{\lambda}=\left(\lambda^{(0)}, \lambda^{(1)}, \ldots, \lambda^{(N)}\right)
$$



We construct dynamics on $\boldsymbol{\lambda}$, inspired by [Borodin-Ferrari 2008], such that when the path evolves by the elementary moves


1 The two-layer Schur process is mapped to a two layer Schur process;
2 the $\lambda_{1}$ marginal of the dynamics corresponds to the recurrence of geometric LPP.
After averaging over $\lambda_{1}^{(0)}, \lambda_{2}^{(0)}$, the law of $R_{1}(j)=\lambda_{1}^{(j)}-\lambda_{1}^{(0)}$ simplifies and can be analytically continued to all $c_{1}, c_{2}$.

## Connection to the Matrix Product Ansatz

The law of the stationary measure $R_{1}$ can be written, in terms of increments $\Delta(j)=R_{1}(j)-R_{1}(j-1) \in \mathbb{N}$ as

$$
\mathbb{P}\left(\bigcap_{j=1}^{N}\left\{\Delta(j)=x_{j}\right\}\right)=\frac{1}{Z(N)} \mathbf{w}^{t}\left(\prod_{j=1}^{N} M\left(x_{j} ; \cdot, \cdot\right)\right) \mathbf{v}
$$

where

$$
M\left(x ; n, n^{\prime}\right)=s_{\lambda^{\prime} / \lambda}(a) \text { where } \begin{cases}n & =\lambda_{1}-\lambda_{2} \\ n^{\prime} & =\lambda_{1}^{\prime}-\lambda_{2}^{\prime} \\ x & =\lambda_{1}^{\prime}-\lambda_{1}^{\prime}\end{cases}
$$

with $\mathbf{w}^{t}=\left(1, c_{1}, c_{1}^{2}, \ldots\right)$ and $\mathbf{v}^{t}=\left(1, c_{2}, c_{2}^{2}, \ldots\right)$.
In other terms, the stationary measure is the Matrix Product State


Hence, another way to interpret the talk is that skew Schur functions provide a representation of the MPA relations for Last Passage Percolation.

## Conclusion

## Summary

The two-layer Schur processes allows to describe the stationary measure of LPP in a strip in terms of reweighted random walks.

The stationary measure is not a priori a Gibbs measure, but becomes so on some enlarged state space.

The method becomes more interesting when applied to more complicated models (log-gamma polymer, KPZ equation) related to Whittaker functions.

## Outlook

- The method applies to other families of symmetric functions [B.-Corwin, in progress]
- Go beyond the stationary measure.


## Thank you for your attention!

