

Calibrating the Medium Effects of Light Clusters in Heavy-Ion Collisions

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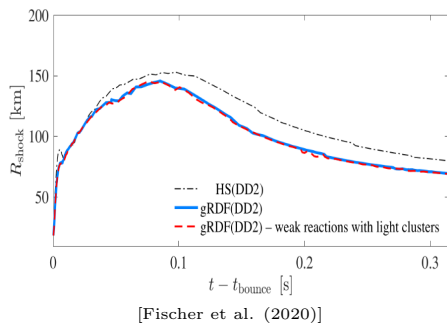
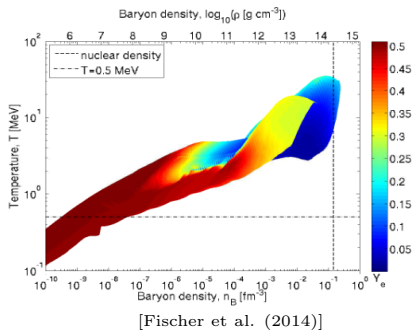
Tiago Custódio¹, Alex Rebillard-Soulié², Rémi Bougault², Diégo Gruyer², Francesca Gulminelli², Tuhin Malik¹, Helena Pais¹ and
Constança Providência¹

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²LPC, Caen

Motivation

- Light nuclei might be present in both Core-Collapse Supernova and Binary Neutron Star Mergers
- Their presence influences the dynamics of these astrophysical events
- Accounting for in-medium modifications to the light clusters is essential to determine their correct abundances



Relativistic Nuclear Field Theory

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- The Lagrangian density for matter made of protons, neutrons and light clusters is:

$$\mathcal{L} = \sum_{b=n,p} \mathcal{L}_b + \sum_{m=\sigma,\omega,\rho} \mathcal{L}_m + \sum_{i=^2\text{H}, ^3\text{H}, ^3\text{He}, ^4\text{He}} \mathcal{L}_i \quad (2)$$

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- The Lagrangian density for the **nucleons** is:

$$\mathcal{L}_b = \bar{\Psi}_b(x) [i\gamma_\mu D_b^\mu - m_b^*] \Psi_b(x) , \quad (3)$$

$$iD_b^\mu = i\partial^\mu - \textcolor{red}{g_{\omega b}}\omega^\mu - \textcolor{red}{g_{\rho b}}\vec{I}_b \cdot \vec{\rho}^\mu \quad (4)$$

$$m_b^* = m_b - \textcolor{red}{g_{\sigma b}}\sigma \quad (5)$$

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- Calibrated to experimental nuclear parameters (e.g. FSU, DD2)

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- Light clusters can be included as point-like independent quasi-particles, in the same way as nucleons, taking into account their corresponding spins

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$$\begin{aligned} \mathcal{L}_i = & \frac{1}{4} (i D_i^\mu \Psi_i^\nu - i D_i^\nu \Psi_i^\mu)^* (i D_{\mu i} \Psi_{\nu i} - i D_{\nu i} \Psi_{\mu i}) \\ & - \frac{1}{2} \Psi_i^{\mu*} (M_i^*)^2 \Psi_{\mu i}, \quad i = {}^2\text{H} \end{aligned} \quad (8)$$

Relativistic Nuclear Field Theory

- Light clusters ($i = {}^2\text{H}, {}^3\text{H}, {}^3\text{He}, {}^4\text{He}$) will have their own cluster-meson couplings:

$$g_{\sigma i} = x_s A_i g_{\sigma N} \quad (9)$$

$$g_{\omega i} = A_i g_{\omega N} \quad (10)$$

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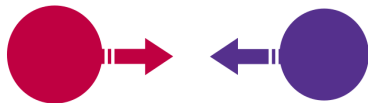
- $x_s(\rho, T)$ is a way of accounting for in-medium modification of the clusters self-energies

INDRA Heavy-Ion Collisions

$^{136,124}\text{Xe} + ^{124,112}\text{Sn}$ (32 MeV/nucleon)

- Only **central collisions** selected (most violent)

