Probing Beyond the Standard Model with Beta Decay and Electron Capture

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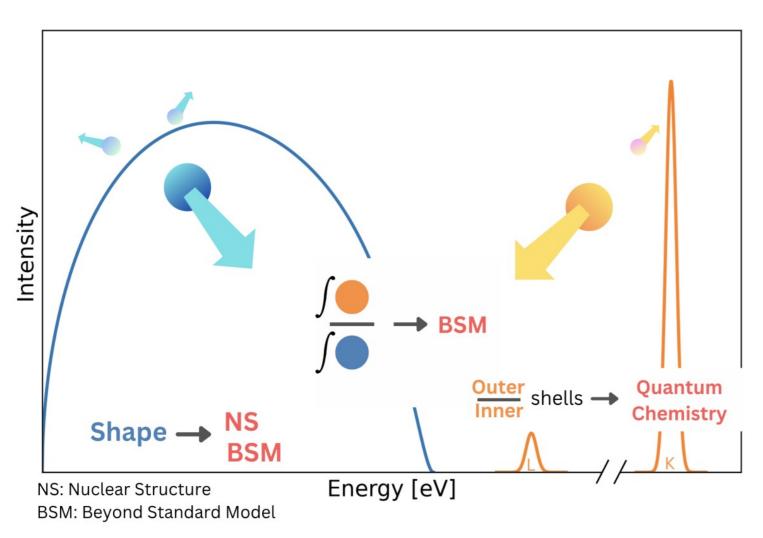
Motivation

Nuclear beta decay and electron capture provide powerful tools to test the Standard Model (SM) and explore new physics. In particular, beta decay can be offers sensitivity to exotic scalar ε_S and tensor ε_T currents at the TeV scale.

This sensitivity appears in the **Fierz interference term**, b_F , which depends linearly on these interactions. Beta-plus decay and electron capture (EC) influence b_F in opposite ways but probe the same nuclear matrix element.

This means the **electron capture to beta-plus branching ratio** is an extremely sensitive probe of new physics with minimal nuclear structure effects and able to provide accuracy of $\sim 10^{-3}$ on b_F

At the experimental level, ASGARD/SALER perform high-precision recoil spectroscopy using Superconducting Tunnel Junctions (STJ) detectors with eV accuracy.



Methodology

EC/β+ ratio shows a dependence of opposite sign on the Fierz interference term while probing the same nuclear matrix element, making it extremely sensitive to new physics.

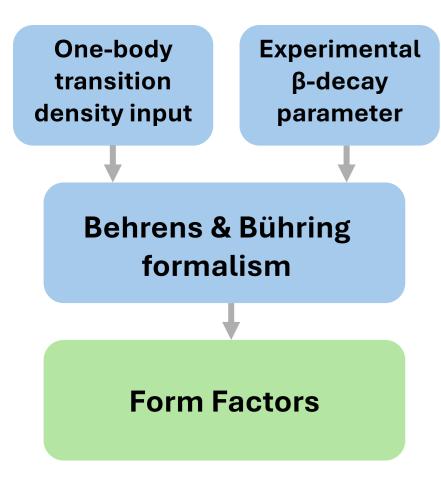
$$\frac{\lambda_{\rm EC}}{\lambda_{\beta^+}} = \frac{f_x}{f_\beta} (1 + \delta_{\rm NS}) \cdot \frac{1 + b_F \frac{m_e}{E_x}}{1 - b_F \left\langle \frac{m_e}{E_e} \right\rangle}$$
Nuclear structure

Nuclear Structures (NS) are calculated using the shell model calculation and H.Behrens & W.Bühring formalism [1] based on a multipole expansion.

$$< f | H_0^{\mathrm{hadr}} L_{\mathrm{lep}}^0 | i > \propto \sum_{L,M} j_L(qR) Y_M^L(\hat{q}) F_L(q^2)$$
 Form factor

The advantages of this decomposition allows all observables to be expressed in terms of form factors for all types of transitions.

The aim of my work was to create a code that would reproduce the calculation of form factors and link them to different observables.



