

# Composite Hybrid Inflation: Walking Dilaton and Waterfall Pions

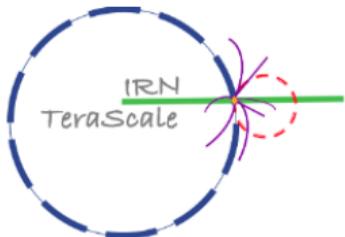
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# Content

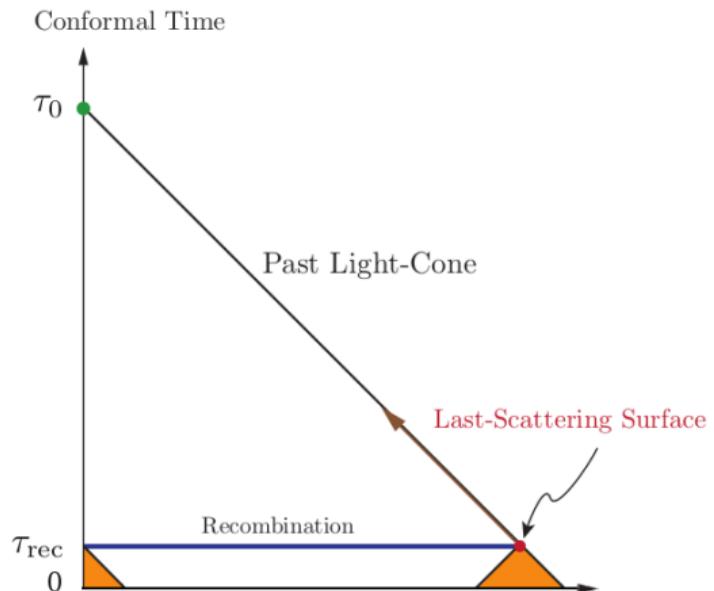
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- Introduction
- The model
- Inflationary Dynamics
- Constraints on the parameter space
- Conclusion

# Introduction : Inflation

Observations from the Cosmic Microwave Background (CMB) :

- Universe is very flat :  $\Omega_k \sim 0$   
     $\Rightarrow$  Flatness problem
- Very small thermal fluctuations, but  
events are causally disconnected  
 $\Rightarrow$  Horizon problem



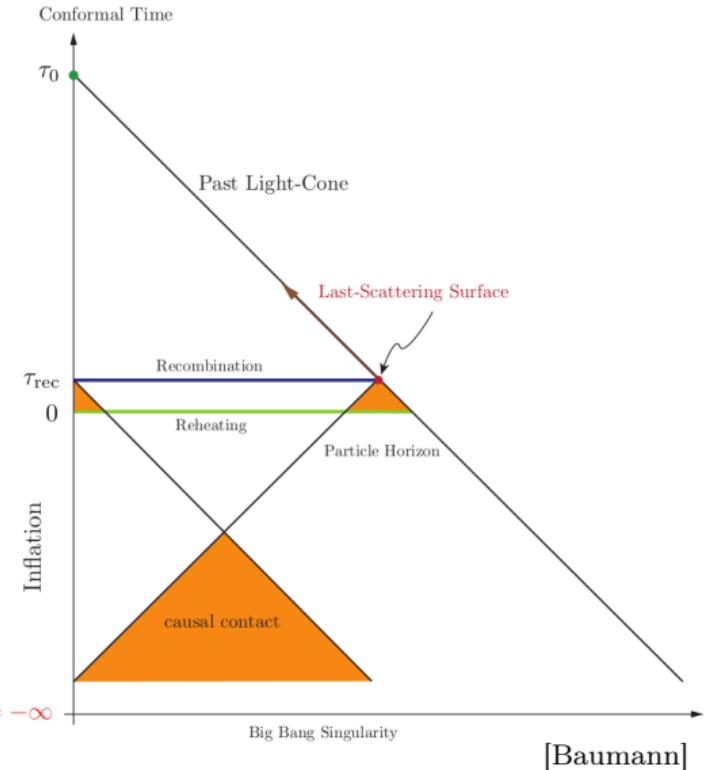
[Baumann]

# Introduction : Inflation

Observations from the Cosmic Microwave Background (CMB) :

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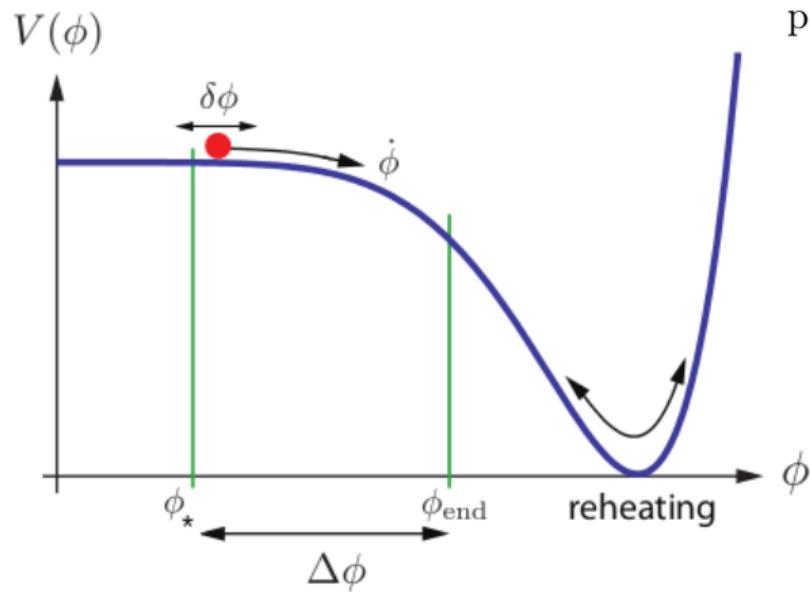
Inflation gives a simple explanation to both problems



