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on behalf of



UNIVERSITÀ

Outline



- Quarkonium Production
- A Large Ion Collider Experiment
 - Experimental apparatus
 - Data-taking mode in Run 3
- Results
 - Charmonium production
 - \circ $\psi(2S)$ -to-J/ ψ cross-section ratio
 - Charmonium polarization
 - \circ J/ ψ -D⁰ associated production
- Conclusions and Outlook











Quarkonium =

bound state of a heavy quark antiquark pair

charmonium: charm + anticharm

bottomonium: bottom + antibottom



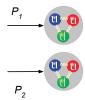
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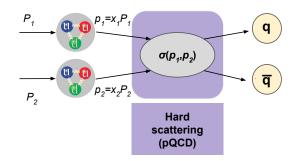
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PRODUCTION
MECHANISM IN
pp COLLISIONS



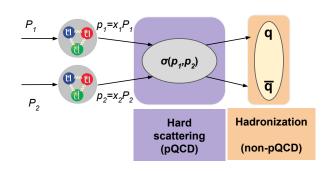
• Hard scattering to produce high-mass quark pair \rightarrow small $\alpha_s \rightarrow$ perturbative QCD calculation possible



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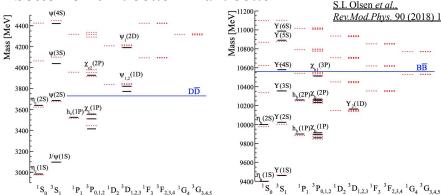
bound state of a heavy quark antiquark pair

PRODUCTION MECHANISM IN pp COLLISIONS



charmonium: charm + anticharm

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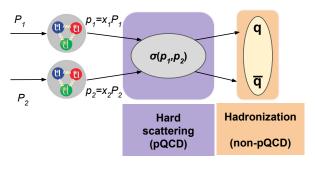


- Hard scattering to produce high-mass quark pair \rightarrow small α_S \rightarrow perturbative QCD calculation possible
- Binding into quarkonium states \rightarrow larger $\alpha_s \rightarrow$ non-perturbative QCD process



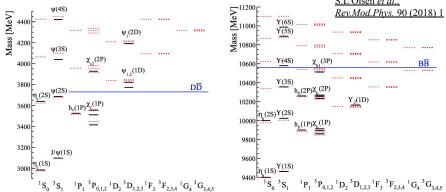
Ouarkonium = bound state of a heavy quark antiquark pair

PRODUCTION MECHANISM IN pp COLLISIONS



charmonium: charm + anticharm

bottomonium: bottom + antibottom



- Hard scattering to produce high-mass quark pair \rightarrow small $\alpha_s \rightarrow$ perturbative QCD calculation possible
- Binding into quarkonium states \rightarrow larger $\alpha_s \rightarrow$ non-perturbative QCD process

The interplay between perturbative and non-perturbative QCD is described by phenomenological models

→ need experimental measurements to constrain the models

pp measurements are also a baseline for studying quark-gluon plasma properties from heavy-ion collisions (see next talk) 7

Phenomenological models for quarkonium production in pp



Color Singlet Model (CSM)

The quarkonium is directly produced in a color-singlet state.

Non-Relativistic QCD (NRQCD)

The quarkonium can be produced in a singlet or octet state, with soft gluon emission to reach the final state.

NRQCD + Color Glass Condensate

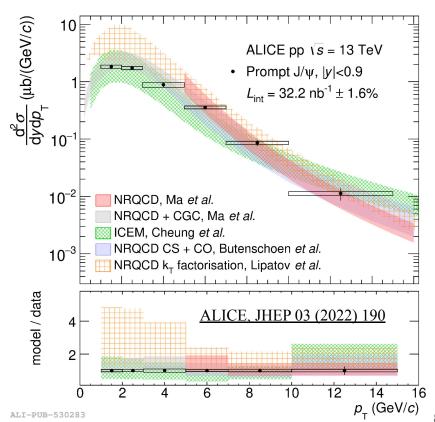
Improvement in the agreement with data at low $p_{\rm T}$ taking into account the gluon saturation regime.

Improved Color Evaporation Model (ICEM)

The quark-antiquark pair evolves into a quarkonium state with a constant probability factor (energy-, momentum- and process-independent).

