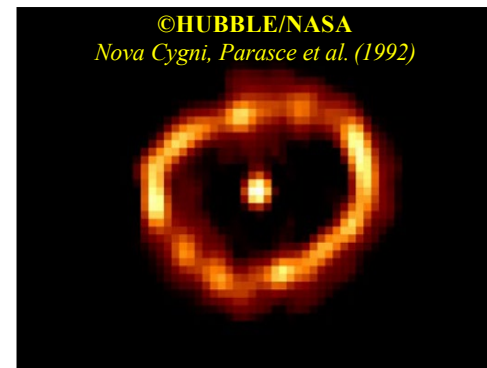
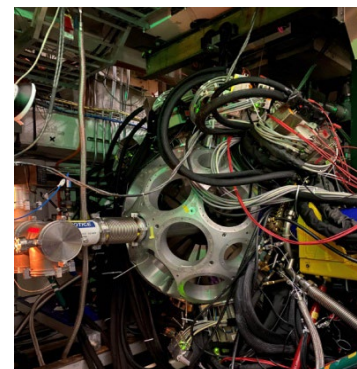
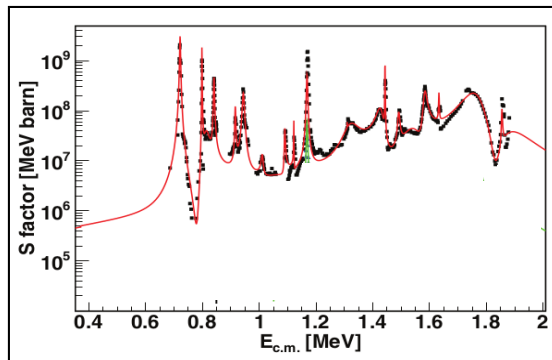
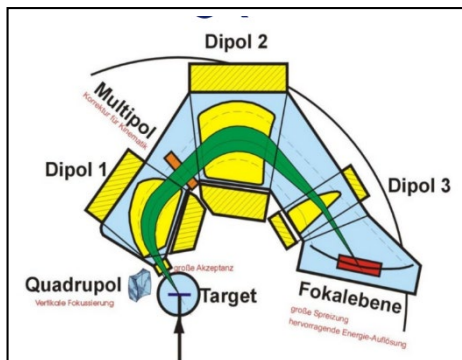
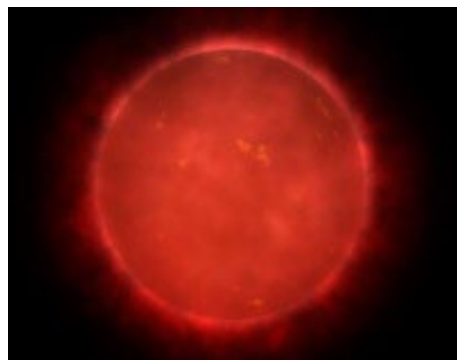




Recent reaction studies of astrophysical interest using transfer reactions

Faïrouz Hammache

Laboratoire de Physique des 2 Infinis Irène Joliot-Curie (IJCLab)

Reaction cross sections in stellar environments

Reactions in quiescent burning stages :

→ $T \sim 0.001 - 0.1$ GK

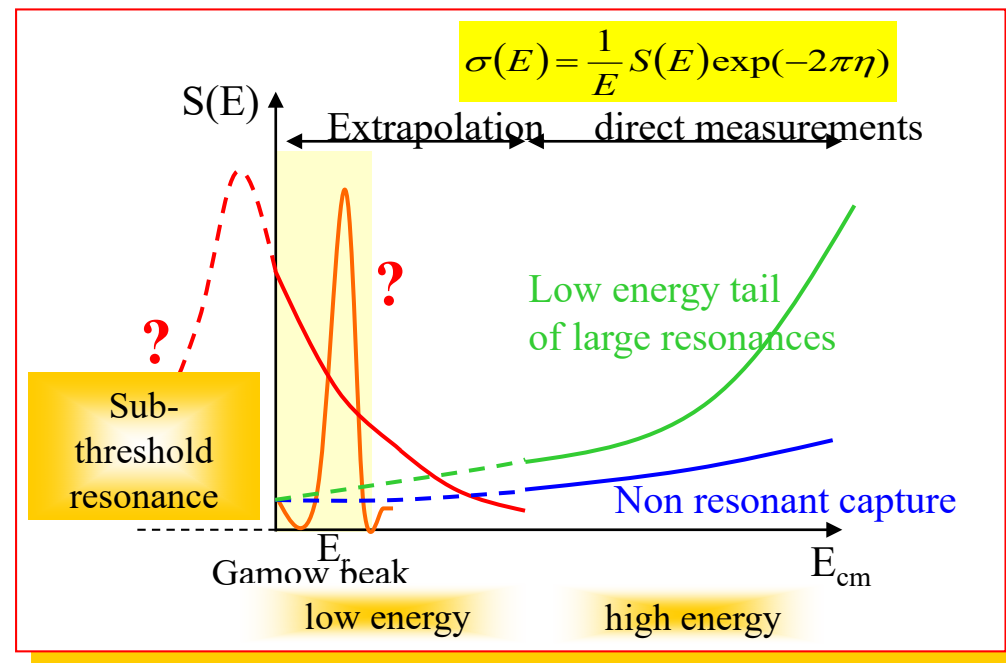
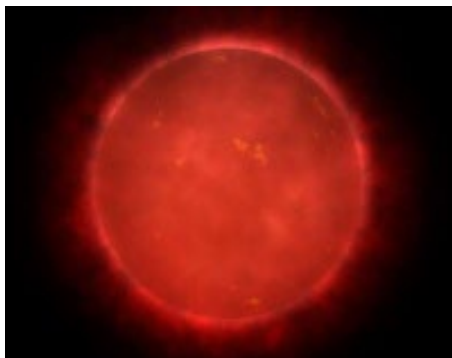
→ **tens keV- MeV** $\ll$ Coulomb barrier E_C

→ $\sigma(E)$ very weak (**0.01 fb- ~100 pb**)

→ Direct measurements are very challenging

→ Often $\sigma(E)$ measured at **high energies** then **extrapolated** at stellar energies

↳ what about **resonances** at very low energy & sub-threshold resonances



Reactions in explosive environments (Novae, X-ray bursts, Supernovae) :

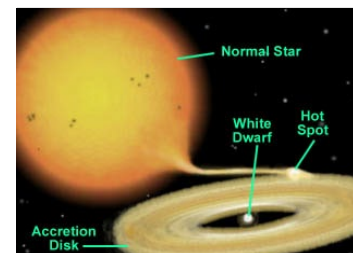
→ $T=0.1-2$ GK

→ **hundreds keV- few MeV**

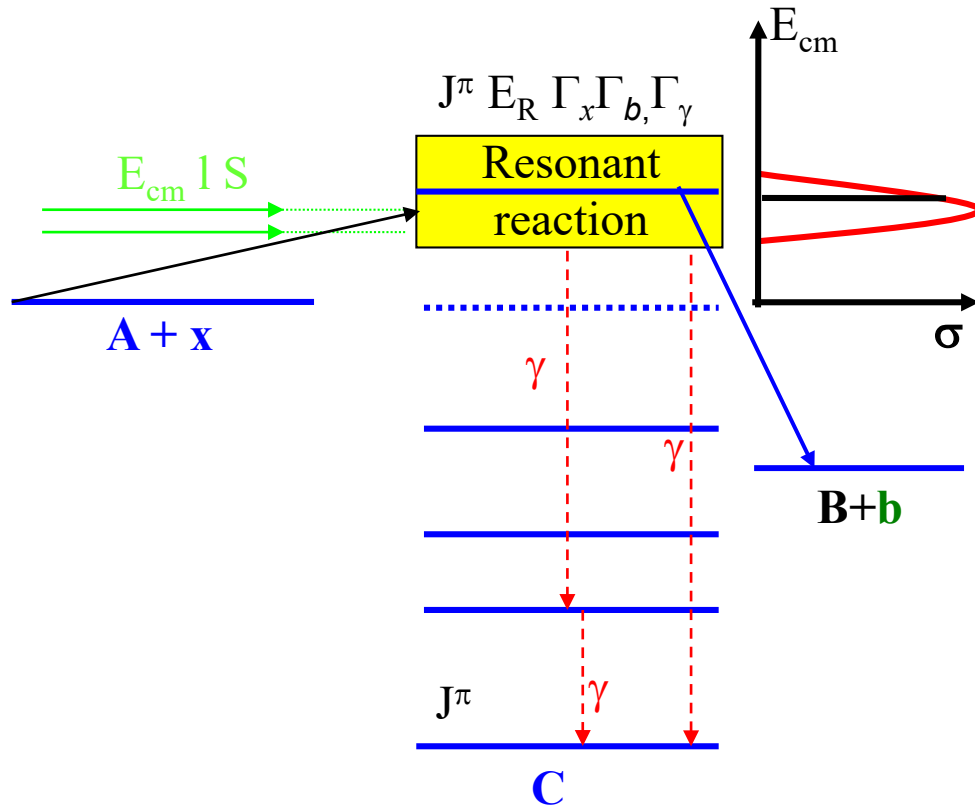
→ σ dominated by **resonances**

→ radioactive nuclei : Low beam intensities ($\sim 10^5 - 10^7$ p/s) for short lived nuclei, **few atoms/cm²** in targets for longed lived ones

↳ Direct measurements challenging



Ingredients needed to evaluate resonant σ & $N_A \langle \sigma v \rangle$



$$\sigma(E) = \pi \hat{\lambda}^2 \frac{2J_c + 1}{(2J_A + 1)(2J_x + 1)} \frac{\Gamma_x \Gamma_y}{(E - E_R)^2 + \Gamma_{tot}^2 / 4} \quad y = b \text{ or } \gamma$$

For narrow resonance cases :

$$\langle \sigma v \rangle = \left(\frac{2\pi}{\mu k T} \right)^{\frac{3}{2}} \hbar (\omega \gamma)_R \exp\left(-\frac{E_R}{k T}\right)$$

$$\rightarrow (\omega \gamma)_R = \frac{2J_c + 1}{(2J_A + 1) \cdot (2J_x + 1)} \frac{\Gamma_x \Gamma_y}{\Gamma_{tot}}$$

➤ The resonant reaction rates can be calculated if the resonant parameters (E_R , J^π , $\Gamma_{x,y}$, Γ , ...) are known