Possible outputs:

- $> \lambda_{EC}/\lambda_{\beta+}$ ratio
- \triangleright Shape factor $\mathcal{C}(E)$
- > β-spectrum, half-life
- Angular parameters
- > Testing the CVC hypothesis
- $\triangleright g_A$ quenching

Context

Effective field theory is a way to interpreting Beyond Standard Model (BSM) physics at scale $\Lambda_{BSM} \gg LHC$.

$$\mathcal{L}_{\text{eff}} \propto \overline{e} \gamma_{\mu} \nu_{L} \cdot \overline{u} \gamma^{\mu} [C_{V} - (C_{A} - 2\epsilon_{R}) \gamma^{5}] d + \epsilon_{S} \overline{e} \nu_{L} \cdot \overline{u} d$$

$$- \epsilon_{P} \overline{e} \nu_{L} \cdot \overline{u} \gamma^{5} d + \epsilon_{T} \overline{e} \sigma_{\mu\nu} \nu_{L} \cdot \overline{u} \sigma^{\mu\nu} (1 - \gamma^{5}) d \quad \text{at quark level}$$

V-A structure in SM.

 ϵ_i are proportional to $(M_W/\Lambda_{BSM})^2$ \rightarrow $\epsilon_i \leq 10^{-4}$ \rightarrow $\Lambda_{BSM} \geq 15 \ TeV$ assuming natural couplings.

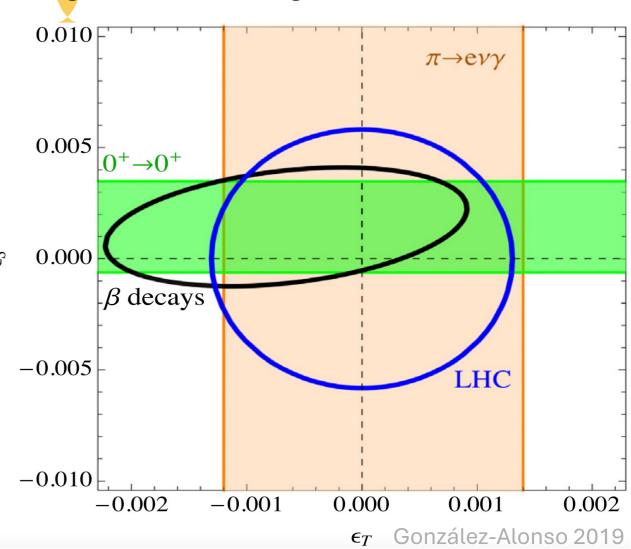
Fierz interference term $m{b}_F$ comes from an interference between the SM and BSM: $\mathcal{M}_{SM} \times \mathcal{M}_{BSM} \approx \mathcal{O}(\epsilon_S, \epsilon_T)$.

$$b_F = \pm \frac{2\gamma}{1 + |\rho|^2} \operatorname{Re}\left(\frac{\epsilon_S g_S}{g_V} + |\rho|^2 \frac{8g_T \epsilon_T}{-2g_A}\right)$$

This term allows the SM to be probed at low energies. b_F depends linearly on the scalar and tensor exotic currents $(\epsilon_S = \epsilon_T = 0 \text{ in SM}).$

There are currently several ... methods for constraining the value of scalar and tensor couplings.

With recent results, β-decay is competitive with the LHC in the search for exotic currents.



Results

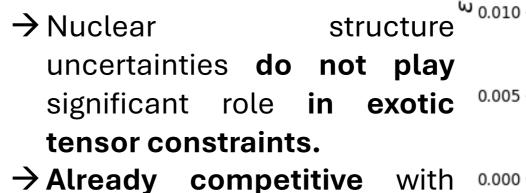
NS calcu<mark>la</mark>tions for 2 data sets: X.Mougeot 2019 [3] & W.Bambynek 1977 [2].