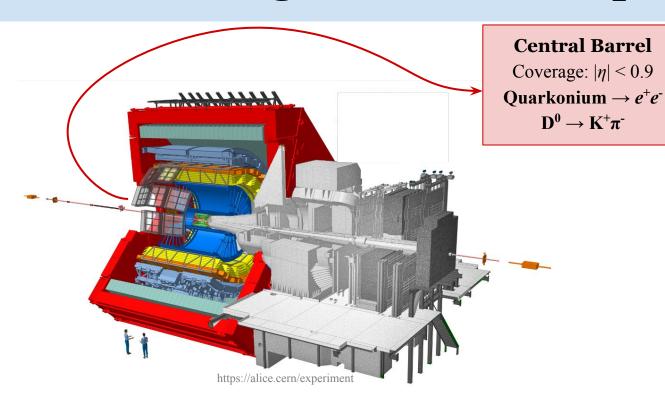
NRQCD with k_T factorization

It also takes into account the transverse momentum of gluons.



A Large Ion Collider Experiment





Inner Tracking System and **Time Projection Chamber**

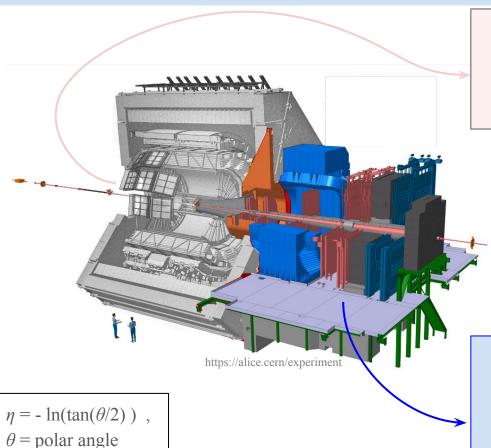
→ particle tracking and identification

Time Of Flight

 \rightarrow particle identification

A Large Ion Collider Experiment





Central Barrel

Coverage: $|\eta| < 0.9$ Quarkonium $\rightarrow e^+e^-$

 $D^0 \rightarrow K^+\pi^-$

 $MCH+MID \Rightarrow inclusive$ charmonium measurements (presented in the next slides)

+ Muon Forward Tracker (MFT) new in Run 3 for prompt/non-prompt separation (see dedicated talk)

Forward Muon Spectrometer

Coverage: $2.5 < \eta < 4.0$

Quarkonium $\rightarrow \mu^+\mu^-$

Muon CHambers

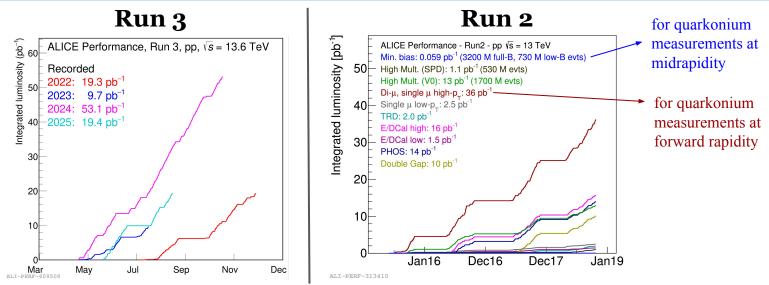
 $\rightarrow \mu^{\pm}$ tracking

Muon IDentifier

 $\rightarrow \mu^{\pm}$ identification

Inclusive quarkonium production: from hardware to software trigger



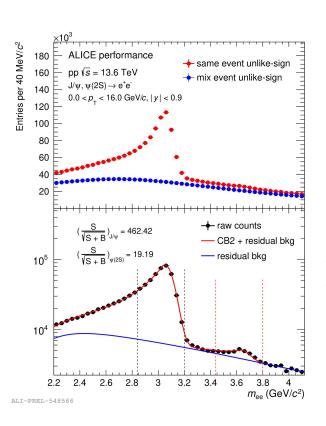


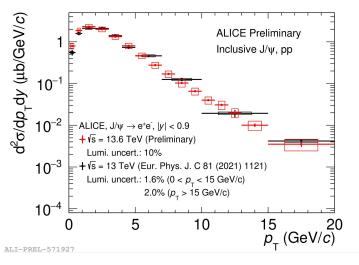
New data-taking mode using **continuous readout + software trigger** (vs Run 2: hardware triggers)

- \rightarrow Possibility of measuring $\psi(2S)$ at midrapidity down to $p_T = 0$
- → Larger data samples → increased statistical precision of the measurements at mid and forward rapidity