Transfer reactions can be used to extract these spectroscopic information

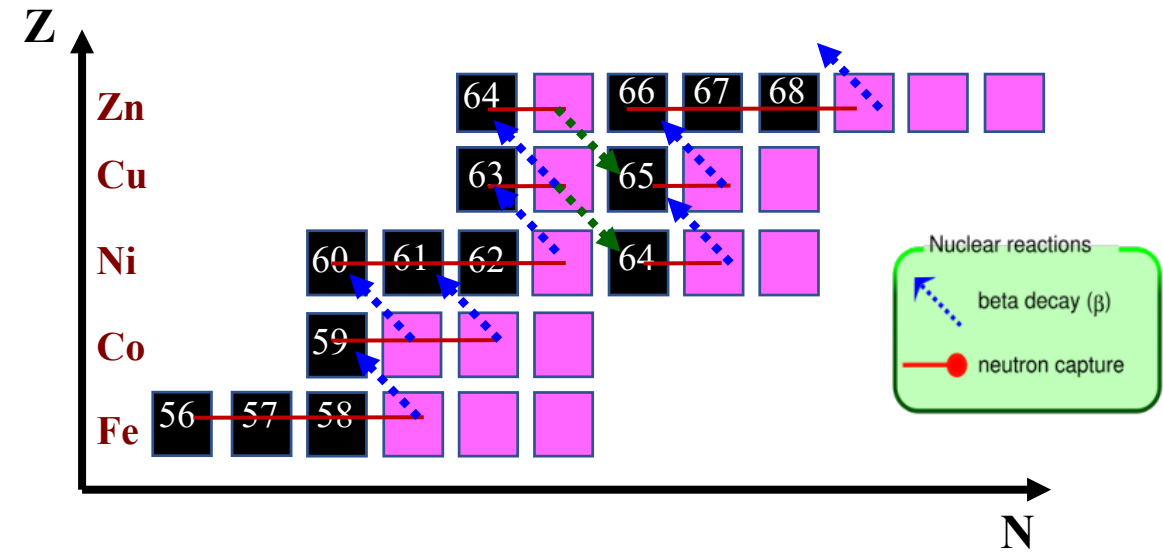
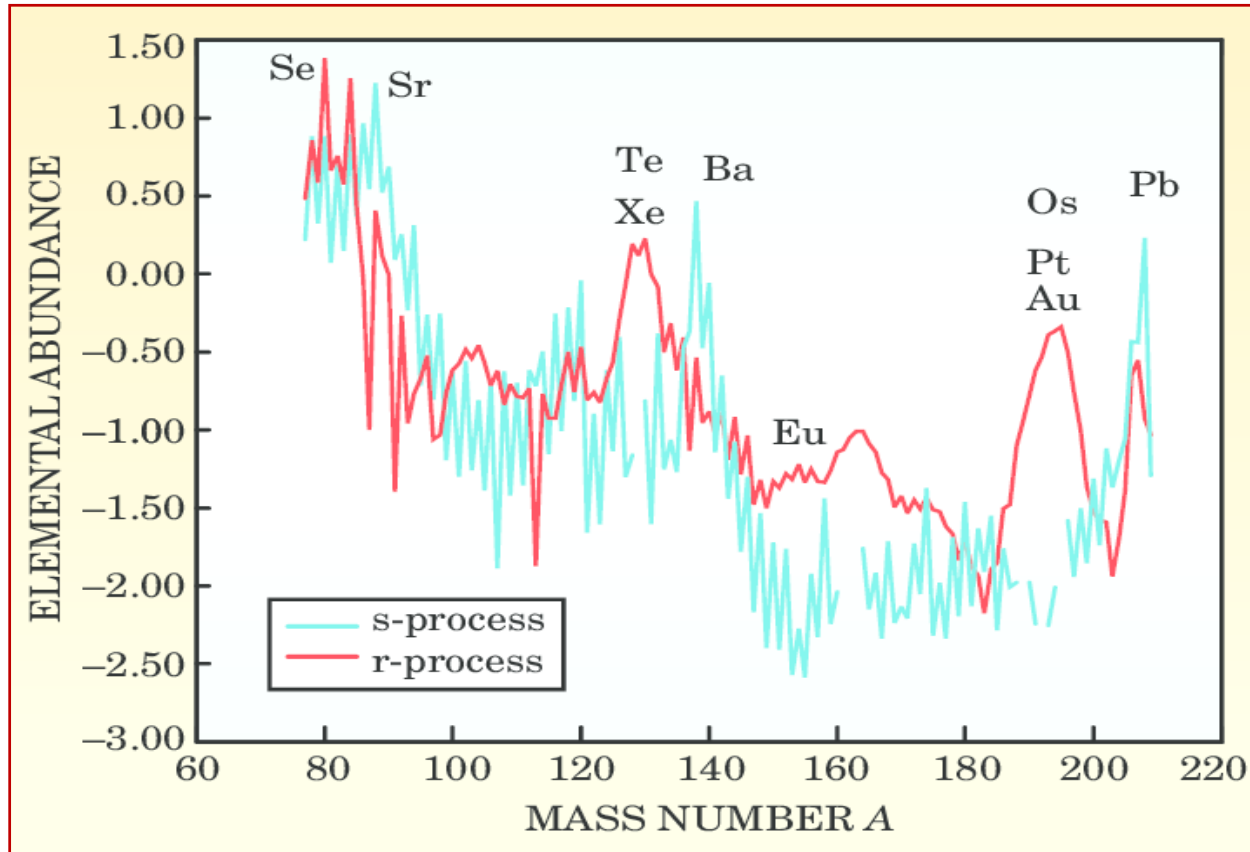
- $E_R = E_x - Q$ is the resonance energy
- $\Gamma_x, \Gamma_b, \Gamma_\gamma$ are the partial decay widths
- $\Gamma_{tot} = \Gamma_x + \Gamma_b + \Gamma_\gamma$ is the total width
- J_A, J_C, J_x : total angular momentum of the nuclei A, C & x

s-process

s-process (s = slow neutron capture process)

$$\tau_{\beta} \ll \tau_n \Leftrightarrow N_n \sim 10^8 - 10^{11} \text{ n/cm}^3$$

→ production of half of the abundance of heavy elements



➤ **Main** component

→ production of $90 < A < 209$ elements in low-mass AGB stars $1-4 M_{\odot}$ ($T \sim 0.1$ GK)

➤ **Weak** component

→ production of $56 < A < 90$ elements in intermediate-mass AGB star ($T \sim 0.3$ GK) & **massive stars** ($T \sim 0.2-1$ GK)

s-process in rotating metal-poor massive stars

Weak s-process → $56 < A < 90$ in **massive stars** $M > 8M_{\odot}$

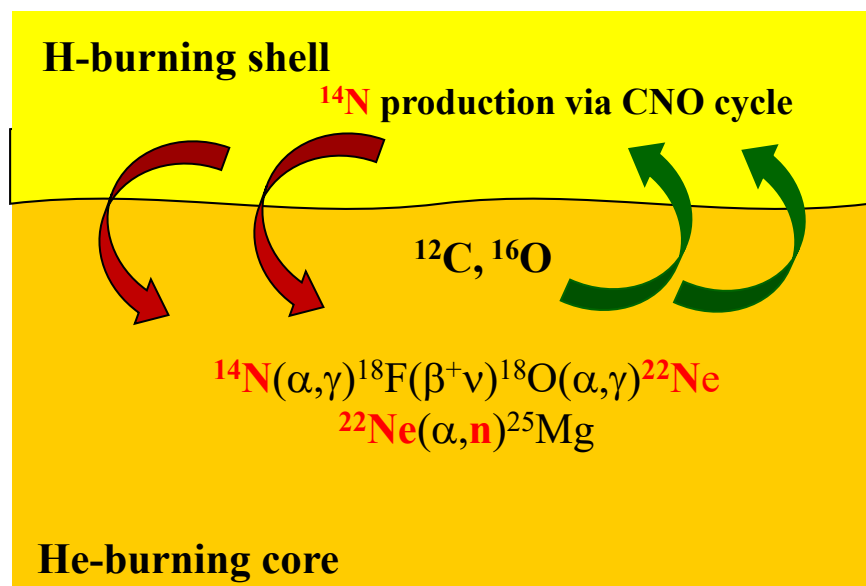
➤ End of core He burning ($T \sim 3 \cdot 10^8$ K, $N_n = 10^6 \text{ cm}^{-3}$)

➤ C-shell burning ($T \sim 1$ GK, $N_n = 10^{11} \text{ cm}^{-3}$)

Main neutron source: $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$



- **Metal-poor** massive stars ($[\text{Fe}/\text{H}] < -1$) → negligible s-process production (low ^{22}Ne & Fe seed abundance)
- With fast rotation induced mixing → ^{22}Ne production in He core strongly enhanced Nishimura+16, Choplin+18



↳ large production of s-elements between Strontium & Barium $90 < A < 140$

- Enhanced weak s-process (es-process) Frischknecht+16

→ Important impact on chemical enrichment in early galaxies.

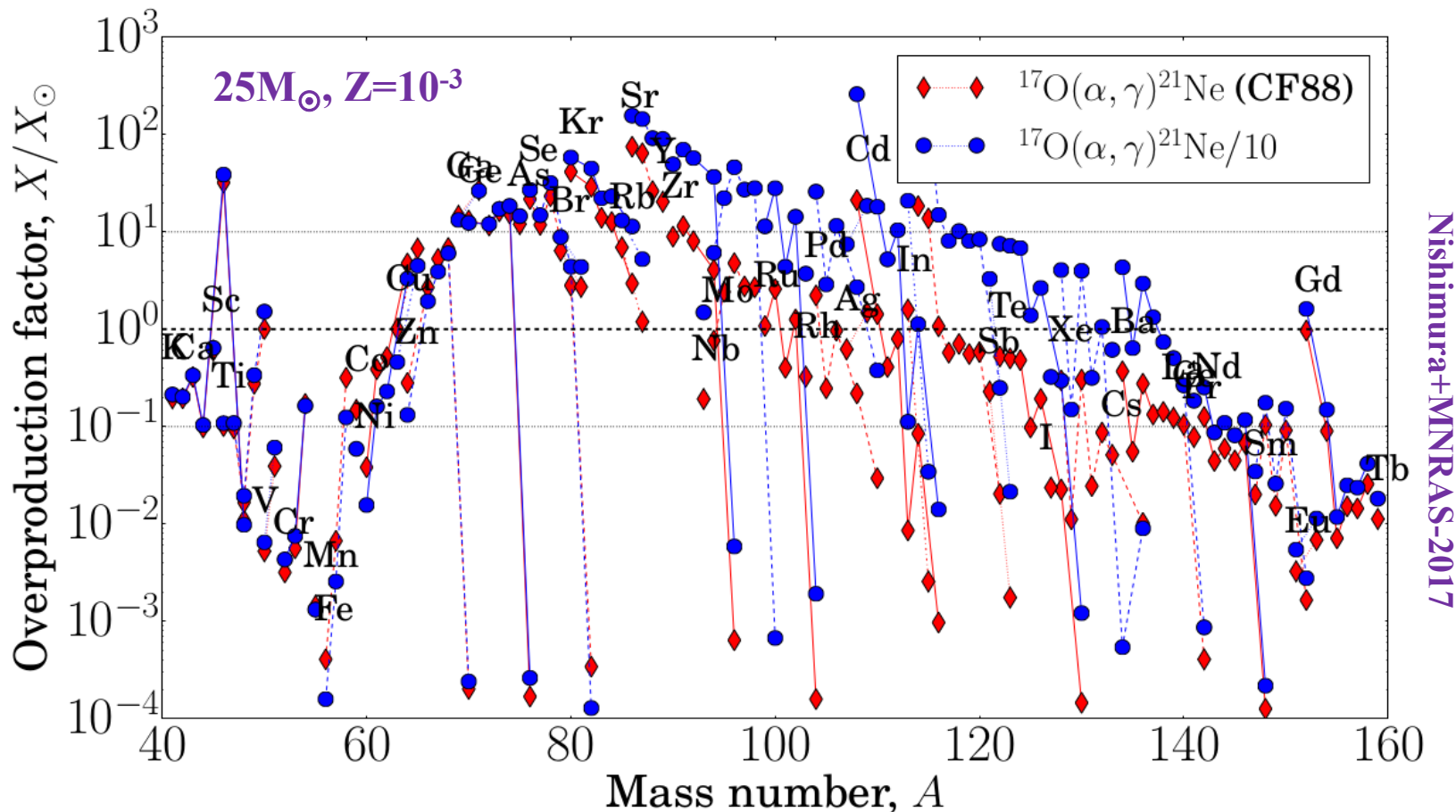
→ Source of heavy elements such as Barium in early universe? Barbuy+14

s-process in rotating metal-poor massive stars

But the final abundances of the enhanced weak s-process strongly depends on:

$^{16}\text{O}(n,\gamma)^{17}\text{O}$ neutron poison effect & $^{17}\text{O}(\alpha,n)/^{17}\text{O}(\alpha,\gamma)$ reaction rate ratio

→ neutron recycling efficiency



Nishimura+MNRAS-2017

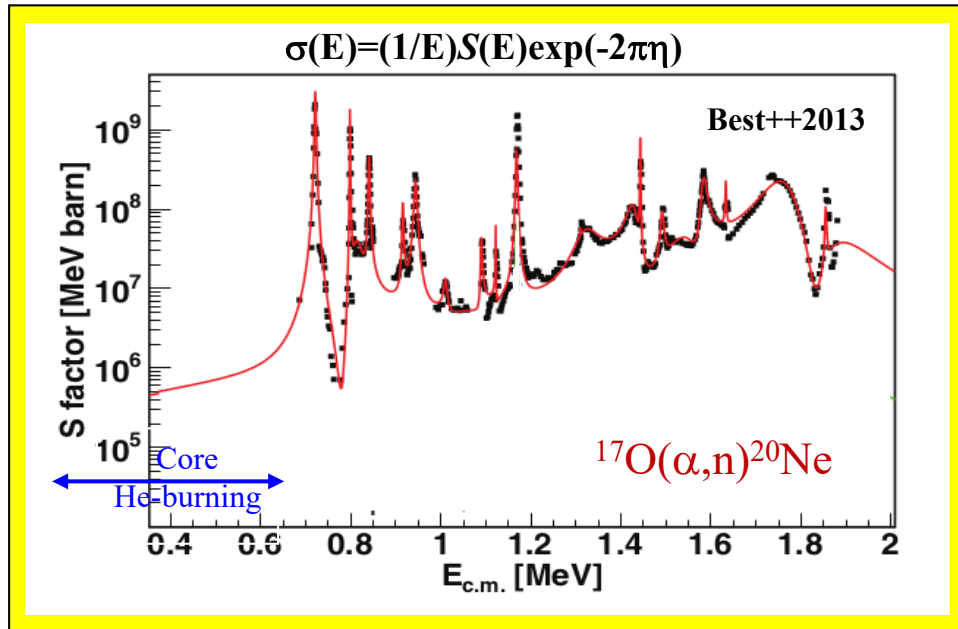
Calculation with $^{17}\text{O}(\alpha,n)^{20}\text{Ne}$ Nacre adopted rate & $^{17}\text{O}(\alpha,\gamma)^{22}\text{Ne}$ CF88 rate

Experimental status on $^{17}\text{O}(\alpha, n)^{20}\text{Ne}$ and $^{17}\text{O}(\alpha, \gamma)^{21}\text{Ne}$