Projectile-target central collision



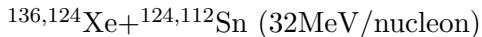
Projectile-target central collision



Angular selection : **mid-velocity products**

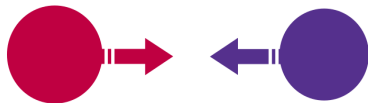
Credits: Alex Rebillard-Soulié

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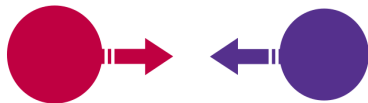
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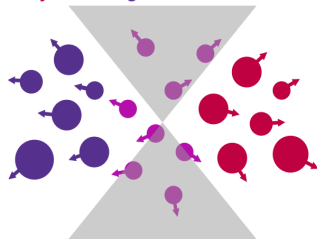
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- **Angular selection** to reduce secondary decays from other sources

Projectile-target central collision



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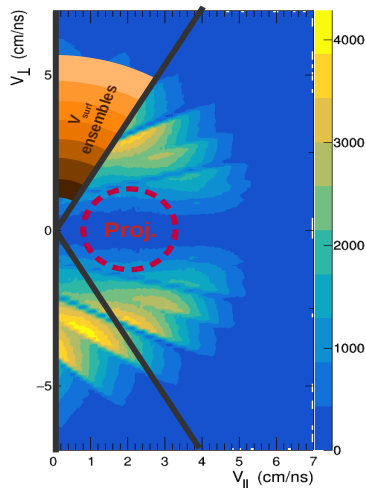


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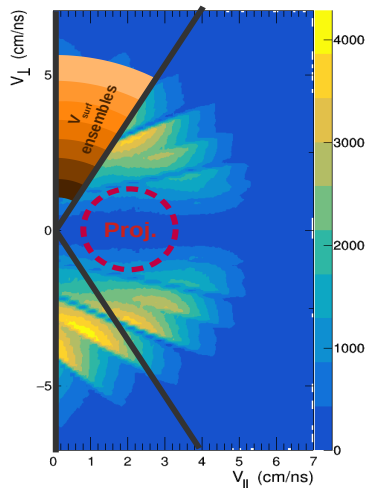
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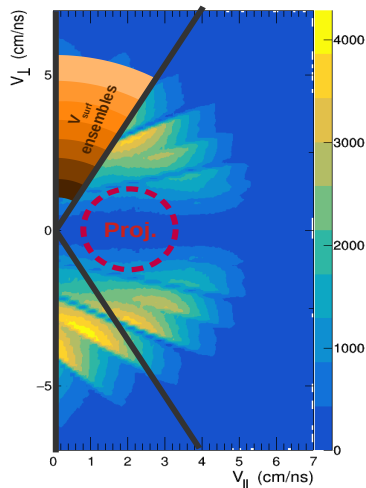
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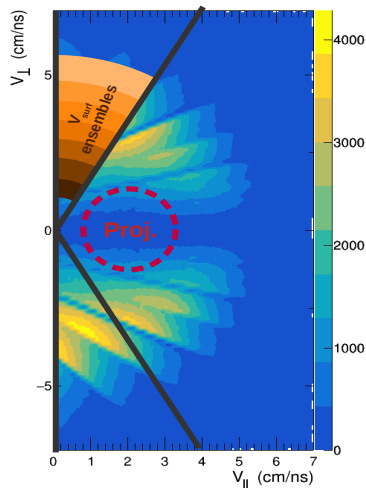
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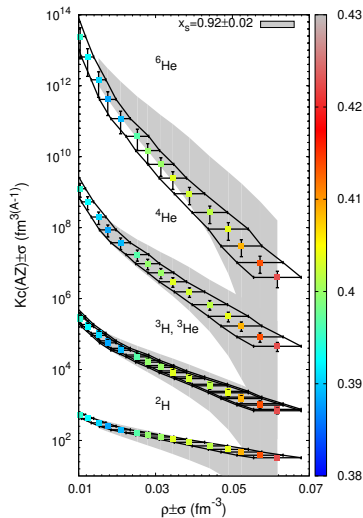
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- Associate a statistical ensemble to each v_{surf} with corresponding particle mass fractions (nucleons and light clusters)