Charmonium production at midrapidity



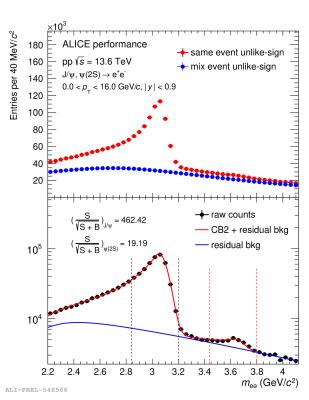


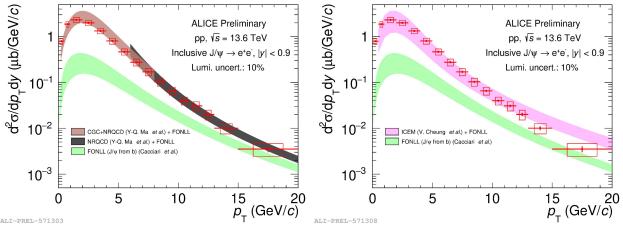


- J/ ψ and ψ (2S) signal extraction performed in differential $p_{\rm T}$ ranges
- J/ ψ Run 3 preliminary cross section at $\sqrt{s} = 13.6$ TeV at midrapidity is **compatible with Run 2** publication at $\sqrt{s} = 13$ TeV

Charmonium production at midrapidity



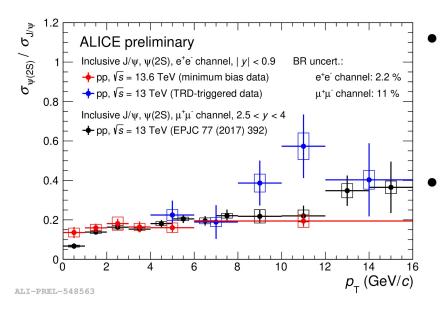




- J/ ψ and ψ (2S) signal extraction performed in differential p_T ranges
- J/ ψ Run 3 preliminary cross-section at $\sqrt{s} = 13.6$ TeV at midrapidity is compatible with Run 2 publication at $\sqrt{s} = 13$ TeV
- Good agreement with NRQCD at high p_T and with NRQCD+CGC at low p_T
- Good agreement with ICEM as well
 + FONLL computation of non-prompt J/ψ production

Charmonium production ratio at midrapidity



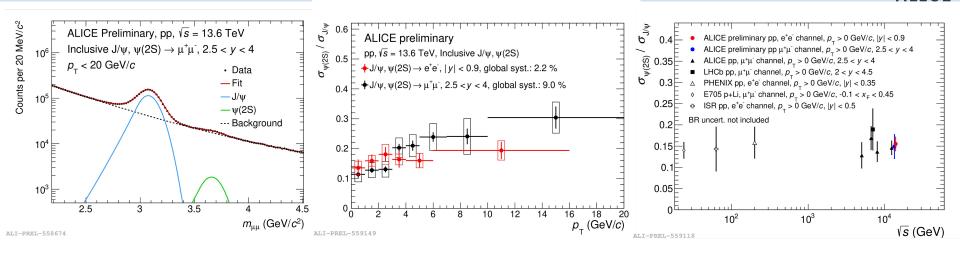


- Interesting to look at the $\psi(2S)$ -to-J/ ψ cross-section ratio
 - It provides access to the fraction of $\psi(2S)$ from J/ψ decay (feed-down factor)
 - Systematic uncertainties are partially cancelled in the ratio
- Midrapidity ratios vs p_T at $\sqrt{s} = 13.6$ TeV and $\sqrt{s} = 13$ TeV \rightarrow both in agreement with forward-rapidity publication at $\sqrt{s} = 13$ TeV

$$f^{\psi(2S)} = \frac{\sigma_{\psi(2S)}}{\sigma_{J/\psi}} \cdot BR_{\psi(2S) \to J/\psi}$$

Charmonium production ratio at forward rapidity

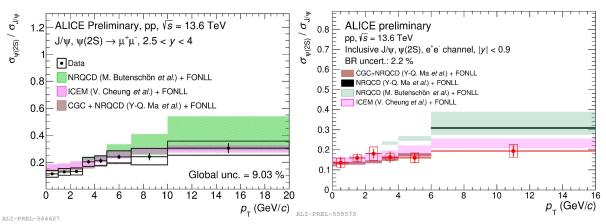




- Preliminary $\psi(2S)$ -to-J/ ψ cross-section ratio extraction vs p_T at forward rapidity at $\sqrt{s} = 13.6$ TeV is consistent with the corresponding measurement at midrapidity
- Integrated result over p_T and y compared to previous measurements \rightarrow no energy dependence