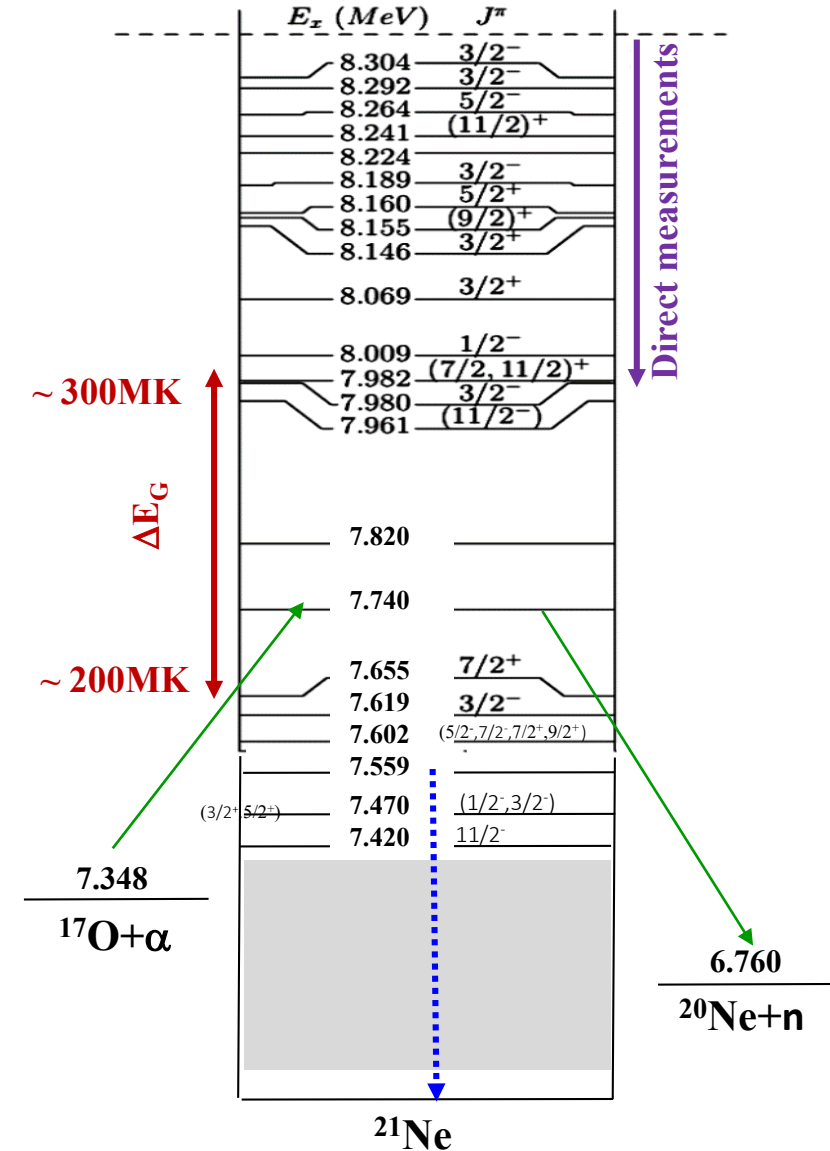
- Core He burning: $T \sim 0.2\text{-}0.3 \text{ GK} \rightarrow E_{\text{c.m.}} \sim 0.297\text{-}0.646 \text{ MeV} \rightarrow E_x = 7.64\text{-}8.00$ in ^{21}Ne
- Shell Carbon burning: $T \sim 1 \text{ GK} \rightarrow E_{\text{c.m.}} \sim 0.783\text{-}1.5 \text{ MeV} \rightarrow E_x = 8.13\text{-}8.85$ in ^{21}Ne

$^{17}\text{O}(\alpha, n)^{20}\text{Ne}$ & $^{17}\text{O}(\alpha, \gamma)^{21}\text{Ne}$ direct measurements:

- Denker++1994, Best++2013
 - Best++2011, Taggart++2019
 - Williams++2022
- No direct measurements @
Core He burning energies ($E_{\text{c.m.}} < 0.630 \text{ MeV}$)



- few **hundreds keV** $\ll E_{\text{Coulomb}}$
- $\sigma(E)$ very weak ($\leq 100 \text{ pb}$)
- Neutron & γ detection: large background

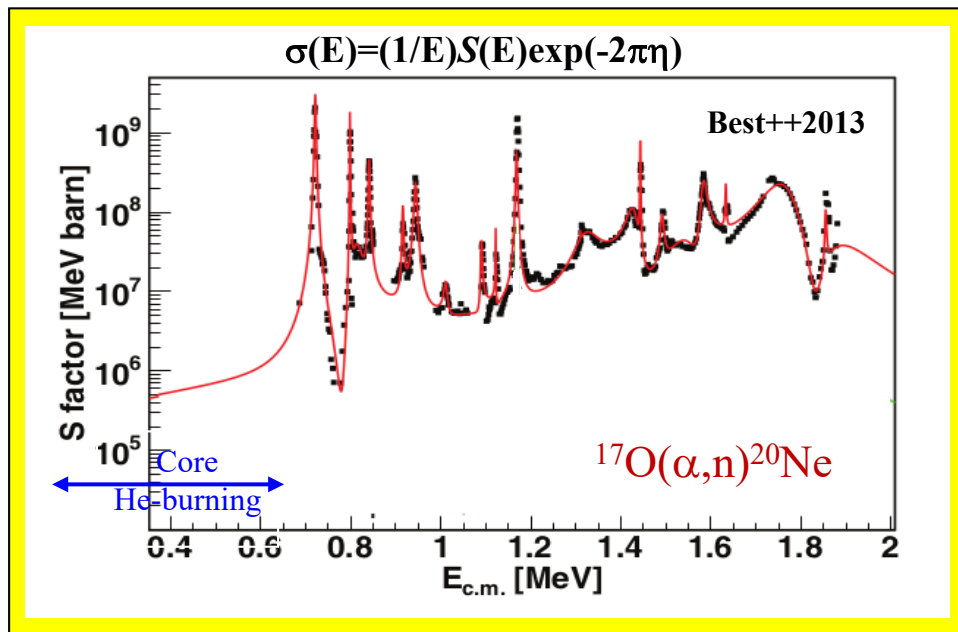


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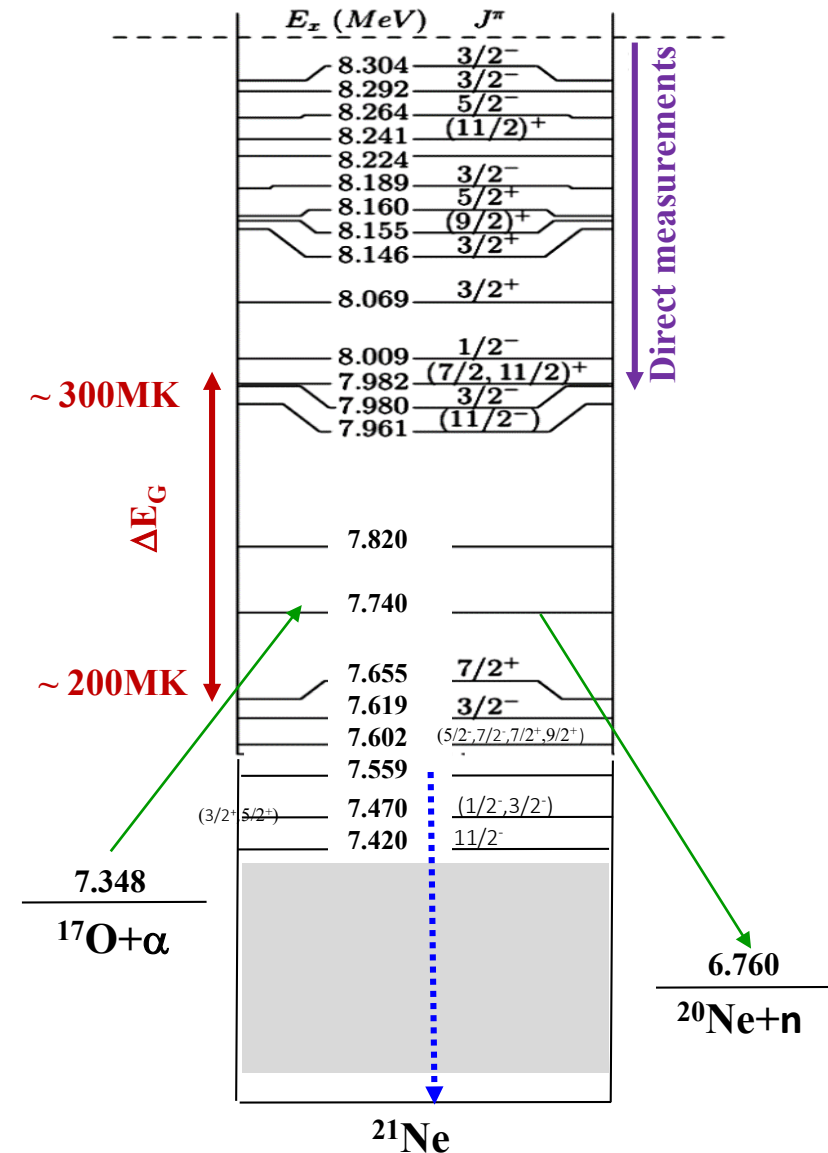
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Core He burning energies ($E_{\text{c.m.}} < 0.630 \text{ MeV}$)



What about Spectroscopy of ^{21}Ne ?

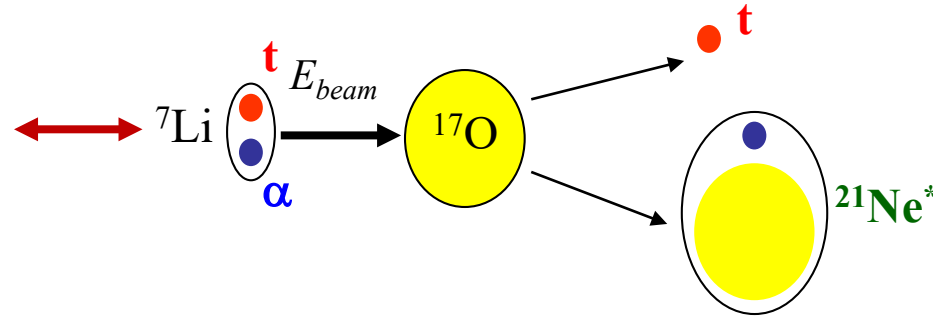
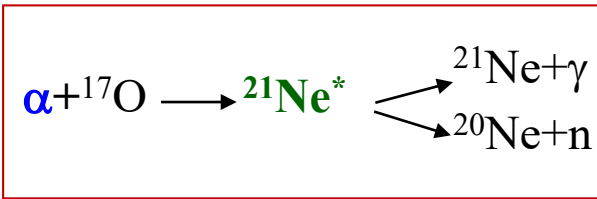
- $\rightarrow$ Known E_x
- $\rightarrow$ Unknown or poorly known Γ_α , Γ_n , $\Gamma_\gamma/\Gamma_{\text{tot}}$
- $\rightarrow$ Few have spin-parity assignments J^π