[Baumann]

# Introduction : Slow-roll Inflation

Scalar fields  $\phi$  dominated by their potential  $V(\phi)$  are good candidates to have an accelerated expansion :



Inflation characterized by slow-roll parameters :

$$\varepsilon = -\frac{\dot{H}}{H^2} \simeq \frac{M_P^2}{2} \left( \frac{V'(\phi)}{V(\phi)} \right)^2 \ll 1 \quad (1)$$

$$\eta = -\frac{\ddot{\phi}}{H\dot{\phi}} \simeq M_P^2 \frac{V''(\phi)}{V(\phi)} \ll 1 \quad (2)$$

# Introduction : Slow-roll Inflation

CMB/Planck observations lead to precision tests for inflationary models :

- Duration :

$$N = \ln \left( \frac{a_{\text{end}}}{a_*} \right) = \int_{\phi_*}^{\phi_{\text{end}}} \frac{d\phi}{\sqrt{2\varepsilon}} \simeq 60$$

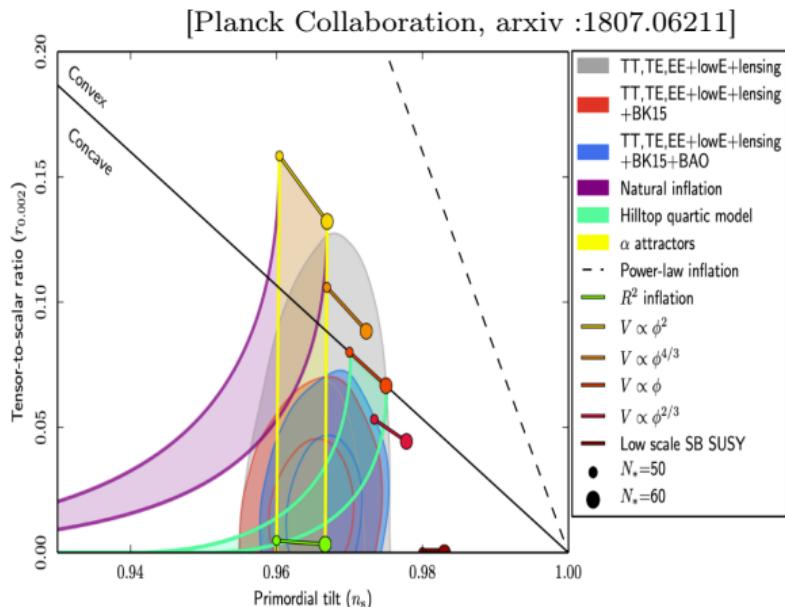
- Scalar to tensor ratio :  $r \simeq 16\varepsilon \leq 0.05$

- Spectral index :

$$n_s \simeq 1 - 6\varepsilon + 2\eta \simeq 0.965$$

- Power spectrum :

$$\mathcal{P}_{\mathcal{R}} \simeq \frac{V(\phi_*)}{24\pi^2 M_P^4 \varepsilon} \simeq 2 \times 10^{-9}$$



# Questions

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Can viable inflationary potential arise from BSM theories like composite models ?

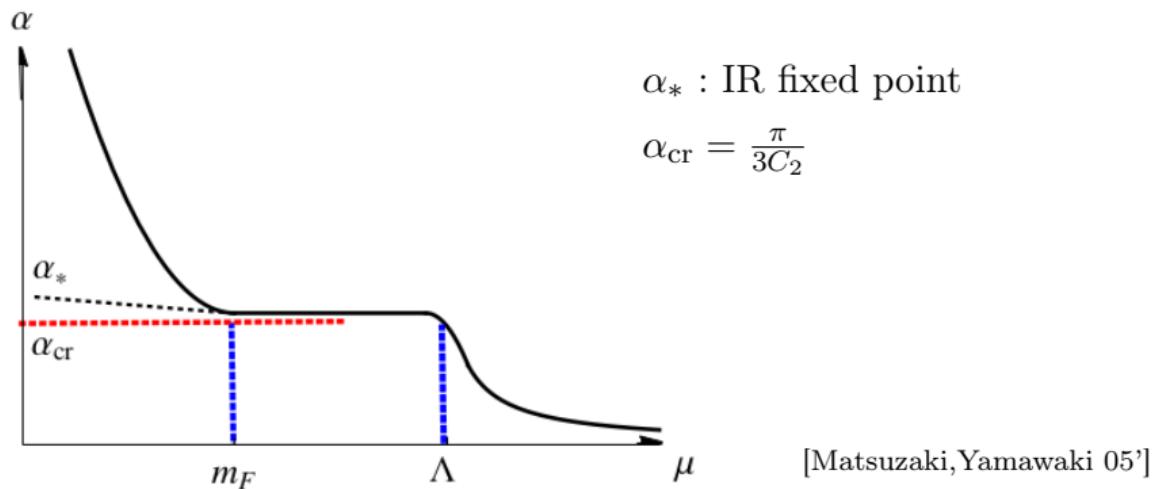
How inflation constrain such theories ? What is the phenomenology ?