ψ(2S)-to-J/ψ ratio: comparison with phenomenological models





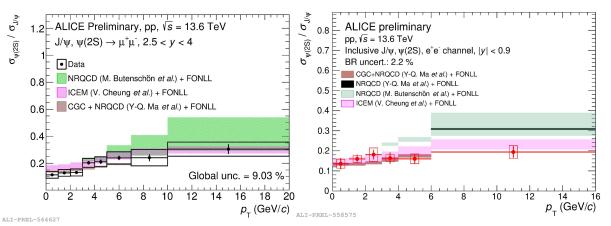
Preliminary ratios vs p_{T} are compared to phenomenological models

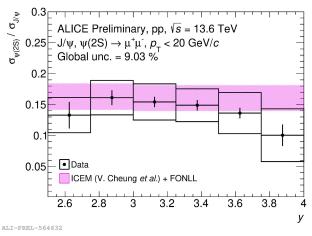
- Both NRQCD and ICEM can describe the ratio at forward rapidity
- At midrapidty, NRQCD overestimates the measurements

Increasing the sample size will allow for better investigation

ψ(2S)-to-J/ψ ratio: comparison with phenomenological models







Preliminary ratios vs p_T are compared to phenomenological models

- Both NRQCD and ICEM can describe the ratio at forward rapidity
- At midrapidty, NRQCD overestimates the measurements

Increasing the sample size will allow for better investigation

Preliminary ratios vs rapidity at forward rapidity are well reproduced by ICEM → **no significant** *y* **dependence**

An improved measurement using the whole data sample will reduce the uncertainties on the experimental points and allow for a better constraint

Polarization of quarkonium states



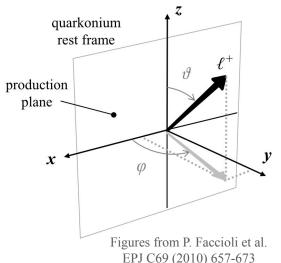
Polarization \rightarrow alignment of particle spin with respect to a quantization axis

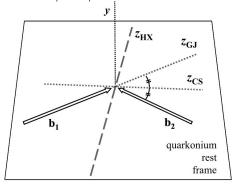
Two-body decay \rightarrow the angular distribution of the decay products reflects the polarization of the quarkonium state

Observable angular distribution:

$$W(\cos\theta,\varphi) \propto \frac{1}{3+\lambda_{\theta}} \cdot (1+\lambda_{\theta}\cos^2\theta + \lambda_{\varphi}\sin^2\theta\cos2\varphi + \lambda_{\theta\varphi}\sin2\theta\cos\varphi)$$

$$(\lambda_{\theta'}, \lambda_{\varphi'}, \lambda_{\theta\varphi}) = (+1, 0, 0) \Rightarrow$$
 Pure transverse polarization $(\lambda_{\theta'}, \lambda_{\varphi'}, \lambda_{\theta\varphi}) = (-1, 0, 0) \Rightarrow$ Pure longitudinal polarization $(\lambda_{\theta'}, \lambda_{\varphi'}, \lambda_{\theta\varphi}) = (0, 0, 0) \Rightarrow$ Zero polarization



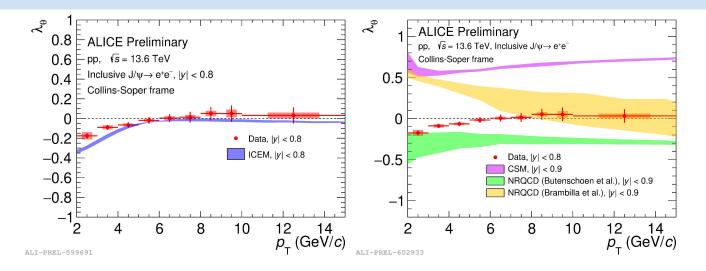


Polarization parameters depend on the reference frame (choice of polarization axis)

- **Helicity frame**: quarkonium direction in the collision center-of-mass frame
- Collins-Soper frame: bisector of the angle of one beam and the opposite of the other beam in the quarkonium rest frame

Charmonium polarization





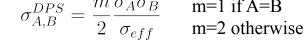
- Huge discrepancy among the models (different models can even predict opposite polarizations)
 ⇒ experimental measurements bring essential constraints to their predictions
- Sizeable J/ ψ longitudinal polarization (negative λ_{θ}) in the Collins-Soper frame at low $p_{\rm T}$ in the dielectron channel at $\sqrt{s} = 13.6 \, {\rm TeV} \rightarrow {\rm qualitatively}$ in agreement with ICEM predictions
- Brambilla's NRQCD \longrightarrow does not reproduce the measured polarization at low $p_{\rm T}$
- Butenschoen's NRQCD \longrightarrow only reproduces the lowest measured $p_{\rm T}$ range
- CSM \longrightarrow fails to reproduce polarization in the whole momentum range