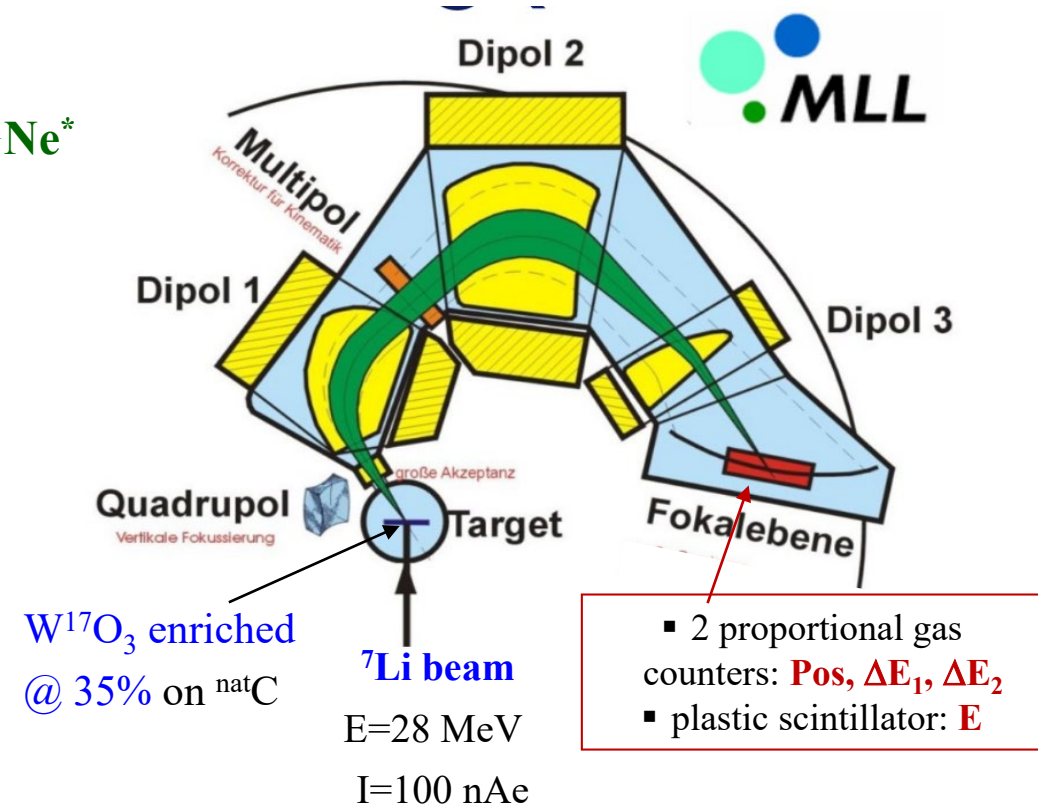
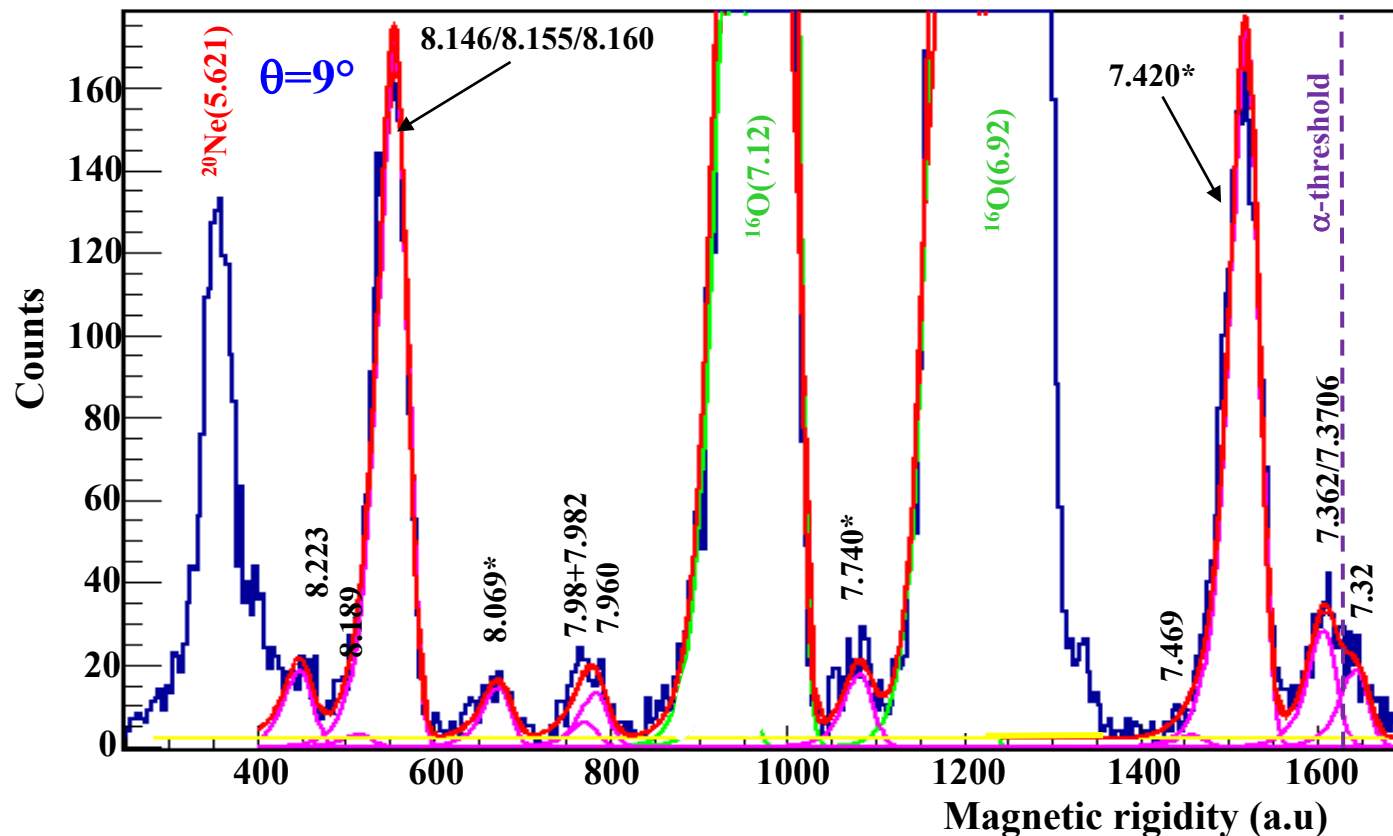
- Neutron transfer reaction $\rightarrow S_n \rightarrow \Gamma_n$ Frost-Schenk++MNRAS2022
- α -transfer reaction $\rightarrow S_\alpha \rightarrow \Gamma_\alpha$ (present work/MLL-exp)

Study of ^{21}Ne states via $^{17}\text{O}(^7\text{Li},t)^{21}\text{Ne}$ α -transfer reaction

Q3D spectrometer

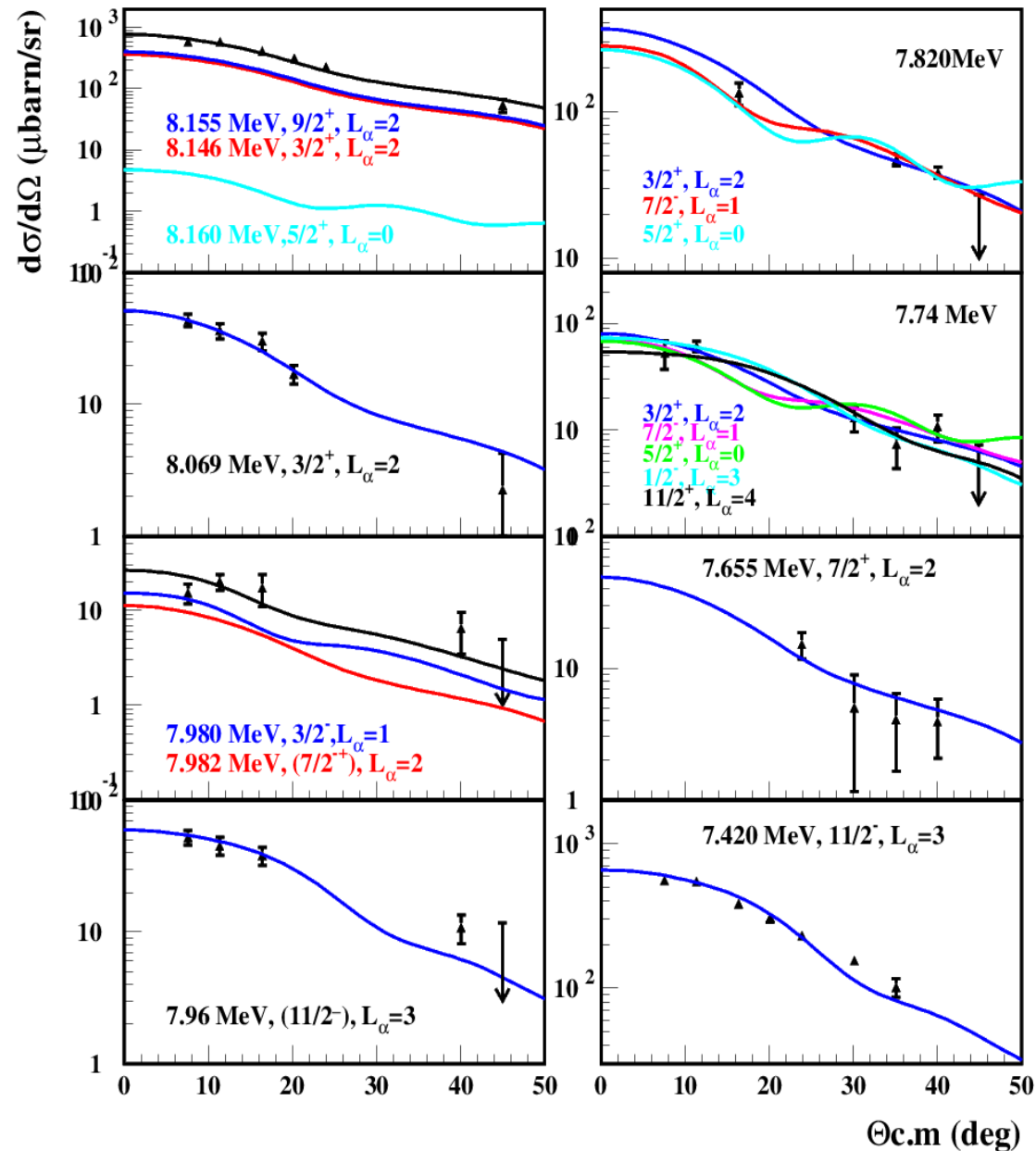


Region of interest



Experimental energy resolution (FWHM) : ~
30 keV (6°) - 71 keV (36°)

FR-DWBA calculations



- From the **normalisation** of the DWBA (Distorted Wave Born Approximation) calculations (FRESCO code) to the data
→ **Spectroscopic factor**

$$\left. \frac{d\sigma}{d\Omega} \right|_{\text{exp}} = C^2 S'_\alpha S_\alpha \left. \frac{d\sigma}{d\Omega} \right|_{\text{FR-DWBA}}$$

$$S'_\alpha = \langle {}^7\text{Li} | t \otimes \alpha \rangle = 1 \quad \text{Kubo et al PRC 1978)}$$

$$S_\alpha = \langle {}^{21}\text{Ne}^* | {}^{17}\text{O} \otimes \alpha \rangle$$

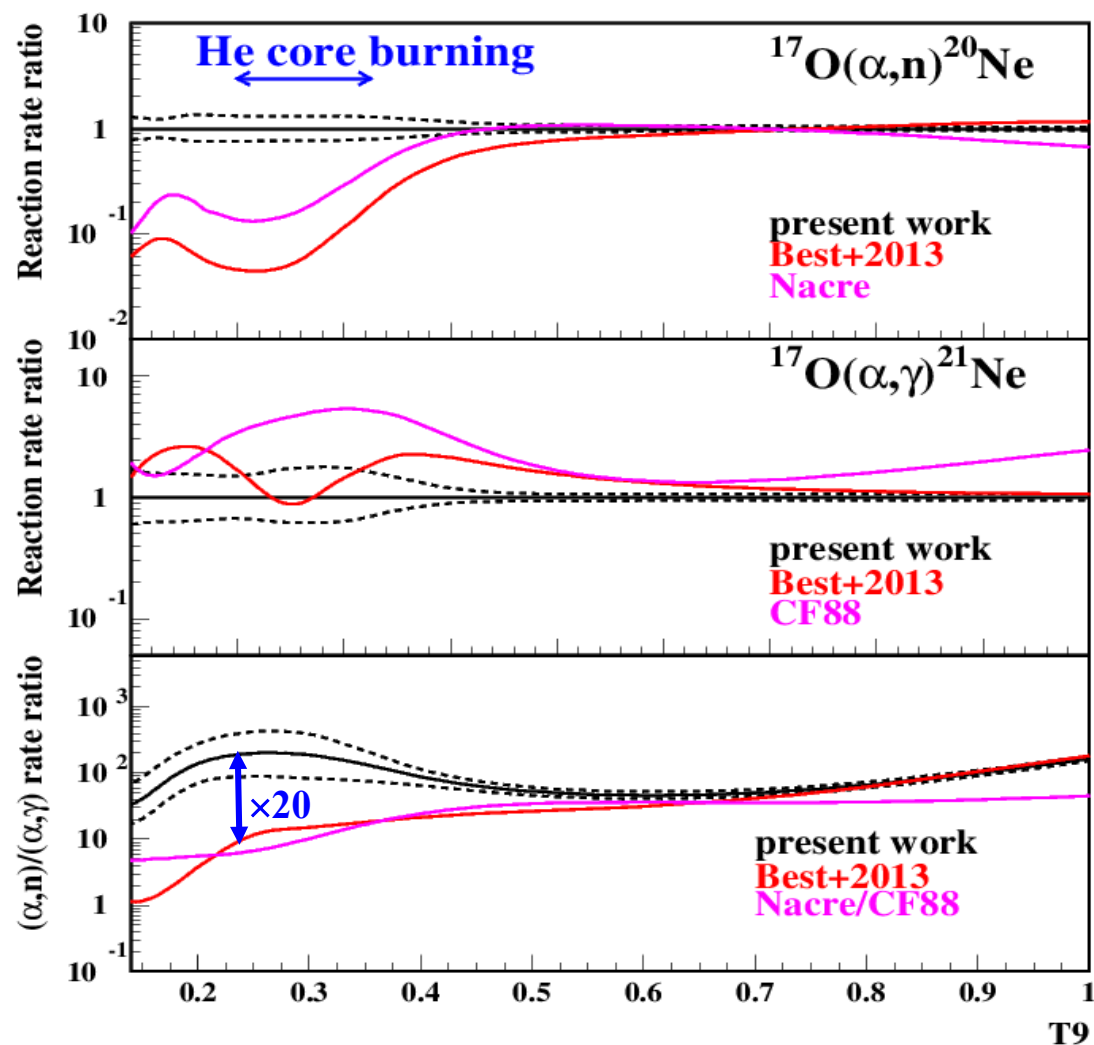
$$S_\alpha \rightarrow \Gamma_\alpha = 2P_l \frac{\hbar^2 R}{2\mu} S_\alpha |\phi(R)|^2$$

P_l = Coulomb & centrifugal penetrability factor
 $\phi(R)$ radial part of the $\alpha + {}^{17}\text{O}$ wave function

• Γ_α uncertainty: 3- 40% (stat), **35%** (optical potentials)

$^{17}\text{O}(\alpha, n)$ & $^{17}\text{O}(\alpha, \gamma)$ new reaction rates & impact on the s-process in low Z massive stars