Credits: Alex Rebillard-Soulié

Calibration of x_s with Kc from HIC

- T and ρ estimated considering an ideal gas of clusters in the grand canonical ensemble
[Bougault et al. (2020)]



Pais et al. (2020)

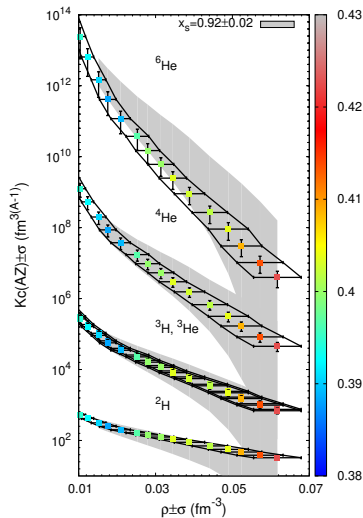
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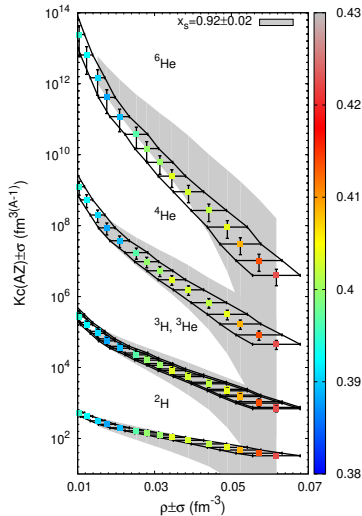


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$$Kc_i = \frac{\rho_i}{\rho_n^{N_i} \rho_p^{Z_i}}$$



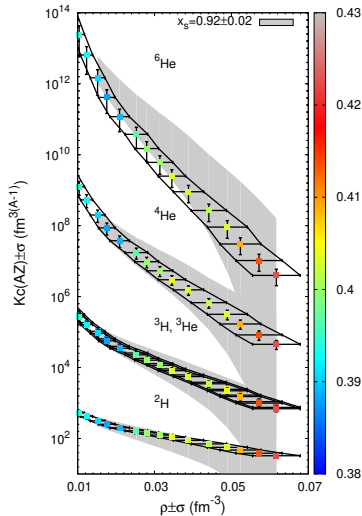
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- x_s has been calibrated by roughly considering the values that best fit the data visually



Pais et al. (2020)

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- The statistical ensembles will be described using RMF theory

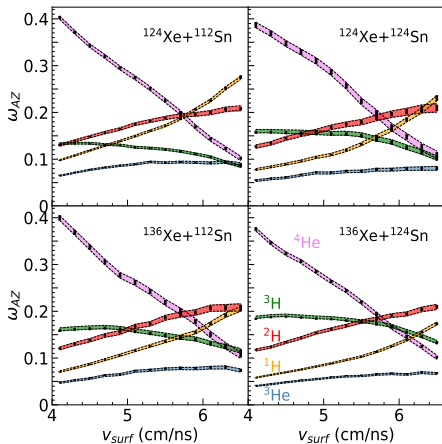
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- For each system and v_{surf} bin, we carry out an independent Bayesian inference on the measured mass fractions ($4N_v = 52$)

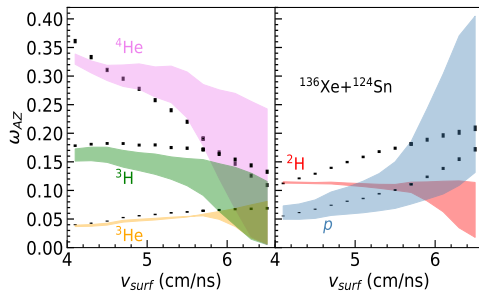
$$p_i(\theta|\{\omega_{AZ}\}) = \frac{p_\theta}{\mathcal{Z}} \mathcal{L}_g(\{\omega_{AZ}\}_i|\theta), \quad \theta = \{T, \rho, x_s(\rho, T)\} \quad (12)$$

