

Philippe's Fest !



*where is ~~Charlie~~ ?
J.-B.*

Playing with random matrices and free probability in noisy many-body systems

(A small — possibly entertaining — exercise)

Denis Bernard (CNRS & LPENS, Paris)

(work done in collaboration with Ludwig Hruza)

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(Standard) random matrices :

with impact on modeling many physical phenomena (say random surfaces, chaotic systems, etc)

– Basic examples :

« Wigner's matrices » $\left\{ \begin{array}{l} M = (M_{ij}) \text{ hermitian, } N \times N \\ \text{with Gaussian entries } \mathbb{E}[M_{ij}M_{ji}] = N^{-1} \end{array} \right.$

« Haar orbits » $\left\{ \begin{array}{l} M = UDU^\dagger \text{ with } U \text{ unitary Haar distributed, } D \text{ diagonal} \\ \text{with HCIZ integral as} \\ \text{cumulant generating function } \mathbb{E}[e^{tr(AM)}] = \int dU e^{tr(AUDU^\dagger)} \end{array} \right.$

- Properties :
- $U(N)$, hence $U(1)^N$ invariance
 - Unstructured matrices (invariance under permutation, in law)
 - Expectation values are « topological »

→ Look at structured matrices ?...

$$M_{i_1 i_2} |M_{i_2 i_3}|^2 M_{i_2 i_3} M_{i_3 i_1} M_{i_1 i_4} M_{i_4 i_1} M_{i_5 i_5} = \text{Diagram} \cdot$$

Some computational experiments...

→ Originates from noisy many-body systems, in particular **QSSEP** (more later...)

– U(1) invariance (by choice): $M_{jk} \equiv_{\text{in law}} e^{i\theta_j} M_{jk} e^{-i\theta_k}$ (→ charge conservation)

$\mathbb{E}[M_{ij}] = 0$, $(i \neq j)$ ← non-trivial graphs are Eulerian

$\mathbb{E}[M_{ii}] = g_1(x)$, $x = i/N$ ← echo of the « structure »

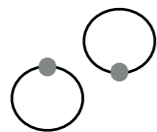
– Scaling of the loop/simple cumulants (→ limiting spectrum and large deviation function)



$\mathbb{E}[M_{ij}M_{ji}] = N^{-1} g_2(x, y)$, $x = i/N, y = j/N, (i \neq j)$

← with smaller 1/N dependence for bigger loops (moments)

– Factorisation of disconnected graphs (at leading order)



$\mathbb{E}[M_{ii}M_{jj}] = g_1(x)g_1(y)$, $(i \neq j)$ (up to 1/N correction)

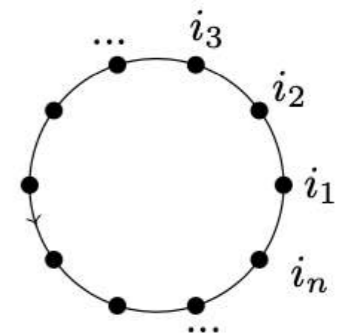
Ensemble of 'structured' random matrices :

– Ensemble of large $N \times N$ matrices M with four defining properties :

- U(1) invariance : $M_{jk} \equiv_{\text{in law}} e^{i\theta_j} M_{jk} e^{-i\theta_k}$

- Scaling of the loop expectation values : ($x_k = i_k/N$, distincts)

$$\mathbb{E}[M_{i_1 i_n} \cdots M_{i_3 i_2} M_{i_2 i_1}] \sim N^{1-n}$$



- Factorisation of product of disconnected graphs (at leading order) :

$$\mathbb{E}[M_{j_1 j_n} \cdots M_{j_2 j_1} \cdot M_{i_1 i_p} \cdots M_{i_2 i_1}] = \mathbb{E}[M_{j_1 j_n} \cdots M_{j_2 j_1}] \mathbb{E}[M_{i_1 i_p} \cdots M_{i_2 i_1}]$$

- Continuity of the loop cumulants at coincident points (at leading order) :

$$N^{n-1} \mathbb{C}[M_{i_1 i_2} M_{i_2 i_3} \cdots M_{i_n i_1}] \text{ continuous at } x_{i_p} = x_{i_q}$$