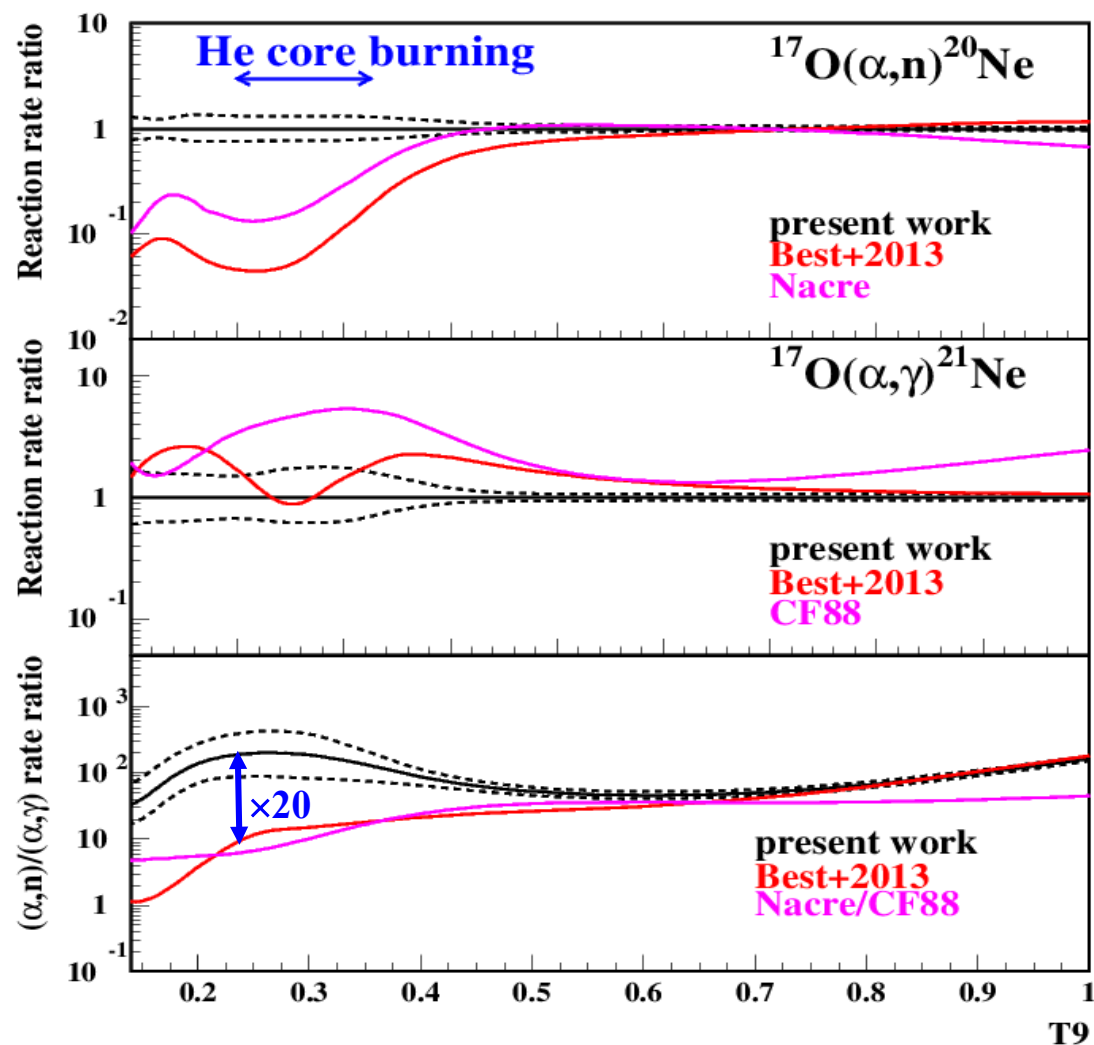
Reaction rates calculations:



→ Better neutron efficiency recycling with a factor of about **20** with the **present rates** than **Best+2013** rates

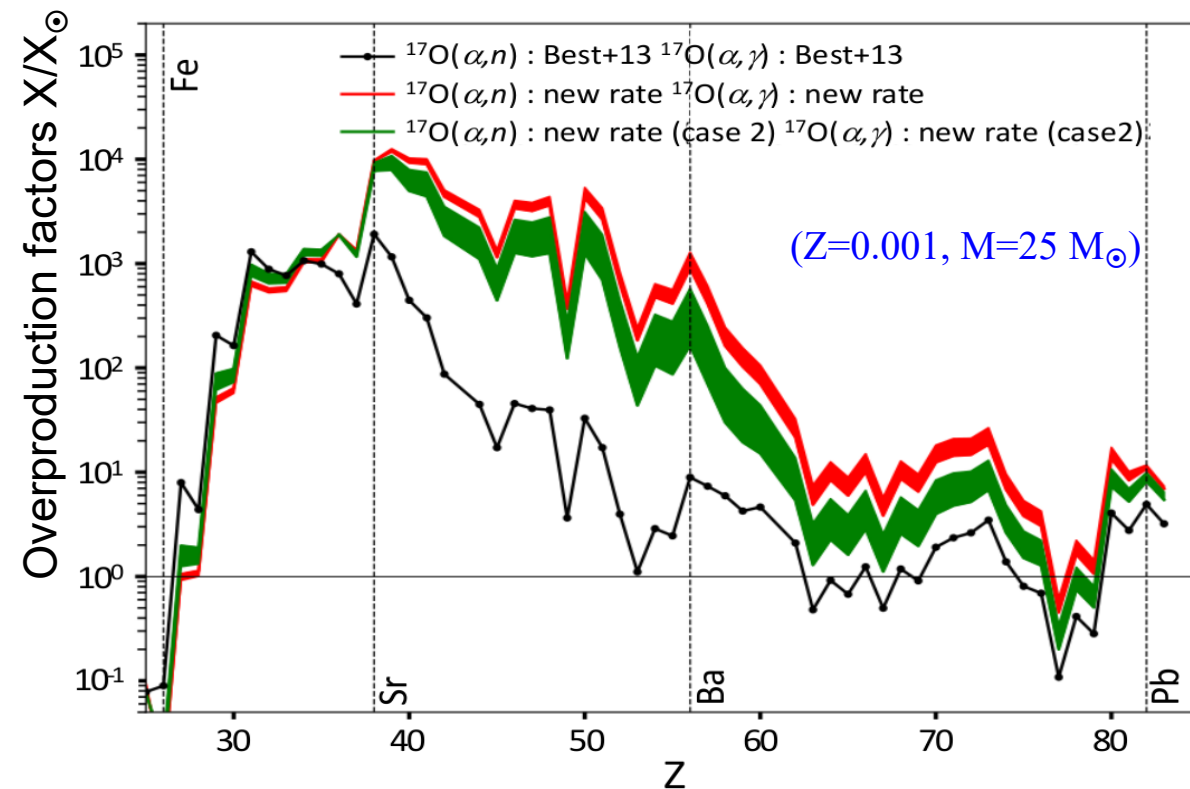
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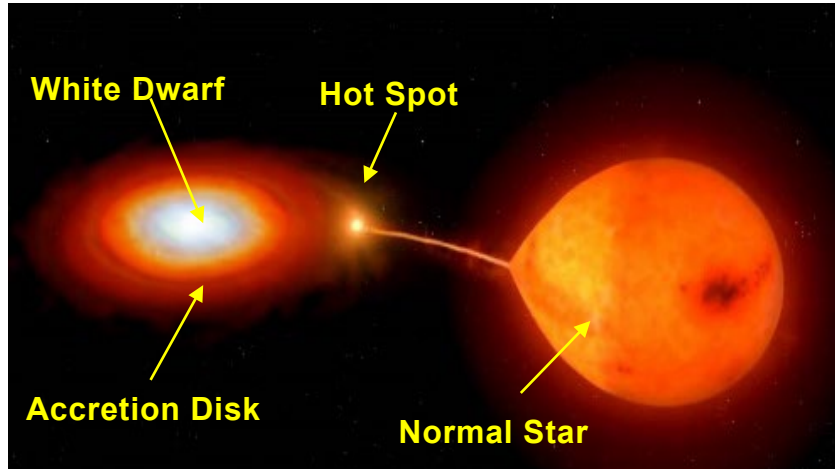
→ Better neutron efficiency recycling with a factor of about 20 with the present rates than Best+2013 rates

Nucleosynthesis calculations:



F. H, P. Adsley, L. Lamia et al., Phys Rev Lett 132, 182701 (2024)

- Large enhancement ($> \text{factor } 30$ (> 20)) of elements $40 < Z < 60$ with the present new rates in comparison to previous rates
- Two order of magnitude ($\sim \times 30$ (case2)) on Barium : largest effect
 - In line with the observation of an enhanced s-process in CEMP-s (Carbon Enhanced Metal poor) stars Hansen+2016



Novae: Thermonuclear runaway at the surface of a white dwarf accreting the H rich material of its companion star

Mechanism well established but

- Ejected mass $\ll$ observed mass
- Amount of mixing between accreted and white dwarf materials?
- Compact star mass

→ Novae produce γ -ray emitters & dust (presolar grains) → key observational constraints

→ soon the first stellar explosive site with all reaction rates based on experimental information J.José+2009

The 3 remaining key reactions with significant uncertainties:

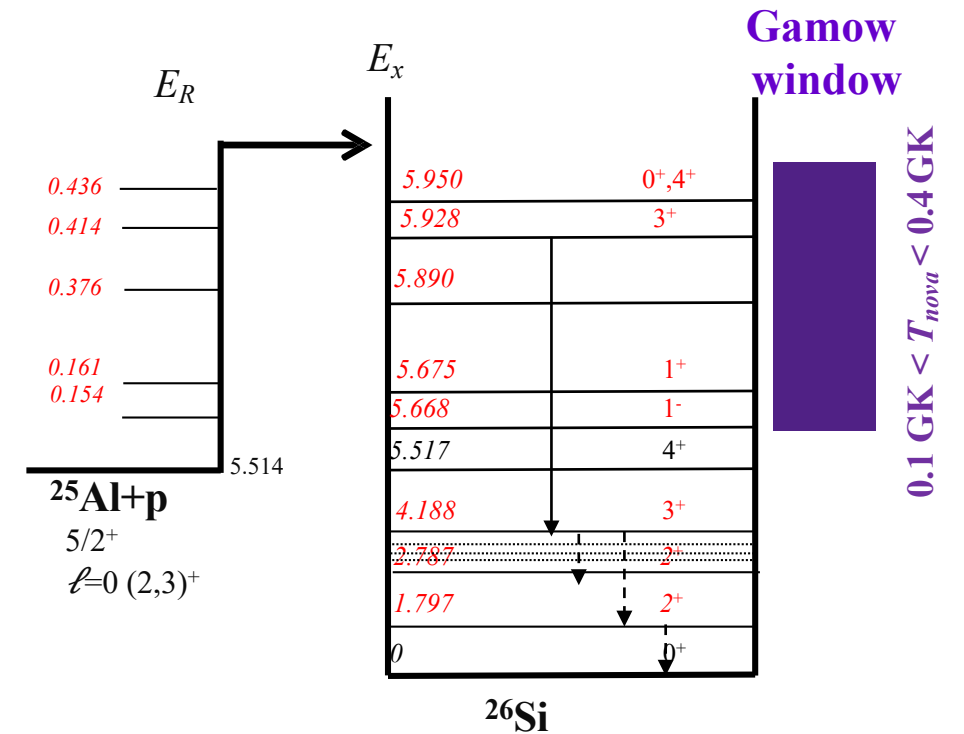
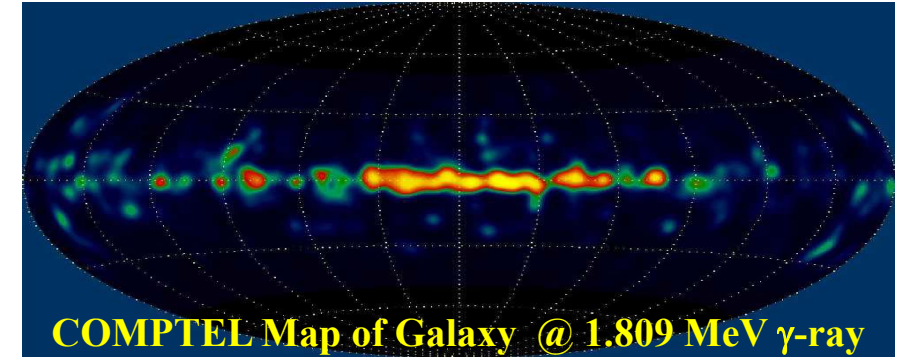
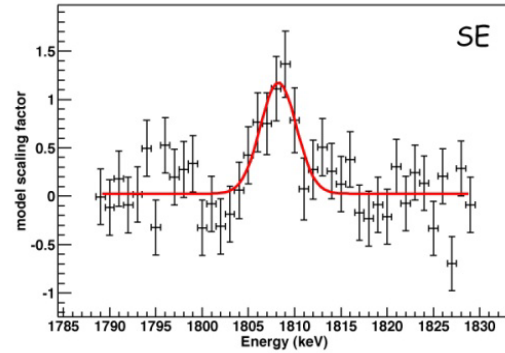
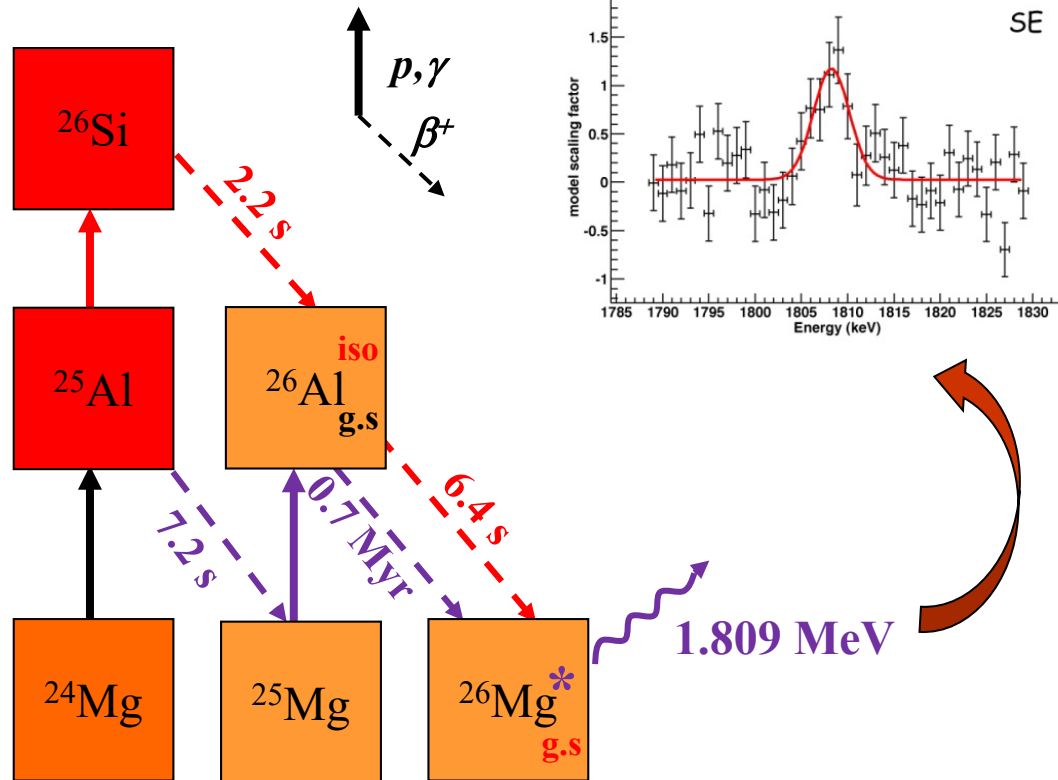
- $^{30}\text{P}(p,\gamma)^{31}\text{S}$: paternity of novae grains Woods+2021? (FRIB)
- $^{18}\text{F}(p,\alpha)^{15}\text{O}$: γ -ray emission ≤ 511 keV Riley+PRC2021 (Split-Pole/ALTO)
- $^{25}\text{Al}(p,\gamma)^{26}\text{Si}$: Classical novae contribution to the observed galactic ^{26}Al γ -ray emission @ 1.809 MeV?