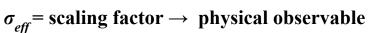
Associated production of charm mesons



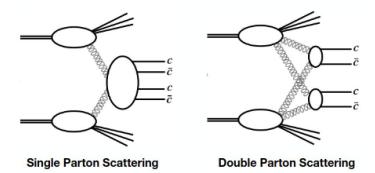
- New insight on the production mechanisms and the role of **Multiple Parton Interactions (MPI) in pp collisions**
- Investigate relative contributions of single parton scattering (SPS) and double parton scattering (DPS)
- Inputs to tune MC generators

$$\sigma_{A,B}^{DPS} = \frac{m}{2} \frac{\sigma_A \sigma_B}{\sigma_{eff}}$$
 m=1 if A=B m=2 otherwise





→ Independent on process and collision energy in the simplest model

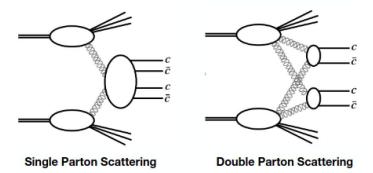


Associated production of charm mesons



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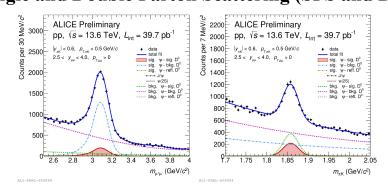
 σ_{eff} = scaling factor \rightarrow physical observable \rightarrow Independent on process and collision energy in the simplest model

J/ψ-D⁰ associated production: contribution from both Single and Double Parton Scattering (SPS and DPS)

 \rightarrow possibility to extract $\sigma_{_{\it eff}}$

Analysis strategy:

- $J/\psi \rightarrow \mu^+\mu^-$ at fwd. rapidity $D^0 \rightarrow K^-\pi^+$ at mid rapidity $1.9 < \Delta y < 4.6$
- Yield extraction using a bidimentional invariant-mass fit



J/ψ-D⁰ associated production



Integrated cross section in the Δy range vs predictions from different MC generators:

PYTHIA 8

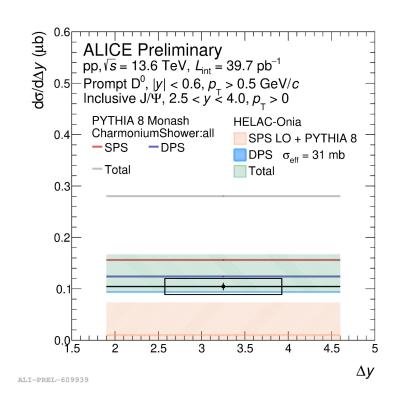
<u>T. Sjöstrand et al., Comput. Phys. Commun. 191 (2015), 159-177</u> (Monash Tune + Charmonium Shower:All)

- both SPS and DPS cross sections computed from the simulation
- overestimates J/ψ -D⁰ cross section

• Helac-Onia

H.-S. Shao, Comput. Phys. Commun. 184 (2013) 2562-2570

- SPS cross section computed from the simulation
- DPS component is based on $\sigma_{\it eff}$ value which is extracted from a fit to the data $\to \sigma_{\it eff}$ = 31 mb

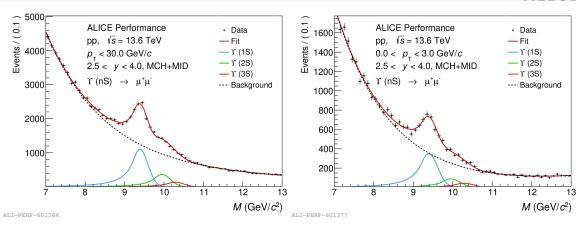


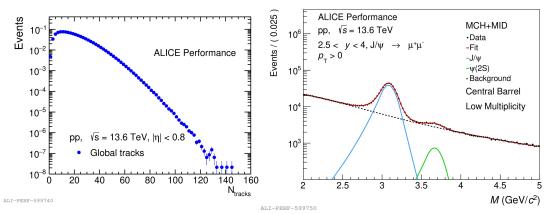
Outlook: Run 3 analyses in progress



Currently working on bottomonium measurements at $\sqrt{s} = 13.6$ TeV

Signal extraction \rightarrow possibility to discriminate $\Upsilon(1S)$, $\Upsilon(2S)$, and $\Upsilon(3S)$ states both in the integrated spectrum (left) and in the $p_{\rm T}$ -differential (right)