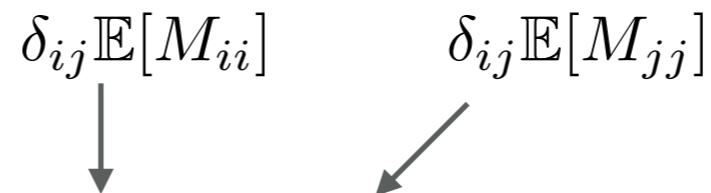
—> Allow to compute the expectation of any graph (at leading order)

The role of non-crossing partitions (i) :

Compute the moments / loop expectation values
(with possible coincident points) via the cumulant expansion

Recall the moment-cumulant formula : $\mathbb{E}[X_1 \cdots X_n] = \sum_{\pi \in P(n)} \mathbb{E}_{\pi}[X_1 \cdots X_n]^c$

– N=1 : $\mathbb{E}[M_{ii}] = g_1(x)$



– N=2 : $\mathbb{E}[M_{ij}M_{ji}] = \mathbb{E}[M_{ij}M_{ji}]^c + \mathbb{E}[M_{ij}]^c \mathbb{E}[M_{ji}]^c$



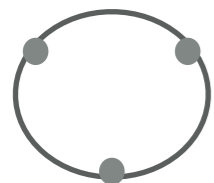
$= N^{-1} [g_2(x, y) + \delta(x - y)g_1(x)g_1(y)]$ (with $\delta_{ij} \rightsquigarrow N^{-1}\delta(x - y)$)

(or $\mathbb{E}[\text{Tr}(h_1 M h_2 M \cdots h_n M)]$, h_k diag)

– N=3 : (idem...)



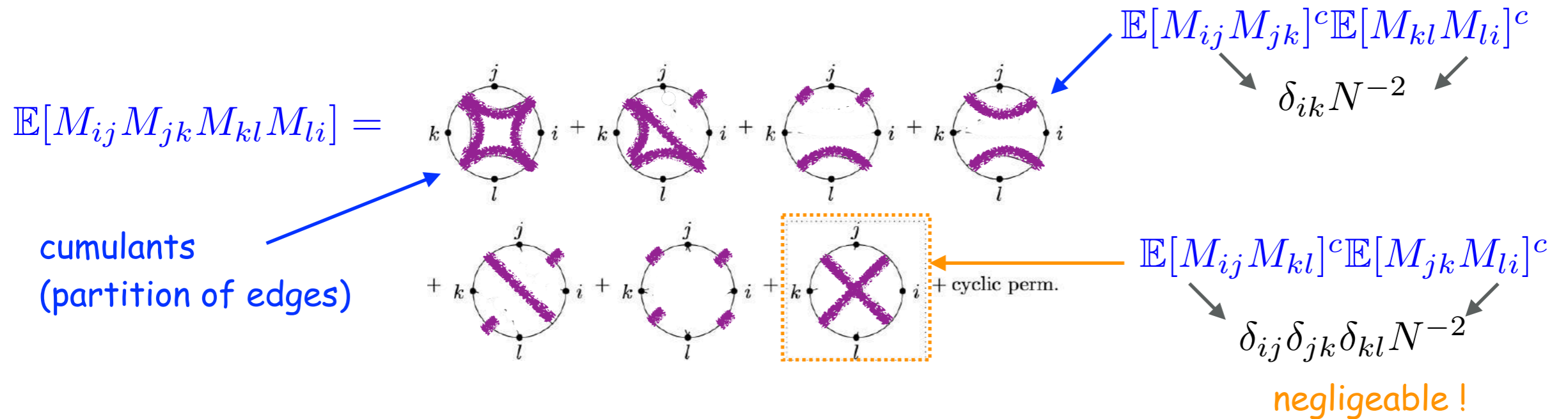
$\mathbb{E}[M_{ij}M_{jk}M_{ki}] = \mathbb{E}[M_{ij}M_{jm}M_{ki}]^c + \mathbb{E}[M_{ij}]^c \mathbb{E}[M_{jk}M_{ki}]^c + \circlearrowleft + \mathbb{E}[M_{ij}]^c \mathbb{E}[M_{jk}]^c \mathbb{E}[M_{ki}]^c$