^{26}Al nucleosynthesis

➤ **Novae contribution to ^{26}Al galactic production 10 – 30 % ?**

➤ Two sequences to produce ^{26}Al ($T_{1/2}=7.2 \cdot 10^5$ yr):

$^{25}\text{Al}(p,\gamma)^{26}\text{Si}$ competes with $^{25}\text{Al}(\beta^+)$



$^{25}\text{Al}(p,\gamma)^{26}\text{Si}$: $\langle \sigma v \rangle_{\text{tot}} = \sum_R \left(\frac{2\pi}{\mu(^{25}\text{Al}, p) k_B T} \right)^{\frac{3}{2}} \hbar^2 \omega \gamma_R \exp \left(-\frac{E_R}{k_B T} \right)$

Factor two uncertainty on resonance strengths
⇒ need to be measured

$^{25}\text{Al}(p,\gamma)^{26}\text{Si}$ case:

Experiment @ FRIB

→ Studied via $d(^{25}\text{Al},n\gamma)^{26}\text{Si}$ @ FRIB/S800 C.Fougères, F.H, N de Séréville, F. de Oliveira+coll

➤ Novel technique: $(d,n\gamma)$ angle-integrated XS measurement (inspired from Kankainen+ (2016) EPJ; $d(^{26}\text{Al},n\gamma)^{27}\text{Si}$)

$$\omega\gamma = \frac{N_\gamma}{N_{\text{target}} I_{\text{beam}} \epsilon_{\text{det.}}^{\text{tot.}} T_{\text{UT}}} \times \frac{\Gamma_p^{\text{s.p.}}}{\sigma_{\text{transfer}}^{\text{DWBA}}} \times \frac{(2J+1)}{(2j+1)(2J_{^{25}\text{Al}}+1)}.$$

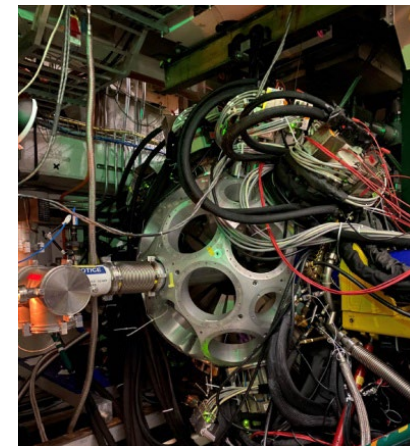
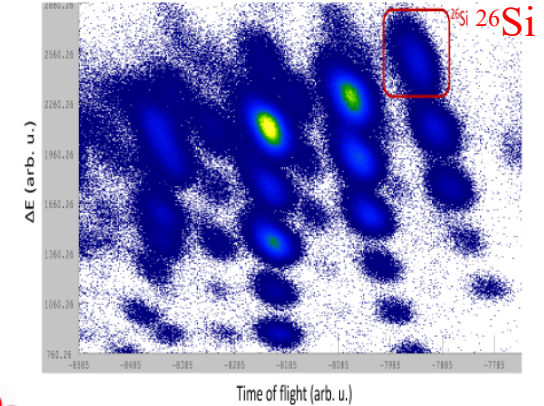
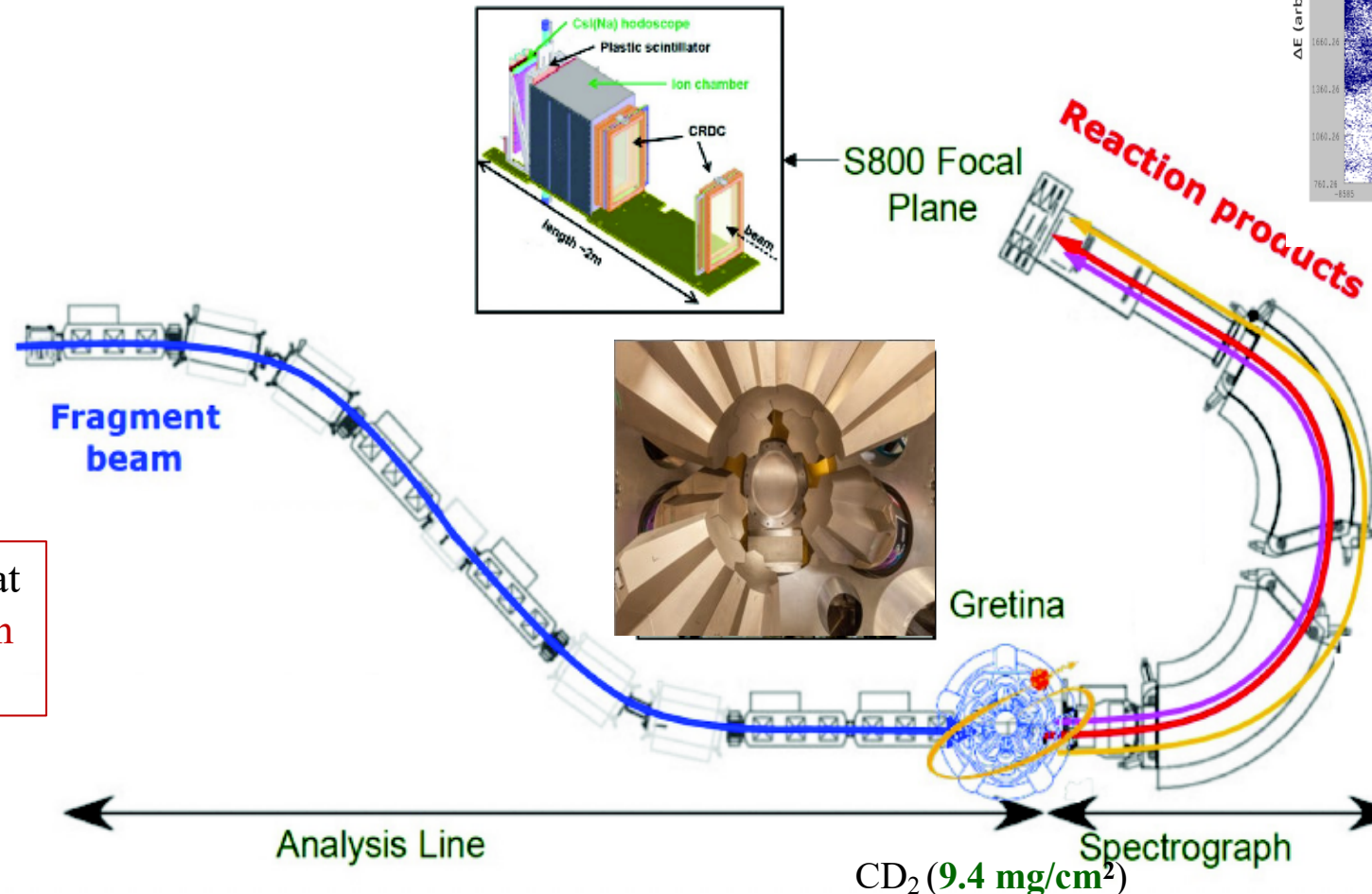
GRETINA&S800@FRIB/ARIS

^{25}Al @ 24 MeV/u

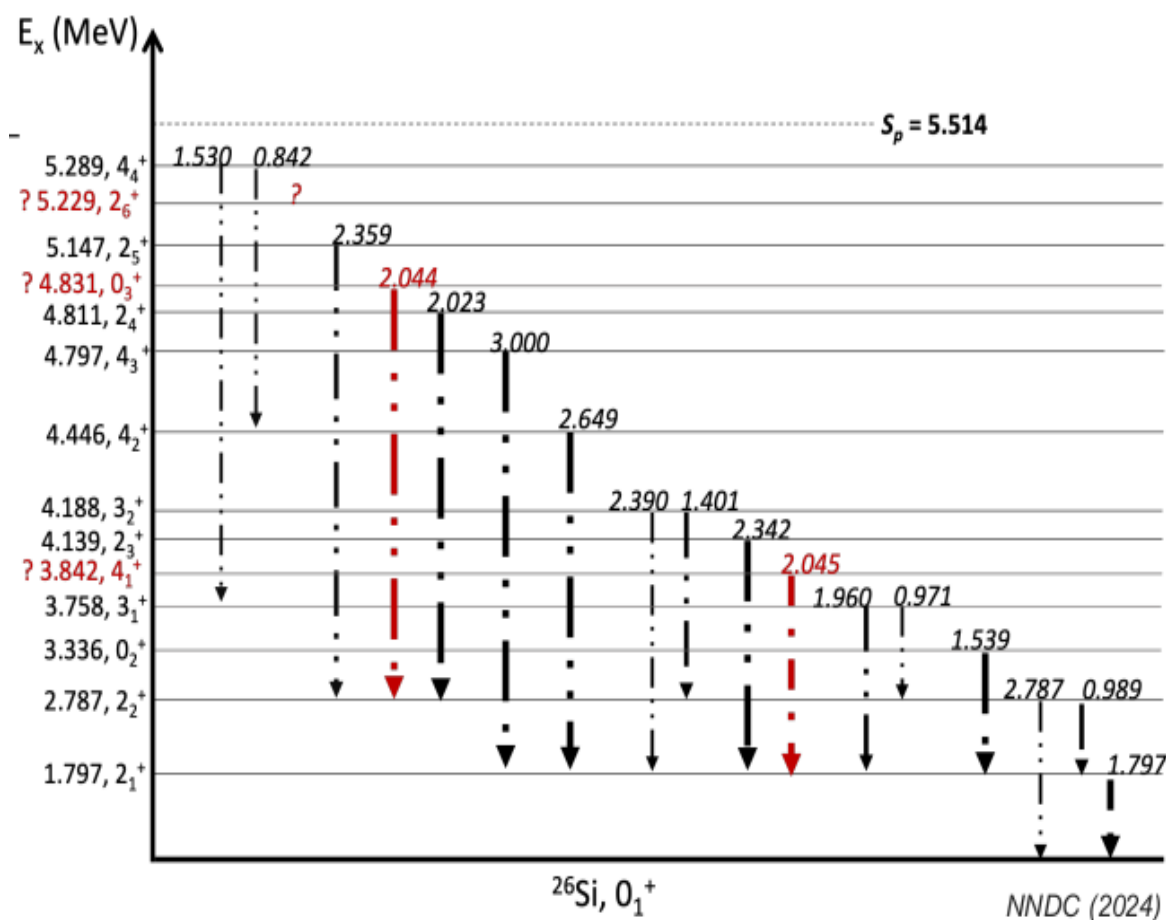
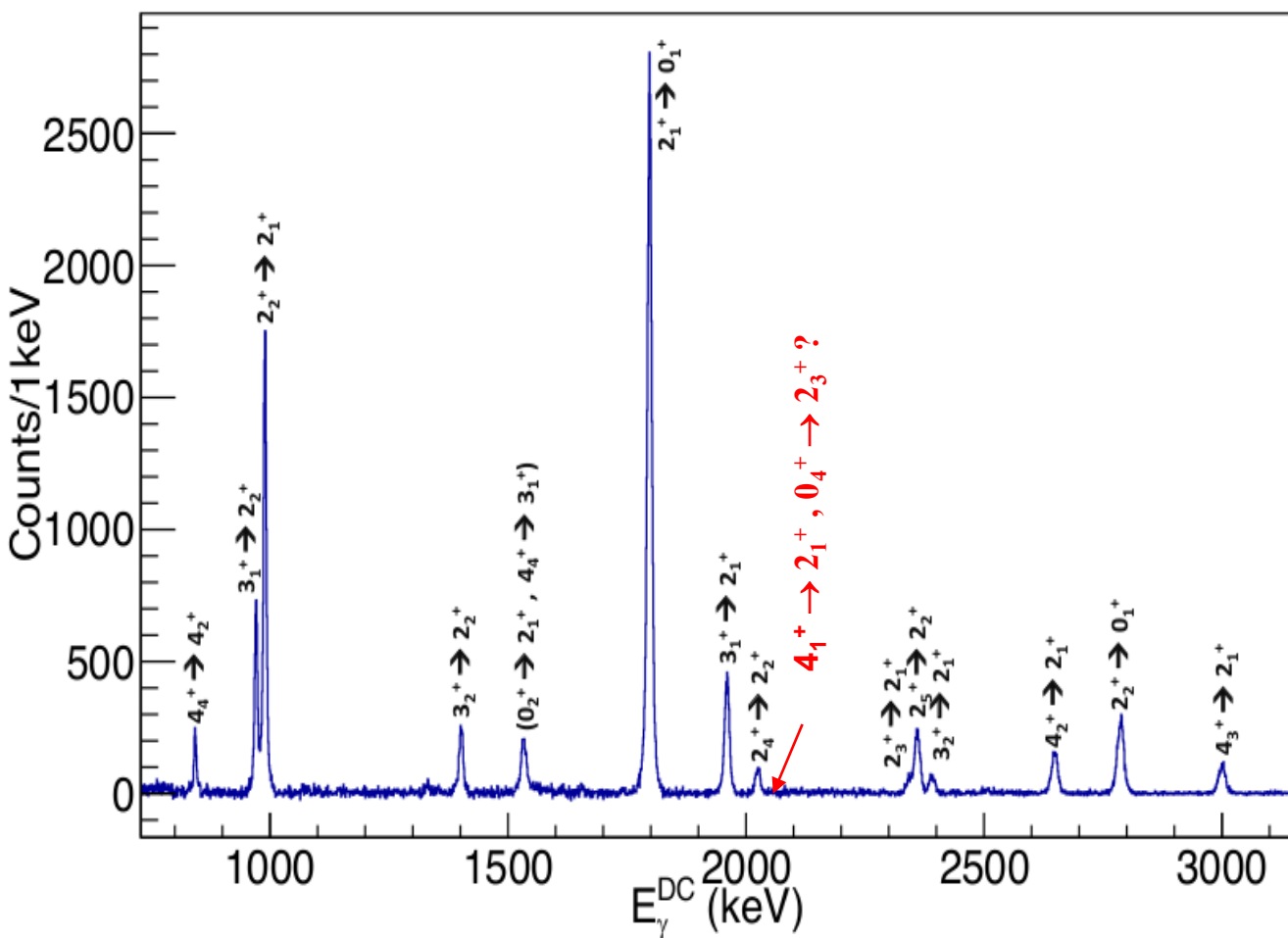
lowest energy ever used @ S800)