# The model

- SU( $N_c$ ) gauge theory,  $N_f$  Dirac fermions (Fundamental)

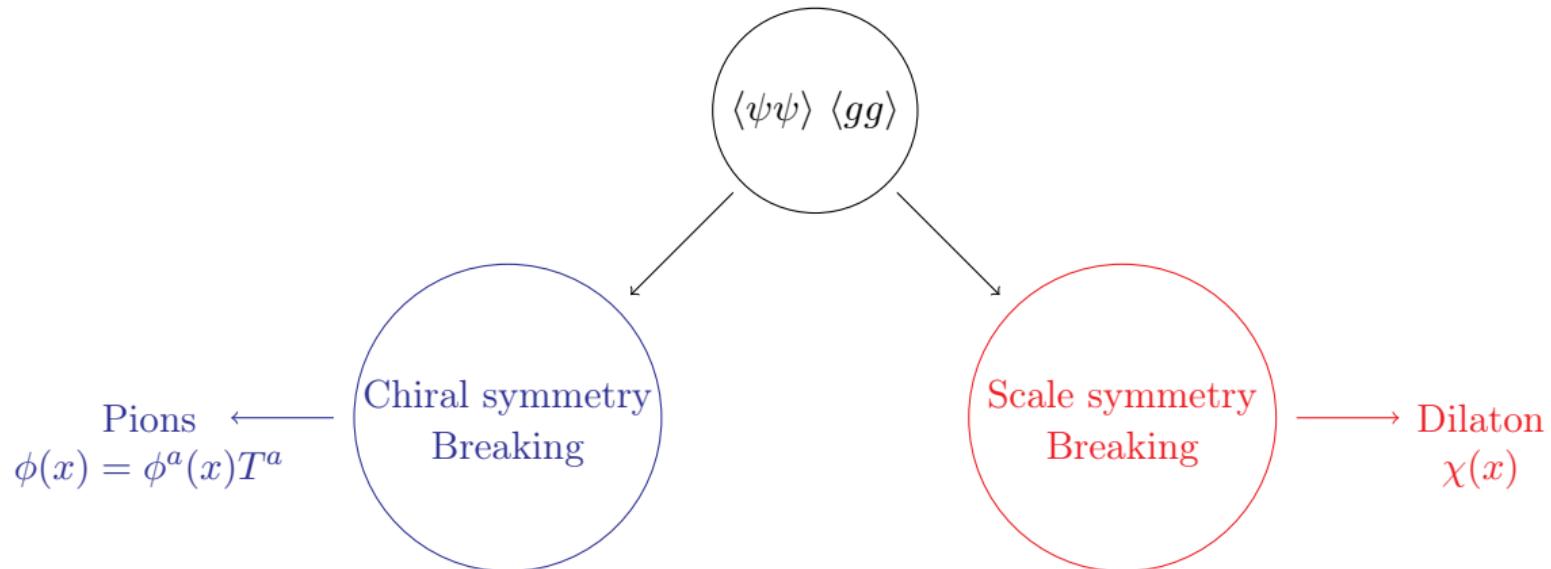
$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + \bar{\psi}_j (i\cancel{D} - m_0) \psi_j \quad (3)$$

- Such theory can exhibit a walking regime :



# The model

- In the IR,  $\mu \ll m_F$ , a fermion and/or a gluon condensate  $\langle\psi\psi\rangle$  (  $\langle gg \rangle$  ) can form.



# The model

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- Chiral Lagrangian with  $M_P = 1$ ,  $U(x) = \exp\left(\frac{i\phi}{f_\phi}\right)$  :

$$\begin{aligned} \frac{\mathcal{L}}{\sqrt{-g}} \supset & \frac{1}{2}R - \frac{1}{2}\partial_\mu\chi\partial^\mu\chi - \frac{f_\phi^2}{2}\left(\frac{\chi}{f_\chi}\right)^2 \text{Tr}\left[\partial_\mu U^\dagger \partial^\mu U\right] - \frac{\lambda_\chi}{4}\chi^4 \left(\log\frac{\chi}{f_\chi} - A\right) \\ & + \frac{\lambda_\chi\delta_1 f_\chi^4}{2}\left(\frac{\chi}{f_\chi}\right)^{3-\gamma_m} \text{Tr}\left[U + U^\dagger\right] + \frac{\lambda_\chi\delta_2 f_\chi^4}{4}\left(\frac{\chi}{f_\chi}\right)^{2(3-\gamma_{4f})} \text{Tr}\left[(U - U^\dagger)^2\right] \\ & - V_0 \end{aligned} \quad (4)$$