$= N^{-2} [g_3(x_1, x_2) + \delta(x_1 - x_2)g_1(x_1)g_2(x_2, x_3) + \circlearrowleft + \delta(x_1 - x_2)\delta(x_2 - x_3)g_1(x_1)g_1(x_2)g_1(x_3)]$

The role of non-crossing partitions (ii) :

– N=4 : the role of non-crossing partitions...



$$= N^{-3} [g_4(x_1, x_2, x_3, x_4) + \delta(x_1, x_2) g_1(x_1) g_3(x_2, x_3, x_4) + \circlearrowleft$$

$$+ \delta(x_1, x_3) g_2(x_1, x_2) g_2(x_3, x_4) + \circlearrowleft$$

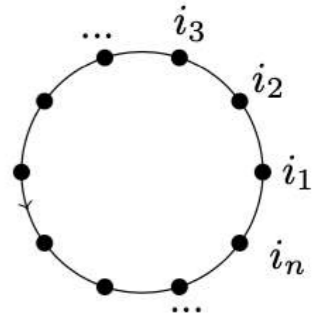
$$+ \dots \text{ "only non-crossing partitions"}]$$

Only the non-crossing partitions contribute. \implies Free probability !

Structured RMT & Loop expectations values:

—→ Dominant role of the loop expectation values.

$$g_n(x_1, x_2, \dots, x_n) := \lim_{N \rightarrow \infty} N^{n-1} \mathbb{E}[M_{i_1 i_2} M_{i_2, i_3} \dots M_{i_n i_1}]^c$$



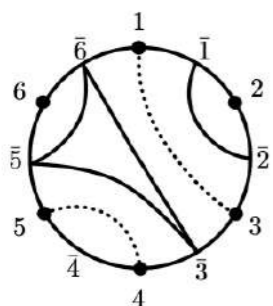
($x_k = i_k/N$, distincts)

- They code for all structures and correlations (at leading order)
- Scaling with N ensures the existence of a large deviation function (formal series)
- We call them the 'local free cumulants'
- They have a natural interpretation as free cumulants in operator valued free probability (see below)

Loop expectations values & non-crossing partition :

– For loop of any size, we have (by similar arguments) :

$$\lim_{N \rightarrow \infty} N^{n-1} \mathbb{E}[M_{i_1 i_2} M_{i_2 i_3} \cdots M_{i_n i_1}] = \sum_{\pi \in NC(n)} g_{\pi}(\vec{x}) \delta_{\pi^*}(\vec{x}) .$$



$$\pi = \{\{\bar{1}, \bar{2}\}, \{\bar{3}, \bar{5}, \bar{6}\}, \{\bar{4}\}\}$$

$$\pi^* = \{\{1, 3\}, \{2\}, \{4, 5\}, \{6\}\}$$

non-crossing partition
(on edges)

its Kreweras dual
(on points)

or
$$N^{-1} \mathbb{E}[\text{Tr}(M^n)] = \sum_{\pi \in NC(n)} \int g_{\pi^*}(\vec{x}) \delta_{\pi}(\vec{x}) d\vec{x}$$

– Operator valued probability alias « conditioned expectations over the diagonal »

e.g.
$$\mathbb{E}[\langle x | M h_1 M \cdots h_n M | x \rangle], \quad h_k \text{ diag}$$

Non-linear transformations

[Thanks to R. Speicher]

The 'axioms' (i)-(iv) are closed under some non-linear transformations

→ A way of constructing new ensembles from simple ones.

Applications (say in ML) ??...

– Polynomial transformations

If M satisfies the axioms, so does $P(M) = \sum_{n>0} a_n M^n$.

(with an explicit expression for the new local free cumulants)

– Point-wise transformations

Given M , define Y with entries $Y_{ij} = M_{ij} f_{ij}(N|M_{ij}|^2)$, with non-linear functions f_{ij} which may be different for each entries. Then

If M satisfies the axioms, so does Y .