Mass Fractions

Present Study (FSU & DD2)

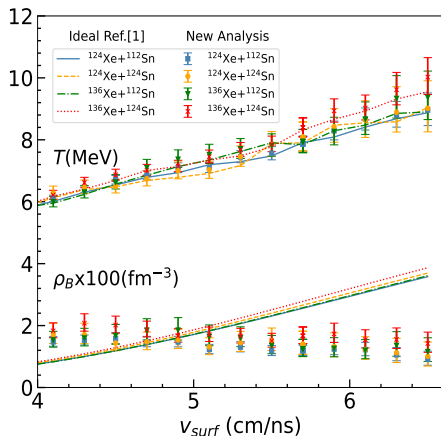
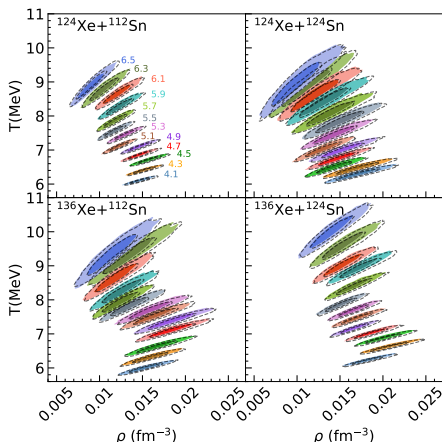


Modified Ideal Gas (FSU)

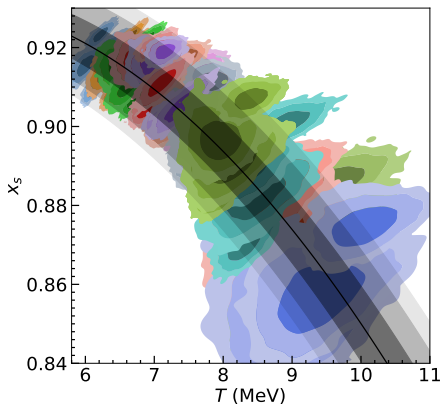


Calibrated Temperatures and Densities

- Temperature evolution similar to the ideal gas estimation
- Results compatible with a single density $\sim 0.015 \text{ fm}^{-3}$: chemical freeze-out density at the surface of the emitting source (?)



Calibrated $x_s(T)$

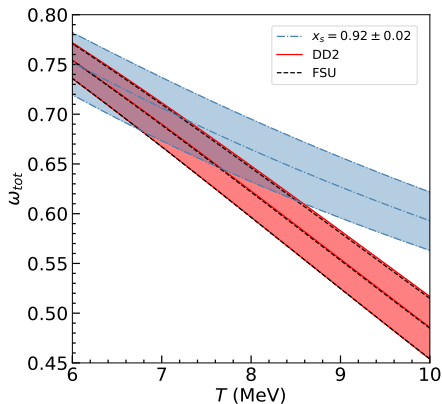


- x_s is temperature dependent
- Interaction weakens with T
- $x_s(T)$ compatible for all four entrance channels
- Limited ρ range cannot provide information on possible x_s dependence on ρ

Parameter	Unit	Median	1σ	2σ
a	MeV^{-2}	-0.00203	± 0.00003	± 0.00006
b	MeV^{-1}	0.01477	± 0.00047	± 0.00093
c		0.90560	± 0.0018	± 0.00355

Table: Parameter estimates a, b, c with $1, 2\sigma$ uncertainties for the quadratic fit
 $x_s = aT^2 + bT + c$

Consequences of $x_s(T)$ for light cluster abundances



- Above $T \sim 8$ MeV abundances are systematically lower than the predictions of modified ideal gas
- Smaller x_s corresponds to weaker cluster- σ coupling, resulting in less bound clusters and, consequently, smaller abundances

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- T shows same increasing behaviour as before but the density turned out to be constant: chemical freeze-out (?)
- x_s shows a dependence on T , weakening the clusters binding and abundances

Thank you!