[Cacciapaglia, Pica, Sannino 20']

# The model

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Kinetic terms

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# The model

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- Chiral Lagrangian with  $M_P = 1$ ,  $U(x) = \exp\left(\frac{i\phi}{f_\phi}\right)$  :

Kinetic terms	Dilaton potential (CW)
$\frac{\mathcal{L}}{\sqrt{-g}} \supset \frac{1}{2}R - \frac{1}{2}\partial_\mu\chi\partial^\mu\chi - \frac{f_\phi^2}{2}\left(\frac{\chi}{f_\chi}\right)^2 \text{Tr}\left[\partial_\mu U^\dagger \partial^\mu U\right] - \frac{\lambda_\chi}{4}\chi^4 \left(\log \frac{\chi}{f_\chi} - A\right)$	
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$- V_0$	

[Cacciapaglia, Pica, Sannino 20']

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[Cacciapaglia, Pica, Sannino 20']

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$- V_0$	Pion mass	4 Fermi

[Cacciapaglia, Pica, Sannino 20']

# The model

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- Lattice data :  $0 < \gamma_m \leq 1$  and  $0 < \gamma_{4f} \leq 1$  [Leung,Love,Bardeen 89'] [LatKMI Collaboration 14', 17']
- Composite theory :  $\max(\delta_1, \delta_2) \lesssim \frac{1}{\lambda_\chi} \left(\frac{f_\phi}{f_\chi}\right)^4$
- De Sitter :  $V_0 > 0$

# The model

- Inflationary potential :

$$V(\phi, \chi) = -\lambda_\chi \delta_1 f_\chi^4 \left( \frac{\chi}{f_\chi} \right)^{3-\gamma_m} \cos \frac{\phi}{f_\phi} - \lambda_\chi \delta_2 f_\chi^4 \left( \frac{\chi}{f_\chi} \right)^{2(3-\gamma_{4f})} \sin^2 \frac{\phi}{f_\phi} + \frac{\lambda_\chi}{4} \chi^4 \left( \log \frac{\chi}{f_\chi} - A \right) + V_0 \quad (5)$$

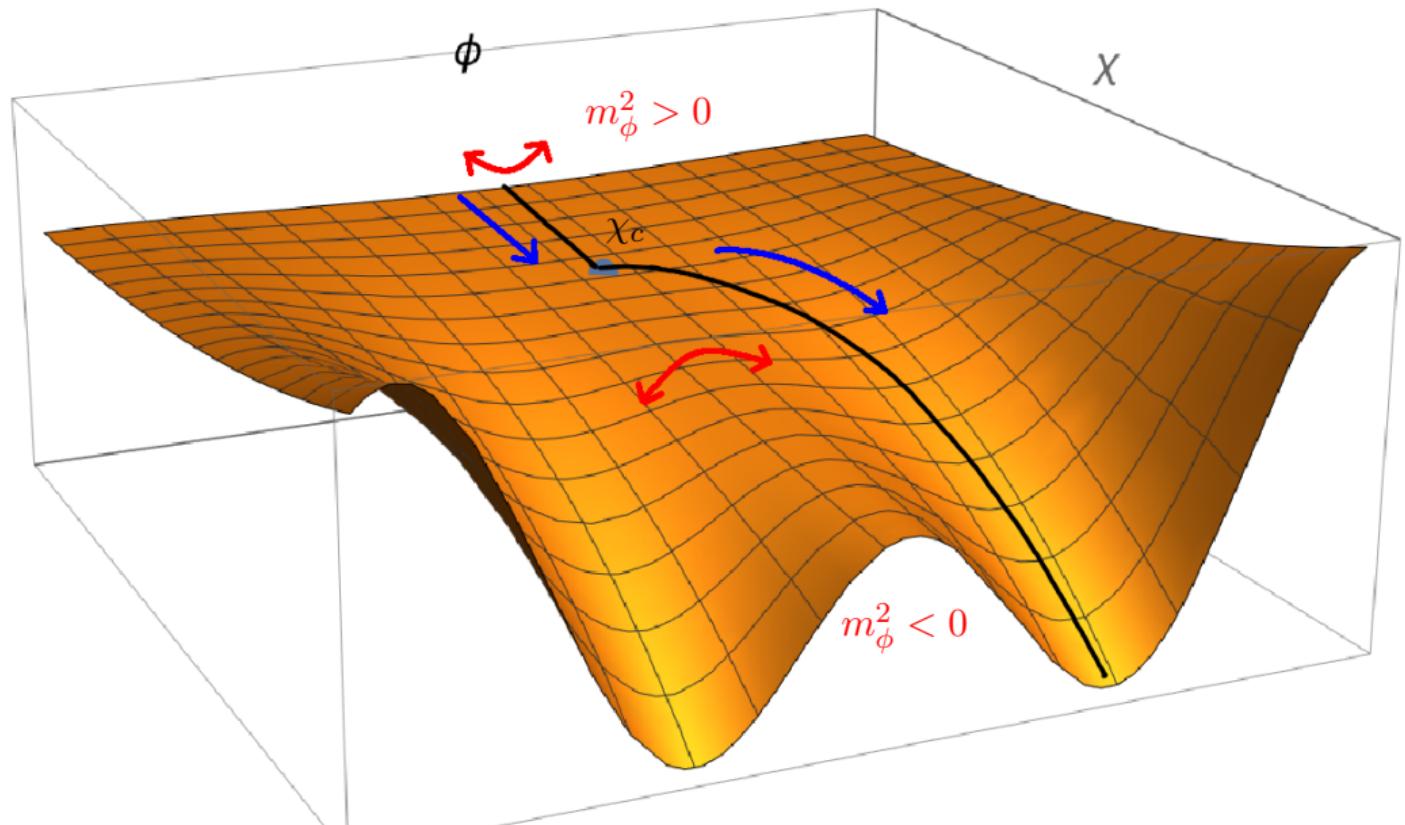
- VEV requirements :  $V(\phi_0, \chi_0) = 0, \quad V_\chi(\phi_0, \chi_0) = 0 \quad \text{and} \quad V_\phi(\phi_0, \chi_0) = 0$