Ongoing measurements of charmonium production as a function of multiplicity at $\sqrt{s} = 13.6 \text{ TeV}$

Multiplicity measured at midrapidity and charmonia at forward rapidity

Conclusions



Several preliminary results in pp collisions at $\sqrt{s} = 13.6$ TeV in the charmonium sector

- J/ ψ cross section \rightarrow compatible with the measurement at $\sqrt{s} = 13$ TeV and phenomenological predictions
- $\psi(2S)$ -to-J/ ψ cross-section ratio \to compatible with measurements at $\sqrt{s} = 13$ TeV and phenomenological models, little to no energy or rapidity dependence
- J/ ψ polarization \rightarrow sizable longitudinal polarization measured at low $p_{\rm T}$: qualitative agreement with the ICEM prediction, disagreement with CSM and NRQCD
- J/ ψ -D⁰ cross section \rightarrow overestimated by PYTHIA, compatible with Helac-Onia if σ_{eff} = 31 mb

Other analyses in progress in pp collisions at $\sqrt{s} = 13.6 \text{ TeV}$:

- Bottomonium cross-section
- Charmonium yield as a function of multiplicity

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Other analyses in progress in pp collisions at $\sqrt{s} = 13.6$ TeV:

- Bottomonium cross-section
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Other important results by ALICE Run 3 in quarkonia discussed at this conference:

- > Separation of prompt and non-prompt charmonia at forward rapidity in pp collisions
 - \rightarrow Beauty production in proton-proton collisions at $\sqrt{s} = 13.6$ TeV by E. Barreau (previous talk)
- Quarkonium measurements in Pb-Pb collisions
- \rightarrow Quarkonium collectivity in Pb-Pb collisions at $\sqrt{s_{NN}} = 5.36$ TeV with ALICE by R. Cerri (next talk)
- Comprehensive overview on quarkonium and heavy-flavor production:
 - → Heavy-flavor and quarkonium measurements from pp to AA collisions by S. Trogolo

THANK YOU FOR YOUR ATTENTION

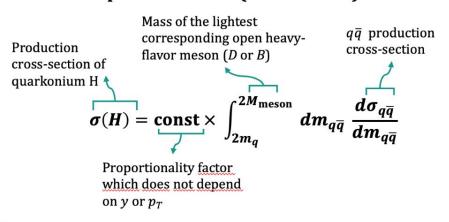


BACKUP

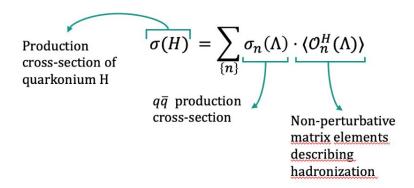
Production of Quarkonia

- ➤ Charm and bottom quarks are produced in the first instants of the partonic collision, i.e. in large energy scatterings: QCD predictions can be computed
- ➤ Their binding into quarkonia is instead characterized by smaller energy scale and cannot be described by perturbative QCD but rather by phenomenological models (non perturbative)

Color Evaporation Model (CEM or ICEM):



Non-Relativistic QCD

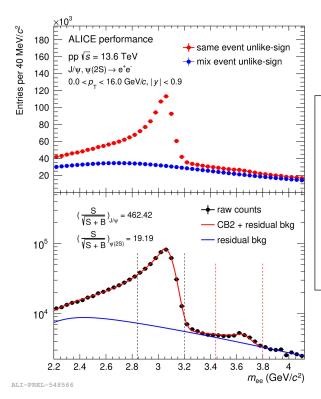


In the case of charmonia:

Contribution of charmonia from b quark decay (non prompt) in the FONLL (Fixed Order + Next to Leading Logarithm) formalism

Charmonium production at midrapidity



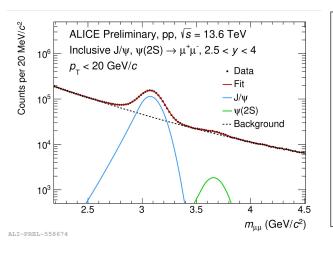


SIGNAL EXTRACTION STRATEGY

- 1. Combine opposite-sign electrons in the same event.
- Subtract opposite-sign electrons taken from mixed-events⇒ remove part of the background.
- Fit the residual distribution with two gaussian functions with asymmetric tails (J/ψ) and $\psi(2S)$, and a phenomenological function for the residual background.