2×10^6 pps, Purity >95%

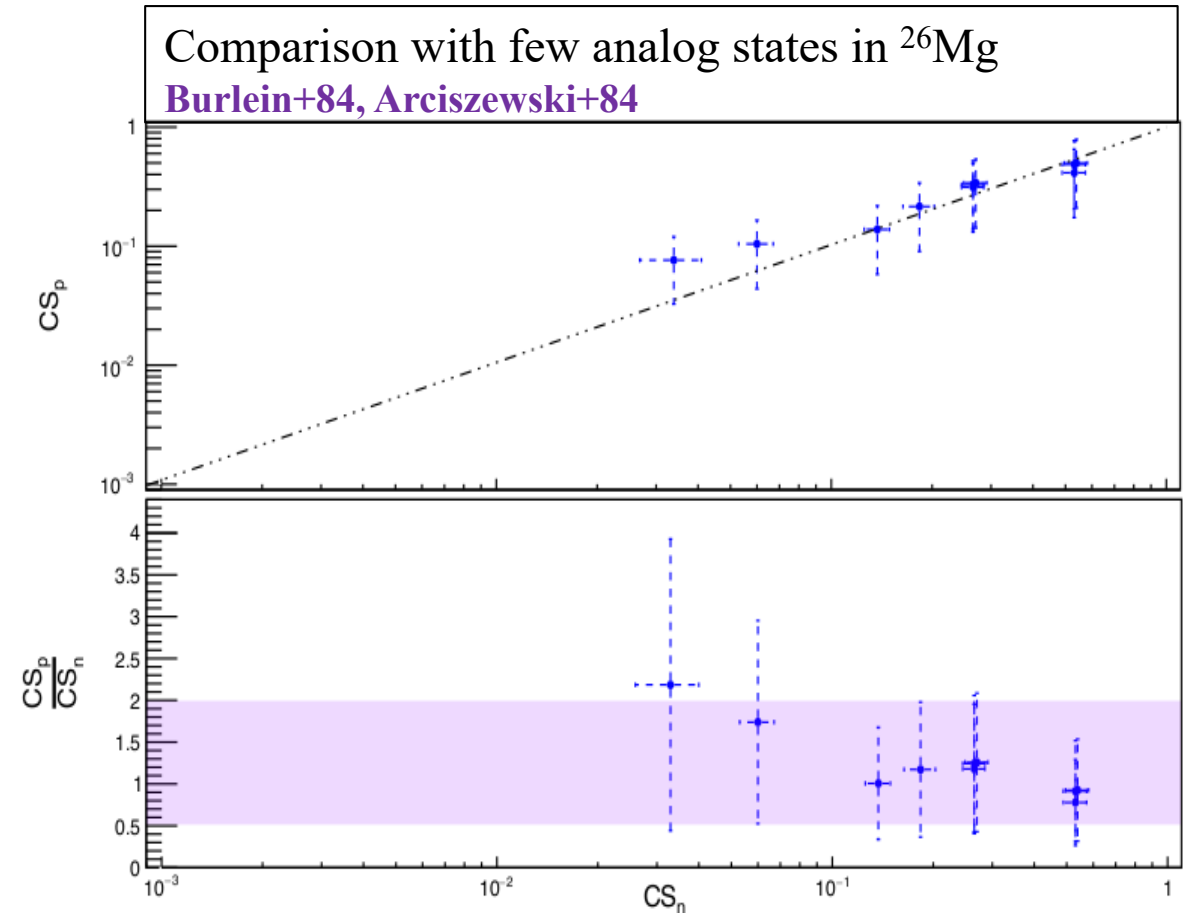
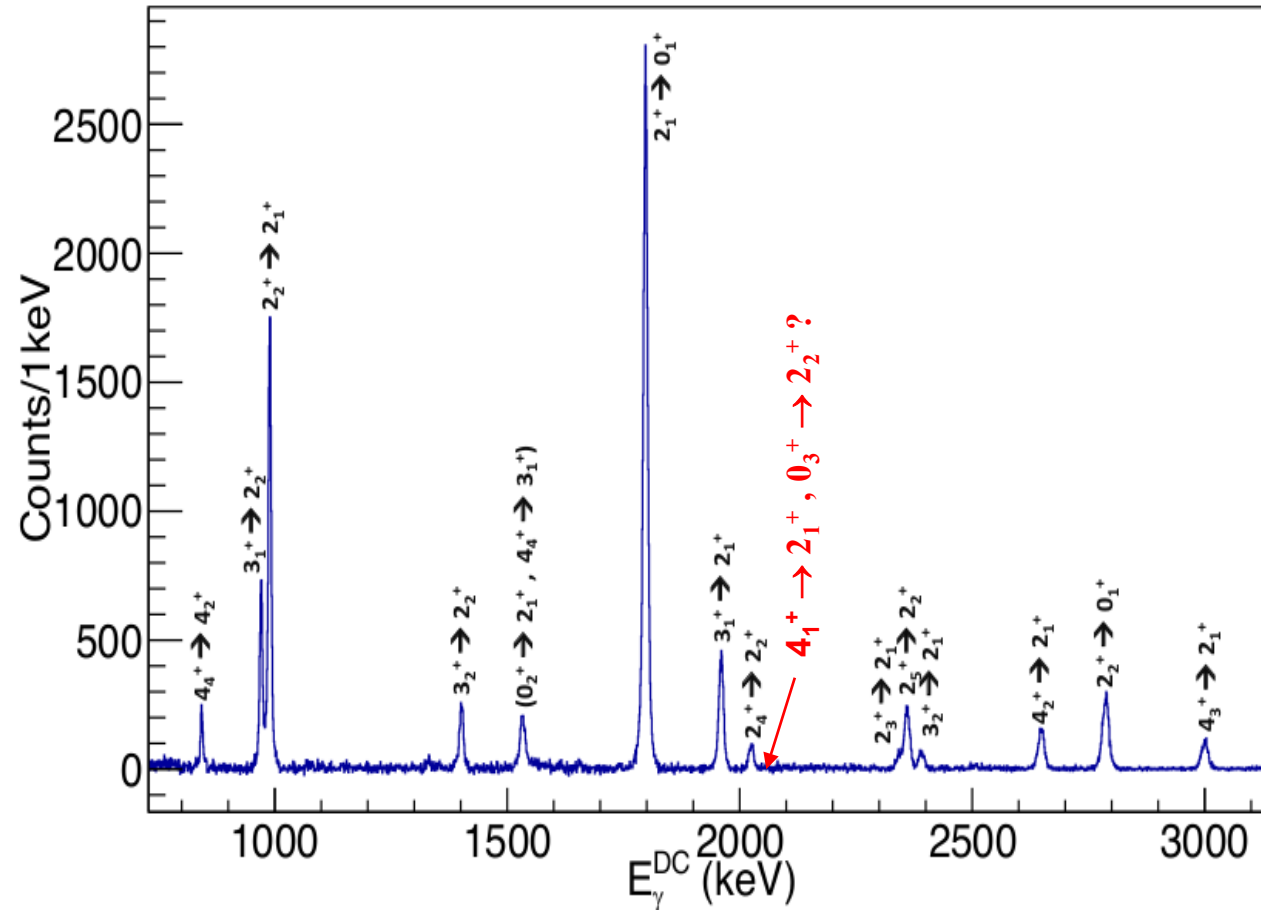
Measurements of ^{26}Si recoils at
S800 in coincidence with γ 's in
GRETINA



➤ γ -ray transitions of 11 states in ^{26}Si below S_p were observed among the 14 referenced

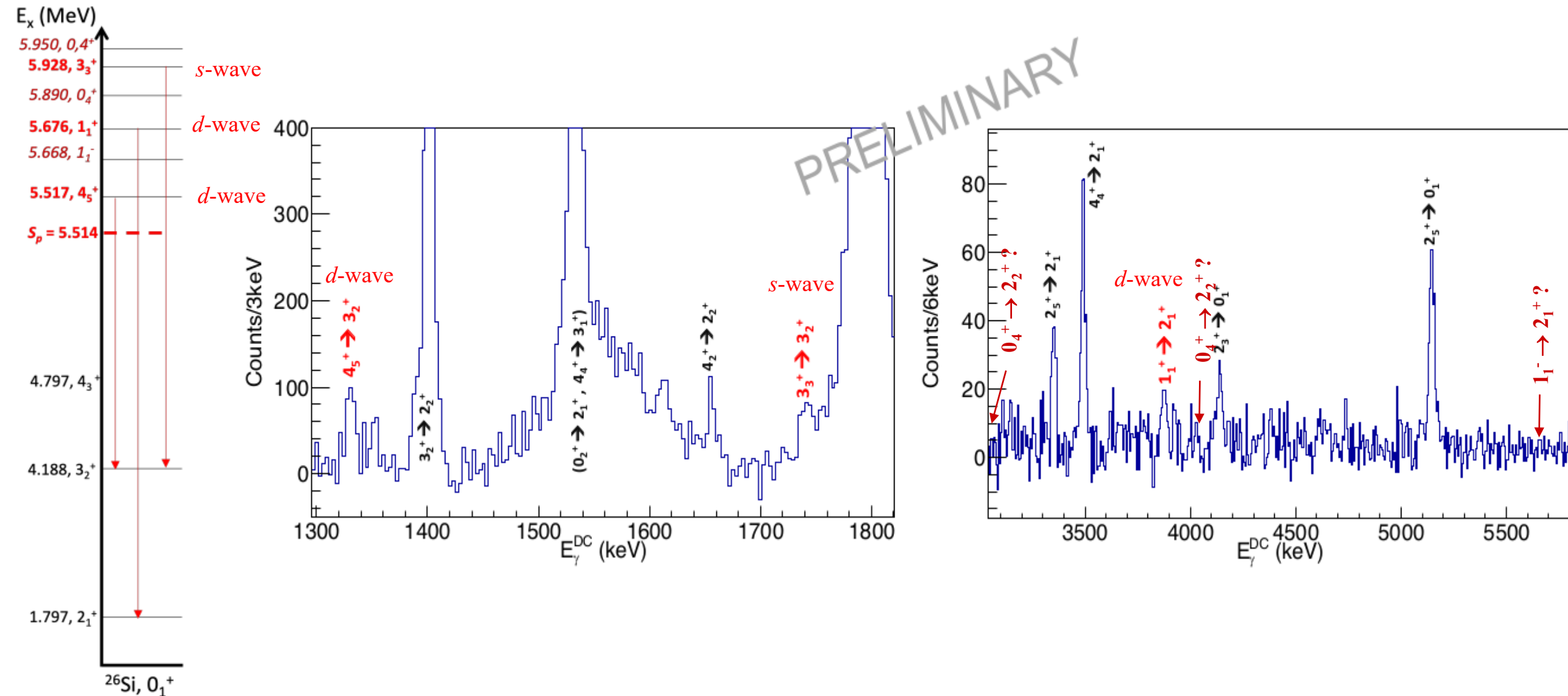


➤ $C^2\text{Sp}$ of the bound states in ^{26}Si determined for the 1st time $\Rightarrow$ direct capture rate component of $^{25}\text{Al}(p,\gamma)^{26}\text{Si}$