$$\chi_0 = f_\chi, \quad \phi_0 = f_\phi \arccos \frac{\delta_1}{2\delta_2} \quad \text{for} \quad 0 < \delta_1 < 2\delta_2$$

$$V_0 = \frac{\lambda_\chi f_\chi^4}{16} \left[ 1 + 2 \frac{\delta_1^2}{\delta_2} (2 + \gamma_m - \gamma_{4f}) - 8\delta_2 (1 - \gamma_{4f}) \right] \quad (6)$$

$$A = \frac{1}{4} \left[ 1 + 2 \frac{\delta_1^2}{\delta_2} (\gamma_m - \gamma_{4f}) - 8\delta_2 (3 - \gamma_{4f}) \right]$$

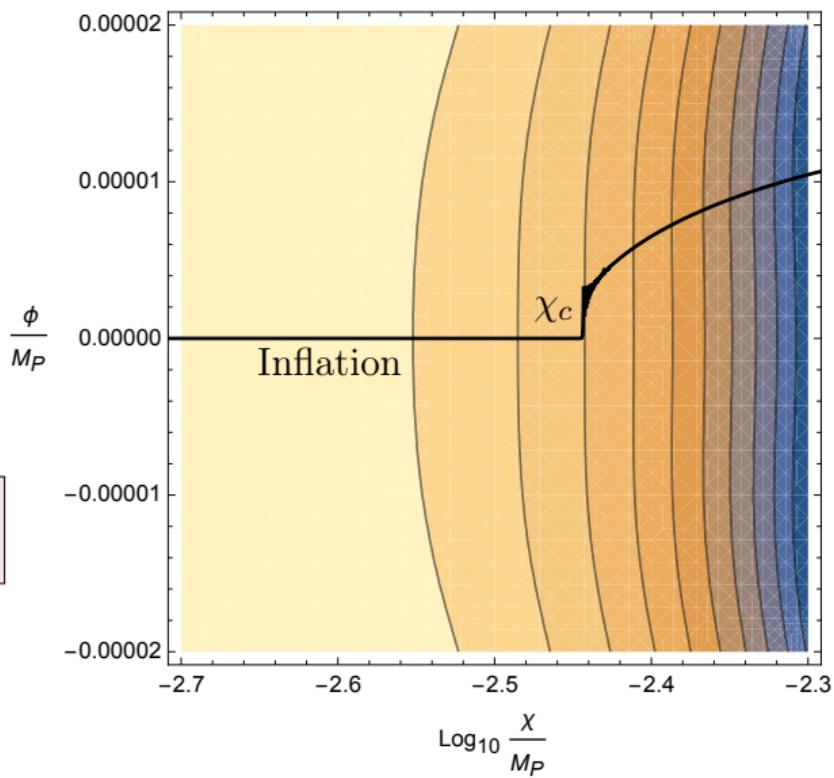
# Inflationary Dynamics : Waterfall



# Inflationary Dynamics : Waterfall

- $\chi_c = \chi(m_\phi^2 = 0) = f_\chi \left( \frac{\delta_1}{2\delta_2} \right)^{\frac{1}{1+\gamma_m - 2\gamma_{4f}}}$
- $\phi(\chi < \chi_c) = 0$
- $\phi(\chi > \chi_c) = f_\phi \cos^{-1} \frac{\delta_1}{2\delta_2} \left( \frac{f_\chi}{\chi} \right)^{3+\gamma_m - 2\gamma_{4f}}$

Tachyonic instability terminates inflation at  $\chi = \chi_c$



# Constraints on the parameter space

- Rapid termination of inflation :

$$\left| m_\phi^2 \right|_{\phi=0, \chi=\chi_c + \Delta\chi} \gg H^2$$

$\implies$  Constrain :  $f_\phi < f_\phi^{\text{waterfall}}$

- Composite consistency :

$$\max(\delta_1, \delta_2) \lesssim \frac{1}{\lambda_\chi} \left( \frac{f_\phi}{f_\chi} \right)^4$$