Charmonium production ratio at forward rapidity





SIGNAL EXTRACTION STRATEGY

- 1. The <u>dimuon</u> invariant-mass spectrum is built by combining opposite-sign muons in the same event
- 2. <u>No subtraction of mixed-event</u> distribution (different strategy from midrapidity)
- 3. The signal is fitted with two gaussian functions with asymmetric tails (J/ψ and $\psi(2S)$), and a phenomenological function for the residual background

Rapidity and Pseudorapidity

Rapidity: measurement of velocity in special relativity, defined as

$$y = \frac{1}{2} \ln \left(\frac{E + p_L c}{E - p_L c} \right)$$

Transformation law:

 $y' = y - y_{\beta}$, $y_{\beta} = \frac{1}{2} \ln \left(\frac{1+\beta}{1-\beta} \right)$ and β is the relative velocity of R' wrt R

Rapidity is complex to measure: it requires particle identification or combined measurement of energy and longitudinal momentum

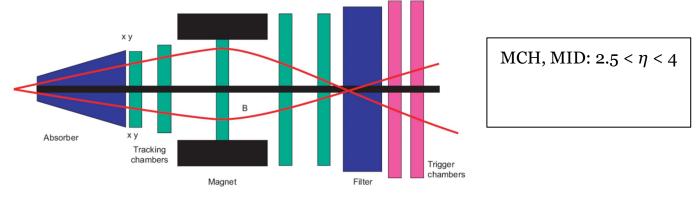
It is more convenient to use pseudorapidity, defined as

$$\eta = -\ln\left(\tan\left(\frac{\theta}{2}\right)\right) = \frac{1}{2}\ln\left(\frac{|\vec{p}| + p_L}{|\vec{p}| - p_L}\right)$$

For a relativistic particle, $|\vec{p}|c \sim E$ and therfore $\eta \approx y$

2) ALICE Muon Spectrometer



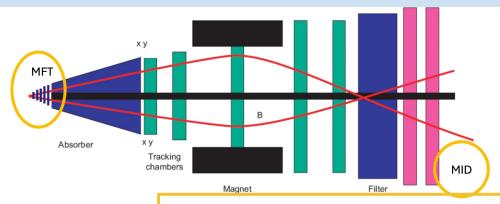


<u>Run 2</u> (2015-2018)

- 1. Hadron Absorber
- 2.Muon (tracking) Chambers (MCH)
- 3. Muon Trigger

2) ALICE Muon Spectrometer





MCH, MID: $2.5 < \eta < 4$

MFT: $2.5 < \eta < 3.6$

Run 2 (2015-2018)

- 1.Hadron Absorber
- 2.Muon (tracking) Chambers (MCH)
- 3. Muon Trigger

<u>Run 3</u> (2022-ongoing)

- O. New Muon Forward Tracker (MFT) before absorber
- 1. Hadron Absorber
- 2. Muon (tracking) Chambers (MCH): new front-end electronics and readout chain
- 3. Muon Identifier (MID) with new front-end electronics and readout chain

Major changes to provide:

- Discrimination between primary and secondary vertices (through MFT)
- Continuous readout (no trigger) through upgrade of front-end and read-out electronics

Muon Tracking Chambers

Made of 5 stations of 2 planes of multi-wire proportional chambers

Both Cathode Pad Chamber (CPC) and Cathode Strip Chamber technologies are used in order to provide 2D information

Muon Tracking Chambers

2 stations of 2 planes of Resistive Plate Chambers (RPC) each

Multi-parton interactions



Multiple parton interactions (MPI) in a single hadronic collision

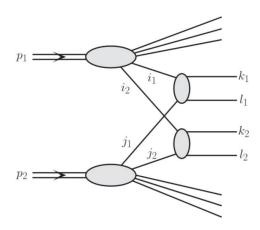
- can explain QGP-like effects in small systems
- inputs to tune MC generators

Simplest case: double-parton-scattering (DPS)

$$\sigma_{A,B}^{DPS} = \frac{m}{2} \frac{\sigma_A \sigma_B}{\sigma_{eff}}$$
 m=1 if A=B m=2 otherwise