(with an explicit expression for the new local free cumulants)

Spectrum of sub-blocks (i)

$$M = \begin{pmatrix} * & * & * & * & * & * & * \\ * & * & * & * & * & * & * \\ * & * & * & * & * & * & * \\ * & * & * & * & * & * & * \\ * & * & * & * & * & * & * \\ * & * & * & * & * & * & * \\ * & * & * & * & * & * & * \end{pmatrix} \rightsquigarrow \hat{M}_I = \begin{pmatrix} 0 & \overline{0} & \overline{0} & \overline{0} & 0 & \overline{0} & 0 \\ 0 & * & * & * & 0 & * & 0 \\ 0 & * & * & * & 0 & * & 0 \\ 0 & * & * & * & 0 & * & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & * & * & * & 0 & * & 0 \\ 0 & \overline{0} & \overline{0} & \overline{0} & 0 & \overline{0} & 0 \end{pmatrix}$$

→ What is the spectral distribution of any (collection of) sub-blocks ?

– For the usual unstructured ensemble is simple or known:

- sub-blocks of Wigner matrices → again Wigner semi-circle law.
- sub-blocks of Haar rotated matrices UDU^*
 - free dilatation of the spectral measure of D

Spectrum of sub-blocks (ii)

→ (Using the moment methods (from which the spectral measure can be deduced...))

Let $F[h](z) := \mathbb{E} N^{-1} \text{Tr} \log(z - M_h)$ with $M_h = h^{\frac{1}{2}} M h^{\frac{1}{2}}$, e.g. $h(x) = \mathbb{I}_{x \in I}$

[D.B.-L. Hruza]

Theorem

$$F[h](z) = \min_{a_z, b_z} \left[\int [\log(z - h(x)b_z(x)) + a_z(x)b_z(x)] dx - F_0[a_z] \right]$$

« generating function of single loop expectation values »

$$F_0[a] = \sum_{n \geq 1} \frac{1}{n} \int d\vec{x} g_n(\vec{x}) a(x_1) \cdots a(x_n)$$

Proof Via nested structured of non-crossing partition or via operator valued free probability

$$b_z(x) = \frac{\delta F_0[a_z]}{\delta a_z(x)} \ \& \ a_z(x) = \frac{h(x)}{z - h(x)b_z(x)}, \quad \rightsquigarrow \quad G[h](z) = \int \frac{dx}{z - h(x)b_z(x)}$$

Origin in / Relation with Q-SSEP :

– Q-SSEP = Quantum Symmetric Simple Exclusion Process : [D.B.-T.Jin, 2019]



$$dH_t = \sqrt{D} \sum_j (c_{j+1}^\dagger c_j dW_t^j + c_j^\dagger c_{j+1} d\overline{W}_t^j) \quad \left\{ \begin{array}{l} + \text{ boundary terms...} \\ + \text{ injection/extraction...} \end{array} \right.$$

– Stochastic many-body quantum system (quadratic but noisy) :

→ « Coherences » alias « matrix of two-point functions »

$$G_{ij} = \langle c_j^\dagger c_i \rangle_t = \text{Tr}(\rho_t c_j^\dagger c_i)$$

random quantum states



Large « structured »
random matrices

– Claim :