$\implies$  Constrain :  $f_\phi > f_\phi^{\text{composite}}$

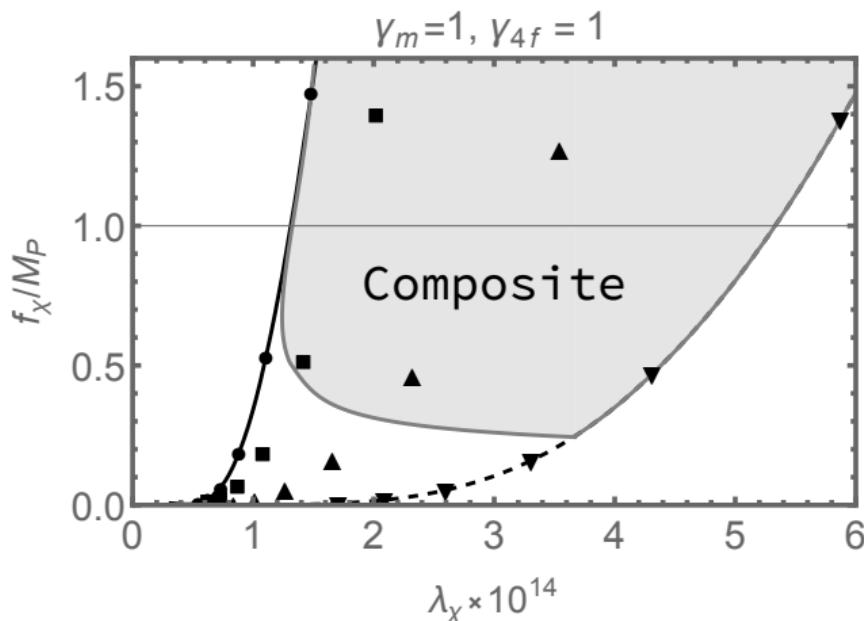
- De Sitter :

$$V_0 > 0$$

$\implies$  Constrain :  $\delta_2 < \delta_2^{\text{dS}}$

## Constraints on the parameter space

- Setting  $\gamma_m = \gamma_{4f} = 1$ , we have 5 free parameters :  $\delta_1, \delta_2, \chi_*, f_\chi, \lambda_\chi$
- CMB observations  $(n_s, N, \mathcal{P}_R)$  allow us to fix 3 parameters :  $\delta_1, \delta_2, \chi_*$



- Predictions :

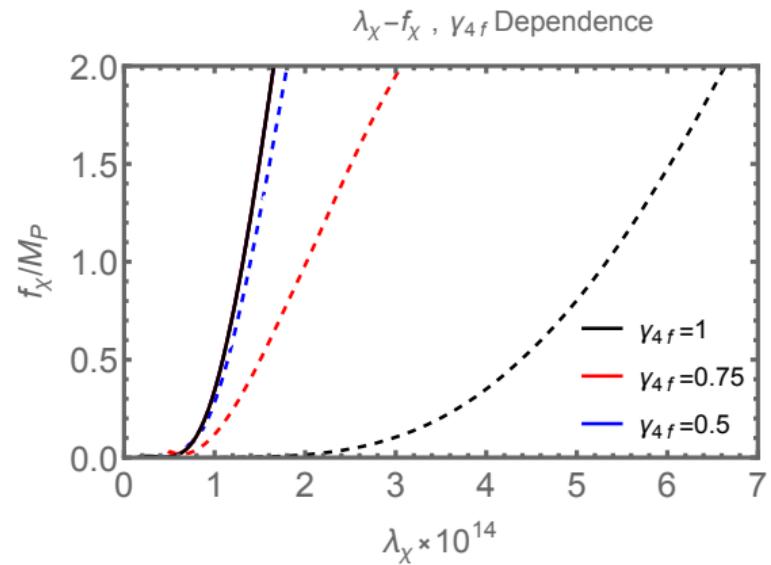
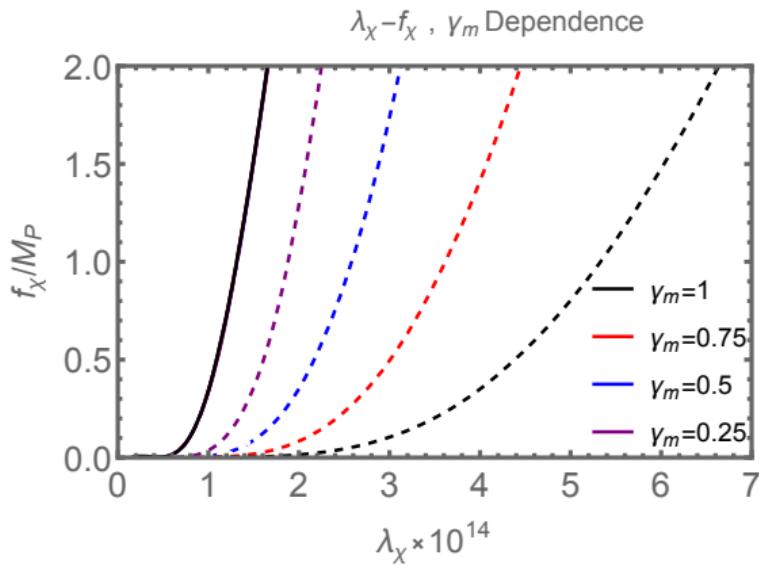
$$m_\chi|_{\chi_0, \phi_0} \sim 10^{11} \text{ GeV}$$