 $\sigma_{\it eff}$ = scaling factor \rightarrow physical observable

Independent on process and collision energy in the simplest model



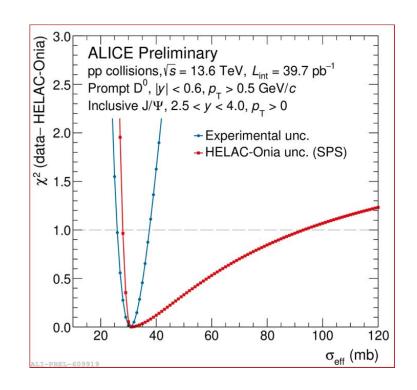
J/ψ -D⁰ associated production



Fit to extract $\sigma_{\it eff}$ using Helac-Onia predictions

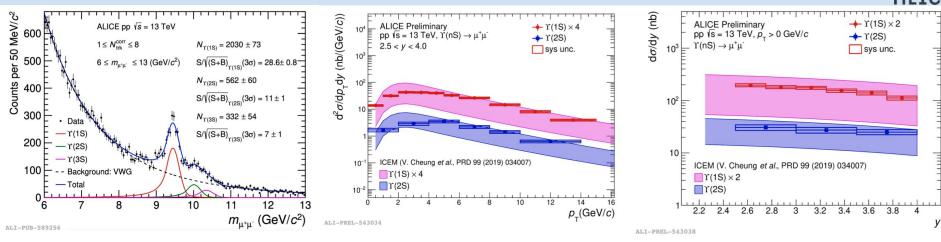
Uncertainties on $\sigma_{\it eff}$ extraction are currently dominated by theoretical prediction

$$\rightarrow \sigma_{eff} = 31 \text{ mb}$$



Bottomonium production in Run 2

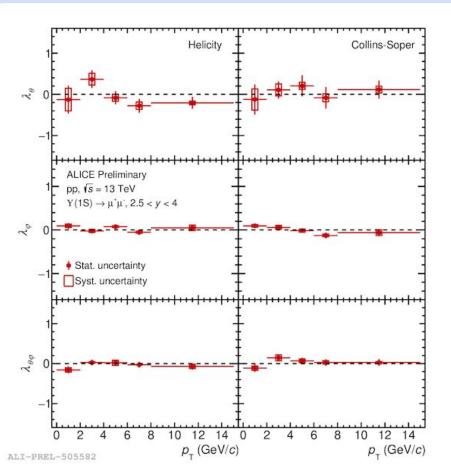




- As for bottomonia, the signal extraction allows for discrimination of $\Upsilon(1S)$, $\Upsilon(2S)$, and $\Upsilon(3S)$ states
- ICEM correctly predicts differential yields of $\Upsilon(1S)$ and $\Upsilon(2S)$ vs p_T and rapidity at 13 TeV at forward rapidity
- Measurements at 13.6 TeV at forward rapidity are ongoing

Bottomonium polarization in Run 2



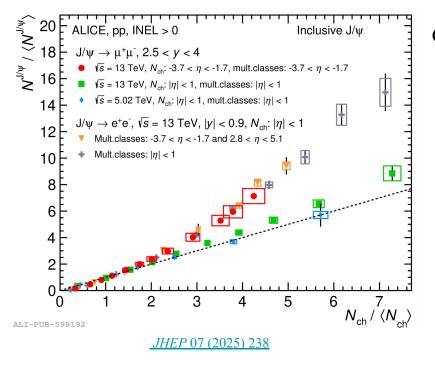


- All Y(1S) polarization coefficients measured at 13 TeV are found to be compatible with 0 within the statistical and systematic uncertainties in both the Helicity and the Collins-Soper frames
- Ongoing measurement at 13.6 TeV with the aim of increasing the sample size and reducing the uncertainties

Charmonium production vs multiplicity (Run 2)

Quarkonium production as a function of the total number of produced particles in a collision (multiplicity)

→ understand the role of multi-parton interactions in quarkonium production in pp collisions



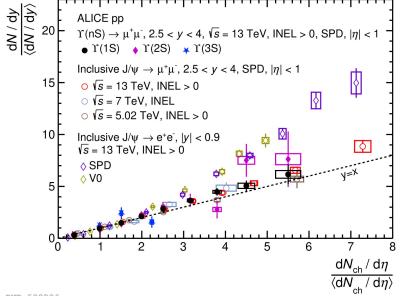
Charmonia at 13 TeV:

- Linear dependence vs multiplicity at low multiplicities
- At high multiplicities:
 - \circ J/ ψ at forward rapidity, multiplicity estimator at midrapidity
 - → linear dependence
 - J/ψ at forward rapidity, multiplicity estimator at forward rapidity
 - → steeper-than-linear dependence
 - \circ J/ ψ at midrapidity with both multiplicity estimators
 - → steeper-than-linear dependence

Bottomonium production vs multiplicity (Run 2)

Quarkonium production as a function of the total number of produced particles in a collision (multiplicity)

→ understand the role of multi-parton interactions in quarkonium production in pp collisions



Bottomonia at 13 TeV:

- Linear dependence vs multiplicity at low multiplicities
- At high multiplicities:
 - Compatible with linear correlation

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