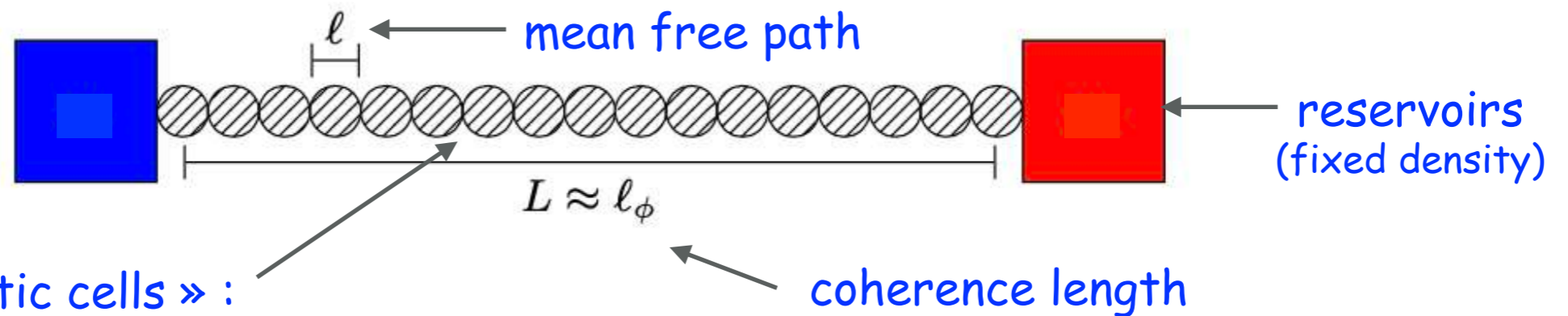
At any time (with the measure E induced by the above SDE),
and hence in particular in the infinite time steady measure :

The matrix of coherences G belongs to the above 'structured' RMT class.

→ Many (physical) consequences : { e.g. long range correlations,
volume law for entanglement mutual information, ...

Application to /Relation with mesoscopic physics :

- « Mesoscopic » (diffusive + coherent) systems out-of-equilibrium



« ballistic cells » :
small scale / fast d.o.f. → noise

- Conjecture : [D.B.-L. Hruza]

If the fast dynamics on the small scale degrees of freedom is ergodic,
If the mean dynamics follows the classical 'macroscopic fluctuation theory',
Then
The matrix of two point functions belongs to the above 'structured' RMT class.

Bon

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!!!

Q-SSEP Spectral density from free probability

- Let $a_z(x) := \mathbb{E}\langle x | \frac{h}{z - G_h} | x \rangle$ then $zh(x)^{-1} = R_0[a_z](x) + a_z(x)^{-1}$
 Let $b_z(x) := R_0[a_z](x) = \delta F_0[a_z] / \delta a_z(x)$ (local R-transform)

- For Q-SSEP:

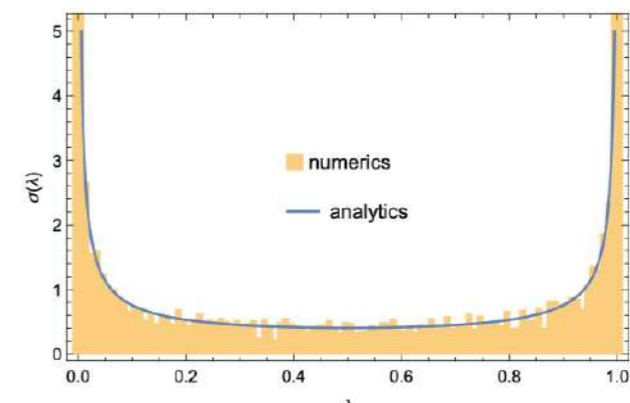
$$\begin{aligned} & [\log(z - b_z(x))]'' = 0, \quad x \in I, \\ & b_z(x)'' = 0, \quad x \notin I, \end{aligned} \quad \text{with } b_z(0) = 0, b_z(1) = 1.$$

because $R_0[a](x)$ is constructed from free cumulants w.r.t. the Lebesgue measure.

- For the complete interval (i.e. for G)

$$d\sigma_{[0,1]}(\lambda) = \frac{d\lambda}{\lambda(1-\lambda)} \frac{1}{\pi^2 + \log^2\left(\frac{1-\lambda}{\lambda}\right)}.$$

which is a Lorentzian distribution for $\nu = \log\left(\frac{1-\lambda}{\lambda}\right)$



Hidden free probability in classical SSEP :

[M. Bauer, D.B., Ph. Biane]

– Large deviation Q-SSEP

$$\mathbb{P}[\mathbf{n}(\cdot) \approx n(\cdot)] \underset{N \rightarrow \infty}{\asymp} e^{-N I_{\text{ssep}}[n]},$$