$$m_\phi|_{\chi_0, \phi_0} \sim 10^{14} \text{ GeV}$$

$$H_{\text{inf}} \sim 10^{10} \text{ GeV}$$

# Constraints on the parameter space

- Dependence on  $\gamma_m$  and  $\gamma_{4f}$  :



# Conclusion

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- Inflation provides a simple way to answer the Flatness and Horizon problems
- Inflation requires a particle to drive the accelerated expansion

⇒ Test if BSM theories can provide a successful framework for inflation

- Study for a composite gauge theory

⇒ Compatibility with CMB observations constrains a lot the parameter space of such theories

- Next : Reheating ? Topological defects ? Gravitational waves production ?  
Non-minimal coupling ?

## Walking Regime

- Beta function :  $\beta^{\text{2loop}}(\alpha) \equiv \mu \frac{\partial \alpha}{\partial \mu} = -b_0 \alpha^2 - b_1 \alpha^3$

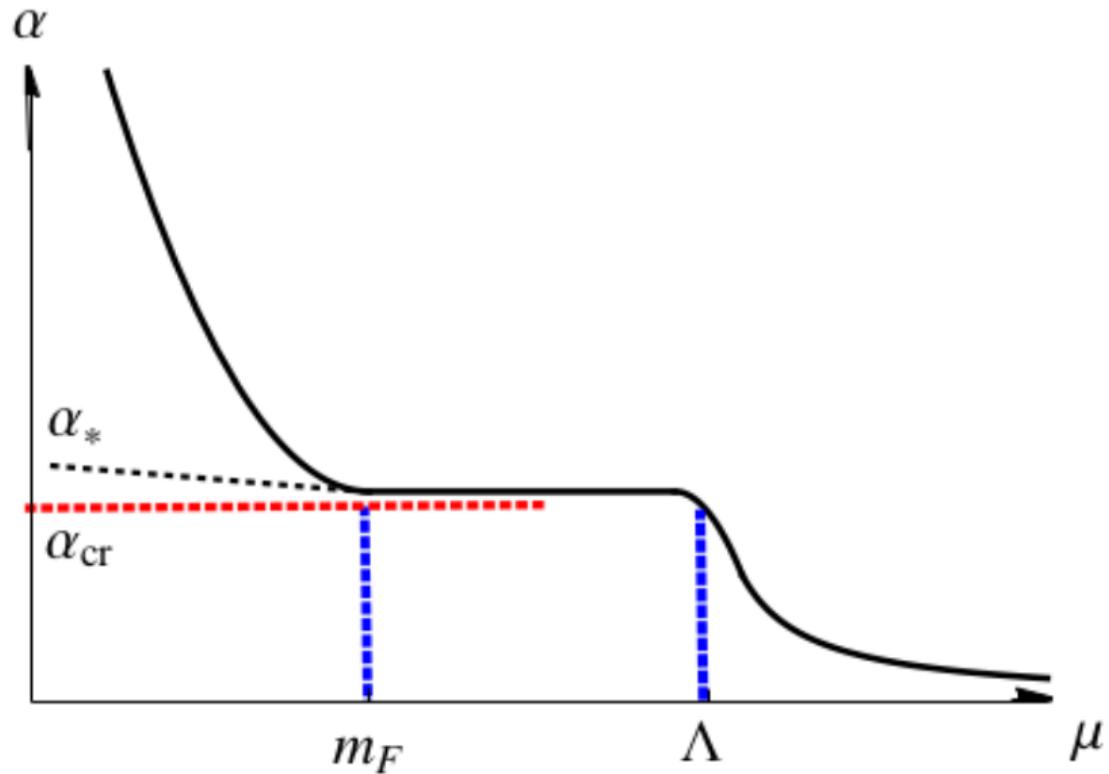
$$b_0 = \frac{1}{6\pi}(11N_c - 2N_f), \quad b_1 = \frac{1}{24\pi^2} \left( 34N_c^2 - 10N_c N_f - 3 \frac{N_c^2 - 1}{N_c} N_f \right)$$

- IR fixed point :  $\alpha_* = -\frac{b_0}{b_1} = \frac{4\pi(11-2n_f)}{N_c(13n_f-34)}$
- Intrinsic scale :  $\Lambda \equiv \mu \exp \left( - \int^{\alpha(\mu^2)} \frac{d\alpha}{\beta^{\text{2loop}}(\alpha)} \right)$ , leads to  $\alpha(\mu^2) = \frac{\alpha_*}{1 + W \left( \frac{1}{e} \left( \frac{\mu}{\Lambda} \right)^{b_0 \alpha_*} \right)}$

$$\implies \alpha(\mu^2)^{\text{IR}} \sim \frac{\alpha_*}{1 + \frac{1}{e} \left( \frac{\mu}{\Lambda} \right)^{b_0 \alpha_*}} \quad \alpha(\mu^2)^{\text{UV}} \sim \frac{1}{b_0 \ln \frac{\mu}{\Lambda}}$$