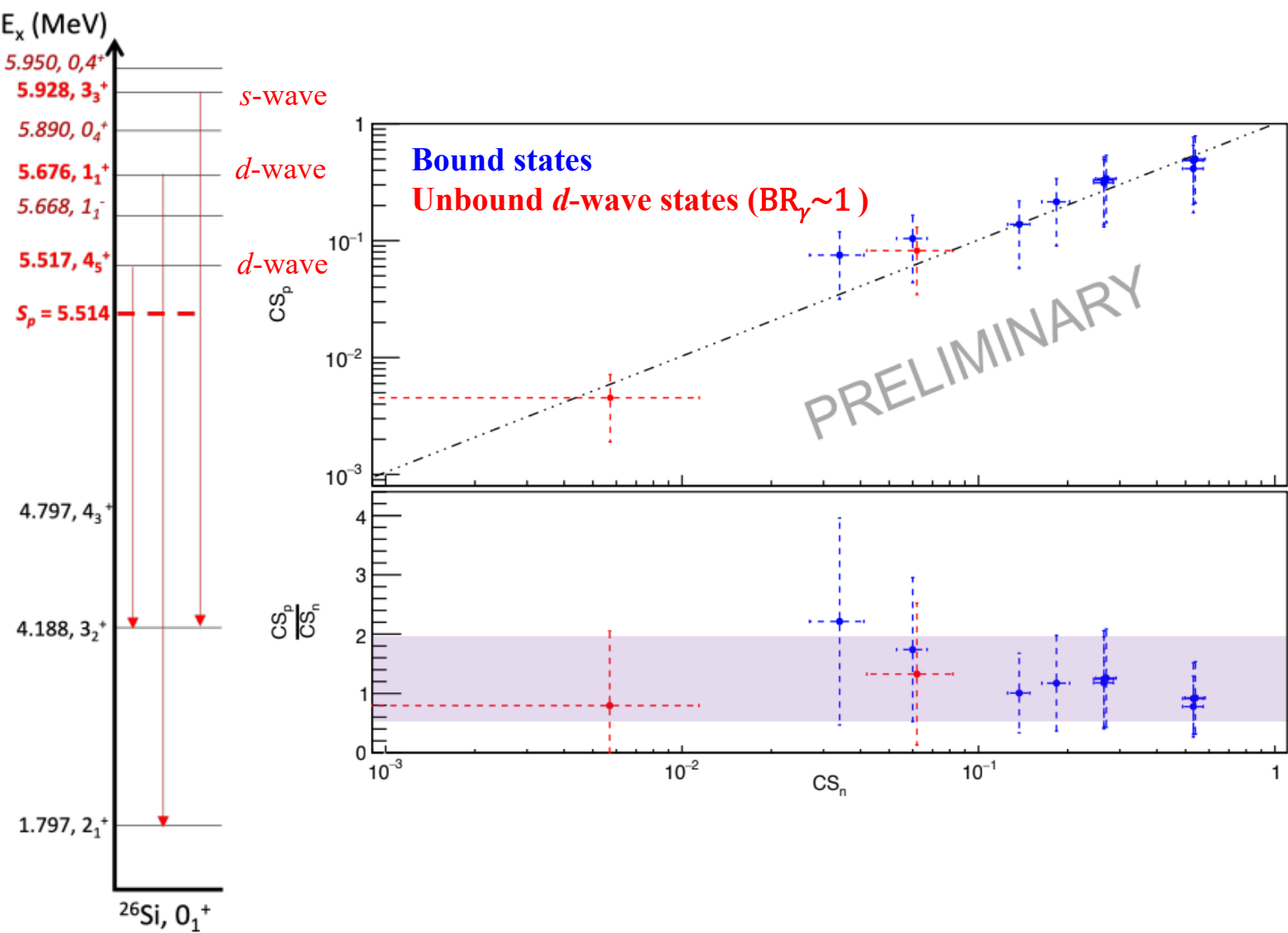


$\rightarrow$ new data on ^{26}Mg : ALTO/SPILT-POLE experiment (Orsay, 2024)

➤ Three resonant states among six of interest are identified in ^{26}Si ($>S_p$)

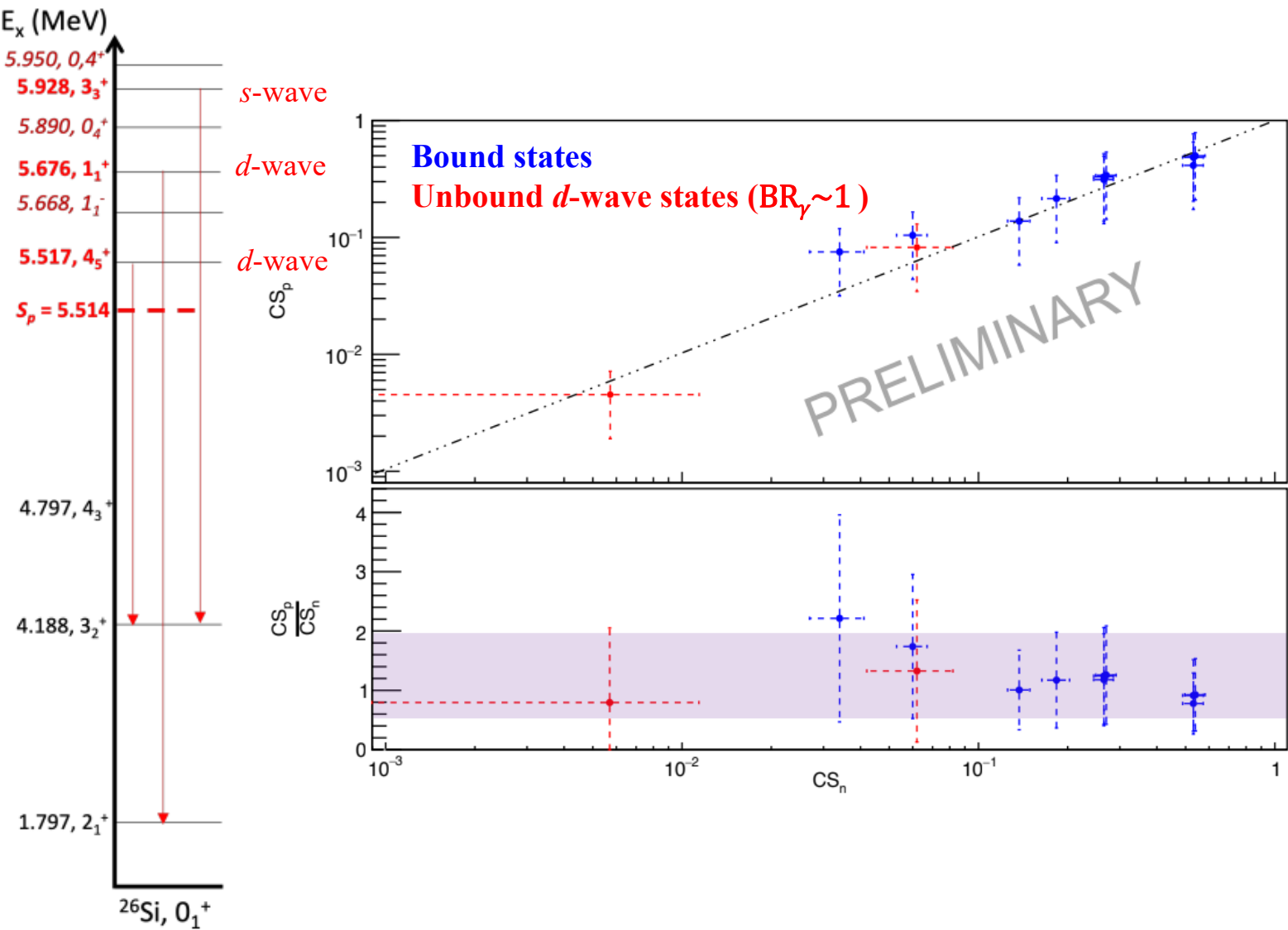


➤ Three resonant states among six of interest are identified in ^{26}Si ($>S_p$)



	Resonance strength (meV)	
	1_1^+	3_3^+
Laird et al (2023)	$< 2.6 \times 10^{-6}$	23(13)
Present work (preliminary)	$2(1) \times 10^{-6}$	35(15)

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Laird et al (2023)	$< 2.6 \times 10^{-6}$	23(13)
Present work (preliminary)	$2(1) \times 10^{-6}$	35(15)

Ongoing analysis

Next:

- Re-evaluation of $^{25}\text{Al}(p,\gamma)^{26}\text{Si}$ rate
- Constrain novae contribution to ^{26}Al galactic amount

Thanks to the collaboration & your attention

PHYSICAL REVIEW LETTERS **132**, 182701 (2024)

Experimental Determination of α Widths of ^{21}Ne Levels in the Region of Astrophysical Interest: New $^{17}\text{O} + \alpha$ Reaction Rates and Impact on the Weak s Process

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Angle-integrated measurement of the $d(^{25}\text{Al}, n\gamma)^{26}\text{Si}$ transfer reaction to probe resonance strengths in $^{25}\text{Al}(p, \gamma)^{26}\text{Si}$ relevant for the production of ^{26}Al in novae FRIB experiment with GREY - S800

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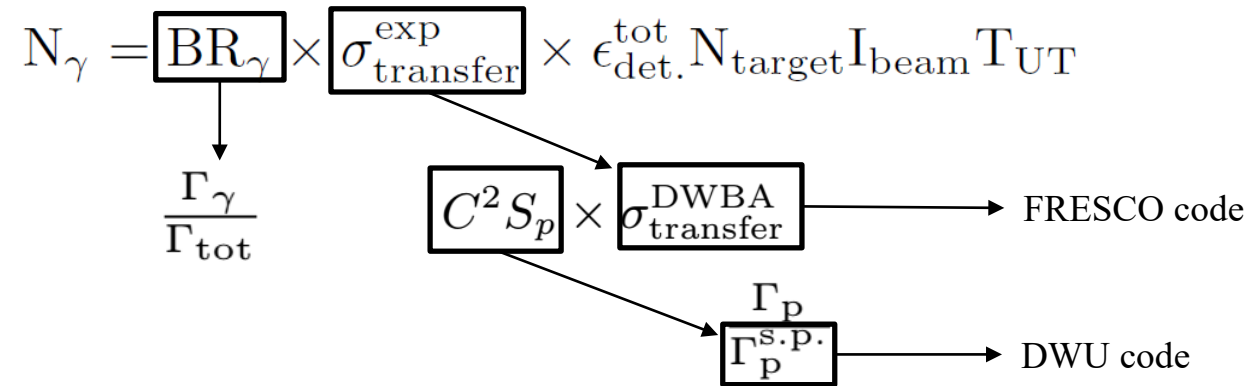
Accessing $\omega\gamma$ via angle-integrated measurement

$$\omega\gamma = \frac{(2J+1)}{(2j+1)(2J_{^{25}\text{Al}}+1)} \frac{\Gamma_p \Gamma_\gamma}{\Gamma_{\text{tot}}}$$

Extension of recent method to measure C^2S_p (NSCL) via $d(^{26}\text{Al}, n\gamma)^{27}\text{Si}$ *Kankainen, EPJ 52 (2016)*

Direct transfer $d(^{25}\text{Al}, n\gamma)^{26}\text{Si}$

Tagging of $^{26}\text{Si}^*$ on γ -ray transitions $\rightarrow$ angle-integrated measurement of cross-section σ_{transfer}



$$\omega\gamma = \frac{N_\gamma}{N_{\text{target}} I_{\text{beam}} \epsilon_{\text{det.}}^{\text{tot.}} T_{\text{UT}}} \times \frac{\Gamma_p^{\text{s.p.}}}{\sigma_{\text{transfer}}^{\text{DWBA}}} \times \frac{(2J+1)}{(2j+1)(2J_{^{25}\text{Al}}+1)}$$

Uncertainties = systematic $\lesssim 30\%$ (from optical potentials) & statistical ($1/\sqrt{N_\gamma}$)