Isotope	$ ho_{exp}$	$ ho_{th}^{I}$	$ ho_{th}^{II}$	$\mid \delta_{NS}^{exp}(\%)$	$\delta^I_{NS}(\%)$	$\delta_{NS}^{II}(\%)$	$\mid (\lambda_{EC}/\lambda_{eta^+})_{exp}$	$(\lambda_{EC}/\lambda_{eta^+})_{th}$
$\overline{\ ^{11}\mathrm{C}^{a,c}}$	-0.7544(8)	-0.7728	-0.6927	0.23	0.30	0.26	0.00225(15)	0.00211(6)
$^{22}\mathrm{Na}^{b,c}$	_	-	-	-2.0	-1.08	0.32	0.1083(9)	0.1097(9)
$^{57}\mathrm{Ni}^d$	_	-	-	-	0.03	-	1.460(47)	1.447(11)
$^{58}\mathrm{Co}^d$	_	-	-	-	0.03	-	5.61(8)	5.63(8)
$^{65}{ m Zn}^e$	-	-	-	-	0.37	-	30.1(5)	29.6(7)
$13\mathrm{N}^{a,c}$	-0.560(1)	-0.5480	-0.5077	0.01	0.03	0.01	0.00168(12)	0.00176(5)
$^{15}\mathrm{O}^{a,c}$	0.630(2)	0.5555	0.5185	0.01	0.02	-0.02	0.00107(6)	0.00091(2)
$^{17}\mathrm{F}^{b,c}$	1.296(1)	1.1572	1.1376	0.48	0.66	0.64	_	0.00135(2)
$^{18}\mathrm{F}^{b,c}$	_ ` `	-	-	-	0.61	0.59	0.030(18)	0.0310(5)
$^{19}\mathrm{Ne}^{b,c}$	-1.602(1)	-1.6338	-1.6197	0.50	0.70	0.69	0.00096(3)	0.00093(1)

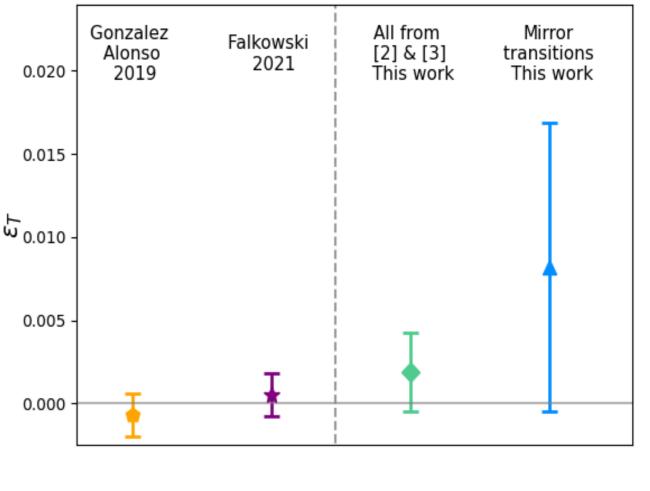
NS corrections are sensitives to the interaction used (I or II).

	$CKPOT^a$	USDB^b	$YSOX^c$	HO^d	$JUN45^e$
Model Space	$1p_{3/2}, 1p_{3/2}$	$1s_{1/2}, 1d_{5/2}$	$1s_{1/2}, 1p_{3/2}, 1p_{1/2}$	N: $1f_{5/2}$, $2p_{3/2}$, $2p_{1/2}$	$1f_{7/2}, 1g_{9/2}$
		$1d_{3/2}$	$1d_{5/2}, 1d_{3/2}$	P: $1f_{7/2}$	$2p_{3/2}, 2p_{3/2}$

For mirror transitions (MT), the **hadronic part is easy** for reasons of symmetry and are "easily" accessible for ab initio calculations.



global dataset.



 $arepsilon_T^{MT} = 0.0082(82)_{exp}(30)_{atom}(8)_{NS} \, (90\%$ CL) $arepsilon_T^{All} = 0.0019(24) \, (90\%$ CL)

Conclusion & Outlook

With nuclear structure corrections at the $\sim 10^{-3}$ level, uncertainties are an order of magnitude lower than typical beta decay searches, and we are already competitive with the global dataset in the search for exotic tensor current.

Experiments aiming to probe the EC/ β + ratio at the 0.1% level are a powerful probe of physics at the tens of TeV scale with minimal nuclear structure uncertainties.

Outlook: ab initio calculations to appropriately estimate nuclear structure uncertainties in a controlled fashion & add condensed matter theory to the Behrens & Bühring formalism to take account of superconducting effects.

[1] : W.Bühring & H.Behrens, Electron Radial Wave Function and Nuclear Beta Decay, 1982.

[2] : W. Bambynek et al. "Orbital electron capture by the nucleus", Reviews of Modern Physics 49.1 (Jan. 1, 1977).

[3]: X. Mougeot, Towards high-precision calculation of electron capture decays, Applied Radiation and Isotopes, 2019.