→ In terms of free probability data

$$I_{\text{ssep}}[n] = \max_{g(\cdot), q(\cdot)} \left(\int_0^1 dx \left[n(x) \log \left(\frac{n(x)}{g(x)} \right) + (1 - n(x)) \log \left(\frac{1 - n(x)}{1 - g(x)} \right) + q(x)g(x) \right] - F_0^{\text{ssep}}[q] \right)$$

with $F_0^{\text{ssep}}[a] = \sum_{n \geq 1} \frac{(-1)^{n-1}}{n} R_n(\mathbb{I}_{[a]}),$

n-th free cumulants of $\mathbb{I}_{[a]}(x) = \int_x^1 dy a(y)$
w.r.t. to the Lebesgue measure on $[0,1]$

– Three steps :

(i) : Q-SSEP /classical SSEP relation

(ii) : Free probability description of the Q-SSEP steady measure

(iii) : A relation between cumulants and « non-coincidants » cumulants
for Bernoulli variables

Emergence of free probability in noisy systems

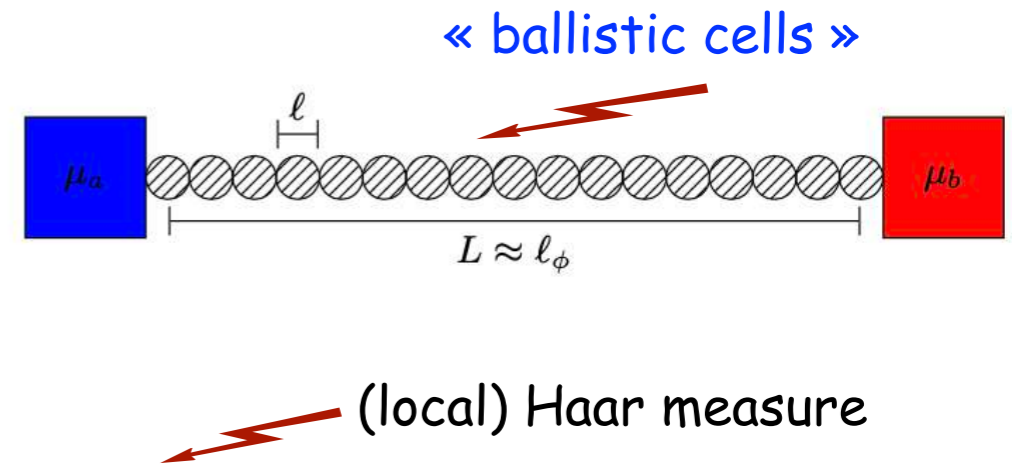
– Coarse-grained description (at mesoscopic scale)

(i) separation of time scales :

fast, closed dynamics within ballistic cells for $t < t_\ell$
 → unitary dynamics within each cells for $t < t_\ell$

(ii) ergodicity of the fast dynamics (→ noise) :

$$\mathbb{E}_t[G_{jk}] := \frac{1}{t_\ell} \int_t^{t+t_\ell} dt' G_{jk}(t') = \text{Tr}(\rho_t [c_i^\dagger c_j] U)$$



– Validity of U(1) sym. + MFT in mean → the « universality three conditions »

● U(1) invariance :

$$G_{jk} \equiv_{\text{in law}} e^{i\theta_j} G_{jk} e^{-i\theta_k}$$

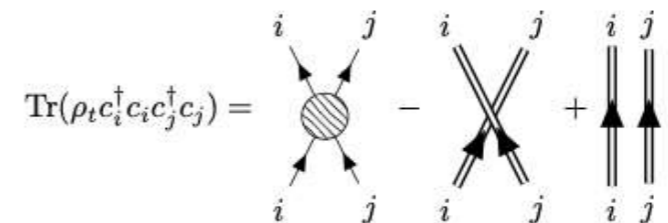
← local conservation and closed unitary dynamics at short time

● Scaling of the loop expectation values :

$$\mathbb{E}[G_{j_1 j_n} \cdots G_{j_3 j_2} G_{j_2 j_1}] \sim N^{1-n}$$

← If mean densities satisfy MFT

If some perturbation theory is valid ($H = H_0 + V$)



● Factorisation of loop expectations :

← closed fast dynamics / cells independence

– Validation / Violation of these three conditions ? ...