$\implies$  Walking regime for  $\mu \ll \Lambda$

## Walking Regime



## Walking Regime

- Miranski scaling :  $m_F \simeq 4\Lambda \exp\left(\frac{-\pi}{\sqrt{\frac{\alpha}{\alpha_{\text{cr}}}-1}}\right)$

$$\implies \beta(\alpha) = \frac{\partial \alpha}{\partial \ln \Lambda} = -\frac{2\alpha_{\text{cr}}}{\pi} \left(\frac{\alpha}{\alpha_{\text{cr}}} - 1\right)^{\frac{3}{2}}$$

- Solving this equation yields :  $\alpha(\mu) = \alpha_{\text{cr}} \left(1 + \frac{\pi^2}{\ln^2 \frac{\mu}{m_F}}\right)$

$\implies$  Walking regime for  $\mu \gg m_F$

- Conclusion : Walking regime for  $m_F \ll \mu \ll \Lambda$  where  $\alpha \sim \alpha_* \sim \alpha_{\text{cr}}$

# Single field inflation

- Corresponds to  $\phi = 0, \delta_2 = 0$  such that :

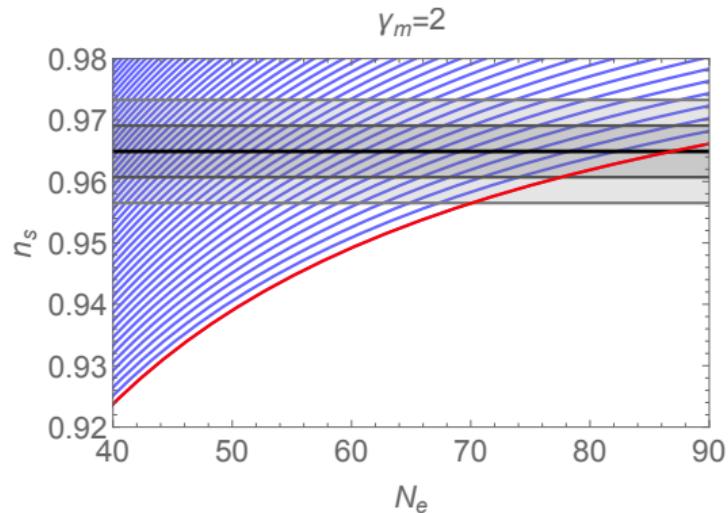
$$V(\phi = 0, \chi) = -\lambda_\chi \delta_1 f_\chi^4 \left( \frac{\chi}{f_\chi} \right)^{3-\gamma_m} + \frac{\lambda_\chi}{4} \chi^4 \left( \log \frac{\chi}{f_\chi} - A^{\text{single}} \right) + V_0^{\text{single}} \quad (7)$$

$$A^{\text{single}} = \frac{1}{4} + \delta_1(\gamma_m - 3) \quad (8)$$

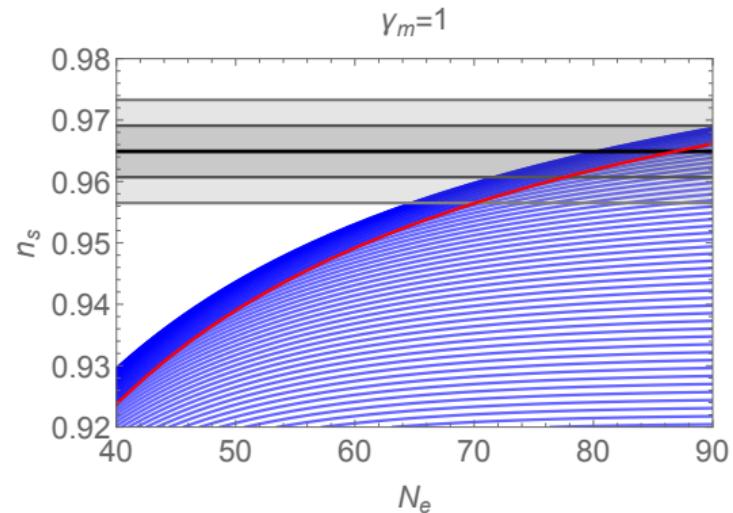
$$V_0^{\text{single}} = \frac{1}{16} \lambda_\chi f_\chi^4 (1 + 4\delta_1(1 + \gamma_m)) \quad (9)$$

$\implies$  CMB observations  $(n_s, N, \mathcal{P}_R)$  fix  $(\lambda_\chi, f_\chi, \chi_*)$

# Single field inflation



Red :  $\delta_1 = 0$



Blue :  $\delta_1 \in [0, 1] \times 10^{-6}$